



Quarkonium Production in Nucleus-nucleus Collisions at LHC

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OUTLINE

- ✿ **Why Quarkonium in heavy-ion physics**
- ✿ **Quarkonium results in Pb-Pb collisions at the LHC**
 - ⦿ **Charmonium**
 - ⦿ **Bottomonium**
- ✿ **Summary**

Where do we stand after 30 years?

A wealth of high-quality data have been accumulated, at various facilities (SPS, RHIC, LHC) for various collision systems

Decisive inputs from LHC results, having access to:

Higher energies

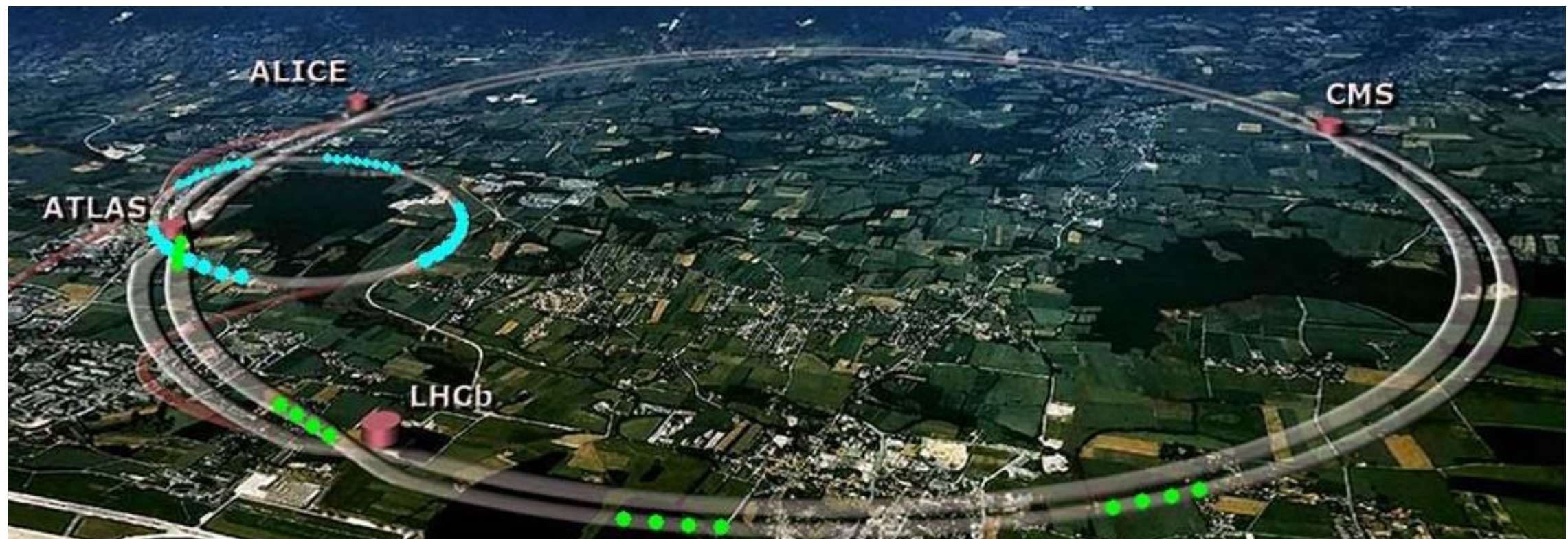
- ▶ stronger quarkonium suppression?

More charm

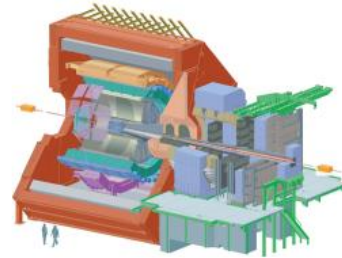
- ▶ larger (re)combination?

More bottom

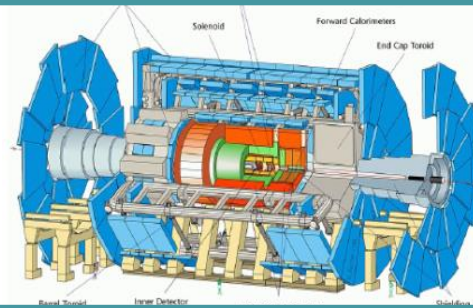
- ▶ Υ can be investigated



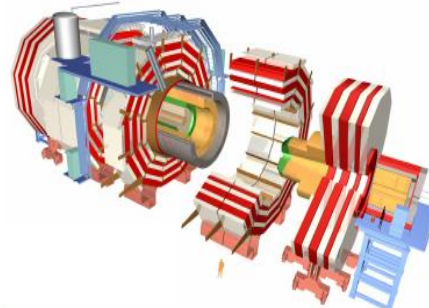
ALICE $J/\psi, \psi(2S) \rightarrow \mu^+\mu^-$
 $\Upsilon \rightarrow \mu^+\mu^-$
 $J/\psi \rightarrow e^+e^-$



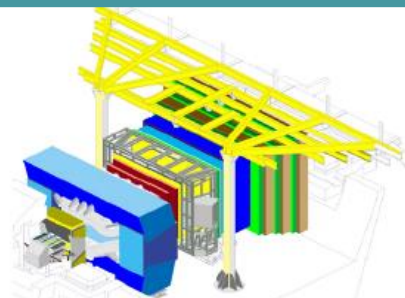
ATLAS $J/\psi \rightarrow \mu^+\mu^-$



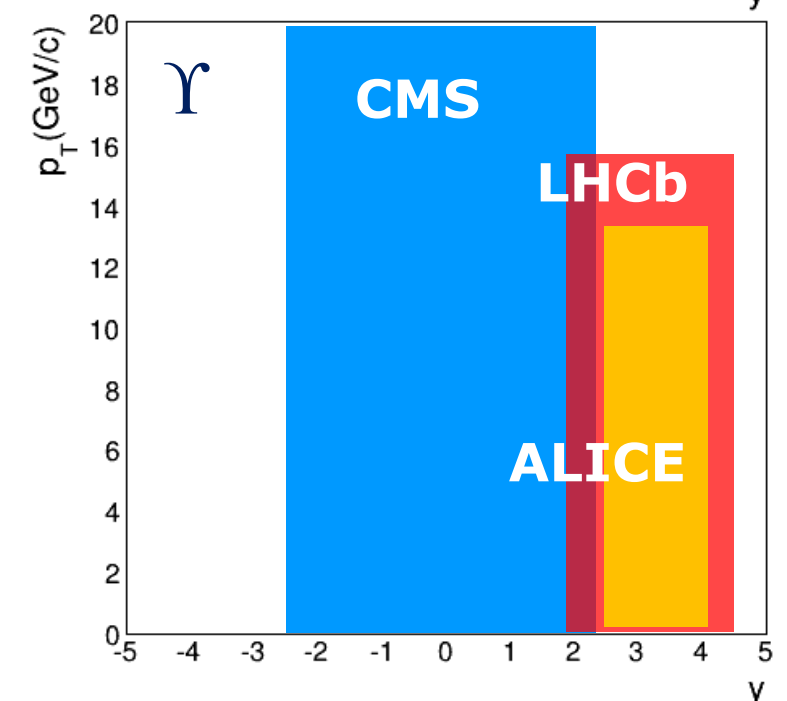
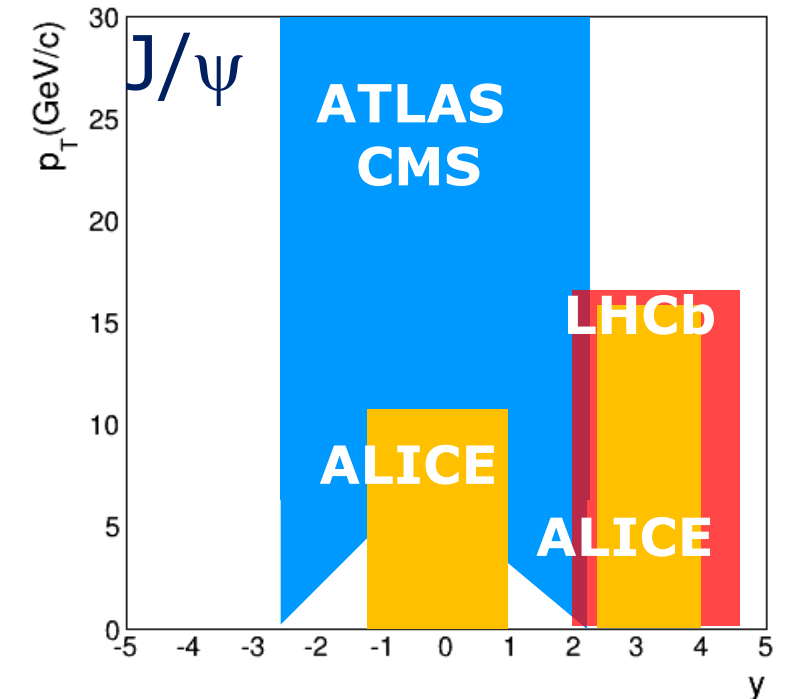
CMS $J/\psi, \psi(2S) \rightarrow \mu^+\mu^-$
 $\Upsilon \rightarrow \mu^+\mu^-$



LHCb $J/\psi, \Upsilon \rightarrow \mu^+\mu^-$



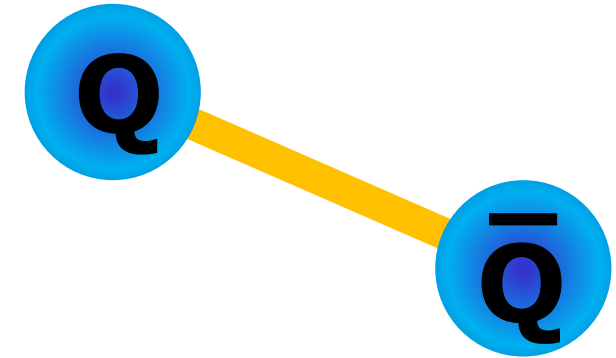
Kinematic coverage of quarkonium measurements



Complementary quarkonium
 results from LHC experiments

Quarkonium is a bound state of Q and \bar{Q} with $m_{Q\bar{Q}} < 2m_D(m_B)$

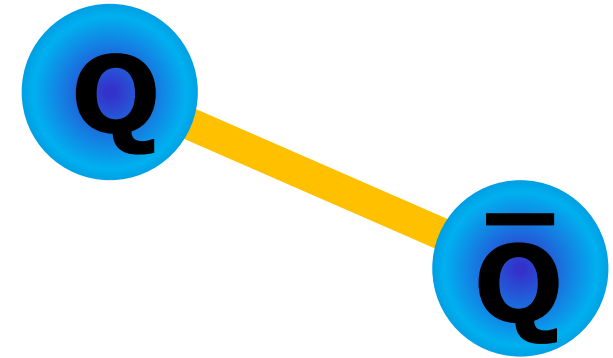
→ Several quarkonium states exist, characterized by different quantum numbers



Quarkonium at T=0

At T=0, the binding of the Q and \bar{Q} quarks can be expressed using the Cornell potential:

$$V(r) = -\frac{\alpha}{r} + kr$$



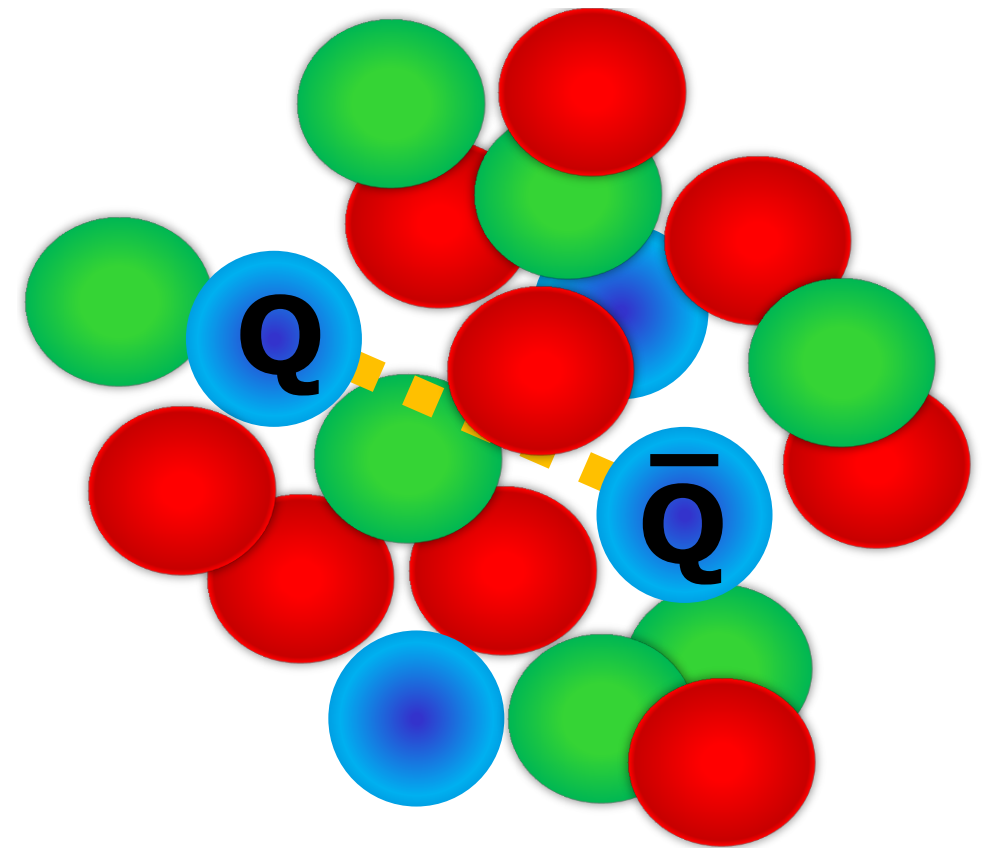
What happens to a $Q\bar{Q}$ pair placed in the QGP?

$$V(r) = -\frac{\alpha}{r} + kr$$

↓

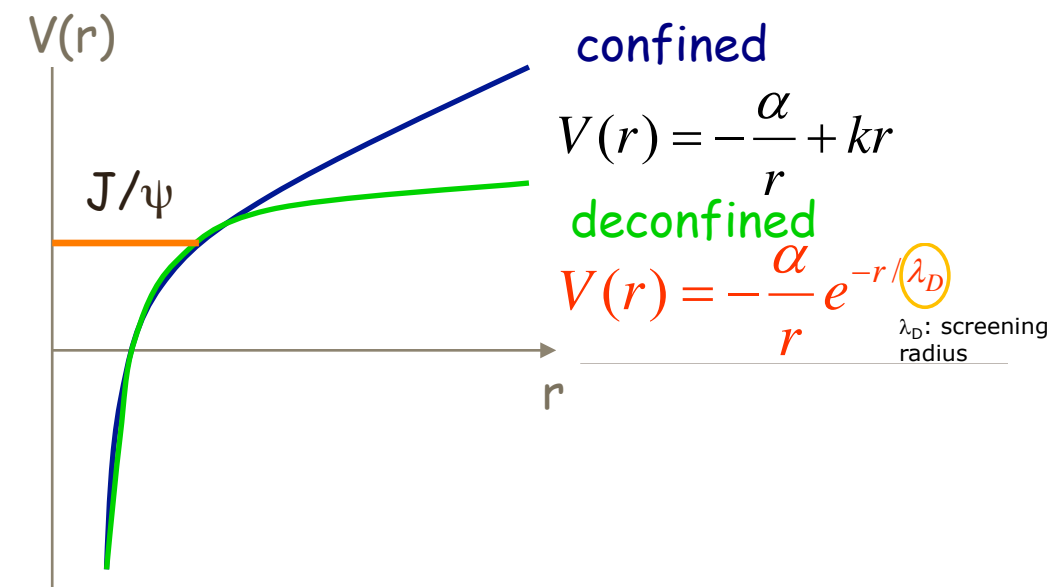
$$V(r) = -\frac{\alpha}{r} e^{-r/\lambda_D}$$

λ_D : screening radius

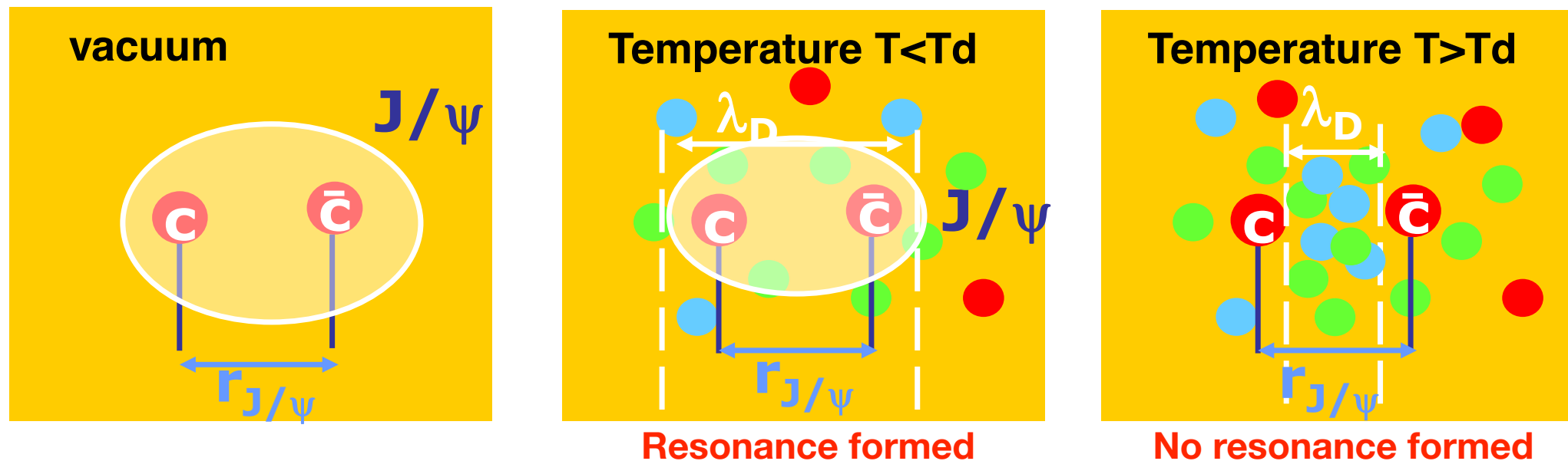


The QGP consists of deconfined colour charges
→ the binding of a $Q\bar{Q}$ pair is subject to the effects of colour screening

- The “confinement” contribution disappears
- The high color density induces a screening of the coulombian term of the potential



The **screening radius** $\lambda_D(T)$ (i.e. the maximum distance which allows the formation of a bound $Q\bar{Q}$ pair) decreases with the temperature T

