

LITTLE CONFORMAL SYMMETRY

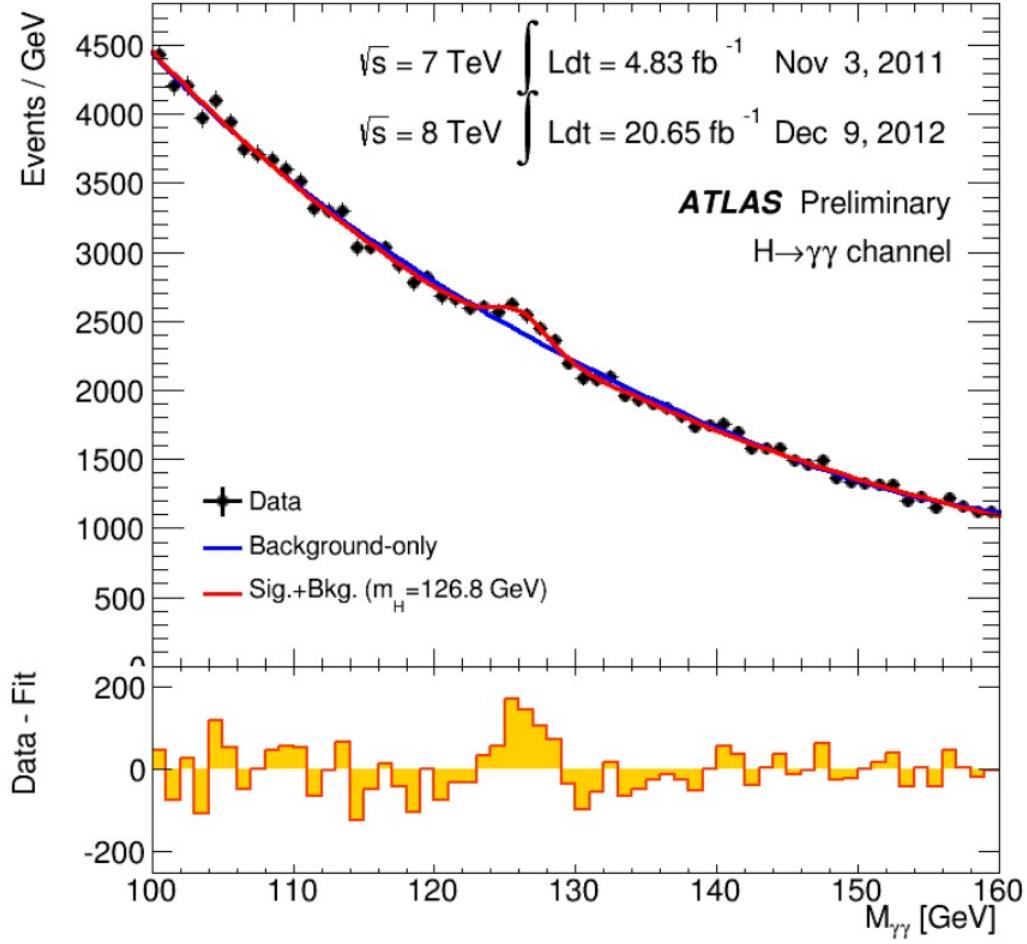
arXiv: 1603.00030



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Invisibles 2017 Workshop
Zurich, Switzerland

In Collaboration with John Terning (UC Davis),
Kitran Colwell (UC Davis)

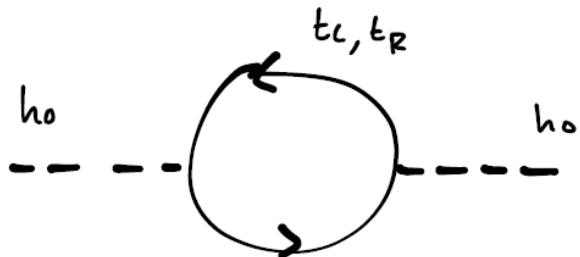
125 GeV Higgs



Amazing!



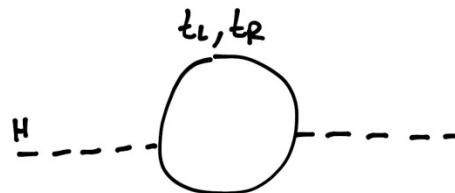
except...



$$\delta m_h^2 = -\frac{N_C}{16\pi^2} |y_t|^2 \times \left[2\Lambda^2 - 6m_t^2 \ln\left(\frac{\Lambda^2 + m_t^2}{m_t^2}\right) + \dots \right]$$

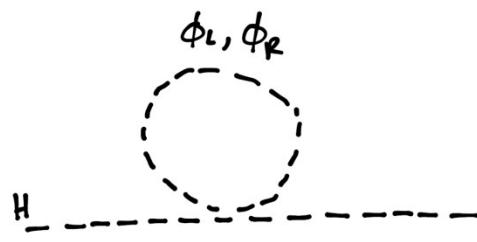
Cancelling the Divergence

SUSY's Claim to Fame



Feynman diagram showing a loop of top quarks (t_L, t_R) connected to a Higgs boson (H). The loop is represented by a solid circle with dashed lines connecting it to the external Higgs and top quark lines.

$$\delta m_h^2 = -\frac{N_c}{16\pi^2} |y_t|^2 \times \left[2\Lambda^2 - 6m_t^2 \ln\left(\frac{\Lambda^2 + m_t^2}{m_t^2}\right) + \dots \right]$$



Feynman diagram showing a loop of scalar fields (ϕ_L, ϕ_R) connected to a Higgs boson (H). The loop is represented by a dashed circle with dashed lines connecting it to the external Higgs and scalar field lines.

$$\delta m_h^2 = \frac{\lambda N}{16\pi^2} \left[2\Lambda^2 - m_L^2 \ln\left(\frac{\Lambda^2 + m_L^2}{m_L^2}\right) - m_R^2 \ln\left(\frac{\Lambda^2 + m_R^2}{m_R^2}\right) + \dots \right]$$

SUSY guarantees such a cancellation

Where are the superpartners?

