



Pseudoscalar Portal to Dark Matter: Beyond Simplified Models

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1611.04593 (with D. Goncalves & P. Machado), 1509.01110

LHC DM WG MEETING 15/12/16



Simplified Models for DM Searches

wish list

- Simple enough as sensible unit within (more) complicated model
- Complete enough to accurately capture relevant LHC physics

(expect them to be *building blocks for UV complete DM theories*)

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Here focus on **Dirac Fermion DM + Spin-0 Mediator**

Scalar Mediator

$$\begin{aligned}\mathcal{L}_s &= \bar{\chi}(i\not{\partial} - m_\chi)\chi + \frac{1}{2}(\partial_\mu s)^2 - \frac{m_s^2}{2}s^2 \\ &- g_\chi s \bar{\chi}\chi - g_{\text{SM}} s \sum_q \frac{y_q}{\sqrt{2}} \bar{q}q\end{aligned}$$

Pseudoscalar Mediator

$$\begin{aligned}\mathcal{L}_a &= \bar{\chi}(i\not{\partial} - m_\chi)\chi + \frac{1}{2}(\partial_\mu a)^2 - \frac{m_a^2}{2}a^2 \\ &- ig_\chi a \bar{\chi}\gamma^5\chi - ig_{\text{SM}} a \sum_q \frac{y_q}{\sqrt{2}} \bar{q}\gamma^5q\end{aligned}$$

(Defined after EWSB)

For PSEUDOSCALAR Mediator, LHC (much) More Sensitive than DM Direct Detection

Strongly Suppressed
Spin Dependent @ Tree-level
Spin Independent @ One-loop

Simplified Models for DM Searches

Complete enough to accurately describe LHC physics ?



The Issue is $SU(2)_L \times U(1)_Y$ Gauge (non) Invariance

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ASSUME DM is SM GAUGE SINGLET

MEDIATOR NEEDS $SU(2)_L \times U(1)_Y$ CHARGE
to couple to SM fermions

Simplified Models for DM Searches

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S -channel pseudoscalar portal:

Mixing between a and Scalar EW Multiplet

$\sin \theta$

New States
(Mediator EW Partners)

New States Only Decouple by Closing Portal: $\sin \theta \sim \frac{\lambda v^2}{M^2 - m_a^2}$

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The Bottom Line:

- Imagine DM from EW Multiplet

DM Phenomenology neglecting other States in Multiplet likely to miss relevant physics

- Here, DM Simpl. Mod. face similar issue

→ Neglecting States in Mediator Multiplet Misses Relevant DM Physics

UV Complete Pseudoscalar Portal(s)

Consistently Complete the Pseudoscalar Portal DM Simplified Model

- $H + S$ (+ DM)
 $\hookrightarrow s + ia$

$$g_s(s + ia) |H|^2 + g_\chi \bar{\chi} (s + ia \gamma^5) \chi + h.c.$$

a can Mix with SM Higgs if CP Violation

\Rightarrow SI DM Direct Detection (scalar portal)

- 2HDM (+ DM)

Berlin, Gori, Lin, Wang, Phys. Rev **D92** (2015) 015005

$$V_{2\text{HDM}}(H_1, H_2) + g_\chi \bar{D}_{\chi_i} H_{1,2} \chi + h.c.$$

\hookrightarrow SU(2) doublet(s)

\Rightarrow Rich(er) DM Sector (+ DM feels SM Gauge Interactions)

$\Rightarrow a, H_0, H^\pm$ (New Scalars)

$$m_{H_0, H^\pm} - m_a \leq \mathcal{O}(\text{few}) \times v$$

- 2HDM + S (+ DM)

(a)

Simplest Completion with DM EW Singlet and Pseudoscalar Mediator

$\Rightarrow a, A, H_0, H^\pm$ (New Scalars)

Simplified Model
Misses Key DM Physics

Stefania Gori's Talk

2HDM + S Completion

(a)

