

*Precision Top Mass
Determination at the LHC
with Jet Grooming*

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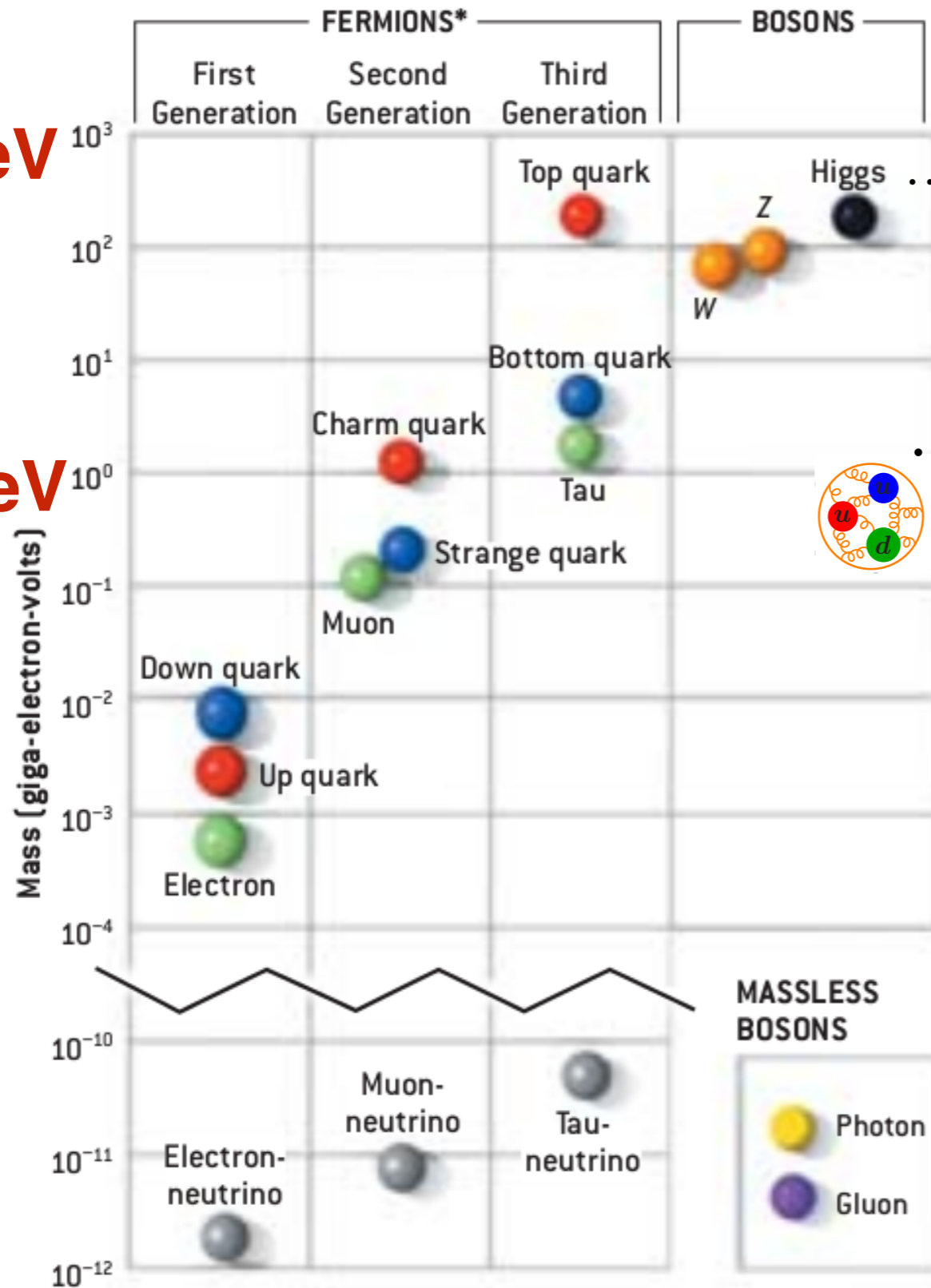
In collaboration with:
Andre Hoang, Aditya Pathak, & Iain Stewart

LHC-TI Meeting, SLAC
Feb. 10th, 2017

Top Quark is special

1 TeV

1 GeV

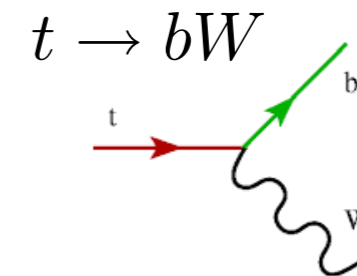


..... Au

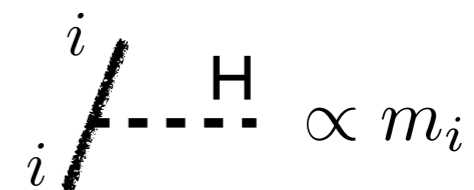
..... proton

- Heaviest known elementary particle

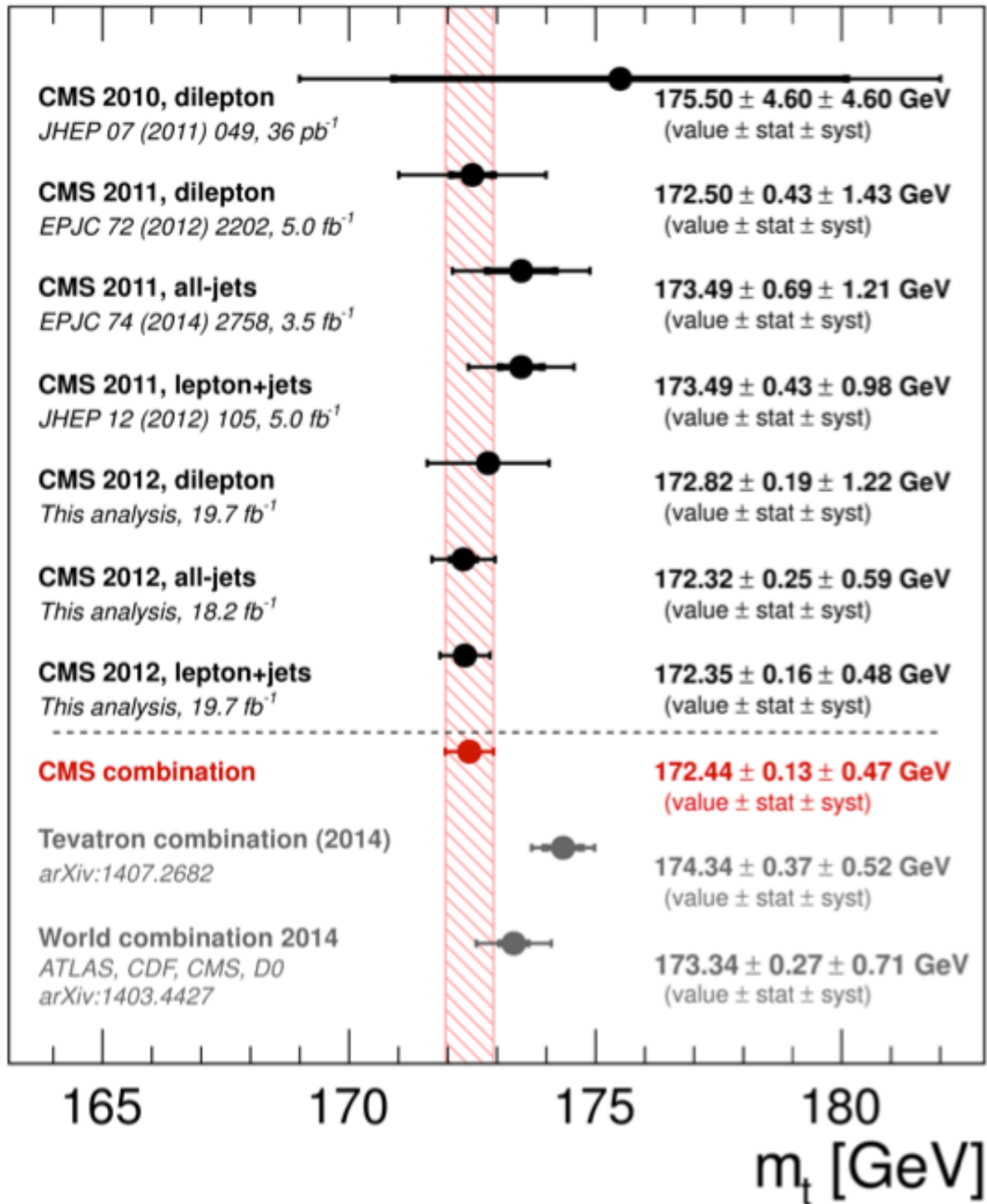
- Decays before hadronization



- Large coupling to the Higgs



Precision Measurements



Tevatron (2014): $m_t = 174.34 \pm 0.64$ GeV

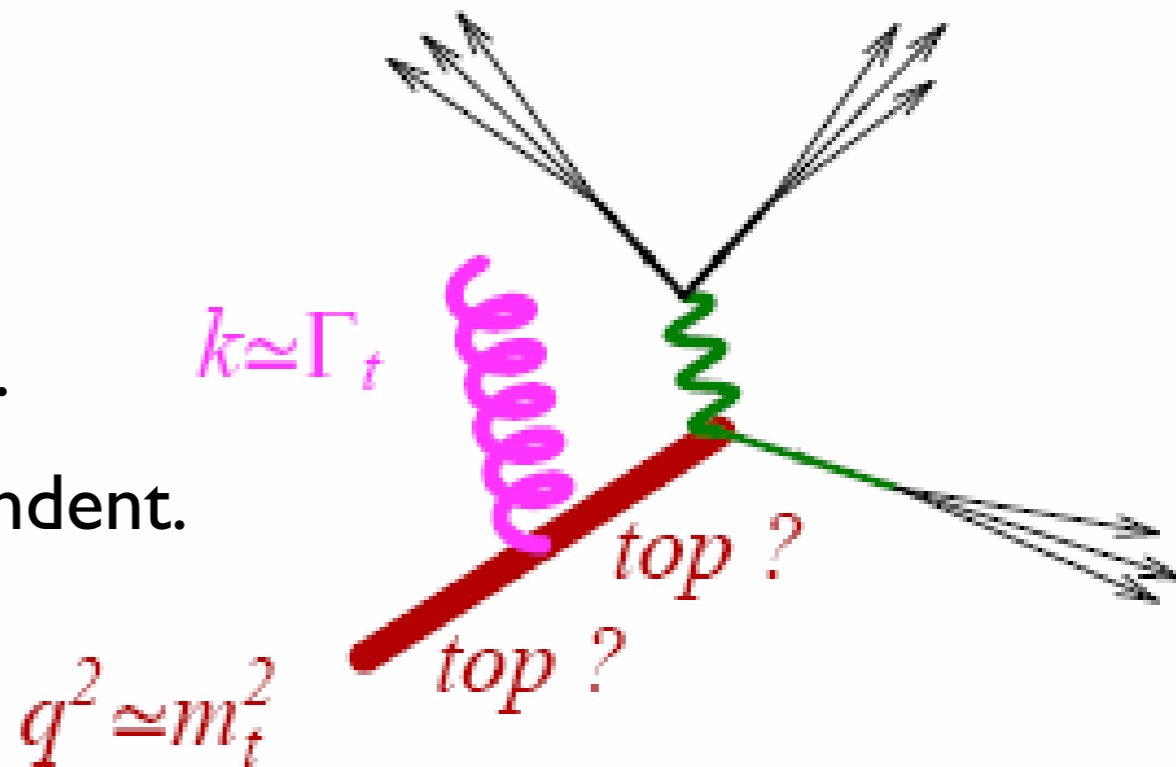
CMS Run 1 (2015): $m_t = 172.44 \pm 0.49$ GeV

ATLAS Run 1 (2016): $m_t = 172.84 \pm 0.70$ GeV

The Top Quark Mass

What is the top mass?

- Top is colored parton.
- Top mass is a parameter of the Lagrangian.
- Top mass is renormalization scheme dependent.



Which top mass?

- Which mass scheme is being measured in kinematic reconstruction methods?
- Mass extracted corresponds to the Monte Carlo (MC) mass.
- How can one relate the MC mass to a well-defined renormalization scheme?

Top Mass Schemes

$$m_t^{\text{pole}} = m_t(R, \mu) + \delta m_t(R, \mu)$$

$$\delta m_t(R, \mu) = R \sum_{n=1}^{\infty} \sum_{k=0}^n a_{nk} \left[\frac{\alpha_s(\mu)}{4\pi} \right]^n \ln^k \frac{\mu}{R}$$

- Different schemes correspond to different coefficients and values for “R”

$\overline{\text{MS}}$	$\overline{m}(\mu),$	$R = \overline{m}(\mu);$
RGI [2]	$m_{\text{RGI}},$	$R = m_{\text{RGI}};$
kinetic [3]	$m_{\text{kin}},$	$R = \mu_f^{\text{kin}};$
1S [4]	$m_{1\text{S}},$	$R = m_{1\text{S}} C_F \alpha_s(\mu)$
PS [5]	$m_{\text{PS}},$	$R = \mu_f^{\text{PS}}.$

“Top Resonance Mass Schemes” $R \sim \Gamma_t$

(Fleming, Hoang, SM, Stewart)

$$m_t^{\text{pole}} = m_t(R, \mu) + \delta m_t(R, \mu)$$

$$\delta m_t(R, \mu) = R \sum_{n=1}^{\infty} \sum_{k=0}^n a_{nk} \left[\frac{\alpha_s(\mu)}{4\pi} \right]^n \ln^k \frac{\mu}{R}$$

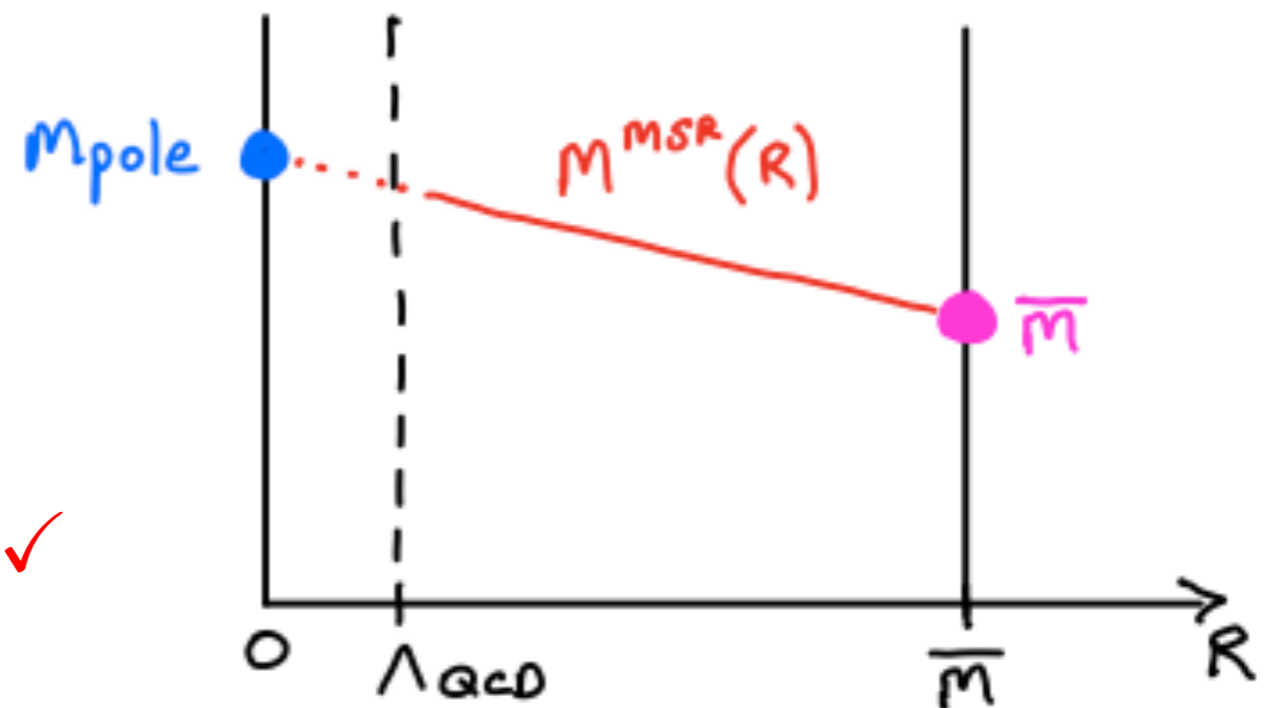
- MSR mass scheme:

(Hoang, Jain, Scimemi, Stewart)

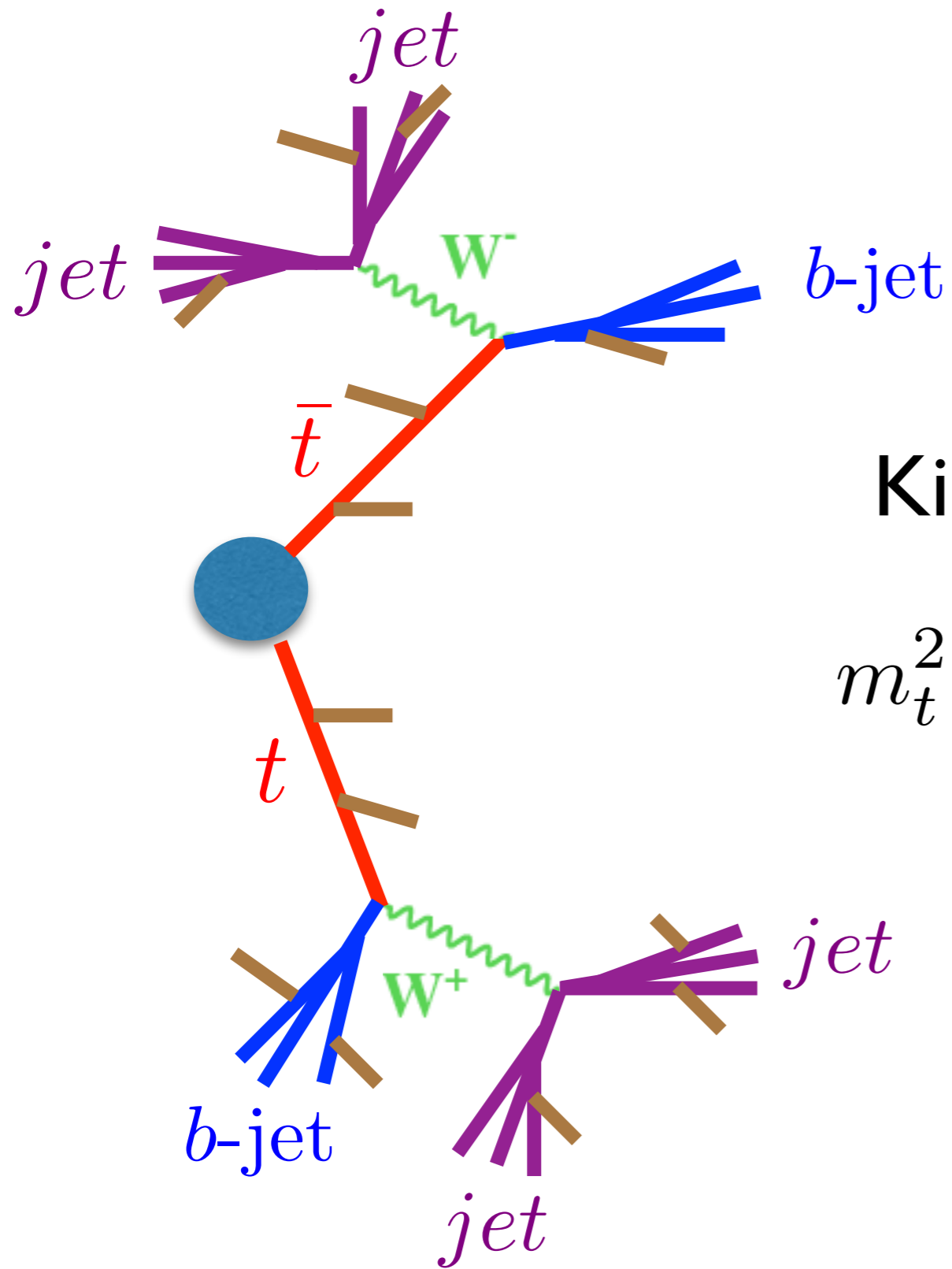
Define using $\overline{\text{MS}}$ coefficients a_{nk}

No ambiguity, $R > \Lambda_{\text{QCD}}$ ✓

Compatible with Breit Wigner, $R \sim \Gamma_t$ ✓



Top Mass from Direct Reconstruction



Kinematic Fit:

$$m_t^2 = p_t^2 = (p_{Jb} + p_{J1} + p_{J2})^2$$

Top Mass from Direct Reconstruction

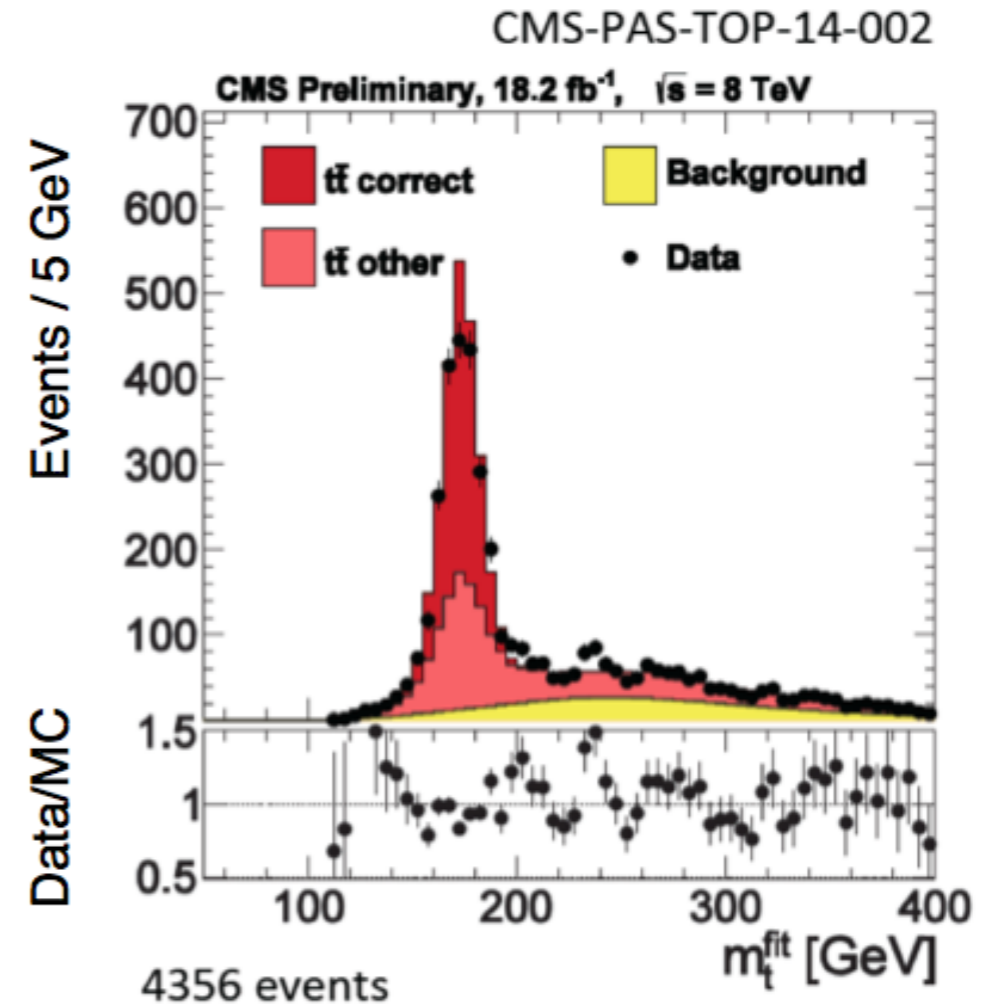
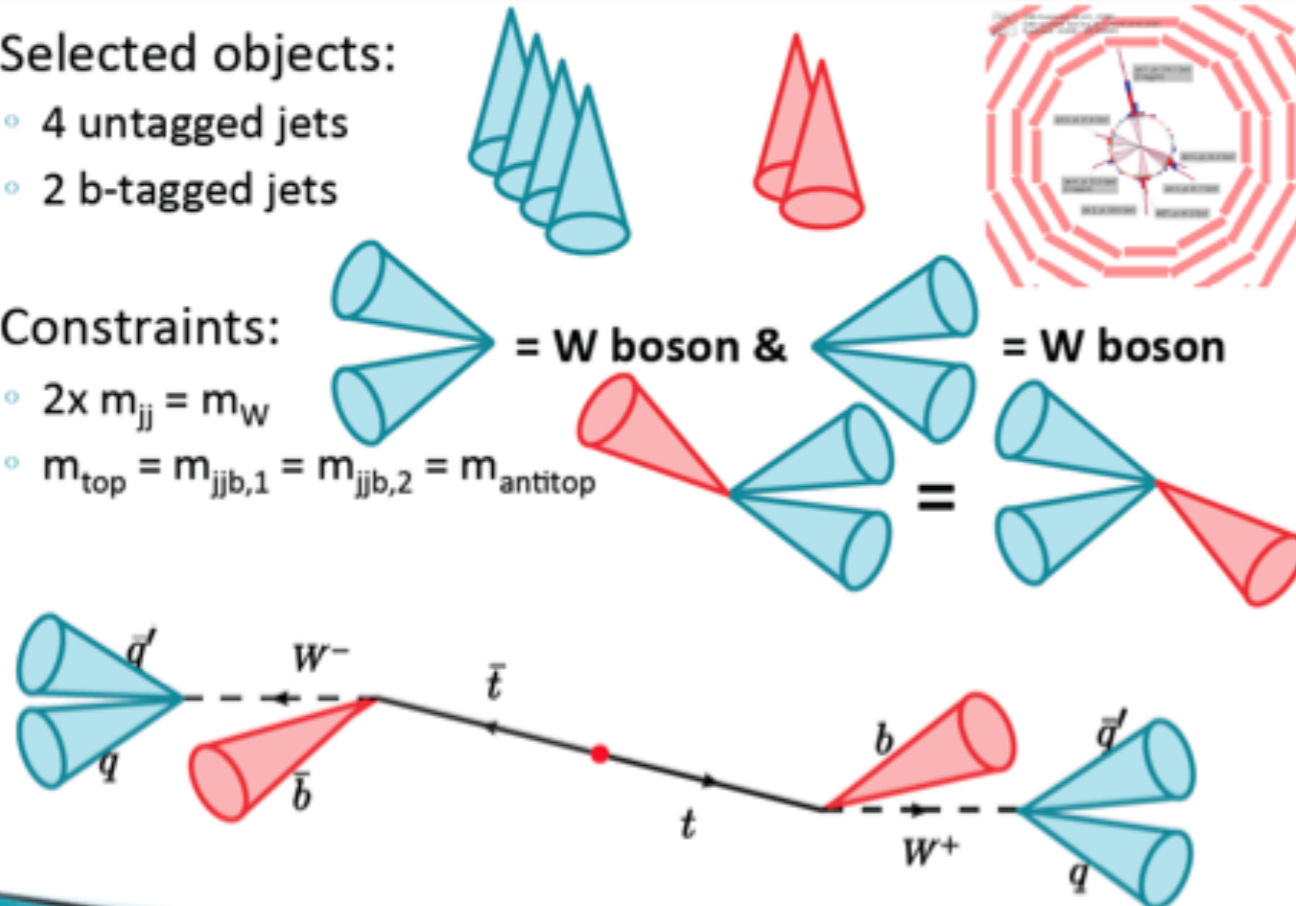
Kinematic Fit