ATLAS SUSY Searches* - 95% CL Lower Limits

Status: August 2016

ATLAS Preliminary

$\sqrt{s} = 7, 8, 13 \text{ TeV}$

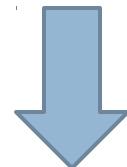
| Model | e, μ, τ, γ | Jets | E_T^{miss} | $\int \mathcal{L} dt [\text{fb}^{-1}]$ | Mass limit | ATLAS Preliminary | | |
|--|--|-------------------------------|---------------------|--|-----------------|--------------------------------|-----------------------------|--|
| | | | | | | $\sqrt{s} = 7, 8 \text{ TeV}$ | $\sqrt{s} = 13 \text{ TeV}$ | |
| Inclusive Searches | MSUGRA/CMSSM | 0-3 $e, \mu/1-2 \tau$ | 2-10 jets/3 b | Yes | 20.3 | \tilde{g}, \tilde{g} | 1.85 TeV | 1507.05525 |
| | $\tilde{g}, \tilde{g} \rightarrow q\bar{q}\chi_1^0$ | 0 | 2-6 jets | Yes | 13.3 | \tilde{g} | 1.35 TeV | ATLAS-CONF-2016-078 |
| | $\tilde{g}, \tilde{g} \rightarrow q\bar{q}\chi_1^0$ [compressed] | mono-jet | 1-3 jets | Yes | 3.2 | \tilde{g} | 608 GeV | 1604.07773 |
| | $\tilde{g}, \tilde{g} \rightarrow q\bar{q}\chi_1^0 + q\bar{q}W^{\pm}\chi_1^0$ | 0 | 2-6 jets | Yes | 13.3 | \tilde{g} | 1.80 TeV | ATLAS-CONF-2016-078 |
| | $\tilde{g}, \tilde{g} \rightarrow q\bar{q}\chi_1^0 + q\bar{q}(\ell/\nu)\chi_1^0$ | 3 e, μ | 4 jets | - | 13.2 | \tilde{g} | 1.83 TeV | ATLAS-CONF-2016-078 |
| | $\tilde{g}, \tilde{g} \rightarrow q\bar{q}W\chi_1^0$ | 2 e, μ (SS) | 0-3 jets | Yes | 13.2 | \tilde{g} | 1.7 TeV | ATLAS-CONF-2016-037 |
| | GMSB ($\tilde{\tau}$ NLSP) | 1-2 $\tau + 0-1 \ell$ | 0-2 jets | Yes | 3.2 | \tilde{g} | 1.6 TeV | ATLAS-CONF-2016-037 |
| | GGM (bino NLSP) | 2 γ | - | Yes | 3.2 | \tilde{g} | 2.0 TeV | 1607.05979 |
| | GGM (higgsino-bino NLSP) | γ | 1 b | Yes | 20.3 | \tilde{g} | 1.37 TeV | 1608.09150 |
| | GGM (higgsino-bino NLSP) | γ | 2 jets | Yes | 13.3 | \tilde{g} | 1.8 TeV | 1507.05493 |
| 1st gen. squarks 1st gen. sleptons | GGM (higgsino NLSP) | 2 e, μ (Z) | 2 jets | Yes | 20.3 | \tilde{g} | 900 GeV | ATLAS-CONF-2016-095 |
| | Gravitino LSP | 0 | mono-jet | Yes | 20.3 | \tilde{g} | 865 GeV | 1503.03290 |
| | | | | | $E^{1/2}$ scale | | | 1502.01518 |
| 1st gen. med. 1st gen. sleptons | $\tilde{g}, \tilde{g} \rightarrow b\bar{b}\chi_1^0$ | 0 | 3 b | Yes | 14.8 | \tilde{g} | 1.89 TeV | ATLAS-CONF-2016-052 |
| | $\tilde{g}, \tilde{g} \rightarrow t\bar{t}\chi_1^0$ | 0-1 e, μ | 3 b | Yes | 14.8 | \tilde{g} | 1.89 TeV | ATLAS-CONF-2016-052 |
| | $\tilde{g}, \tilde{g} \rightarrow b\bar{b}\chi_1^0$ | 0-1 e, μ | 3 b | Yes | 20.1 | \tilde{g} | 1.37 TeV | 1407.0800 |
| 2nd gen. squarks direct production | $\tilde{b}_1 \tilde{b}_1, \tilde{b}_1 \rightarrow b\bar{b}\chi_1^0$ | 0 | 2 b | Yes | 3.2 | \tilde{b}_1 | 840 GeV | 1608.08772 |
| | $\tilde{b}_1 \tilde{b}_1, \tilde{b}_1 \rightarrow b\bar{b}\chi_1^0$ | 2 e, μ (SS) | 1 b | Yes | 13.2 | \tilde{b}_1 | 325-689 GeV | ATLAS-CONF-2016-037 |