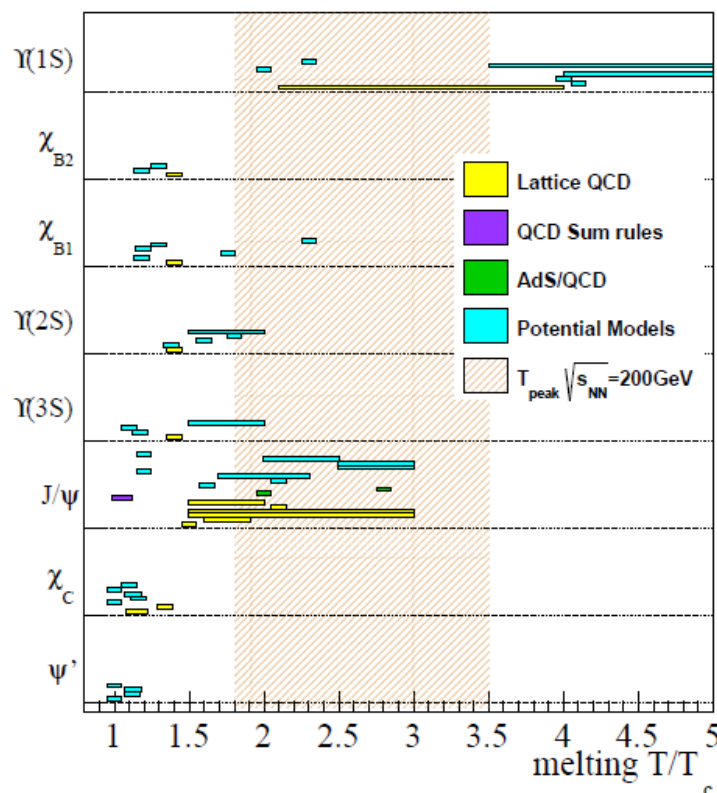


Quarkonium dissociation when

$$r_{\text{Debye}} < r_{\text{Quarkonium}}$$

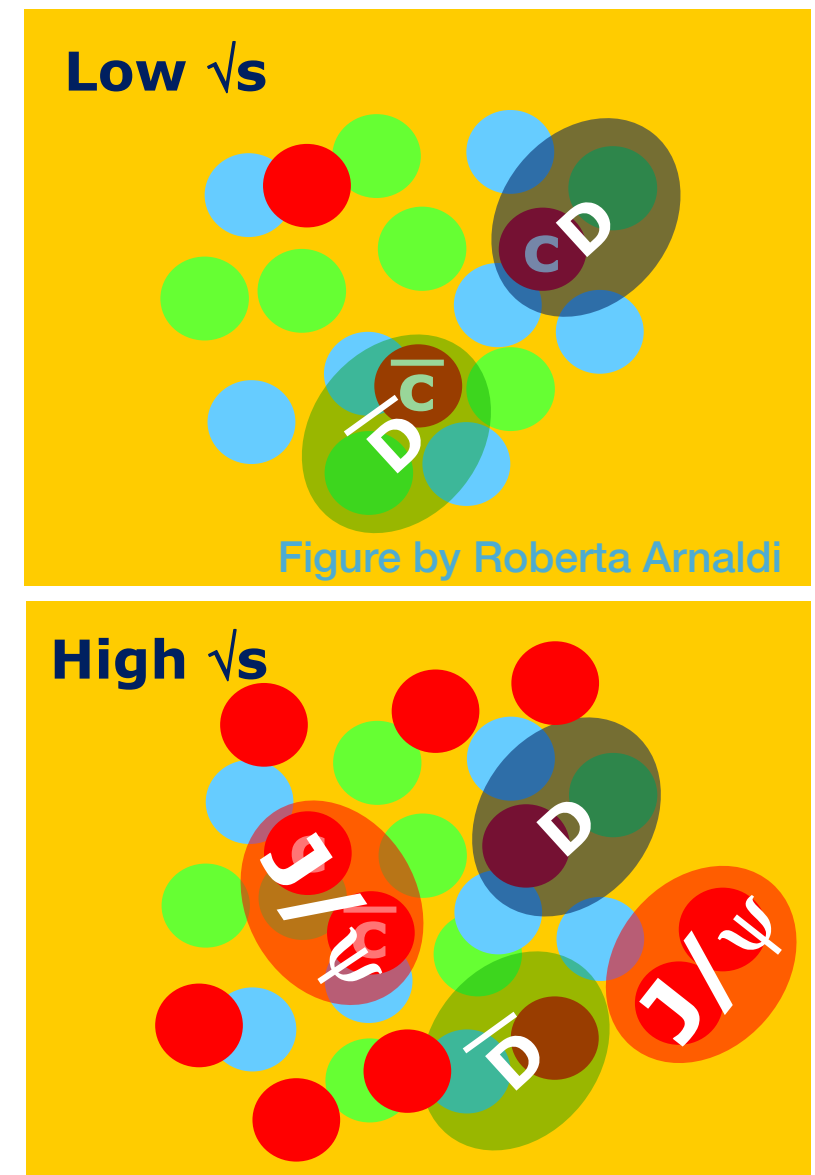
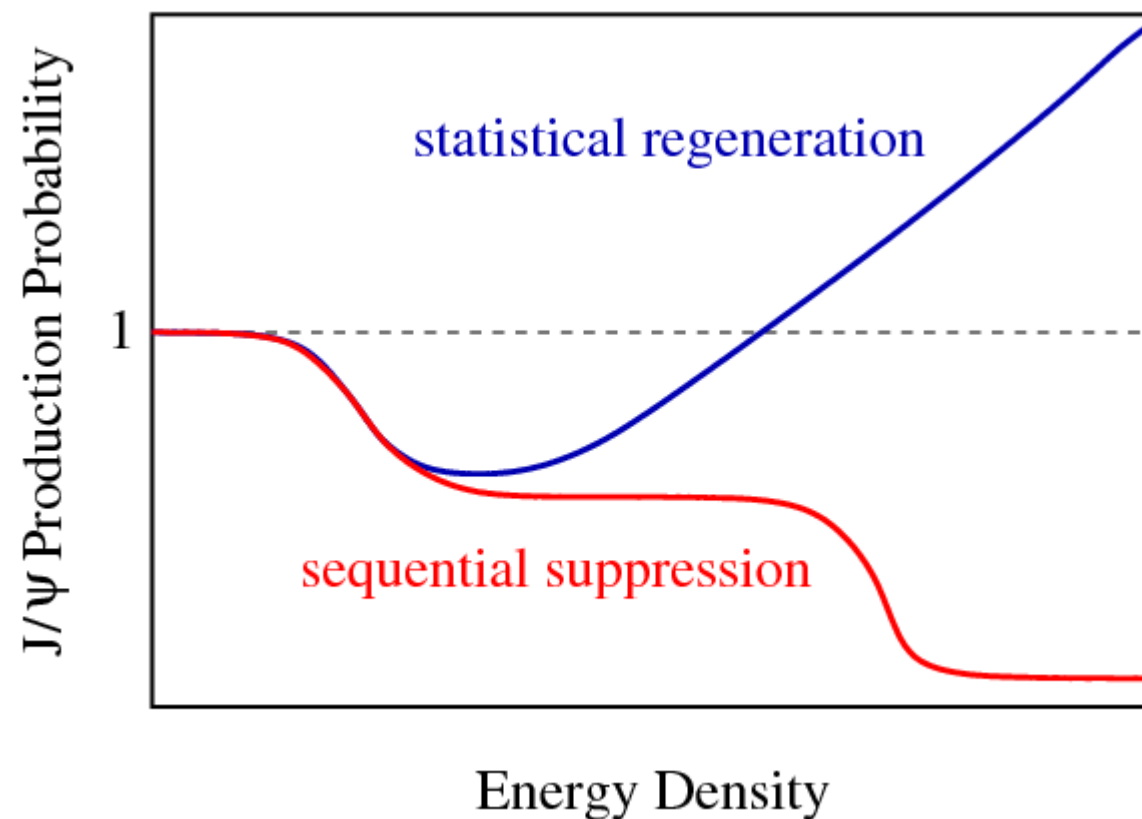
Differences in the binding energies of the quarkonium states lead to a sequential melting of the states with increasing temperature

Thermometer of the initial QGP temperature



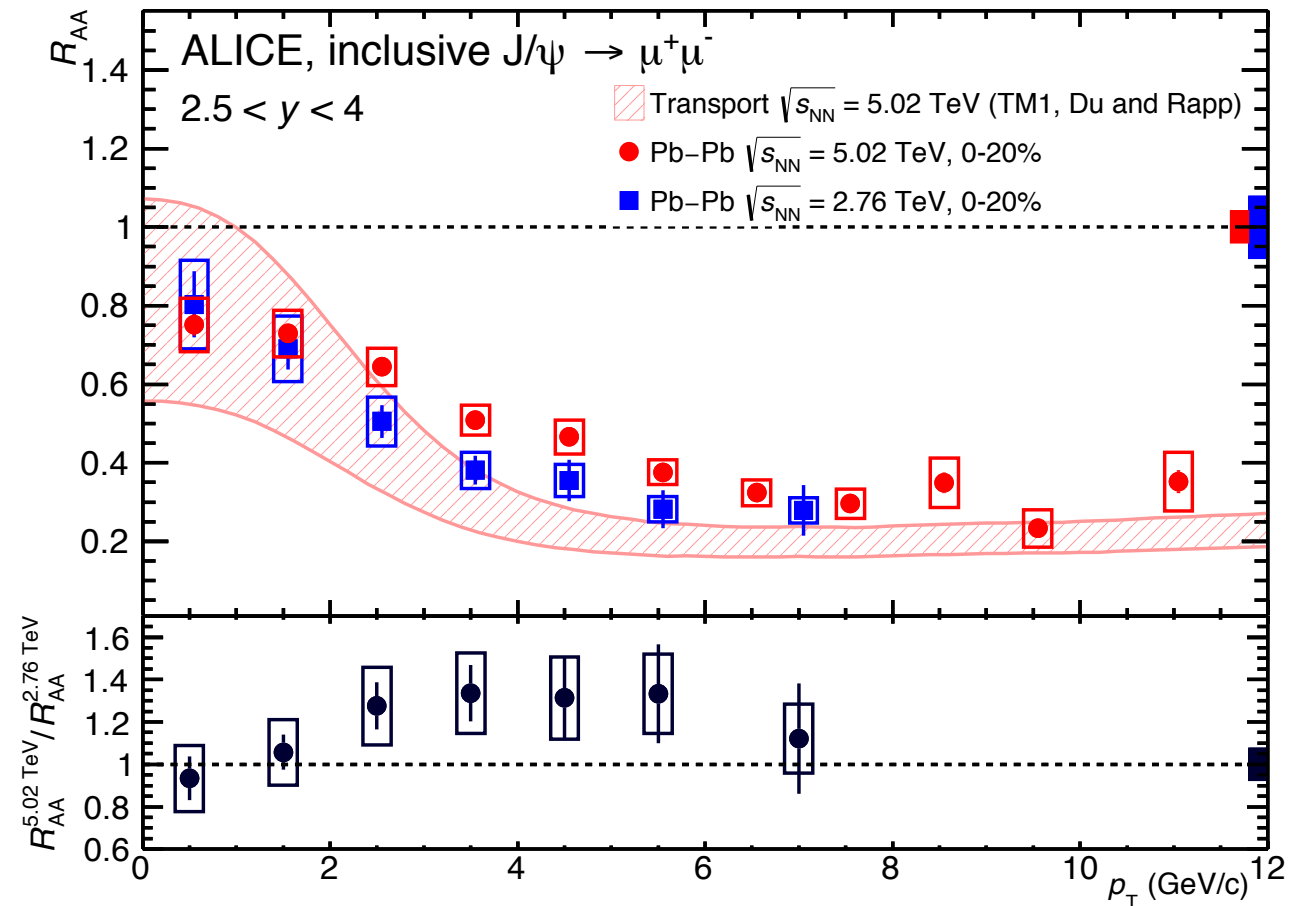
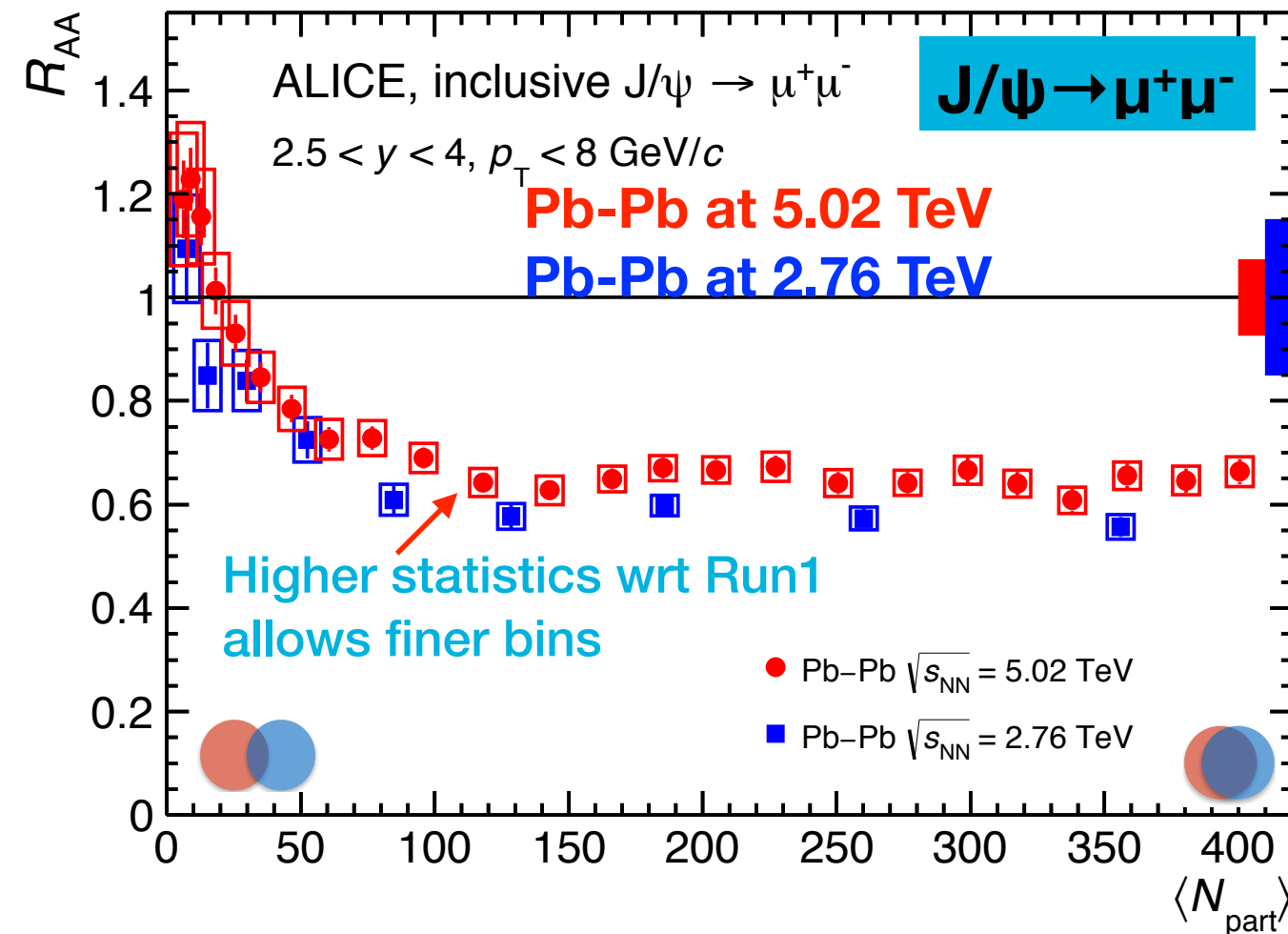
From suppression to regeneration

Central AA	SPS 20 GeV	RHIC 200 GeV	LHC 2.76 TeV	LHC 5 TeV
$N_{c\bar{c}}/\text{event}$	~ 0.2	~ 10	~ 75	~ 115



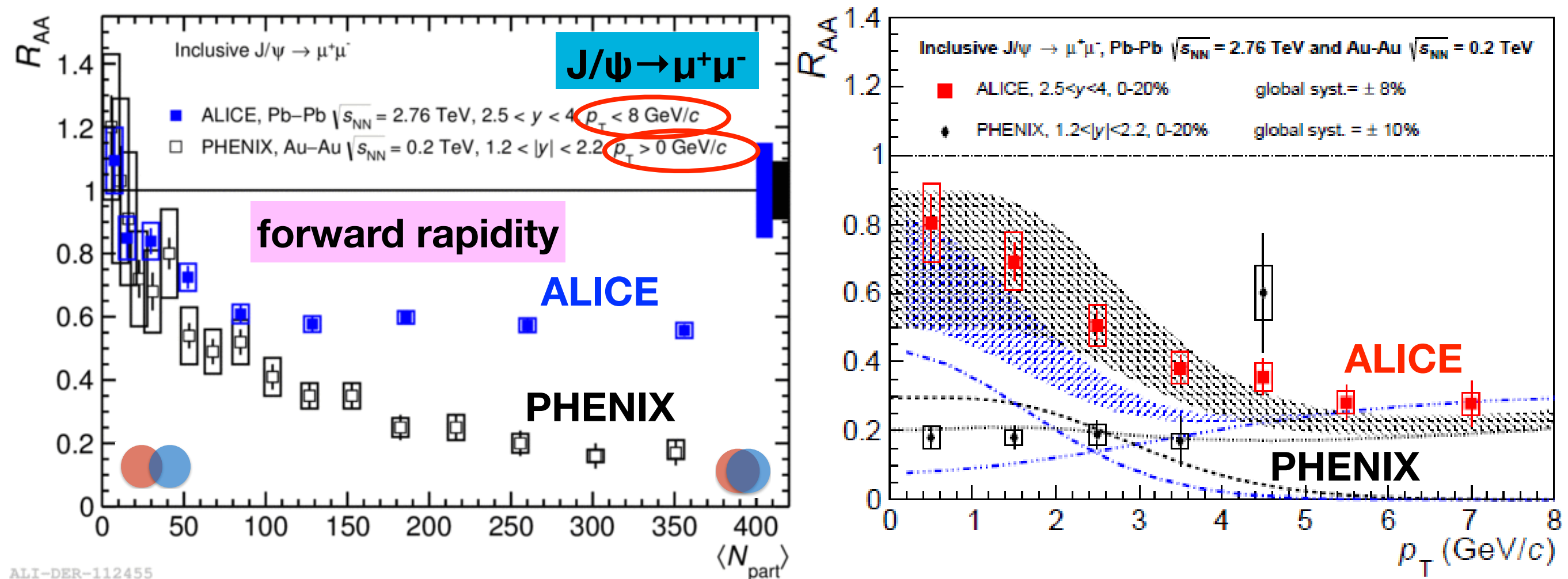
An enhancement via (re)combination of $c\bar{c}$ pairs producing quarkonia can take place at hadronization or during QGP stage

Although the “screening+recombination” picture is conceptually simple and attractive, a realistic description implies a **sophisticate treatment** (ex. In-medium formation of quarkonium, Heavy quark diffusion, ...)



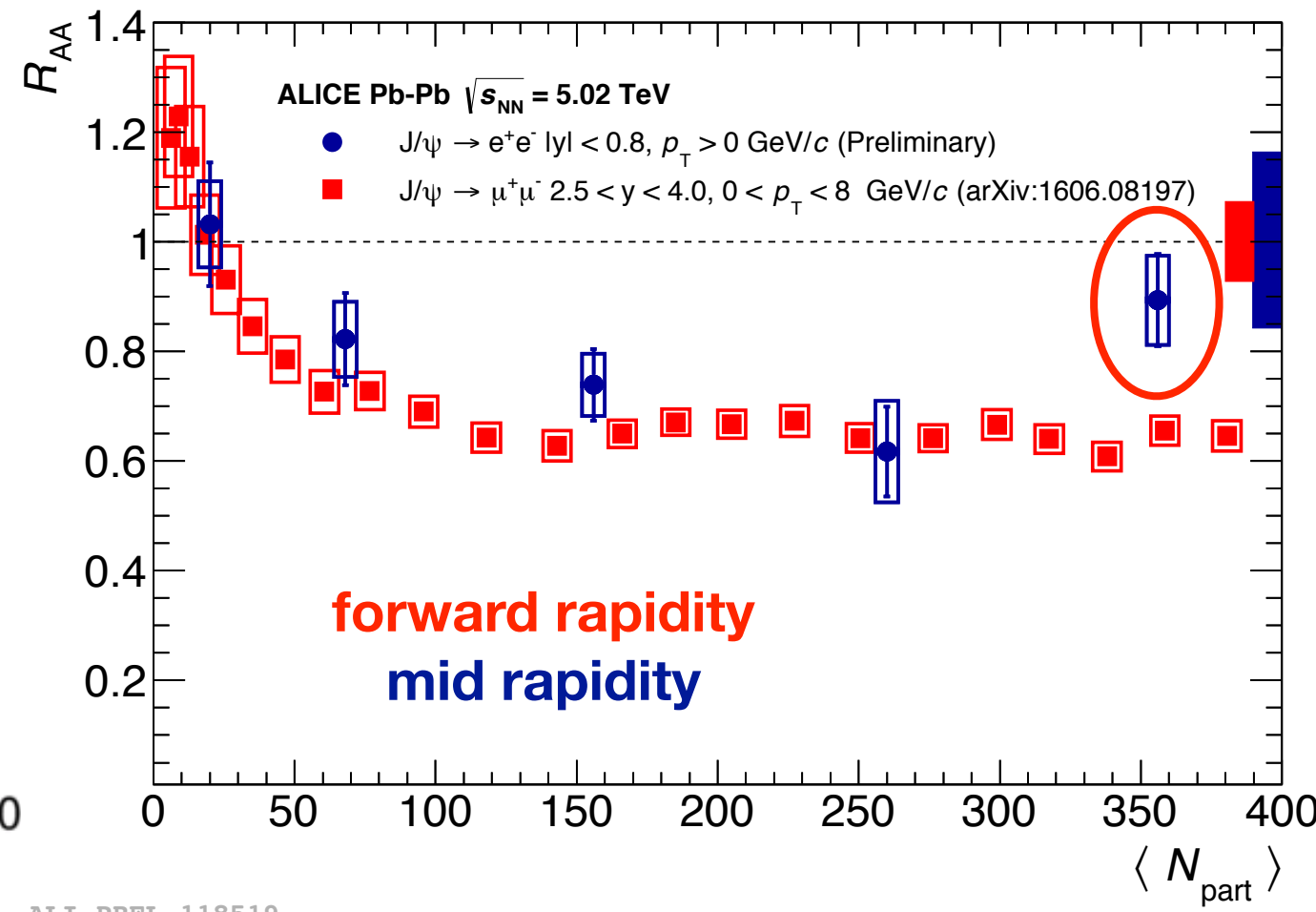
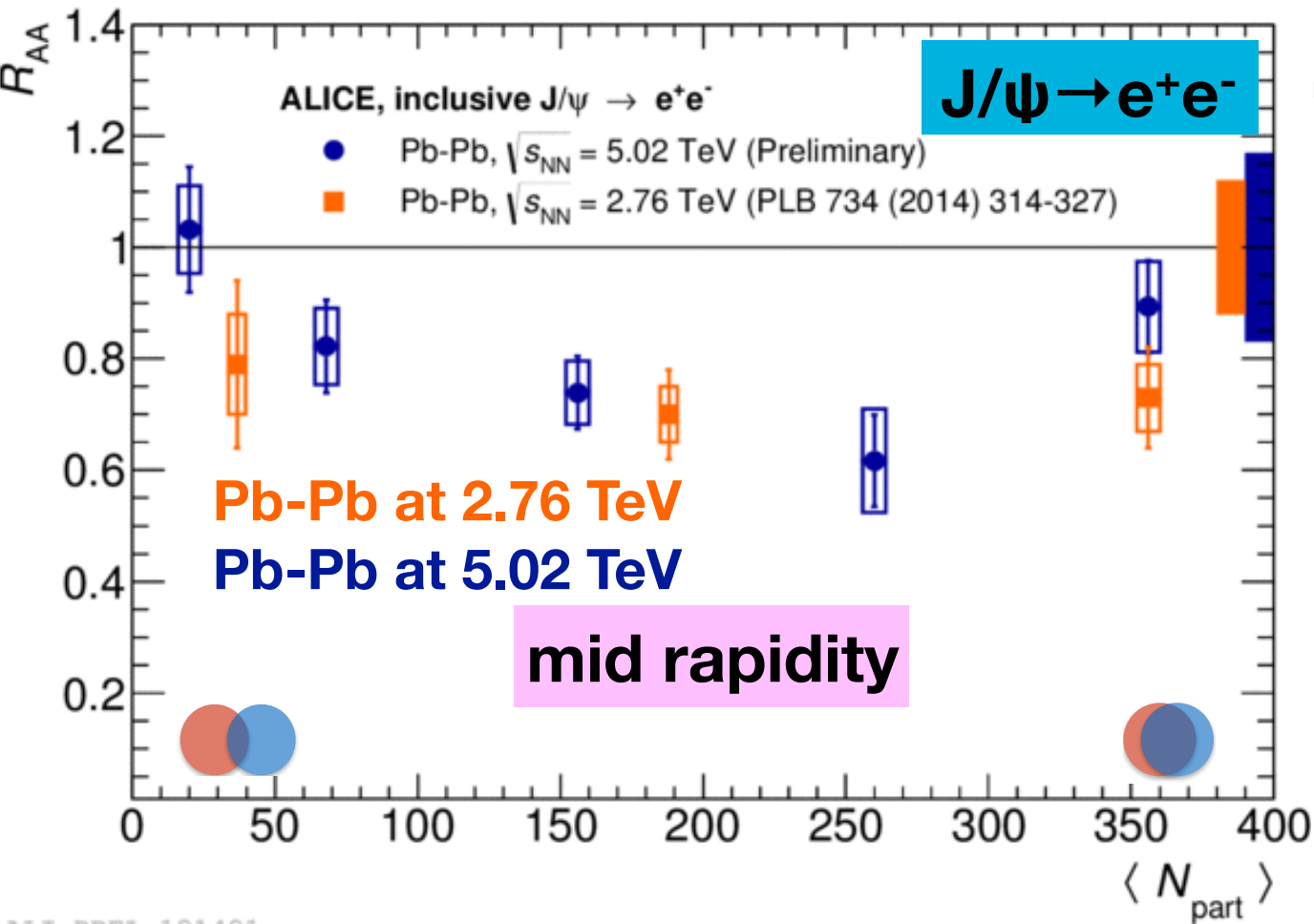
- Clear J/ψ suppression with **almost no centrality dependence** above $N_{part} \sim 100$
 - Suppression insensitive to the collision centrality in semi-central and central collisions → **indication of regeneration**
- Hint of an increase of R_{AA} at 5.02 TeV wrt 2.76 TeV is observed between 2-6 GeV/c
- J/ψ is **less suppressed at low p_T than at high p_T** → **hint of the $c\bar{c}$ recombination?**
 (as expected in regeneration models: regeneration contribution important at low p_T)