▶ Selected objects:

- 4 untagged jets
- 2 b-tagged jets

▶ Constraints:

- $2 \times m_{jj} = m_W$
- $m_{\text{top}} = m_{jjb,1} = m_{jjb,2} = m_{\text{antitop}}$



all-jets channel at 8 TeV

Use Monte Carlo Simulations for templates

Determine best fit value for Monte Carlo top-mass parameter

CMS: $m_t^{\text{MC}} = 172.44 \pm 0.49$

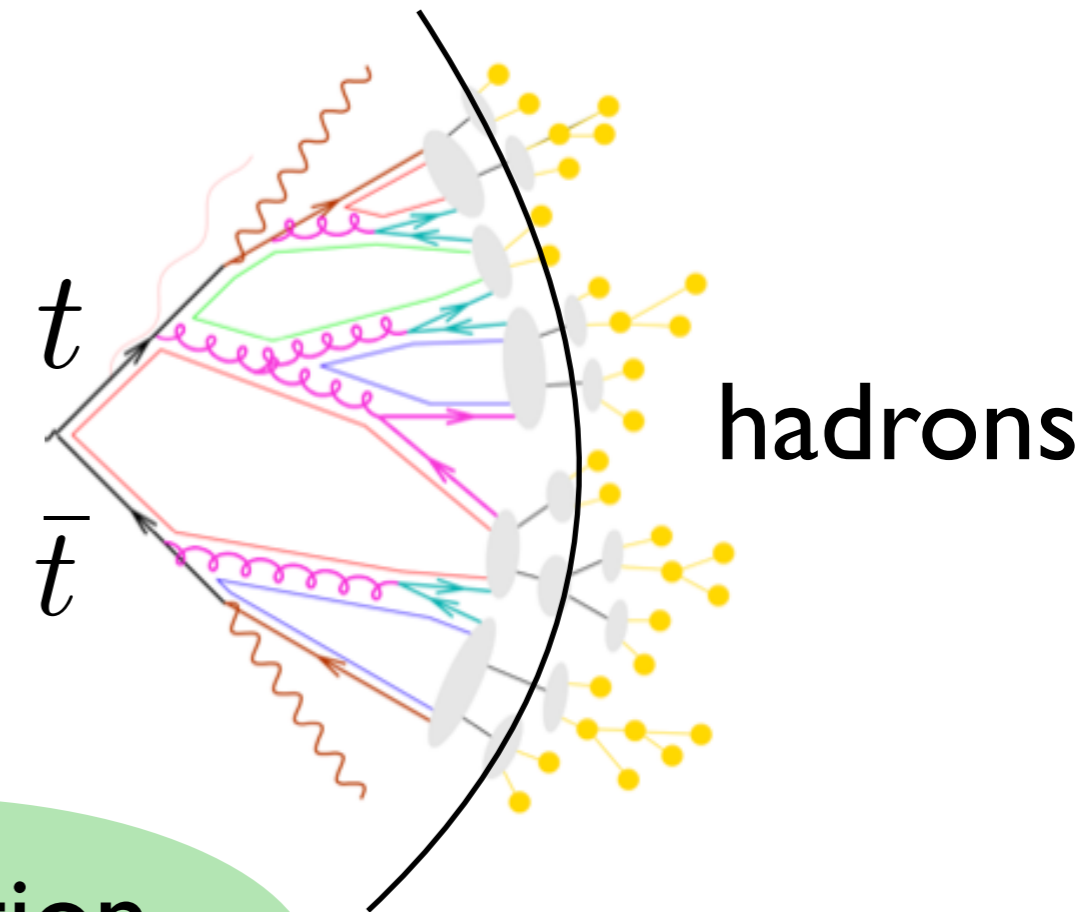
ATLAS: $m_t^{\text{MC}} = 172.84 \pm 0.70$

$m_t^{\text{pole}}, \bar{m}_t, m_t^{\text{MSR}}, \dots$

Theory (QFT)

Simulation
(Monte Carlo)

Experiment



$\Lambda^{\text{shower}} = 1 \text{ GeV}$

m_t^{MC}

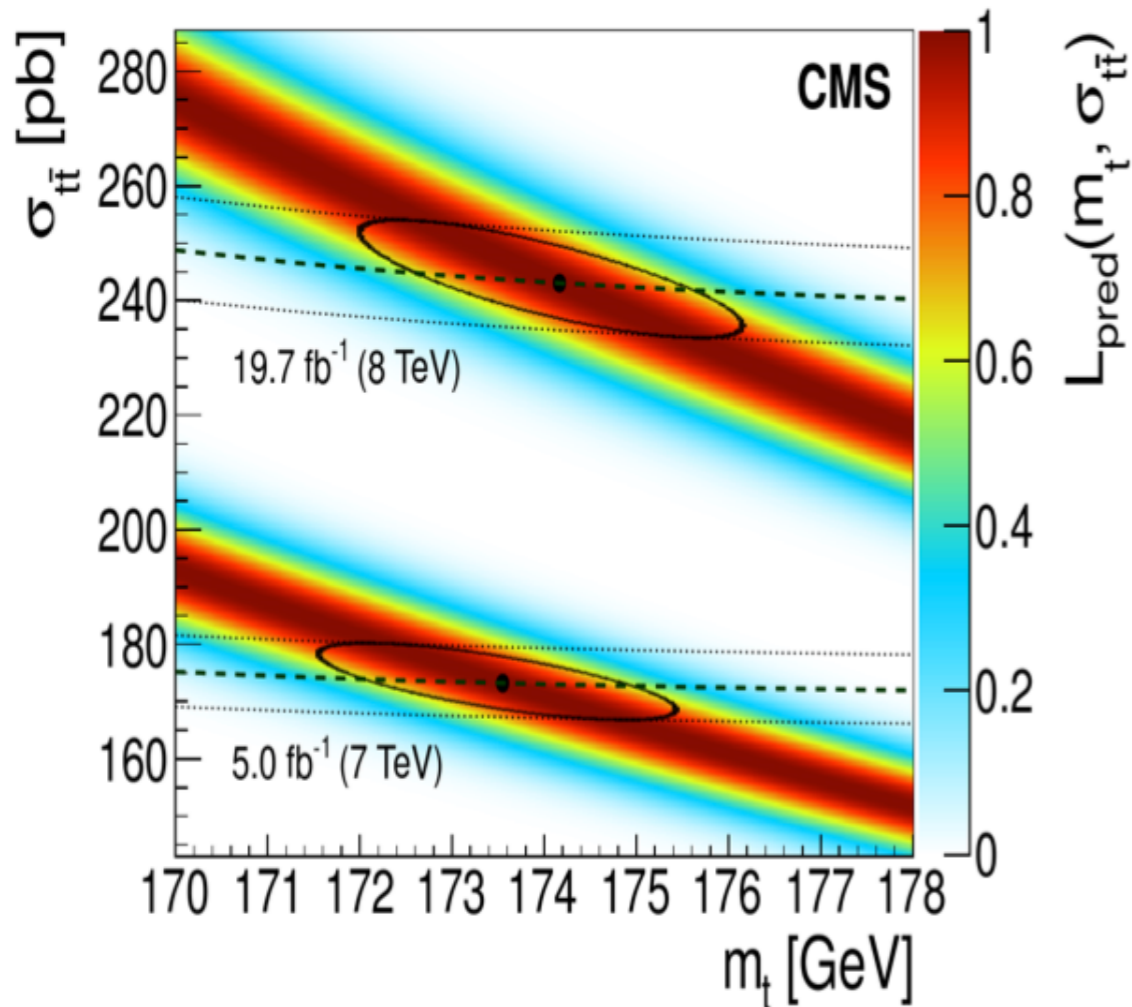
No Ambiguity ✓

Breit Wigner ✓

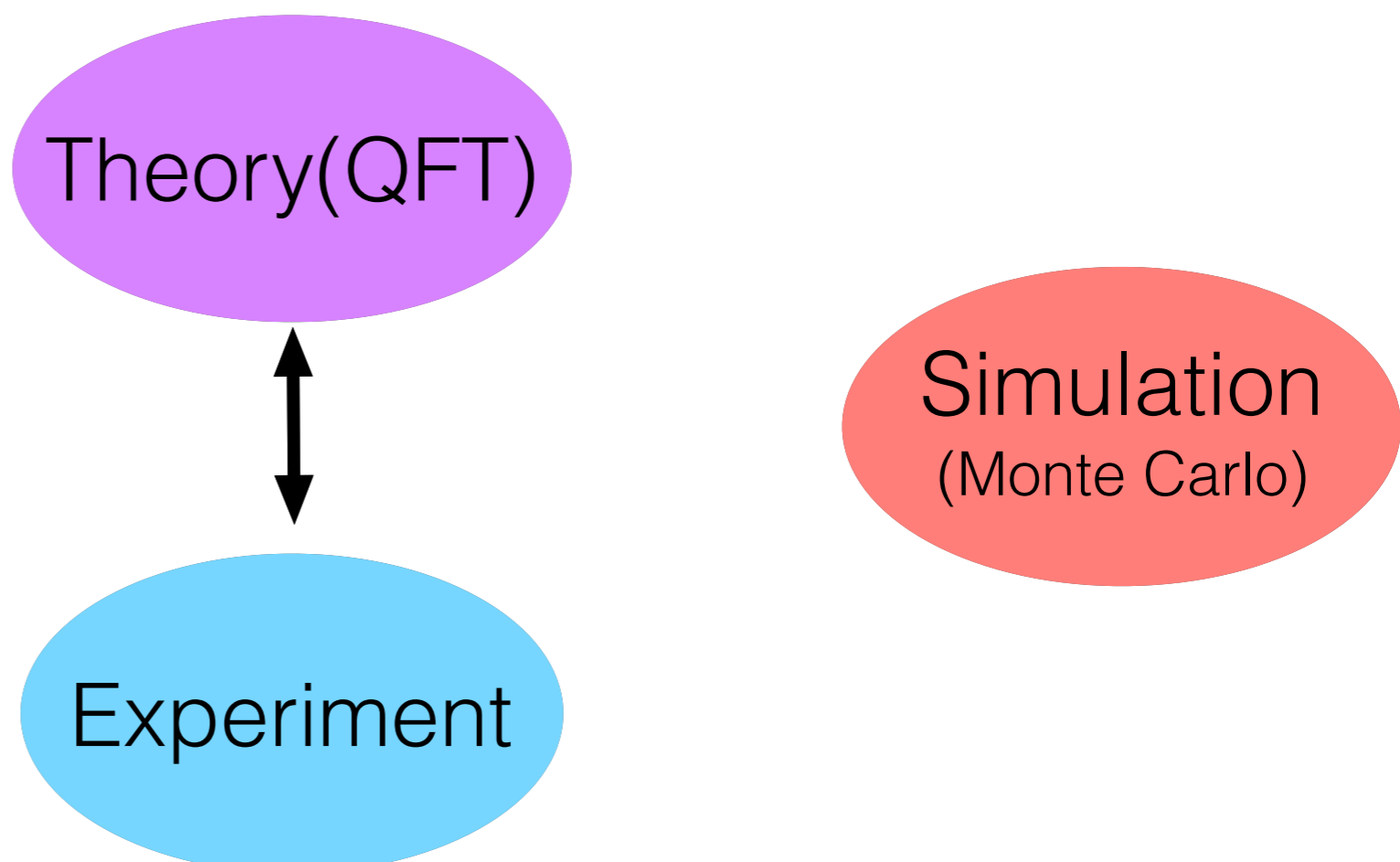
Definition ?

Top Pole Mass from Total Cross-Section

- Clean interpretation of top mass scheme is possible from cross-section measurements



$$\sigma^{\text{exp}}(pp \rightarrow t\bar{t}) = \sigma_{t\bar{t}}^{\text{th}}(m_t)$$



	m_t [GeV]
NNPDF3.0	$173.8^{+1.7}_{-1.8}$
MMHT2014	$174.1^{+1.8}_{-2.0}$
CT14	$174.3^{+2.1}_{-2.2}$

NNLO+NNLL
(Czakon,Fielder,Mitov)

Threshold Scan

(Fadin, Khoze; Peskin, Strassler; Hoang, Manohar, Stewart, Teubner; Kuhn, Martinez, Miquel; Beneke, Kiyo, Schuller; Kniehl, Penin,...)

- Clean interpretation of top mass scheme is possible from Threshold Scan

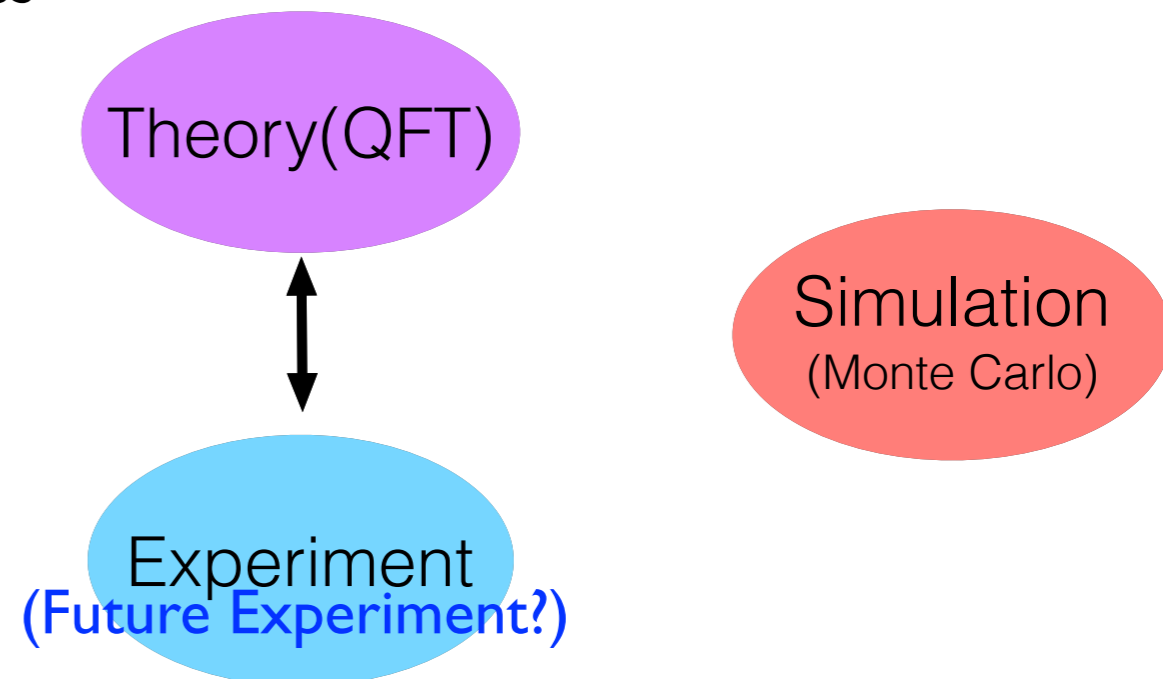
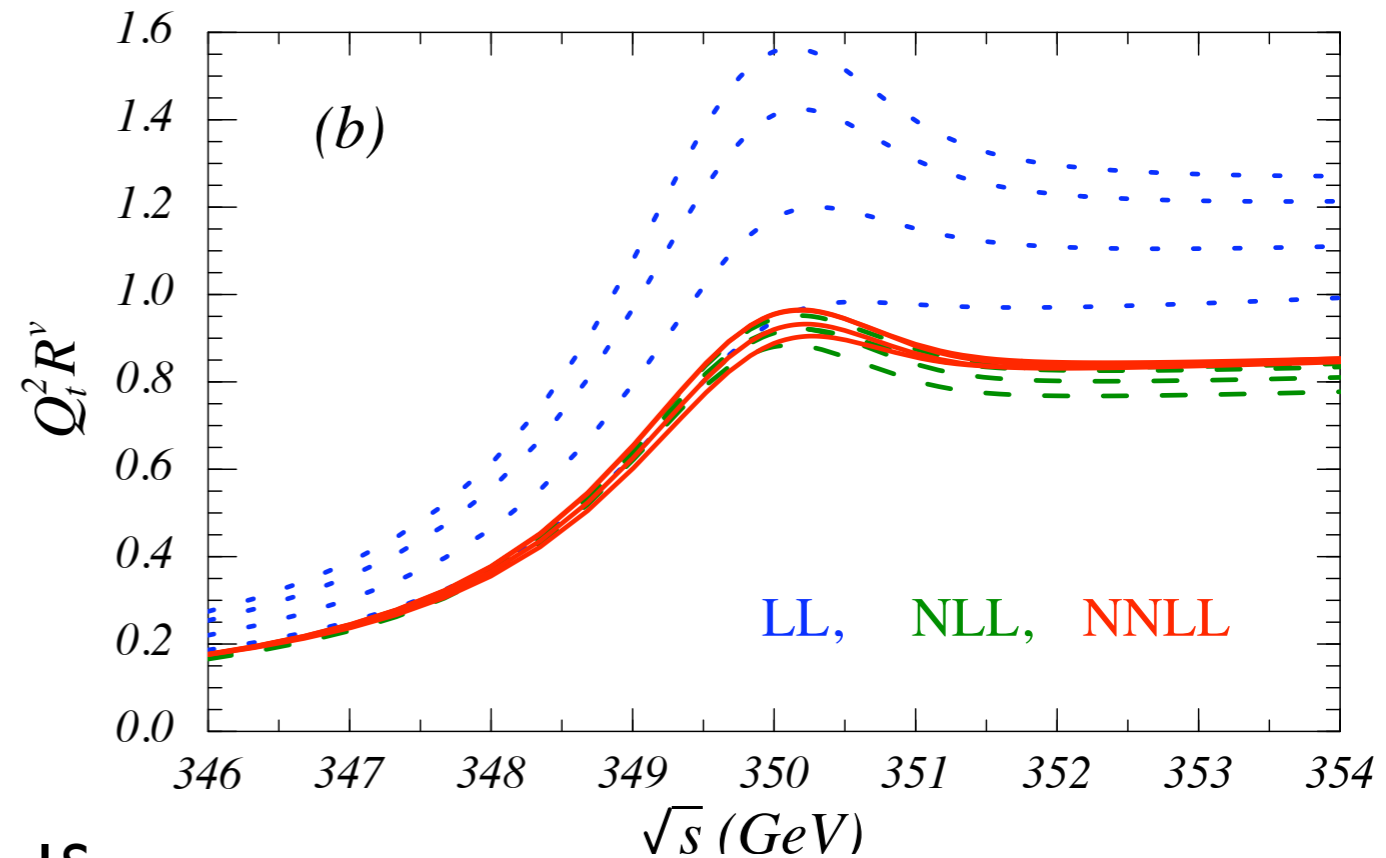
Top pair-production at threshold

- Shape of cross-section sensitive to top mass.
- Top width provides IR cutoff.
- Non-perturbative effects are small.

