KJE @ 2nd CERN-ECFA-NuPECC Workshop on the LHeC Divonne-les-Bains, 1-3 Sep, 2009

Global NLO analysis of nPDFs: EPS09 + reanalysis for the LHeC

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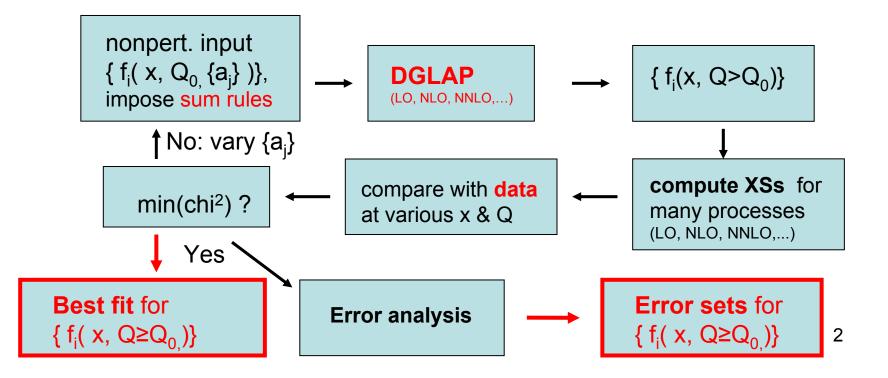
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- EPS09 JHEP 0904:065,2009, arXiv:0902.4154 [hep-ph] https://www.jyu.fi/fysiikka/en/research/highenergy/urhic/nPDFs
- renanalysis with LHeC pseudodata included

Global analysis of PDFs & nPDFs

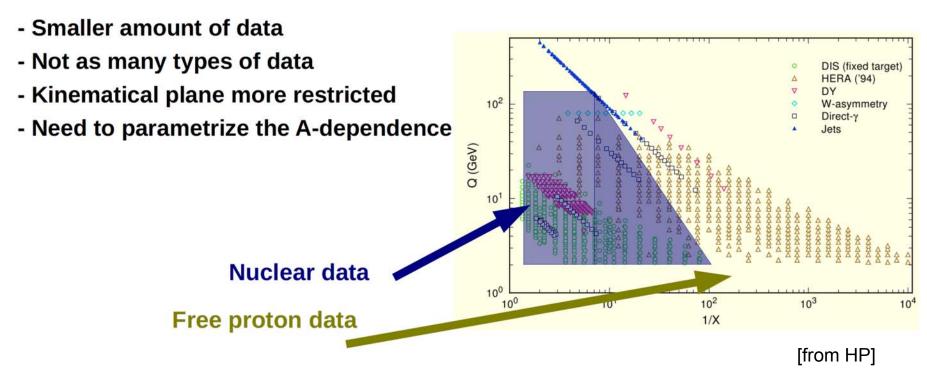
- test of pQCD & factorization $\sigma_{AB \to h+X} = \sum_{i,j} f_i^A(Q^2) \otimes \hat{\sigma}_{ij \to h+X} \otimes f_j^B(Q^2) + \mathcal{O}(Q^2)^{-n}$
- procedure to find a universal set $\{f_i(x,Q^2)\}$ at all x and $Q^2 >> \Lambda^2_{QCD}$
 - 1. Iterate until best fit, best set of initial parameters {a_i} found
 - 2. Perform error analysis for quantifying the PDF uncertainties



Challenges in extracting the nonperturbative input for PDFs and nPDFs

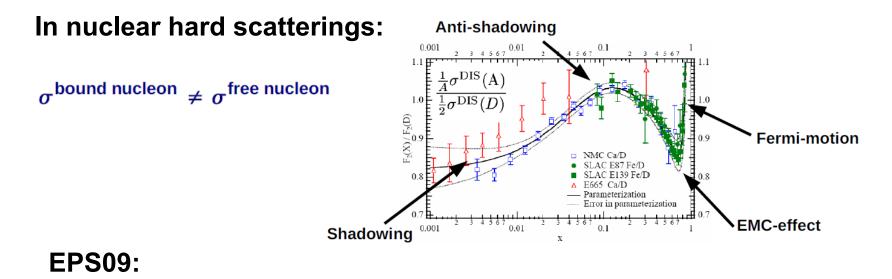
- data (=constraints) are not at one fixed Q² but correlated in x and Q² [fig.]
- how to account for the exp. statistical/systematic/normalization errors?
- parameter space 15-30 d & XSs require multi-d numerical integrations
- → need very fast DGLAP solver and XS solvers
 - we (EPS09) have them now in NLO [see Hannus PhD thesis]

Further challenges in nPDF analysis vs. that of free-proton PDFs:



For hadron collisions, excellent global fits [CTEQ, MRST/W,...]

