

Probing Velocity Dependent Self Interacting Dark Matter

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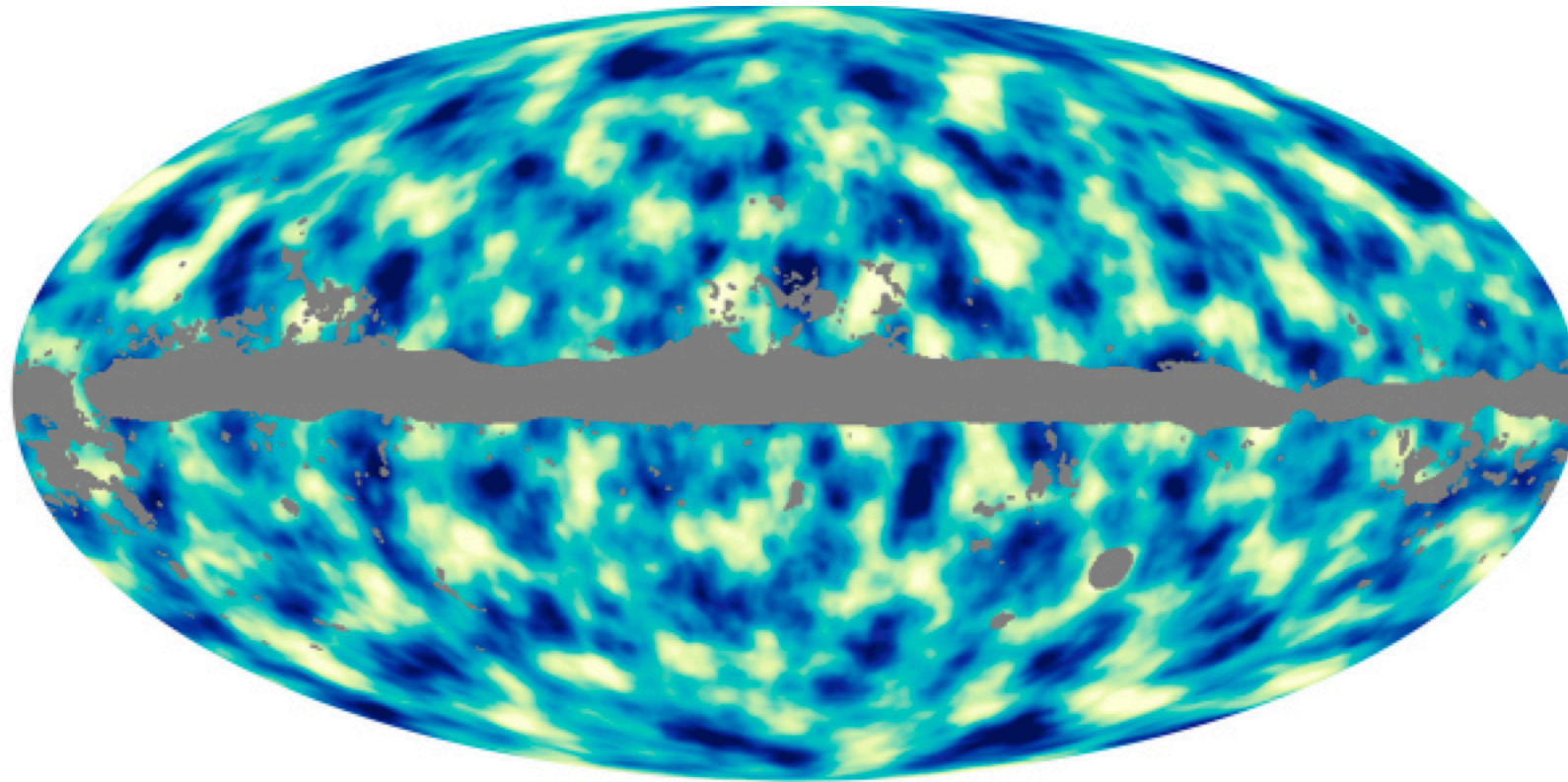
Outline

1. CDM Small Scale Potential Problems
 - DM Self-Interaction (SIDM) as possible solution
2. Current constraints on SIDM
3. Velocity Dependent SIDM (vdSIDM)
4. Enhanced ν flux from DM annihilation

Indirect Probes (ν) on vdSIDM models

Collisionless CDM

Extremely successful at large scales



Date: 02 April 2013

Satellite: Planck

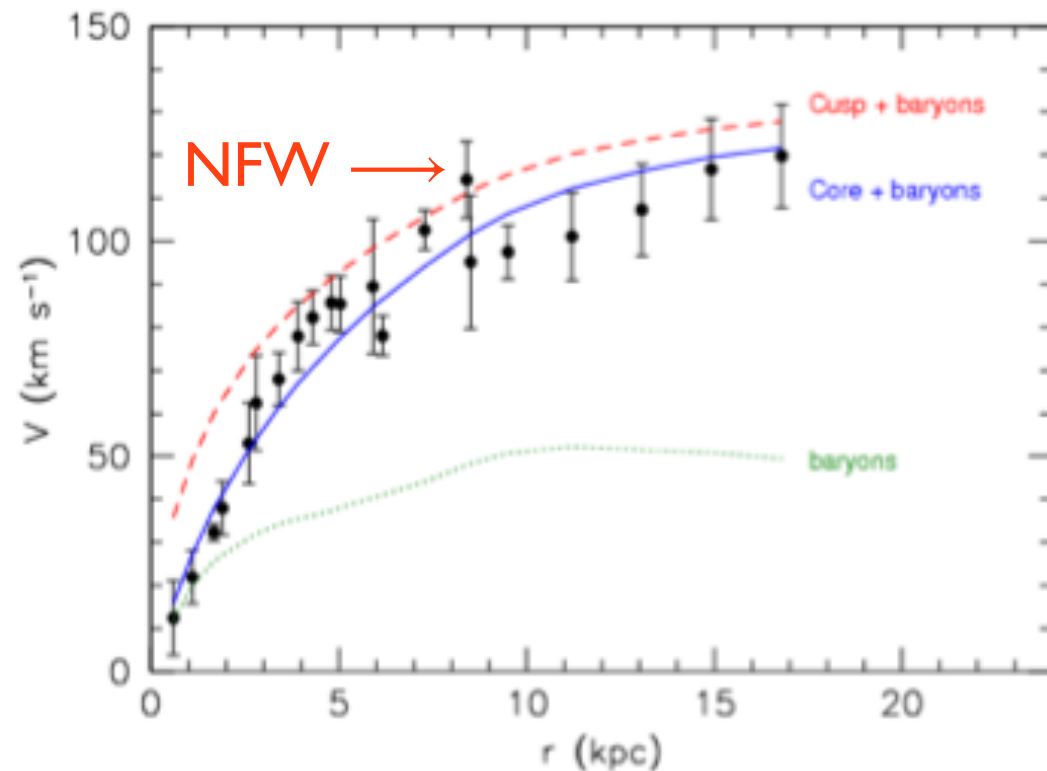
Depicts: All-sky map of dark matter distribution in the Universe

Copyright: ESA and the Planck Collaboration

CDM simulations fit very well large scale
observations

CDM Potential Problems

at small scale structure formation



Core / Cusp

CDM: too much DM ~ few Kpc

Majority of gal rot curves: better fit by
cored profile

(Weinberg et al., arXiv:1306.0913)

Data: F568-3 (SSDS)



Too Big to Fail

9 "classic" massive SIM DM subhalos

Boylan-Kolchin et al. (MNRAS 415, 2011)

(Weinberg et al., arXiv:1306.0913)

