Theoretical perspectives on $R(D)$ and $R(D^*)$

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New physics scale and flavor

- SM cannot be the full story — past theoretical prejudices haven't been confirmed
- Are measures of fine tuning misleading, and NP is order of magnitude heavier?
- New physics at a TeV — MFV probably useful approximation to its flavor structure

  $\Updownarrow$

  New physics at $10^{1-2}\,\mathrm{TeV}$ — less strong flavor suppression, MFV less motivated
- Strong SM suppressions (GIM, CKM, loops, chiral) $\Rightarrow$ sensitive to very high scales

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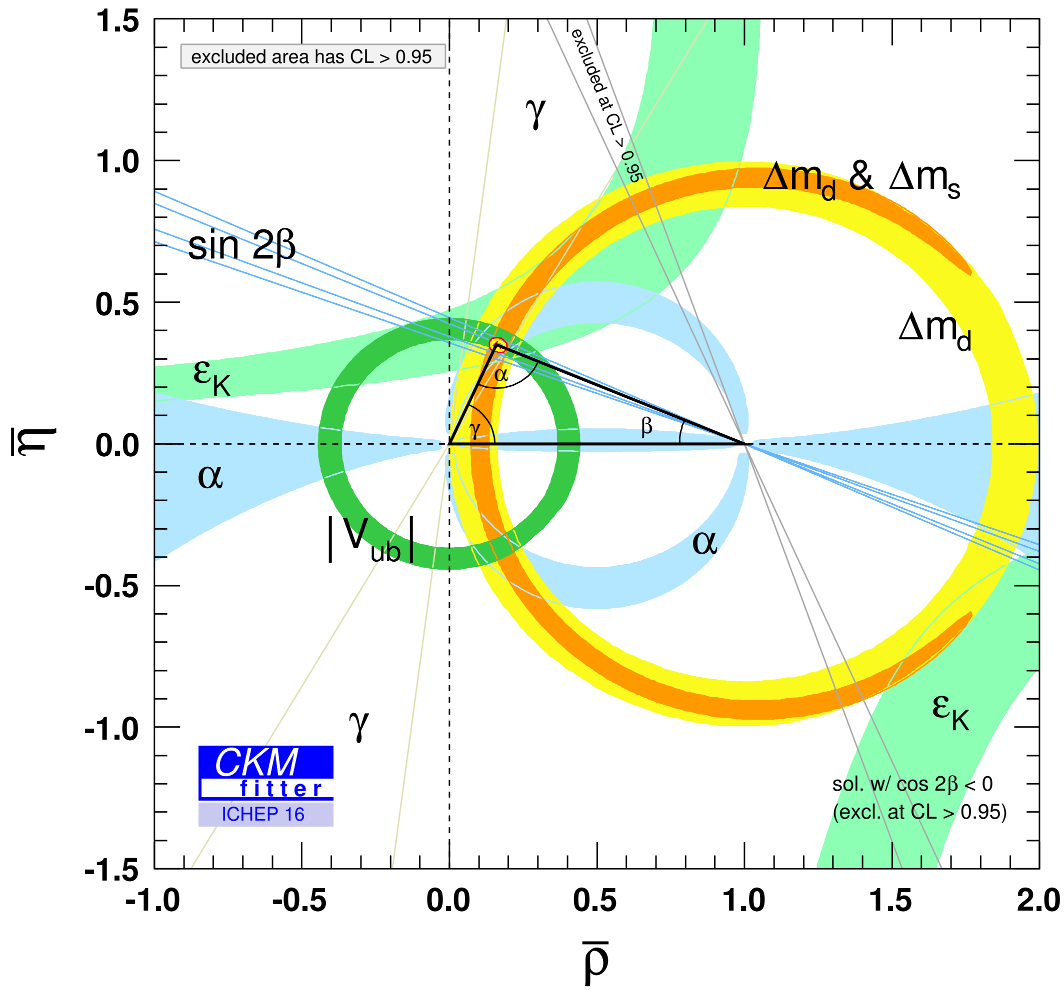
- Future: $\dfrac{\text{(Belle II data set)}}{\text{(Belle data set)}} \sim \dfrac{\text{(LHCb lifetime)}}{\text{(LHCb now)}} \sim \dfrac{\text{(ATLAS \& CMS 3/ab)}}{\text{(ATLAS \& CMS now)}} \sim 50-100$
- Conservatively: increases in mass scales probed $\sqrt[4]{50} \sim 2.7$

  (for dim-6 contributions to B decays, H couplings, etc.)

New questions for $100\times$ more data? New theory ideas? Data always motivated theory progress!

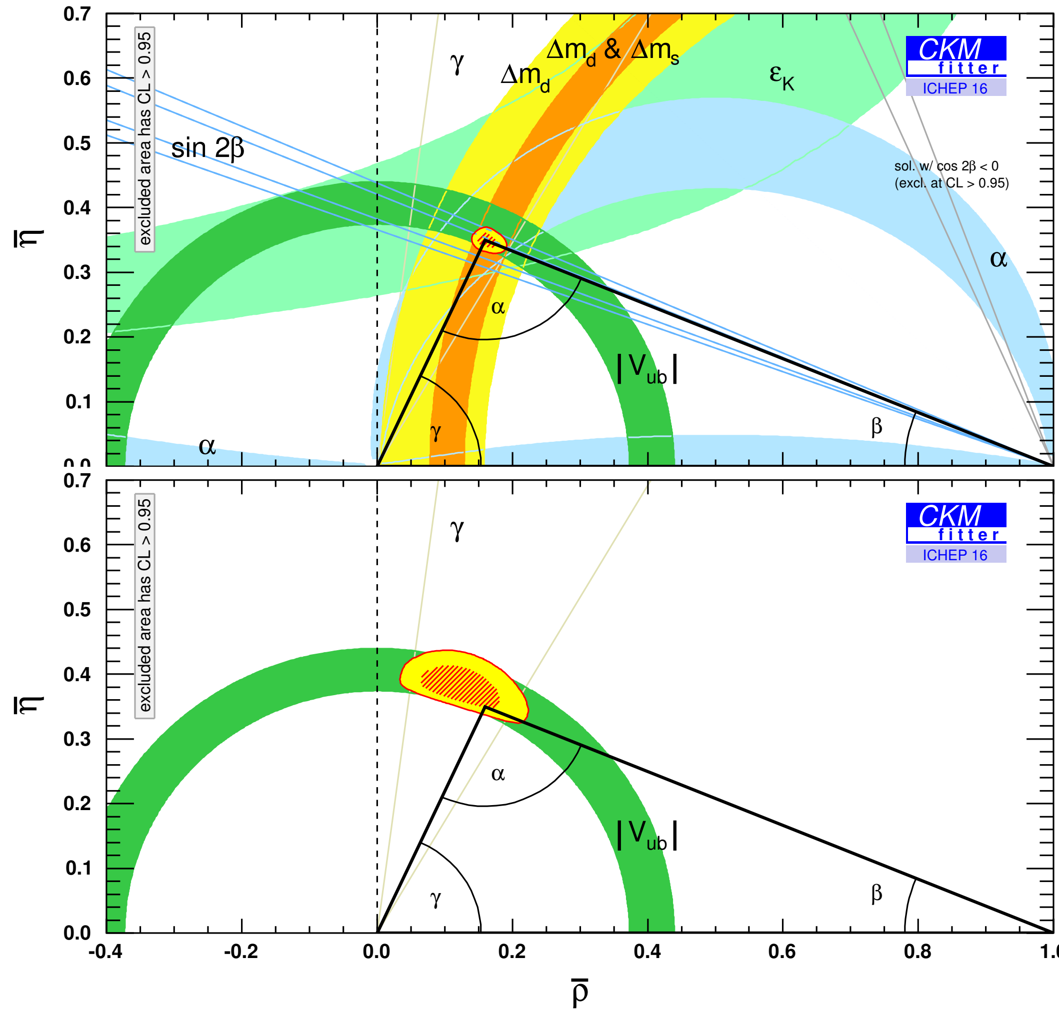
CKM fit: SM vs. NP constraints

- SM dominates CP viol. $\Rightarrow$ KM Nobel
- The implications of the consistency often overstated



CKM fit: SM vs. NP constraints

- SM dominates CP viol. $\Rightarrow$ KM Nobel
- The implications of the consistency often overstated
- Much larger allowed region if the SM is not assumed
- Tree-level (mainly V_{ub} & γ) vs. loop-dominated measurements crucial



- In loop (FCNC) processes NP / SM $\sim 20\%$ is still allowed (mixing, $B \to X\ell^+\ell^-$, $X\gamma$, etc.)

Often discussed tensions with the SM

- Intriguing tensions — could become the first clear evidence for NP
  - R_K and R_{K^*}
  - $R(D)$ and $R(D^*)$
  - P_5' and other angular distributions
  - $B_s \to \phi\mu^+\mu^-$ rate
  - $(g-2)_\mu$
  - ϵ'/ϵ

  Only $R(D^{(*)})$ is permissible at inVisibles — at least one ν in the final state :–)

  Uncertainties? What if theory uncertainty of hadronic model dependent parts is set to 100%?

- I am working on $R(D^{(*)})$, b/c theory can be improved a lot, indep't of current data

  What are the smallest deviations from the SM that can be unambiguously established?

  Likely lead (at least) to resolving the 20-some yr inclusive / exclusive $|V_{cb}|$ tension

Outline

- Use $B \to D^{(*)} l\bar{\nu}$ to refine $B \to D^{(*)} \tau\bar{\nu}$, lattice independent, improvable

  [Bernlochner, ZL, Papucci, Robinson, 1703.05330]

- MFV models, leptoquarks [Freytsis, ZL, Ruderman, 1506.08896]

  Suppress e & μ instead of enhancing τ ? [Freytsis, ZL, Ruderman, soon]

- $B \to D^{**} \ell\bar{\nu}$ in the SM and $R(D^{**})$ [Bernlochner, ZL, 1606.09300.]

  $B \to D^{**} \ell\bar{\nu}$ for arbitrary new physics [soon]

- Fully differential distributions [Robinson, ZL, Papucci, 1610.02045]

  Developing `Hammer` MC (Helicity Amplitude Module for Matrix Element Reweighting) [Bernlochner, Duell, Robinson, ZL, Papucci]

'When you think you can finally forget a topic, it's just about to become important'

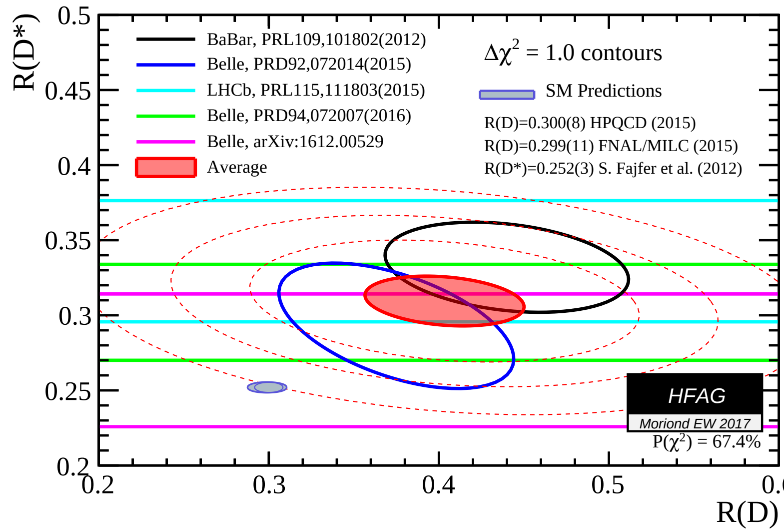
The tension with the SM

- BaBar, Belle, LHCb: $R(X) = \dfrac{\Gamma(B \to X\tau\bar\nu)}{\Gamma(B \to X(e/\mu)\bar\nu)}$

  $R(D) = 0.403 \pm 0.047\,, \quad R(D^*) = 0.310 \pm 0.017$

  June 5 @ FPCP: LHCb $\tau \to \nu 3\pi$ analysis for $R(D^*)$

  4.1σ from SM predictions — robust due to heavy quark symmetry + lattice QCD (only D so far)