| | $\tilde{b}_1 \tilde{b}_1, \tilde{b}_1 \rightarrow b\bar{b}\chi_1^0$ | 0-2 e, μ | 1-2 b | Yes | 4.7/13.3 | \tilde{b}_1 | 200-720 GeV | 1209.2102, ATLAS-CONF-2016-077 |
| | $\tilde{t}_1 \tilde{t}_1, \tilde{t}_1 \rightarrow W\tilde{b}\chi_1^0$ or \tilde{t}_1^0 | 0-2 e, μ | 0-2 jets/1-2 b | Yes | 4.7/13.3 | \tilde{t}_1 | 90-198 GeV | 1506.08618, ATLAS-CONF-2016-077 |
| | $\tilde{t}_1 \tilde{t}_1, \tilde{t}_1 \rightarrow \ell\bar{\ell}\chi_1^0$ | 0 | mono-jet | Yes | 3.2 | \tilde{t}_1 | 90-323 GeV | 1604.07773 |
| | $\tilde{t}_1 \tilde{t}_1$ (natural GMSB) | 2 e, μ (Z) | 1 b | Yes | 20.3 | \tilde{t}_1 | 150-600 GeV | 1403.5222 |
| | $\tilde{t}_2 \tilde{t}_2, \tilde{t}_2 \rightarrow \tilde{t}_2 + Z$ | 3 e, μ (Z) | 1 b | Yes | 13.3 | \tilde{t}_2 | 290-700 GeV | ATLAS-CONF-2016-038 |
| | $\tilde{t}_2 \tilde{t}_2, \tilde{t}_2 \rightarrow \tilde{t}_2 + k$ | 3 e, μ | 6 jets + 2 b | Yes | 20.3 | \tilde{t}_2 | 320-620 GeV | 1506.08816 |
| | $\tilde{e}_1 \tilde{e}_1, \tilde{e}_1 \rightarrow \ell\bar{\ell}\chi_1^0$ | 2 e, μ | 0 | Yes | 20.3 | \tilde{e}_1 | 90-335 GeV | 1403.5204 |
| | $\tilde{e}_1 \tilde{e}_1, \tilde{e}_1 \rightarrow \ell\bar{\ell}\ell\bar{\ell}\nu\nu$ | 2 e, μ | 0 | Yes | 13.3 | \tilde{e}_1 | 640 GeV | ATLAS-CONF-2016-095 |
| EW direct | $\tilde{e}_1 \tilde{e}_1, \tilde{e}_1 \rightarrow \tau\bar{\tau}\chi_1^0$ | 2 τ | - | Yes | 14.8 | \tilde{e}_1 | 580 GeV | ATLAS-CONF-2016-095 |
| | $\tilde{e}_1^0 \tilde{e}_1^0 \rightarrow \tilde{e}_1 \tilde{e}_1, \nu\bar{\nu}, \tilde{\nu}\bar{\nu}$ | 3 e, μ | 0 | Yes | 13.3 | $\tilde{e}_1^0, \tilde{e}_1^0$ | 1.0 TeV | ATLAS-CONF-2016-095 |
| | $\tilde{e}_1^0 \tilde{e}_1^0 \rightarrow W\tilde{b}\chi_1^0$, $\tilde{e}_1^0 \tilde{e}_1^0 \rightarrow h\tilde{b}\ell/W\ell/\tau\tau/yy$ | 2-3 e, μ | 0-2 jets | Yes | 20.3 | $\tilde{e}_1^0, \tilde{e}_1^0$ | 425 GeV | 1403.5204, 1402.7029 |
| | $\tilde{e}_1^0 \tilde{e}_1^0 \rightarrow W\tilde{b}\chi_1^0$, $\tilde{e}_1^0 \tilde{e}_1^0 \rightarrow h\tilde{b}\ell/W\ell/\tau\tau/yy$ | 4 e, μ, γ | 0-2 b | Yes | 20.3 | $\tilde{e}_1^0, \tilde{e}_1^0$ | 270 GeV | 1501.07110 |
| | $\tilde{e}_1^0 \tilde{e}_1^0 \rightarrow \ell\bar{\ell}\chi_1^0$ | 4 e, μ | 0 | Yes | 20.3 | $\tilde{e}_1^0, \tilde{e}_1^0$ | 635 GeV | 1405.5088 |
| | GGM (wino NLSP) weak prod. | 1 $e, \mu + \gamma$ | - | Yes | 20.3 | \tilde{W} | 115-370 GeV | 1507.05493 |
| | GGM (bino NLSP) weak prod. | 2 γ | - | Yes | 20.3 | \tilde{W} | 590 GeV | 1507.05493 |
| | Direct $\tilde{\chi}_1^0 \tilde{\chi}_1^0$ prod., long-lived $\tilde{\chi}_1^0$ | Disapp. trk | 1 jet | Yes | 20.3 | $\tilde{\chi}_1^0$ | 270 GeV | 1310.03875 |
| | Direct $\tilde{\chi}_1^0 \tilde{\chi}_1^0$ prod., long-lived $\tilde{\chi}_1^0$ | dE/dx trk | - | Yes | 18.4 | $\tilde{\chi}_1^0$ | 495 GeV | 1508.05332 |