CDM simulations predict too much mass
in halos and subhalos central regions

Self Interacting Dark Matter

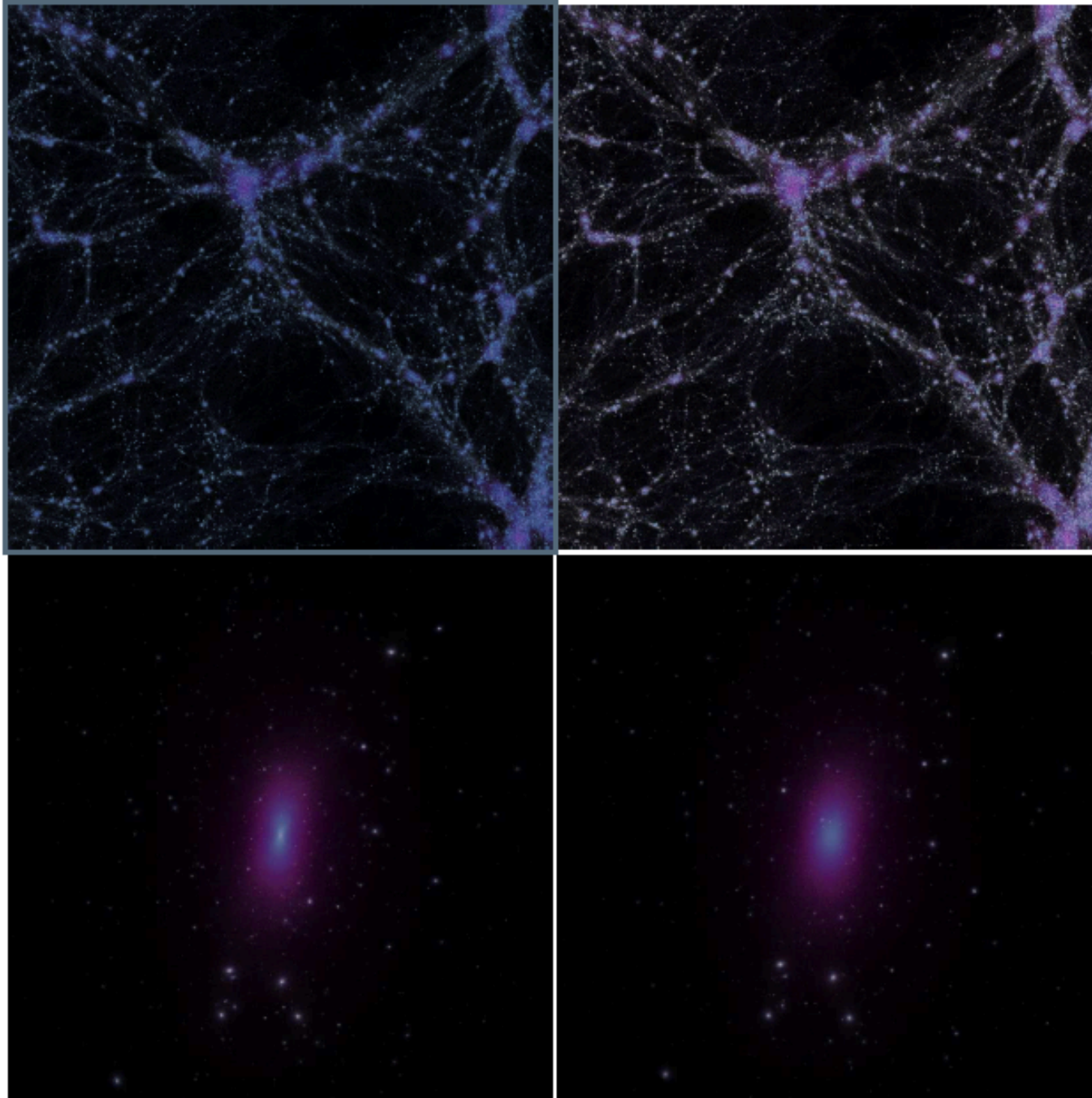
SIDM solves Small Scale Potential Problems

(Spergel and P. Steinhardt, PRL **84**, 2000)

DM scatters before reaching center of galaxy

$$\begin{aligned}\frac{\sigma_{\chi\chi}}{m_{\chi}} &= 8 \times 10^{-(25-22)} \text{ cm}^2/\text{GeV} \\ &= 4.5 - 450 \text{ cm}^2/\text{g}\end{aligned}$$

SIDM Simulations



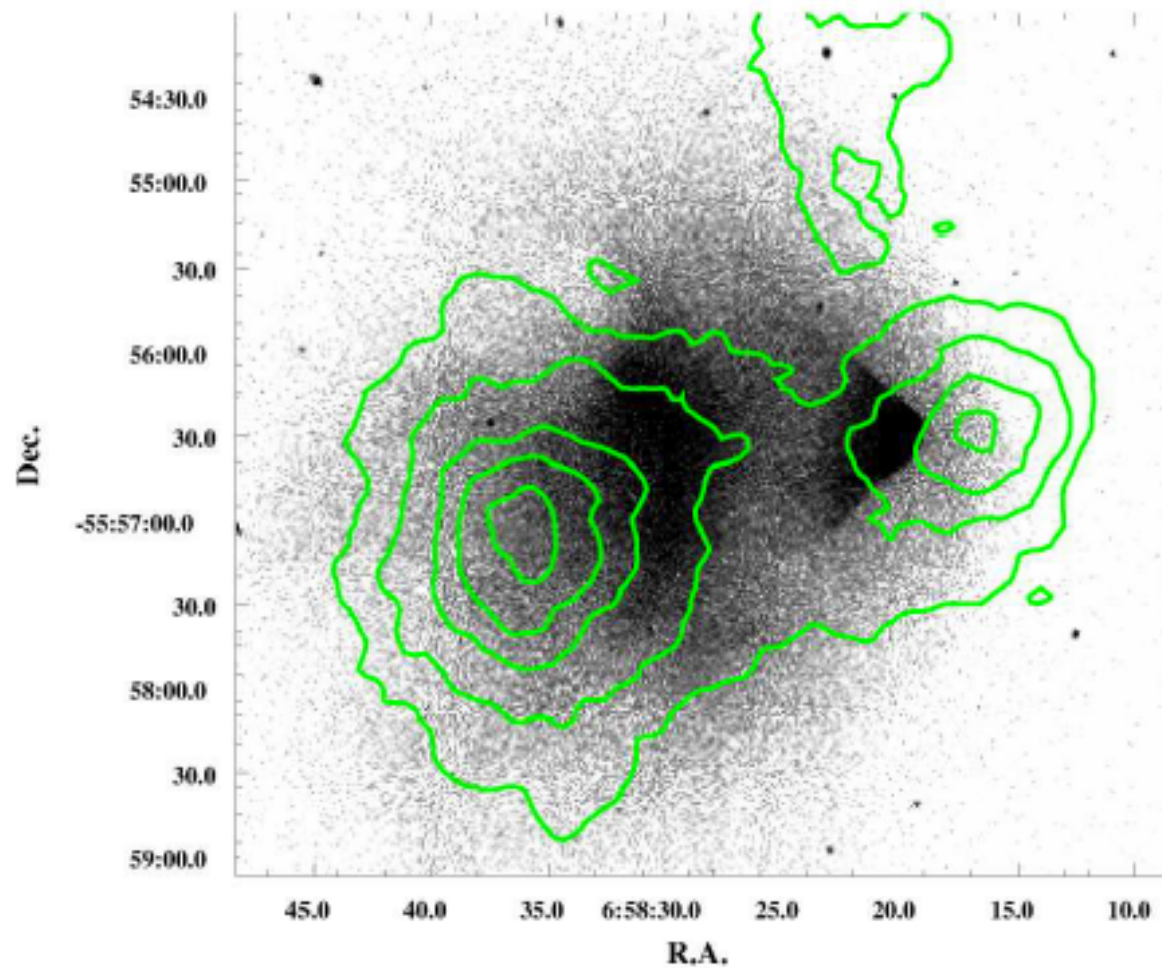
CDM and SIDM simulations

(M. Rocha et al., MNRAS 430, 2013)

- SIDM numerical simulation
 - constant density cores: much reduced central density
 - subhalo content is modestly reduced

$$\frac{\sigma_{\chi\chi}}{m_{\chi}} \simeq 1 \text{ cm}^2/\text{g} \Rightarrow \text{central density is TOO LOW}$$
$$\simeq 0.1 \text{ cm}^2/\text{g} \Rightarrow \text{consistent}$$

SIDM Constraints



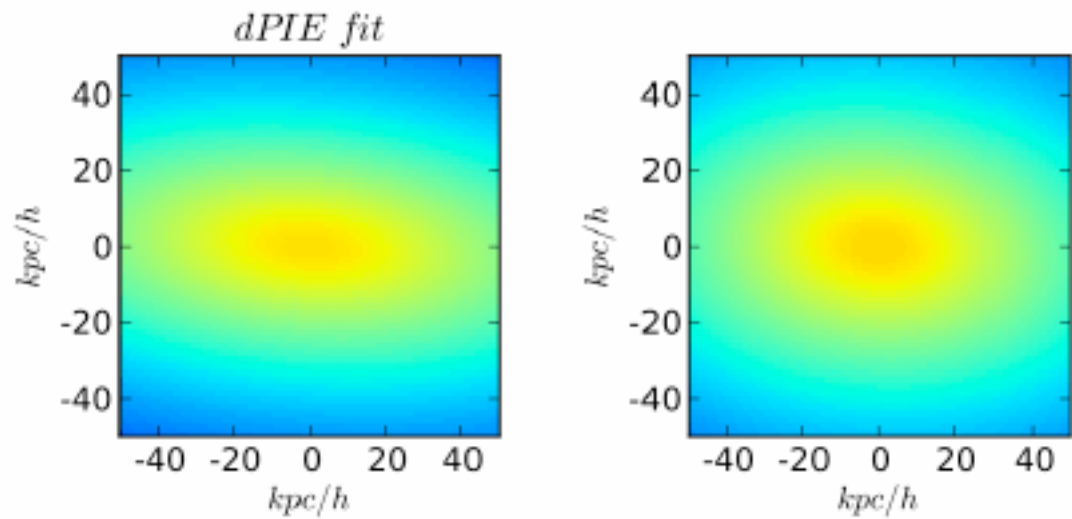
Bullet Cluster

(S. Randall et al., ApJ 679, 2008)