Low- p_T J/ψ : ALICE vs PHENIX



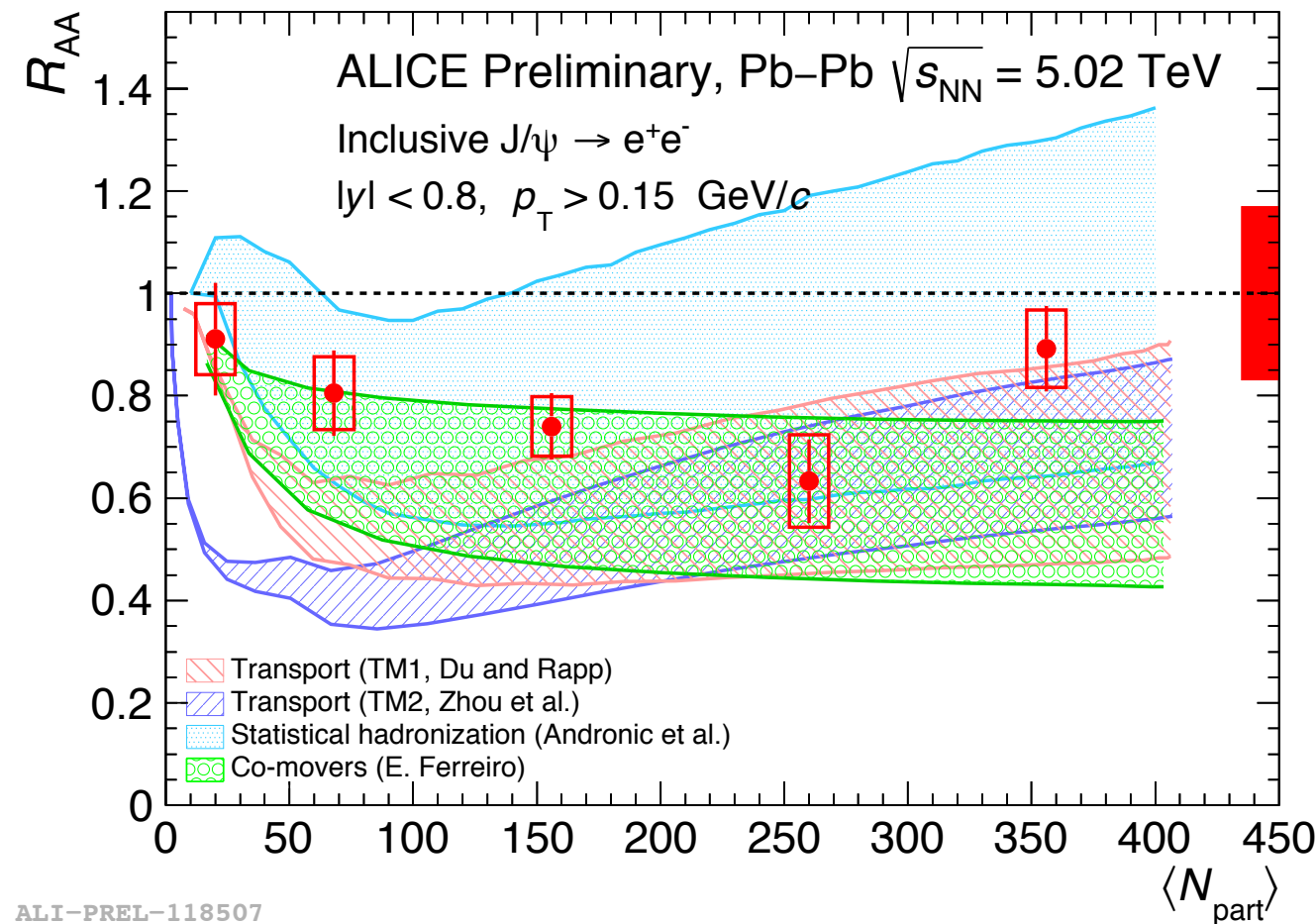
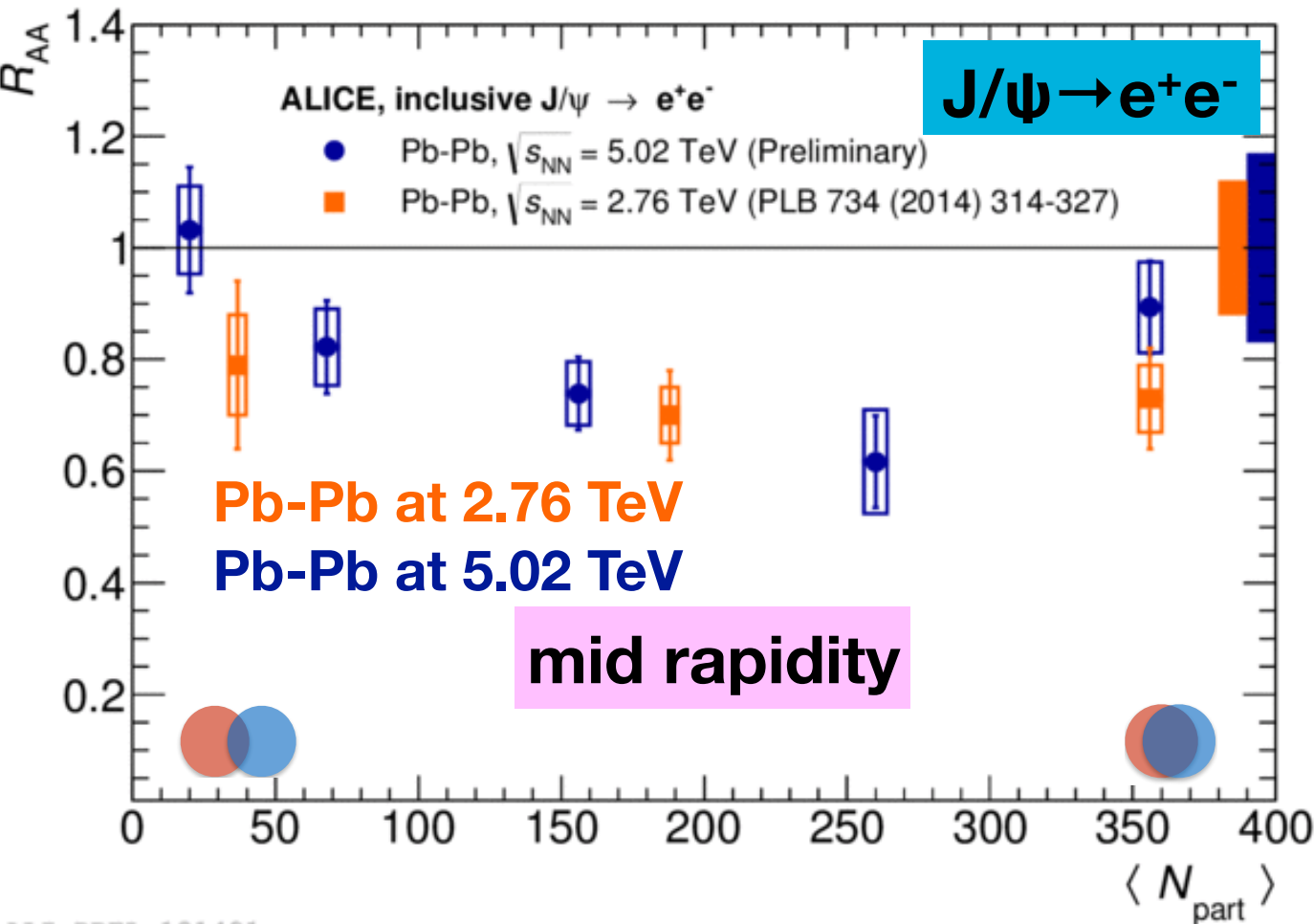
- ▶ Results vs centrality dominated by low- p_T J/ψ
 - Systematically larger R_{AA} values for central events at LHC
 - R_{AA} increases at low- p_T at LHC
- ▶ Possible interpretation:
 - RHIC energy \rightarrow suppression effects dominate
 - LHC energy \rightarrow suppression + regeneration

Low- p_T J/ψ : central vs forward rapidity



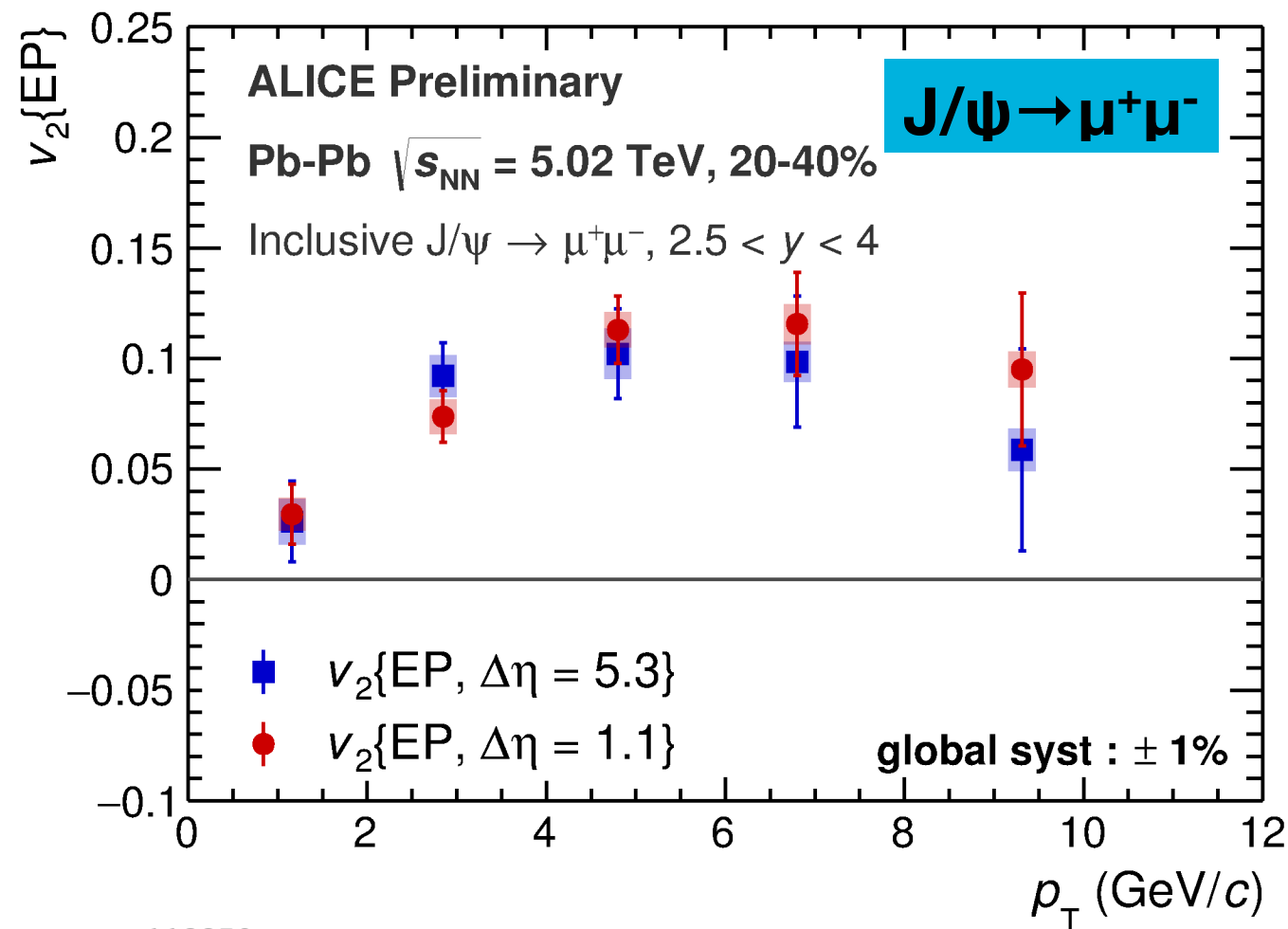
- ▶ **No significant $\sqrt{s_{NN}}$ -dependence** of R_{AA} (5.02 vs 2.76 TeV) as the observations at forward-rapidity
- ▶ Central Pb-Pb: **hints for a weaker suppression at mid-rapidity** with respect to forward-rapidity results at $\sqrt{s_{NN}}=5.02$ TeV \rightarrow expected in a (re)generation scenario (fluctuation cannot be excluded)

Low- p_T J/ψ : central vs forward rapidity

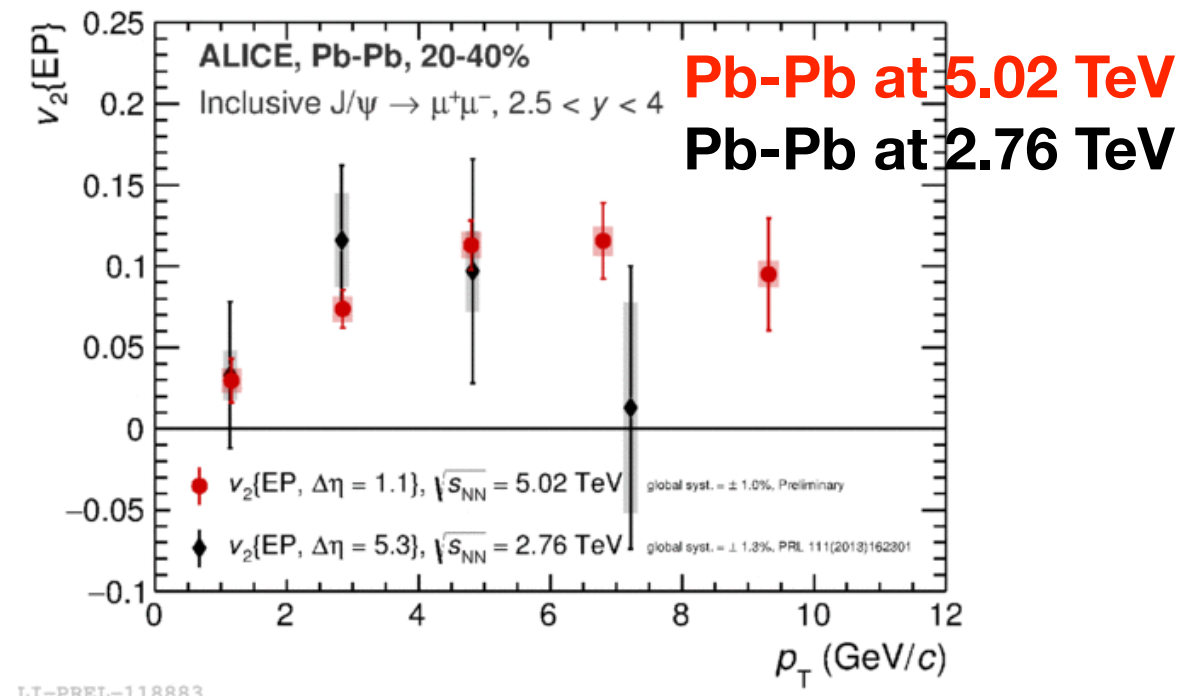


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- ▶ Central Pb-Pb: **hints for a weaker suppression at mid-rapidity** with respect to forward-rapidity results at $\sqrt{s_{NN}}=5.02$ TeV \rightarrow expected in a (re)generation scenario (fluctuation cannot be excluded)
- ▶ **Transport and statistical models have large uncertainties** (shadowing +open charm cross section)

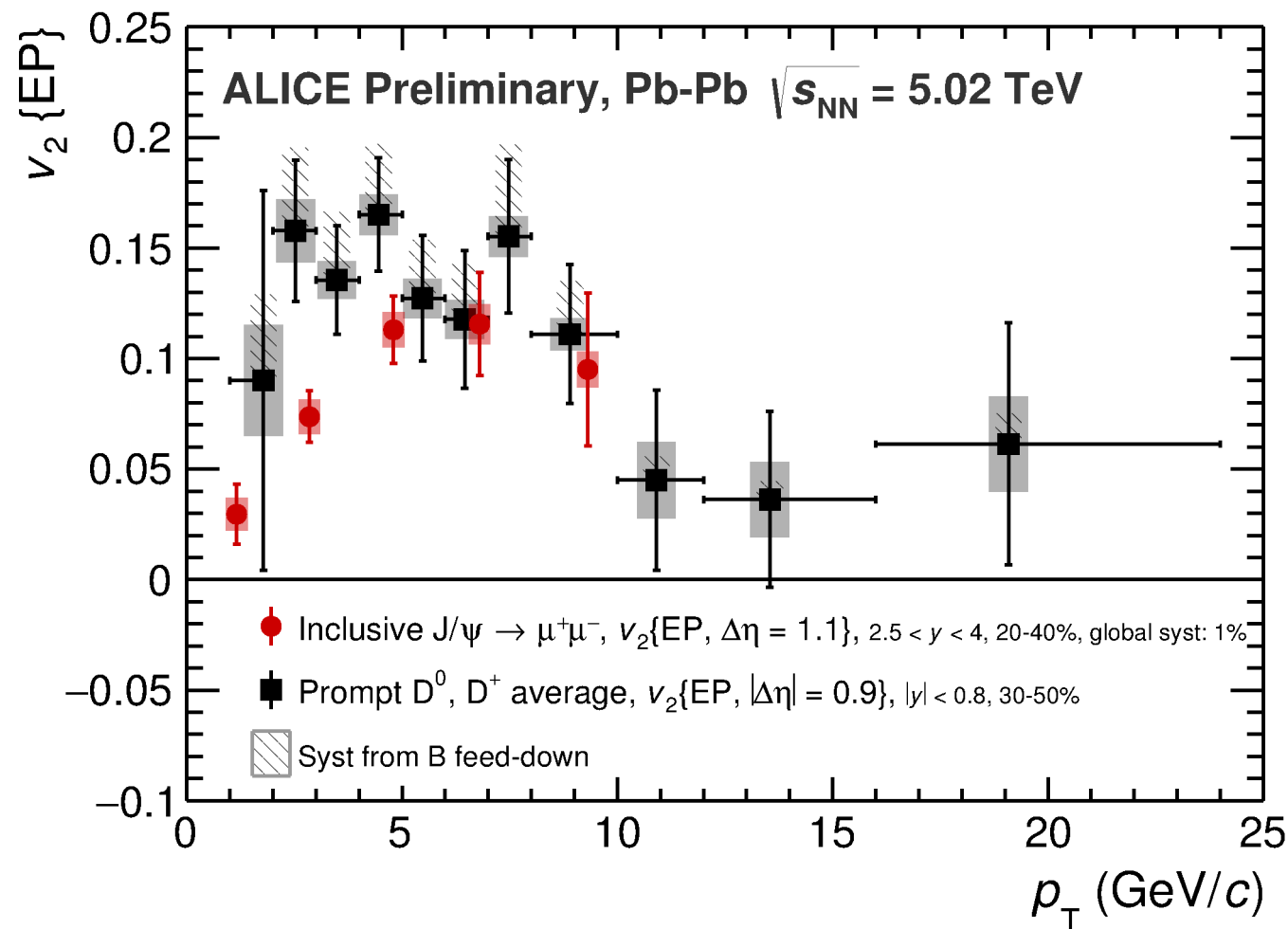
J/ψ v₂ results



p_T (GeV/c)	0-2	2-4	4-6	6-8	8-12
$\Delta\eta = 1.1$	2.2σ	6.3σ	7.4σ	5.0σ	2.8σ
$\Delta\eta = 5.3$	1.4σ	6.2σ	5.0σ	3.3σ	1.3σ



- ▶ The contribution of J/ψ from (re)combination could lead to an **elliptic flow** signal at LHC → hints observed in run-1 results
- ▶ From hint to evidence for a non-zero v₂ signal, maximum for 4 < p_T < 6 GeV/c 20-40% centrality
- ▶ A significant fraction of observed J/ψ comes from charm quarks which thermalized in the QGP