| | Stable, stopped \tilde{g} Hadron | 0 | 1-5 jets | Yes | 27.9 | \tilde{g} | 850 GeV | 1310.65884 |
| Long-lived particles | Stable \tilde{g} R-hadron | trk | - | - | 3.2 | \tilde{g} | 1.58 TeV | 1608.05129 |
| | Metastable \tilde{g} R-hadron | dE/dx trk | - | - | 3.2 | \tilde{g} | 1.57 TeV | 1604.045820 |
| | GMSB, stable $\tilde{t}_1, \tilde{b}_1 \rightarrow \tilde{t}_1 \tilde{b}_1$, $\tilde{t}_1 \tilde{b}_1 \rightarrow \ell\bar{\ell}\chi_1^0$ | 1-2 μ | - | - | 19.1 | \tilde{t}_1, \tilde{b}_1 | 537 GeV | 1411.6705 |
| | GMSB, $\tilde{t}_1 \rightarrow \tilde{t}_1 \tilde{b}_1$, long-lived \tilde{t}_1^0 | 2 γ | - | Yes | 20.3 | \tilde{t}_1^0 | 440 GeV | 1409.5542 |
| | gg, $\tilde{t}_1 \rightarrow ee/ee/\mu\mu/\mu\mu$ | displ. $\nu\nu/\mu\mu/\mu\nu$ | - | - | 20.3 | \tilde{t}_1 | 1.0 TeV | 1504.05162 |
| | gg, $\tilde{t}_1 \rightarrow ee/ee/\mu\mu/\mu\mu$ | displ. vts + jets | - | - | 20.3 | \tilde{t}_1 | 1.0 TeV | 1504.05162 |
| | LNV $p_T \rightarrow \tilde{\nu}_T + X, \tilde{\nu}_T \rightarrow e\bar{\nu}_T/\mu\bar{\nu}_T/\tau\bar{\nu}_T$ | - | - | - | 3.2 | $\tilde{\nu}_T$ | 1.9 TeV | 1607.08779 |
| | Bi-linear RPV CMSSM | 2 e, μ (SS) | 0-3 b | Yes | 20.3 | \tilde{g}, \tilde{g} | 1.45 TeV | 1404.2500 |
| | $\tilde{e}_1 \tilde{e}_1, \tilde{e}_1 \rightarrow W\tilde{b}\chi_1^0 \rightarrow ee\nu_e, ee\nu_e, \mu\mu\nu_e$ | 4 e, μ | - | Yes | 13.3 | \tilde{e}_1 | 1.14 TeV | ATLAS-CONF-2016-075 |
| | $\tilde{e}_1 \tilde{e}_1, \tilde{e}_1 \rightarrow W\tilde{b}\chi_1^0 \rightarrow \tau\tau\nu_\tau, \tau\tau\nu_\tau$ | 3 $e, \mu + \tau$ | - | Yes | 20.3 | \tilde{e}_1 | 450 GeV | 1405.5088 |
| RPV | $\tilde{g}, \tilde{g} \rightarrow q\bar{q}q\bar{q}$ | 0 | 4-5 large- R jets | - | 14.8 | \tilde{g} | 1.08 TeV | ATLAS-CONF-2016-057 |
| | $\tilde{g}, \tilde{g} \rightarrow q\bar{q}\chi_1^0, \tilde{g}, \tilde{g} \rightarrow q\bar{q}q\bar{q}$ | 0 | 4-5 large- R jets | - | 14.8 | \tilde{g} | 1.55 TeV | ATLAS-CONF-2016-057 |
| | $\tilde{g}, \tilde{g} \rightarrow t\bar{t}, \tilde{b}_1 \rightarrow b\bar{b}x$ | 1 e, μ | 8-10 jets/0-4 b | - | 14.8 | \tilde{g} | 1.73 TeV | ATLAS-CONF-2016-094 |
| | $\tilde{g}, \tilde{g} \rightarrow t\bar{t}, \tilde{b}_1 \rightarrow b\bar{b}x$ | 1 e, μ | 8-10 jets/0-4 b | - | 14.8 | \tilde{g} | 1.4 TeV | ATLAS-CONF-2016-094 |
| | $\tilde{t}_1 \tilde{b}_1, \tilde{t}_1 \rightarrow b\bar{b}x$ | 0 | 2 jets + 2 b | - | 15.4 | \tilde{t}_1 | 410 GeV | ATLAS-CONF-2016-022, ATLAS-CONF-2016-084 |
| | $\tilde{t}_1 \tilde{b}_1, \tilde{t}_1 \rightarrow b\bar{b}x$ | 2 e, μ | 2 b | - | 20.3 | \tilde{t}_1 | 450-510 GeV | ATLAS-CONF-2016-015 |
| | Scalar charm, $\tilde{c} \rightarrow \tilde{c}\tilde{c}^0$ | 0 | 2 c | Yes | 20.3 | \tilde{c} | 0.4-1.0 TeV | 1501.01325 |
| <p>*Only a selection of the available mass limits on new states or phenomena is shown.</p> | | | | | | | | |



