Effective Field Theories at one loop

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Based on:

arXiv:1704.04504 (to appear in EPJC)

In collaboration with: Alejandro Celis, Avelino Vicente, Javier Virto

JHEP 1609 (2016) 156 [arXiv:1607.02142]

In collaboration with: Jorge Portolés and Pedro Ruíz-Femenía

Motivation: Effective Field Theories

• EFTs offer a **model independent approach**: Physics above a given scale mapped into Wilson coefficients of higher dimensional operators

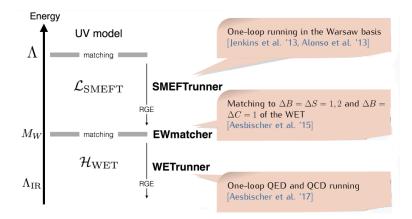
$$\mathcal{L}_{\text{SMEFT}} = \mathcal{L}_{\text{SM}} + \frac{1}{\Lambda} \sum_{k} C_{k}^{(5)} Q_{k}^{(5)} + \frac{1}{\Lambda^{2}} \sum_{k} C_{k}^{(6)} Q_{k}^{(6)} + \mathcal{O}\left(\frac{1}{\Lambda^{3}}\right)$$

- LHC + EWPD data can be used to do global fits of the Wilson coefficients of the SMEFT
 - \longrightarrow Any UV model should satisfy the resulting bounds
- EFTs at one-loop
 - Some contributions to the EFT are only generated at one-loop. Moreover, with the increasing precision in some observables, one-loop corrections can be important
 - Ongoing efforts to extend the EFT analyses to NLO See for instance CERN Yellow Report 4
 - Loop corrections to the EFT also important in Flavor Physics
 e.g. Feruglio, Paradisi, Pattori, 1705.00929; Pruna, Signer, 1408.3565

DsixTools: The SMEFT toolkit

Celis, JF, Vicente, Virto, 1704.04504

A Mathematica package for the handling of the d = 6 SMEFT

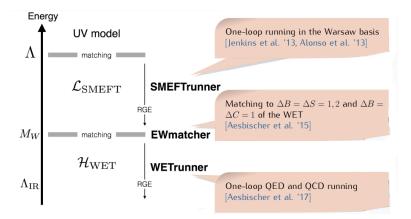


- Modular structure ⇒ Each module can be used independently
- DsixTools is expected to be extended with more modules

DsixTools: The SMEFT toolkit

Celis, JF, Vicente, Virto, 1704.04504

A Mathematica package for the handling of the d = 6 SMEFT



 A (simple) systematic framework to match any UV theory to its EFT is highly desirable!

Conventional matching vs functional approach

Two general approaches for the construction of the EFT:

• Matching the diagrammatic computation of Green Functions with light particles in the external legs in the UV theory and in the EFT

$$\phi_L$$
 ϕ_H + ϕ_H +

• Integrate out the heavy fields, and extract the local contributions relevant for the Wilson coefficients

$$e^{iS} = \mathcal{N} \int \mathcal{D}\phi_L \mathcal{D}\phi_H \exp\left[i \int dx \,\mathcal{L}(\phi_L, \phi_H)\right]$$

Functional integration techniques are **more powerful** than the matching procedure when one aims to determine the **full EFT**

One-loop effective action

Expand the UV Lagrangian around the classical fields

$$\mathcal{L} = \mathcal{L}^{\text{tree}}(\hat{\eta}) + \left(\eta^{\dagger} \frac{\delta \mathcal{L}}{\delta \eta^{*}} + \frac{\delta \mathcal{L}}{\delta \eta} \eta\right)_{\eta = \hat{\eta}} + \frac{1}{2} \eta^{\dagger} \frac{\delta \mathcal{L}}{\delta \eta^{*} \delta \eta} \eta + \mathcal{O}(\eta^{3})$$
$$= 0 \quad (\text{EOM}) \qquad \equiv O \quad \stackrel{\text{fluctuation}}{\text{operator}}$$

From the fluctuation operator we get the one-loop effective action

$$e^{iS^{1loop}} = \mathcal{N} \int \mathcal{D}\eta \exp\left[i \int dx \frac{1}{2} \eta^{\dagger} O \eta\right] \qquad \eta \equiv \begin{pmatrix} \eta_H \\ \eta_L \end{pmatrix}$$

is form of the fluctuation operator:

Generic form of the fluctuation operator:

$$O = \begin{pmatrix} \Delta_H & X_{LH}^{\dagger} \\ X_{LH} & \Delta_L \end{pmatrix} \qquad \qquad \Delta_H, \Delta_L: \text{ heavy and light loops} \\ X_{HL}: \text{ heavy-light loops} \end{cases}$$

Our aim: compute the one-loop heavy particle effects and extract their contributions to the Wilson coefficients of the EFT

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Two important simplifications JF, Portolés, Femenía, 1607.02142

The effect of the heavy-light loops can be shifted to the heavy part

$$\eta \to P \eta \qquad O = \begin{pmatrix} \Delta_H & X_{LH}^{\dagger} \\ X_{LH} & \Delta_L \end{pmatrix} \longrightarrow P^{\dagger} O P = \begin{pmatrix} \widetilde{\Delta}_H & 0 \\ 0 & \Delta_L \end{pmatrix}$$

- Loops with heavy fields receive contributions from the soft $(p \sim m_L)$ and the hard $(p \sim m_H)$ momentum regions but only the hard region contributes to the EFT
 - The method of expansion by regions allows to separate each region by Taylor expanding the loop integrand over the parameters that are small there, and integrating over the full domain. Beneke, Smirnov '98

$$S_{H} = \underbrace{S_{H}^{\text{hard}}}_{p \sim m_{H} \gg m_{L}} + \underbrace{S_{H}^{\text{soft}}}_{p \sim m_{L} \ll m_{H}}$$
Pure short distance! $p \sim m_{L} \ll m_{H}$
Heavy propagator expanded in p/m_{H}
(same as loops with tree EFT vertices)

expanded in p/m_H

Long story short

We developed a functional method for the one-loop matching that clarifies the treatment of heavy-light loops and provides important simplifications

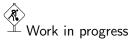
$$\int d^d x \, \mathcal{L}_{\rm EFT}^{\rm 1loop} = S_H^{\rm hard}$$

- A shift in the field allows to easily isolate the heavy-field contributions
- **Compute only the hard part**, the rest cancel in the matching (also true for diagrammatic matching!)

$$S_H^{\text{soft}} + S_L - S_{\text{EFT}}^{\text{tree}} = 0$$

- $S_H^{\rm hard}$ can be easily computed using functional techniques. Mostly algebraic problem... but a large amount of algebra is involved
- Automation becomes essential for any realistic model

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A Mathematica package for the matching and RGE evolution from the new physics scale to the scale of low energy observables

Manual: arXiv:1704.04504 Website: https://dsixtools.github.io/

New features are coming soon Comments (including critical ones!), questions and suggestions are welcome!

Thanks!