- Tension: $R(D^{(*)})$ vs. $\mathcal{B}(b \to X\tau^+\nu) = (2.41 \pm 0.23)\%$ (LEP) [Freytsis, ZL, Ruderman]

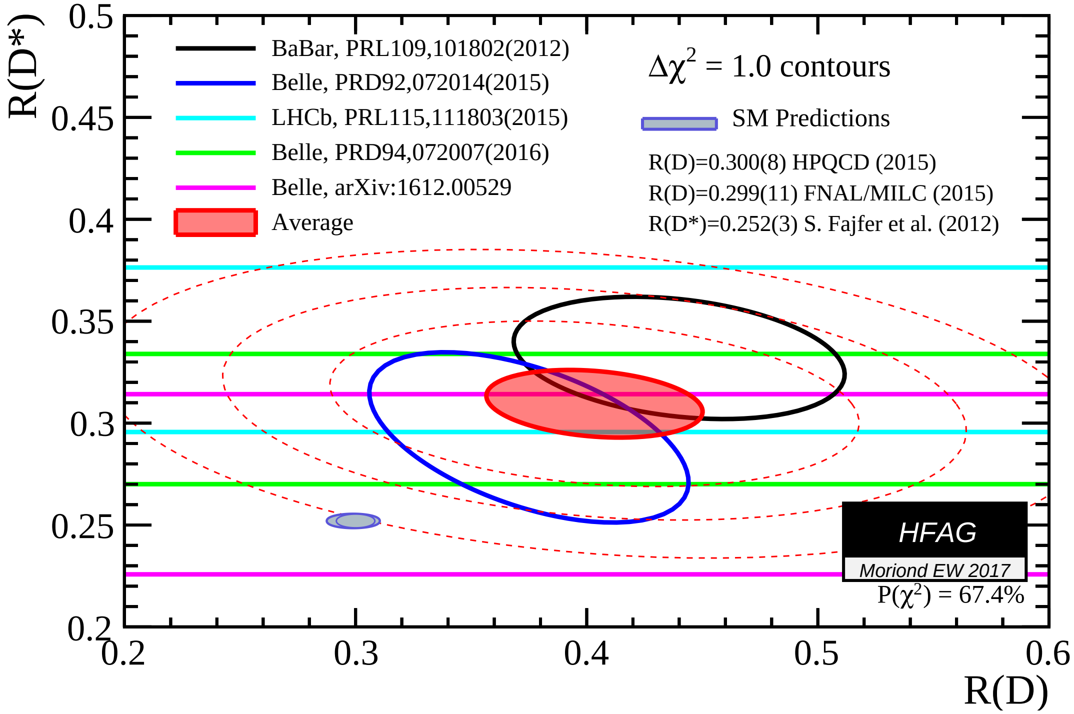
  SM: $R(X_c) = 0.223 \pm 0.004$ — no $\mathcal{B}(B \to X\tau\bar\nu)$ measurement since LEP

  Imply NP at a fairly low scale (leptoquarks, W', etc.), likely visible at the LHC

- Will become clear one way or another: forthcoming LHCb result + Belle II

- Experimental precision will improve a lot + theory uncertainty also improvable

Refining SM predictions



Can it be a theory issue?

Measured spectra for e & μ final states

- 4 functions: q^2 spectra in D & D^* + two q^2-dependent angular distrib. in D^*, $R_{1,2}$
  All form factors = Isgur-Wise function $+\Lambda_{\rm QCD}/m_{c,b} + \alpha_s$ corrections

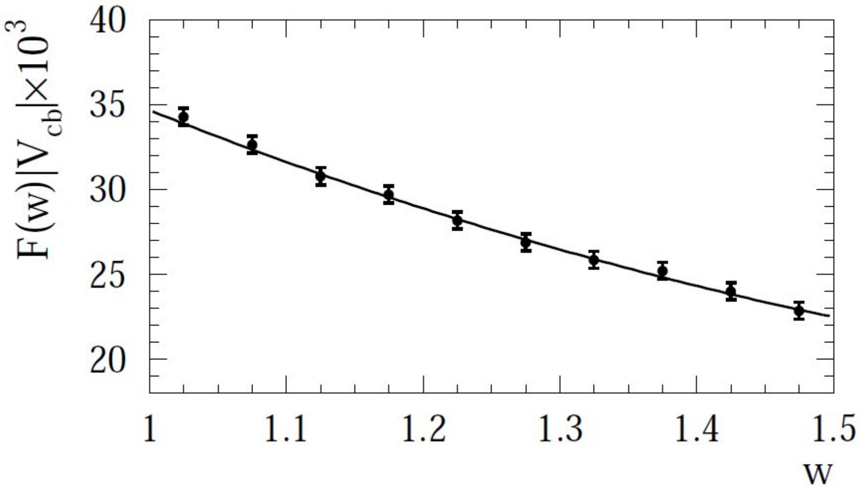
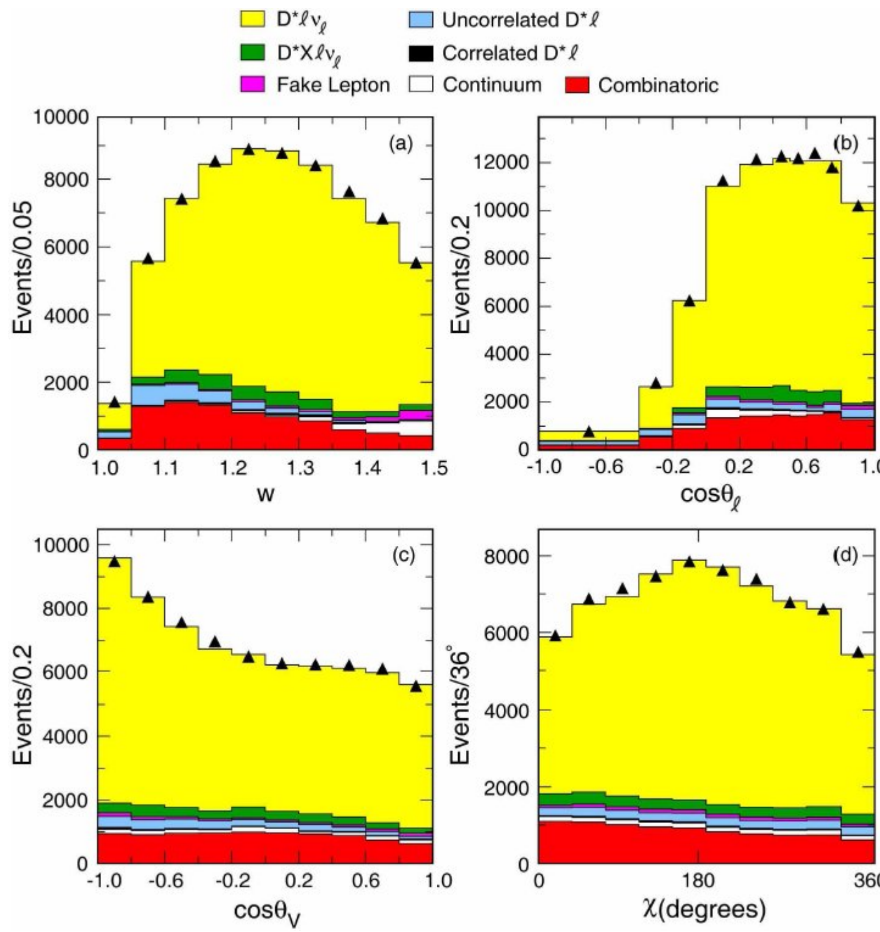


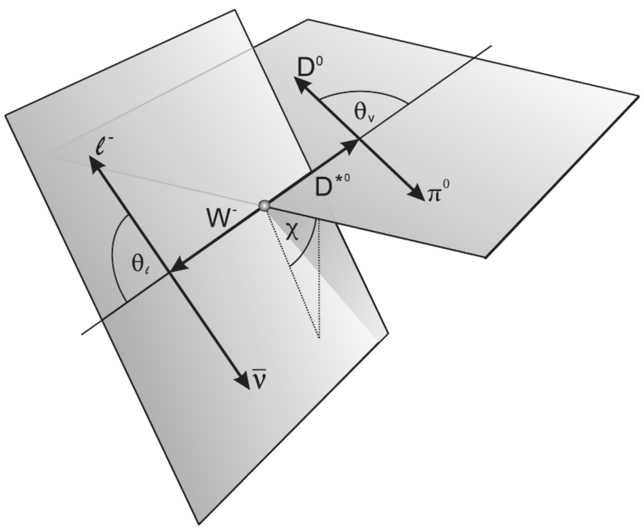
FIG. 6: The measured w dependence of $\mathcal{F}(w)|V_{cb}|$ (data points) compared to the theoretical function with the fitted parameters (solid line). The experimental uncertainties are too small to be visible.



[Plot from BaBar 0705.4008; only Belle unfolded 1510.03657, 1702.01521]

Available for the first time

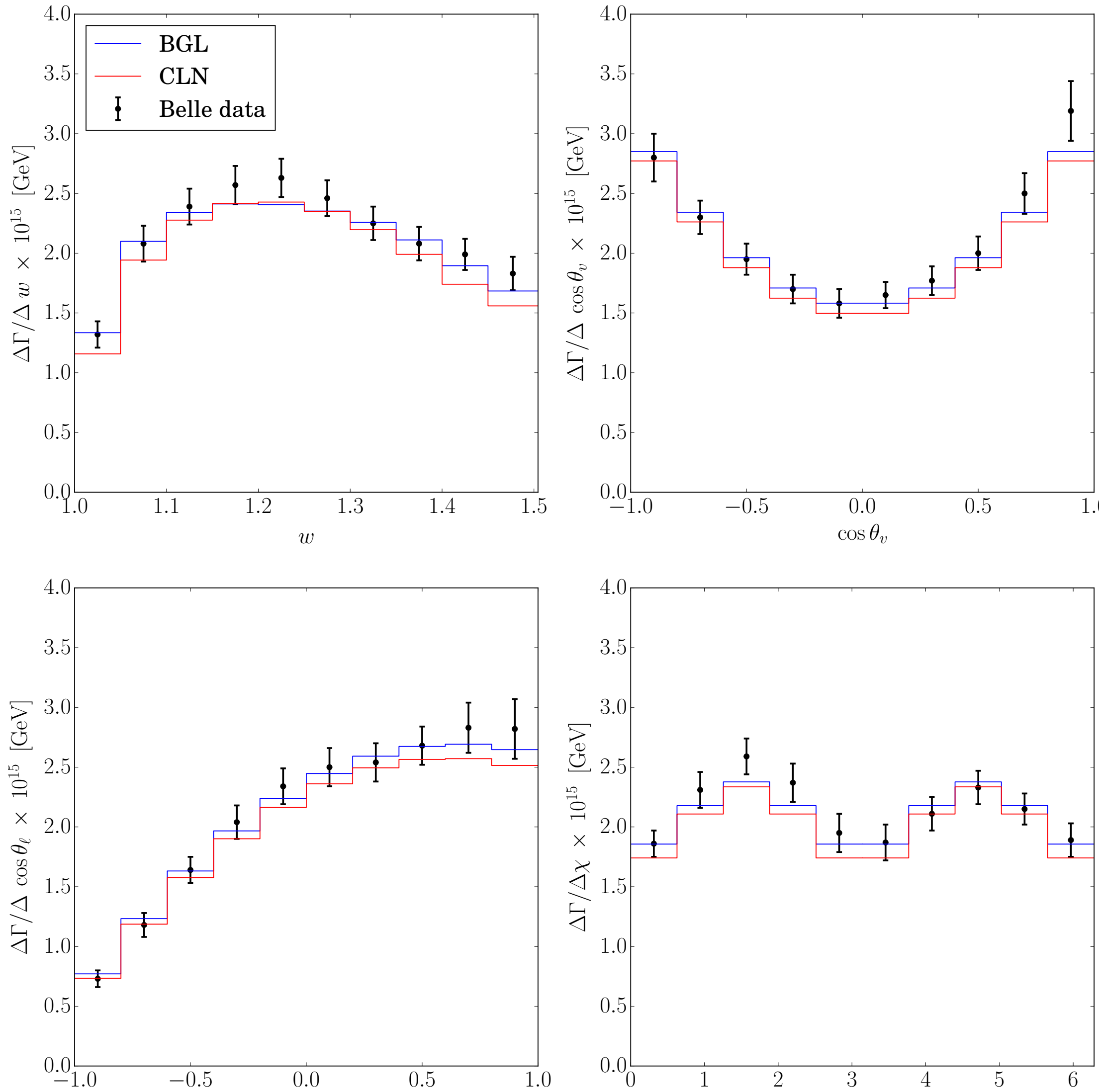
- Belle published their unfolded $B \to D^* l \bar{\nu}$ results [1702.01521]