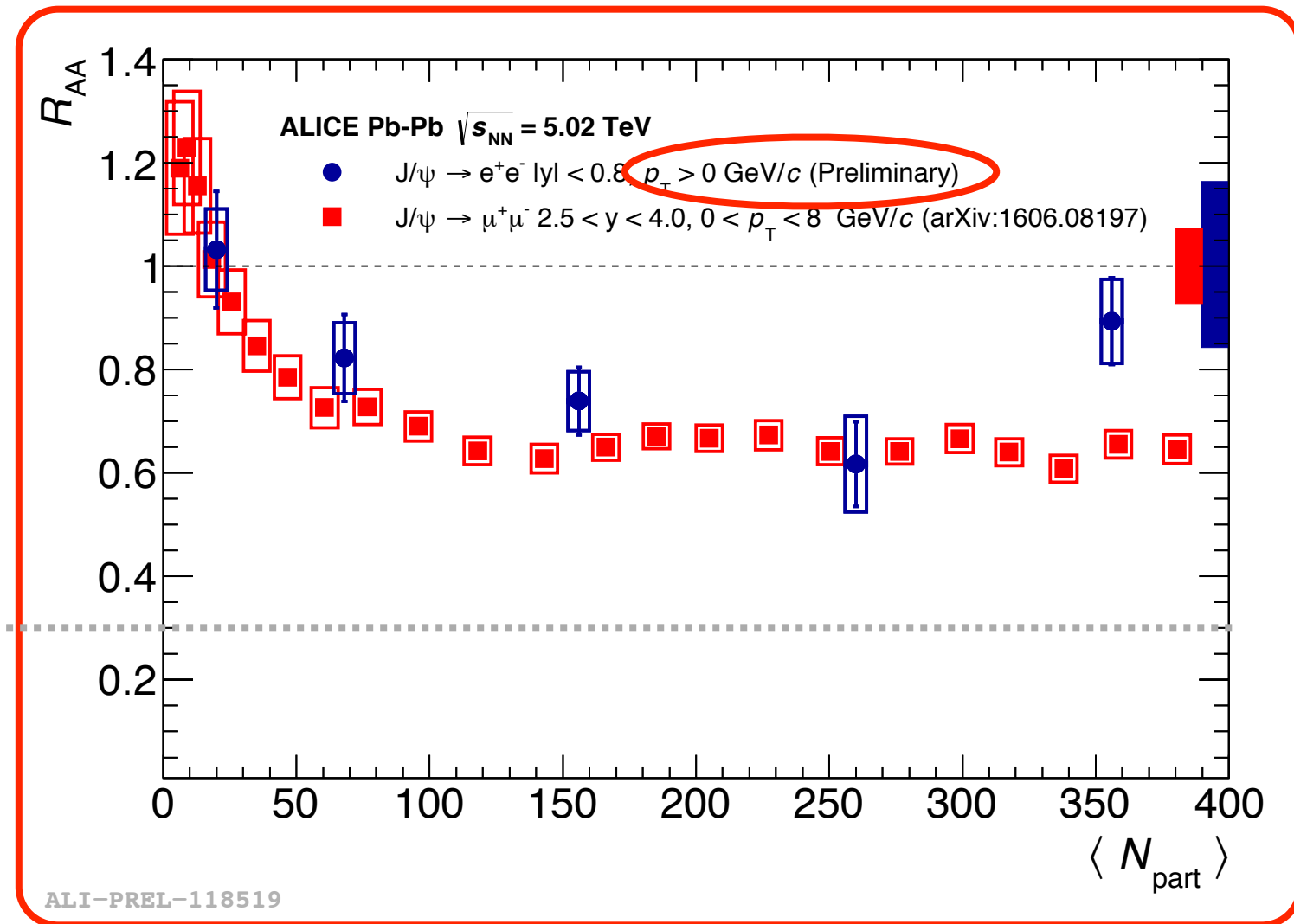
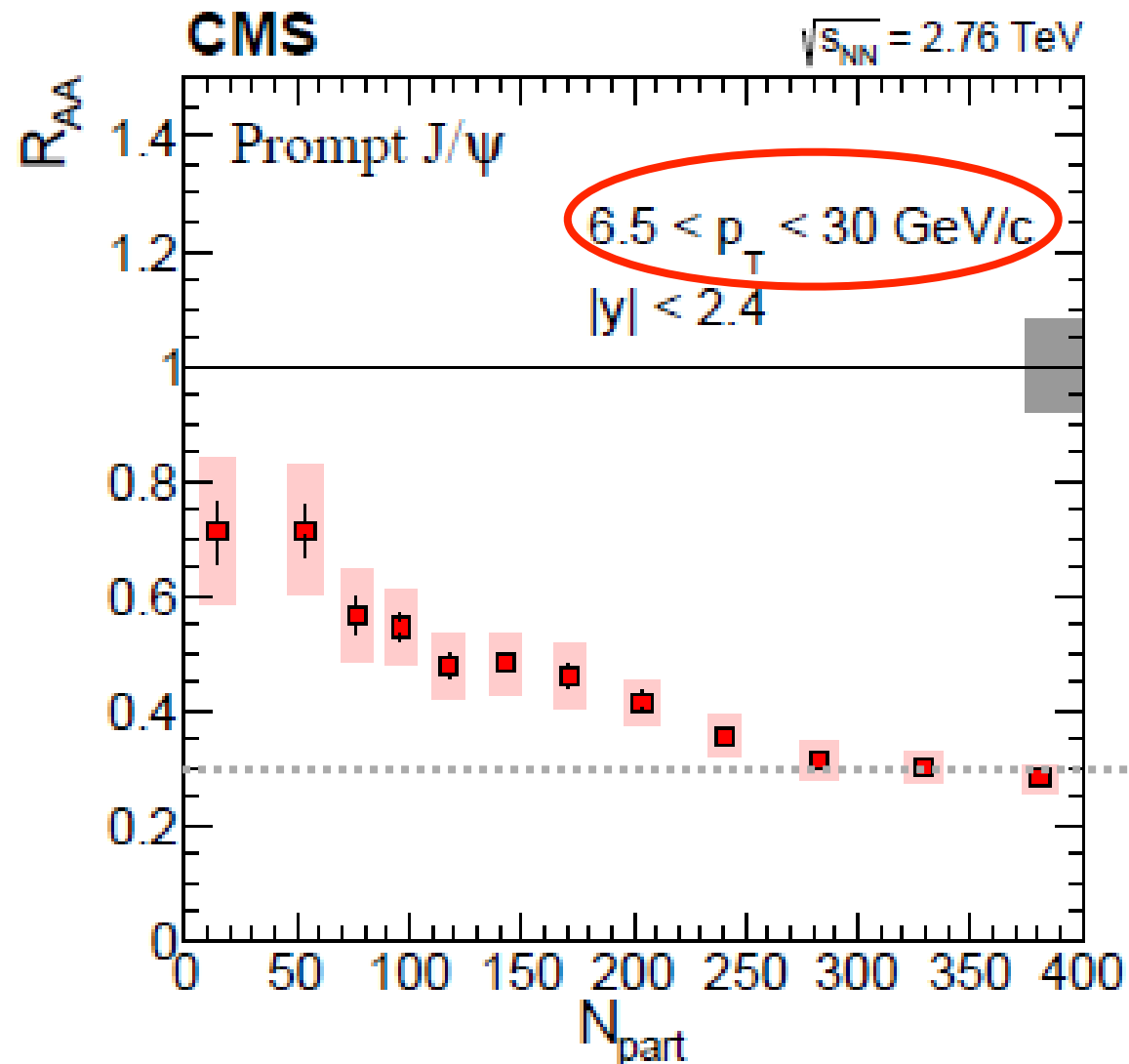


Comparison **closed vs open charm**
 → Learn about **light vs heavy**
quark flow

- ▶ The contribution of J/ψ from (re)combination could lead to an **elliptic flow** signal at LHC → hints observed in run-1 results
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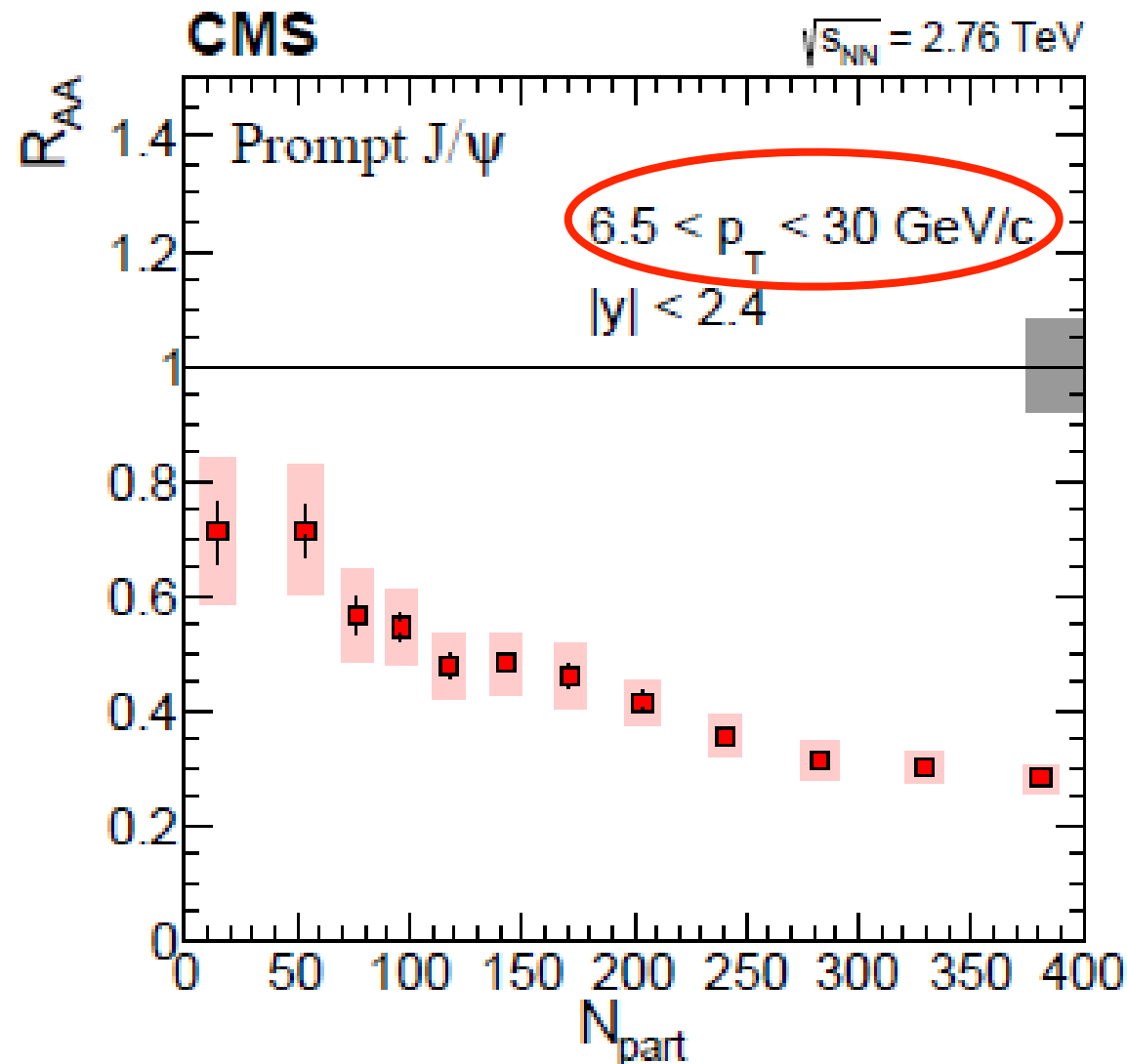
High- p_T J/ψ

V. Khachatryan et al. (CMS), arXiv:
1610.00613

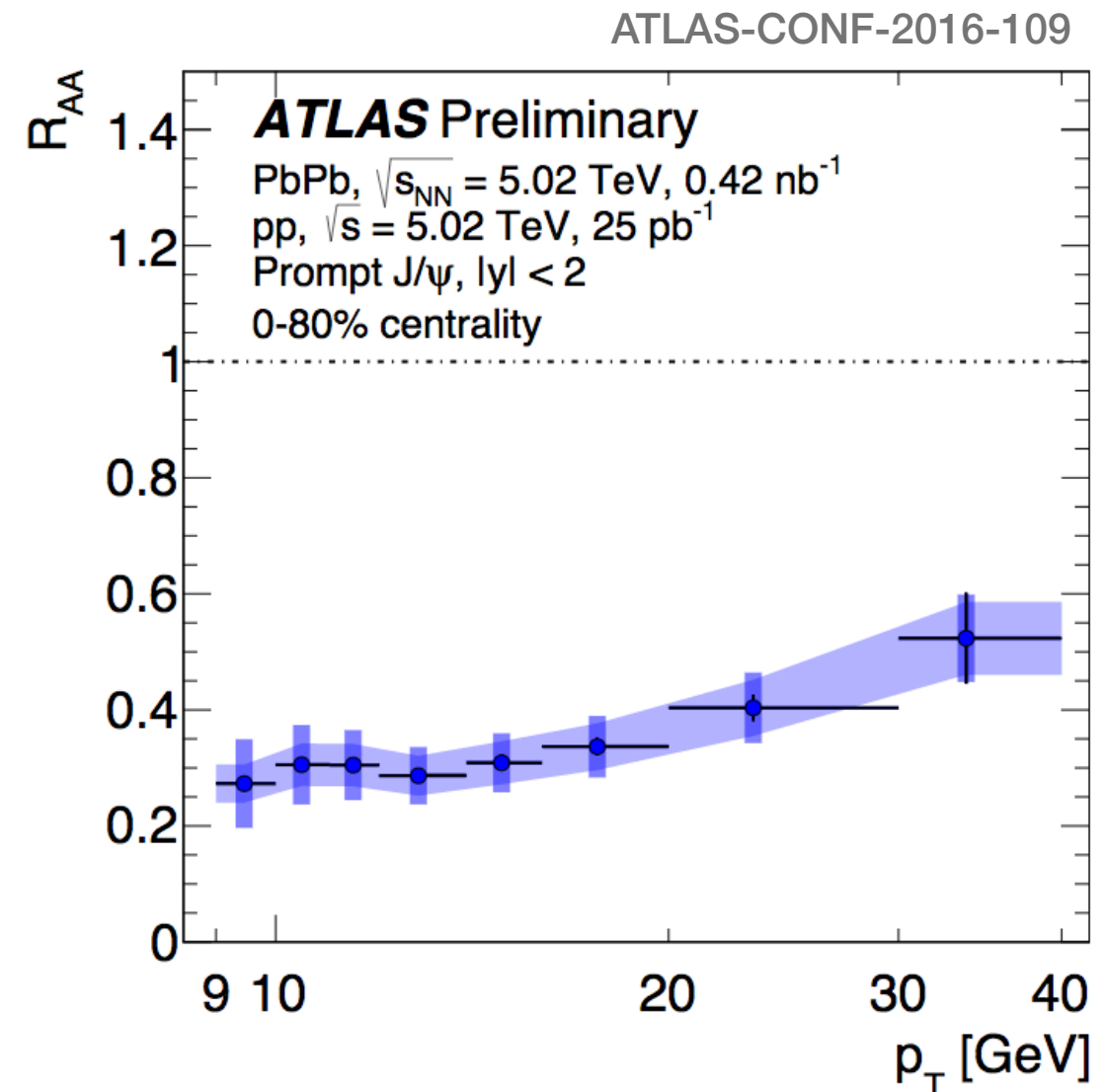


- Striking difference with respect to low- p_T J/ψ
- Suppression increases with centrality at high p_T , down to $R_{AA} \sim 0.3$

V. Khachatryan et al. (CMS), arXiv:
1610.00613

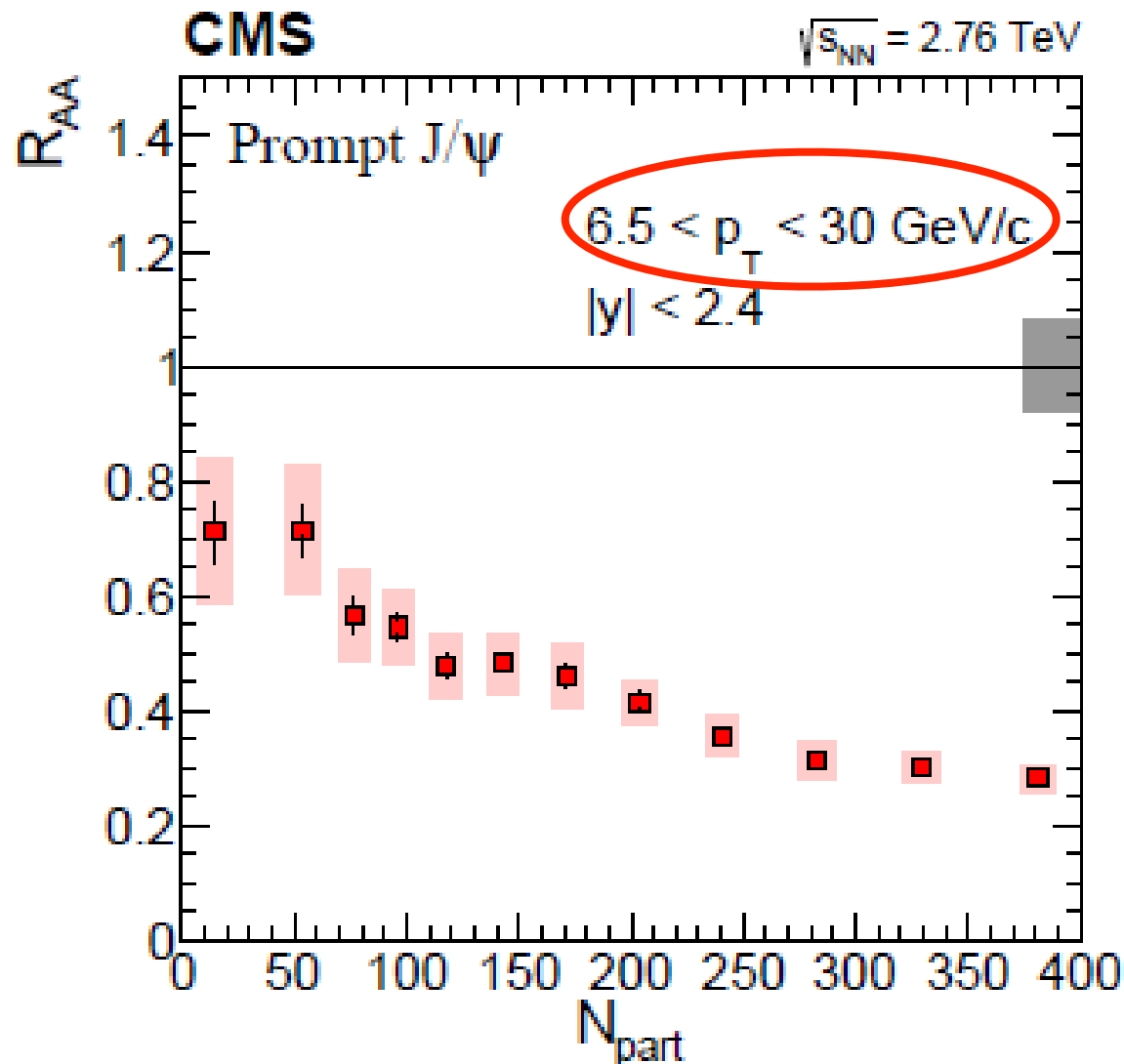


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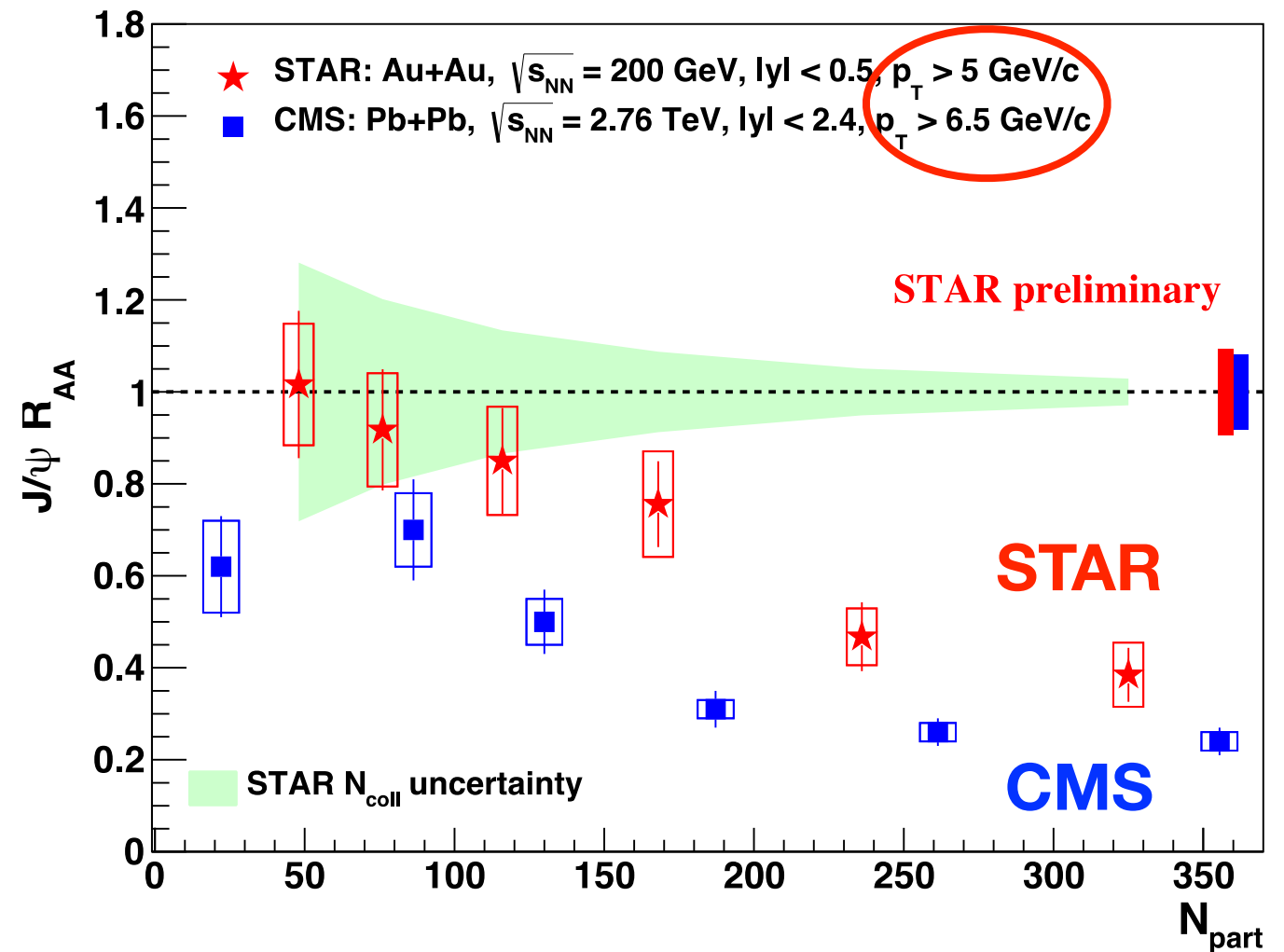


- ▶ R_{AA} increases for $p_T > 20$ GeV/c
- ▶ Related to energy loss effects, rather than dissociation?

V. Khachatryan et al. (CMS), arXiv: 1610.00613

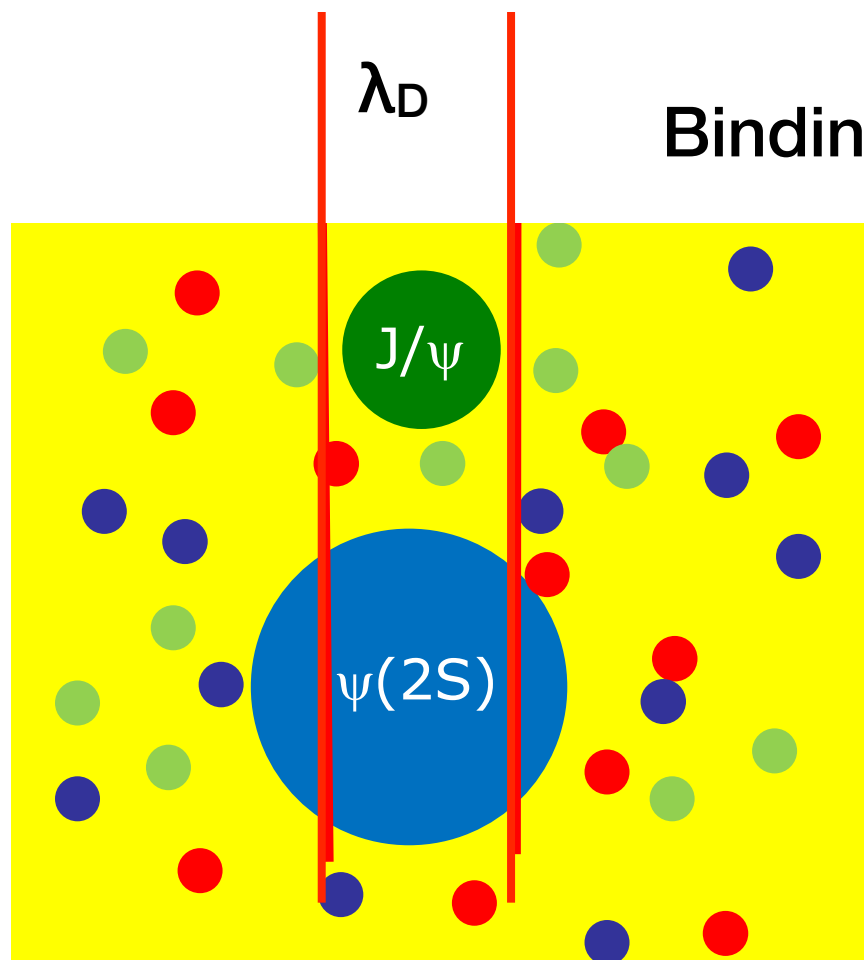


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- ▶ R_{AA} increases for $p_T > 20$ GeV/c
- ▶ Related to energy loss effects, rather than dissociation?
- ▶ $R_{AA}^{LHC} < R_{AA}^{RHIC} \rightarrow$ weak regeneration?