Alternative Theory: Introduce a new Gauge Boson

$$\delta m_h^2 \ni \text{---} \circ \text{---} + \text{---} z' \text{---}$$

$$\delta m_h^2 \ni (-2N_C y_t^2 + 3C_2(S) g_N^2) \frac{\Lambda^2}{16\pi^2}$$



$$0 = -2N_C y_t^2 + 3C_2(S) g_N^2$$

A cancellation is possible if y_t
and g_N satisfy this relationship.



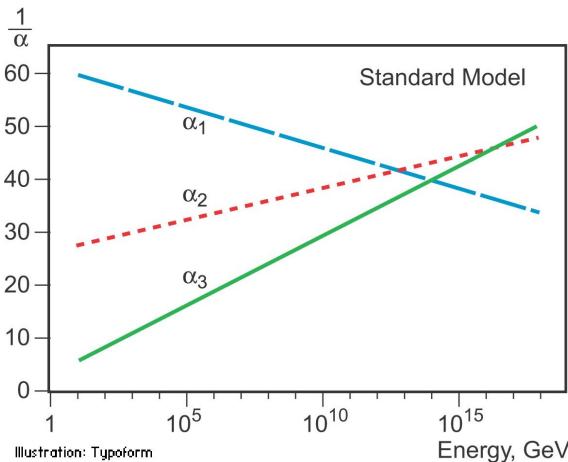
Some Problems....

- How is this not just different fine tuning?

$$0 = -2N_C y_t^2 + 3C_2(S) g_N^2$$



- y_t and g_N run, so a cancellation at one scale will be spoiled



[www.nobelprize.org/
nobel_prizes/physics/
laureates/2004/](http://www.nobelprize.org/nobel_prizes/physics/laureates/2004/)

Some Problems....

- Why would the cutoffs for both loops be the same?

$$\delta m_h^2 \ni (-2N_C y_t^2 + 3C_2(S) g_N^2) \frac{\Lambda^2}{16\pi^2}$$

↓

$$\delta m_h^2 \ni -2N_C y_t^2 \Lambda_t^2 + 3C_2(S) g_N^2 \Lambda_N^2$$



Depends on high energy theory

Solution: Little Conformal Symmetry

What if this relation is the result
of an IR symmetry?

$$0 = -2N_C y_t^2 + 3C_2(S) g_N^2$$

M. J. G. Veltman, "The Infrared-Ultraviolet Connection," Acta Phys. Polon. B **12** (1981) 437.

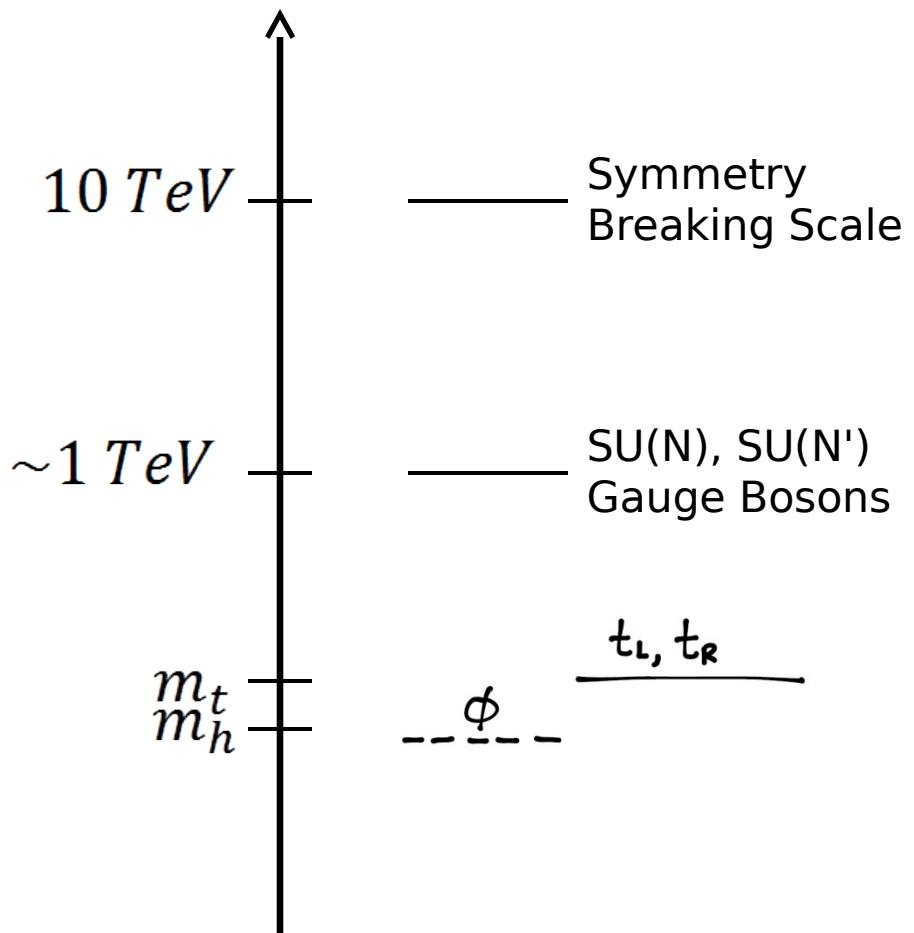
- Impose Conformal Symmetry to derive this relationship between couplings
- Allows for a naturally small Higgs mass with superpartners > 10 TeV, similar to Little Higgs models

N. Arkani-Hamed, A. G. Cohen, T. Gregoire, and J. G. Wacker,
Phenomenology of electroweak symmetry breaking from theory space,
[*hep-ph/0202089*](https://arxiv.org/abs/hep-ph/0202089)

Little Conformal Symmetry: A Simple Toy Model

$$\mathcal{L} \ni -y_t \bar{t}_R \phi t_L$$

| | copies | $SU(N)$ |
|--------|----------------|---------|
| t_L | $2 \times N_c$ | 1 |
| t_R | N_c | □ |
| ϕ | 2 | □ |



Conclusions

- Fixed points can produce a cancellation of the Higgs mass quadratic divergence with a new $SU(N)$ gauge boson
- This cancellation prevents the Higgs mass from being sensitive to new physics up to the 10 TeV scale
- More complicated and realistic models can be explored
 - Add another gauge group: Embedding in QCD
 - Neutral Naturalness: Competing with QCD

Thank You!