Theorists can use it — impossible in the past

BGL = Boyd, Grinstein, Lebed, ’95–97

CLN = Caprini, Lellouch, Neubert, ’97



[1703.08170]

Basics of $B \to D^{(*)}\ell\bar{\nu}$

- Only Lorentz invariance: 6 functions of q^2, only 4 measurable with $e,\ \mu$ final states

$$\langle D|\,\bar{c}\gamma^\mu b\,|\overline{B}\rangle \;=\; f_+(q^2)(p_B+p_D)^\mu + \Big[f_0(q^2)-f_+(q^2)\Big]\frac{m_B^2-m_D^2}{q^2}\,q^\mu$$

$$\langle D^*|\,\bar{c}\gamma^\mu b\,|\overline{B}\rangle \;=\; -ig(q^2)\,\epsilon^{\mu\nu\rho\sigma}\,\varepsilon^*_\nu\,(p_B+p_{D^*})_\rho\,q_\sigma$$

$$\langle D^*|\,\bar{c}\gamma^\mu\gamma^5 b\,|\overline{B}\rangle \;=\; \varepsilon^{*\mu}f(q^2) + a_+(q^2)\,(\varepsilon^*\cdot p_B)\,(p_B+p_{D^*})^\mu + a_-(q^2)\,(\varepsilon^*\cdot p_B)\,q^\mu$$

Two form factors involving $q^\mu = p_B^\mu - p_{D^{(*)}}^\mu$ do not contribute for $m_l = 0$

- HQET constraints: 6 functions $\Rightarrow 1$ in $m_{c,b} \gg \Lambda_{\rm QCD}$ limit $+\,3$ at $\mathcal{O}(\Lambda_{\rm QCD}/m_{c,b})$

$$\langle D|\,\bar{c}\gamma^\mu b\,|\overline{B}\rangle \;=\; \sqrt{m_B m_D}\,\Big[h_+(v+v')^\mu + h_-(v-v')^\mu\Big] \qquad\qquad w = v_B\cdot v'_{D^{(*)}}$$

$$\langle D^*|\,\bar{c}\gamma^\mu b\,|\overline{B}\rangle \;=\; i\sqrt{m_B m_{D^*}}\,h_V\,\varepsilon^{\mu\nu\alpha\beta}\,\epsilon^*_\nu v'_\alpha v_\beta$$

$$\langle D^*|\,\bar{c}\gamma^\mu\gamma^5 b\,|\overline{B}\rangle \;=\; \sqrt{m_B m_{D^*}}\,\Big[h_{A_1}(w+1)\epsilon^{*\mu} - h_{A_2}(\epsilon^*\cdot v)v^\mu - h_{A_3}(\epsilon^*\cdot v)v'^\mu\Big]$$

$m_{c,b} \gg \Lambda_{\rm QCD}$ limit: $h_+ = h_V = h_{A_1} = h_{A_3} = \xi(w)$ and $h_- = h_{A_2} = 0$

- Constrain all 4 functions from $B \to D^{(*)}l\bar{\nu}$ $\Rightarrow$ $\mathcal{O}(\Lambda_{\rm QCD}^2/m_{c,b}^2\,,\ \alpha_s^2)$ uncertainties

Form factor expansion details

- Expand form factors to order $\varepsilon_{c,b} = \Lambda_{\rm QCD}/(2m_{c,b})$ and α_s (new results for tensor ff)

$$f_i(w) = \xi(w)\bigg[1 + \varepsilon_c\, f_i^{(c,1)}(w) + \varepsilon_b\, f_i^{(b,1)}(w) + \alpha_s\, f_i^{(\alpha_s)}\Big(\frac{m_c}{m_b}, w\Big) + \mathcal{O}(\varepsilon_{c,b}^2\,, \alpha_s^2)\bigg]$$

Absorbed $\xi(w) \to \xi(w) + 2(\varepsilon_c + \varepsilon_b)\chi_1(w)$, so only $\chi_{2,3}$ and $\eta = \xi_3/\xi$ remain

Known for SM terms since the early 90s, but not written down for others before

The $\alpha_s\, \varepsilon_{c,b}$ terms are known, should be included if NP established
Expect that fit readjusts subleading Isgur-Wise functions $\Rightarrow$ modest impacts

- $\chi_{2,3}$ & η calculated in QCD sum rules — parametrize: [ZL, Neubert, Nir, '92–93]

$1/m$ Lagrangian: $\hat\chi_2^{\rm ren}(1) = -0.06 \pm 0.02$ $\quad \hat\chi_2^{\prime\,\rm ren}(1) = 0 \pm 0.02$ $\quad \hat\chi_3^{\prime\,\rm ren}(1) = 0.04 \pm 0.02$

$1/m$ current: $\eta(1) = 0.62 \pm 0.2\,, \quad \eta'(1) = 0 \pm 0.2$ (Luke's thm. $\Rightarrow \hat\chi_3(1) = 0$)

Central values match what CLN used, these uncertainties $>$ in original papers

Inputs and $|V_{cb}|$ fits

- Lattice QCD: $B \to D$ at $w = 1,\ 1.08,\ 1.16$
  $B \to D^*$ at $w = 1$

- Analyticity-based constraints on shapes of form factors

  BGL: no HQET relations in parametrization, treat 3 form factors as unrelated

  CLN: use HQET + QCD sum rules for $\mathcal{O}(\Lambda_{\rm QCD}/m_{c,b})$, no uncertainties assigned
  more caveats in practical implementations

- Fewer fit parameters in CLN, used by all experimental measurements since '97

  Used also in theory papers (except lattice) to derive SM predictions for $R(D^{(*)})$

Bigi, Gambino, Schacht, 1703.06124, $|V_{cb}|_{\rm BGL} = (41.7^{+2.0}_{-2.1}) \times 10^{-3}$

Grinstein & Kobach, 1703.08170, $|V_{cb}|_{\rm BGL} = (41.9^{+2.0}_{-1.9}) \times 10^{-3}$

Belle, 1702.01521, $|V_{cb}|_{\rm CLN} = (37.4 \pm 1.3) \times 10^{-3}$ $(38.2 \pm 1.5$ in 1703.06124)

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Consider 7 different fit scenarios

- All calculations of subleading $\Lambda_{\rm QCD}/m_{c,b}$ Isgur-Wise functions model dependent

  Only $R(D)$ calculated in LQCD — all others did not include uncertainties properly

- Theory [CLN] & exp papers: $R_{1,2}(w) = \underbrace{R_{1,2}(1)}_{\rm fit} + \underbrace{R'_{1,2}(1)}_{\rm fixed}(w-1) + \underbrace{R''_{1,2}(1)}_{\rm fixed}(w-1)^2/2$

  In HQET: $R_{1,2}(1) = 1 + \mathcal{O}(\Lambda_{\rm QCD}/m_{c,b}\,, \alpha_s)$ $\qquad R^{(n)}_{1,2}(1) = 0 + \mathcal{O}(\Lambda_{\rm QCD}/m_{c,b}\,, \alpha_s)$

  Sometimes calculations using QCD sum rule predictions for $\Lambda_{\rm QCD}/m_{c,b}$ corrections are called the HQET predictions

- Our fits:

Fit	QCDSR	Lattice QCD $\mathcal{F}(1)$	Lattice QCD $f_{+,0}(1)$	Lattice QCD $f_{+,0}(w>1)$	Belle Data
L$_{w=1}$	—	+	+	—	+
L$_{w=1}$+SR	+	+	+	—	+
NoL	—	—	—	—	+
NoL+SR	+	—	—	—	+
L$_{w\geq1}$	—	+	+	+	+
L$_{w\geq1}$+SR	+	+	+	+	+
th:L$_{w\geq1}$+SR	+	+	+	+	—

Aside: Fit details

- Standard choice to minimize range of expansion param' z_* in unitarity constraints:

$$z_*(w) = \frac{\sqrt{w+1} - \sqrt{2}\,a}{\sqrt{w+1} + \sqrt{2}\,a}\,, \qquad a = \left(\frac{1+r_D}{2\,\sqrt{r_D}}\right)^{1/2}$$

- Parametrize similar to CLN — wanted to start with fit comparable to prior results

$$\frac{\mathcal{G}(w)}{\mathcal{G}(w_0)} \simeq 1 - 8a^2\rho_*^2\, z_* + \Big(V_{21}\rho_*^2 - V_{20}\Big)\, z_*^2$$

Translate this to $\xi(w)/\xi(w_0)$ to be able to simultaneously fit $B \to D$ and $B \to D^*$

Uncertainty in z_*^2 term may be sizable — we checked that fit results are stable if constraint between the slope and the curvature is relaxed

Keep uncertainties and correlations in form factor ratios ($\Lambda_{\rm QCD}/m$ Isgur-Wise fn's)

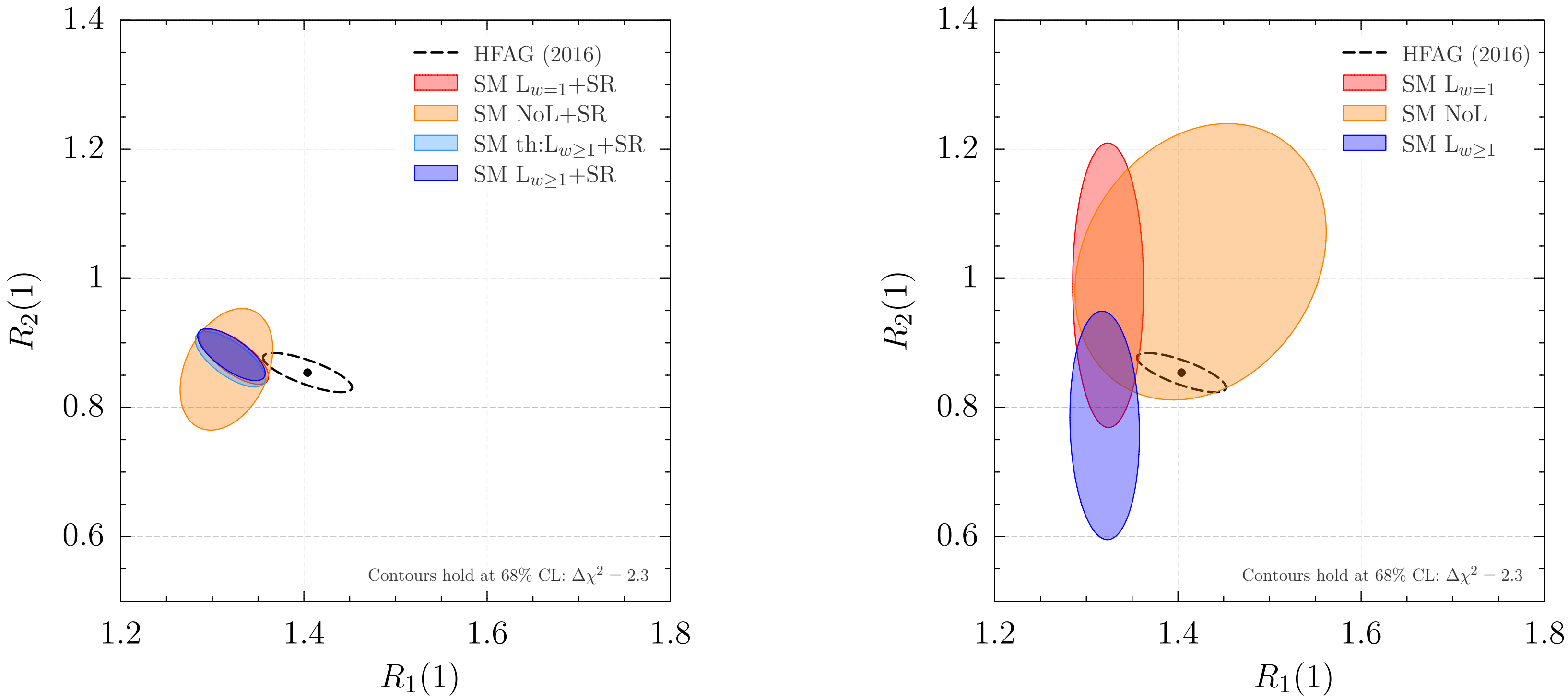
- In progress: study systematically orders/constraints in fit, HQET corrections, etc.