$\psi(2s)$ in Pb-Pb



Binding energy $\sim (2m_D - m_\psi) \rightarrow \psi(2S) \sim 60 \text{ MeV}, J/\psi \sim 640 \text{ MeV}$

- ▶ Expect **much stronger dissociation** effects for the **weakly bound $\psi(2S)$ state**
- ▶ Effect of re-combination on $\psi(2S)$ more subtle
→ important when the system is more diluted
⇒ **Important to test models!**

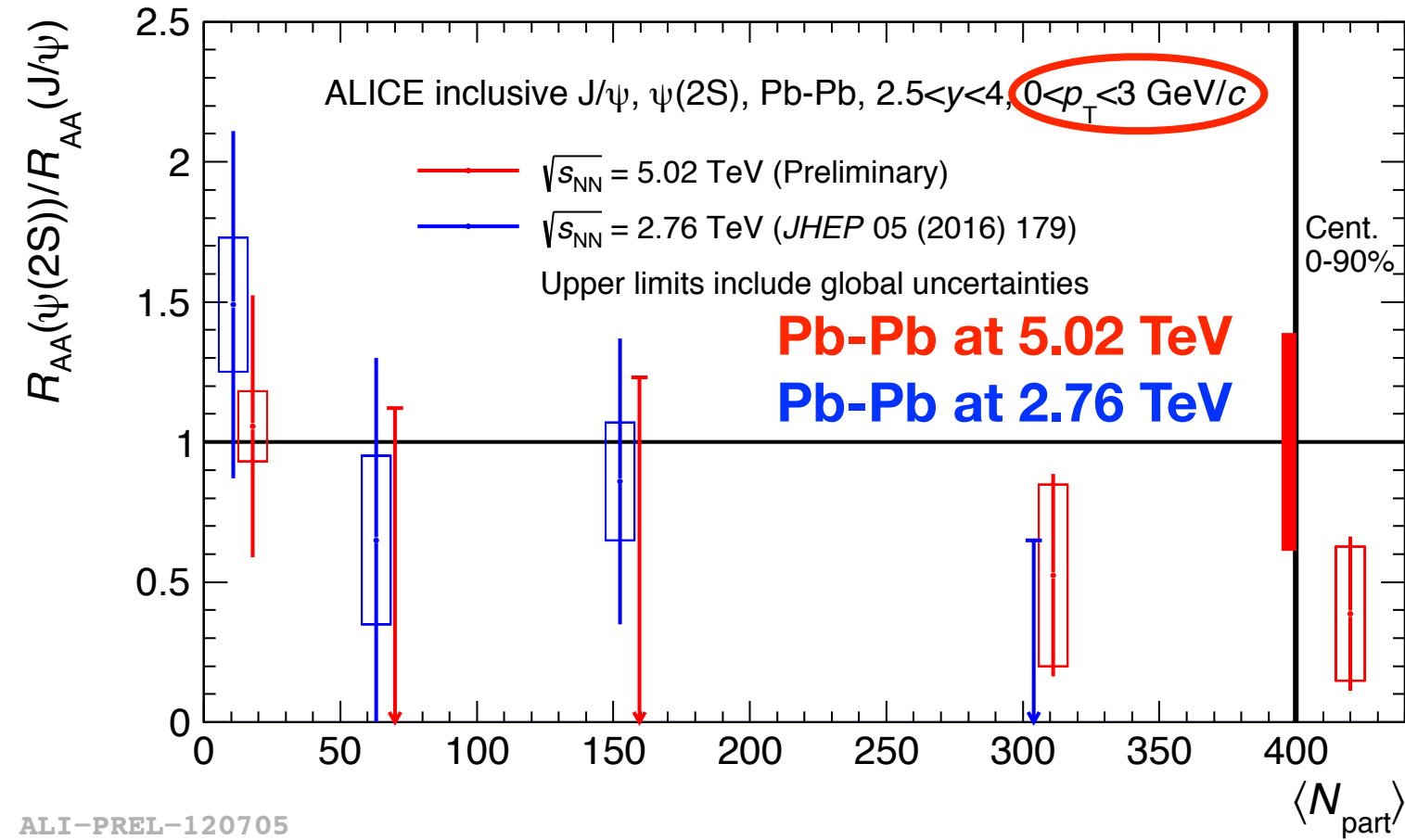
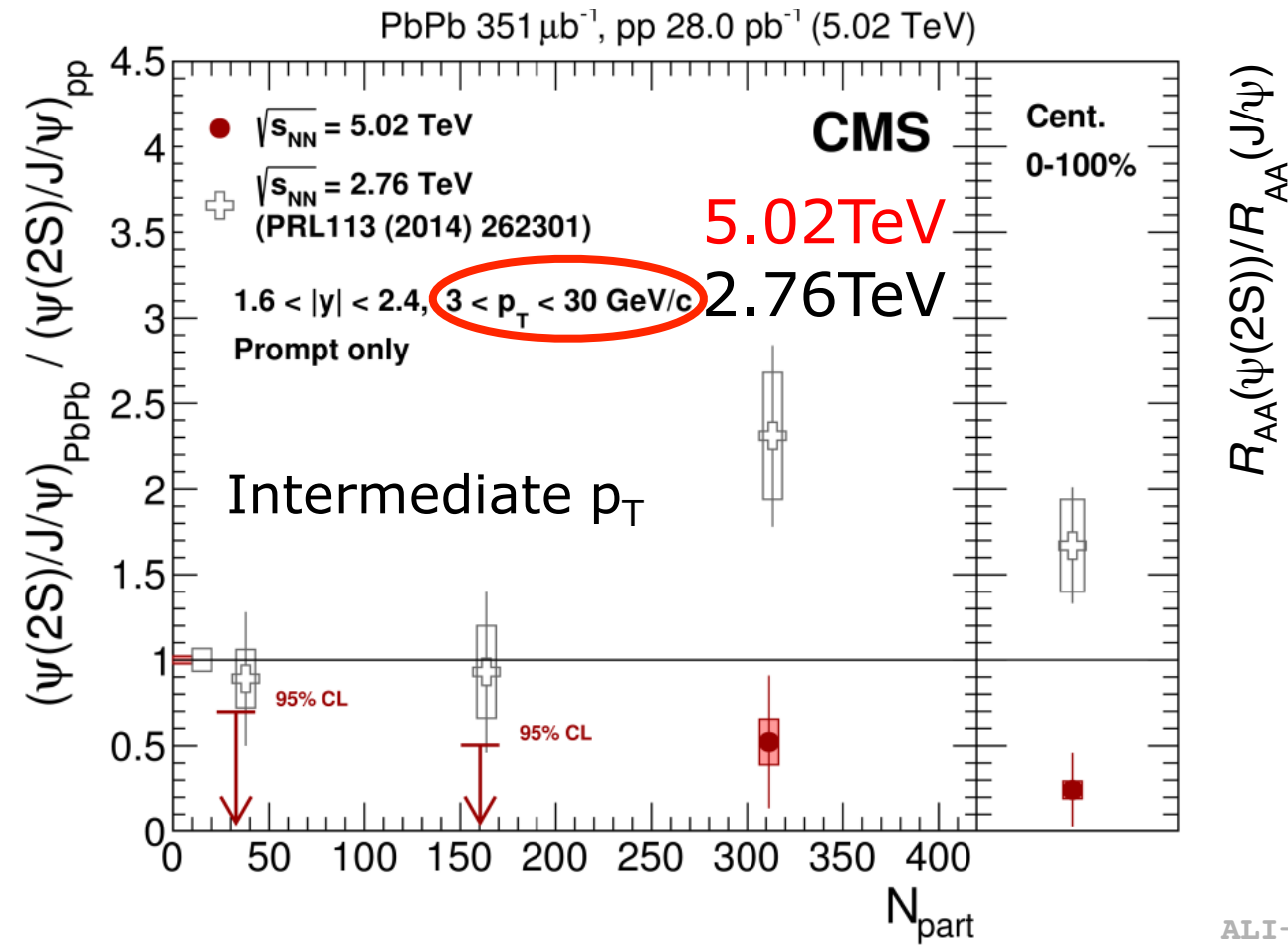
state	J/ψ	χ_c	$\psi(2S)$
Mass(GeV)	3.10	3.51	3.69
ΔE (GeV)	0.64	0.22	0.05
r_o (fm)	0.50	0.72	0.90

state	$Y(1S)$	$Y(2S)$	$Y(3S)$
Mass(GeV)	9.46	10.0	10.36
ΔE (GeV)	1.10	0.54	0.20
r_o (fm)	0.28	0.56	0.78

(Digal, Petrecki, Satz PRD 64(2001) 0940150)

Double ratios $\psi(2s)/J/\psi$

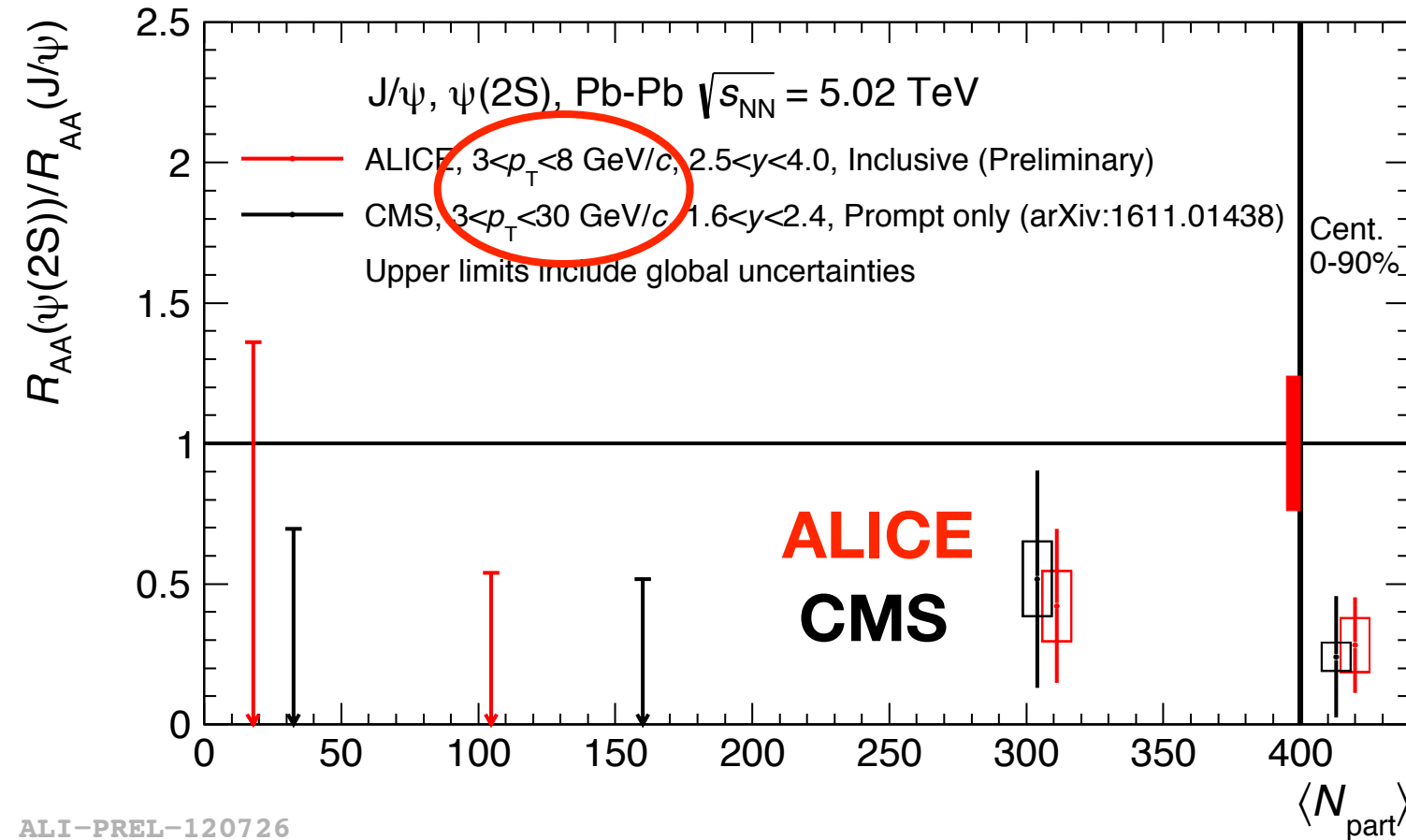
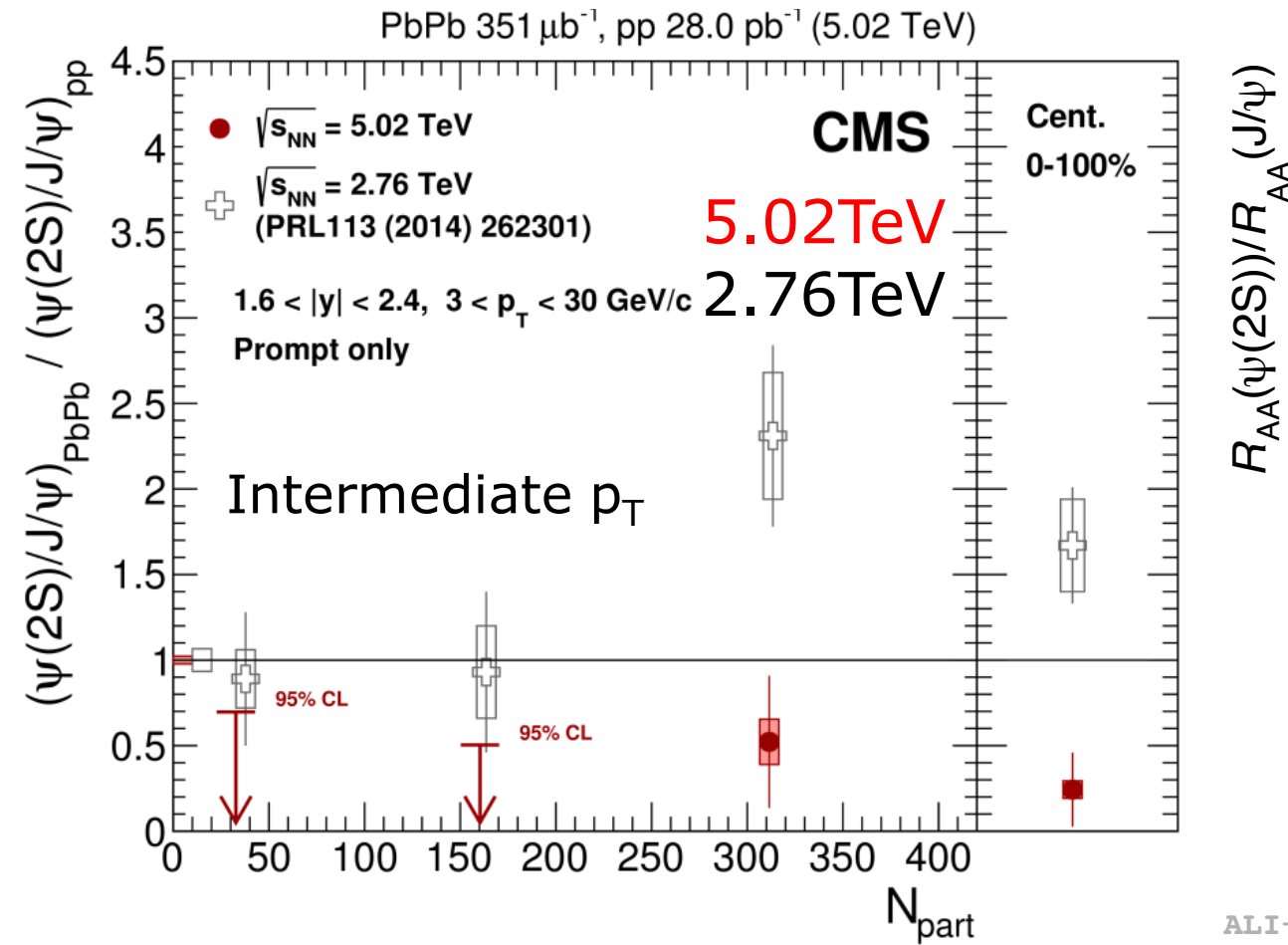
V. Khachatryan et al. (CMS), arXiv:
1611.01438



- ▶ $(\psi(2S)/J/\psi)_{\text{PbPb}} / (\psi(2S)/J/\psi)_{\text{pp}} \rightarrow \ll 1$ in a dissociation scenario
- ▶ CMS (intermediate p_{T}), **enhancement** to **suppression** for increasing $\sqrt{s_{\text{NN}}}$
- ▶ ALICE extends down to $p_{\text{T}}=0$, suppression is seen
- ▶ Proposed mechanism (Rapp arXiv:1609.04868) for enhancement: **$\psi(2S)$ regeneration mainly occurring later**, when radial flow is already built-up

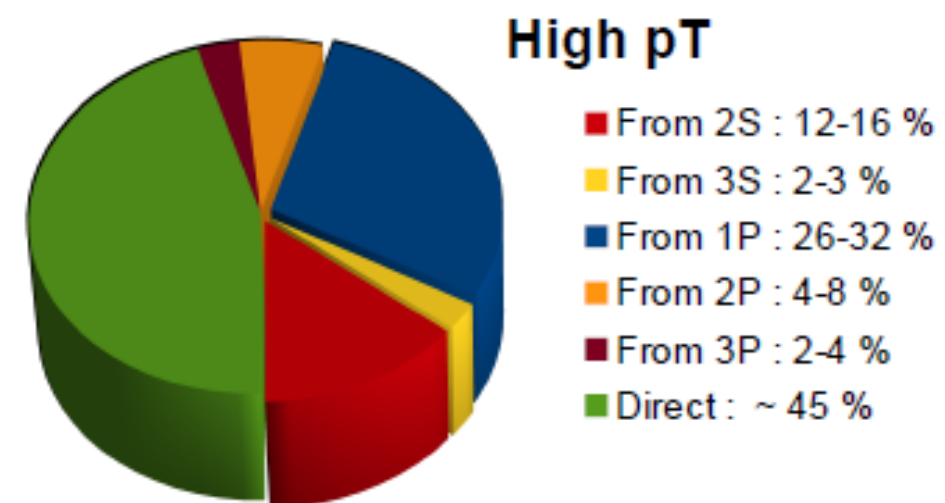
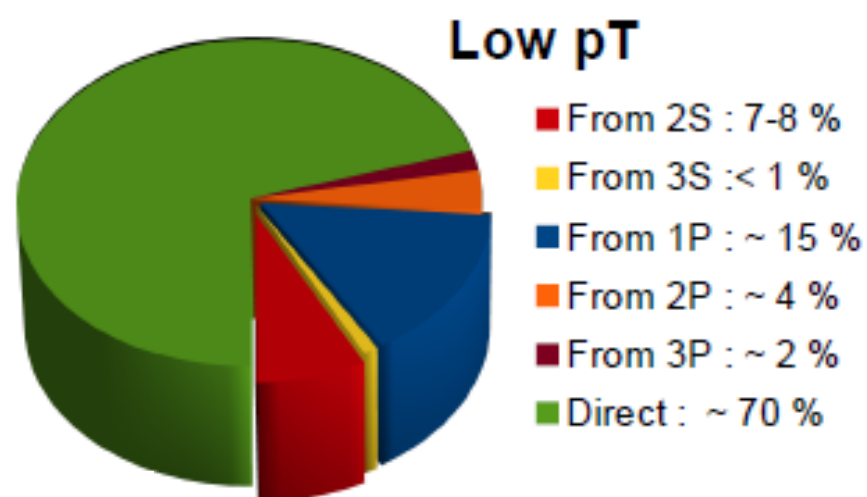
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- ▶ $(\psi(2S)/J/\psi)_{\text{PbPb}} / (\psi(2S)/J/\psi)_{\text{pp}} \rightarrow \ll 1$ in a dissociation scenario
- ▶ CMS (intermediate p_{T}), **enhancement** to **suppression** for increasing $\sqrt{s_{\text{NN}}}$
- ▶ ALICE extends down to $p_{\text{T}}=0$, suppression is seen
- ▶ Good compatibility at $\sqrt{s_{\text{NN}}}=5.02 \text{ TeV}$ in the common p_{T} range
- ▶ Proposed mechanism (Rapp arXiv:1609.04868) for enhancement: **$\psi(2S)$ regeneration mainly occurring later**, when radial flow is already built-up