Simple Renormalizable Pseudoscalar Portal to DM

Ipek, McKeen, Nelson, Phys. Rev **D90** (2014) 055021

Visible Sector

$$\begin{aligned}
 V_{2\text{HDM}} = & \mu_1^2 |H_1|^2 + \mu_2^2 |H_2|^2 - \mu^2 [H_1^\dagger H_2 + \text{h.c.}] \\
 & + \frac{\lambda_1}{2} |H_1|^4 + \frac{\lambda_2}{2} |H_2|^4 + \lambda_3 |H_1|^2 |H_2|^2 \\
 & + \lambda_4 |H_1^\dagger H_2|^2 + \frac{\lambda_5}{2} \left[(H_1^\dagger H_2)^2 + \text{h.c.} \right]
 \end{aligned}$$

Dark Sector

$$V_{\text{Dark}} = m_\chi \bar{\chi} \chi + \frac{1}{2} (\partial_\mu a_0)^2 + \frac{m_{a_0}^2}{2} a_0^2 + ig_\chi a_0 \bar{\chi} \gamma^5 \chi$$

Portal

$$V_{\text{portal}} = i \kappa a_0 H_1^\dagger H_2 + \text{h.c.}$$

Nomura, Thaler, Phys. Rev **D79** (2009) 075008

$$H_j = \begin{pmatrix} \phi_j^+ \\ \frac{v_j + h_j + i \eta_j}{\sqrt{2}} \end{pmatrix} \quad \begin{aligned} H^\pm &= -s_\beta \phi_1^\pm + c_\beta \phi_2^\pm & A_0 &= -s_\beta \eta_1 + c_\beta \eta_2 \\ h &= -s_\alpha h_1 + c_\alpha h_2 & H_0 &= -c_\alpha h_1 - s_\alpha h_2 \end{aligned}$$

↪ 125 GeV Higgs

$$-\mathcal{L}_{\text{Yuk}} = Y_{1,2}^u \bar{Q}_L q_R^u \tilde{H}_{1,2} + Y_{1,2}^d \bar{Q}_L q_R^d H_{1,2} + Y_{1,2}^\ell \bar{L}_L \ell_R H_{1,2} + \text{h.c.}$$

Assume Natural Flavour Conservation in \mathcal{L}_{Yuk}

To Keep in Mind: **2 Scalar Doublets** \Rightarrow **New Constraints**

EW Precision Observables

$$m_{H^\pm} \simeq m_{H_0} \quad \text{or} \quad m_{H^\pm} \simeq m_{A_0} \quad (\text{appr. custodial symmetry})$$

Higgs Signal Strengths

Flavour Physics

...

2HDM + S Completion

(a)

Seyda Ipek's Talk

Simple Renormalizable Pseudoscalar Portal to DM

Ipek, McKeen, Nelson, *Phys. Rev D* **90** (2014) 055021

Visible Sector

$$\begin{aligned}
 V_{2\text{HDM}} &= \mu_1^2 |H_1|^2 + \mu_2^2 |H_2|^2 - \mu^2 [H_1^\dagger H_2 + \text{h.c.}] \\
 &+ \frac{\lambda_1}{2} |H_1|^4 + \frac{\lambda_2}{2} |H_2|^4 + \lambda_3 |H_1|^2 |H_2|^2 \\
 &+ \lambda_4 |H_1^\dagger H_2|^2 + \frac{\lambda_5}{2} \left[(H_1^\dagger H_2)^2 + \text{h.c.} \right]
 \end{aligned}$$

Dark Sector

$$V_{\text{Dark}} = m_\chi \bar{\chi} \chi + \frac{1}{2} (\partial_\mu a_0)^2 + \frac{m_{a_0}^2}{2} a_0^2 + ig_\chi a_0 \bar{\chi} \gamma^5 \chi$$