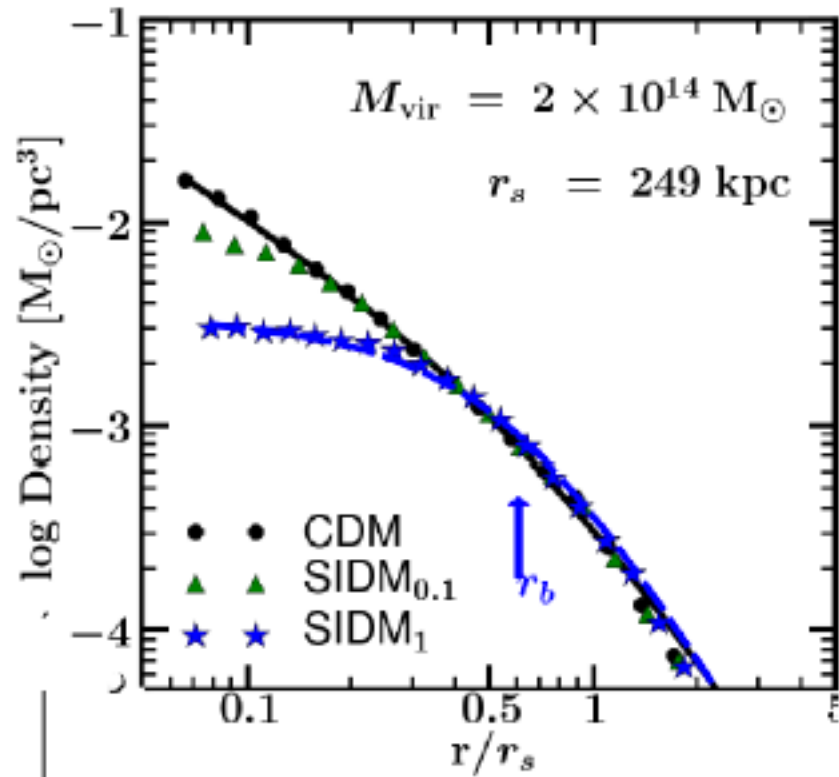
Xray image and lensing contour

$$\frac{\sigma_{\chi\chi}}{m_{\chi}} < 1.25 \text{ cm}^2/\text{g}$$

Constraints from Simulations



(A. Peter et al., arXiv:1208.3026)
CMB and SIDM 0.1



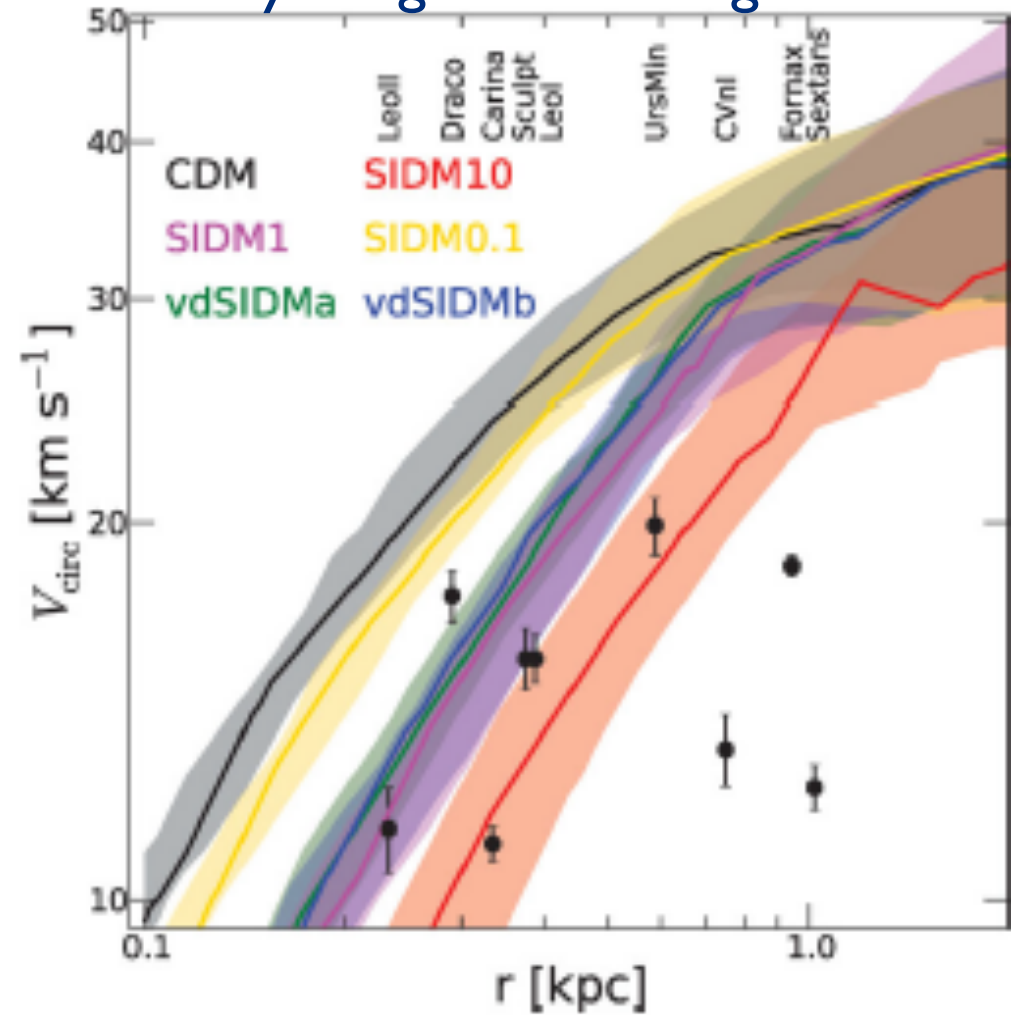
DM halos surface densities

(A. Peter et al., MNRAS 430, 2013)

Milky Way Dwarfs Kinematics

(Zavala et al., MNRAS 431, 2013)