- \rightarrow factorization theorem seems to hold well
- → PDF uncertainties have been quantified & error sets available



Q1: How well does factorization work for nuclear collisions -- can we find process-independent NLO nPDFs such that $\sigma_{AB \to h+X} = \sum_{i=i} f_i^A(Q^2) \otimes \hat{\sigma}_{ij \to h+X} \otimes f_j^B(Q^2) + \mathcal{O}(Q^2)^{-n}$

Q2: How large are the nPDF uncertainties?

Q3: Release error sets for the nPDFs?

Progress in the global nPDF analyses

1998:	EKS98	[Eskola, Kolhinen, Ruuskanen, Salgado] LO			
2001:	HKM	[Hirai, Kumano, Miyama]	LO + error estimates		
2004 :	HKN04	[Hirai, Kumano, Nagai]	LO + error est.		
2004 :	nDS	[de Florian, Sassot]	NLO		
2007:	EKPS	[Eskola, Kolhinen, Paukkunen, Salgado]	LO + error est.		
2007:	HKN07	[Hirai, Kumano, Nagai]	NLO + error est.		
2008:	EPS08	[Eskola, Paukkunen, Salgado]	LO, w. RHIC data		
2009:	EPS09	[Eskola, Paukkunen, Salgado]	NLO, w. RHIC data,		
			+ nPDF error sets!		

New elements in EPS09:

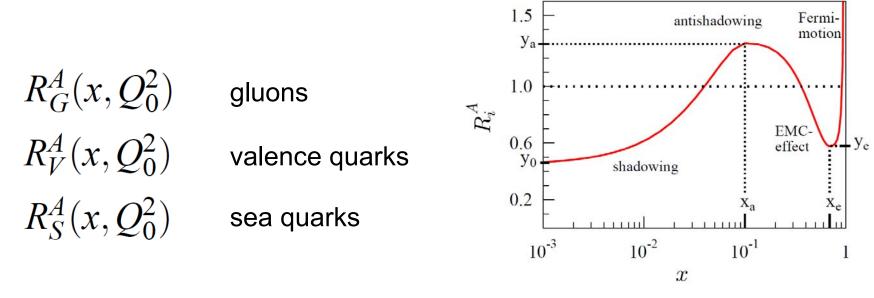
- with I+A DIS & DY in p+A, include also RHIC pi0 data from d+Au
- error analysis now at similar sophistication level as in CTEQ
- release the best fit + **30 error PDF sets** (for each A!) for general use
- → users can now study the propagation of nPDF-uncertainties into any hard XS!

EPS09 global analysis framework

• define the bound-proton PDFs vs. fixed free-p PDFs:

$$f_i^A(x,Q^2) \equiv R_i^A(x,Q^2) f_i^{\text{CTEQ6.1M}}(x,Q^2)$$

- MSbar & zero-mass variable flavor-number sceme
- isospin symmetry: $u_{p,n} = d_{n,p}$
- conserve momentum& baryon number
- initial parametrization at Q₀=1.3 GeV:



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A-dependence is in the parameters, e.g. y_a=y_a(C) (A/12) ^{pa}

The data in EPS09				
• 3 types				
DIS: I+A → I+X DY: p+A → I ⁺ I ⁻ +X pi0: d+Au → pi0+X at RHIC				
Altogether • 929 data points • 32 data sets				
$ \begin{array}{c} & \text{Drell-Yan} \\ & \text{SLAC DIS} \\ & \text{SLAC DIS} \\ & \text{NMC & EMC DIS} \\ & \text{PHENIX } \pi^0 \eta = 0.0 \end{array} $				
1.0				
$0.1 = BRAHMS h^{-} \eta = 2.2 = BRAHMS h^{-} \eta = 2.2 = BRAHMS h^{-} \eta = 3.2 = BRAHMS h^{-} \eta$				
$\begin{array}{cccccccccccccccccccccccccccccccccccc$				