Physics well understood

- NRQCD is the appropriate EFT.
- Well-defined relation to short distance mass; eg. \overline{MS} mass
- Precision:

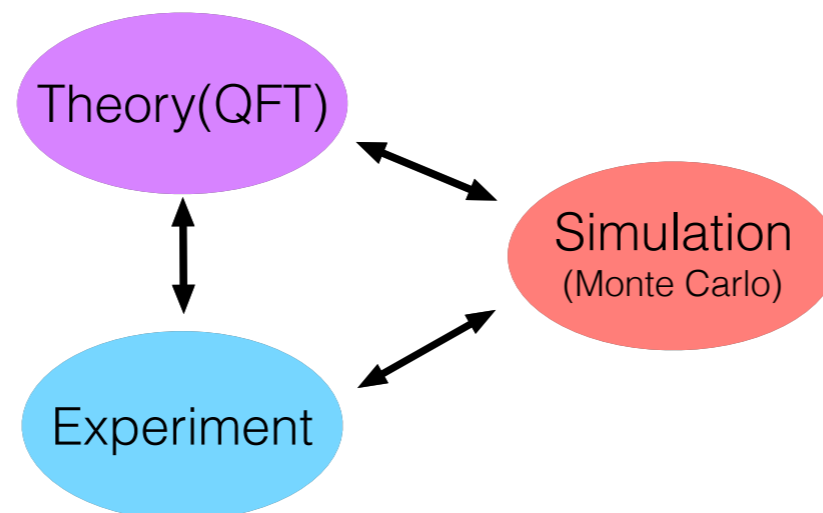
$$\delta m_t \sim 100 \text{ MeV}$$



Improving Top Mass Measurement at the LHC

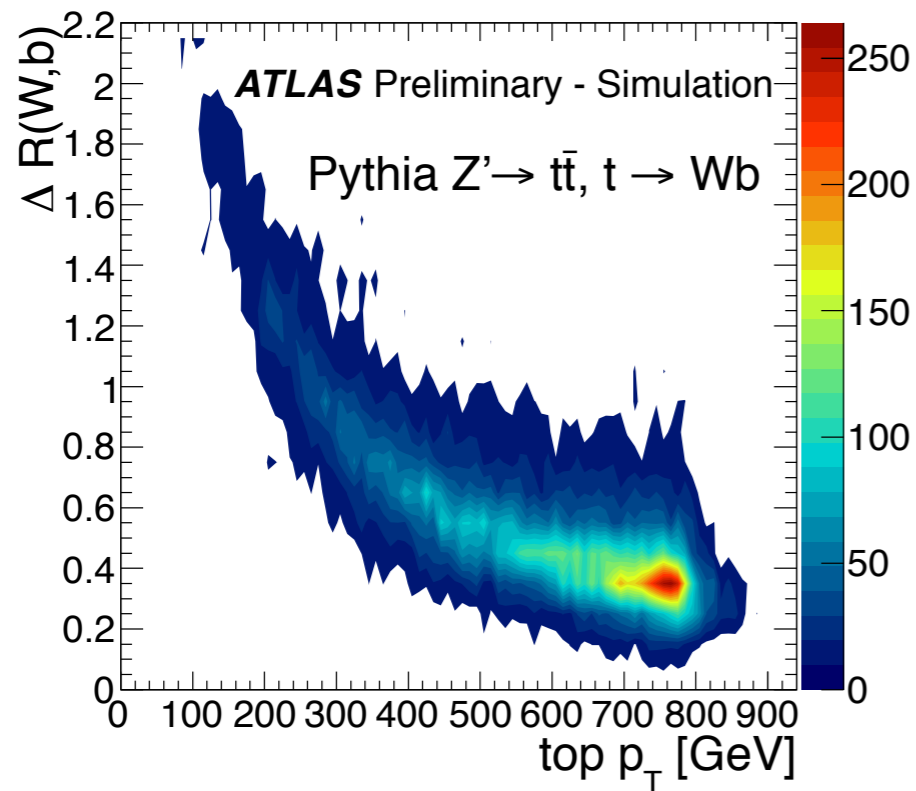
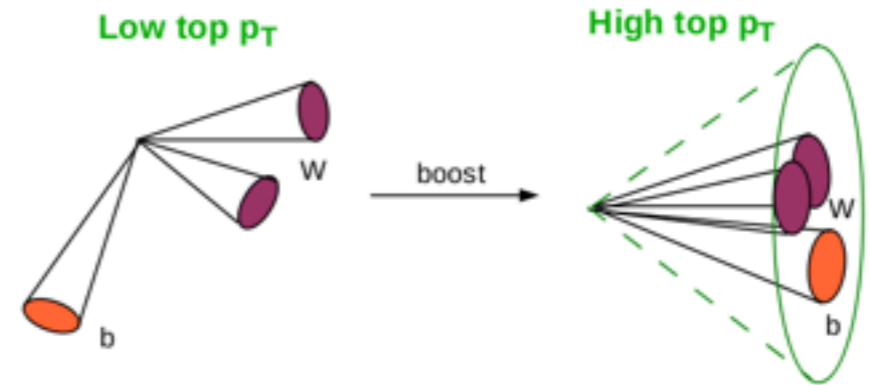
$$M_t^{\text{peak}} = m_t + (\text{nonperturbative effects}) + (\text{perturbative effects})$$

- Observable must be kinematically sensitive to the Top Mass
- Observable must be theoretically tractable
- Observable must have well-defined relation to top mass scheme
- Control contamination from hadron collider environment



Boosted Top Quarks

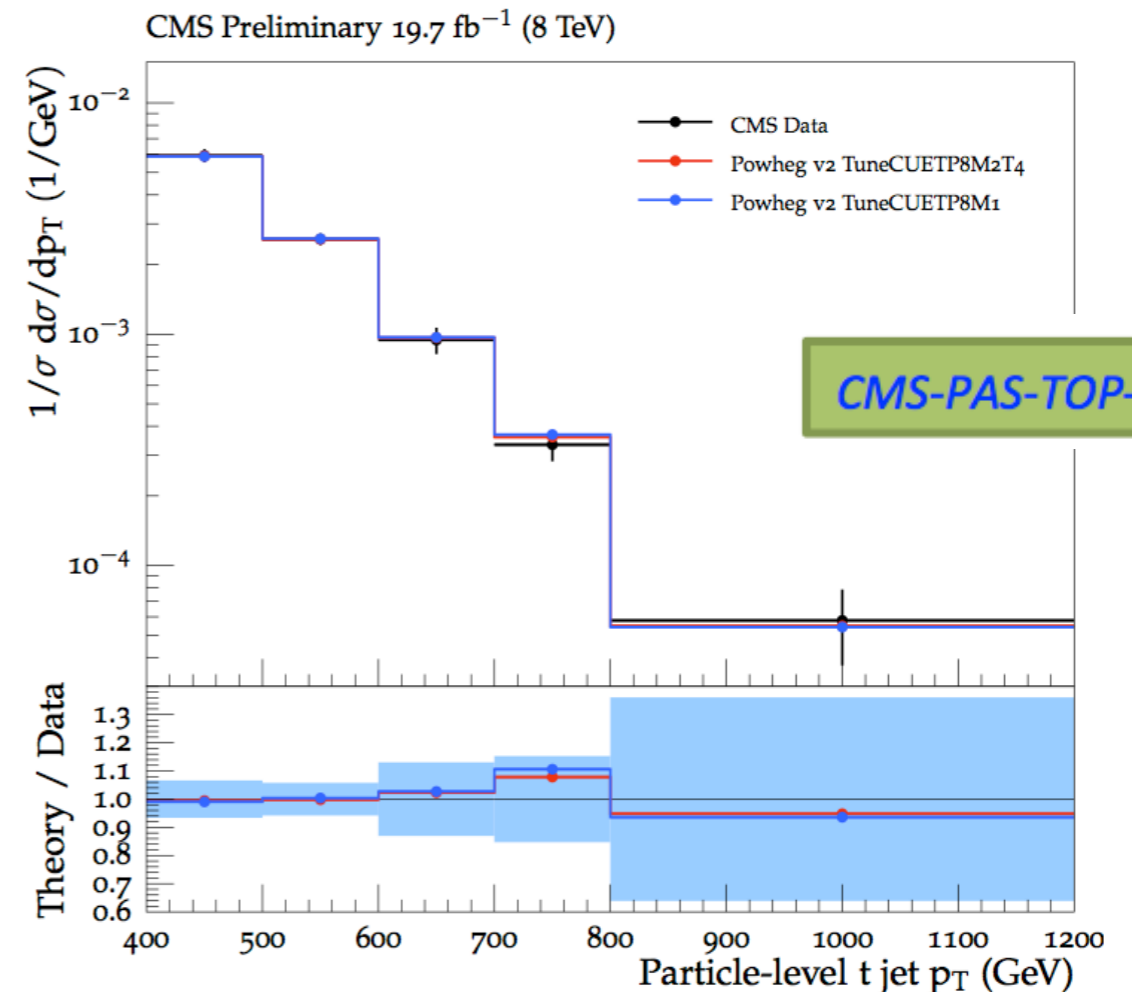
- Boosted top quarks provide first simplification
- Decay products contained in a “fat” jet



(a) $t \rightarrow Wb$

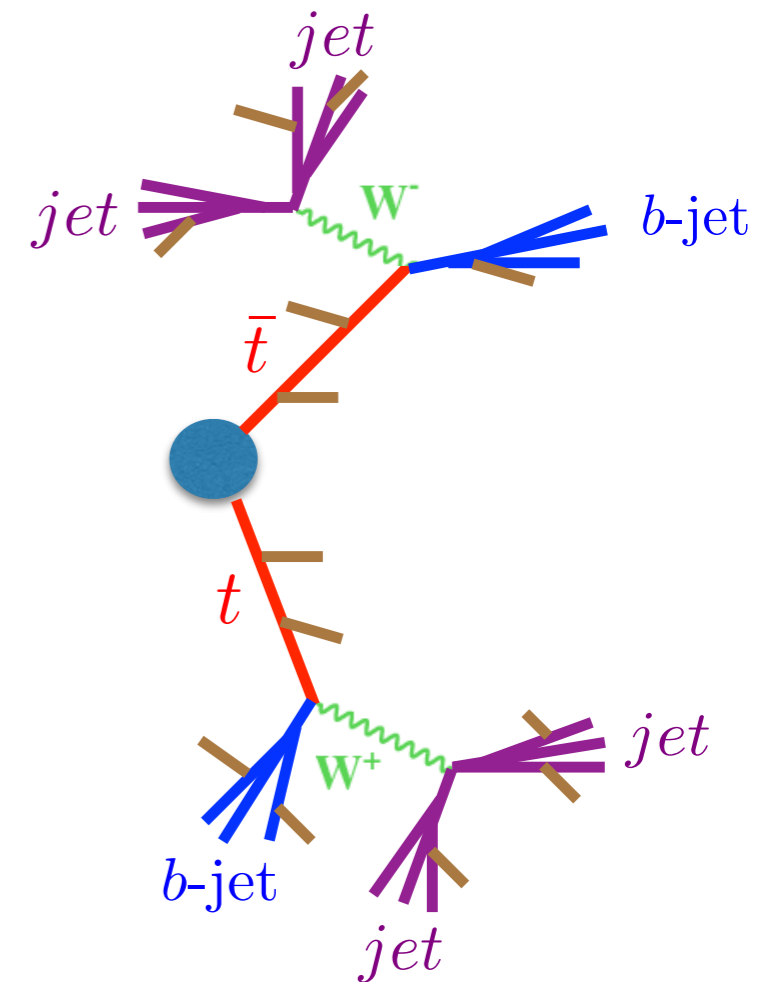
Rule of thumb for decay products:

$$\Delta R(W, b) \sim \frac{2m_t}{p_T}$$



Theory Issues at Hadron Collider

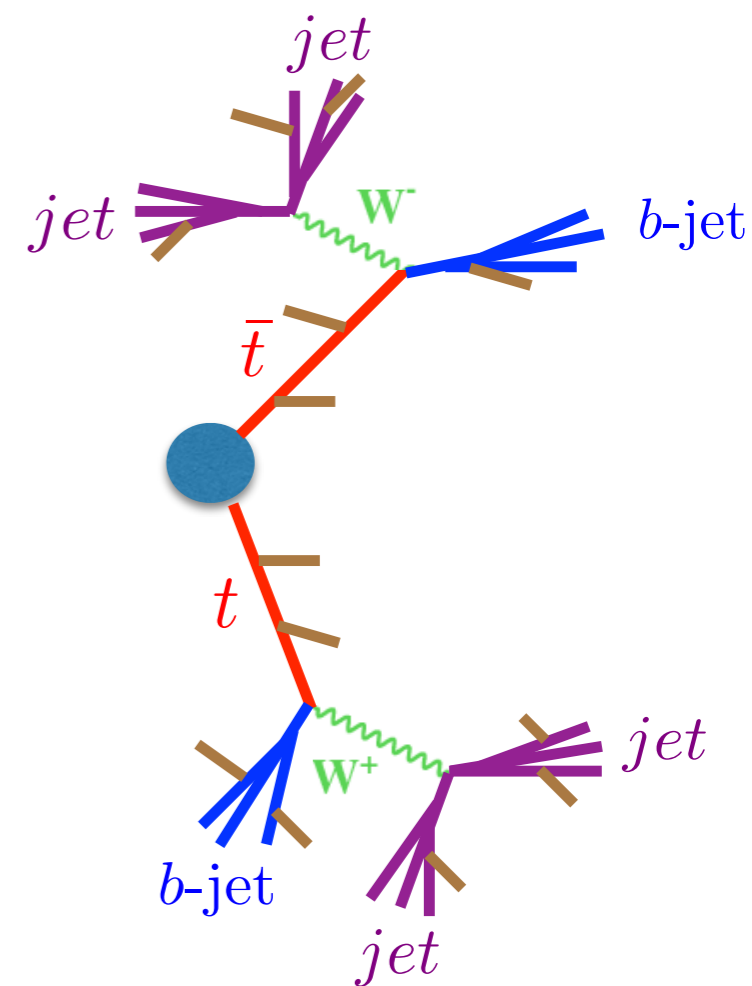
- Jet observable
- Clear relation to top mass scheme
- Initial state radiation
- Final state radiation
- Beam remnant
- Parton distribution functions
- Color reconnection
- Underlying events, pile up
- Summing large logs



Theory Issues at Hadron Collider

- Jet observable ★
- Clear relation to top mass scheme ★
- Initial state radiation
- Final state radiation ★
- Beam remnant
- Parton distribution functions
- Color reconnection ★
- Underlying events, pile up
- Summing large logs ★

First
 $e^+e^- \rightarrow t\bar{t}X$
and the issues ★



Top Mass From Jet Distributions

(Fleming, Hoang, SM, Stewart)

$$e^+ e^- \rightarrow t \bar{t} X$$

- Top jet hemisphere mass distribution sensitive to top mass:

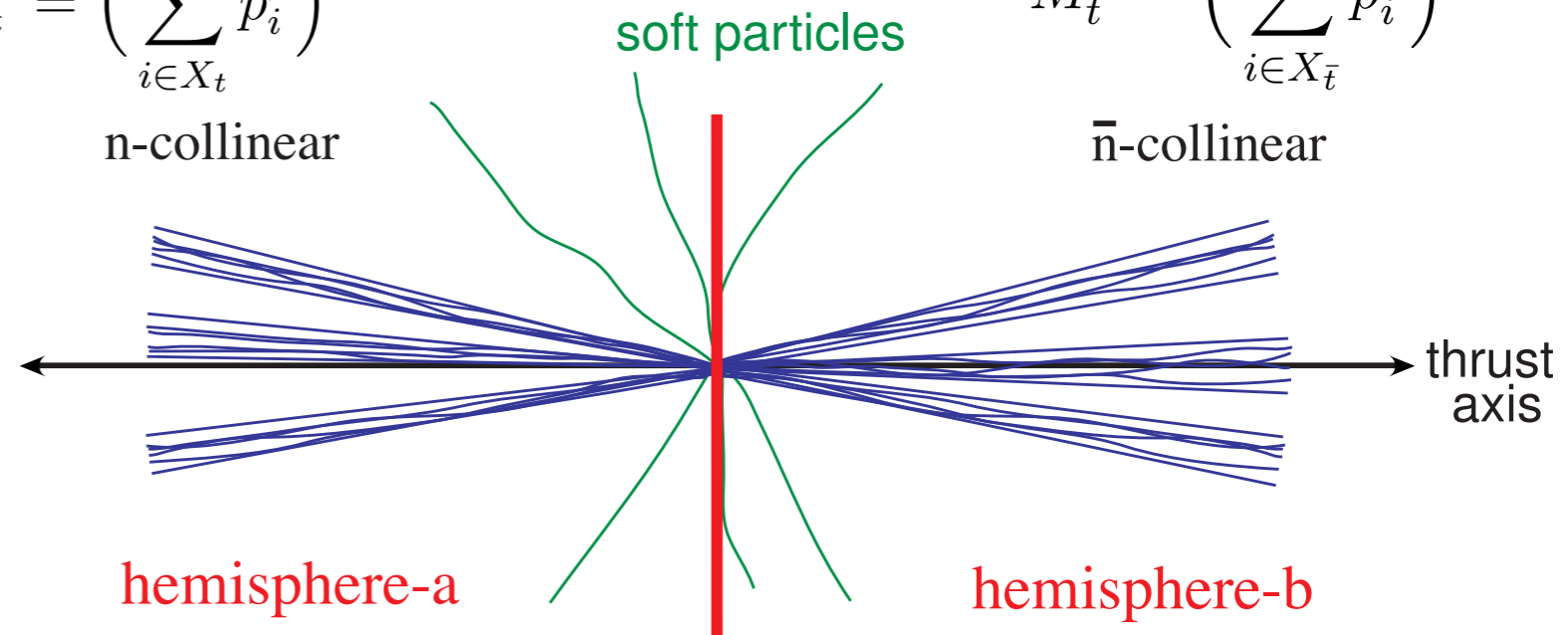
$$\frac{d^2 \sigma}{dM_t^2 dM_{\bar{t}}^2}$$

$$M_t^2 = \left(\sum_{i \in X_t} p_i^\mu \right)^2$$

n-collinear

$$M_{\bar{t}}^2 = \left(\sum_{i \in X_{\bar{t}}} p_i^\mu \right)^2$$

\bar{n} -collinear



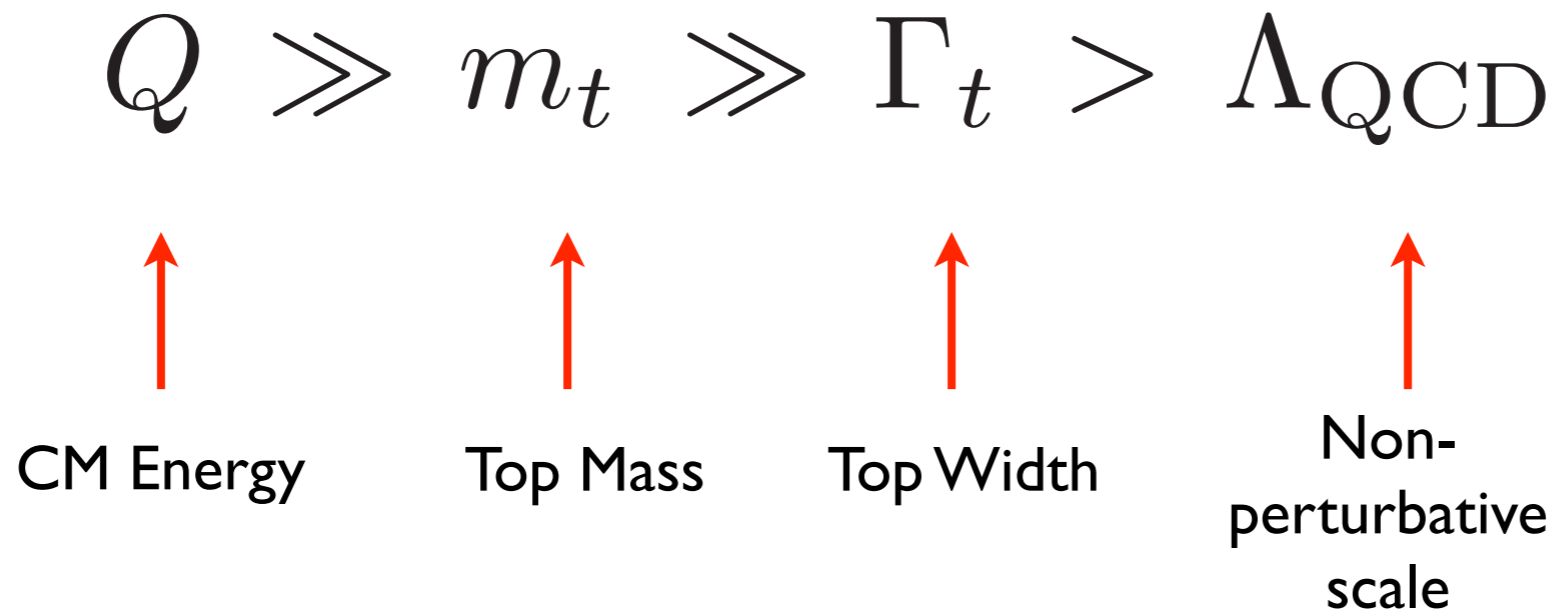
- Peak region:

$$\frac{M_{t,\bar{t}}^2 - m^2}{m} \sim \Gamma \ll m$$

Top Mass From Jet Distributions

$$e^+e^- \rightarrow t\bar{t}X$$

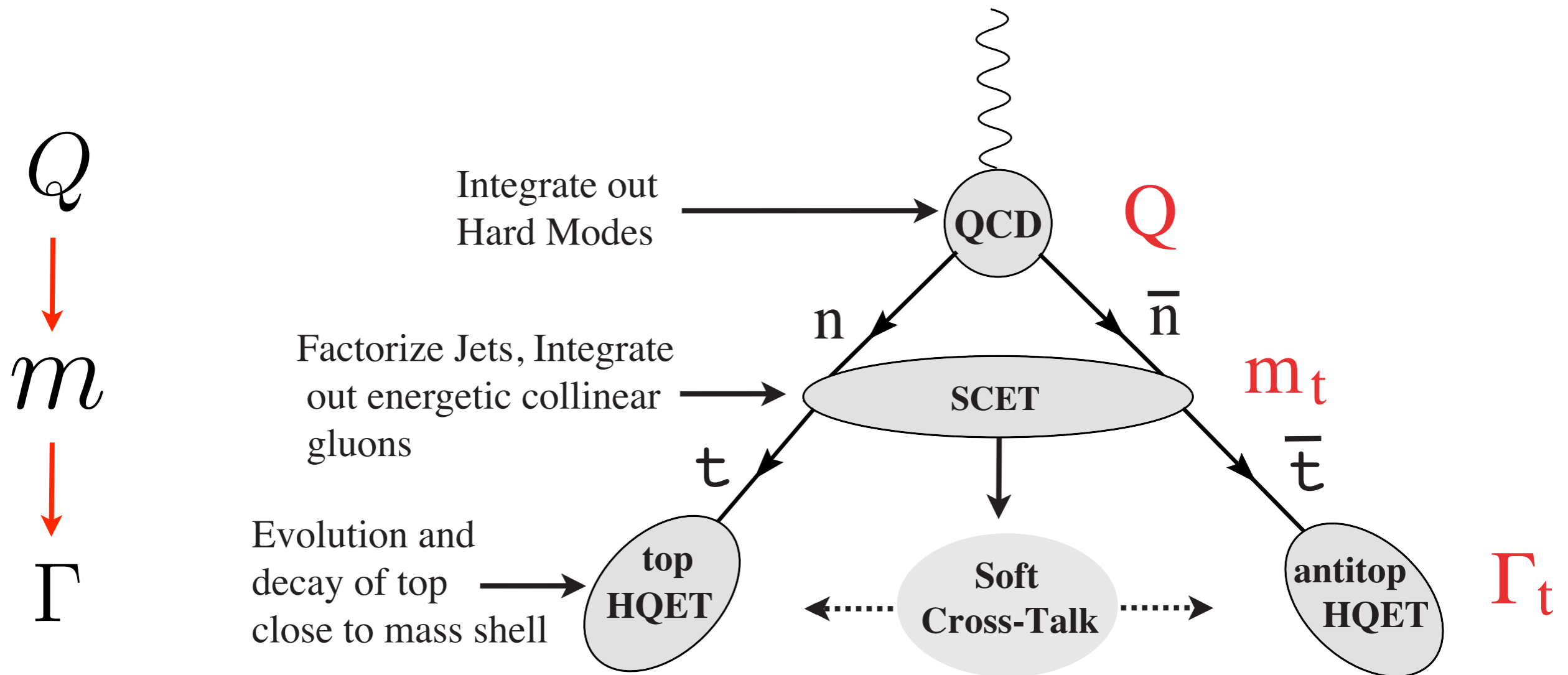
Boosted top quark pair-production:

$$Q \gg m_t \gg \Gamma_t > \Lambda_{\text{QCD}}$$


CM Energy Top Mass Top Width Non-perturbative scale

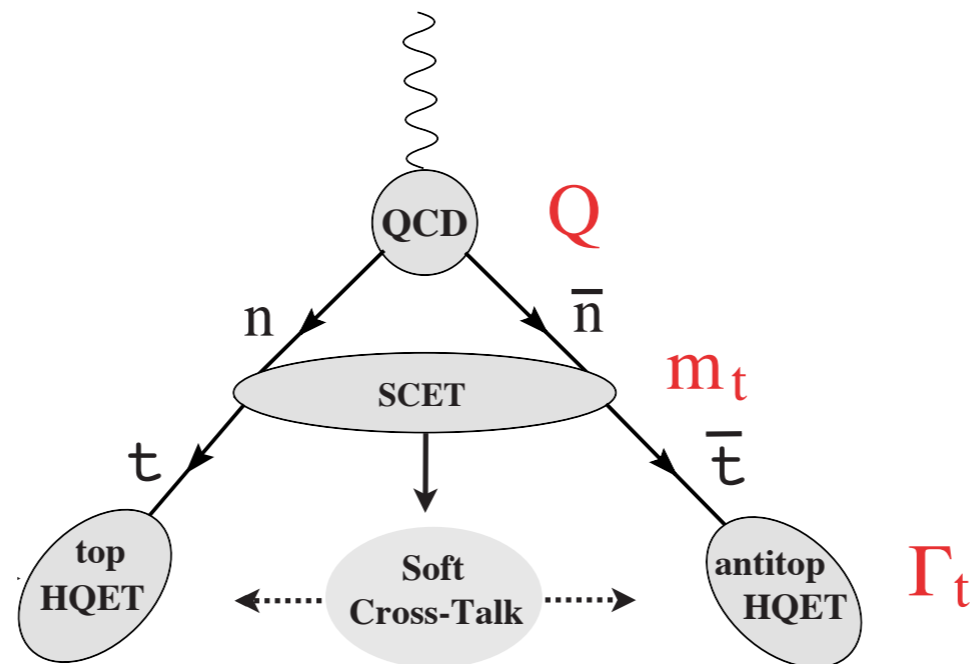
Four scales in the problem!