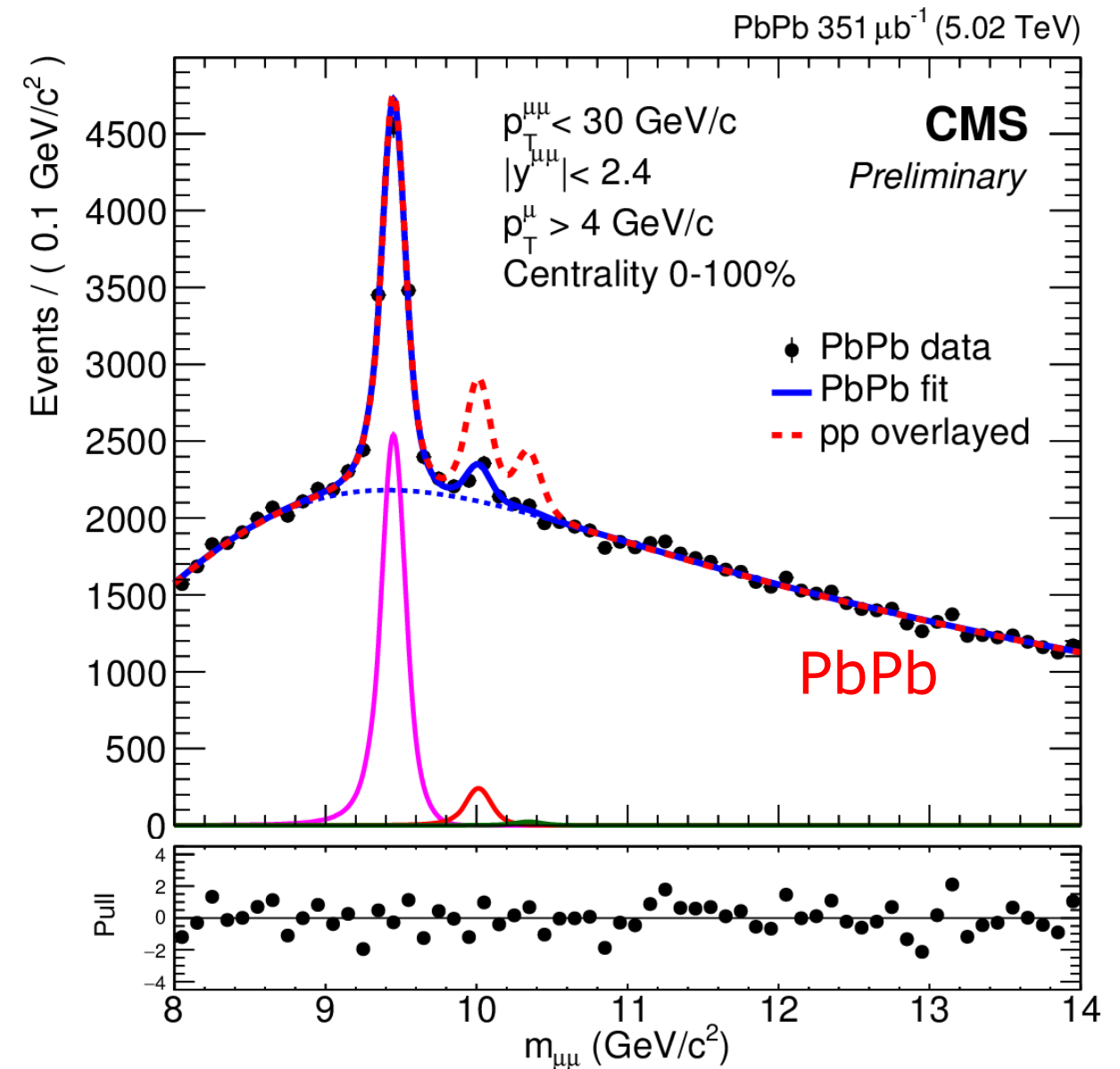
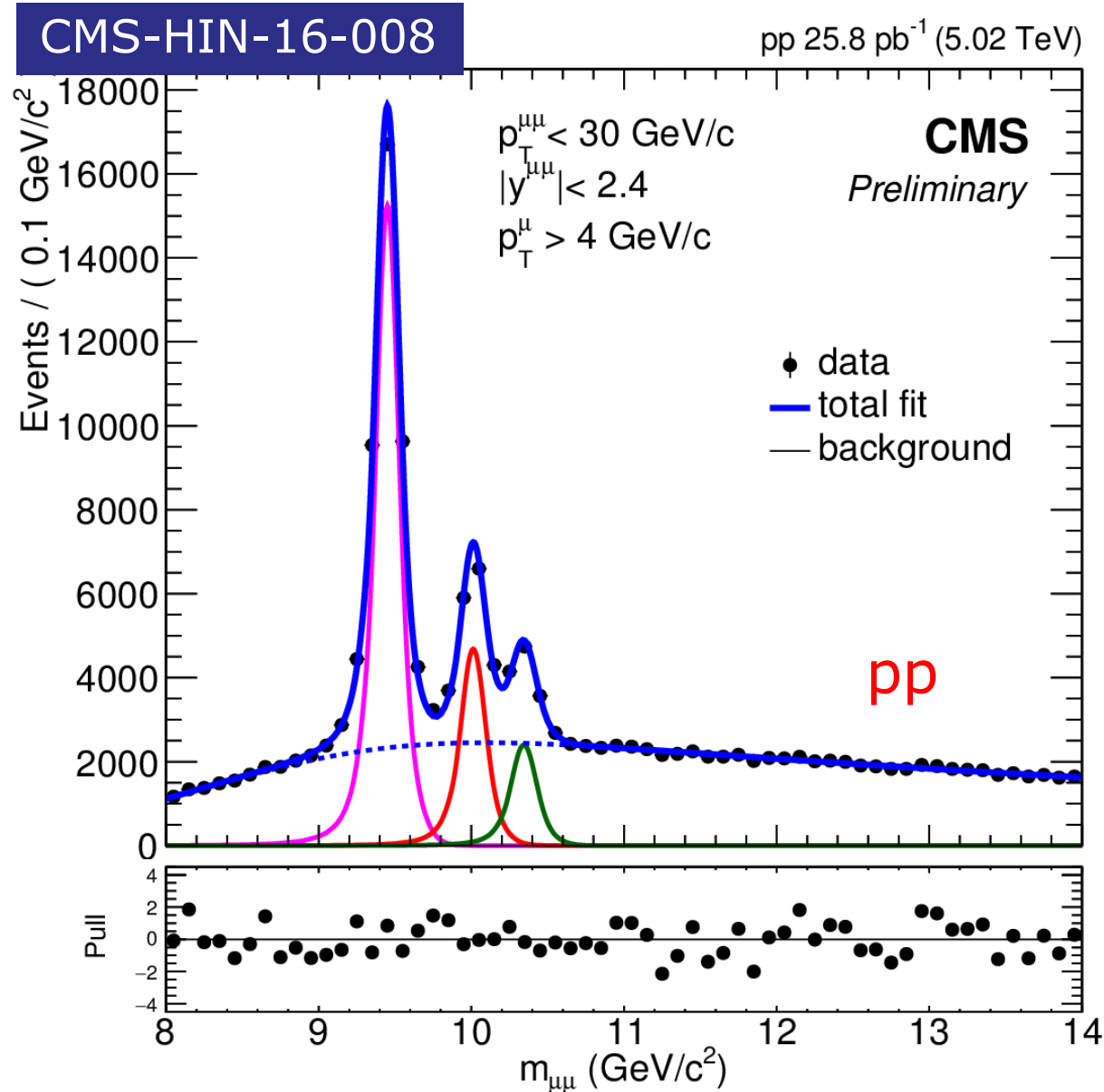
- Re-combination effects not strong → simpler interpretation?
- $Y(1S)$ very strongly bound, $E_b=(2m_B-m_{Y(1S)}) \sim 1100$ MeV → probe of hot QGP
- Together with $Y(2S)$ ($E_b \sim 500$ MeV) and $Y(3S)$ ($E_b \sim 200$ MeV) → provide (very) different sensitivity to the medium
- Caveats
 - 1) Realistic theory description anyway not straightforward
 - 2) The feed-down structure of the bottomonium sector is not trivial → has an impact on the interpretation of the result



A. Andronic et al., Eur. Phys. J. C 76 (2016) 107

Bottomonium (sequential) suppression

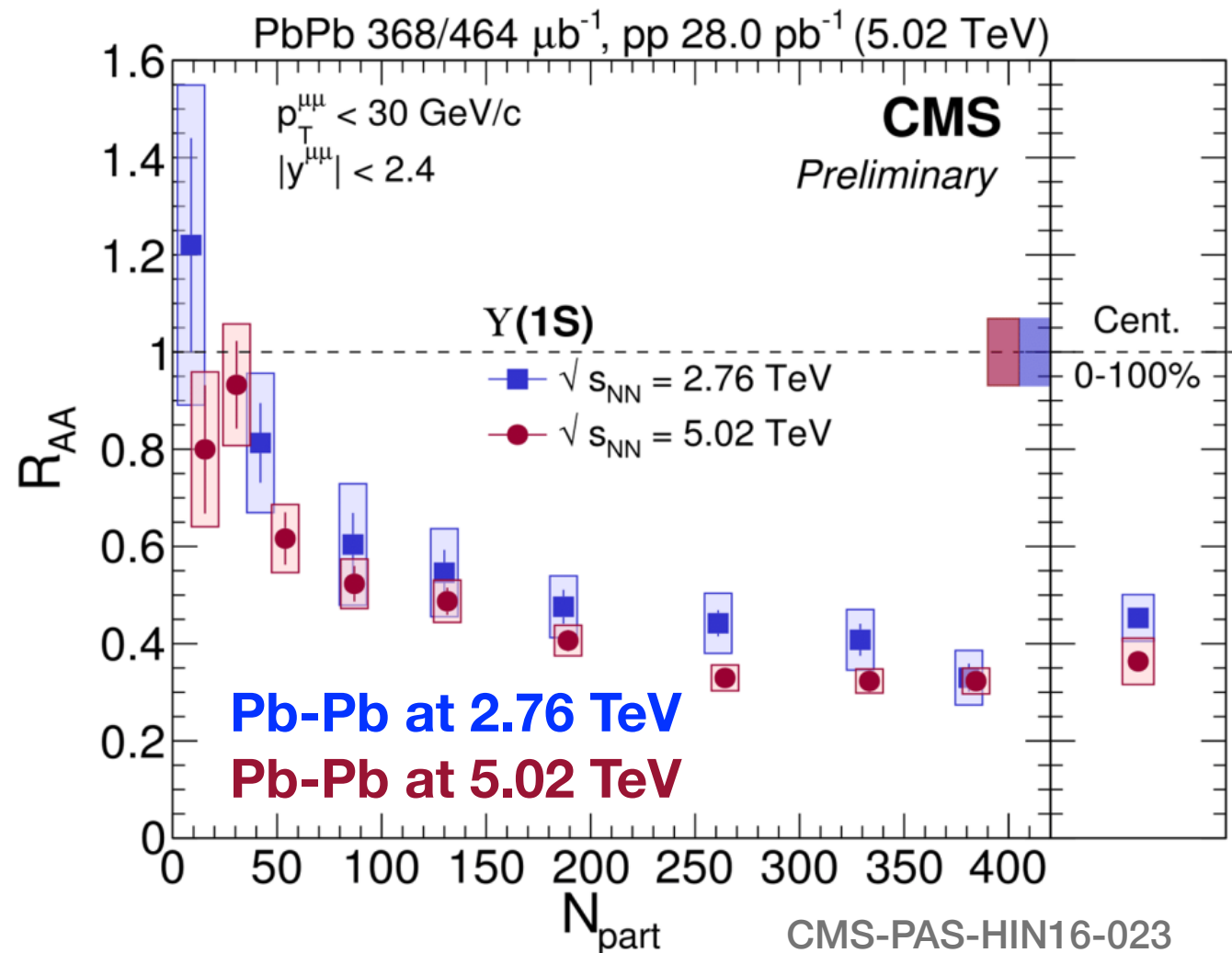
- Quarkonia in HI at the LHC - Sequential suppression?



- Recent **CMS results at $\sqrt{s}=5.02$ TeV** confirm the Y(2S,3S) suppression relative to the strongly bound Y(1S)!

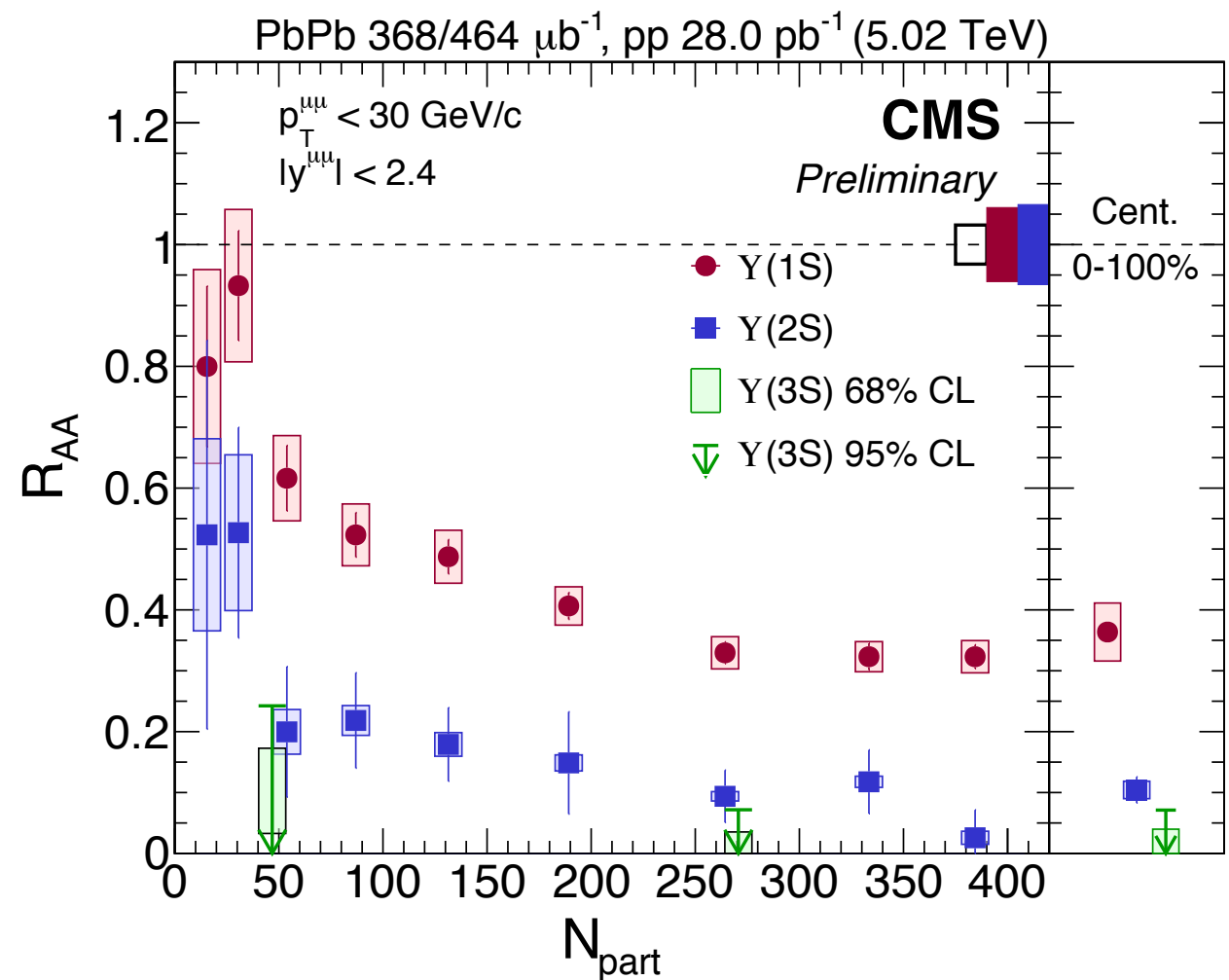
R_{AA} of $Y(1S)$ and $Y(2S)$

- $\sqrt{s_{NN}} = 2.76$ TeV, strong centrality dependence, up to factor ~ 2 and ~ 8 suppression for $Y(1S)$ and $Y(2S)$, respectively



CMS-PAS-HIN16-023

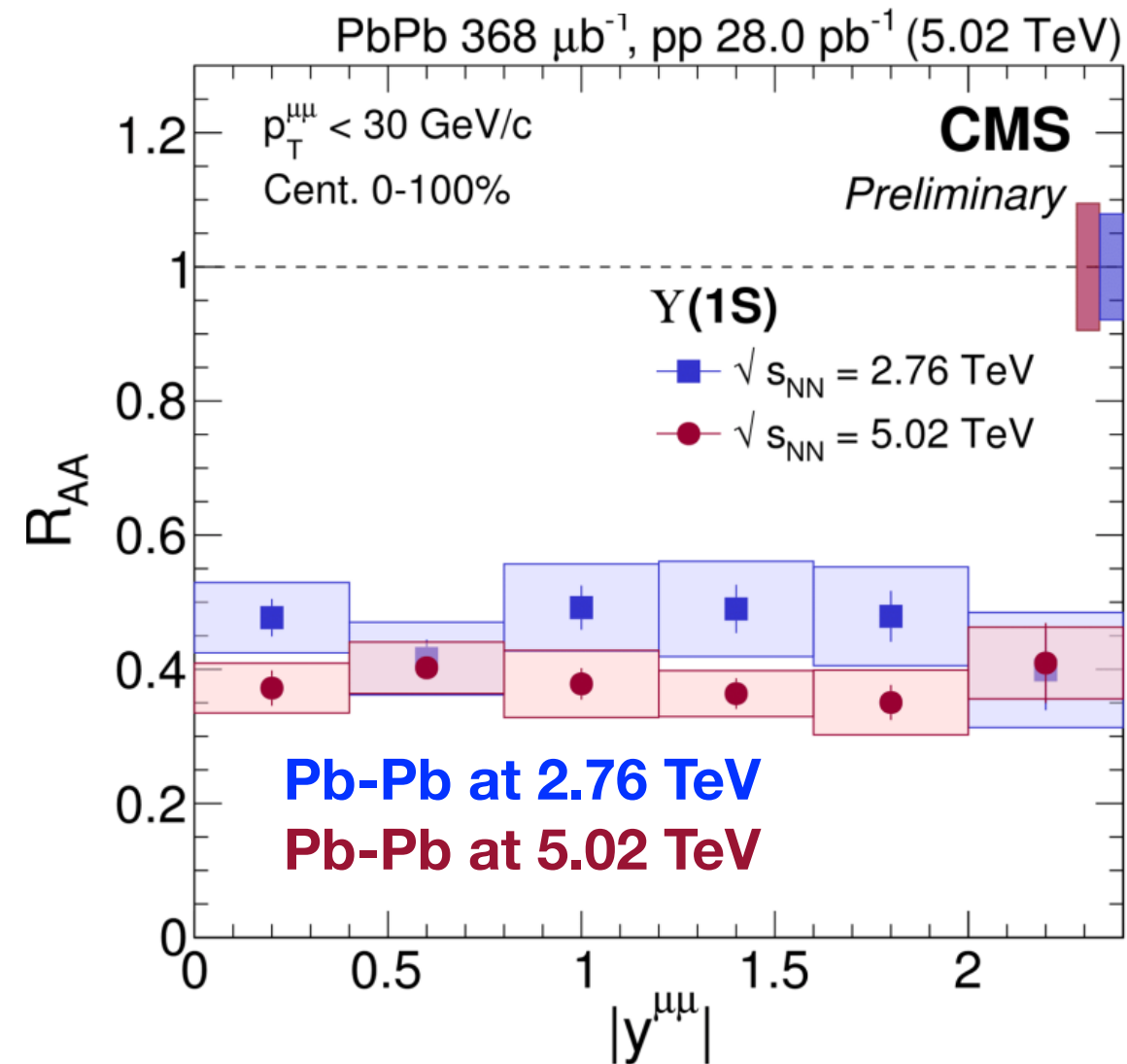
V. Khachatryan et al. (CMS), arXiv:
1611.01510



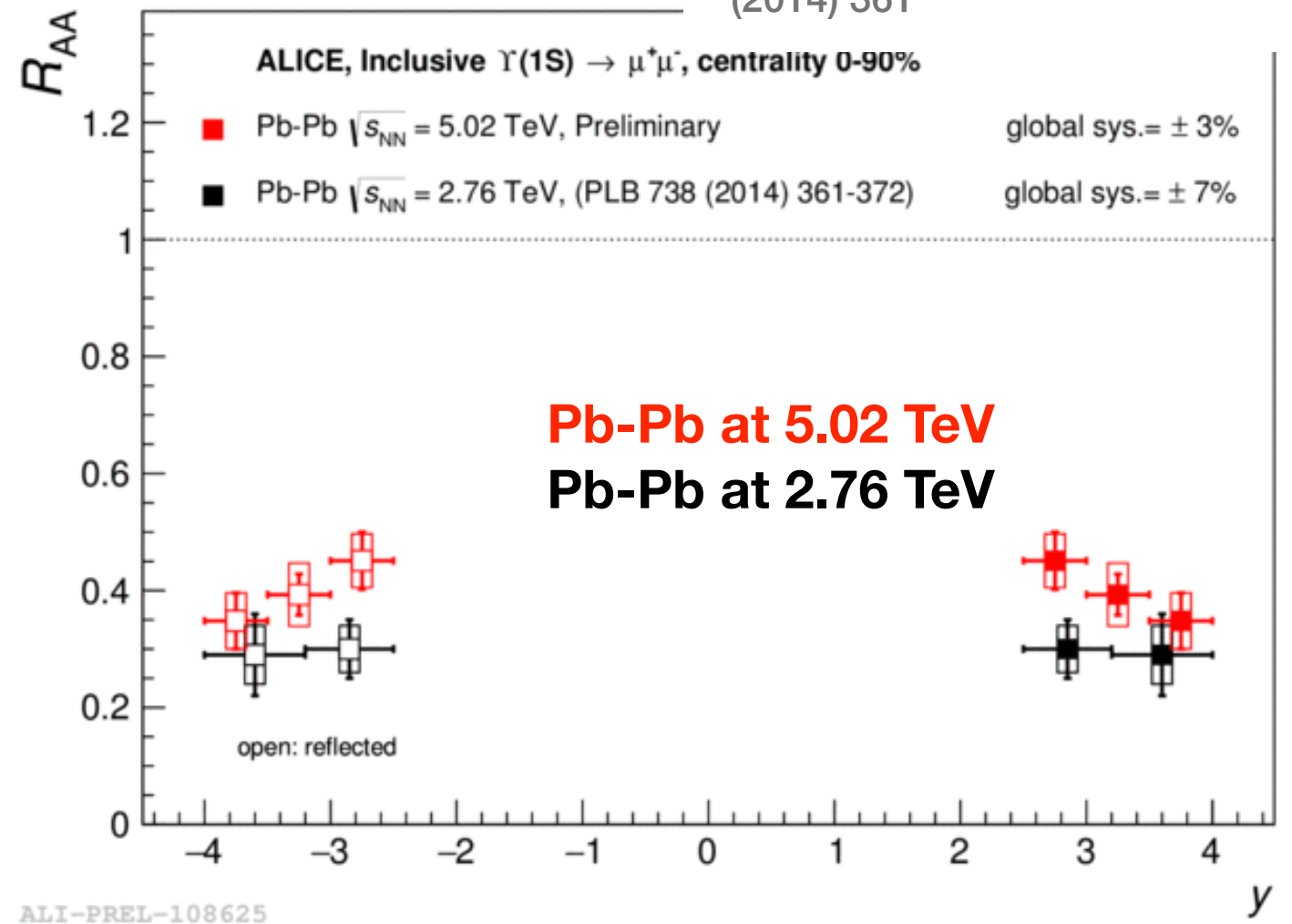
- New CMS results at $\sqrt{s_{NN}} = 5.02$ TeV
→ Indications for slightly stronger suppression