Experimental inputs and self-consistency

- Experimental inputs: $B \to D l\bar{\nu}: \ \mathrm{d}\Gamma/\mathrm{d}w$ (Only Belle published fully corrected distributions)

  $B \to D^* l\bar{\nu}: \ \mathrm{d}\Gamma/\mathrm{d}w,\ R_1(w),\ R_2(w)$

  Model-dependent inputs in SM predictions for $R_{1,2}$ in all exp. fits & theory papers

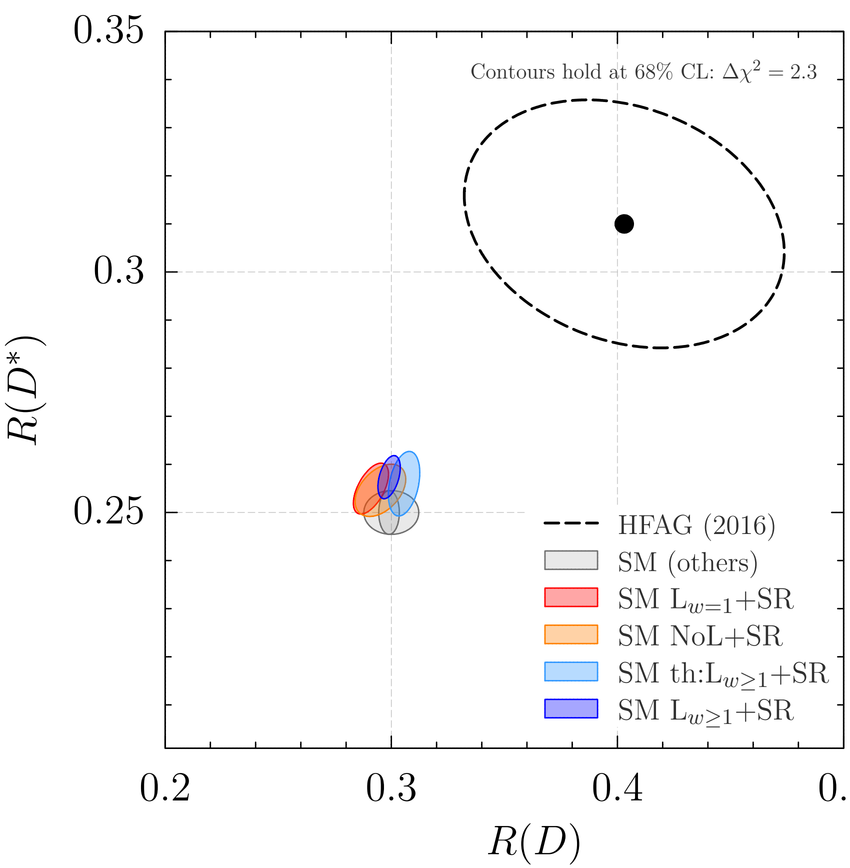


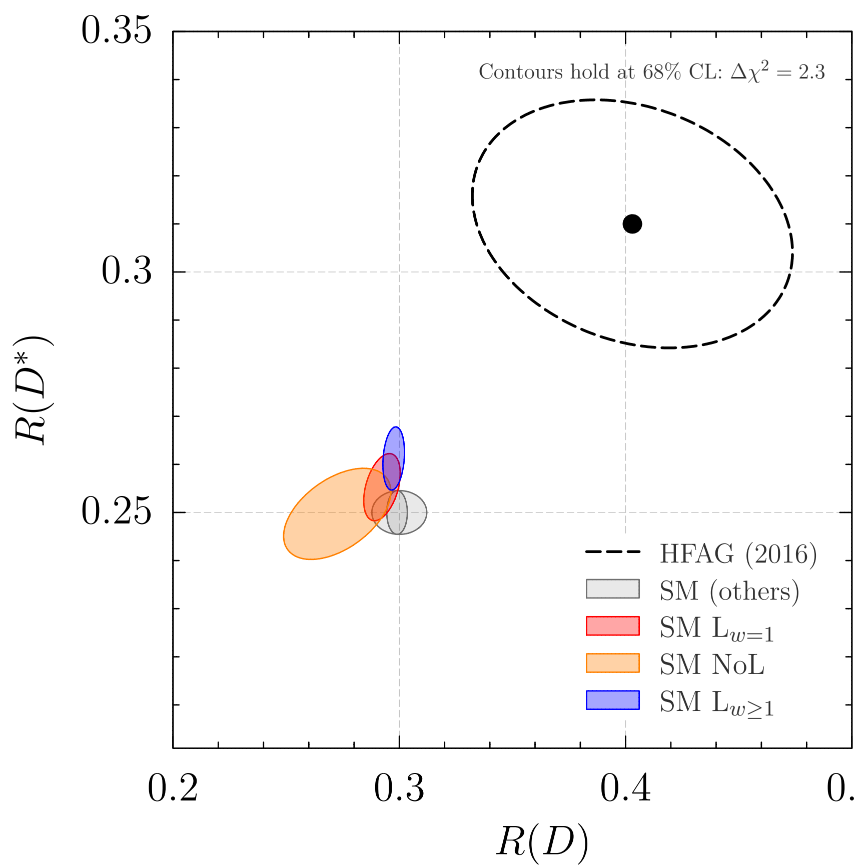
- Mild tension for $R_1(1)$ — may affect $|V_{cb}|$ from $B \to D^{(*)} l\bar{\nu}$, long standing issues

  In $1S$ scheme: $R_1(1) \simeq 1.34 - 0.12\,\eta(1)\,,\qquad R_2(1) \simeq 0.98 - 0.42\,\eta(1) - 0.54\,\hat{\chi}_2(1)$

Our SM predictions for $R(D)$ and $R(D^*)$

- Significance of the tension is (surprisingly) stable across our fit scenarios:





- Fit just a quadratic polynomial in z_*: consistent results

Summary of SM predictions

- Small variations: heavy quark symmetry & phase space leave little wiggle room

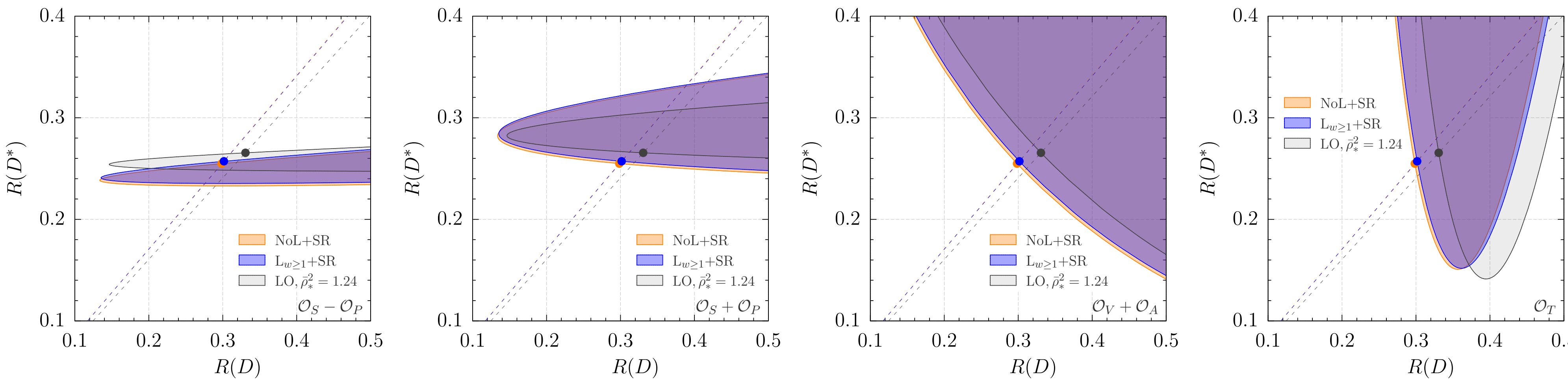
Scenario	$R(D)$	$R(D^*)$	Correlation
$\mathrm{L}_{w=1}$	0.292 ± 0.005	0.255 ± 0.005	41%
$\mathrm{L}_{w=1}$+SR	0.291 ± 0.005	0.255 ± 0.003	57%
NoL	0.273 ± 0.016	0.250 ± 0.006	49%
NoL+SR	0.295 ± 0.007	0.255 ± 0.004	43%
$\mathrm{L}_{w\geq 1}$	0.298 ± 0.003	0.261 ± 0.004	19%
$\mathrm{L}_{w\geq 1}$+SR	$\mathbf{0.299 \pm 0.003}$	$\mathbf{0.257 \pm 0.003}$	44%
th:$\mathrm{L}_{w\geq 1}$+SR	0.306 ± 0.005	0.256 ± 0.004	33%
Data [HFAG]	0.403 ± 0.047	0.310 ± 0.017	-23%
Lattice [FLAG]	0.300 ± 0.008	—	—
Bigi, Gambino '16	0.299 ± 0.003	—	—
Fajfer et al. '12	—	0.252 ± 0.003	—

- Tension between our "$\mathrm{L}_{w\geq 1}$+SR" fit and data is 3.9σ, with p-value $= 11.5 \times 10^{-5}$

  (close to HFAG: 3.9σ, with p-value $= 8.3 \times 10^{-5}$)

Impact on new physics effects

- Add only one NP operator to the SM at a time: $O_S - O_P,\ O_S + O_P,\ O_V + O_A,\ O_T$



- Not all $1/m$ corrections in literature, some $\mathcal{O}(1/m)$ form factors had 100% uncert. (i.e., tensor currents vanishing in heavy quark limit)
- Shifts from gray regions non-negligible — if one seriously wanted to fit a NP model

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New physics options

Consider fits to redundant set of operators

- Likely tree-level: different fermion orderings convenient to understand mediators