A Symphony of Effective Field Theories



$$Q \gg m_t \gg \Gamma_t > \Lambda_{\text{QCD}}$$

Factorization Formula



Hard modes
integrated out



“Hard” collinear modes
integrated out



Evolution and decay
of top quarks close
to mass shell



Soft cross-talk



$$\frac{d\sigma}{dM_t^2 dM_{\bar{t}}^2} = \sigma_0 H_Q(Q, \mu_m) H_m\left(m_J, \frac{Q}{m_J}, \mu_m, \mu\right) \times \int dl^+ dl^- B_+\left(\hat{s}_t - \frac{Ql^+}{m_J}, \Gamma_t, \mu\right) B_-\left(\hat{s}_{\bar{t}} - \frac{Ql^-}{m_J}, \Gamma_t, \mu\right) S(l^+, l^-, \mu)$$

Top Mass From Jet Distributions

$$\frac{d\sigma}{dM_t^2 dM_{\bar{t}}^2} = \sigma_0 H_Q(Q, \mu_m) H_m\left(m_J, \frac{Q}{m_J}, \mu_m, \mu\right) \\ \times \int d\ell^+ d\ell^- B_+\left(\hat{s}_t - \frac{Q\ell^+}{m_J}, \Gamma_t, \mu\right) B_-\left(\hat{s}_{\bar{t}} - \frac{Q\ell^-}{m_J}, \Gamma_t, \mu\right) S(\ell^+, \ell^-, \mu)$$

- HQET Lagrangian determines dynamics of top-jet functions B:

$$\mathcal{L}_{\pm} = \bar{h}_{v_{\pm}} \left(i v_{\pm} \cdot D_{\pm} - \delta m + \frac{i}{2} \Gamma_t \right) h_{v_{\pm}} \quad \delta m = m_{\text{pole}} - m$$

- EFT power counting defines the “top resonance” mass schemes:

$$\delta m \sim \hat{s}_t \sim \hat{s}_{\bar{t}} \sim \Gamma$$

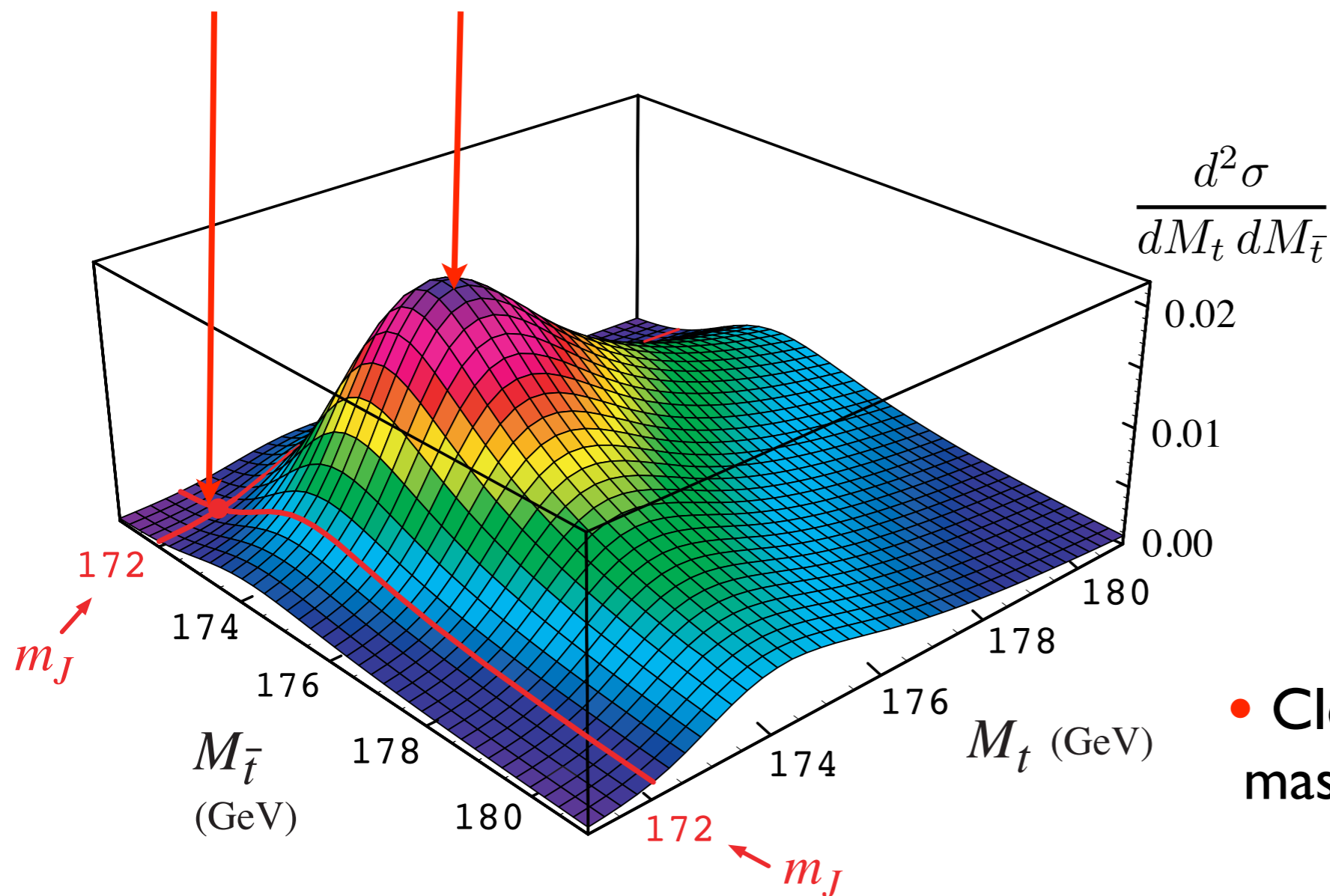
- Note power counting of EFT breaks down for MSbar mass:

$$\delta \bar{m} \sim \alpha_s \bar{m} \gg \Gamma$$

Top Mass From Jet Distributions

$$\frac{d\sigma}{dM_t^2 dM_{\bar{t}}^2} = \sigma_0 H_Q(Q, \mu_m) H_m\left(m_J, \frac{Q}{m_J}, \mu_m, \mu\right) \\ \times \int d\ell^+ d\ell^- B_+\left(\hat{s}_t - \frac{Q\ell^+}{m_J}, \Gamma_t, \mu\right) B_-\left(\hat{s}_{\bar{t}} - \frac{Q\ell^-}{m_J}, \Gamma_t, \mu\right) S(\ell^+, \ell^-, \mu)$$

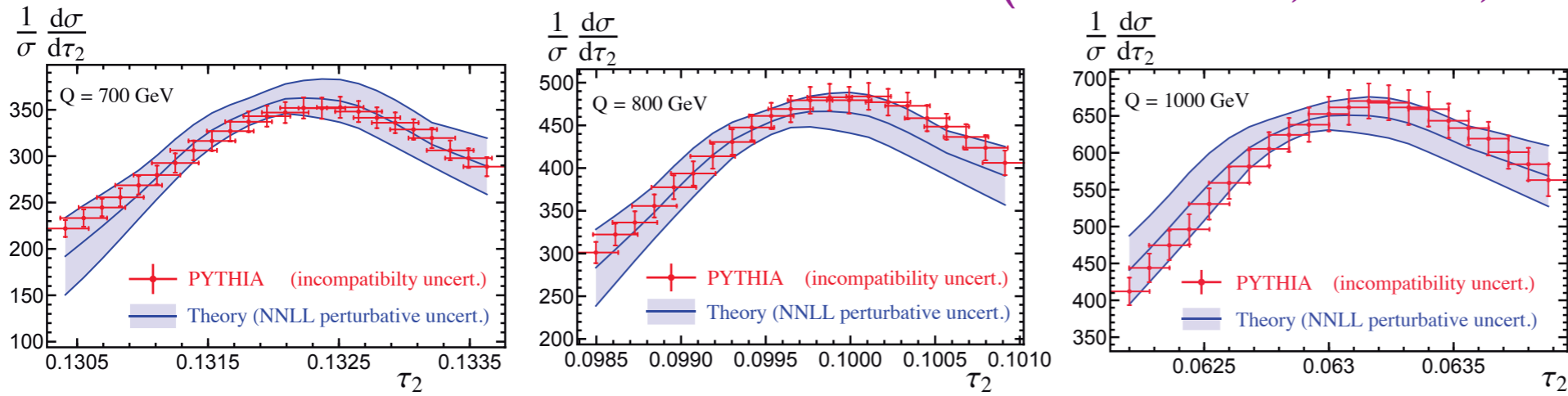
$$m_J = M_{\text{peak}} - \Gamma(\alpha_s + \dots) - \frac{Q\Lambda_{QCD}}{m}$$



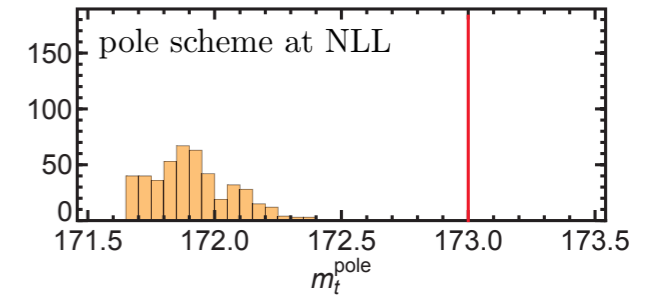
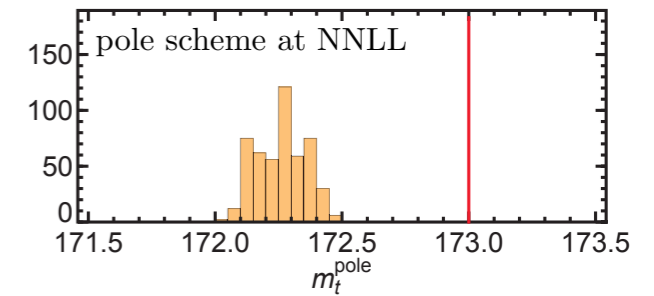
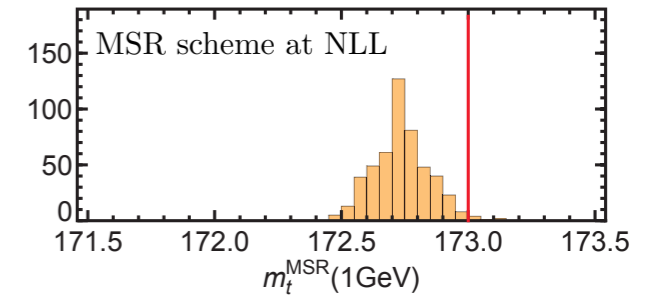
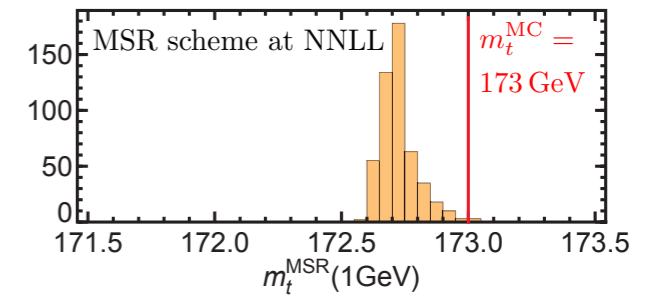
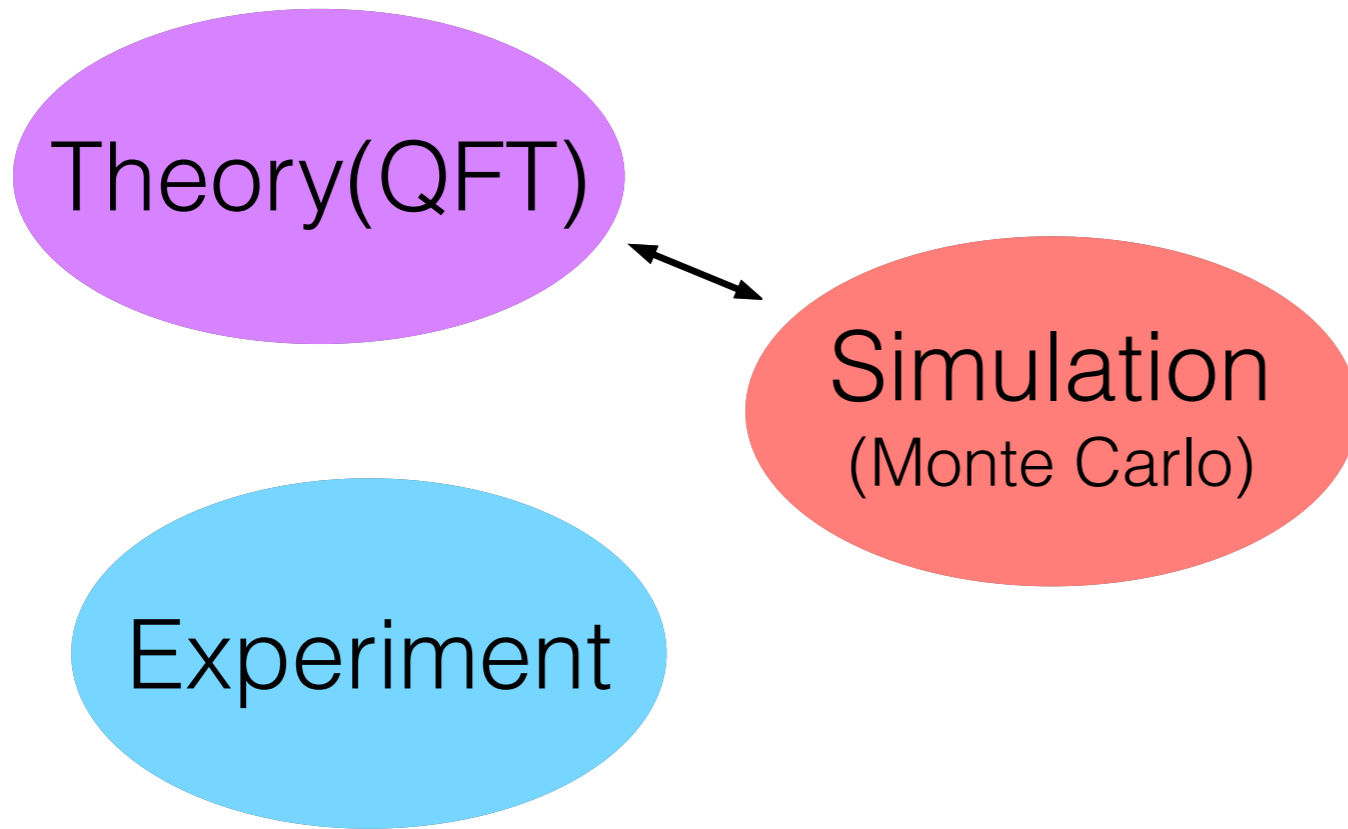
- Clear relation between top mass and peak position

Top Mass Calibration of Monte Carlo

(Butenschoen, Dehnadi, Hoang, Mateu, Preisser, Stewart)

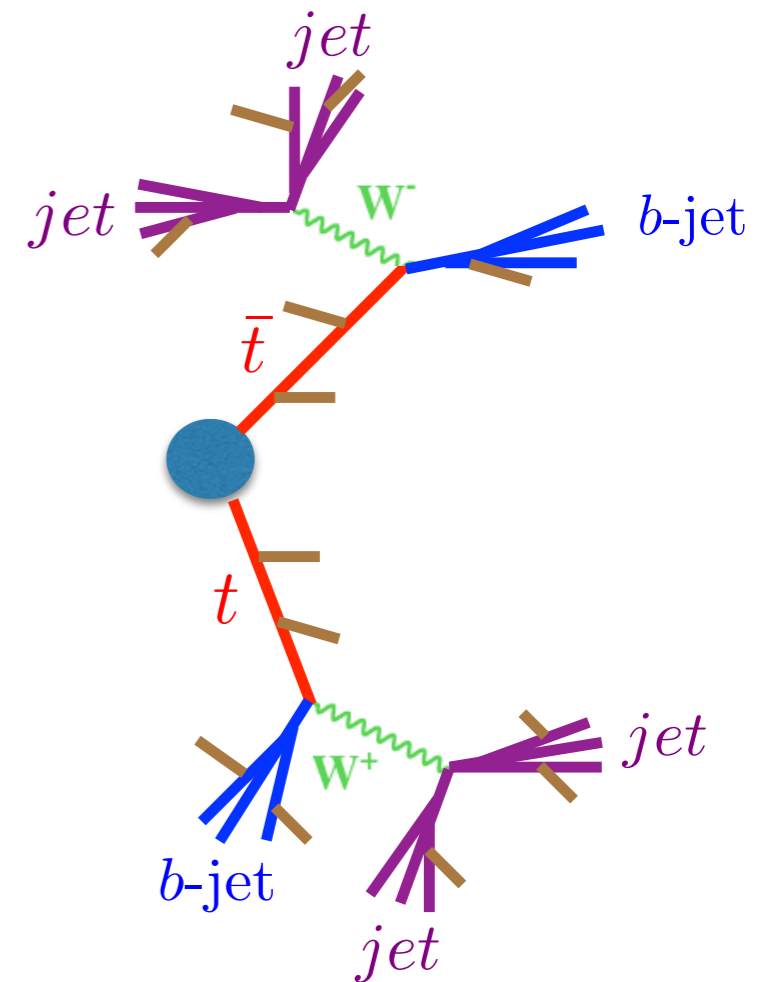


- Monte Carlo top mass is found to be closer to the MSR mass than the pole mass.



Theory Issues at Hadron Collider

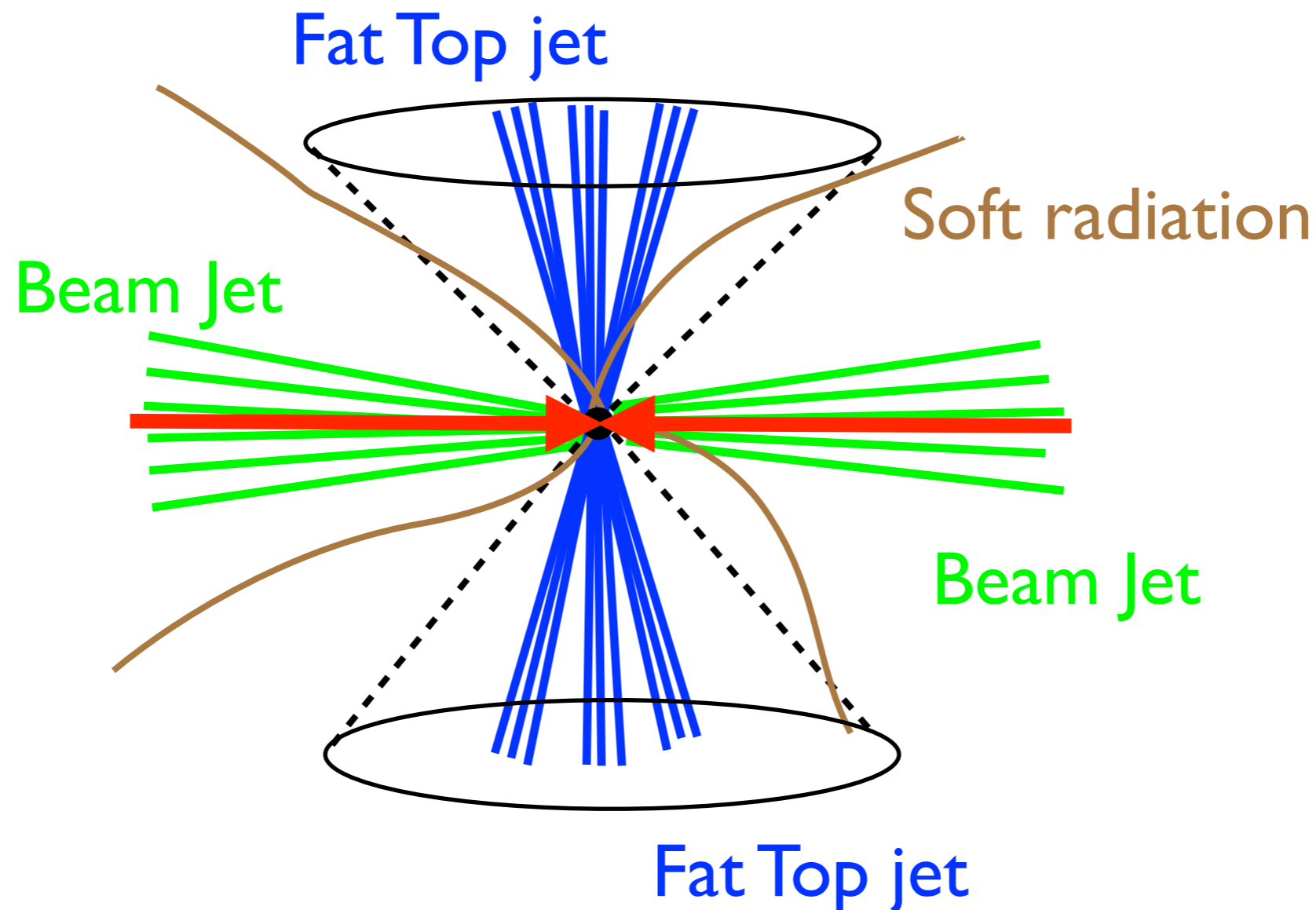
- Jet observable
- Clear relation to top mass scheme
- Initial state radiation
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Factorization for Boosted Tops at Hadron Colliders

(Hoang, SM, Pathak, Stewart)

- Lepton collider methods can be extended to hadron colliders.
- Make use of the 2-Jettiness event shape. (Stewart, Tackmann, Waalewijn)



$$\frac{d^2\sigma}{dM_{J_1}^2 dM_{J_2}^2 d\mathcal{T}^{\text{cut}}} = \text{tr} \left[\hat{H}_{Q_m} \hat{S}(\mathcal{T}^{\text{cut}}, R, \dots) \otimes F \right] \otimes J_B \otimes J_B \otimes \mathcal{I} \otimes f f$$

Factorization for Boosted Tops at Hadron Colliders

(Hoang, SM, Pathak, Stewart)

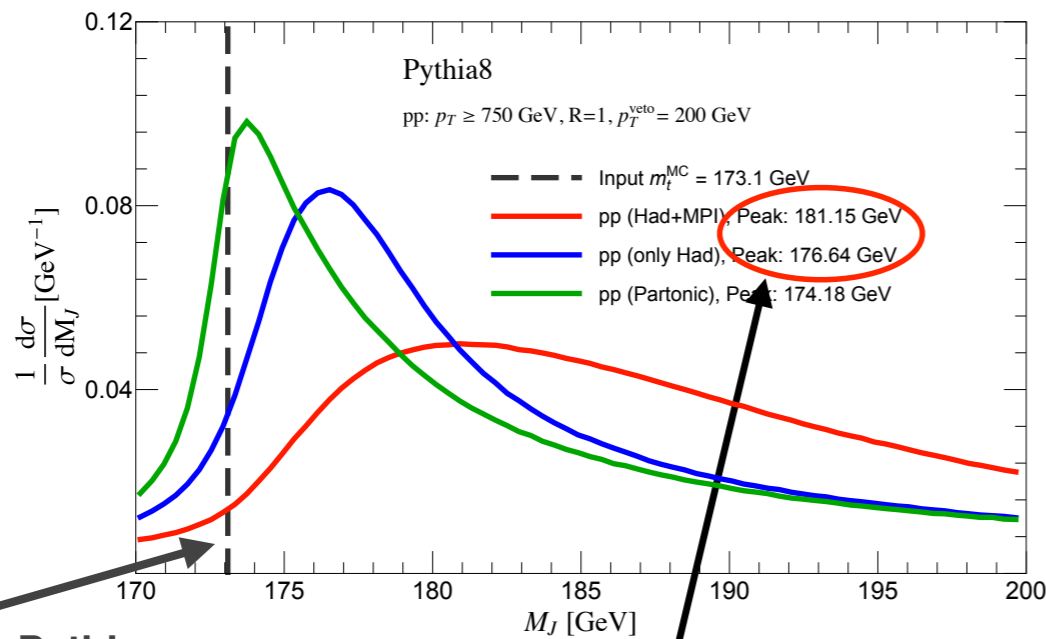
$$\frac{d^2\sigma}{dM_{J_1}^2 dM_{J_2}^2 d\mathcal{T}^{\text{cut}}} = \text{tr} \left[\hat{H}_{Qm} \hat{S}(\mathcal{T}^{\text{cut}}, R, \dots) \otimes F \right] \otimes J_B \otimes J_B \otimes \mathcal{I} \otimes f f$$

Hard function
Soft function
Top HQET jet functions
ISR
PDFs

Top HQET jet functions
(identical to lepton collider!)