R_{AA} vs y of $Y(1S)$

CMS-PAS-HIN16-023

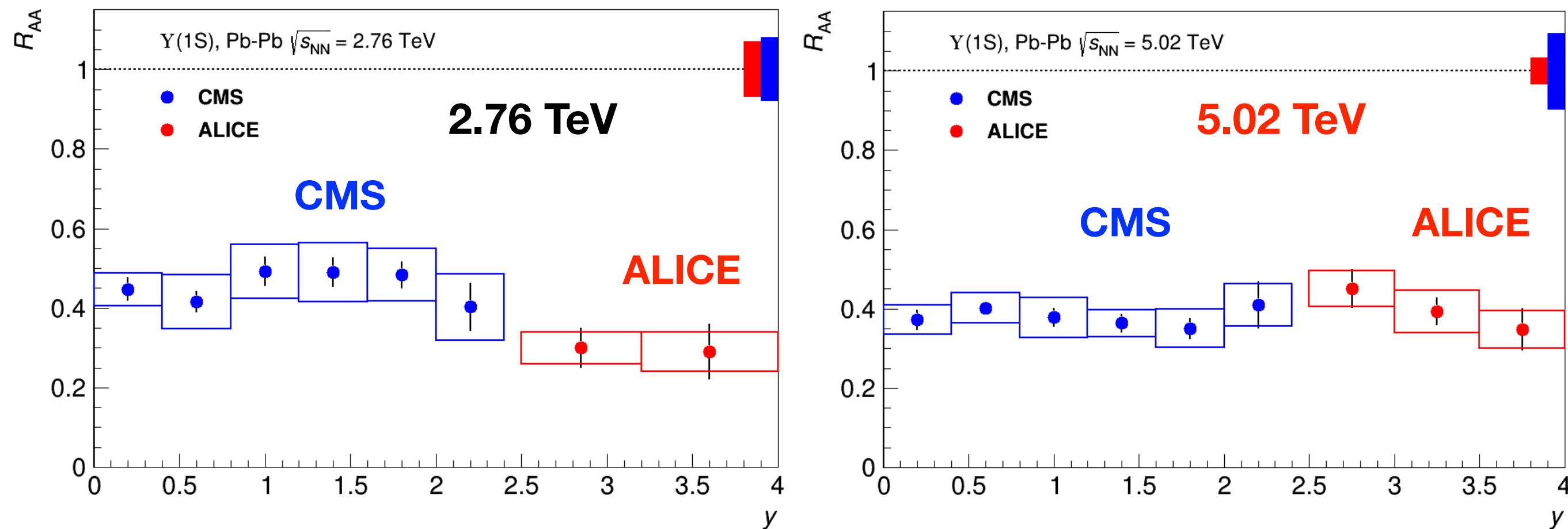


B. Abelev et al., (ALICE) PLB738 (2014) 361



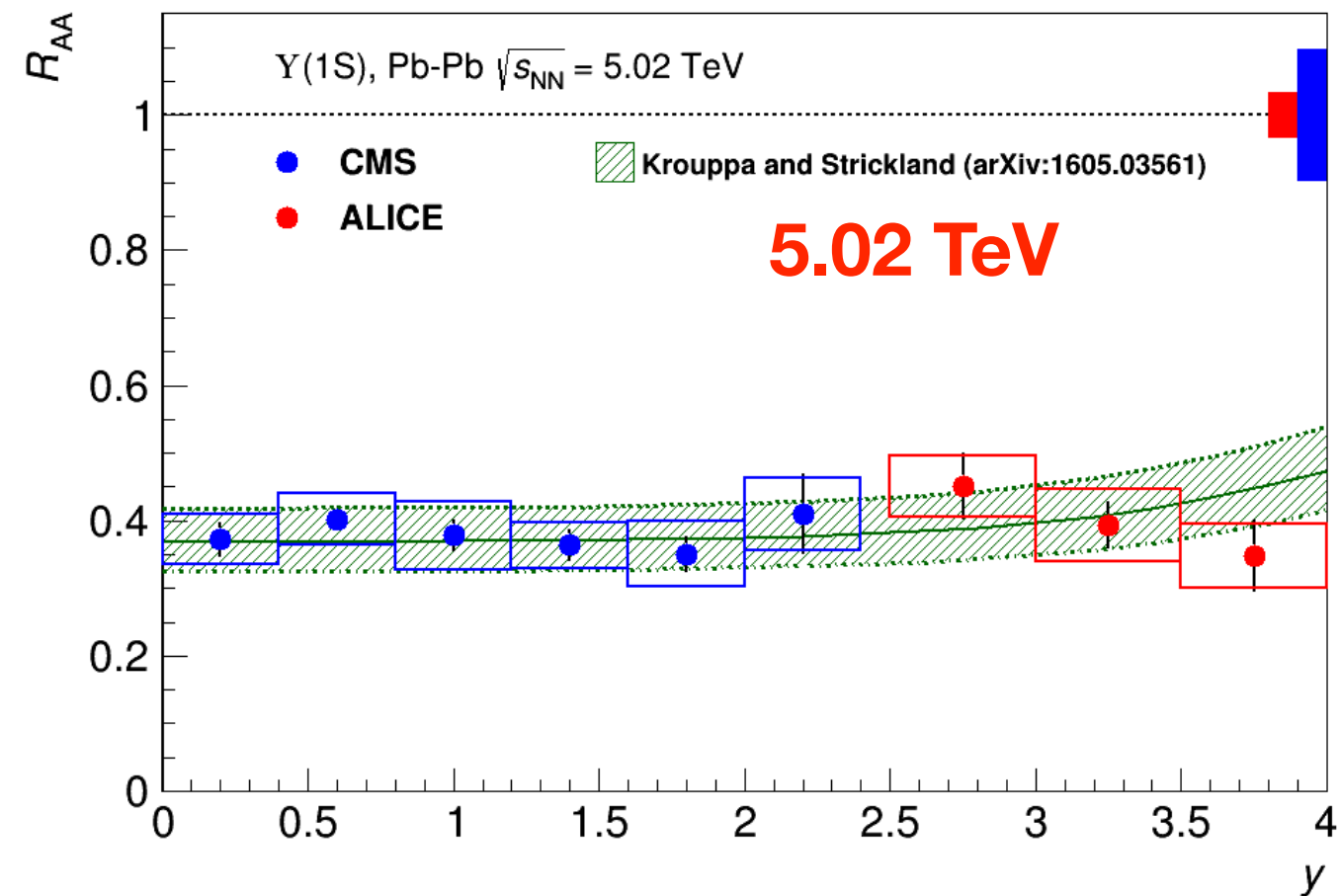
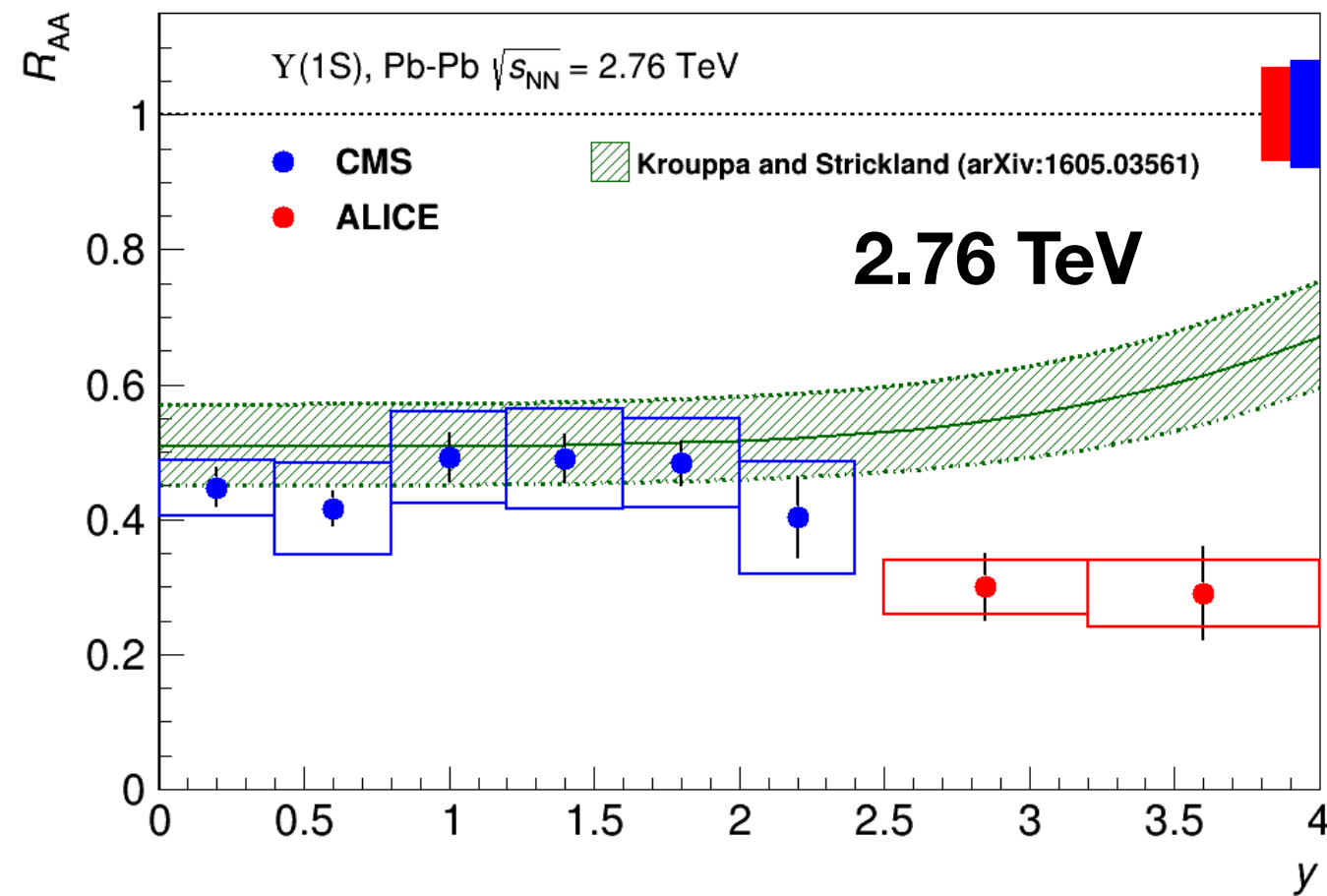
- CMS \rightarrow hints for **more suppression** at $\sqrt{s_{NN}} = 5.02 \text{ TeV}$
- ALICE \rightarrow hints for **less suppression** at $\sqrt{s_{NN}} = 5.02 \text{ TeV}$
- Compare R_{AA} vs y for the two experiments in a single plot

R_{AA} vs y of $Y(1S)$



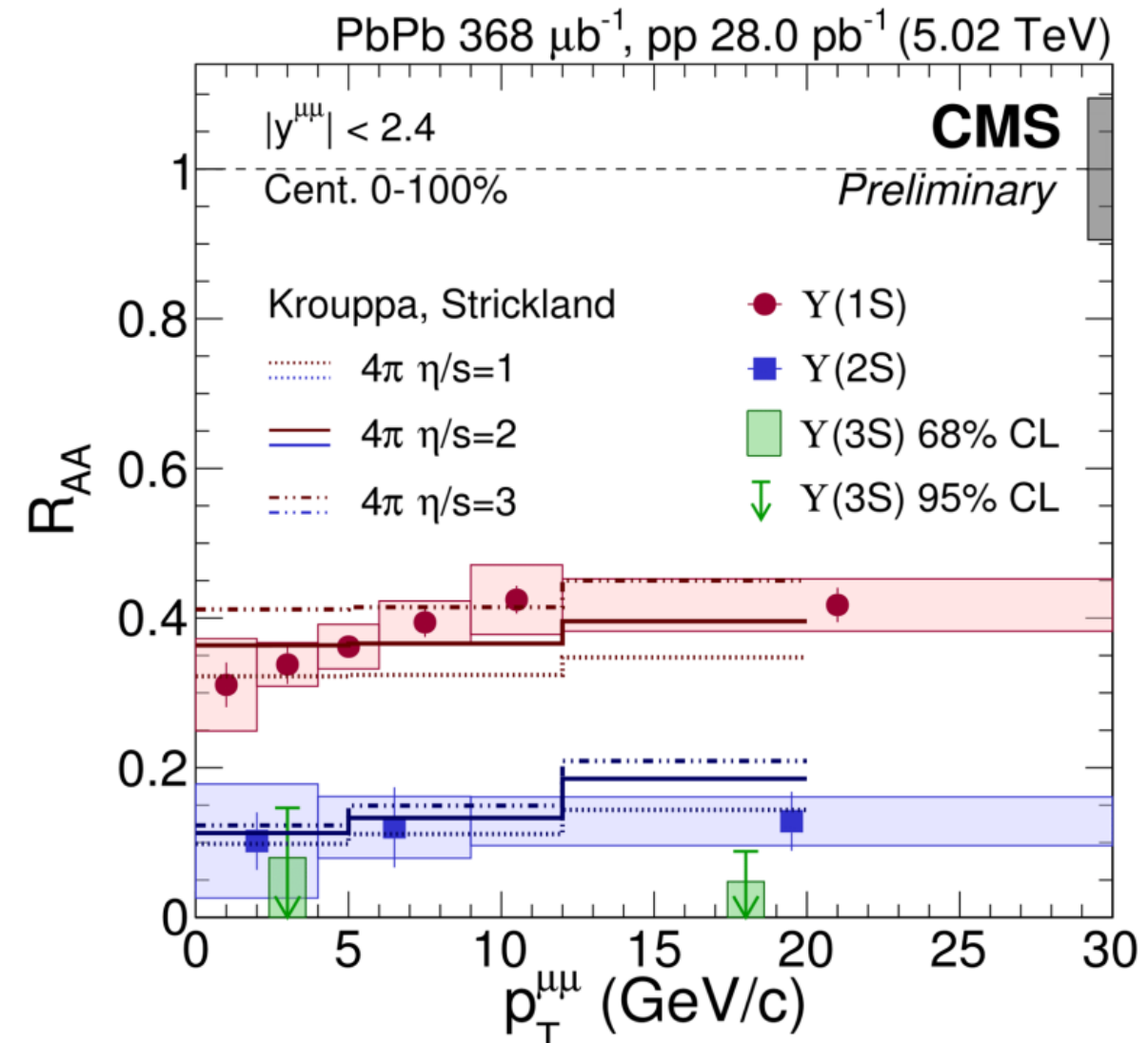
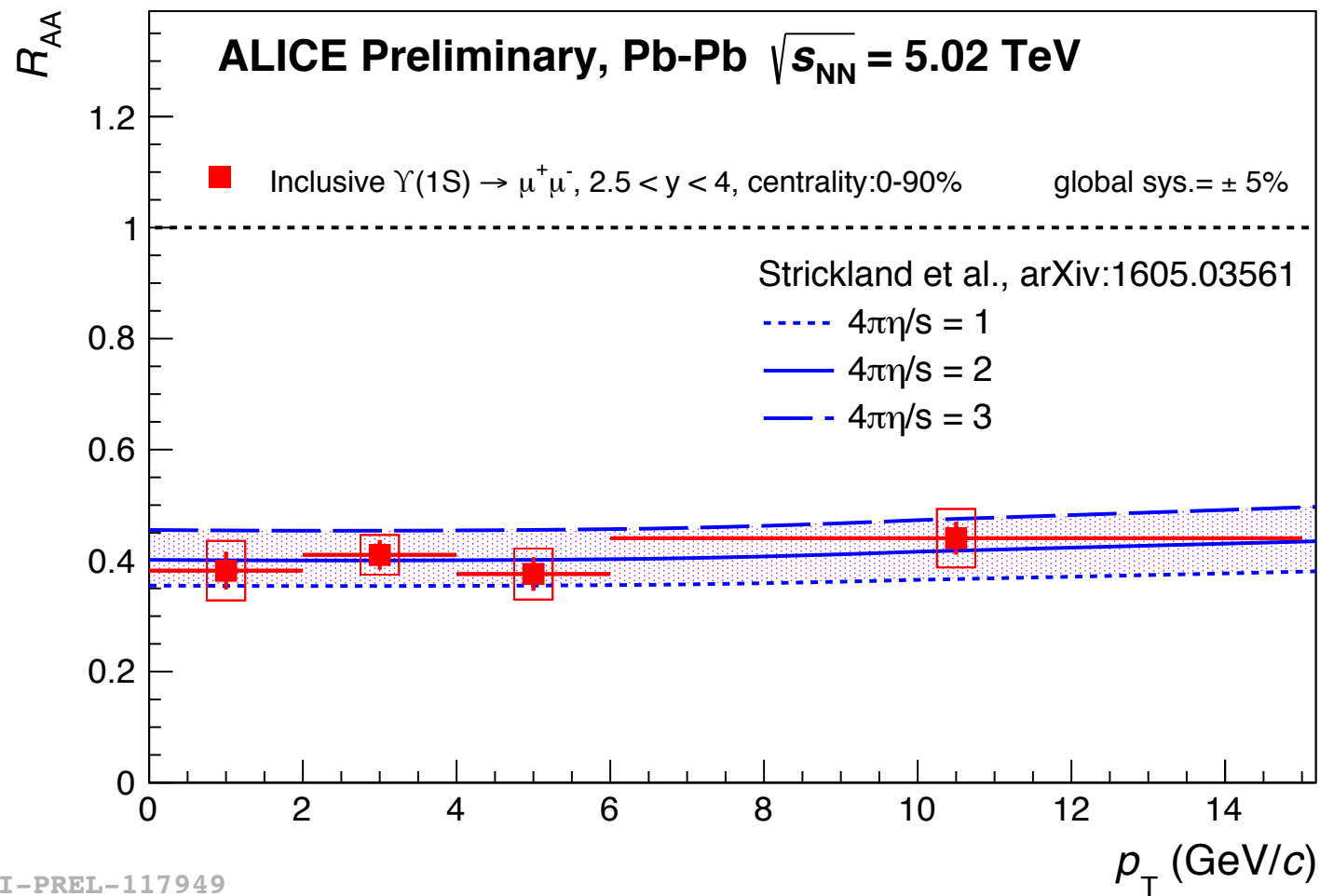
- Suppression increases with rapidity at $\sqrt{s_{NN}}=2.76$ TeV
- Suppression constant vs rapidity at $\sqrt{s_{NN}}=5.02$ TeV
- $\sqrt{s_{NN}}=2.76$ TeV: typical features of a (re)generation pattern, which seems to vanish at $\sqrt{s_{NN}}=5.02$ TeV
- Systematic uncertainties not negligible

R_{AA} vs y of $Y(1S)$



- Suppression increases with rapidity at $\sqrt{s_{NN}}=2.76$ TeV
- Suppression constant vs rapidity at $\sqrt{s_{NN}}=5.02$ TeV
- $\sqrt{s_{NN}}=2.76$ TeV: typical features of a (re)generation pattern, which seems to vanish at $\sqrt{s_{NN}}=5.02$ TeV
- Systematic uncertainties not negligible
- Can the rapidity-dependence of CNM effects play a role? Not likely

R_{AA} vs p_T of $Y(1S)$



CMS-PAS-HIN16-023

- Both CMS and ALICE measure weak or no dependence of R_{AA} vs p_T
- Fair agreement with theoretical model (Strickland)

- Charmonia (J/ψ , $\psi(2S)$)

Firm evidence for J/ψ elliptic flow and strong re-generation effects

→ Charm quarks thermalization in the deconfined medium

- Bottomonia ($Y(1S)$, $Y(2S)$, $Y(3S)$)

Suppression effects strongly correlated with binding energy

→ Evidence for resonance melting in a hot QGP

Thank you for your attention!

Extra Slides

Experimental evidence for direct $\Upsilon(1S)$ suppression?

- Direct $\Upsilon(1S)$ suppression implies QGP temperatures at least $\sim 2 T_c$
- Experimental evidence for direct $\Upsilon(1S)$ suppression needs control over
 - Feed-down from S and P bottomonium states

Recent LHCb results imply a $\sim 30\%$ effect at (fairly) low p_T in pp

- Size of CNM effects \rightarrow weak but not precisely known

- Starting from CMS results and assuming all the remaining Pb-Pb $\Upsilon(1S)$ are direct

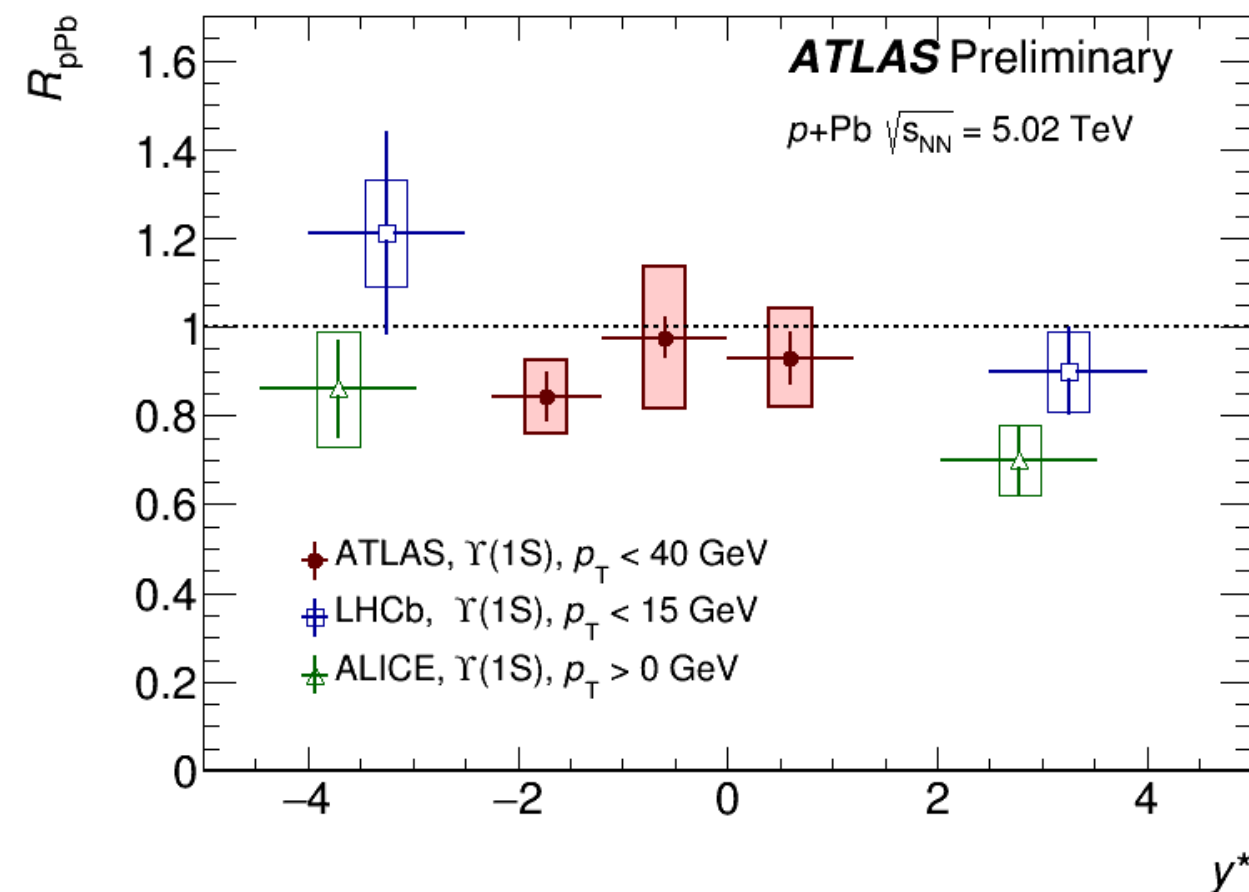
$$R_{AA}^{\text{incl}} \Upsilon(1S) \sim 0.36$$

$$R_{AA}^{\text{direct}} \Upsilon(1S) \sim 0.36/0.7 = 0.51$$

CNM effects (-1σ level)

$$\rightarrow (R_{pA} - 1\sigma)^2 \sim 0.8^2 = 0.64$$

- Experimental indication for direct $\Upsilon(1S)$ suppression!