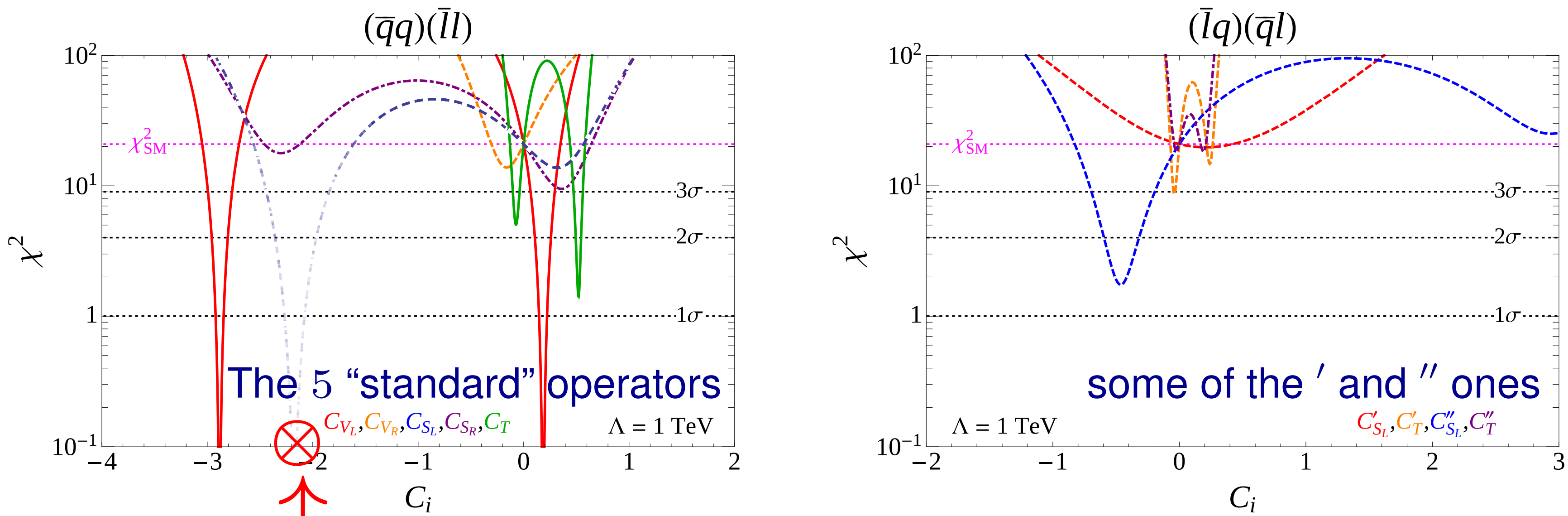
Usually only the first 5 operators considered, related by Fierz

from dim-6 terms, others from dim-8 only ⇓

	Operator	Fierz identity	Allowed Current	$\delta\mathcal{L}_{\rm int}$
$\mathcal{O}_{V_L}$	$(\bar c\gamma_\mu P_L b)\,(\bar\tau\gamma^\mu P_L\nu)$		$(1,3)_0$	$(g_q\bar q_L\boldsymbol{\tau}\gamma^\mu q_L + g_\ell\bar\ell_L\boldsymbol{\tau}\gamma^\mu\ell_L)W'_\mu$
$\mathcal{O}_{V_R}$	$(\bar c\gamma_\mu P_R b)\,(\bar\tau\gamma^\mu P_L\nu)$			
$\mathcal{O}_{S_R}$	$(\bar c P_R b)\,(\bar\tau P_L\nu)$		$\rangle(1,2)_{1/2}$	$(\lambda_d\bar q_L d_R\phi + \lambda_u\bar q_L u_R i\tau_2\phi^\dagger + \lambda_\ell\bar\ell_L e_R\phi)$
$\mathcal{O}_{S_L}$	$(\bar c P_L b)\,(\bar\tau P_L\nu)$			
$\mathcal{O}_T$	$(\bar c\sigma^{\mu\nu} P_L b)\,(\bar\tau\sigma_{\mu\nu} P_L\nu)$			
$\mathcal{O}'_{V_L}$	$(\bar\tau\gamma_\mu P_L b)\,(\bar c\gamma^\mu P_L\nu)$	$\longleftrightarrow\ \mathcal{O}_{V_L}\langle$	$(3,3)_{2/3}$	$\lambda\,\bar q_L\boldsymbol{\tau}\gamma_\mu\ell_L\boldsymbol{U}^\mu$
			$\rangle(3,1)_{2/3}$	$(\lambda\,\bar q_L\gamma_\mu\ell_L + \tilde\lambda\,\bar d_R\gamma_\mu e_R)U^\mu$
$\mathcal{O}'_{V_R}$	$(\bar\tau\gamma_\mu P_R b)\,(\bar c\gamma^\mu P_L\nu)$	$\longleftrightarrow\ -2\mathcal{O}_{S_R}$		
$\mathcal{O}'_{S_R}$	$(\bar\tau P_R b)\,(\bar c P_L\nu)$	$\longleftrightarrow\ -\frac{1}{2}\mathcal{O}_{V_R}$		
$\mathcal{O}'_{S_L}$	$(\bar\tau P_L b)\,(\bar c P_L\nu)$	$\longleftrightarrow\ -\frac{1}{2}\mathcal{O}_{S_L} - \frac{1}{8}\mathcal{O}_T$	$(3,2)_{7/6}$	$(\lambda\,\bar u_R\ell_L + \tilde\lambda\,\bar q_L i\tau_2 e_R)R$
$\mathcal{O}'_T$	$(\bar\tau\sigma^{\mu\nu} P_L b)\,(\bar c\sigma_{\mu\nu} P_L\nu)$	$\longleftrightarrow\ -6\mathcal{O}_{S_L} + \frac{1}{2}\mathcal{O}_T$		
$\mathcal{O}''_{V_L}$	$(\bar\tau\gamma_\mu P_L c^c)\,(\bar b^c\gamma^\mu P_L\nu)$	$\longleftrightarrow\ -\mathcal{O}_{V_R}$		
$\mathcal{O}''_{V_R}$	$(\bar\tau\gamma_\mu P_R c^c)\,(\bar b^c\gamma^\mu P_L\nu)$	$\longleftrightarrow\ -2\mathcal{O}_{S_R}$	$(\bar 3,2)_{5/3}$	$(\lambda\,\bar d^c_R\gamma_\mu\ell_L + \tilde\lambda\,\bar q^c_L\gamma_\mu e_R)V^\mu$
$\mathcal{O}''_{S_R}$	$(\bar\tau P_R c^c)\,(\bar b^c P_L\nu)$	$\longleftrightarrow\ \frac{1}{2}\mathcal{O}_{V_L}\langle$	$(\bar 3,3)_{1/3}$	$\lambda\,\bar q^c_L i\tau_2\boldsymbol{\tau}\ell_L\boldsymbol{S}$
$\mathcal{O}''_{S_L}$	$(\bar\tau P_L c^c)\,(\bar b^c P_L\nu)$	$\longleftrightarrow\ -\frac{1}{2}\mathcal{O}_{S_L} + \frac{1}{8}\mathcal{O}_T$	$\rangle(\bar 3,1)_{1/3}$	$(\lambda\,\bar q^c_L i\tau_2\ell_L + \tilde\lambda\,\bar u^c_R e_R)S$
$\mathcal{O}''_T$	$(\bar\tau\sigma^{\mu\nu} P_L c^c)\,(\bar b^c\sigma_{\mu\nu} P_L\nu)$	$\longleftrightarrow\ -6\mathcal{O}_{S_L} - \frac{1}{2}\mathcal{O}_T$		

[Freytsis, ZL, Ruderman, 1506.08896]

Fits to a single operator

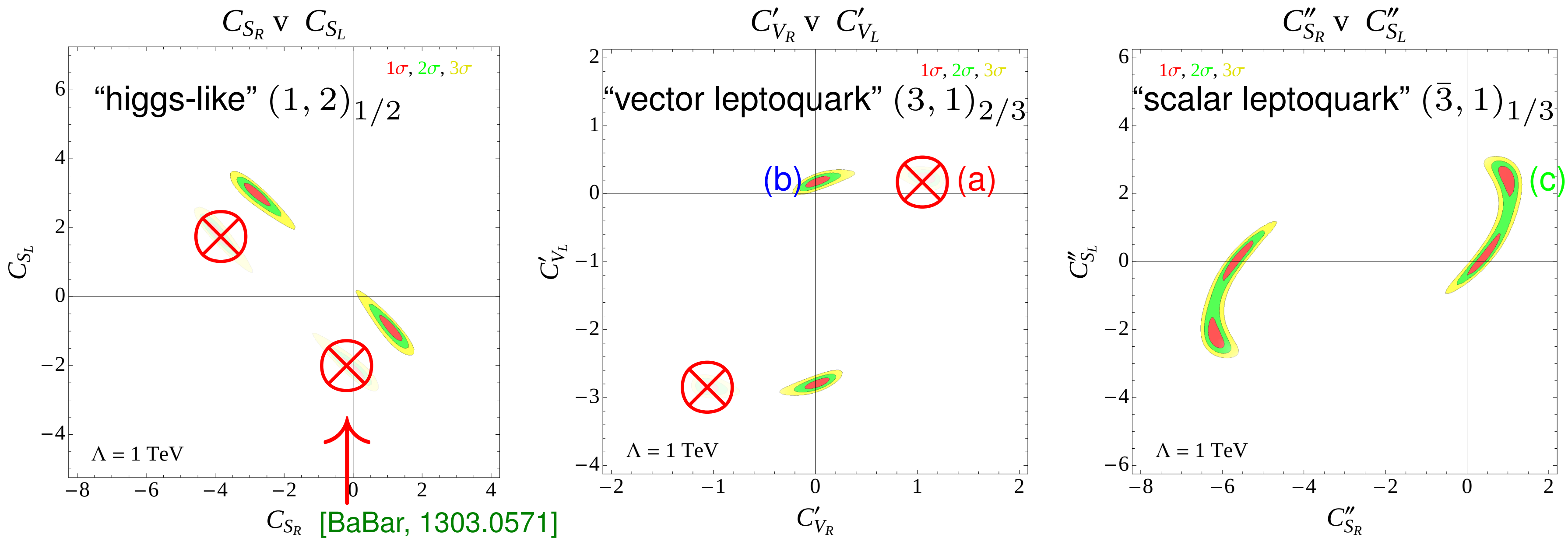


Ruled out by the BaBar q^2 spectrum [1303.0571]

- Large coefficients, $\Lambda = 1\,\mathrm{TeV}$ in plots $\Rightarrow$ fairly light mediators (obvious: 20–30% of a tree-level rate)

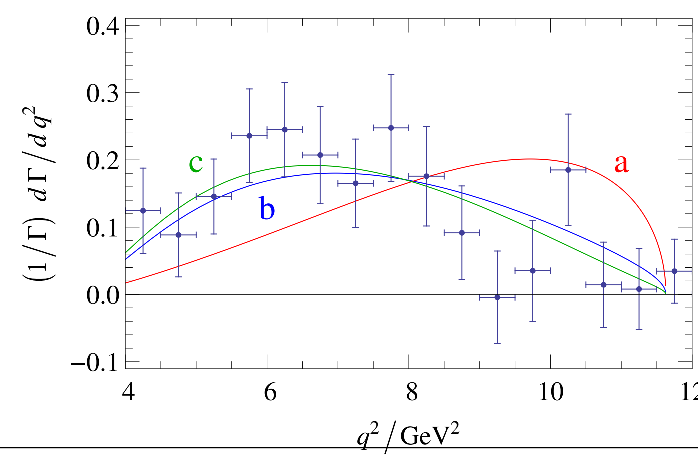
In HQET limit, we confirmed the “classic” paper [Goldberger, hep-ph/9902311]

Fits to two operators



The $\otimes$ solution are ruled out by the q^2 spectrum

Operator coefficients	
$C'_{V_L} = 0.24$	$C'_{V_R} = 1.10$
$C'_{V_L} = 0.18$	$C'_{V_R} = -0.01$
$C''_{S_R} = 0.96$	$C''_{S_L} = 2.41$



Operator fits $\to$ viable MFV models?

- Good fits for several mediators: scalar, "Higgs-like" $(1,2)_{1/2}$
  vector, "W'-like" $(1,3)_0$
  "scalar leptoquark" $(\bar{3},1)_{1/3}$ or $(\bar{3},3)_{1/3}$
  "vector leptoquark" $(3,1)_{2/3}$ or $(3,3)_{2/3}$

  We did not try to fit any of the other anomalies simultaneously

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- Which BSM scenarios can be MFV? [Freytsis, ZL, Ruderman, 1506.08896]

  Viable leptoquarks: scalar $S(\mathbf{1},\mathbf{1},\bar{\mathbf{3}})$ or vector $U_\mu(\mathbf{1},\mathbf{1},\mathbf{3})$

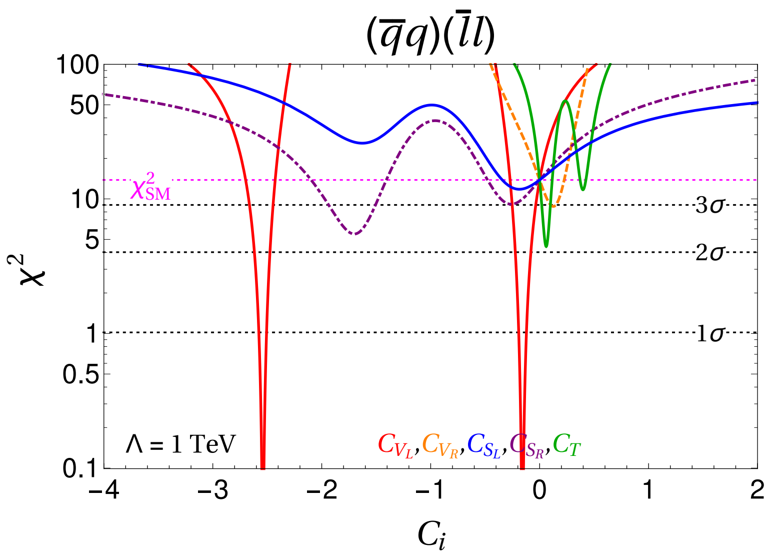
  Bounds: $b \to s\nu\bar{\nu}$, D^0 & K^0 mixing, $Z \to \tau^+\tau^-$, LHC contact int., $pp \to \tau^+\tau^-$, etc.