X-ray image and lensing contour

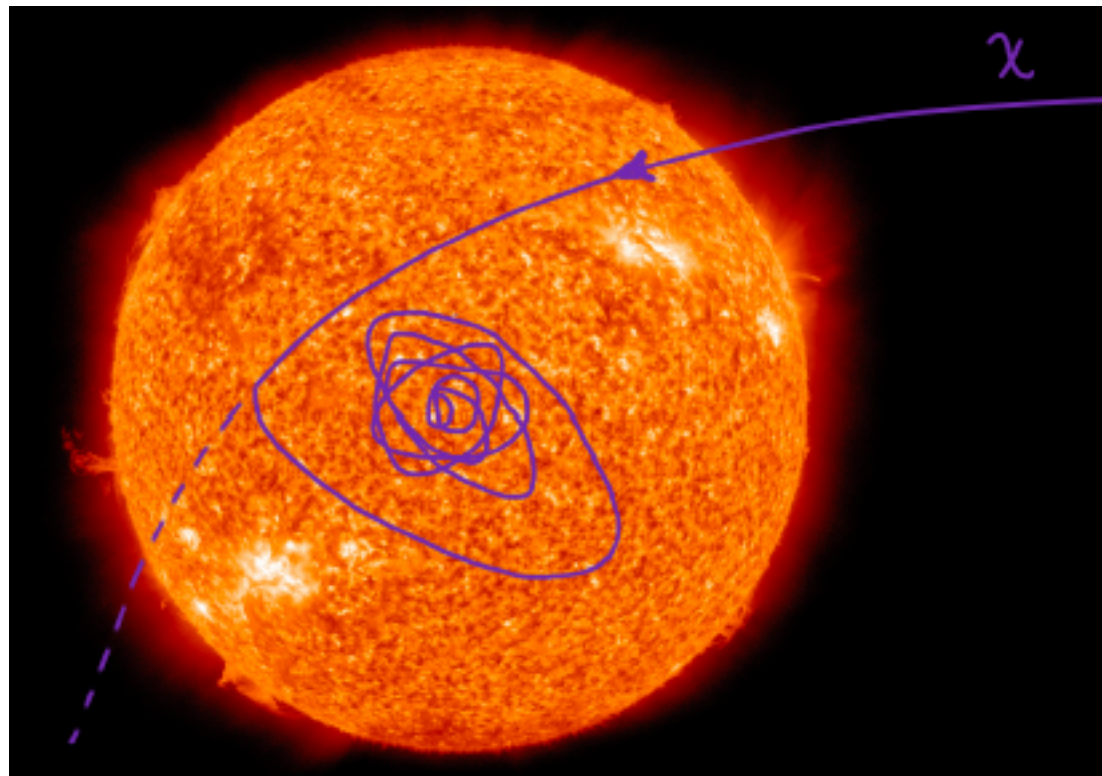


SIDM will be effective if

$$0.1 < \frac{\sigma_{\chi\chi}}{m_{\chi}} < 1 \text{ cm}^2/\text{g}$$

Probing SIDM with neutrinos

Self-Interaction enhances DM capture in the Sun



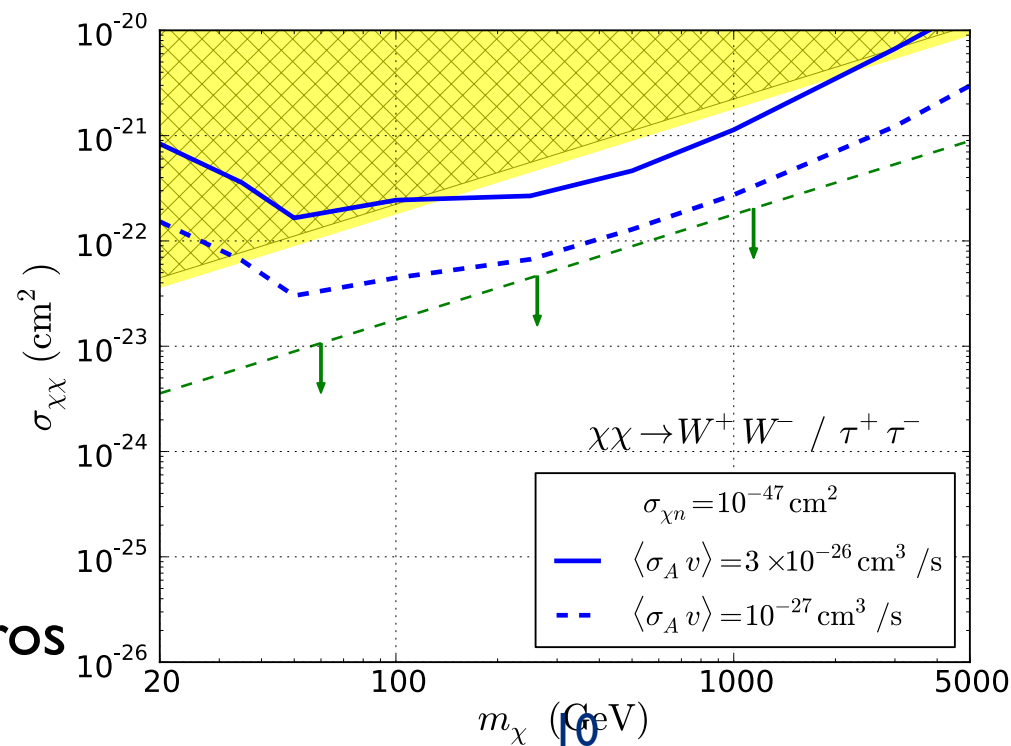
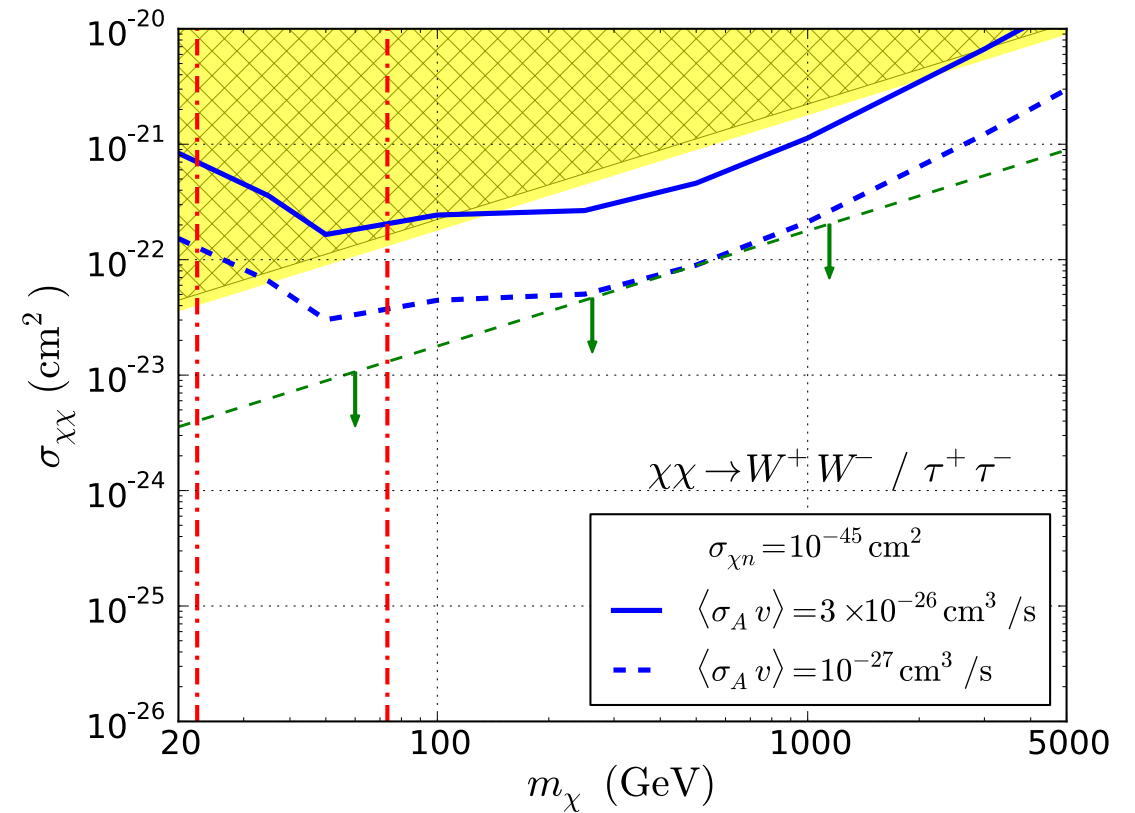
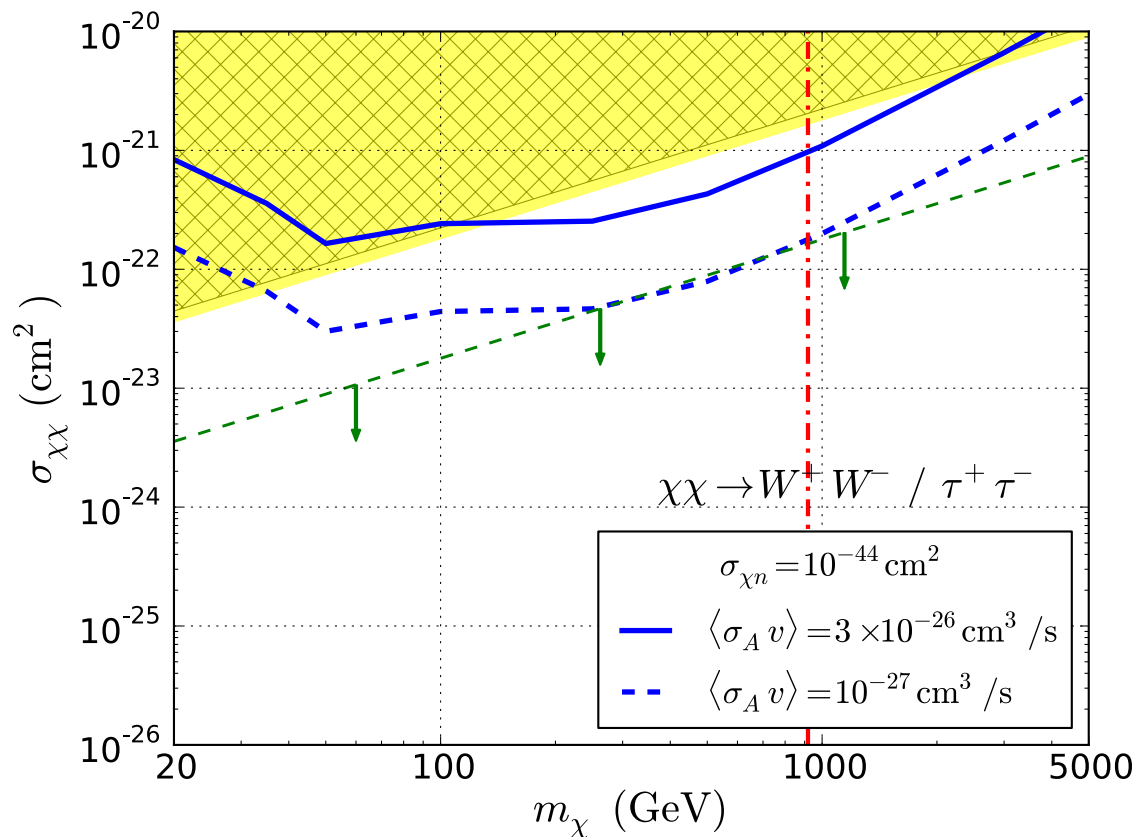
ν flux from DM annihilation will also be enhanced