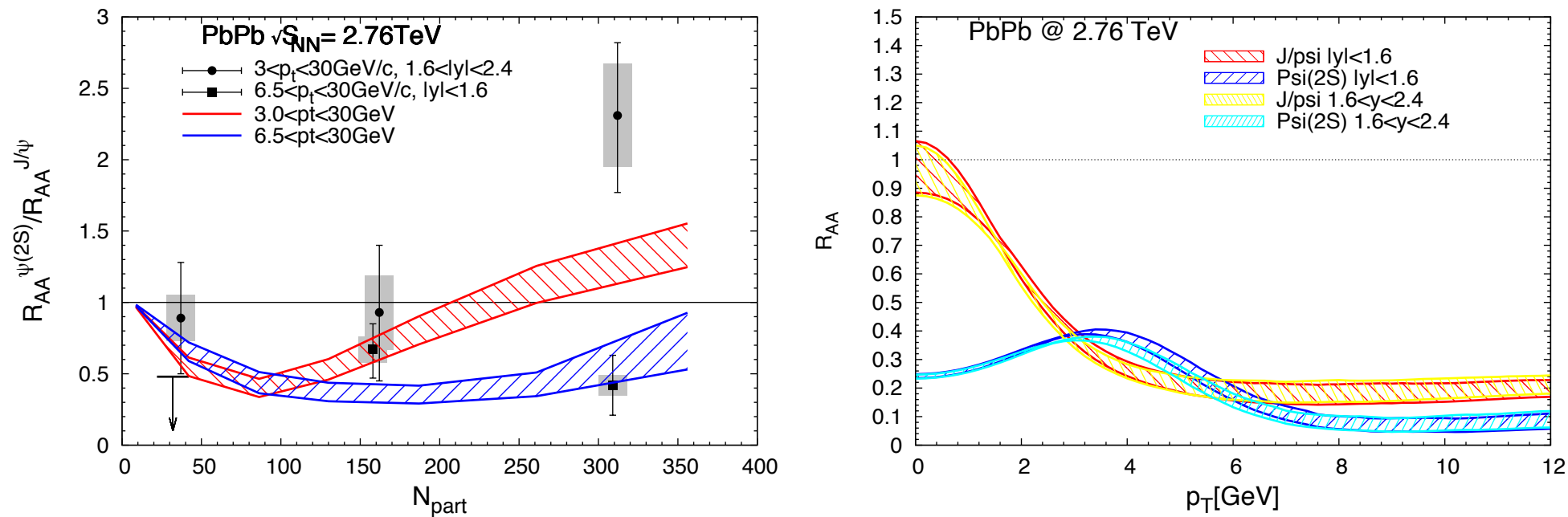


Figure 1. Charmonium production in Pb-Pb collisions at $\sqrt{s_{NN}}=2.76\text{ TeV}$ within the kinetic rate equation approach. Left panel: centrality dependence of the double ratio $R_{AA}(\psi(2S))/R_{AA}(J/\psi)$ for $p_T > 6.5\text{ GeV}$ and $|y| < 1.6$ (blue band) as well as $p_T > 3\text{ GeV}$ and $1.6 < |y| < 2.4$ (red band), compared to CMS data [7]. Right panel: p_T dependence of the individual J/ψ and $\psi(2S)$ R_{AA} 's for central collisions. A 10 % shadowing is assumed in the p_T -spectra according to EPS09 NLO [12].

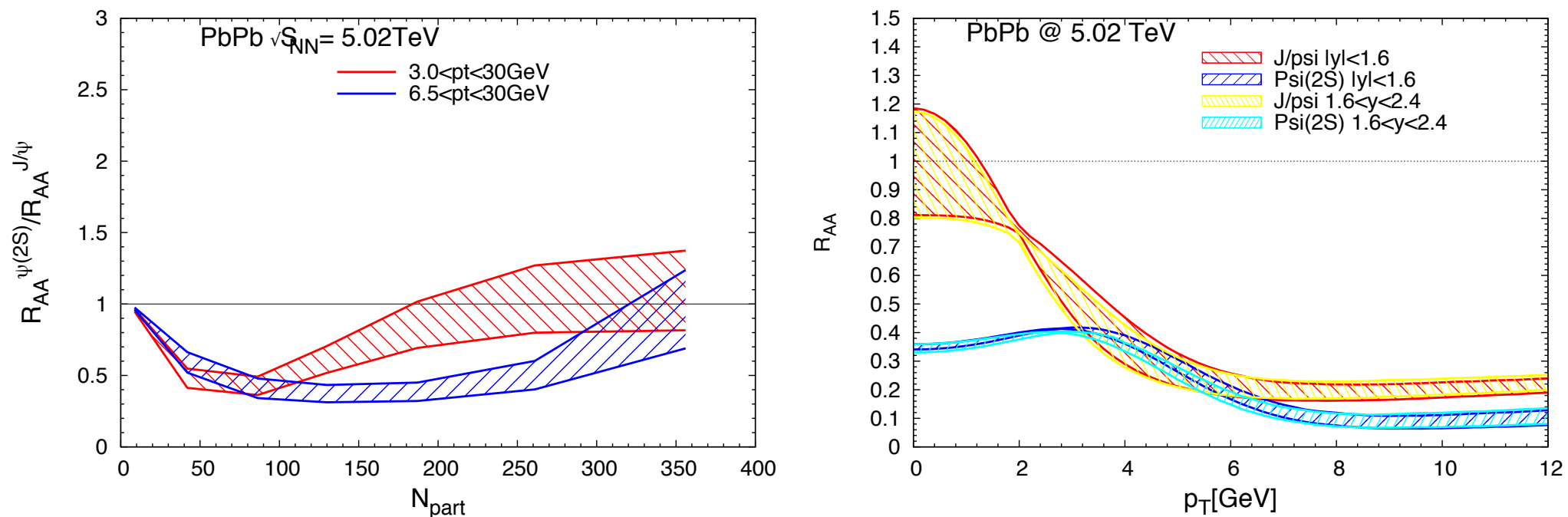
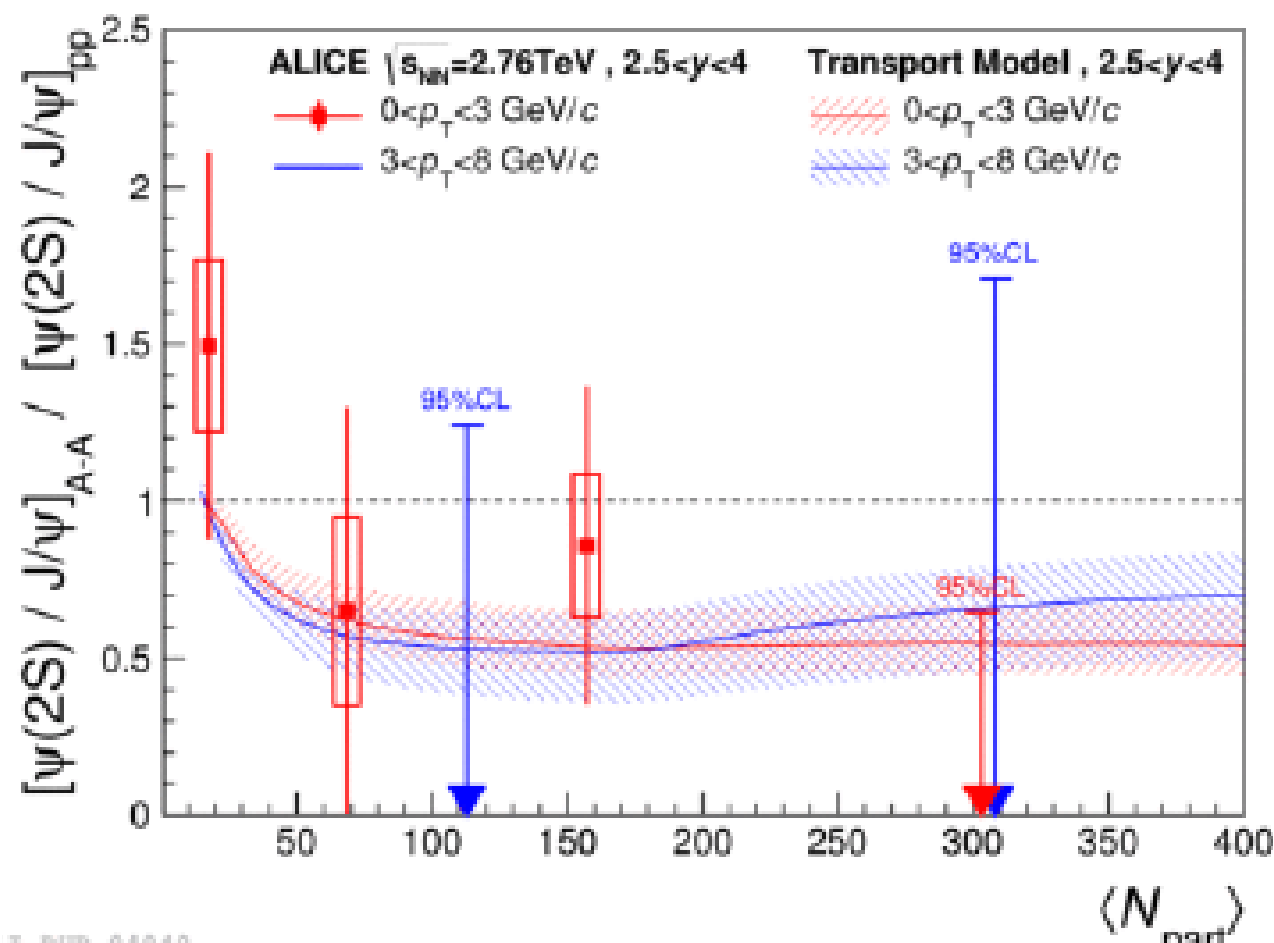
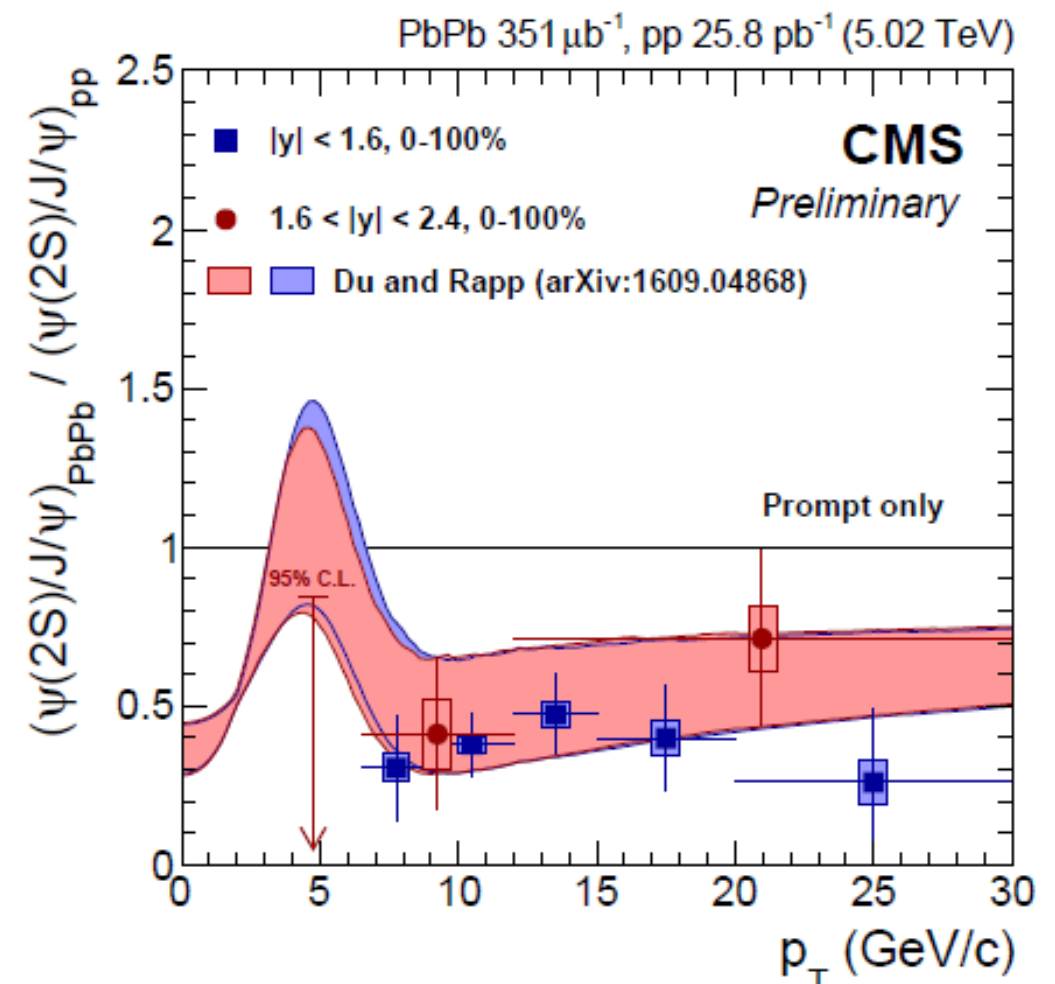
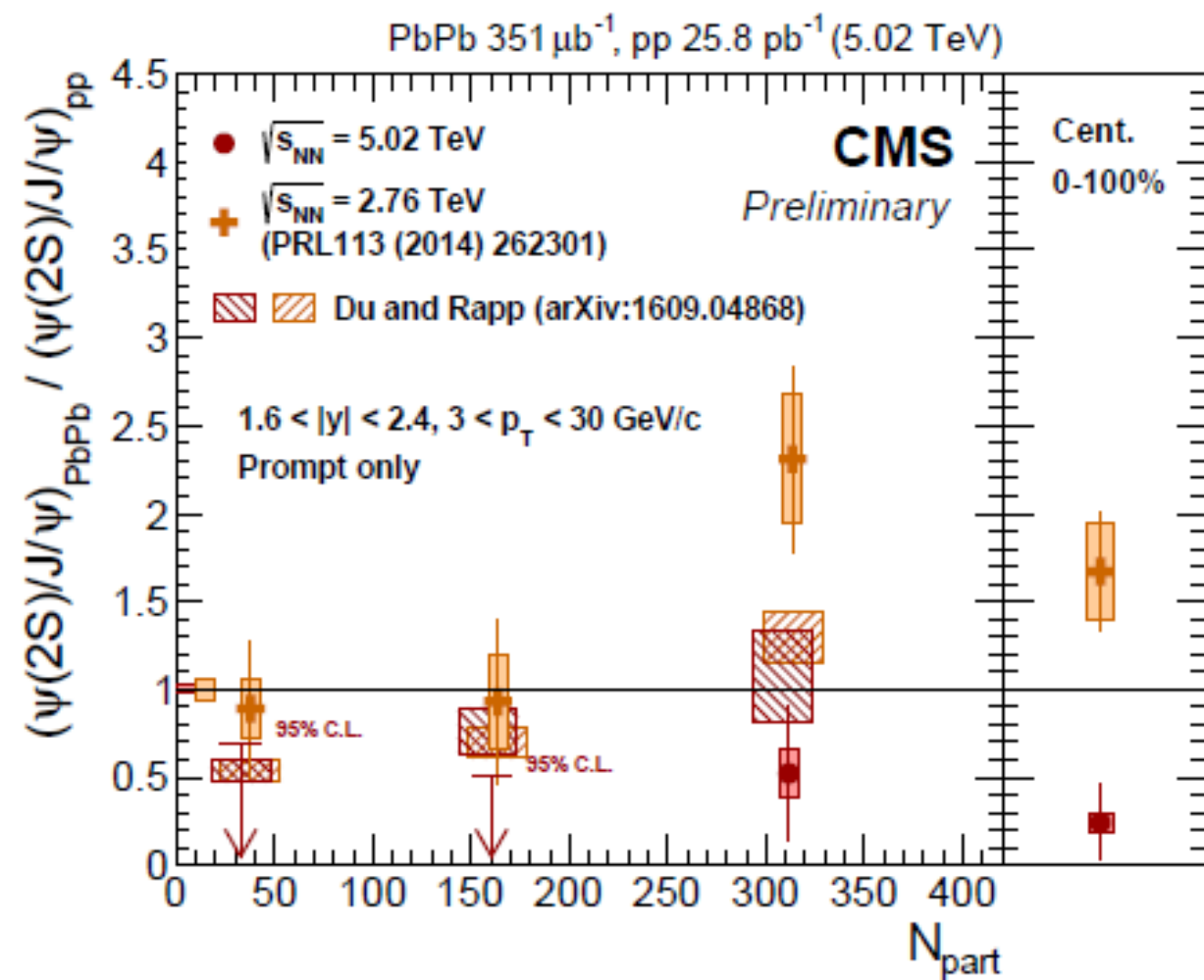


Figure 2. Same as Fig. 1 but for Pb-Pb collisions at $\sqrt{s_{NN}}=5.02\text{ TeV}$.



- $\psi(2S)$ regeneration occurring at higher p_T due to larger flow push
- Smart ad-hoc explanation for the enhancement at 2.76 TeV, still needed?
- Quality of ALICE results should improve in run-2 in order to give valuable input

increase in suppression for forward rapidities, which is due to the increased plateau halfwidth used in the initial conditions.

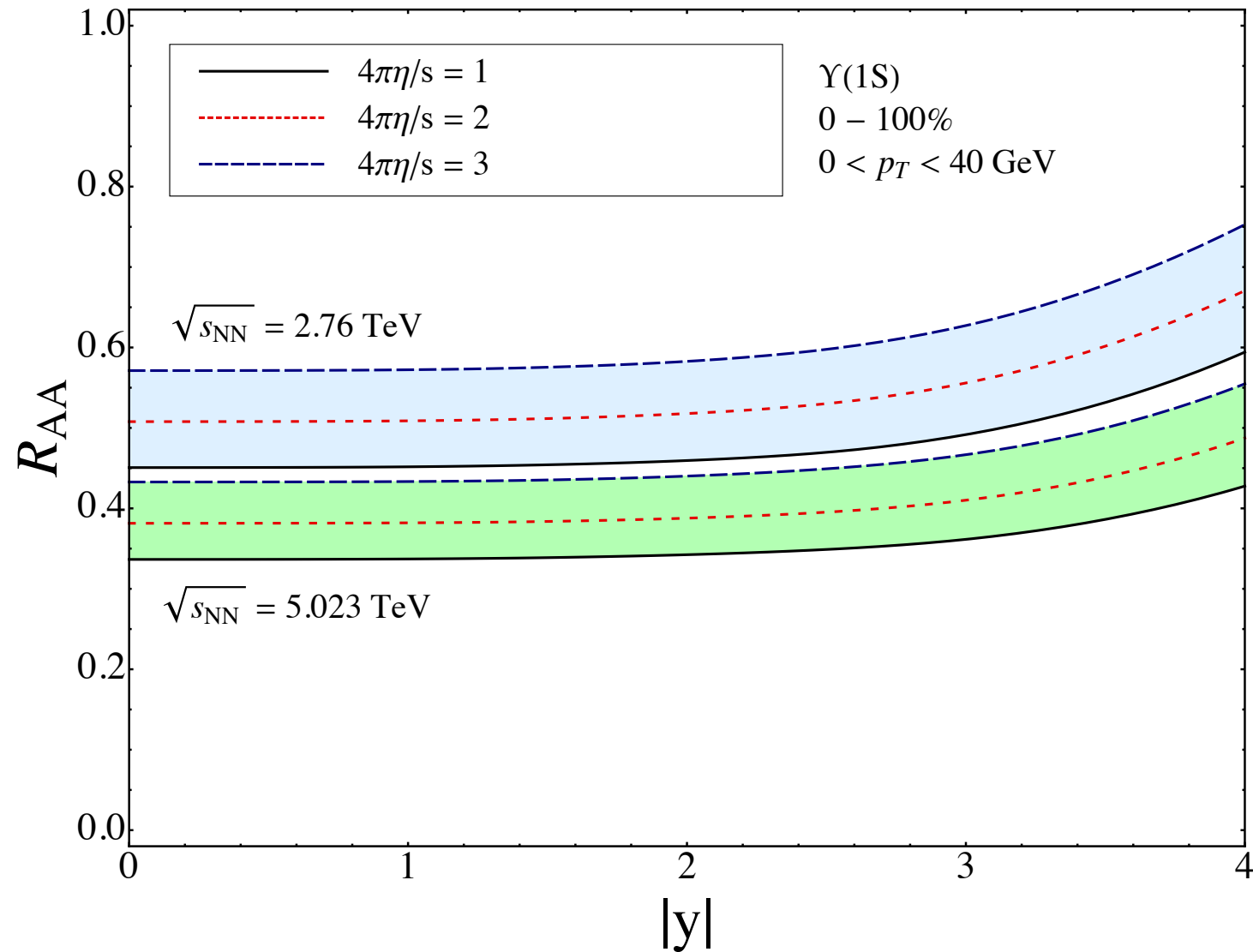


Figure 4. (Color online) Inclusive $Y(1S)$ state calculated with feed down contributions from excited states. Here we show a comparison between $\sqrt{s_{NN}} = 2.76$ TeV and $\sqrt{s_{NN}} = 5.023$ TeV collision energies.