Portal

$$V_{\text{portal}} = i \kappa a_0 H_1^\dagger H_2 + \text{h.c.}$$

Nomura, Thaler, *Phys. Rev D* **79** (2009) 075008

Mixing

Physical States

$$A = c_\theta A_0 + s_\theta a_0 \quad , \quad a = c_\theta a_0 - s_\theta A_0$$

$$H_j = \begin{pmatrix} \phi_j^+ \\ \frac{v_j + h_j + i \eta_j}{\sqrt{2}} \end{pmatrix} \quad \begin{aligned} H^\pm &= -s_\beta \phi_1^\pm + c_\beta \phi_2^\pm \\ h &= -s_\alpha h_1 + c_\alpha h_2 \end{aligned} \quad \begin{aligned} A_0 &= -s_\beta \eta_1 + c_\beta \eta_2 \\ H_0 &= -c_\alpha h_1 - s_\alpha h_2 \end{aligned}$$

$$V_{\text{Dark}} \supset ig_\chi (c_\theta a + s_\theta A) \bar{\chi} \gamma^5 \chi$$

$$\begin{aligned}
 V_{\text{portal}} &= \frac{(m_A^2 - m_a^2) s_{2\theta}}{2v} (c_{\beta-\alpha} H_0 - s_{\beta-\alpha} h) \\
 &\times [aA (s_\theta^2 - c_\theta^2) + (a^2 - A^2) s_\theta c_\theta]
 \end{aligned}$$

2HDM + S Completion

(a)

NEW STATES A H^\pm H_0
 multi-mediator scenario \leftarrow
 scenario

Mass Splittings among A H^\pm H_0 bounded by 2HDM unitarity

Ginzburg, Ivanov, *Phys. Rev D* **72** (2005) 115010

$$m_i - m_j \leq \mathcal{O}(\text{few}) \times v$$

1) $m_{A,H_0,H^\pm} \sim m_a$

2 mediators scenario (affects e.g. Relic Density)

H^\pm H_0 phenomenology as in 2HDM

2) $m_{A,H_0,H^\pm} \gg m_a$

NEW STATES DO NOT (completely) DECOUPLE
 (if $\sin \theta$ is kept fixed)

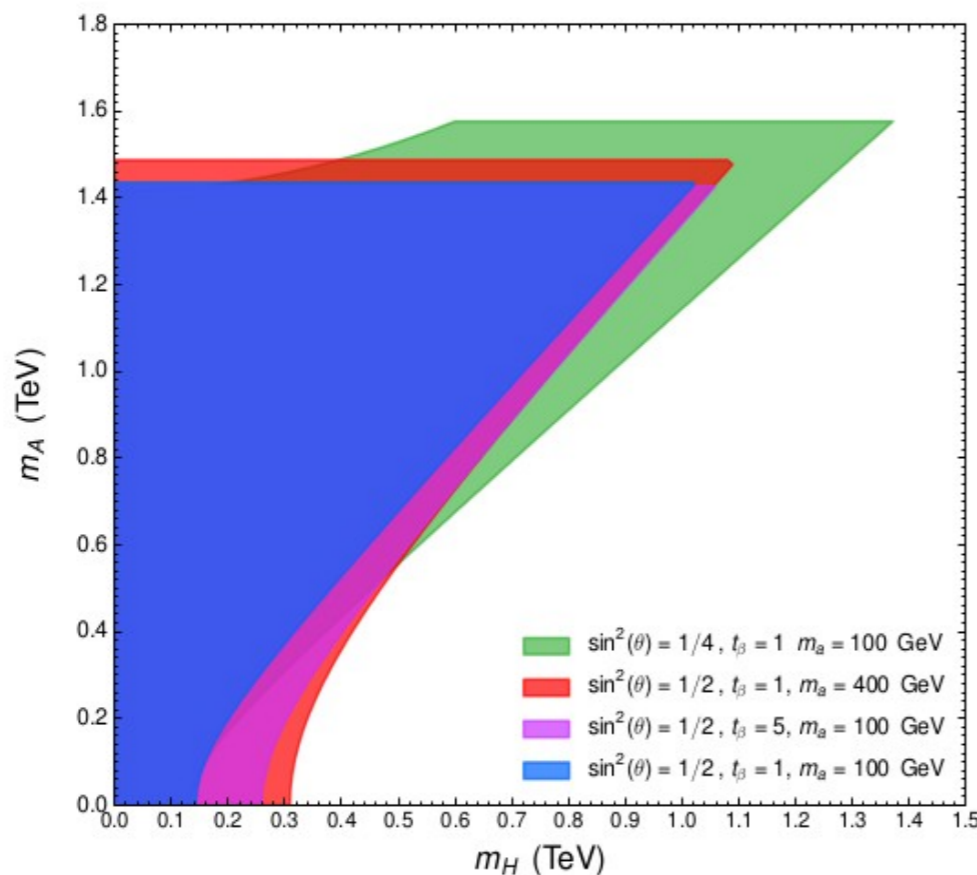
Assume $[c_{\beta-\alpha} = 0 \quad m_{H^\pm} = m_{H_0}]$