Independently probe SI interesting $\sigma_{\chi\chi}/m_{\chi\chi}$ region





1. determine enhanced ν flux (simulation)
2. compare predictions with IceCube results

Probing SIDM models

$W^+W^- / \tau^+\tau^-$ - Spin Independent



Region above
blue curve:
excluded at
90% CL

-  Bullet Cluster
-  Halo Shapes
-  SIDM too low
-  LUX (to the left or between lines)

SIDM Constraints from IceCube

- SIDM is severely constrained if annihilates into WW

$$\frac{\sigma_{\chi\chi}}{m_{\chi}} < 0.6 \text{ cm}^2/\text{g} \quad \text{if} \quad \langle\sigma v\rangle = 3 \times 10^{-26} \text{ cm}^3/\text{s}$$

$$\frac{\sigma_{\chi\chi}}{m_{\chi}} < 0.1 \text{ cm}^2/\text{g} \quad \text{if} \quad \langle\sigma v\rangle = 1 \times 10^{-27} \text{ cm}^3/\text{s}$$

most SIDM effective models are ruled out

- $b\bar{b}$ analysis independently confirms bullet cluster results

SIDM can solve CDM potential small scale problems if:
Annihilation produces lower energy neutrinos
→ Self-scattering is velocity dependent

vdSIDM Enhances DM Capture and Annihilation

$$\dot{N} = \Gamma_C + \Gamma_{\chi\chi} - \Gamma_A$$

$$\Gamma_{\chi\chi} \propto n_{\chi}^{\text{halo}} N_{\chi} \sigma_{\chi\chi}(\mathbf{v}_{\text{rel}}) v_{\text{rel}}$$

$\sigma_{\chi\chi}(\mathbf{v}_{\text{rel}}) \rightarrow$ possible Sommerfeld enhanced

$$\Gamma_A \rightarrow \parallel$$

Equilibrium among capture and annihilation rates

\Rightarrow maximum annihilation rate

vdSIDM Model

Tulin, Yu & Zurek

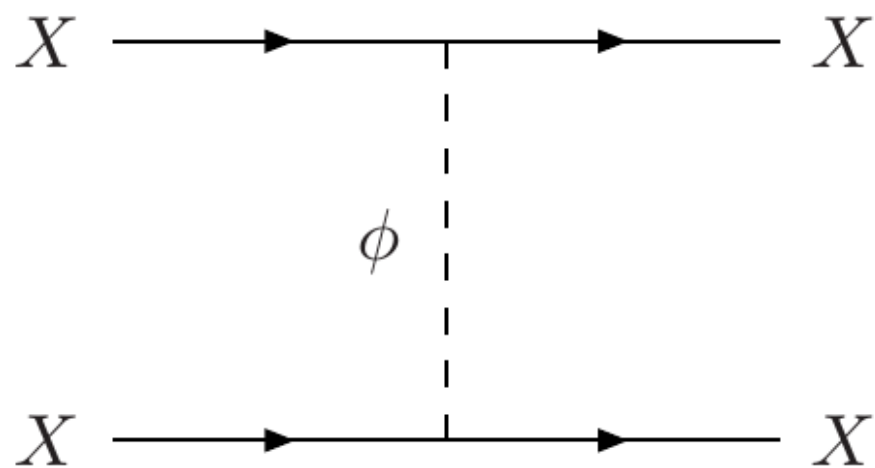
PRD 87 (2013)

DM Elastic Scattering

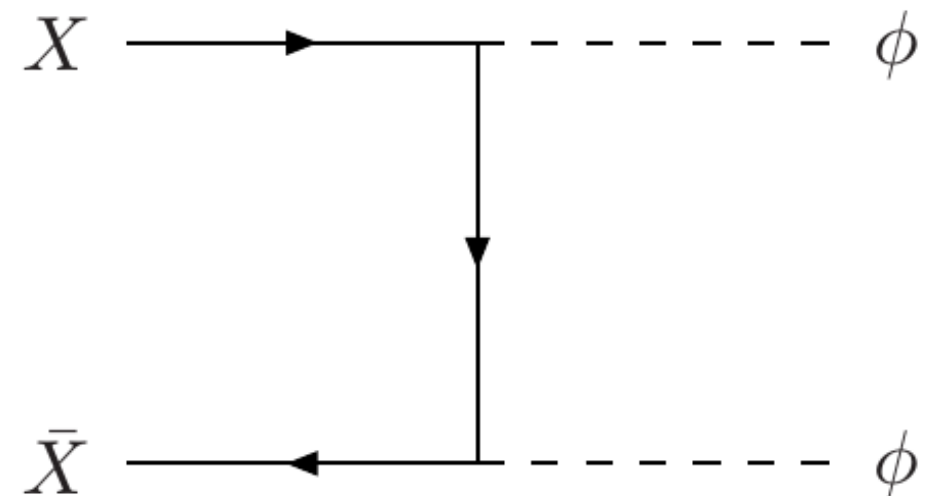
$$\mathcal{L}_{\text{int}} = g_{\chi} \bar{\chi} \gamma^{\mu} \chi \phi_{\mu}$$

$\chi \equiv$ DM Fermion

$\phi \equiv$ vector mediator



Self Interaction



DM Annihilation

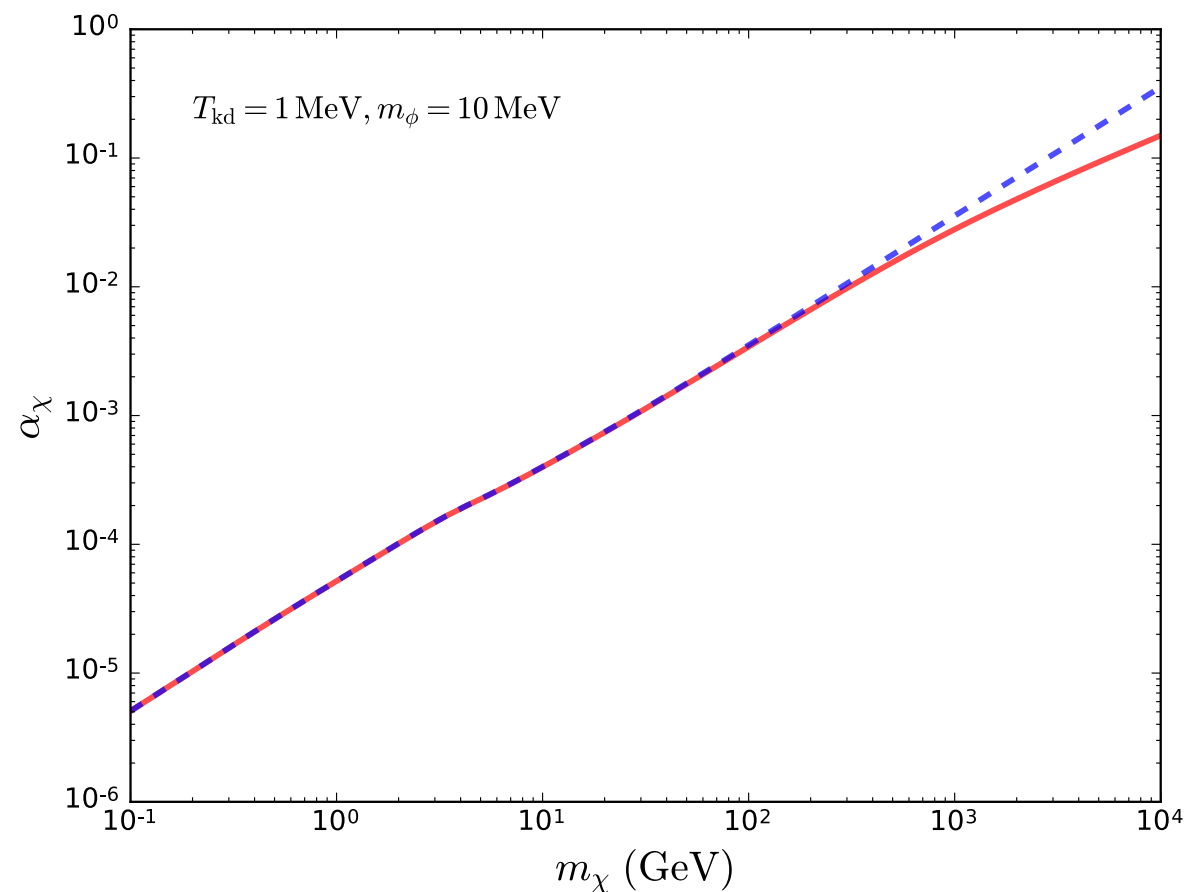
vdSIDM Model

Parameters: $\alpha_\chi = \frac{g_\chi}{4\pi}$, m_χ , m_ϕ

$m_\phi \sim 1 - 100 \text{ MeV}$

to solve small scale problems

(Tulin, Yu, Zurek - PRD 87 (2013))