- Jet mass spectrum is quite sensitive to contamination:

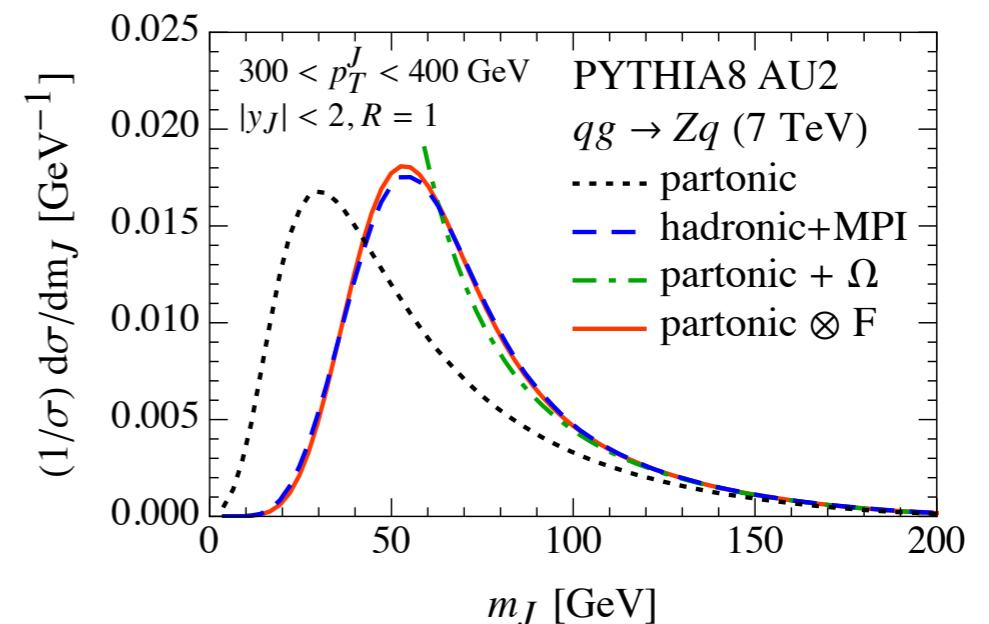
Effect of UE/MPI



Significant contamination

- Same soft model for hadronization can describe UE

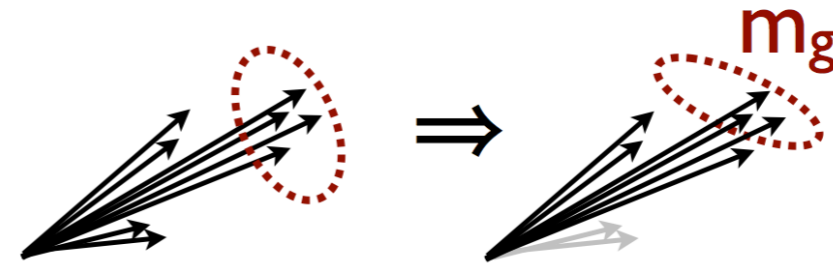
(Stewart, Tackmann, Waalewijn)



Soft Drop

(Larkoski, Marzani, Soyez, Thaler)

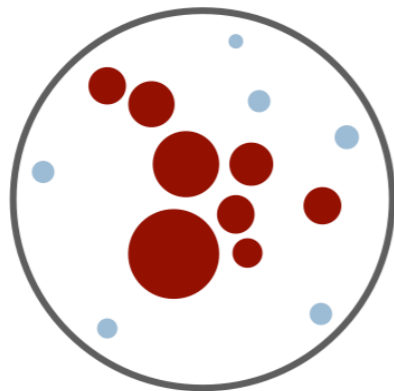
- Soft drop grooming reduces sensitivity for jet mass spectrum to soft contamination.



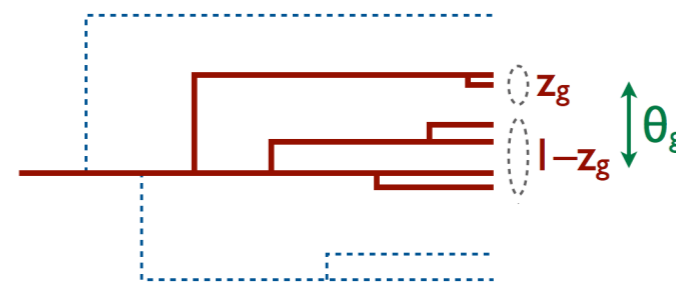
$$\frac{\min(p_{Ti}, p_{Tj})}{p_{Ti} + p_{Tj}} > z_{\text{cut}} \left(\frac{\Delta R_{ij}}{R_0} \right)^\beta$$

$$z > z_{\text{cut}} \theta^\beta$$

Groomed jet

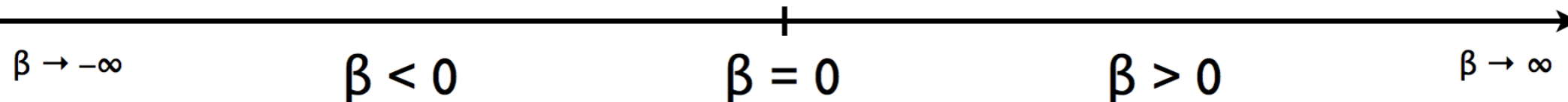


Groomed Clustering tree



More Grooming

Less Grooming

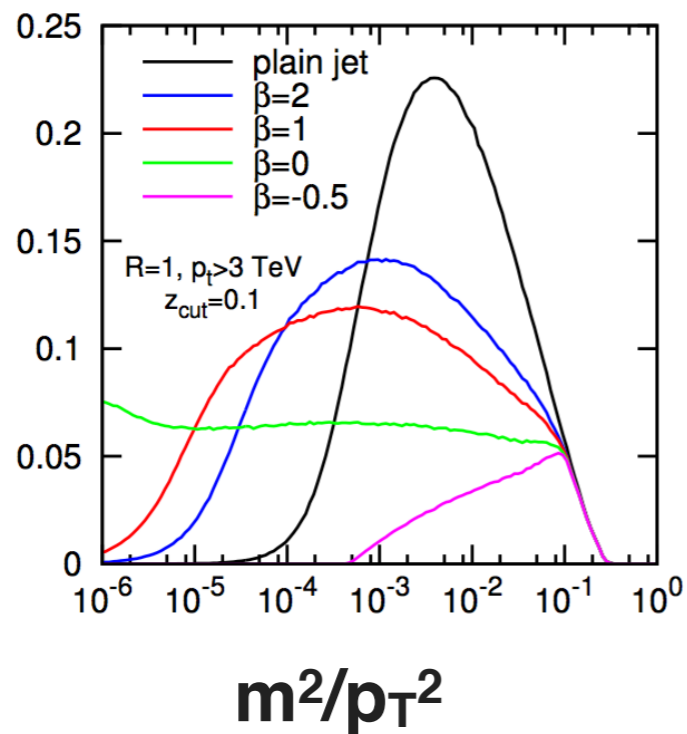


Soft Drop

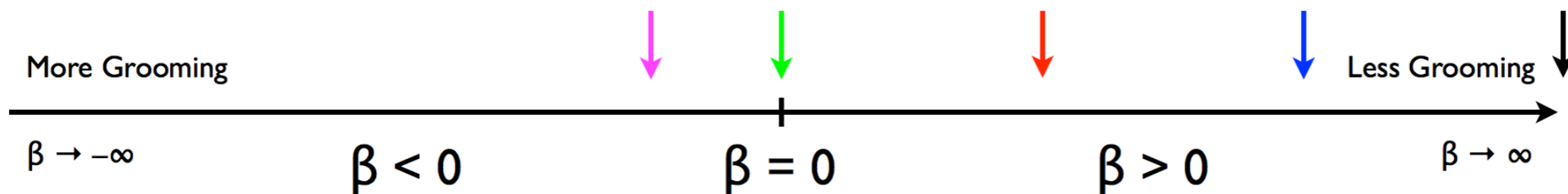
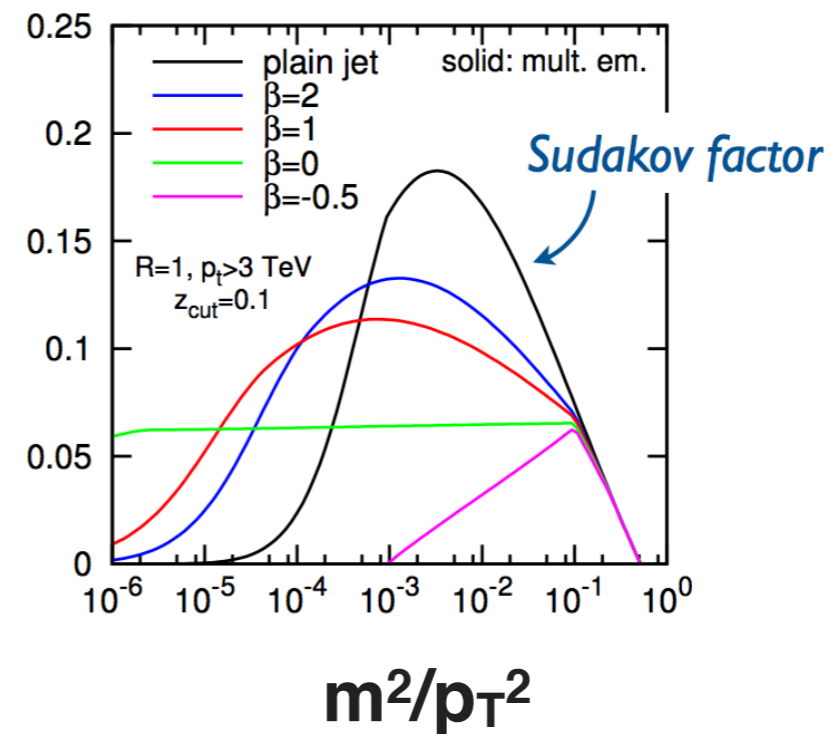
(Larkoski, Marzani, Soyez, Thaler)

- Jet mass spectrum for massless partons:

Pythia8, partonic

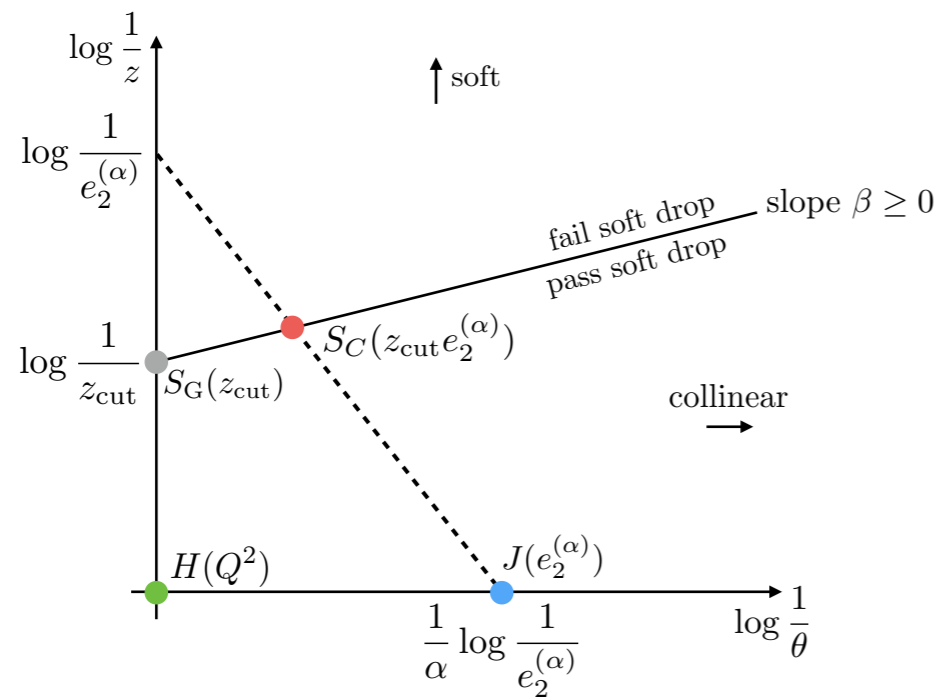
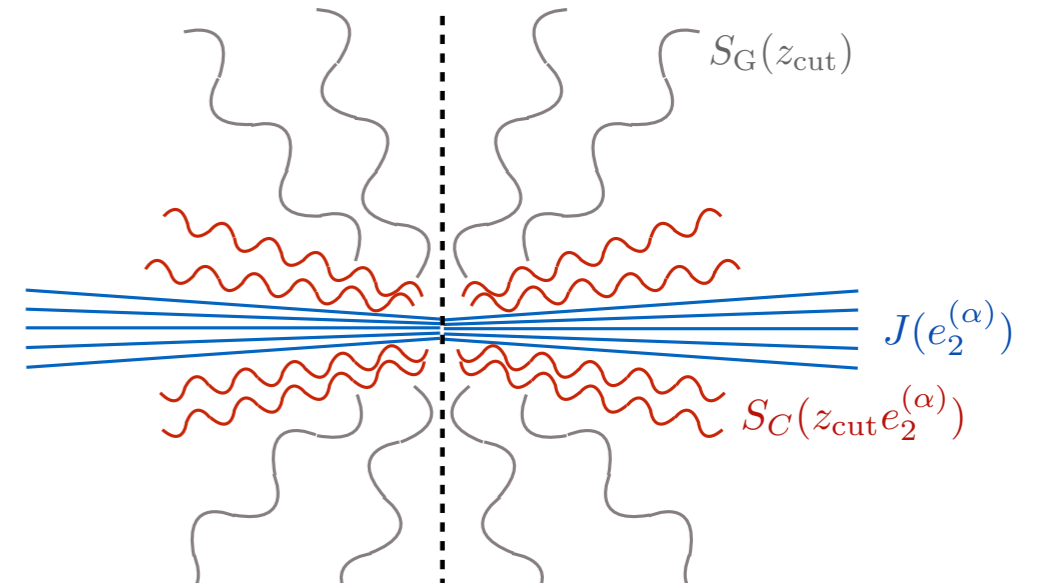
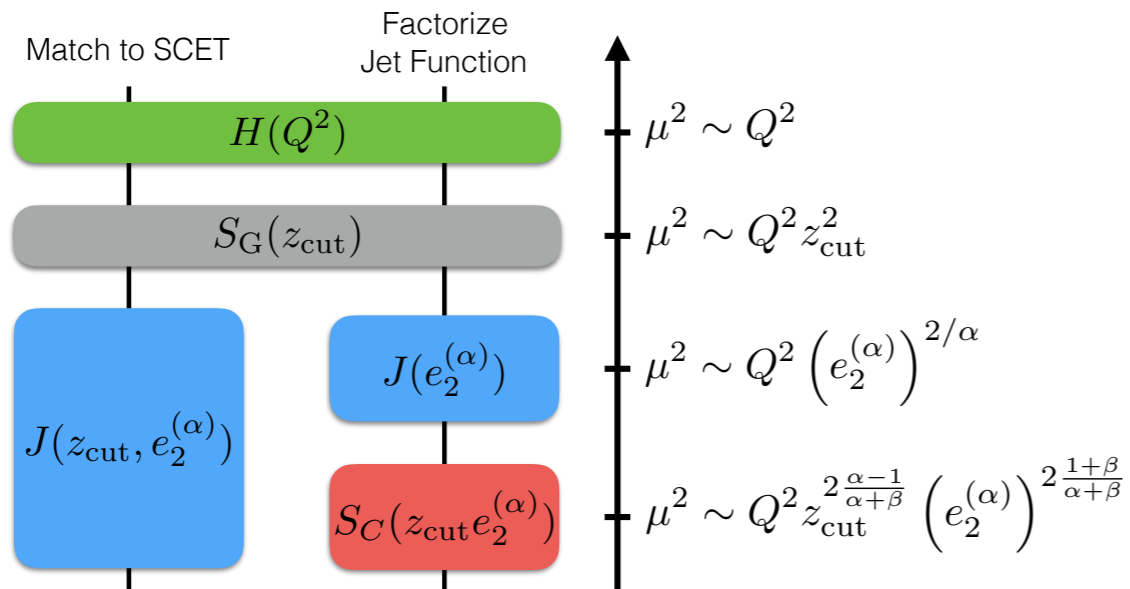


Pert QCD at \sim NLL



Soft Drop Factorization

(Fyre, Larkoski, Marzani, Schwartz, Yan)



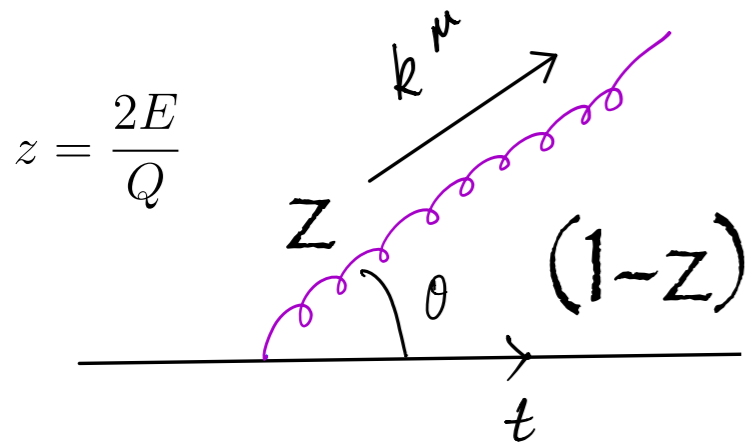
- Factorization

$$\frac{d^2\sigma}{de_{2,L}^{(\alpha)} de_{2,R}^{(\alpha)}} = H(Q^2) S_G(z_{\text{cut}}) \left[S_C(z_{\text{cut}} e_{2,L}^{(\alpha)}) \otimes J(e_{2,L}^{(\alpha)}) \right] \left[S_C(z_{\text{cut}} e_{2,R}^{(\alpha)}) \otimes J(e_{2,R}^{(\alpha)}) \right]$$

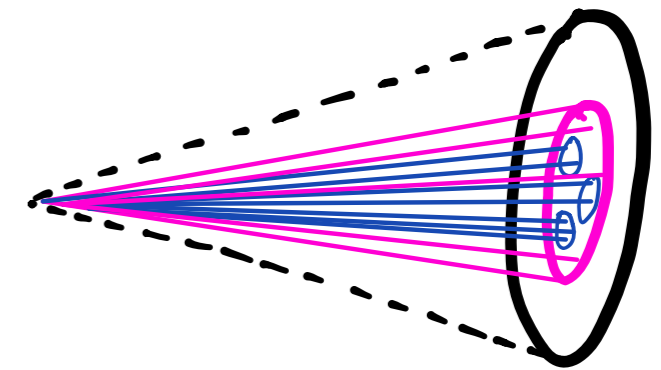
- Wide angle soft radiation groomed away
- No non-global logs

Top Jet Mass with Soft Drop

Top Jet Mass Constraints with Soft Drop



$$k_j^\mu = (k^+, k^-, k_\perp) = (E(1 - \cos \theta), E(1 + \cos \theta), k_\perp)$$



- Top jet mass constraint in the peak region:

$$\frac{M_{t,\bar{t}}^2 - m^2}{m} \sim \Gamma \ll m \quad \longrightarrow \quad z \left[(1 - \cos \theta) + \frac{m^2}{Q^2} (1 + \cos \theta) \right] \sim \frac{2m\Gamma_t}{Q^2}$$

- Soft Drop constraint:

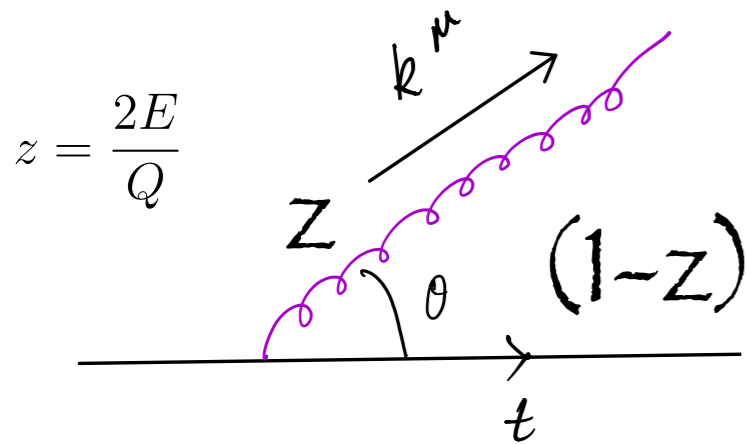
$$z > z_{\text{cut}} \theta^\beta \quad \longrightarrow \quad \left(\frac{Q}{2m} \right)^\beta \gg \frac{\Gamma_t}{m_t} \left(\frac{Q}{2m} \right)^\beta \gg z_{\text{cut}} \gg \frac{2m_t\Gamma_t}{Q^2}$$

keep decay products

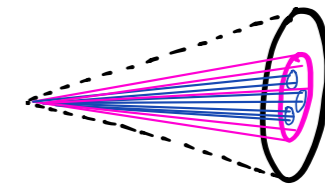
keep ucollinear radiation that satisfies jet mass constraint

removes usoft radiation that satisfies jet mass constraint (usoft contamination)

Soft Drop Constraints: Light Grooming



$$k_j^\mu = (k^+, k^-, k_\perp) = (E(1 - \cos \theta), E(1 + \cos \theta), k_\perp)$$



- Soft Drop constraint (soft radiation at wide angles are groomed away):

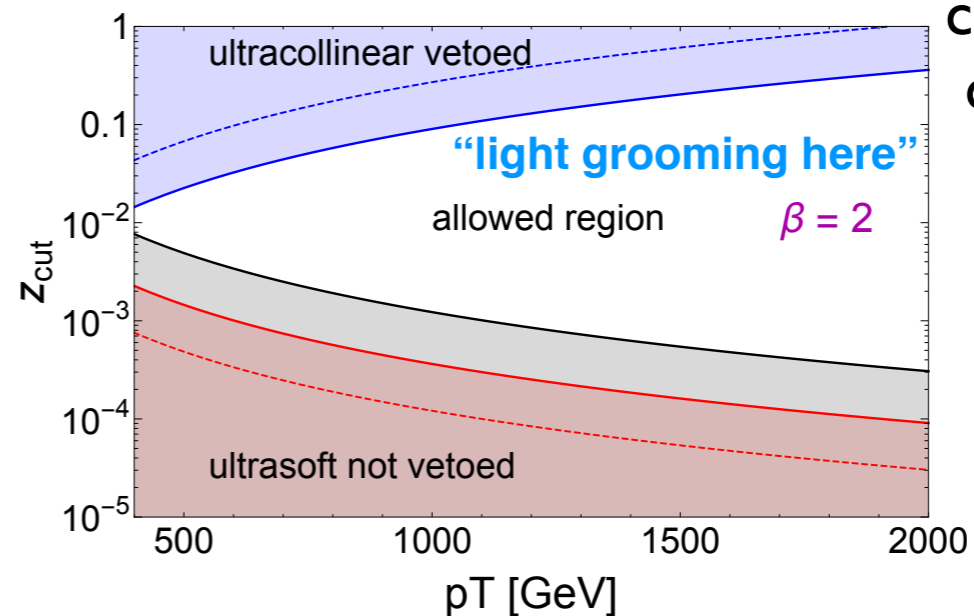
$$z > z_{\text{cut}} \theta^\beta \longrightarrow \left(\frac{Q}{2m}\right)^\beta \gg \frac{\Gamma_t}{m_t} \left(\frac{Q}{2m}\right)^\beta \gg z_{\text{cut}} \gg \frac{2m_t \Gamma_t}{Q^2}$$