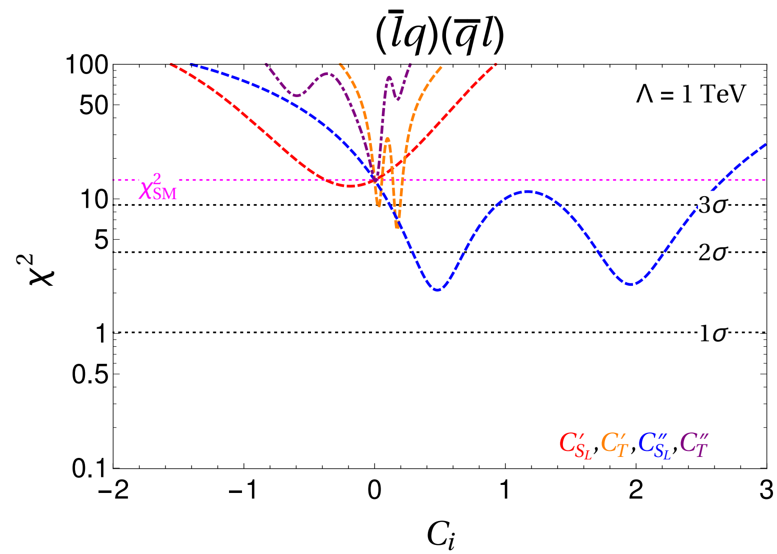
  In this case there is no $bb\tau\tau$ coupling [See Greljo's talk yesterday for many other options]

How odd scenarios may be viable?

- All papers enhance the τ mode compared to the SM

  Can one suppress the e and μ modes instead? [Freytsis, ZL, Ruderman, to appear]





- Unique viable option: modify the SM four-fermion operator

  Good fit with: $V_{cb}^{(\rm exp)} \sim V_{cb}^{(\rm SM)} \times 0.9 \qquad V_{ub}^{(\rm exp)} \sim V_{ub}^{(\rm SM)} \times 0.9$

- Many relevant constraints, some of the strongest from ϵ_K and B mixing

What about $e - \mu$ (non)universality?

- How well is the difference of the e and μ rates constrained?

Parameters	De sample	$D\mu$ sample	combined result
ρ_D^2	$1.22 \pm 0.05 \pm 0.10$	$1.10 \pm 0.07 \pm 0.10$	$1.16 \pm 0.04 \pm 0.08$
$\rho_{D^*}^2$	$1.34 \pm 0.05 \pm 0.09$	$1.33 \pm 0.06 \pm 0.09$	$1.33 \pm 0.04 \pm 0.09$
R_1	$1.59 \pm 0.09 \pm 0.15$	$1.53 \pm 0.10 \pm 0.17$	$1.56 \pm 0.07 \pm 0.15$
R_2	$0.67 \pm 0.07 \pm 0.10$	$0.68 \pm 0.08 \pm 0.10$	$0.66 \pm 0.05 \pm 0.09$
$\mathcal{B}(D^0\ell\overline{\nu})(\%)$	$2.38 \pm 0.04 \pm 0.15$	$2.25 \pm 0.04 \pm 0.17$	$2.32 \pm 0.03 \pm 0.13$
$\mathcal{B}(D^{*0}\ell\overline{\nu})(\%)$	$5.50 \pm 0.05 \pm 0.23$	$5.34 \pm 0.06 \pm 0.37$	$5.48 \pm 0.04 \pm 0.22$
χ^2/n.d.f. (probability)	416/468 (0.96)	488/464 (0.21)	2.0/6 (0.92)

[BaBar, 0809.0828 — similar results in Belle, 1010.5620]

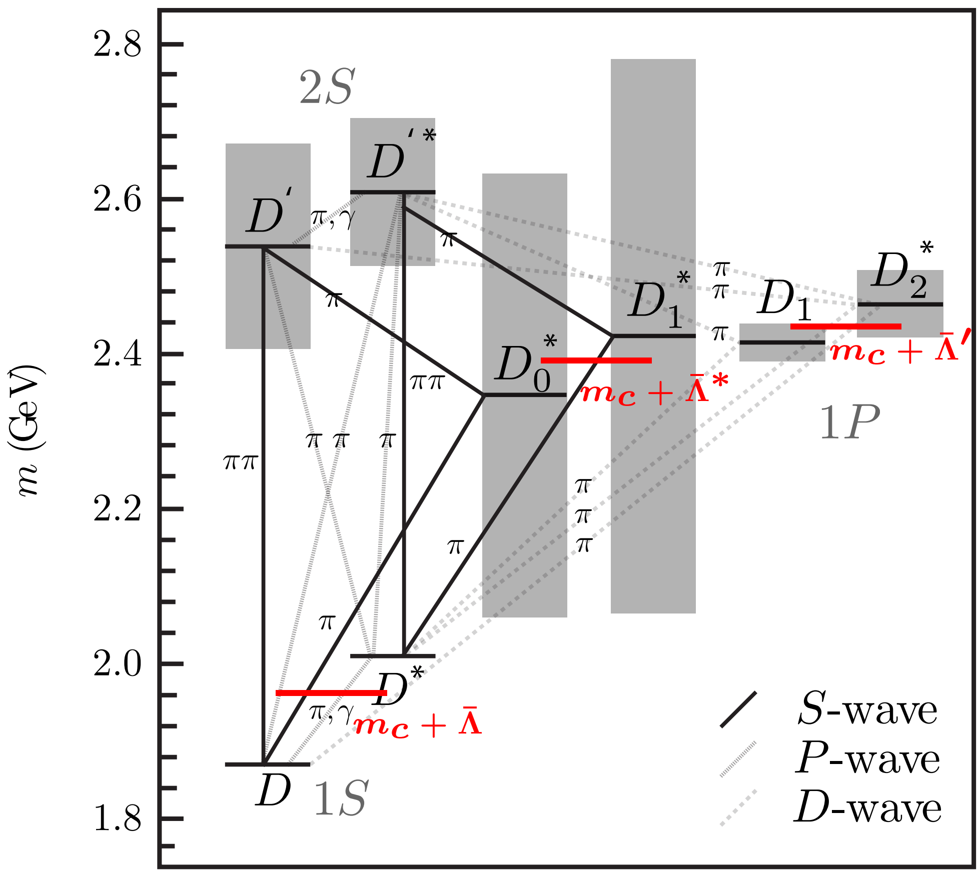
- 10% difference allowed... some wrong statements...

Γ_1	$e^+\nu_e$ anything	$(10.86 \pm 0.16)\%$
Γ_2	$\overline{p}e^+\nu_e$ anything	$< 5.9 \times 10^{-4}$
Γ_3	$\mu^+\nu_\mu$ anything	$(10.86 \pm 0.16)\%$
Γ_4	$\ell^+\nu_\ell$ anything	$(10.86 \pm 0.16)\%$

- How much better can difference be constrained?

  Reaching the 1% level on ratio might be possible (but challenging) at Belle II

$B \to D^{**}\tau\bar{\nu}$



Particle	$s_l^{\pi_l}$	J^P	m (MeV)	Γ (MeV)
D_0^*	$\frac{1}{2}^+$	0^+	2330	270
D_1^*	$\frac{1}{2}^+$	1^+	2427	384
D_1	$\frac{3}{2}^+$	1^+	2421	34
D_2^*	$\frac{3}{2}^+$	2^+	2462	48

Parameter	$\bar{\Lambda}$	$\bar{\Lambda}'$	$\bar{\Lambda}^*$
Value [GeV]	0.40	0.80	0.76

Why bother...?

- $B \to D^{**}\tau\bar\nu$: rates to narrow D_1, D_2^* measurable? No predictions

  In $B_s \to D_s^{**}\ell\bar\nu$ case, all 4 D_s^{**} states are narrow $\Rightarrow$ LHCb?

- Largest syst. uncertainty in $R(D^{(*)})$
- May matter for tensions between inclusive and exclusive $|V_{cb}|$ and $|V_{ub}|$ determinations
- Complementary sensitivity to NP
- Complementary experimentally
- Decay rates not too small

	$R(D)$ [%]	$R(D^*)$ [%]	Correlation
$D^{(*(*))}\ell\nu$ shapes	4.2	1.5	0.04
D^{**} composition	1.3	3.0	-0.63
Fake D yield	0.5	0.3	0.13
Fake ℓ yield	0.5	0.6	-0.66
D_s yield	0.1	0.1	-0.85
Rest yield	0.1	0.0	-0.70
Efficiency ratio f^{D^+}	2.5	0.7	-0.98
Efficiency ratio f^{D^0}	1.8	0.4	0.86
Efficiency ratio $f_{\rm eff}^{D^{*+}}$	1.3	2.5	-0.99
Efficiency ratio $f_{\rm eff}^{D^{*0}}$	0.7	1.1	0.94
CF double ratio g^+	2.2	2.0	-1.00
CF double ratio g^0	1.7	1.0	-1.00
Efficiency ratio $f_{\rm wc}$	0.0	0.0	0.84
$M^2_{\rm miss}$ shape	0.6	1.0	0.00
$o'_{\rm NB}$ shape	3.2	0.8	0.00
Lepton PID efficiency	0.5	0.5	1.00
Total	7.1	5.2	−0.32

[Belle, 1507.03233]

Some model independent results

- At $w \equiv v \cdot v' = 1$, the $\mathcal{O}(\Lambda_{\rm QCD}/m_{c,b})$ matrix element is determined by masses and leading order Isgur-Wise function [Leibovich, ZL, Stewart, Wise, hep-ph/9703213, hep-ph/9705467]

  Kinematic range: $1 \leq w \lesssim 1.3$ and in the τ case $1 \leq w \lesssim 1.2$

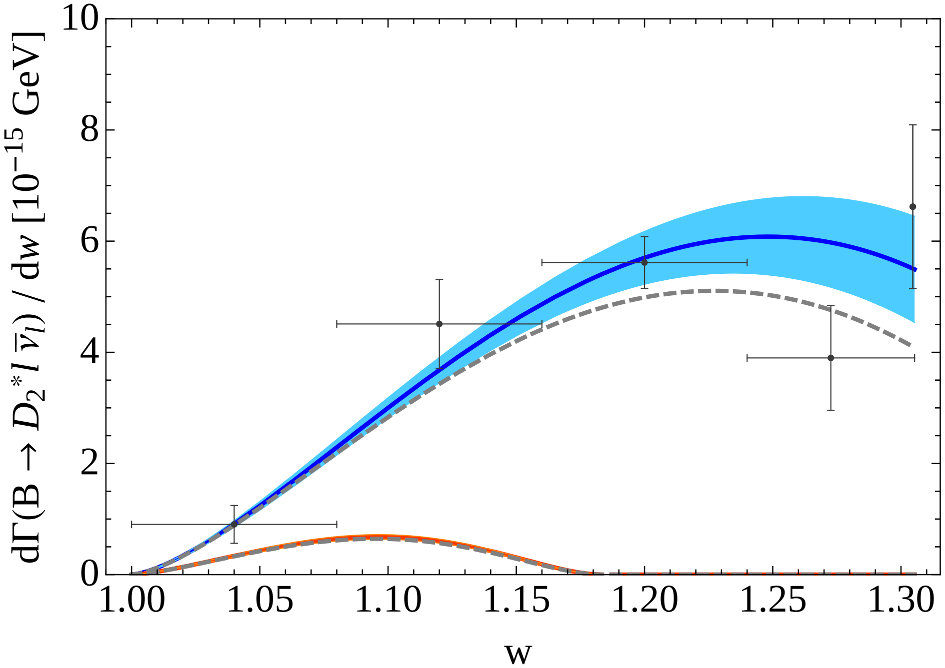
  Meson masses: $$m_{H_\pm} = m_Q + \bar\Lambda^H - \frac{\lambda_1^H}{2m_Q} \pm \frac{n_\mp\, \lambda_2^H}{2m_Q} + \ldots \qquad n_\pm = 2J_\pm + 1$$