Assuming Ω_{DM} is set by thermal freeze-out

Coupling to SM

Kaplinghat, Tulin & Yu

PRD 89 (2014)

ϕ mediator couples to SM through γ or Z mixing

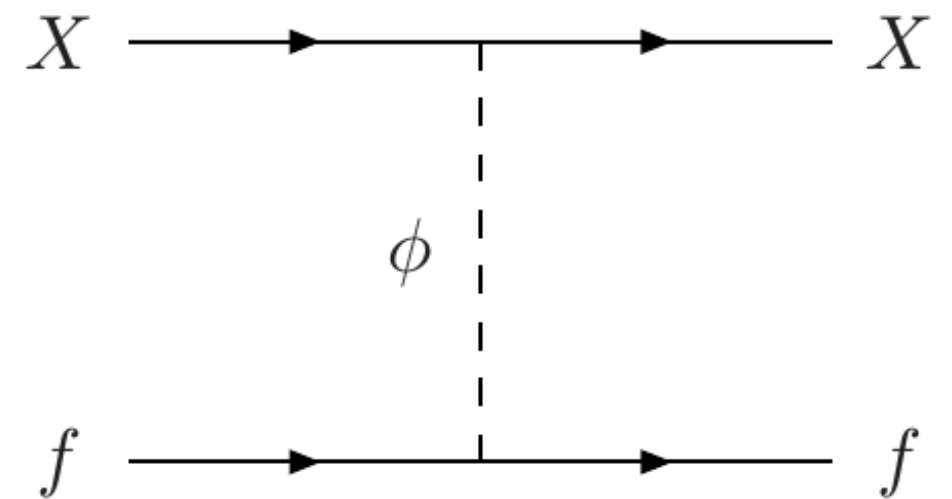
$$\mathcal{L}_{\text{mix}} = \frac{\epsilon_\gamma}{2} \phi_{\nu\mu} \mathbf{F}^{\mu\nu} + m_Z^2 \epsilon_Z \phi_\mu \mathbf{Z}^\mu$$

$$\mathcal{L}_{\text{int}} = e \phi_\mu (\epsilon_p \bar{\mathbf{p}} \gamma^\mu \mathbf{p} + \epsilon_n \bar{\mathbf{n}} \gamma^\mu \mathbf{n})$$

$$\epsilon_p = \epsilon_\gamma + 0.05 \epsilon_Z$$

$$\epsilon_n = -0.6 \epsilon_Z$$

$$\epsilon_\gamma \text{ and } \epsilon_Z \ll 1$$



DM - nucleon scattering

Capture in the Sun

Scattering with Sun's Nuclei

$$\Gamma_C \propto n_\chi n_N \sigma_{\chi N}$$

$$\sigma_{\chi N}(q^2 = 0) = 16\pi\alpha_{\text{em}}\alpha_\chi \frac{\mu_{\chi N}^2}{m_\phi^2} \left[\epsilon_p Z + \epsilon_n (A - Z)^2 \right]$$

$m_\phi \sim 1 - 100 \text{ MeV}$ is about same order as momentum transfer

suppression factor:

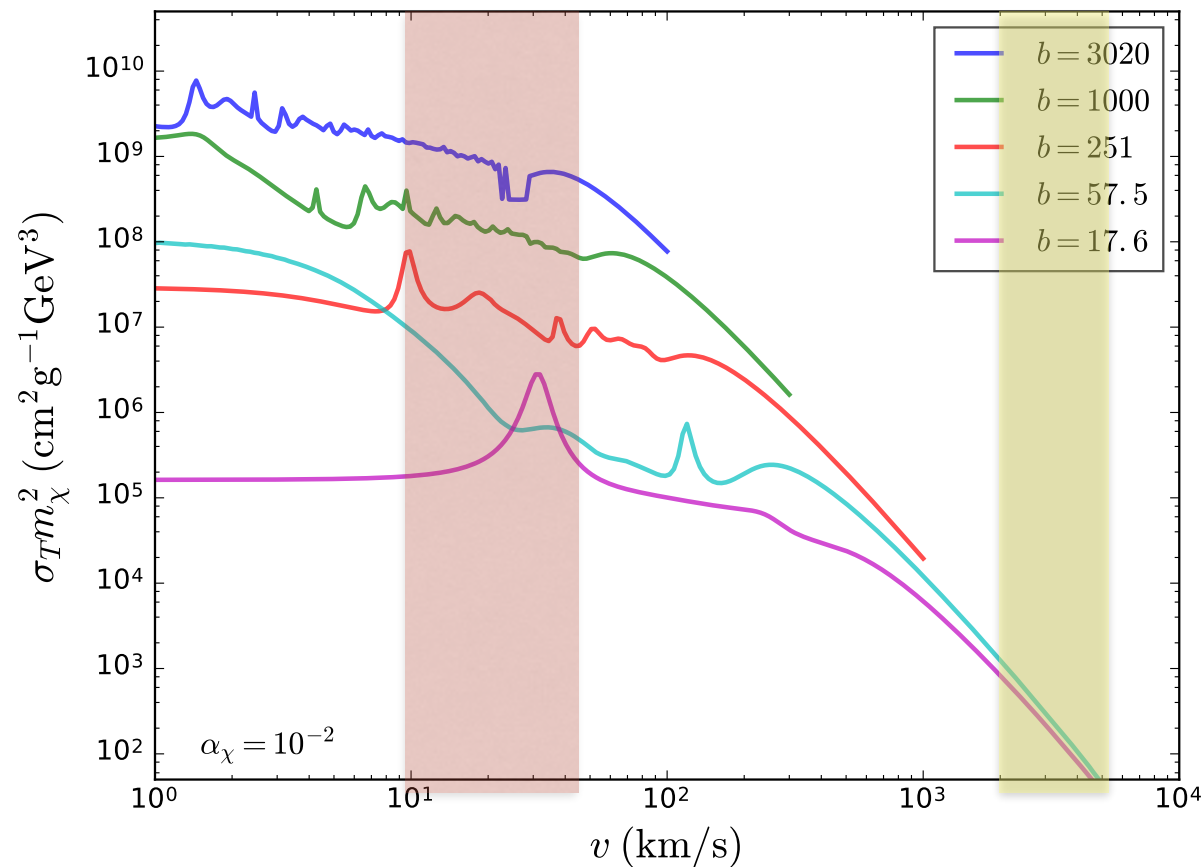
$$\sigma_{\chi N} = \sigma_{\chi N}(q^2 = 0) \times \frac{m_\phi^4}{(m_\phi^2 + q^2)^2}$$

Capture in the Sun

DM Self Scattering

non relativistic limit \Rightarrow Yukawa potential

$$\mathbf{V}(\mathbf{r}) = \pm \frac{\alpha_\chi}{r} \exp(-m_\phi r) \quad - \rightarrow \text{attractive } (\chi\bar{\chi})$$
$$+ \rightarrow \text{repulsive } (\chi\chi \text{ or } \bar{\chi}\bar{\chi})$$



$$\mathbf{b} = \alpha_\chi m_\chi / m_\phi$$

Self Interacting σ :
partial wave method

Tulin, Yu & Zurek
PRD 87 (2013)