$M^2 = \mu^2 / (s_\beta c_\beta)$

$$\Delta_a^2 \equiv m_A^2 - m_a^2 \quad \Delta_H^2 = M^2 - m_{H^\pm}^2 + 2m_W^2 - m_h^2/2$$

$$\Lambda_\pm = \left[\frac{\Delta_H^2}{v^2} - \frac{\Delta_a^2 (1 - c_{4\theta})}{8v^2} \pm \sqrt{\frac{\Delta_H^4}{v^4} + \frac{\Delta_a^4 (1 - c_{4\theta})}{8v^4}} \right]$$

Pert. Unitarity $\rightarrow |\Lambda_\pm| \leq 8\pi$

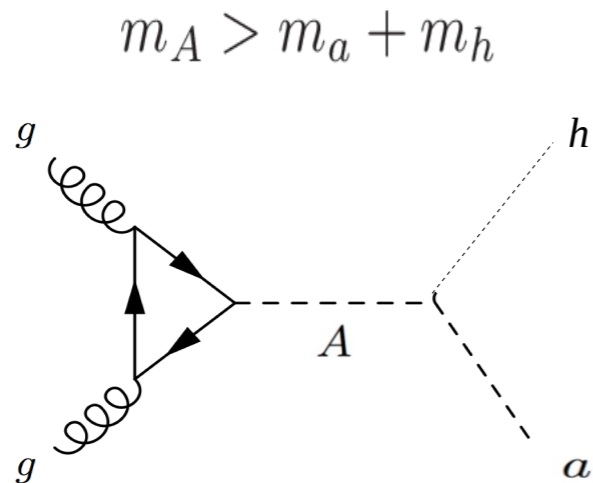


2HDM + S Completion

(a)

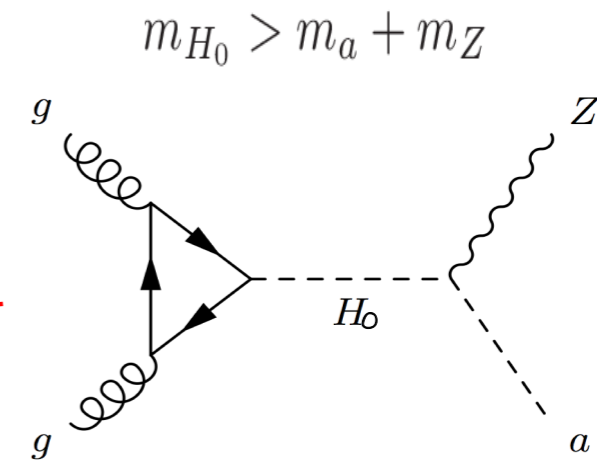
LHC SIGNATURES $m_{A,H_0,H^\pm} \gg m_a$

JMN, Phys. Rev **D93** (2016) 031701



(Resonant) mono-h/mono-Z
(mono-W)

(Assume a Decays invisibly)



$$\cancel{E}_T^{\max} \sim \frac{1}{2m_A} \sqrt{(m_A^2 - m_h^2 - m_a^2)^2 - 4m_h^2 m_a^2}$$