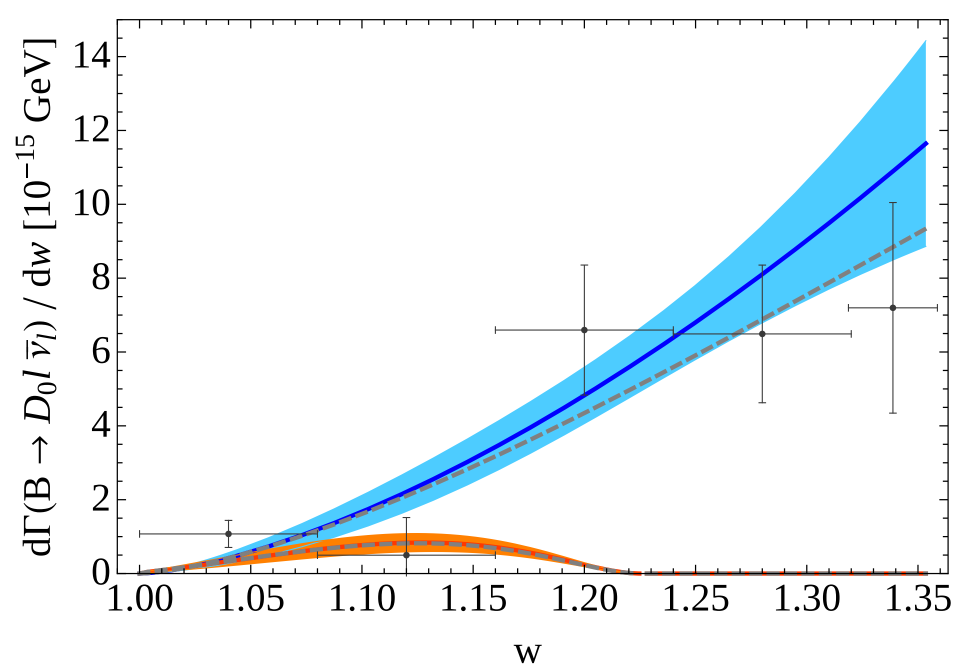
  For example:
$$\frac{\langle D_1(v',\epsilon)|V^\mu|B(v)\rangle}{\sqrt{m_{D_1}\, m_B}} = f_{V_1}\epsilon^{*\mu} + (f_{V_2} v^\mu + f_{V_3} v'^\mu)(\epsilon^* \cdot v)$$
$$\sqrt{6}\, f_{V_1}(w) = (1-w^2)\,\tau(w) - 4\,\frac{\bar\Lambda' - \bar\Lambda}{m_c}\,\tau(w) + \mathcal{O}\bigg(\frac{w-1}{m_{c,b}}\bigg) + \ldots$$

- These “known” $\mathcal{O}(\Lambda_{\rm QCD}/m_{c,b})$ terms are numerically very important

- No expressions in the literature for $B \to D^{**}\,\tau\bar\nu$ rates at all — fixing this...

Predictions for spectra





Rates for $\boldsymbol{e}$, $\boldsymbol{\mu}$ vs. $\boldsymbol{\tau}$

[Data from Belle, 0711.3252]

- Study all uncertainties, including effects neglected in LLSW
- As for $B \to D^{(*)} \ell \bar{\nu}$, heavy quark symmetry relates the extra form factor in the τ mode to those with e, μ — finalizing the uncertainties

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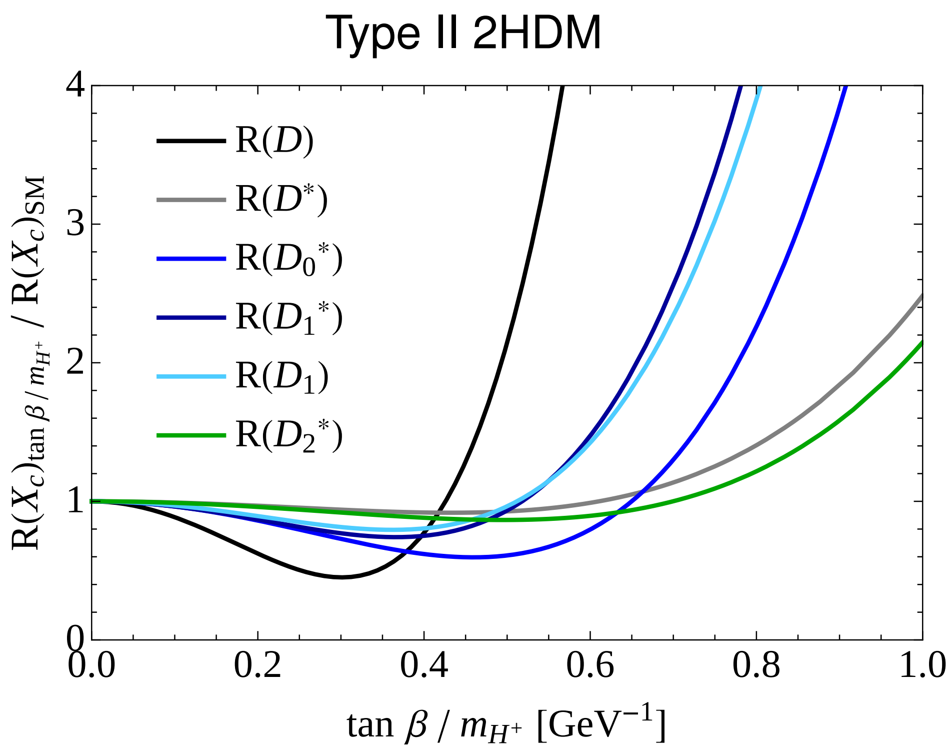
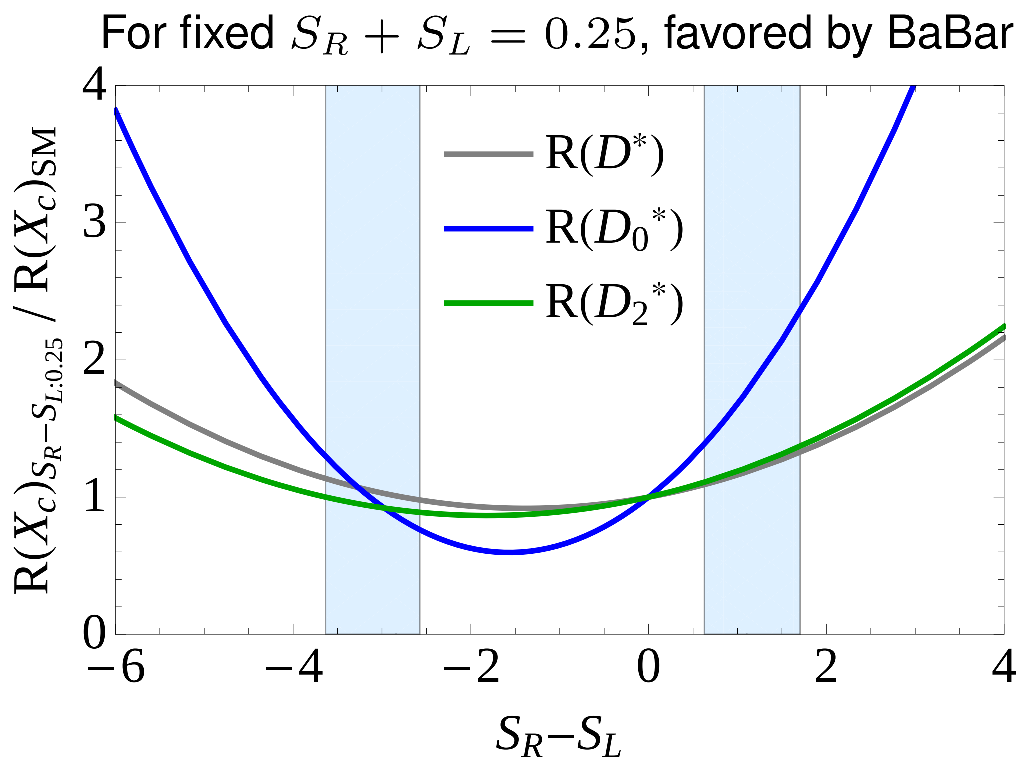
Complementary sensitivities to NP

- Complementary sensitivities

[Bernlochner & ZL, 1606.09300]





Different patterns in two blue bands

- 2HDM just for illustration — explore influence of all possible non-SM operators

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Final comments

Conclusions

- $B \to D^{(*)} \tau \bar{\nu}$: amusing if NP shows up in an operator w/o much SM suppression
- SM predictions can be systematically improved with more data
- There are good operator fits, and (somewhat) sensible MFV leptoquark models

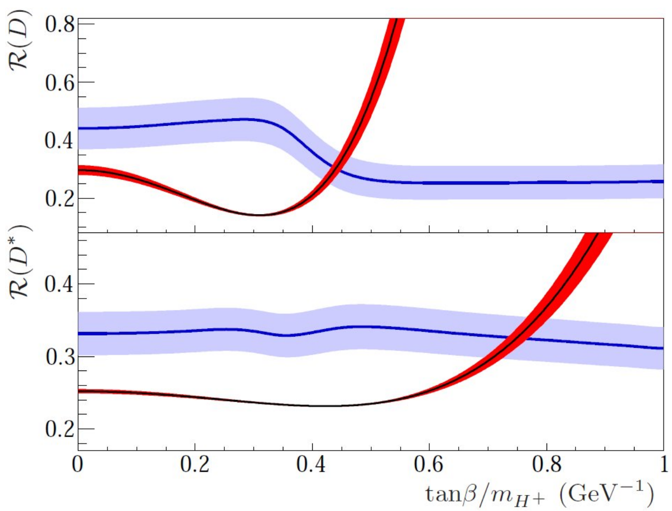
  (Fairly wild scenarios still viable)
- Measurements will improve in the next decade by nearly an order of magnitude

  (Even if central values change, plenty of room for significant deviations from SM)
- More theory progress to come, will impact measurements and sensitivity to BSM

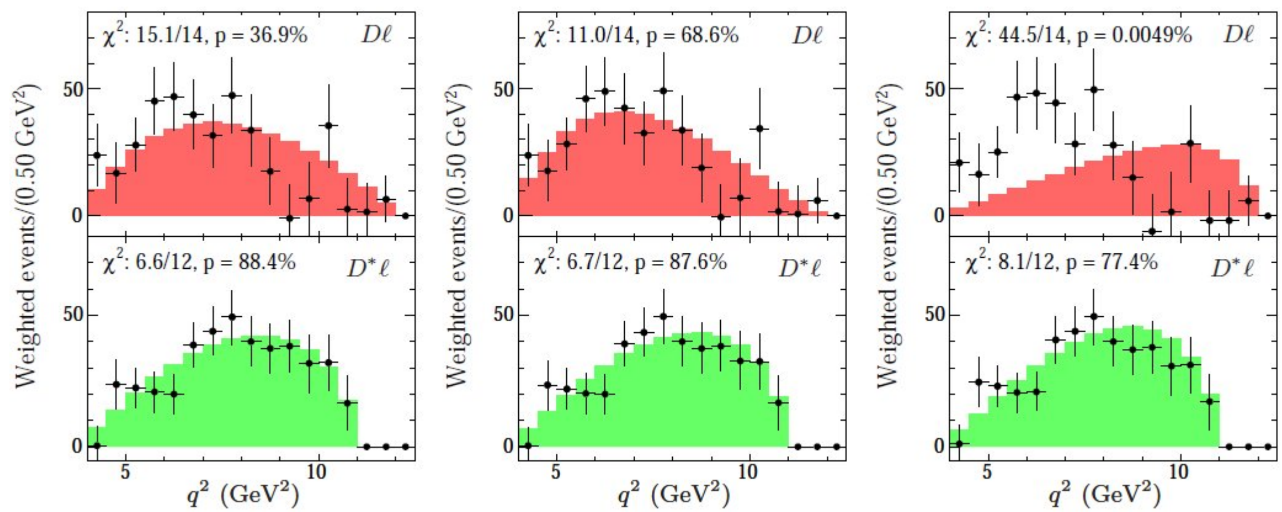
Bonus slides

BaBar statements from q^2 spectrum results

- BaBar studied consistency of rates with 2HDM, and $\mathrm{d}\Gamma/\mathrm{d}q^2$ with several models