Sommerfeld effect does not
play significant role

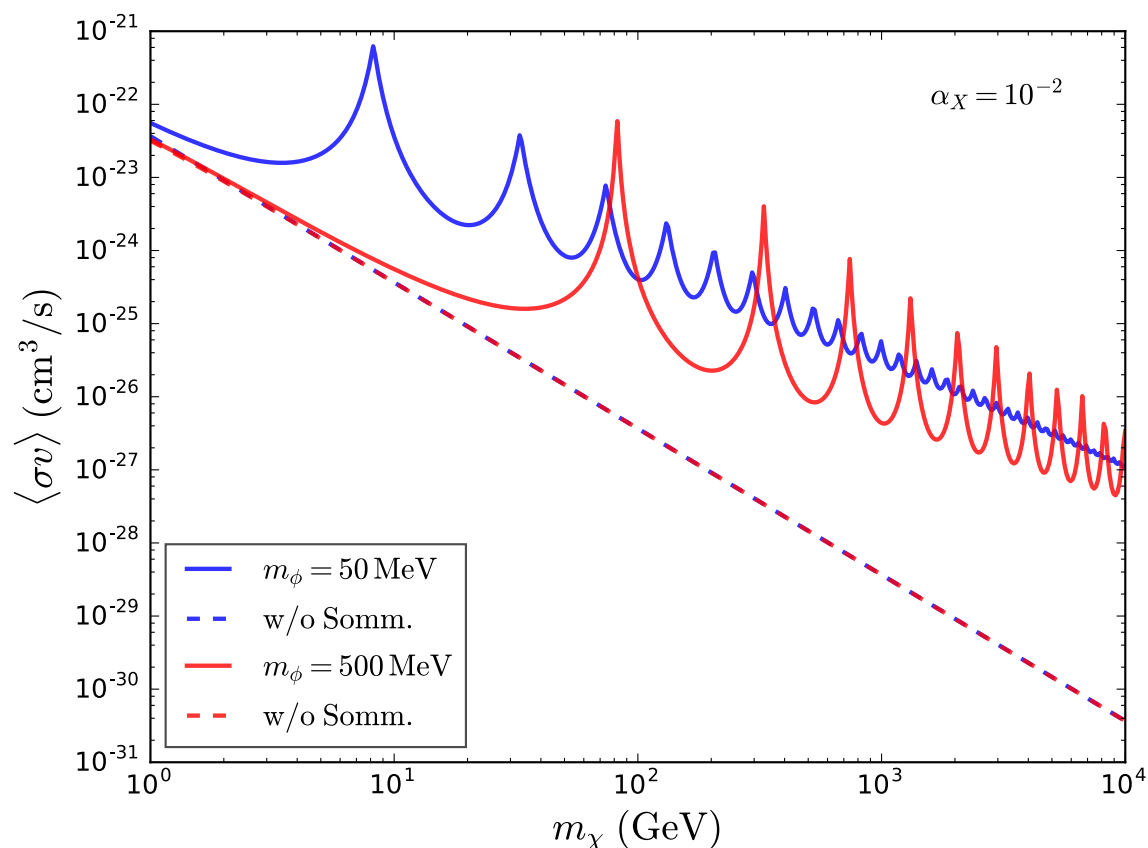
Annihilation in the Sun

$$\Gamma_A = \frac{1}{2} \langle \sigma_A \mathbf{v} \rangle V_{\text{eff}}$$

$$\langle \sigma_A \mathbf{v} \rangle = \frac{1}{2} \left(\frac{m_\chi}{\pi T_\chi} \right)^{3/2} \int \mathbf{S} (\sigma_{\text{a}\mathbf{v}})^{\text{tree}} \mathbf{v}^2 e^{-\frac{m_\chi \mathbf{v}^2}{4T_\chi}} d\mathbf{v}$$

$\mathbf{S} \equiv$ Sommerfeld factor

Annihilation XS



Annihilation Channel:

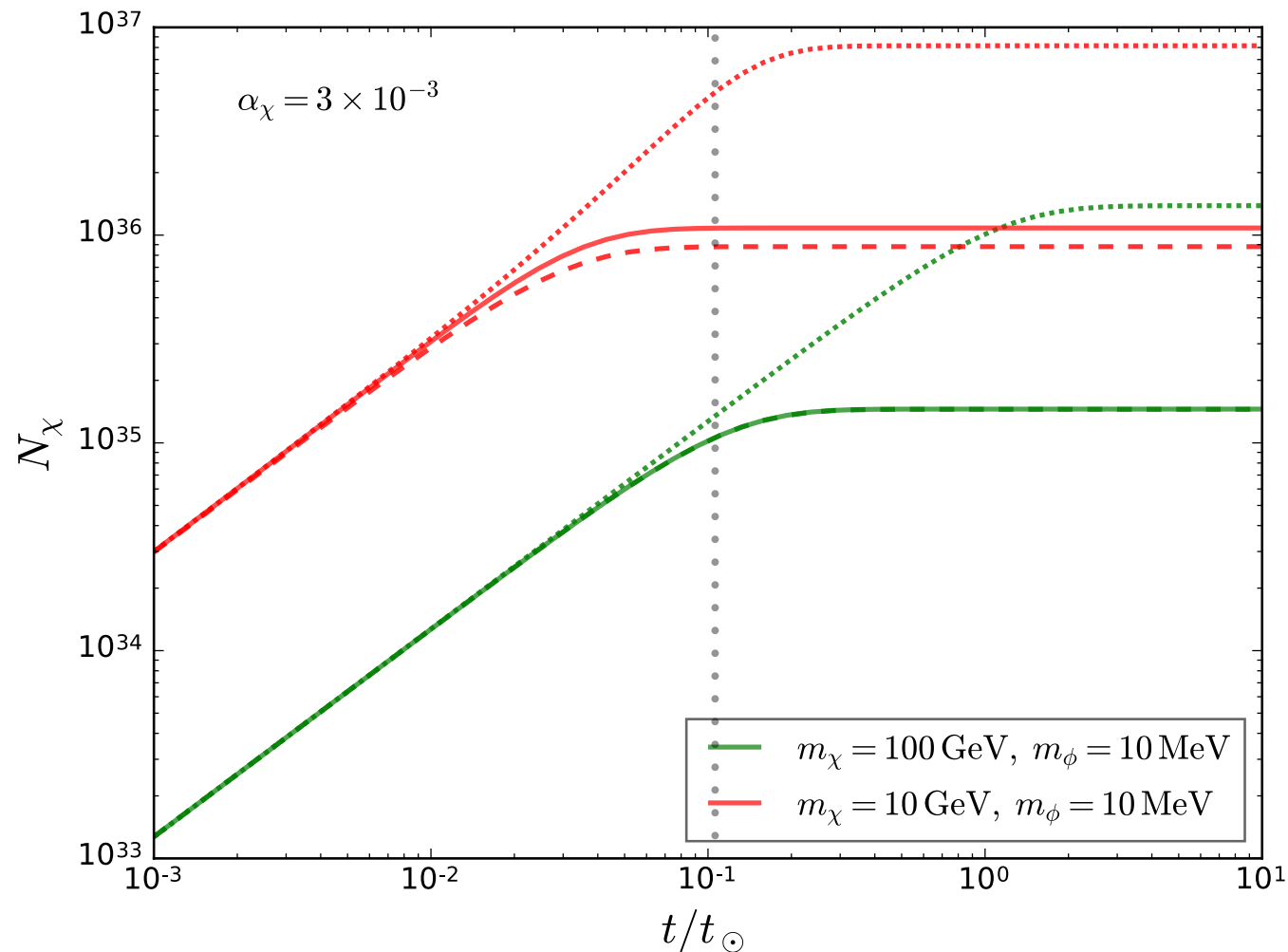
$$\chi \bar{\chi} \rightarrow \phi \phi \rightarrow 2\nu_1 2\bar{\nu}_1$$

Branching Ratio:

$$\text{BR}(\epsilon_\gamma = \epsilon_Z) = \frac{6}{7}$$

$$\text{BR}(\epsilon_\gamma = 0, \epsilon_Z) = 1$$

N_χ Time Evolution in the Sun



Enhancements on Γ_A :

- Self Capture
- Sommerfeld effect on annihilation

Suppression on Γ_A :

- momentum transfer suppression $\sigma_{\chi N}$

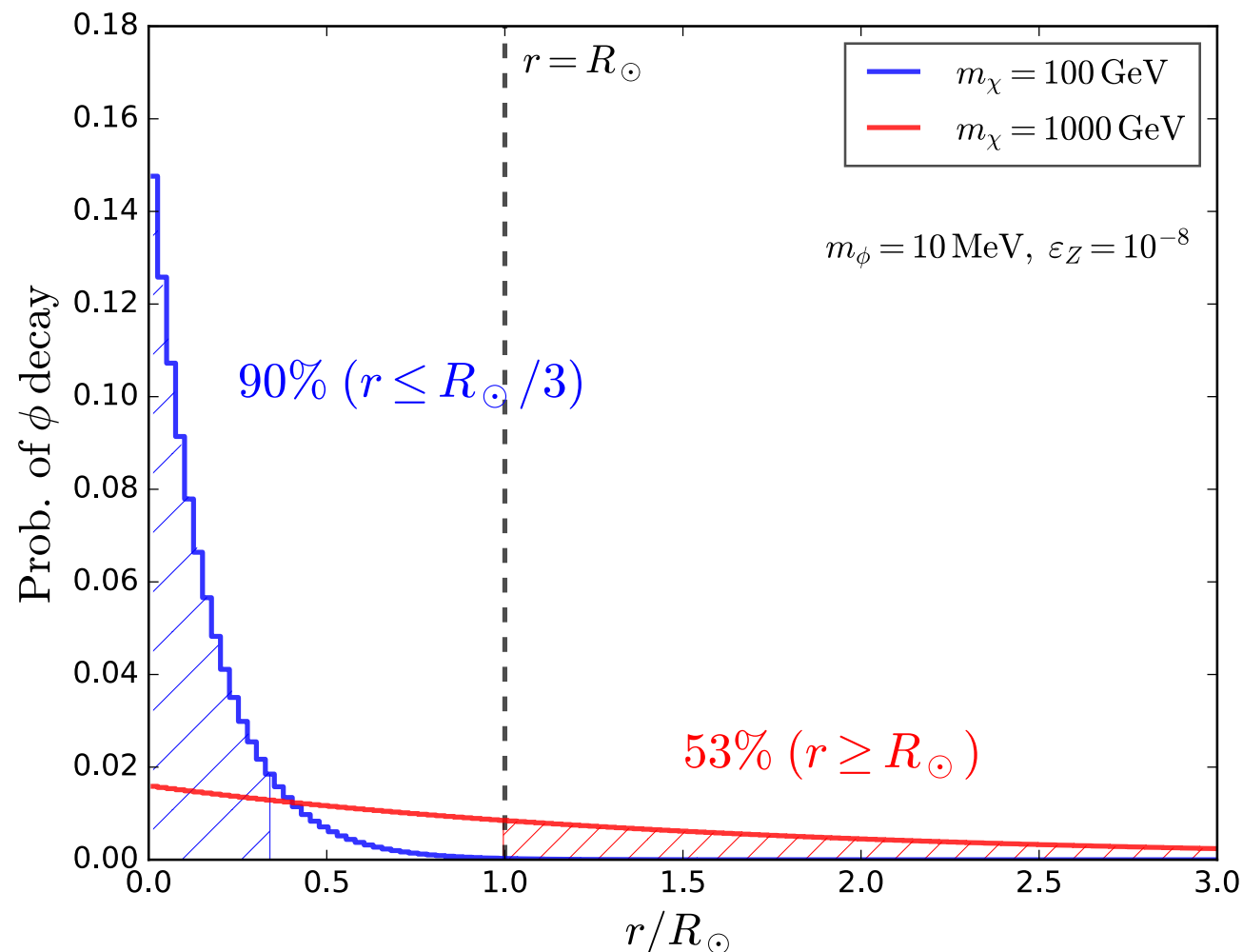
- - - Sommerfeld (NO Self Capture)