Experiment	Process	Nuclei	Data points	χ^2 LO	χ^2 NLO	Weight
SLAC E-139	DIS	He(4)/D	21	6.5	7.3	1
NMC 95, re.	DIS	He/D	16	14.5	15.6	5
NMC 95	DIS	Li(6)/D	15	23.6	16.8	1
NMC 95, Q^2 dep.	DIS	Li(6)/D	153	162.2	157.0	1
SLAC E-139	DIS	Be(9)/D	20	9.6	12.2	1
NMC 96	DIS	$\mathrm{Be}(9)/\mathrm{C}$	15	3.8	3.8	1
SLAC E-139	DIS	C(12)/D	7	4.1	3.2	1
NMC 95	DIS	C/D	15	15.0	13.8	1
NMC 95, Q^2 dep.	DIS	C/D	165	141.8	142.0	1
NMC 95, re.	DIS	\dot{C}/D	16	19.3	20.5	1
NMC 95, re.	DIS	C/Li	20	30.3	28.4	1
FNAL-E772	DY	C/D	9	7.5	8.3	1
SLAC E-139	DIS	Al(27)/D	20	10.9	12.5	1
NMC 96	DIS		15	6.0	5.8	1
NMC 90	D15	Al/C	19	0.0	0.0	1
SLAC E-139	DIS	Ca(40)/D	7	5.0	4.1	1
FNAL-E772	DY	Ca/D	9	2.9	3.4	15
NMC 95, re.	DIS	Ca/D	15	25.4	24.7	1
NMC 95, re.	DIS	Ca/Li	20	23.9	19.6	1
NMC 96	DIS	Ca/C	15	6.0	6.0	1
SLAC E-139	DIS	Fe(56)/D	26	19.1	23.9	1
FNAL-E772	DY	Fe/D	9	2.1	2.2	15
NMC 96	DIS	$\overline{\rm Fe}/{\rm C}$	15	11.0	10.8	1
FNAL-E866	DY	Fe/Be	28	20.9	21.7	ĩ
		,	-	2010		Ŷ
CERN EMC	DIS	Cu(64)/D	19	13.4	14.8	1
SLAC E-139	DIS	Ag(108)/D	7	3.8	2.9	1
NMC 96	DIS	Sn(117)/C	15	9.6	9.1	1
NMC 96, Q^2 dep.	DIS	Sn/C	144	80.2	82.8	10
rune oo, q dep.	210	<i>S</i> H / C		00.2	01.0	(x=0.0125 only
FNAL-E772	DY	W(184)/D	9	7.0	6.7	10
FNAL-E866	DY	W/Be	28	27.3	24.2	1
SLAC E-139	DIS	Au(197)/D	21	11.6	13.8	1
RHIC-PHENIX	π^0 prod.	dAu/pp	20	7.3	6.3	20
	-	/ = =	15	6.00		
NMC 96	DIS	$\rm Pb/C$	15	6.90	7.2	1
Total			929	738.6	731.3	

Generalized Chi2 in EPS09

usual definition
$$\chi_N^2(\{z\}) \equiv \sum_{i \in N} \left[\frac{D_i - T_i(\{z\})}{\sigma_i} \right]^2$$
now define
$$\chi^2(\{z\}) \equiv \sum_N w_N \chi_N^2(\{z\})$$
$$\chi_N^2(\{z\}) \equiv \left(\frac{1 - f_N}{\sigma_N^{\text{norm}}} \right)^2 + \sum_{i \in N} \left[\frac{f_N D_i - T_i(\{z\})}{\sigma_i} \right]^2$$

1. Weights w_N for data sets N

r

which provide important constraints but whose #data points is small

- 2. Treatment of the additional overall normalization errors σ^{norm} given by the experiment (also in e.g. CTEQ):
 - f_N is determined by minimizing the chi²_N for each data set N, for each parameter-set candidate during the overall chi² minimization
 - "penalty factor" $(...)^2$ accounts for the fact that $f_N = 1$ should be the optimal value \Leftrightarrow the normalization given by the experiment
 - f_N is an output of the analysis