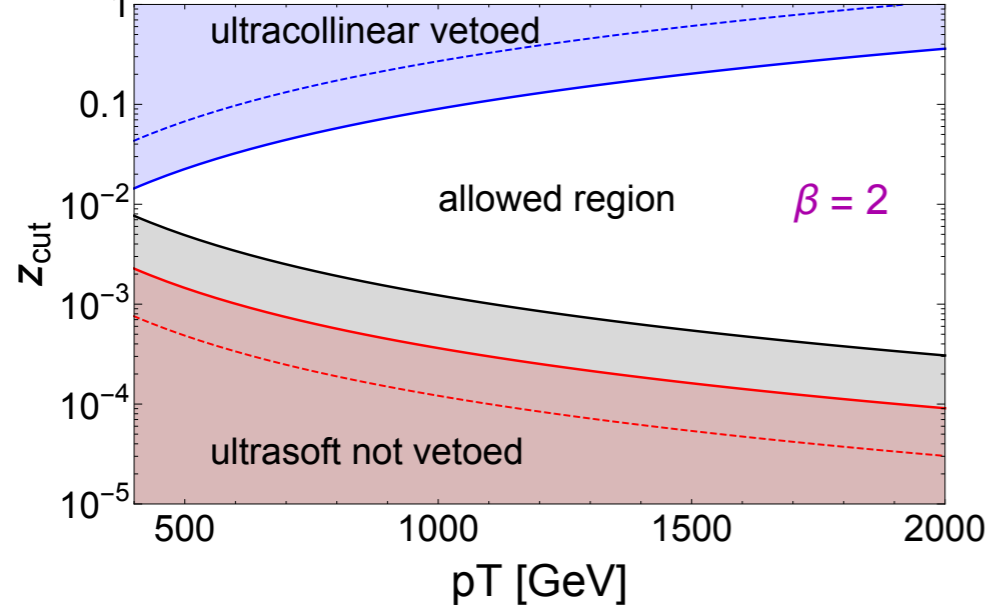
keep decay products

keep ucollinear radiation,
not touch top mass

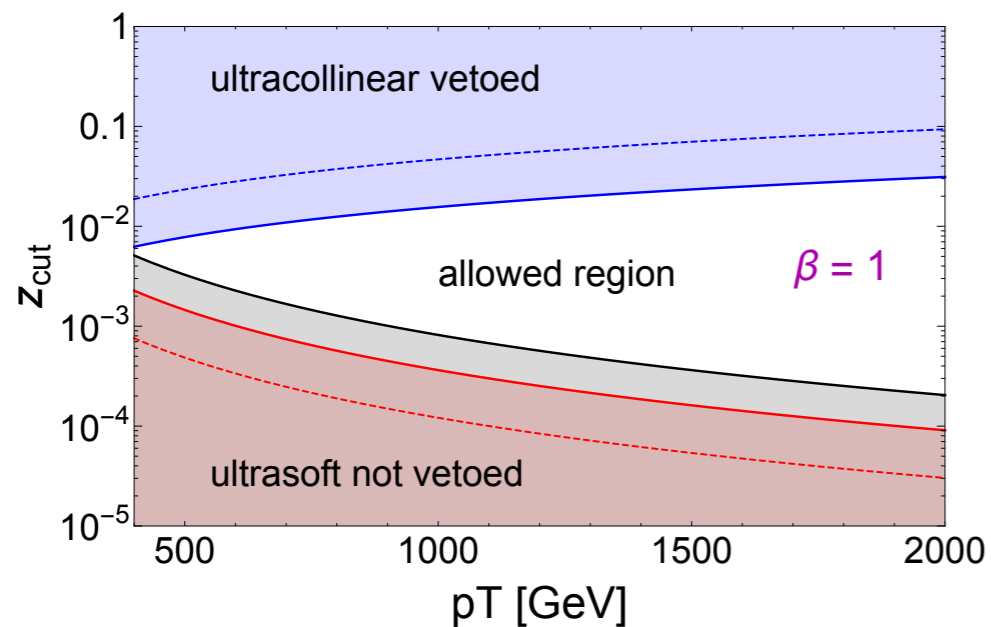
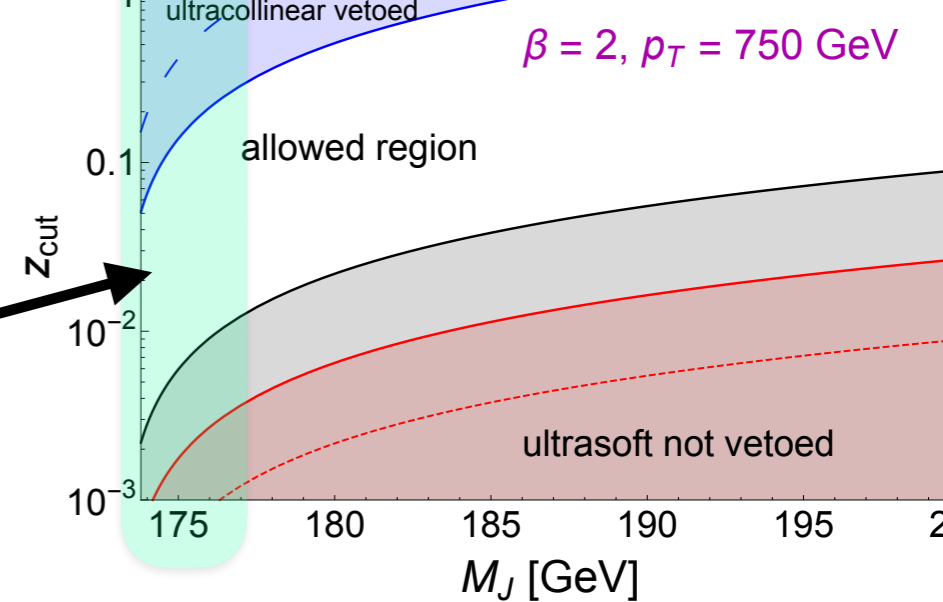
removes usoft radiation
that satisfies jet mass
constraint (usoft
contamination)

- Soft Drop parameter ranges are restricted to “light grooming”:
- Greater allowed Soft Drop parameter range at higher pT:

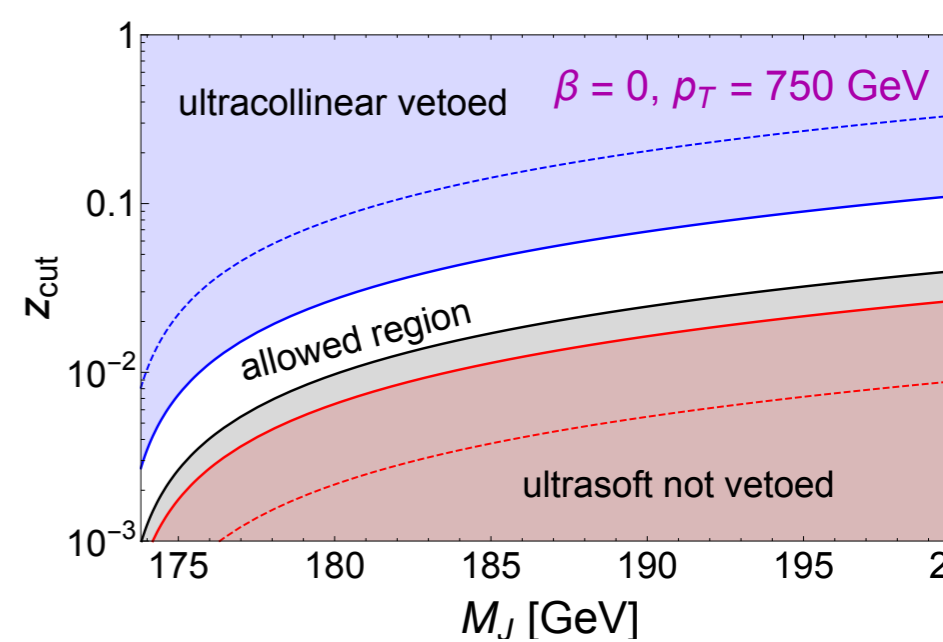
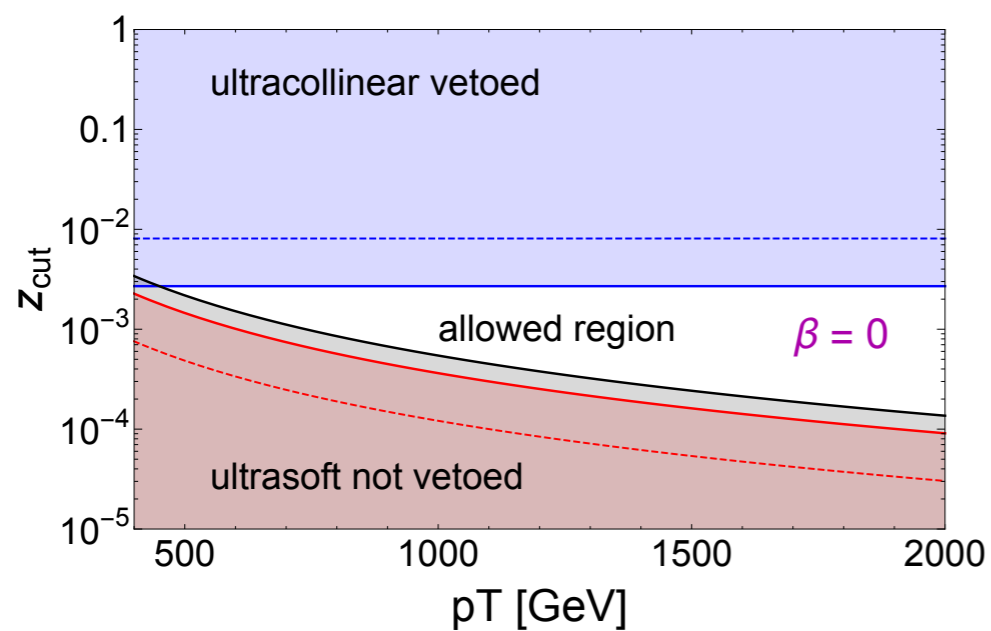
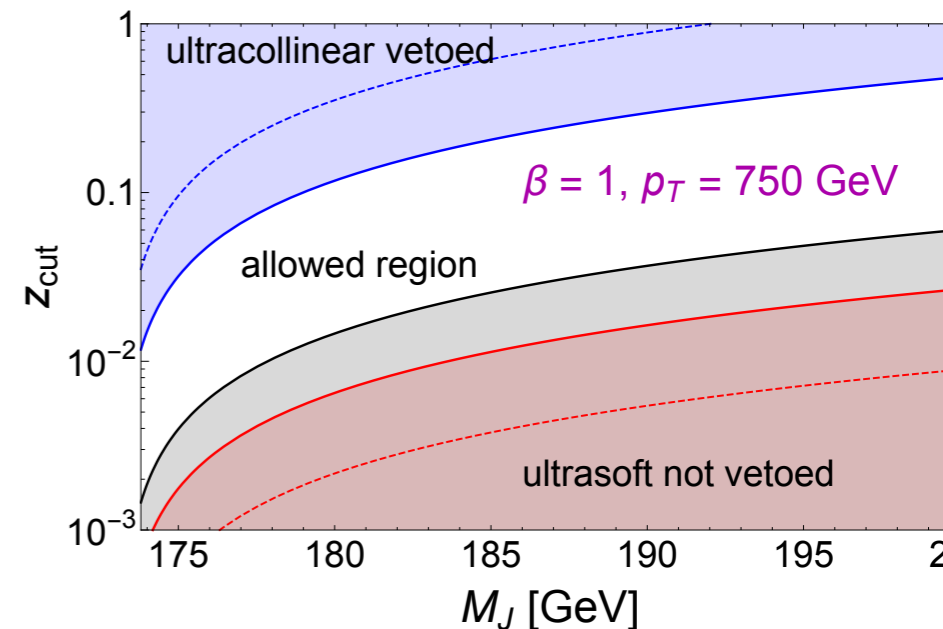




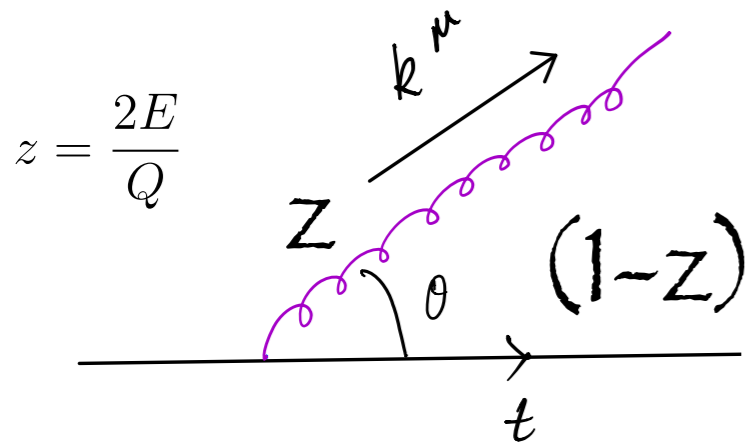
Peak Region



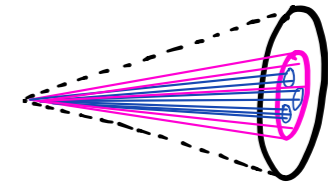
More grooming



Soft Drop Constraints: Collinear Soft and Global Soft Modes



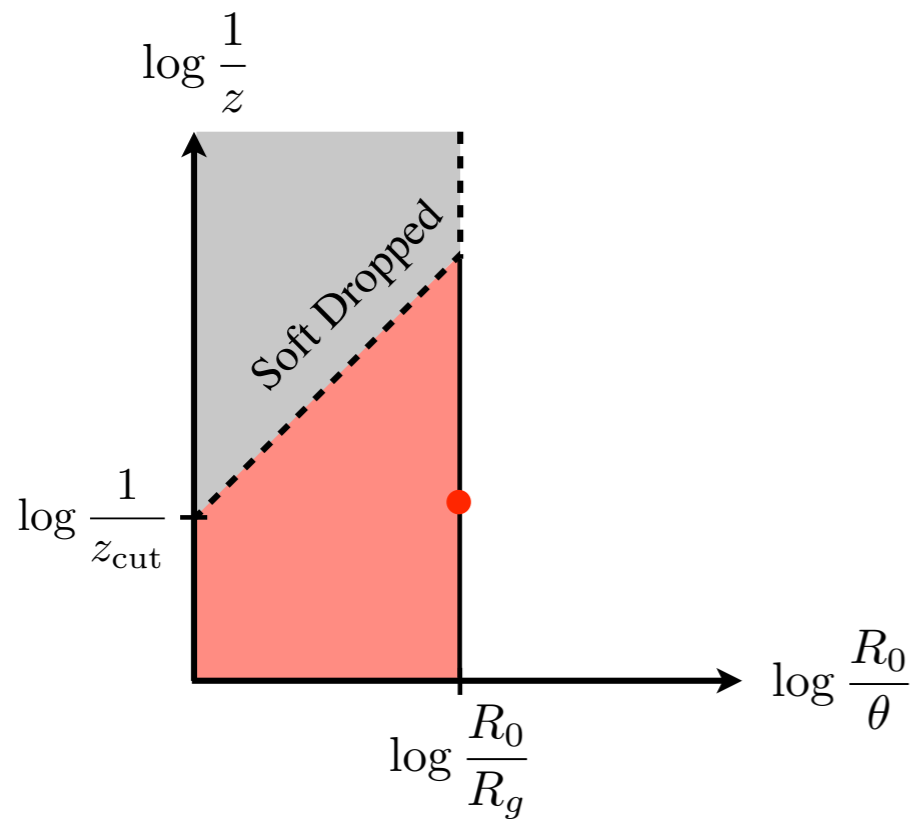
$$k_j^\mu = (k^+, k^-, k_\perp) = (E(1 - \cos \theta), E(1 + \cos \theta), k_\perp)$$



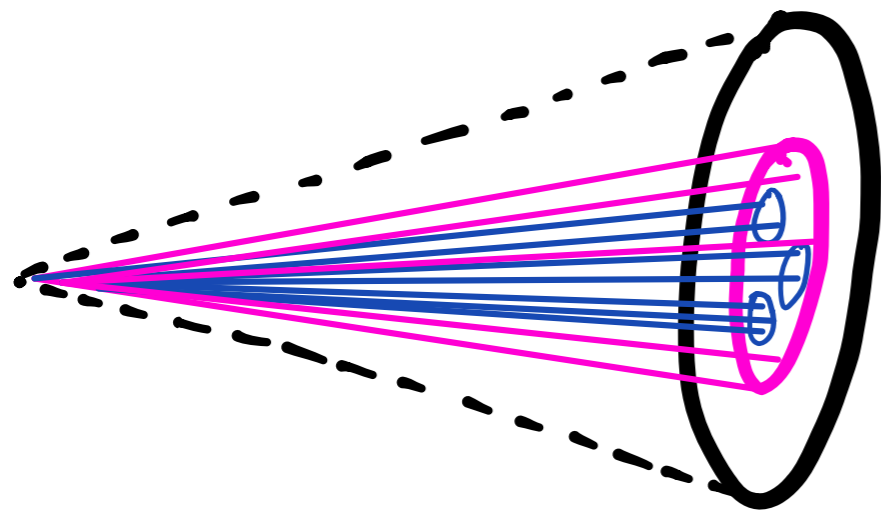
- Ultra-Collinear modes trivially pass soft drop (only constrained by invariant mass of jet):
- Ultra-Soft modes trivially fail soft drop and don't contribute to jet mass.
- A “collinear-soft” mode arises at the boundary that “saturates” soft drop and jet mass constraints.
- “Global-soft” modes at large angles with energy comparable to $\sim Q z_{\text{cut}}$ are also sensitive to the soft drop constraint.

Groomed Jet Radius: R_g

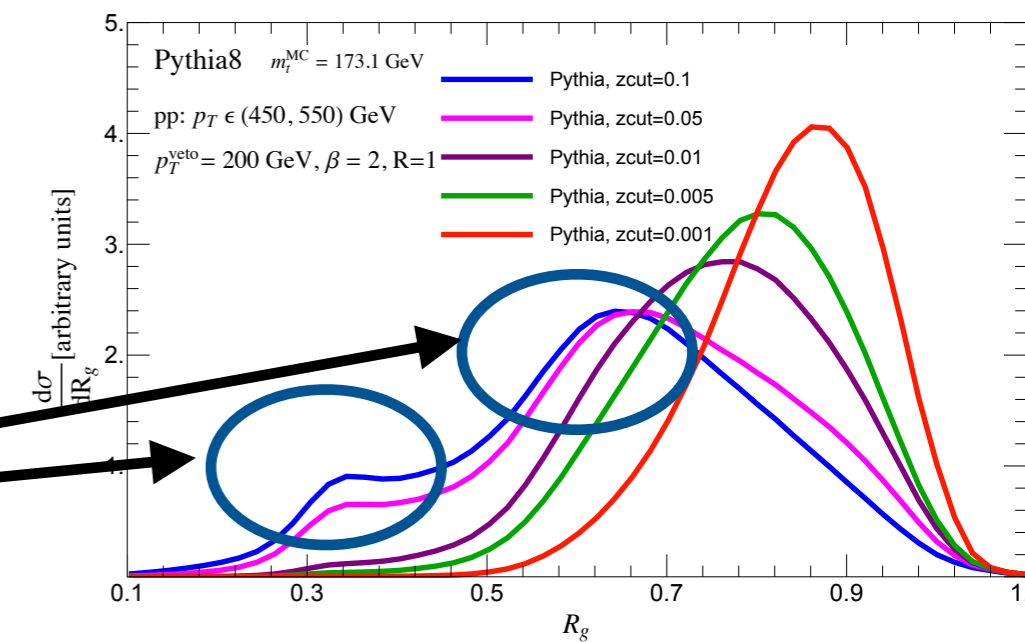
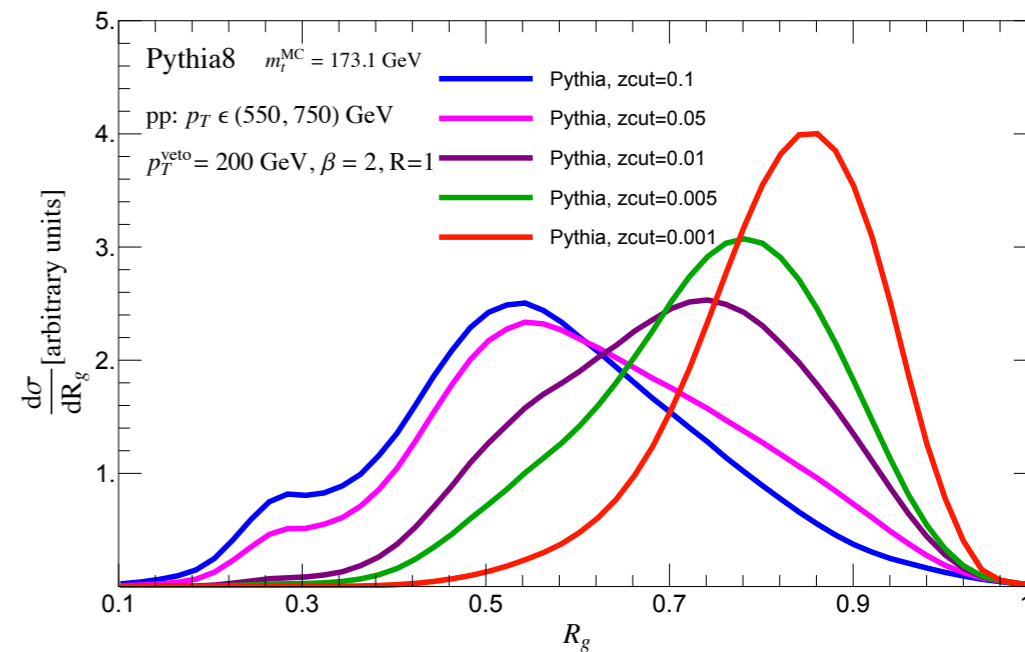
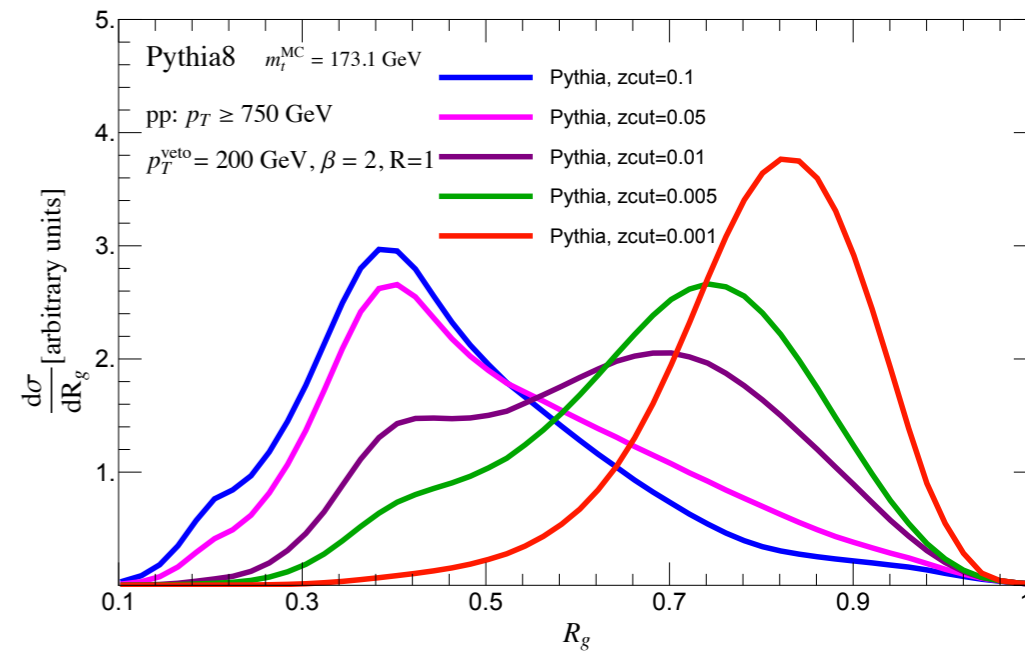
(Larkoski, Marzani, Soyez, Thaler)