— Sommerfeld + Self Capture

..... Only Self Capture

ν Production and Propagation

ϕ lifetime is important

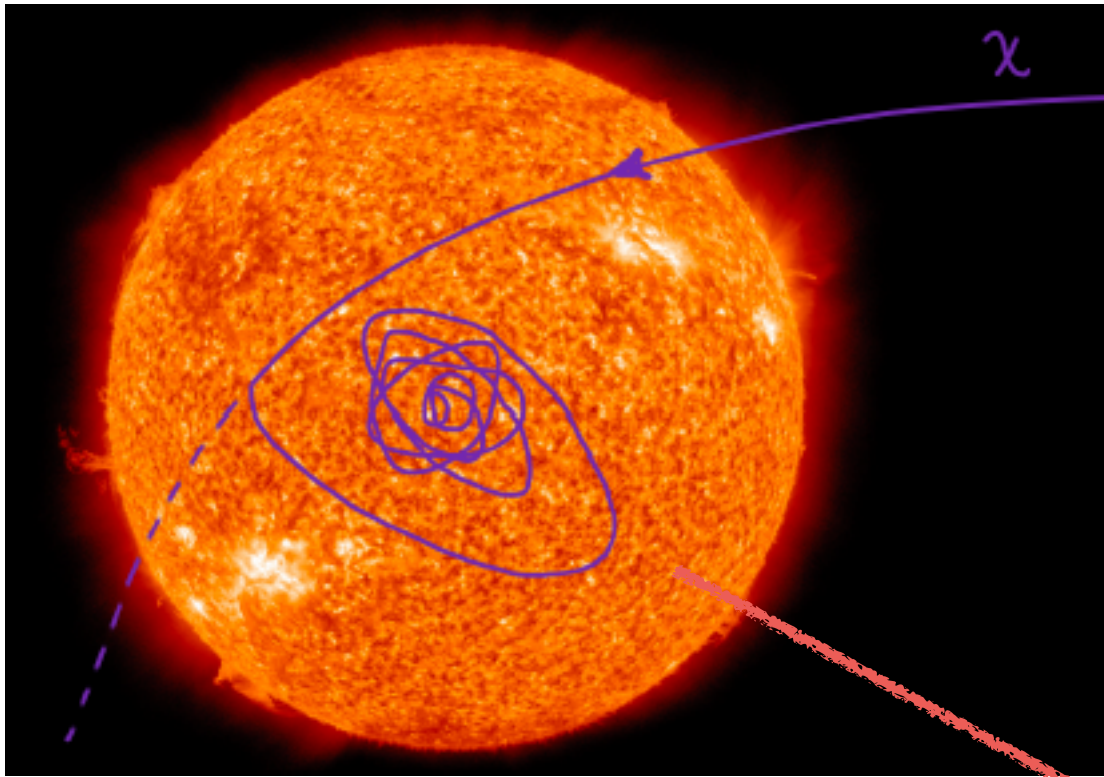


Standard Wimp: ν production in Sun's core

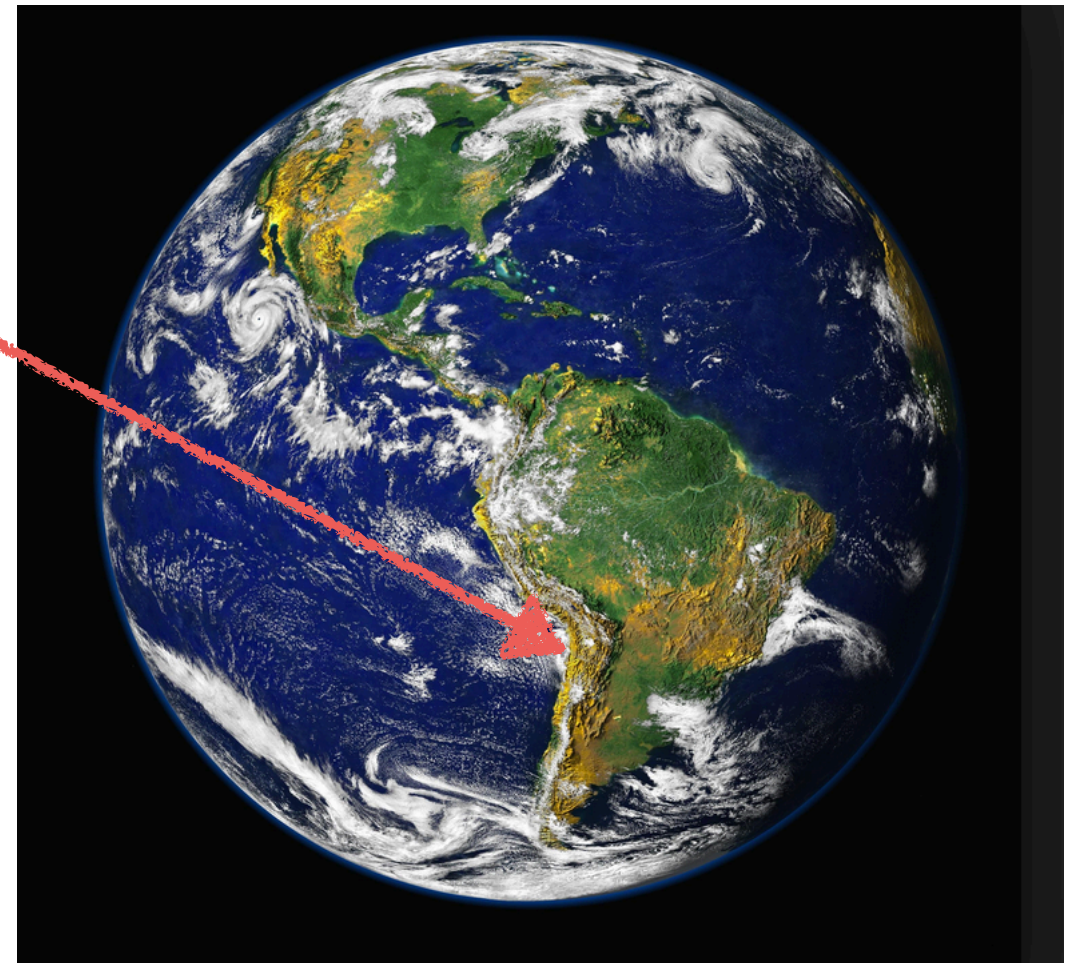
vdSIDMS: ν production beyond the core

$$\tau_\phi = 1 \text{ s} \left(\frac{10^{-10}}{\epsilon_Z} \right)^2 \left(\frac{m_\phi}{10 \text{ MeV}} \right)$$

Further enhances expected neutrino signal for some values of parameter space



Results: arXiv:1711.02052
IA, D.S.Robertson



Conclusions

- SIDM is severely constrained if annihilates into WW
- SIDM annihilation into $b\bar{b}$ confirms bullet cluster results

I.A, C. P. de Los Heros & Denis S. Robertson JCAP **02**, 2014

- if $\sigma_{\chi\chi}$ is velocity dependent, there will be enhancements on annihilation rate and a significant enhancement on the neutrino production

stay tuned:
arXiv:1706.XXXXX
IA, D.S.Robertson