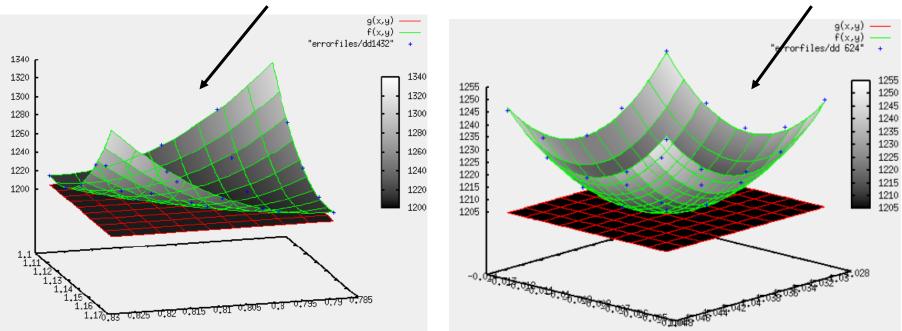
Q: How to estimate the uncertainties in the nPDFs ? How to propagate these into hard XSs ?

A: Error analysis in EPS09 via Hessian method

• expand around the minimum, this defines the Hessian matrix H

$$\chi^{2} - \chi^{2}_{0} \approx \sum_{ij} \frac{1}{2} \frac{\partial^{2} \chi^{2}}{\partial a_{i} \partial a_{j}} (a_{i} - a_{i}^{0}) (a_{j} - a_{j}^{0}) \equiv \sum_{ij} H_{ij} (a_{i} - a_{i}^{0}) (a_{j} - a_{j}^{0})$$

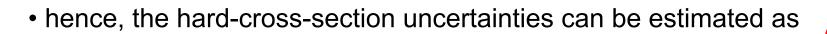
parameters are <u>correlated</u>, H must be diagonalized and <u>correlations removed</u>



O(15) parameters, O(100) pairs

 eigenvectors of H provide a basis {z_i} of uncorrelated parameter combinations, where

$$\chi^2 \approx \chi_0^2 + \sum_i z_i^2 +$$



60 ⊽×°

20

-10

-5

Eigendirections

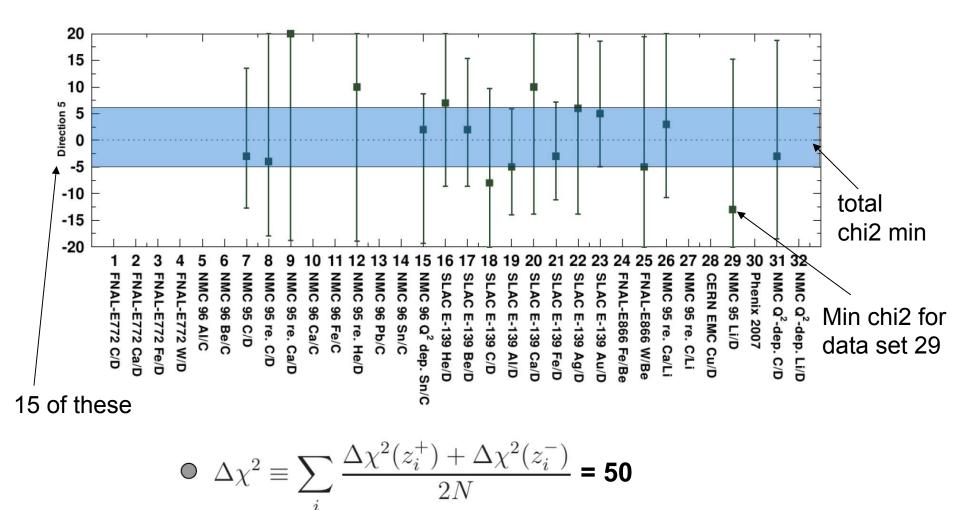
$$(\Delta X)^2 = \left(\frac{\partial X}{\partial z_1} \cdot \delta z_1\right)^2 + \left(\frac{\partial X}{\partial z_2} \cdot \delta z_2\right)^2 + \cdots$$

- but in the uncertainty estimates, how large should δz_k i.e. $\Delta \chi^2$ be, in order to correctly reflect the uncertainties in the <u>data ?</u>
 - \rightarrow No unique procedure exists for this
 - → use a "90 % confidence criterion" similar to CTEQ

5

We take $\Delta \chi^2 = 50$ based on requiring the χ^2 -contribution of each data set to stay within its 90% confidence range.