Cancellation from an IR Fixed Point

- Quadratic Divergence Cancelled

$$\alpha_t = \frac{3C_2(\phi)}{2N_C} \alpha_N$$

- Top Yukawa at fixed point

M E. Machacek and M. T. Vaughn,
Nucl Phys. B 236. 221 (1983)

$$\alpha_t = \frac{B_N}{A} \alpha_N$$

M E. Machacek and M. T.
Vaughn, Nucl Phys. B 222. 83
(1983)

$$\beta(\alpha_N) \ni \left(\text{loop diagram} + m \text{ loop} + m \text{ (dashed loop)} \right) + \left(\text{loop diagram} + \text{loop diagram} + \dots \right)$$

- β –function coefficients depend on the matter content of the UV theory

T. Banks and A. Zaks, “On the Phase Structure of Vector-Like Gauge Theories with Massless Fermion,” *Nucl. Phys. B* 196 (1982) 189.

Little Conformal Symmetry: More Realistic Toy Model

Problem:

$$\beta(\alpha_t) = \frac{1}{2\pi} (A \alpha_t^2 - B_N \alpha_N \alpha_t + B_3 \alpha_3 \alpha_t + \delta m_{\eta_2}^2 \bar{\alpha}_t - B_1 \alpha_1 \alpha_t)$$

We still want these to cancel

QCD

$SU(2)_L$

$U(1)_Y$

Too big! Drags $\beta(\alpha_t)$ negative quickly

Idea: Embed $SU(3) \subset SU(N)$

More Specifically ...

$$SU(3) \subset SU(3)_L \times SU(3)_R \subset SU(N) \times SU(N')$$

- We have two new gauge couplings instead of one: α_N and $\alpha_{N'}$,
- Add another constraint to ensure both gauge groups are at a fixed point

More Realistic Toy Model

| | $SU(N)$ | $SU(N')$ |
|-------|---------|----------|
| t_R | □ | 1 |
| H | □ | □ |
| q_L | 1 | □ |

Constraints

$$\delta m_h^2 = 0$$

$$\dot{\alpha}_t = 0$$

$$\dot{\alpha}_N = 0$$

$$\dot{\alpha}_{N'} = 0$$

new constraint & new variable



Theory with Realistic Symmetry Breaking



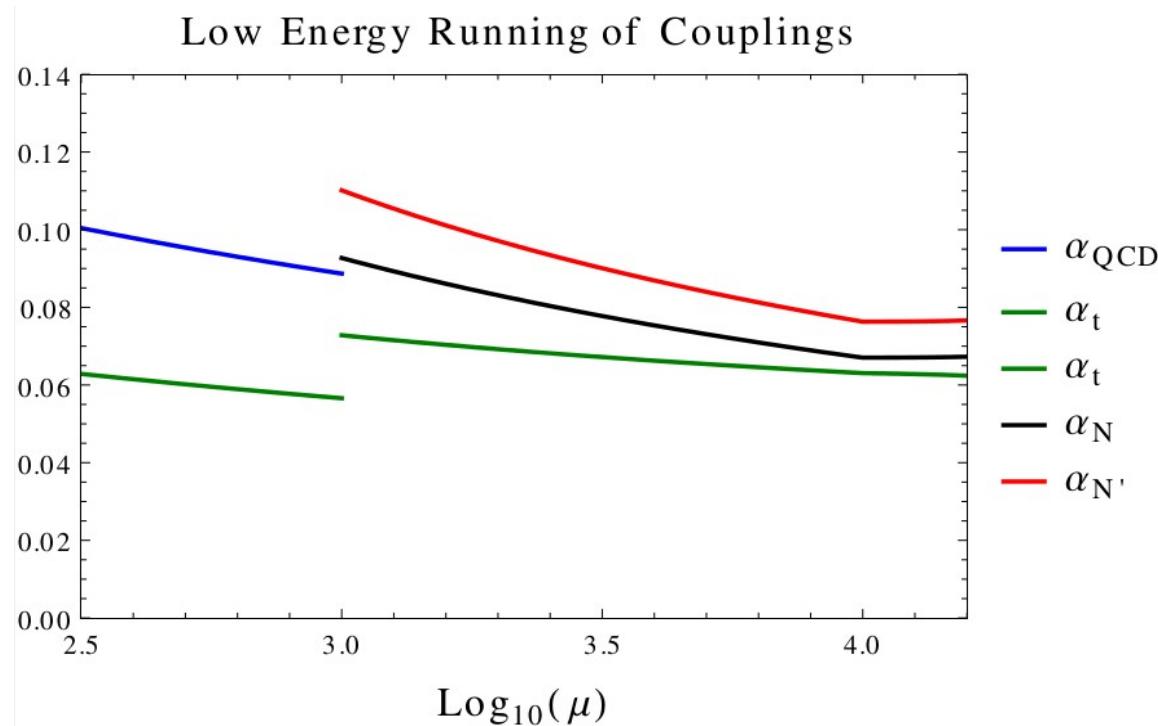
Realistic symmetry breaking scenario:

$$\frac{1}{\alpha_3} \geq \frac{1}{\alpha_N} + \frac{1}{\alpha_{N'}}$$



Realistic top quark Yukawa coupling:

$$\alpha_t^{\bar{MS}}(m_t) \text{ within } \pm 5.65 \text{ %}$$



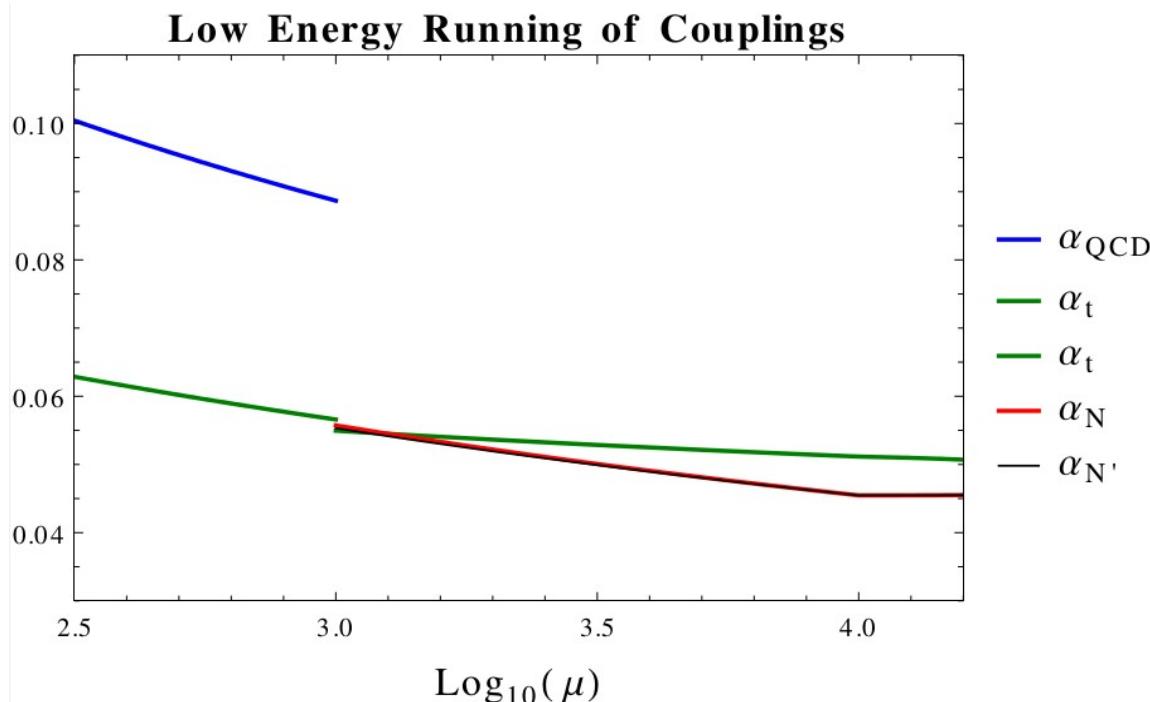
Theory with Realistic Yukawa Coupling

✗ Realistic symmetry breaking scenario:

$$\frac{1}{\alpha_3} \geq \frac{1}{\alpha_N} + \frac{1}{\alpha_{N'}}$$

✓ Realistic top quark Yukawa coupling:

$$\alpha_t^{\bar{MS}}(m_t) \text{ within } \pm 5.65 \text{ %}$$



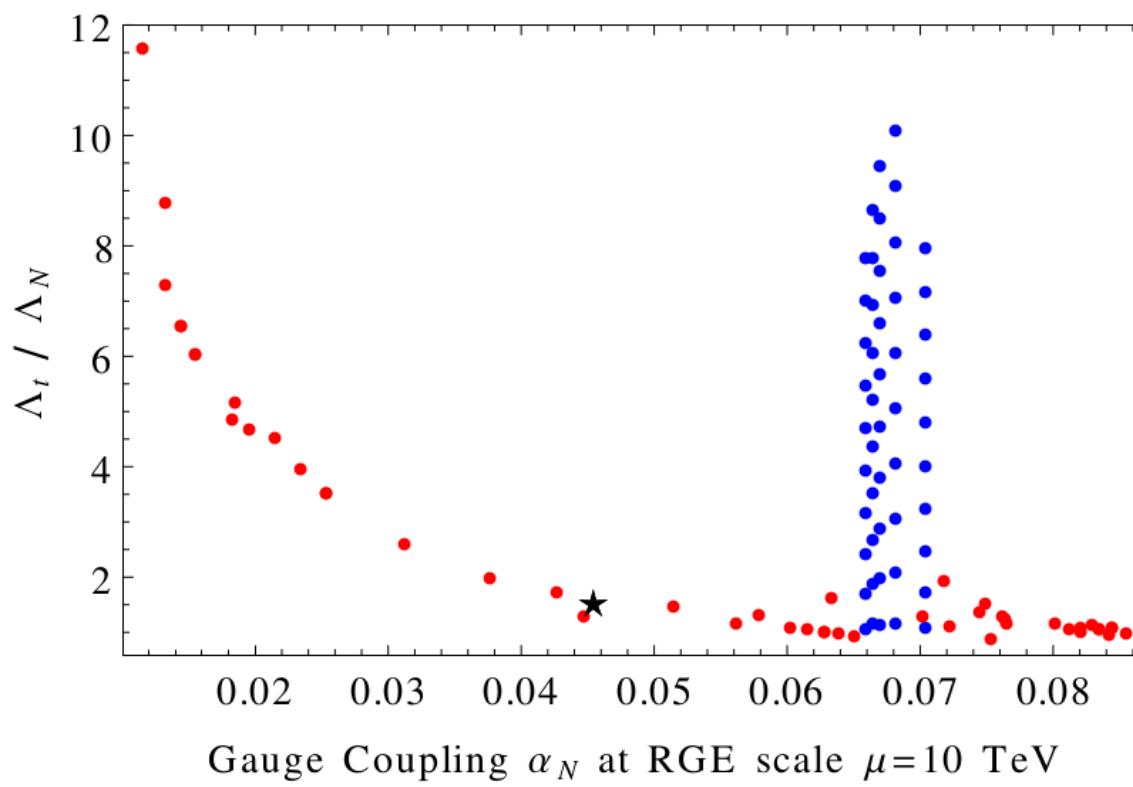
Little Conformal Symmetry UV Theories

Realistic symmetry breaking scenario:

$$\frac{1}{\alpha_3} \geq \frac{1}{\alpha_N} + \frac{1}{\alpha_{N'}}$$

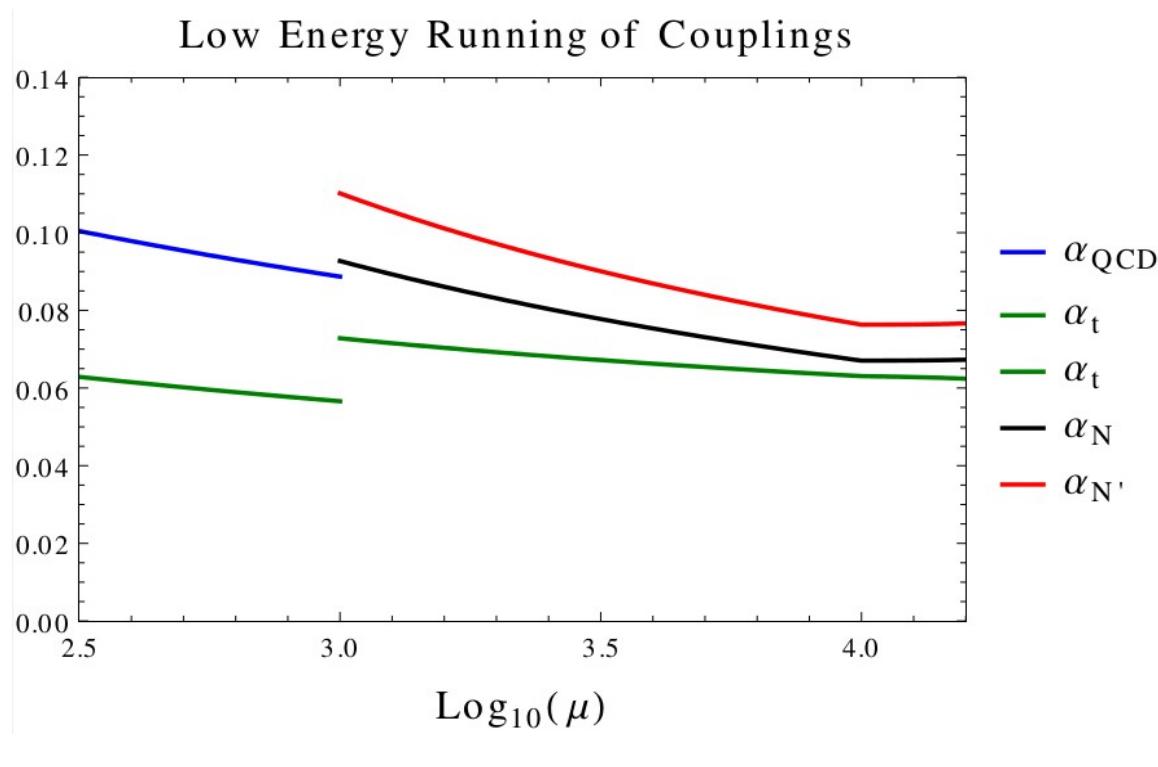
Realistic top quark Yukawa coupling:

$$\alpha_t^{\bar{MS}}(m_t) \text{ within } \pm 5.65 \text{ %}$$



Top Yukawa and QCD

- The top Yukawa coupling is still being affected by QCD

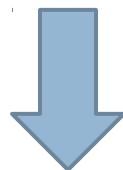


- Small improvements require more complicated matter content
- We need more control over $\beta(\alpha_t)$

Neutral Naturalness

$$\delta m_h^2 \ni \text{---} \circ \text{---} + \text{---} \circ \overset{\Phi}{\curvearrowleft}$$

$$\delta m_h^2 \ni (-2 N_C \alpha_t + \lambda) \frac{\Lambda^2}{16 \pi^2}$$



$$0 = -2 N_C \alpha_t \Lambda_t^2 + \lambda \Lambda_\Phi^2$$