Lower
 p_T

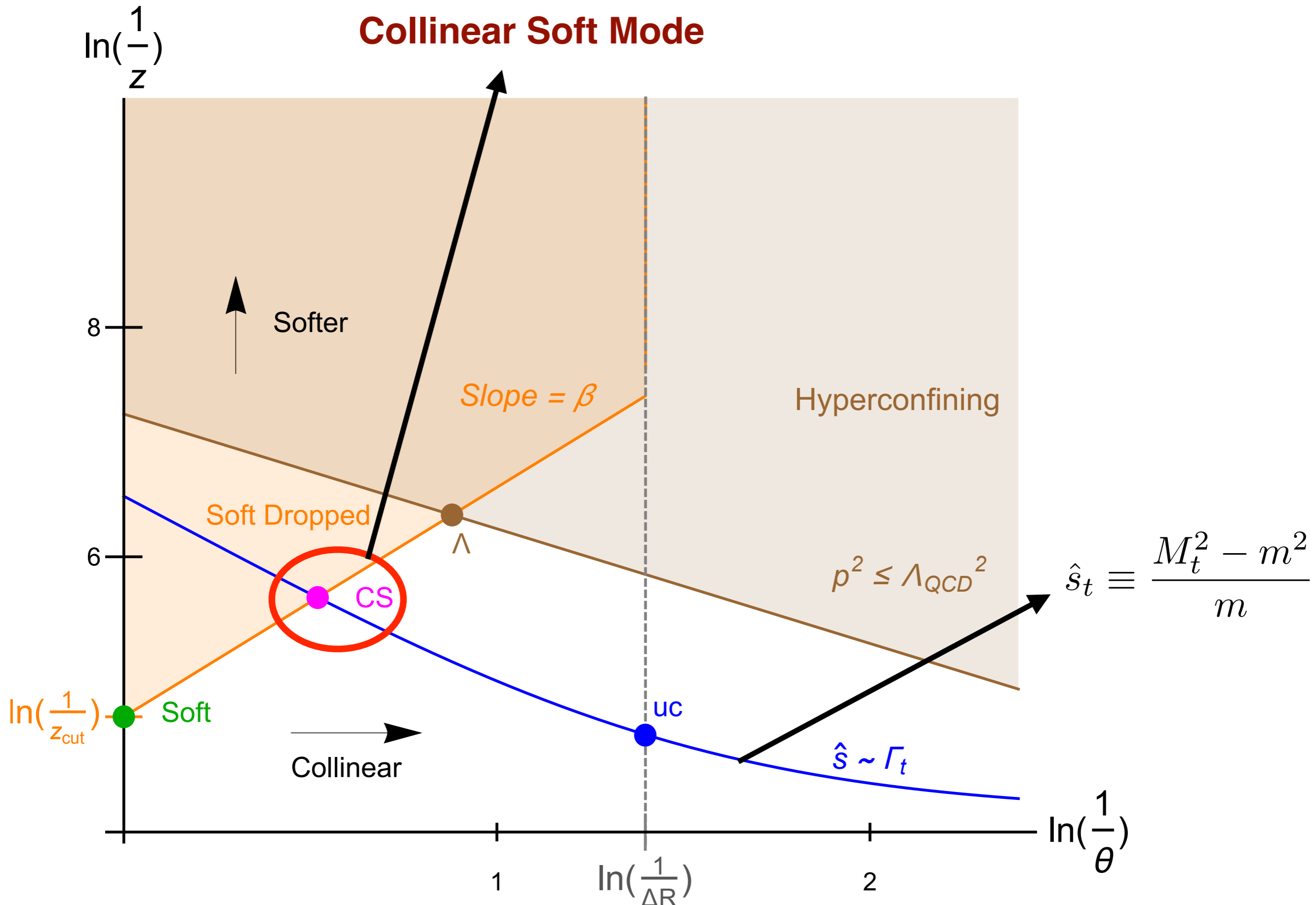


from decay products



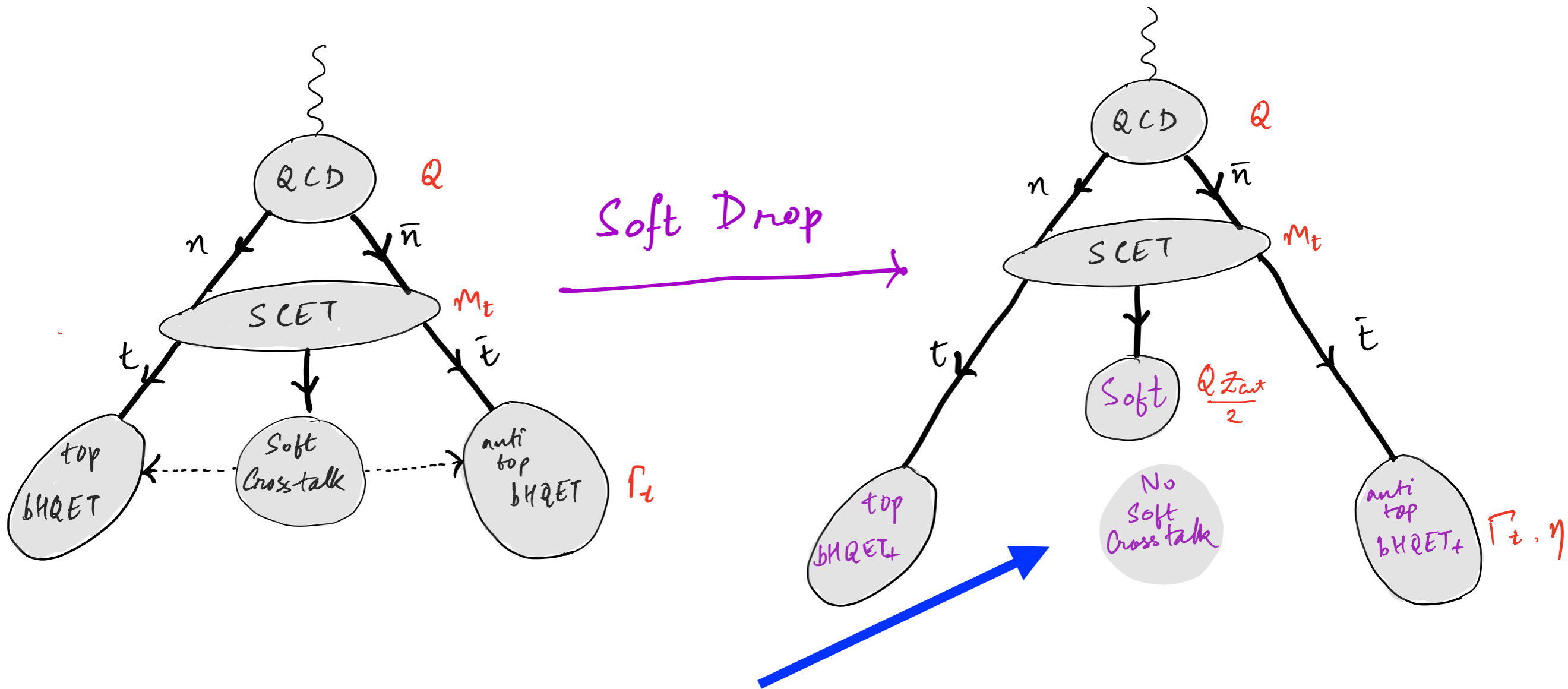
Top Jet Mass with Soft Drop

Collinear Soft Mode



Factorization with Soft Drop

$$\frac{d\sigma}{dM_J} = N \int dl dk J_B \left(\hat{s}_t - \frac{Ql}{m}, \Gamma_t, \delta m \right) S_C \left[\left(l - \frac{m}{Q} k \right)^{\frac{1+\beta}{2+\beta}} \left(2^\beta Q z_{\text{cut}} \right)^{\frac{1}{1+\beta}}, \beta \right] F_C(k)$$

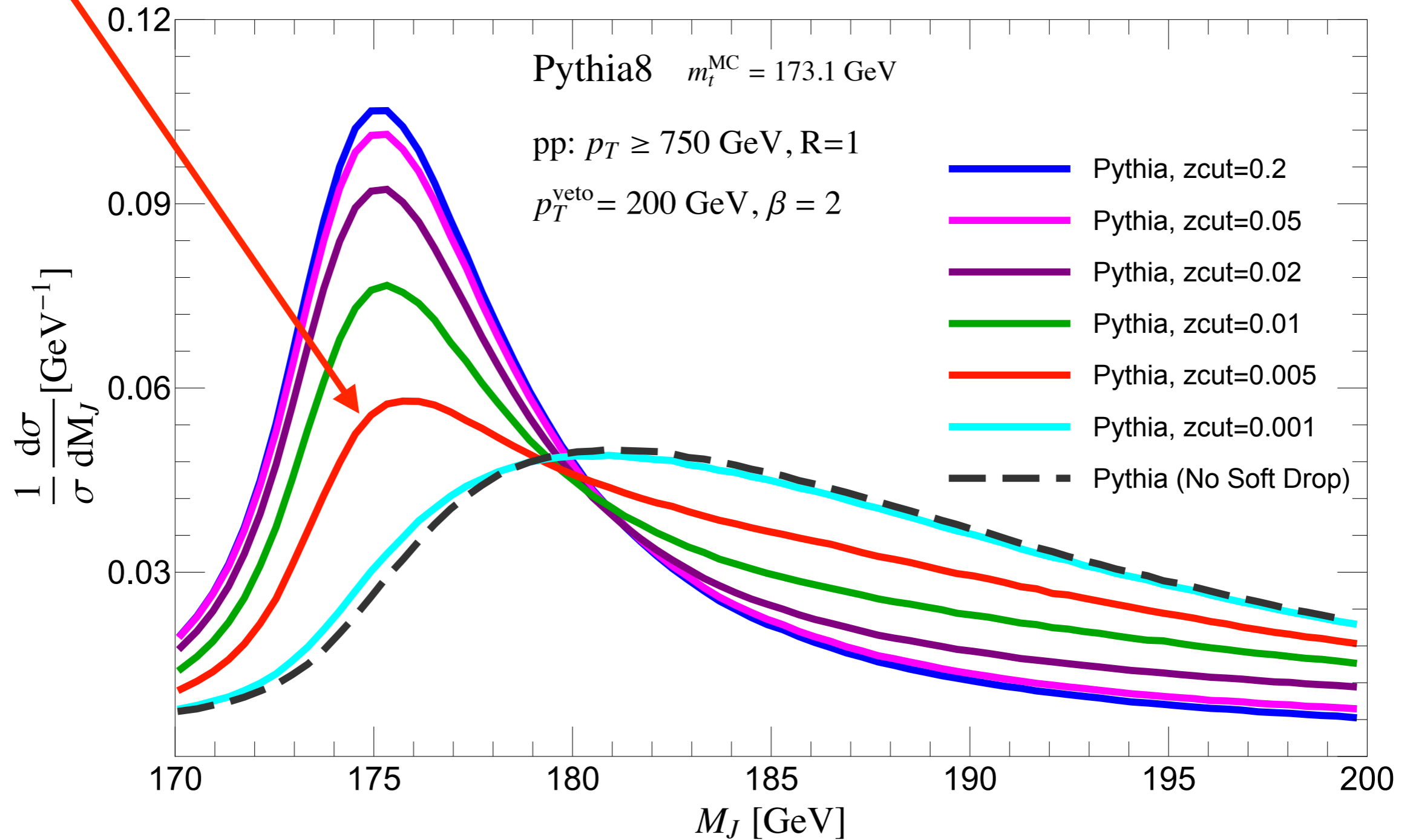


Now includes semi leptonic decays!

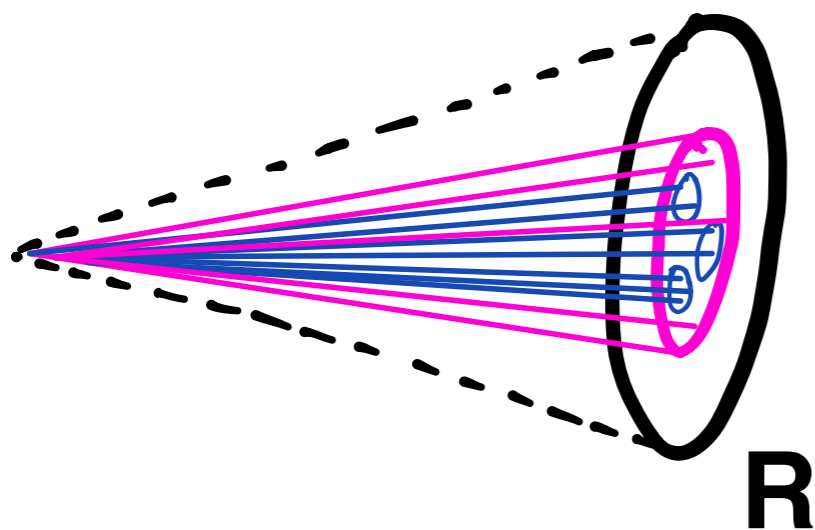
z_{cut} dependence

predict transition for “light Soft Drop” ✓

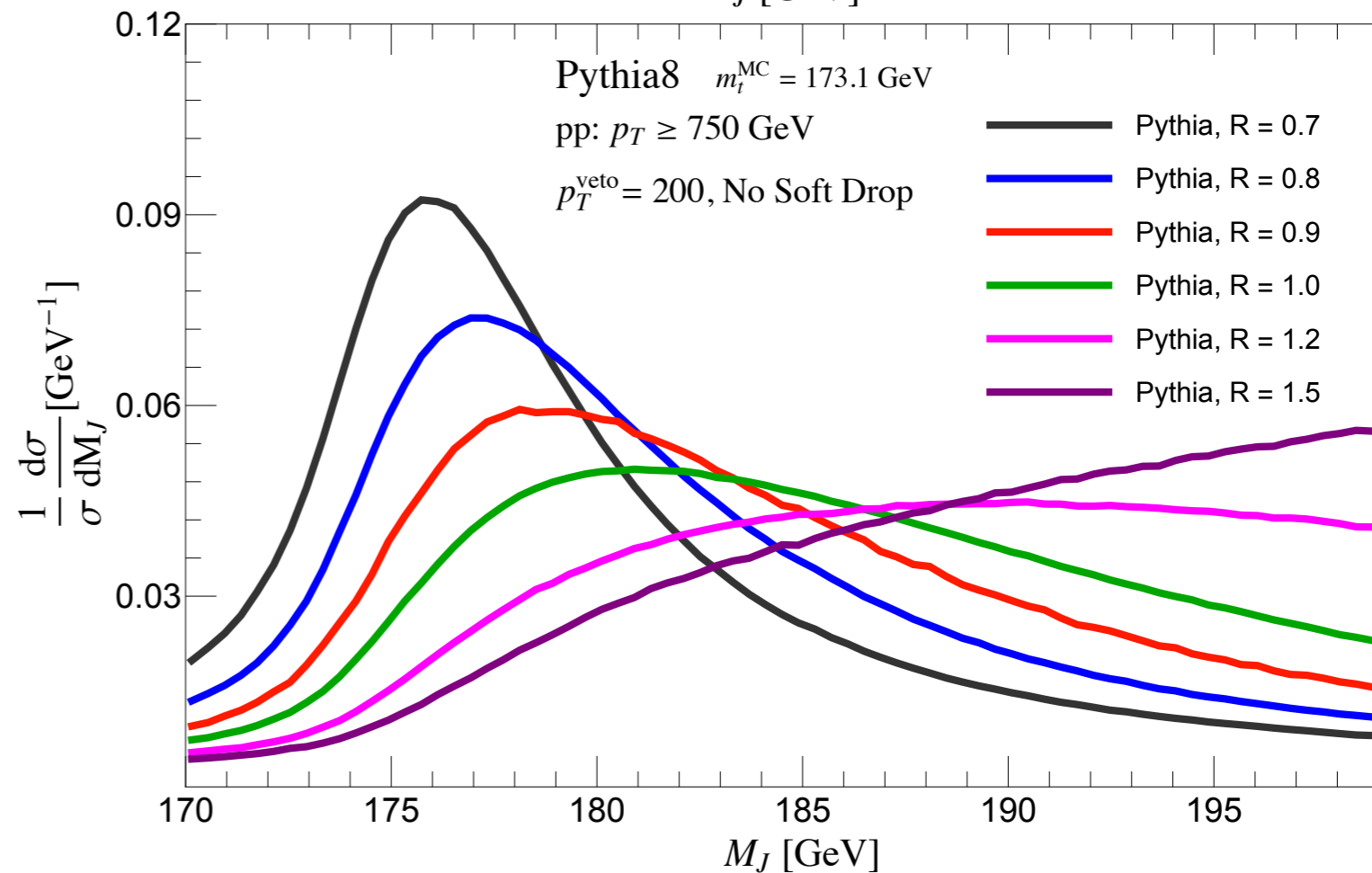
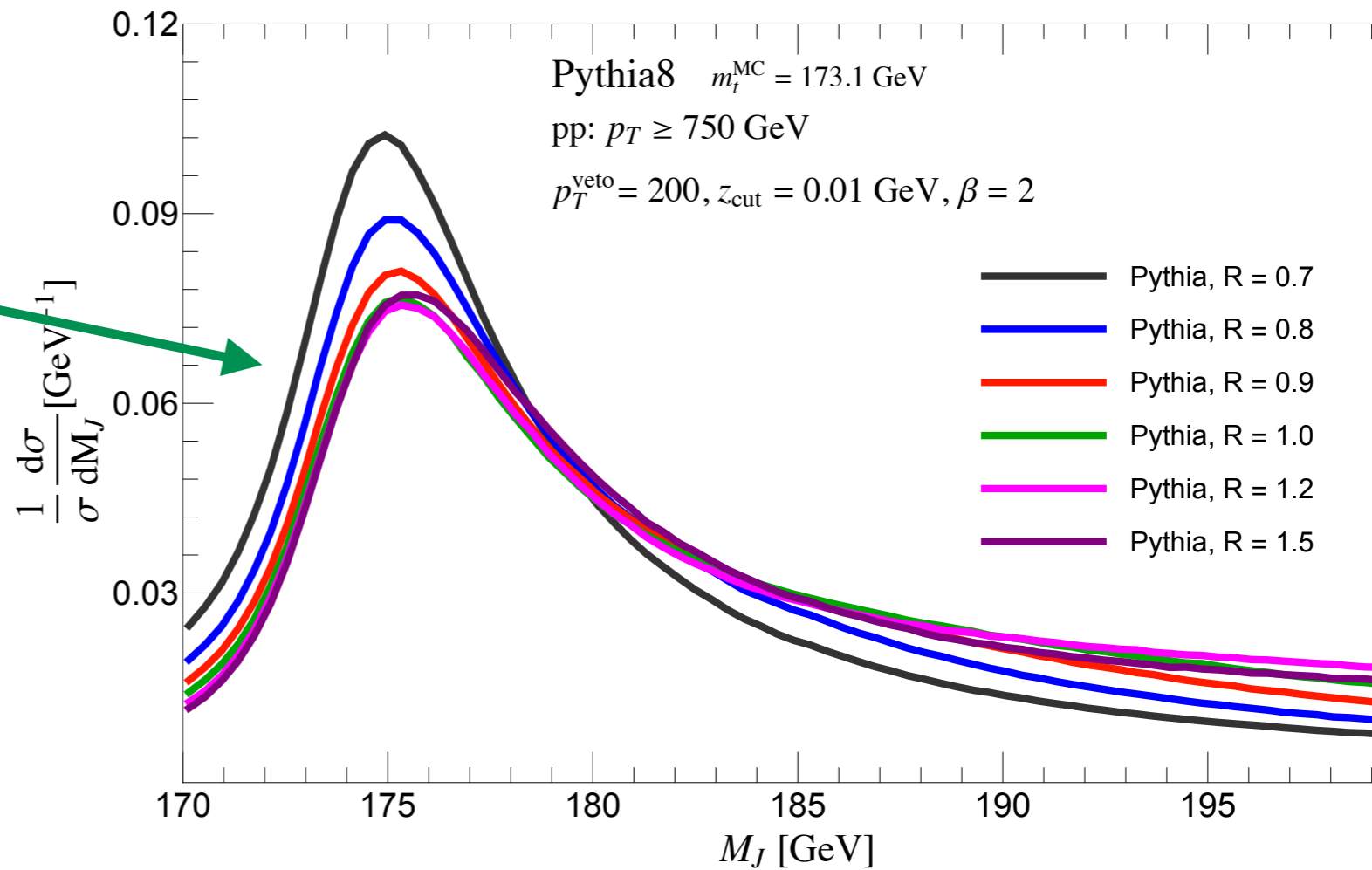
most contamination
is removed