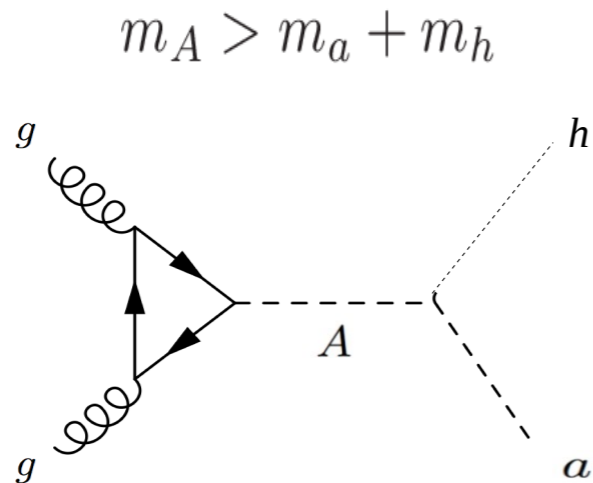
$$\cancel{E}_T^{\max} \sim \frac{1}{2m_{H_0}} \sqrt{(m_{H_0}^2 - m_Z^2 - m_a^2)^2 - 4m_Z^2 m_a^2}$$

2HDM + S Completion

(a)

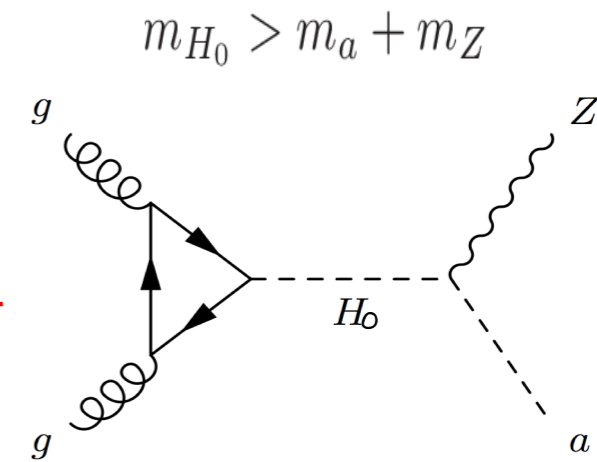
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JMN, Phys. Rev **D93** (2016) 031701



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$$\cancel{E}_T^{\max} \sim \frac{1}{2m_{H_0}} \sqrt{(m_{H_0}^2 - m_Z^2 - m_a^2)^2 - 4m_Z^2 m_a^2}$$