[PRL 109 (2012) 101802, arXiv:1205.5442]



[PRD 88 (2013) 072012, arXiv:1303.0571]

- Found that type-II 2HDM gave nearly as bad fit to the data as the SM

- $\mathrm{d}\Gamma/\mathrm{d}q^2$ has additional discriminating power (no other distribution measured yet)

- No public info on bin-to-bin correlations, eyeball which solutions are (dis)favored

Survey of MFV model

- Scalars: Need $C_{S_L}/C_{S_R} \sim \mathcal{O}(1)$
  Hard to avoid y_c suppression or $\mathcal{O}(1)$ coupling to 1st generation
- Vectors: Rescaling the SM operator (O_{V_L}) gives good fit to the data
  Flavor singlet excluded by LHC, simplest charges don't work w/o assumptions
  If dynamics allows $W'\bar{Q}^3_L Q^3_L$, but not $W'\bar{Q}^i_L Q^i_L$, viable models exist; beyond MFV [Greljo, Isidori, Marzocca, 1506.0170]

- Leptoquarks: Viable MFV models exist

  Simplest choices — leptoquarks could be electroweak $SU(2)_L$ singlets or triplets:

$$\text{scalars: } S \sim (\bar{\mathbf{3}},\mathbf{1},\mathbf{1})\,, \quad (\mathbf{1},\bar{\mathbf{3}},\mathbf{1})\,, \quad (\mathbf{1},\mathbf{1},\bar{\mathbf{3}})$$

$$\text{vectors: } U_\mu \sim (\mathbf{3},\mathbf{1},\mathbf{1})\,, \quad (\mathbf{1},\mathbf{3},\mathbf{1})\,, \quad (\mathbf{1},\mathbf{1},\mathbf{3})$$

- Possibly viable: $S(\mathbf{1},\mathbf{1},\bar{\mathbf{3}})$ and $U_\mu(\mathbf{1},\mathbf{1},\mathbf{3}) \Rightarrow$ consider in more detail

  Both can be electroweak singlets or triplets

Excluding MFV scalars and vectors

- Scalars: Need comparable values of C_{S_L} and C_{S_R}

  If $H^\pm$ flavor singlet, $C_{S_L} \propto y_c$, so cannot fit $R(D^{(*)})$ keeping y_t perturbative

  If $H^\pm$ is charged under flavor (combination of Y-s, to couple to quarks & leptons), to generate $C_{S_L} \sim C_{S_R}$, some $\mathcal{O}(1)$ coupling to 1st generation quarks unavoidable
  Bounds on $4q$ or $2q2\ell$ operators exclude it

- Vectors: Rescaling the SM operator (O_{V_L}) gives good fit to the data

  Flavor singlet w/ W-like couplings: $m_{W'} \gtrsim 1.8\,\text{TeV} \iff 0.2 \sim g^2|V_{cb}|(1\,\text{TeV}/m_{W'})^2$

  Couplings to u, d suppressed for $(\bar{\mathbf{3}}, \mathbf{3}, \mathbf{1})$ and $(\bar{\mathbf{3}}, \mathbf{1}, \mathbf{3})$ under $U(3)_Q \times U(3)_u \times U(3)_d$
  $(\bar{\mathbf{3}}, \mathbf{3}, \mathbf{1})$: $b \to c$ transitions suppressed by y_c, too small
  $(\bar{\mathbf{3}}, \mathbf{1}, \mathbf{3})$: can fit data if $y_b = \mathcal{O}(1)$, but excluded by tree-level FCNC via W'^0

  (If dynamics allows $W'\bar{Q}_L^3 Q_L^3$, but not $W'\bar{Q}_L^i Q_L^i$, viable models exist; beyond MFV [Greljo, Isidori, Marzocca, 1506.0170])

MFV leptoquarks

- Assign charges under flavor sym.: [viable MFV LQs: Freytsis, ZL, Ruderman]

$$U(3)_Q \times U(3)_u \times U(3)_d$$

- Simplest choices — leptoquarks could be electroweak $SU(2)_L$ singlets or triplets:

scalars: $S \sim (\mathbf{\bar{3}},\mathbf{1},\mathbf{1})\,, \quad (\mathbf{1},\mathbf{\bar{3}},\mathbf{1})\,, \quad (\mathbf{1},\mathbf{1},\mathbf{\bar{3}})$

vectors: $U_\mu \sim (\mathbf{3},\mathbf{1},\mathbf{1})\,, \quad (\mathbf{1},\mathbf{3},\mathbf{1})\,, \quad (\mathbf{1},\mathbf{1},\mathbf{3})$

$S(\mathbf{\bar{3}},\mathbf{1},\mathbf{1})$ and $U_\mu(\mathbf{3},\mathbf{1},\mathbf{1})$ give large $pp \to \tau^+\tau^-$, excluded by Z' searches

$S(\mathbf{1},\mathbf{\bar{3}},\mathbf{1})$ and $U_\mu(\mathbf{1},\mathbf{3},\mathbf{1})$ give y_c suppressed $B \to D^{(*)}\tau\bar{\nu}$ contributions
$\Rightarrow$ too large couplings, or too light leptoquarks

- Possibly viable: $S(\mathbf{1},\mathbf{1},\mathbf{\bar{3}})$ and $U_\mu(\mathbf{1},\mathbf{1},\mathbf{3})$ $\Rightarrow$ consider in more detail

Both can be electroweak singlets or triplets

The $S(1,1,\bar{3})$ scalar LQ

- Interactions terms for electroweak singlet:

$$\mathcal{L} = S(\lambda Y_d^\dagger\, \bar{q}_L^c i\tau_2 \ell_L + \tilde{\lambda} Y_d^\dagger Y_u\, \bar{u}_R^c e_R)$$
$$= S_i(\lambda y_{d_i} V_{ji}^*\, \bar{u}_{Lj}^c e_L - \lambda y_{d_i} \bar{d}_{Li}^c \nu_L + \tilde{\lambda} y_{d_i} y_{u_j} V_{ji}^*\, \bar{u}_{Rj}^c e_R)$$

Integrating out S, contribution to $R(X_c)$ via: $(m_{S_3} \neq m_{S_1} = m_{S_2})$

$$-\frac{V_{cb}^*}{m_{S_3}^2}\Big(\lambda^2 y_b^2\, \mathcal{O}''_{S_R} + \lambda\tilde{\lambda} y_c y_b^2\, \mathcal{O}''_{S_L}\Big)$$

[electroweak triplet has no $\tilde{\lambda}$ term]

- Can fit $R(D^{(*)})$ data if $y_b = \mathcal{O}(1)$ Check $Z\tau^+\tau^-$ constraints, etc.

- Leptons: (i) τ alignment, charge LQ and 3rd gen. leptons opposite under $U(1)_\tau$
  (ii) lepton MFV, $(\mathbf{1},\bar{\mathbf{3}})$ under $U(3)_L \times U(3)_e$ [constraints differ]

- LHC Run 1 bounds on pair-produced LQ decaying to $t\tau$ or $b\nu$, $m_{S_3} \gtrsim 560\,\mathrm{GeV}$

Constraints from $b \to s\nu\bar{\nu}$

- With three Yukawa spurion insertions, one can write:

$$\delta\mathcal{L}' = \lambda' S Y_d^\dagger Y_u Y_u^\dagger \, \bar{q}_L^c i\tau_2 \ell_L$$

- Generates four-fermion operator:

$$\frac{V_{tb}^* V_{ts}}{2m_{S_3}^2} \, y_t^2 y_b^2 \, \lambda'\lambda \, (\bar{b}_L \gamma^\mu s_L \, \bar{\nu}_L \gamma_\mu \nu_L)$$

- Current limits on $B \to K\nu\bar{\nu}$ imply: $\lambda'/\lambda \lesssim 0.1$ — some suppression of λ' required

- Electroweak singlet vector LQ is the only one of the four models w/o this constraint

  (E.g., vector triplet has $\lambda' \, \bar{q}_L Y_u Y_u^\dagger Y_d \, \boldsymbol{\tau} \gamma_\mu \ell_L \boldsymbol{U}^\mu$ term)

- If central values & patterns change, more "mainstream" MFV models may fit

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Many signals, tests, consequences

- LHC: several extensions to current searches would be interesting
  - Extend $\tilde{t}$ and $\tilde{b}$ searches to higher prod. cross section
  - Search for $t \to b\tau\bar{\nu}$, $c\tau^+\tau^-$ nonresonant decays
  - Search for states on-shell in t-channel, but not in s-channel
  - Search for $t\tau$ resonances

- Low energy probes:
  - Firm up $B \to D^{(*)}\tau\bar{\nu}$ rate and kinematic distributions; Cross checks w/ inclusive
  - Smaller theor. error in $[\mathrm{d}\Gamma(B \to D^{(*)}\tau\bar{\nu})/\mathrm{d}q^2]/[\mathrm{d}\Gamma(B \to D^{(*)}l\bar{\nu})/\mathrm{d}q^2]$ at same q^2
  - Improve bounds on $\mathcal{B}(B \to K^{(*)}\nu\bar{\nu})$
  - $\mathcal{B}(D \to \pi\nu\bar{\nu}) \sim 10^{-5}$ possible, maybe BES III; enhanced $\mathcal{B}(D \to \mu^+\mu^-)$
  - $\mathcal{B}(B_s \to \tau^+\tau^-) \sim 10^{-3}$ possible

Not excluded?

- LQ pair production
- top decays
- t-channel non-resonant l^+l^- production
- LEP $Z \to l^+l^-$, HERA LQ production
- $c\bar{c}e^+e^-$ contact interaction / compositness
- $B - \overline{B}$ mixing, $K - \overline{K}$ mixing, $D - \overline{D}$ mixing
- $B \to X_s\nu\bar{\nu}$, $K \to \pi\nu\bar{\nu}$
- $D \to l^+l^-$ at tree level
- $B^- \to \mu\bar{\nu}$ at tree level
- $B_s \to \mu^+\mu^-$ and $K_L \to \mu^+\mu^-$ at one loop

- Strongest constraint from ϵ_K:

$$|\epsilon_K|_{\rm SM} = \frac{G_F^2\, m_W^2\, m_K f_K^2}{6\sqrt{2}\,\pi^2\,\Delta m_K}\,\hat{B}_K\,\kappa_\epsilon |V_{cb}|^2\,\lambda^2\,\bar{\eta}\Big[|V_{cb}|^2(1-\bar{\rho})\eta_{tt}S_0(x_t) + \eta_{ct}S_0(x_c,x_t) - \eta_{cc}x_c\Big]$$

$|\epsilon_K|_{\rm exp} = (2.23 \pm 0.01) \times 10^{-3}$ vs. $|\epsilon_K|_{\rm SM} = (1.81 \pm 0.28) \times 10^{-3}$ [Brod & Gorbahn, 2011]

  – Uncertainties big enough to allow for $5-10\%$ enhancement of $|V_{cb}|$

  – The $R(D^{(*)})$ excess may shrink and be significant; can also make cocktails...

- Even an enhancement much smaller than today can become 5σ in the future

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