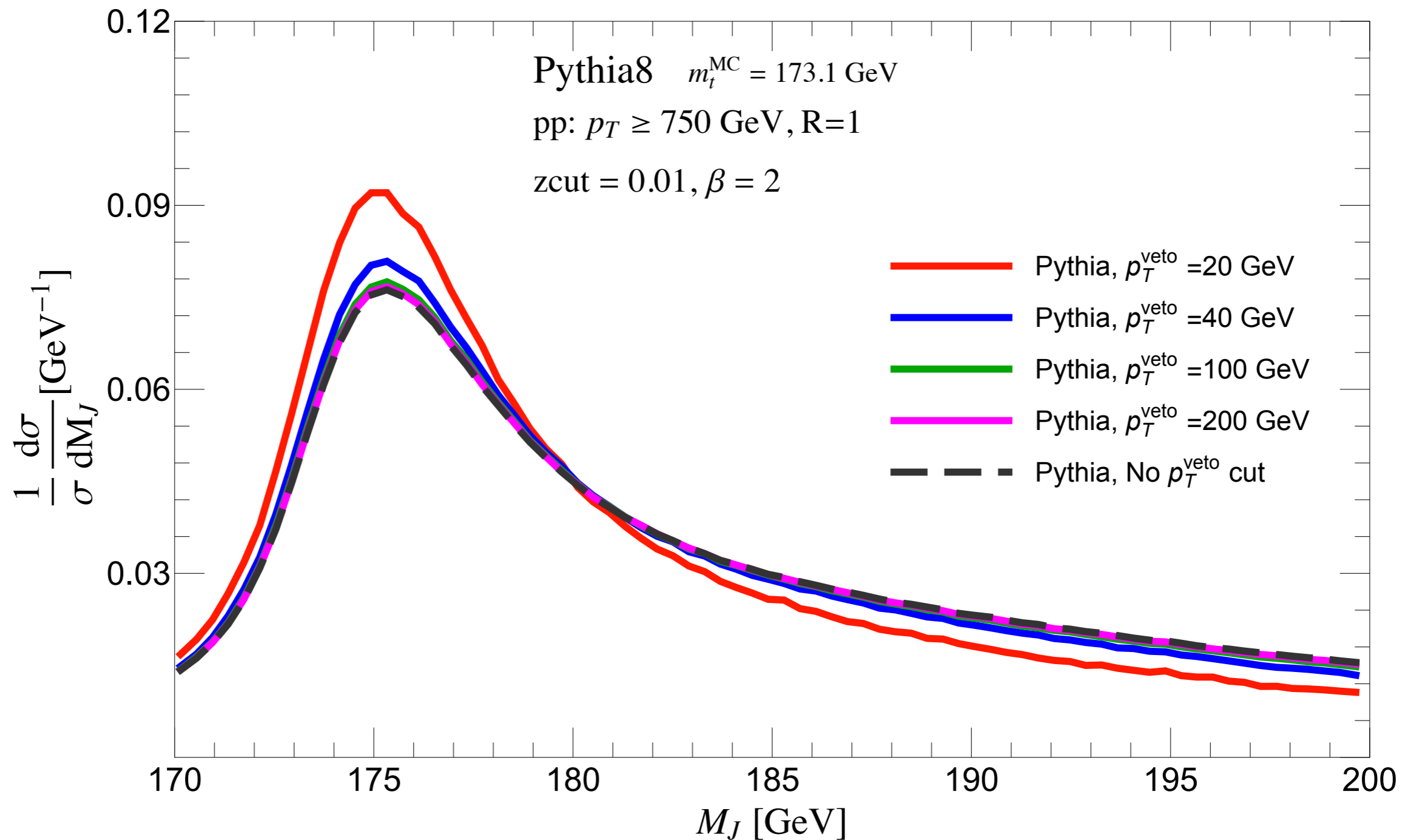
**predict:
independent of
Jet Radius**



**Without
Soft Drop
(huge):**

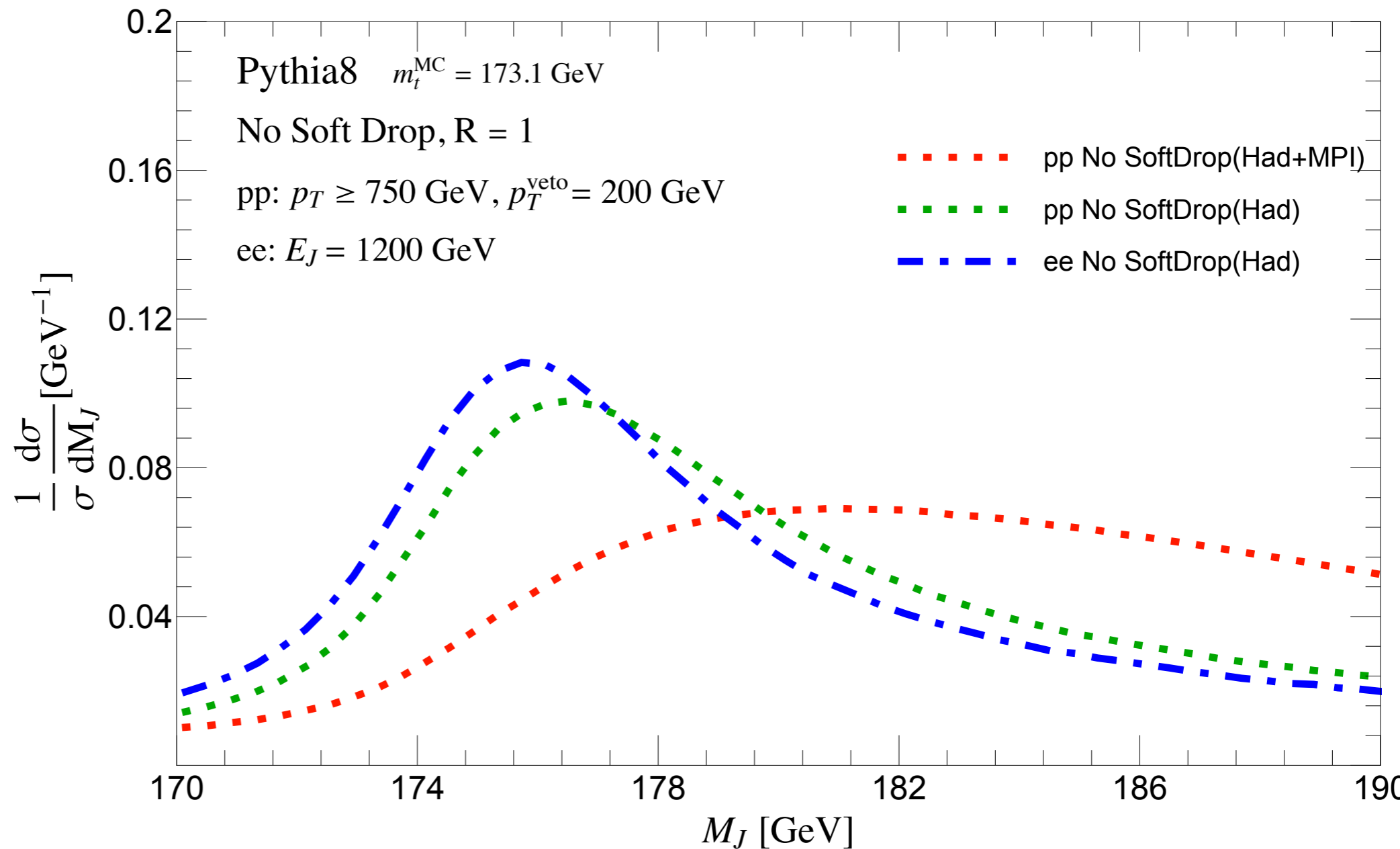


Predict independent of cutoff on radiation outside the jet (“jet veto”):



Soft Drop prediction: Same Result for e^+e^- and pp collisions

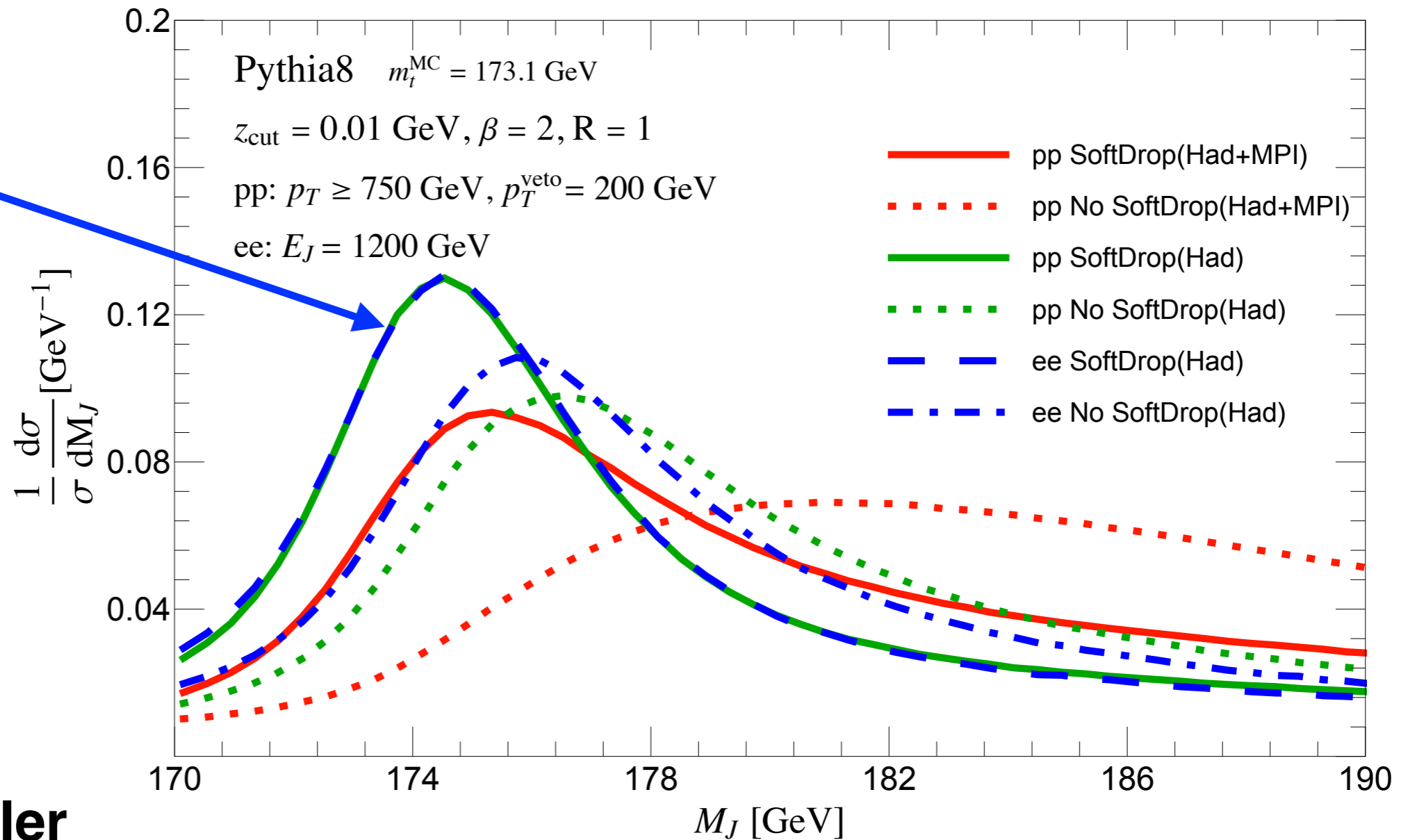
**Without
Soft Drop
(differ):**



Soft Drop prediction: Same Result for e^+e^- and pp collisions

With
Soft Drop:

Great!



much smaller
contamination

Compare Simulations to Our Theory

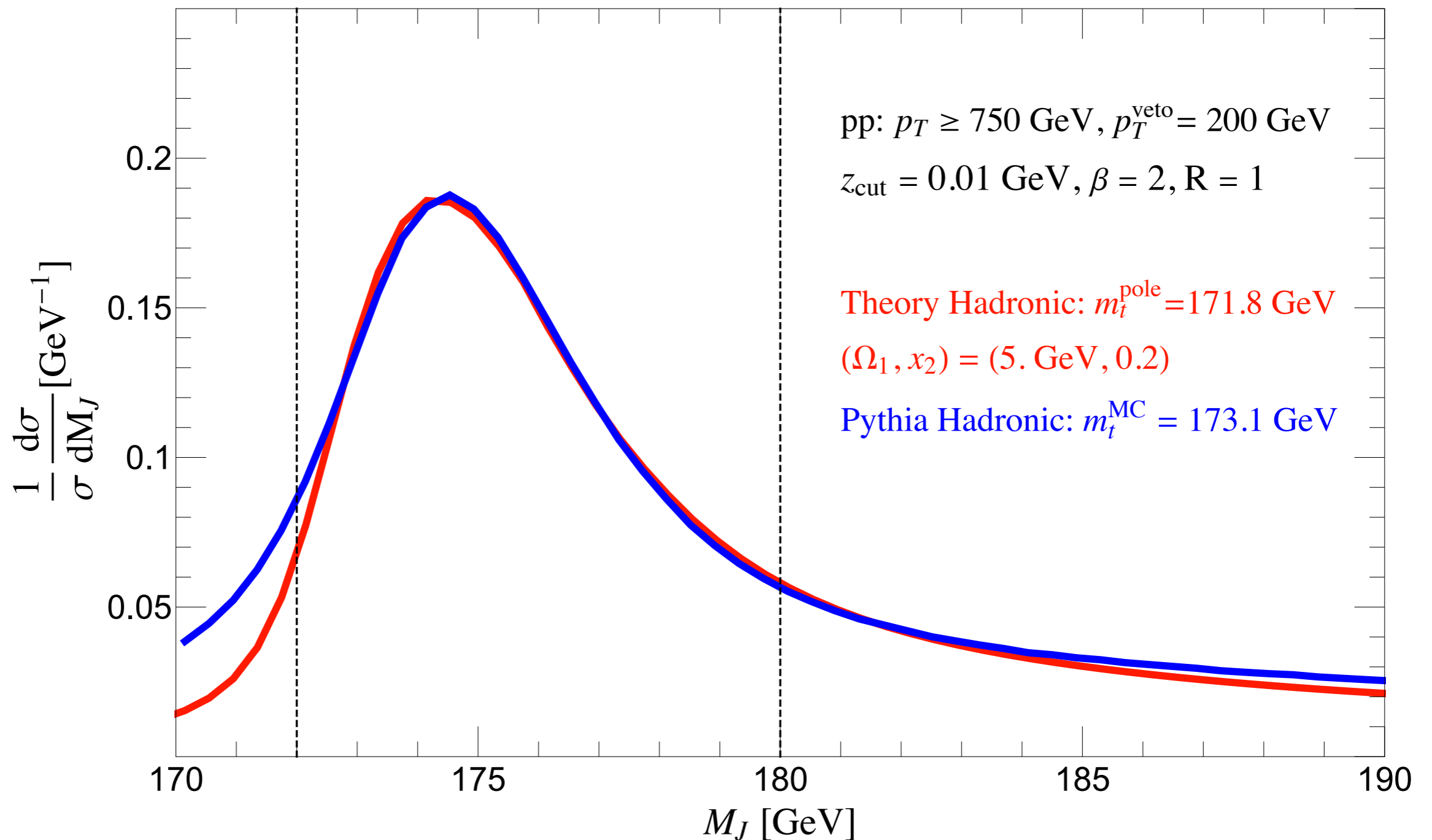
(preliminary)

Pythia Simulation vs. Theory (with Soft Drop)

**without
contamination:**

$$m_t^{\text{pole}} = 171.8 \text{ GeV}$$

$$m_t^{\text{MC}} = 173.1 \text{ GeV}$$



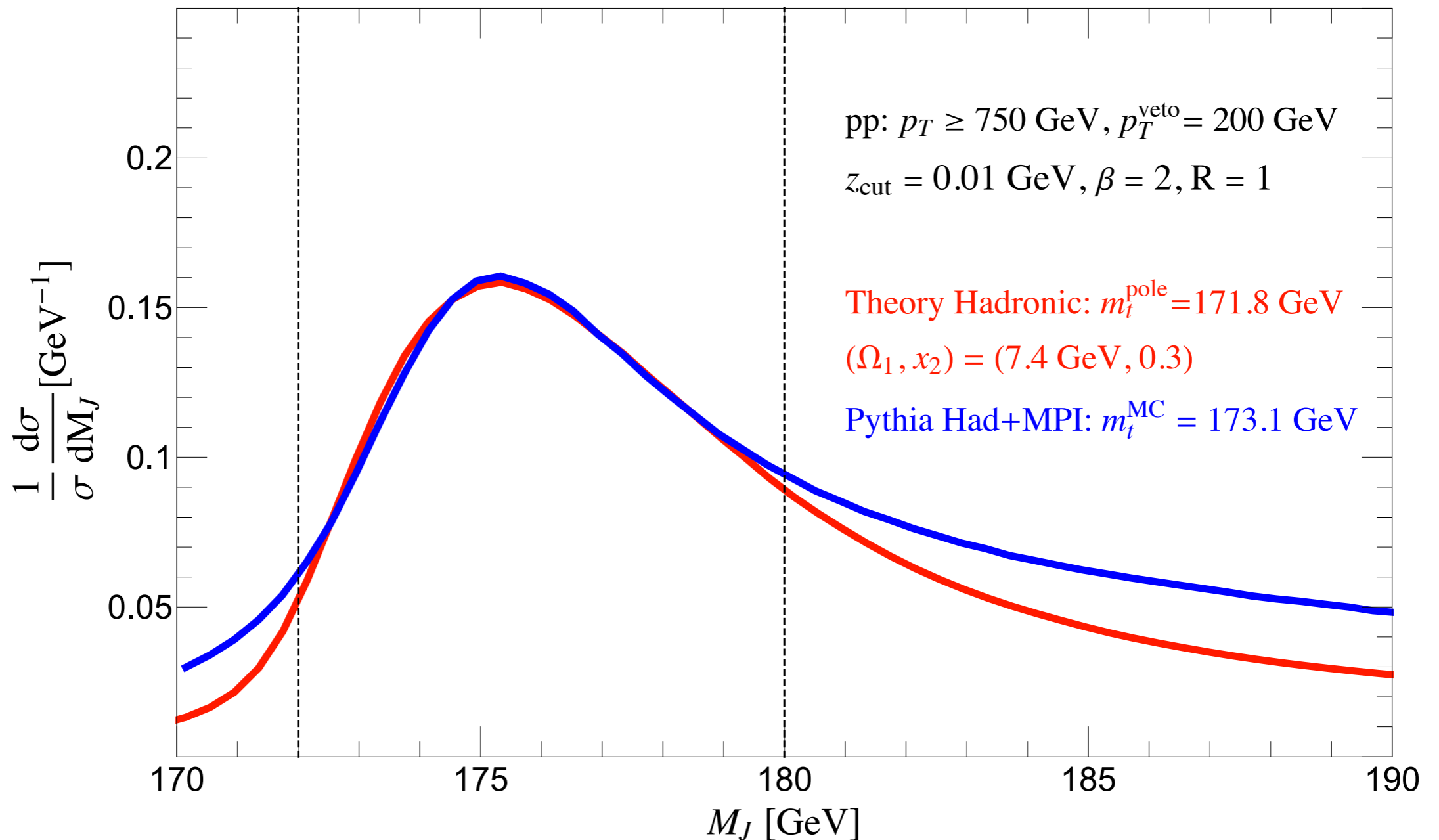
Pythia Simulation vs. Theory (with Soft Drop)

with
contamination:

$$m_t^{\text{pole}} = 171.8 \text{ GeV}$$

Same!

$$m_t^{\text{MC}} = 173.1 \text{ GeV}$$

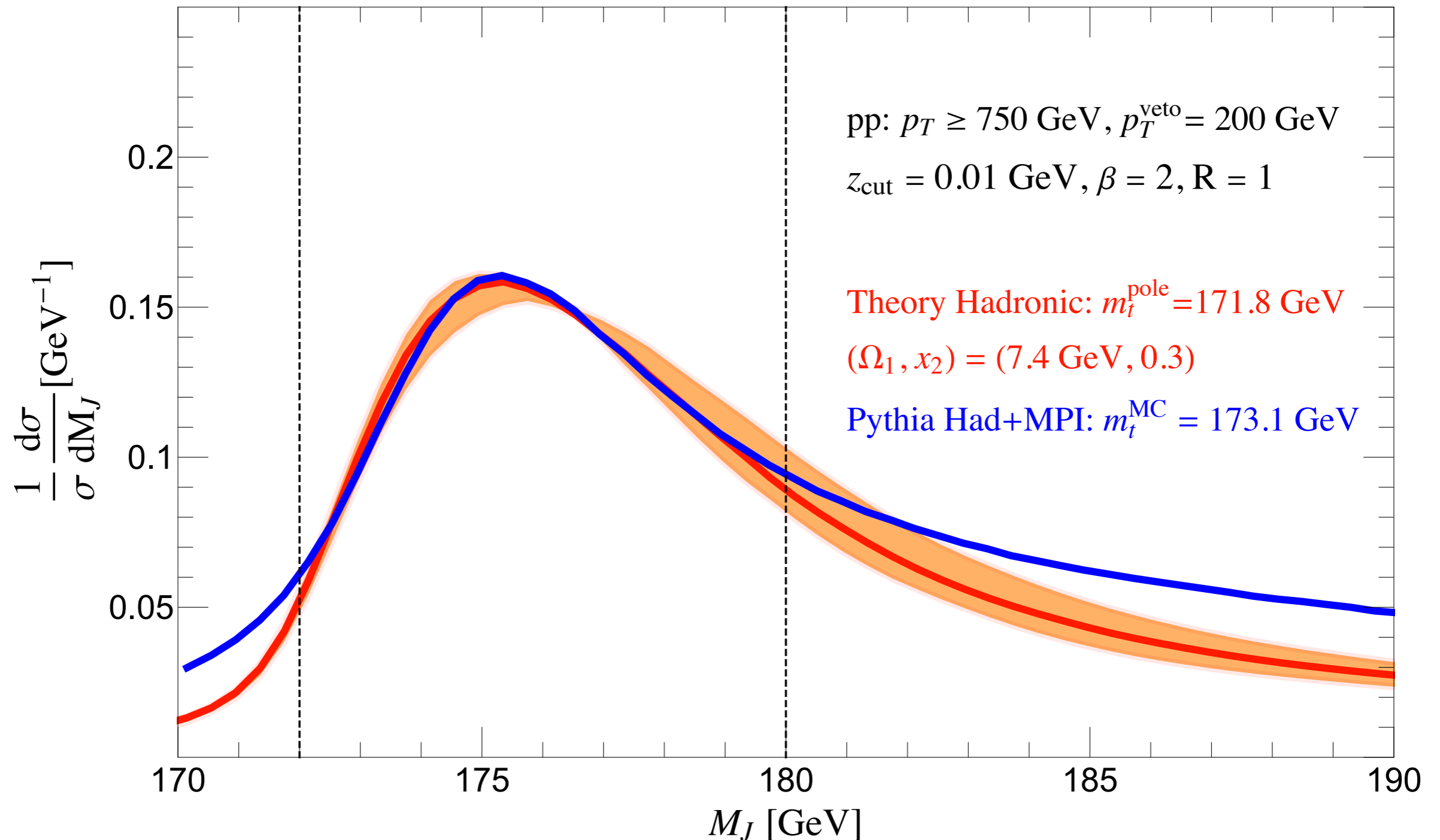


Dominant change is expected: Ω_1 (hadronization)

Pythia Simulation vs. Theory (with Soft Drop)

**Add uncertainties from
scale variation:**

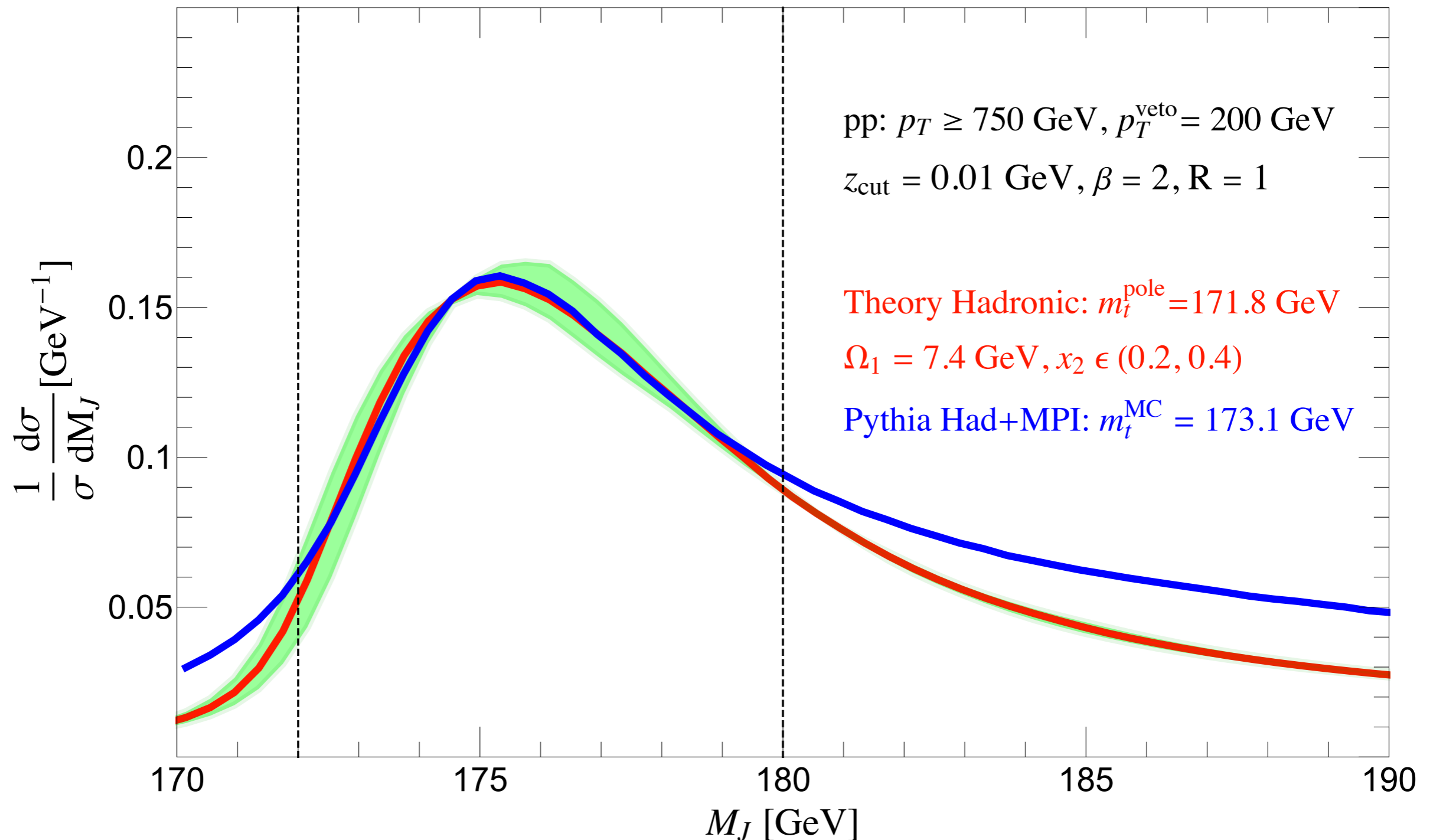
Translation of theory uncertainties to
the fit parameters is in progress.



Pythia Simulation vs. Theory (with Soft Drop)

Testing sensitivity to higher moments:

$$\Omega_1 = \int dk k F^{\text{model}}(k) \quad x_n = \frac{\Omega_n^c}{(\Omega_1)^n}$$



Summary

- Dominant uncertainty in top mass is identifying the renormalization scheme.
- Requires a QFT factorization framework, with clear top mass scheme information, to be compared to simulation or data.
- Discussed a new factorization framework for boosted top quarks.
- Exploited soft drop grooming techniques to reduce sensitivity to contamination.