Assumed Gluon Fusion production of H_0 A

for 2HDM Type II & $t_\beta \gg 1 \Rightarrow \bar{b}b$ - associated

mono-h needs multi-mediator **2HDM + S**

~~2HDM~~

$\bar{b}b + \cancel{E}_T + Z, h$

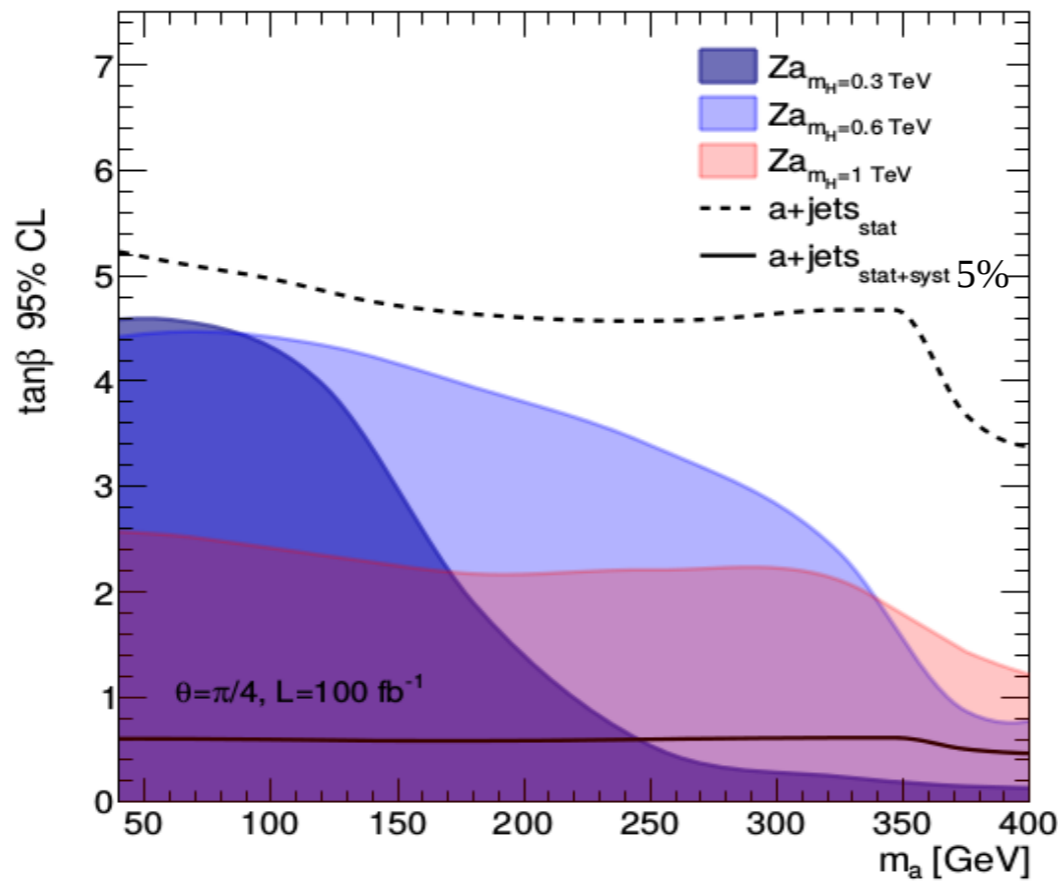
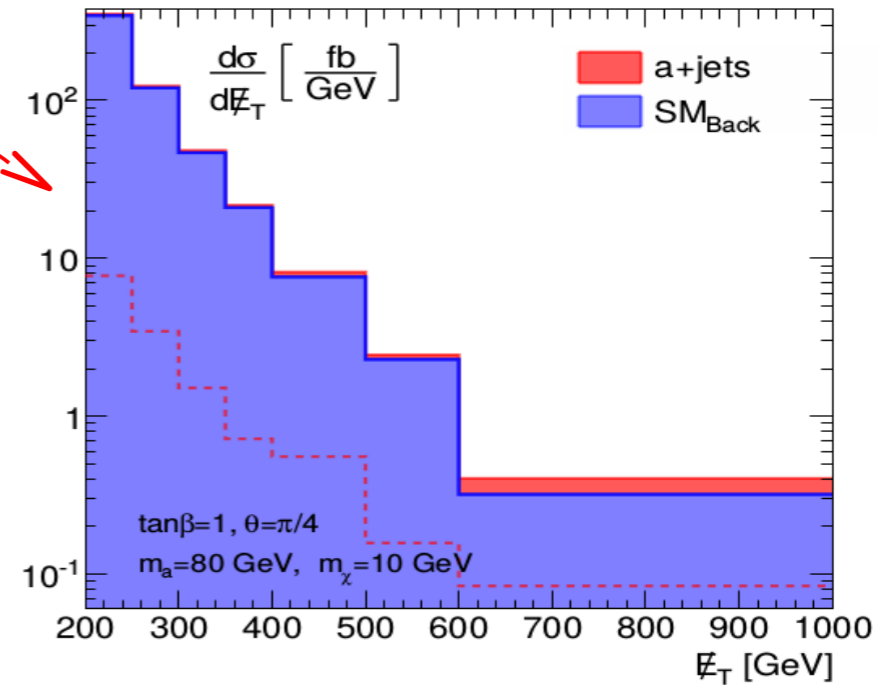
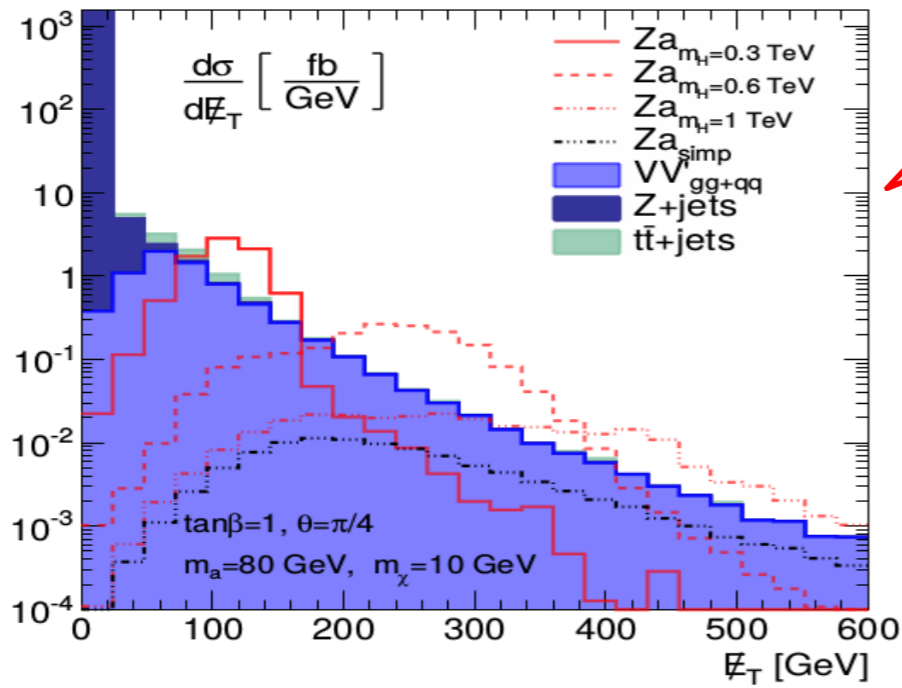
Fairbairn, JMN, Tunney, work in progress

mono-Z is 2HDM generic **2HDM + S**
2HDM

2HDM + S Completion

(a)

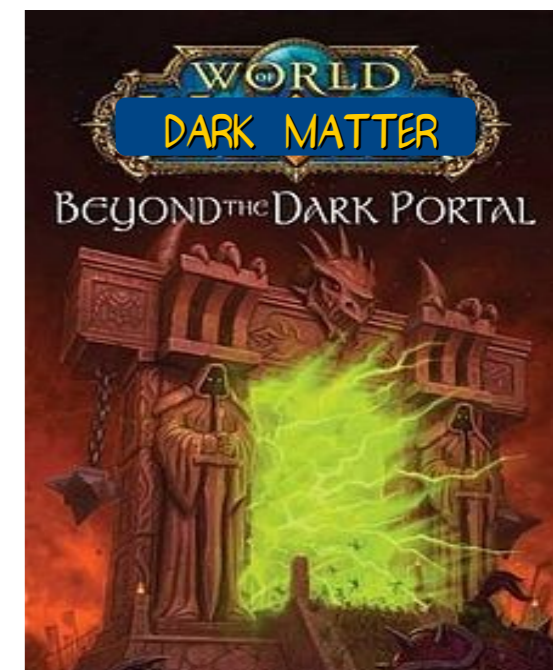
Compare $Z + \cancel{E}_T$ to Jets + \cancel{E}_T @ LHC



Sensitivity Comparison

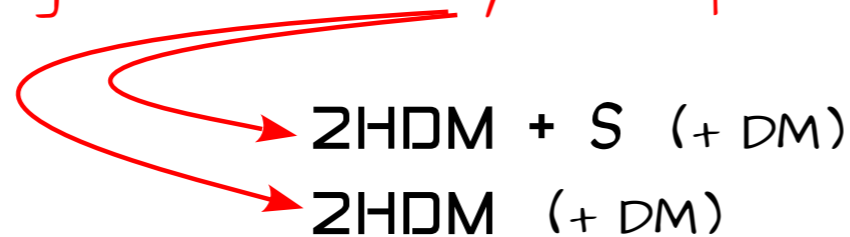
Goncalves, Machado, JMN, ArXiv:1611.04593

Summary & Thoughts



- Pseudoscalar Portal to DM:

Strong Case to Go Beyond Simplified Model



- Both 2HDM (+ DM), 2HDM + S (+ DM) Capture Scalar/Pseudoscalar Portal
- Resonant *mono-Z* Signature Generic in 2HDM (+ DM), 2HDM + S (+ DM)
Scalar/Pseudoscalar Portal for $m_\phi \gg m_{\text{med}}$
- Consistent Completions Can Point to Reasonable Ranges of Couplings
within Simplified Model $\rightarrow g_{\text{SM}}$