 \bigcirc



[from HP]

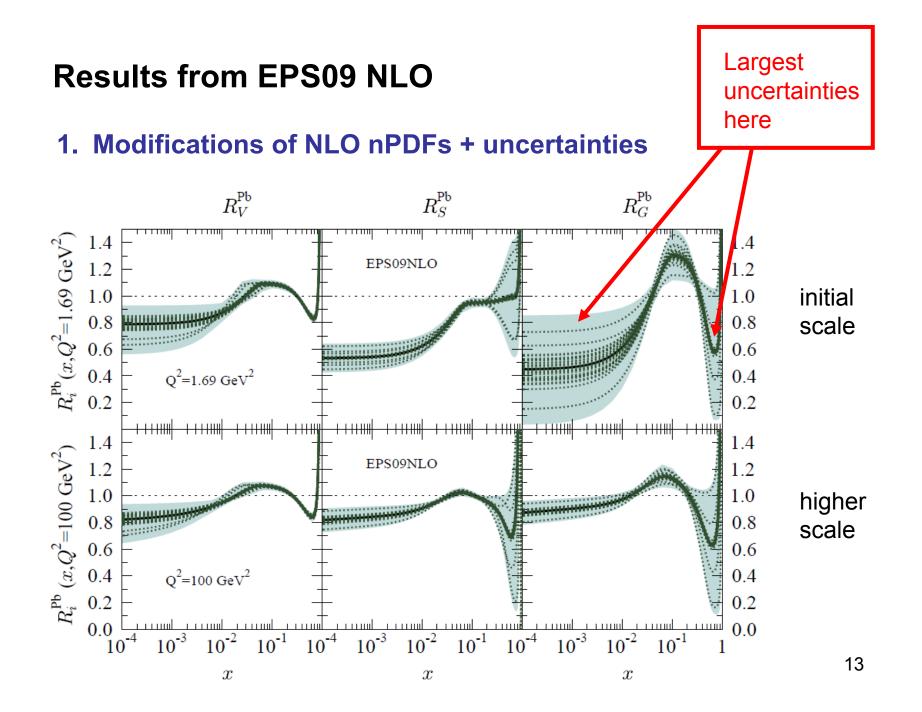
The EPS09 package

= the central set, **best fit S⁰ + 30 nPDF error sets** S_k **±** where each S_k **±** is obtained by displacing the fit parameters to the +/direction along the eigendirections z_k such that chi2 grows by 50

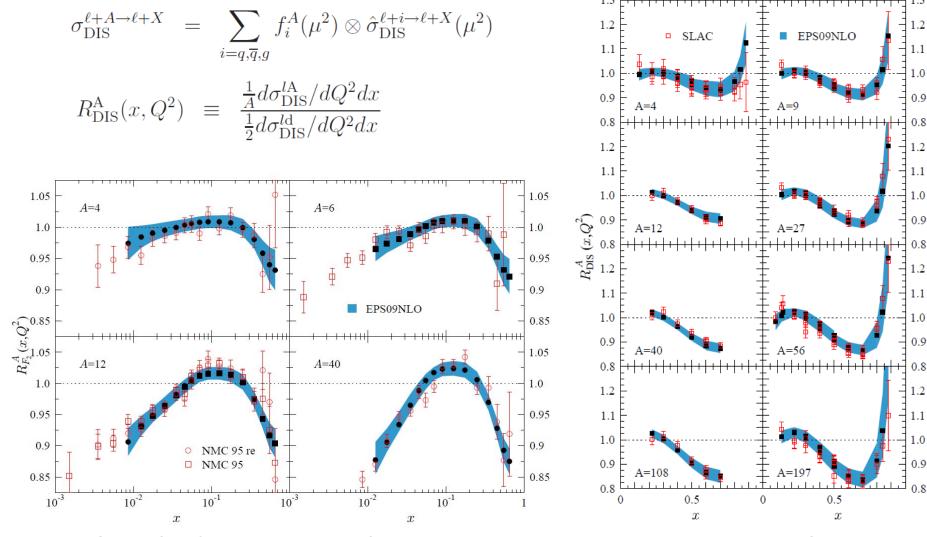
the user can then

 compute the propagation of the nPDF-originating uncertainties into the cross section X of his/her interest as follows:

$$(\Delta X^{+})^{2} \approx \sum_{k} \left[\max \left\{ X(S_{k}^{+}) - X(S^{0}), X(S_{k}^{-}) - X(S^{0}), 0 \right\} \right]^{2}$$
$$(\Delta X^{-})^{2} \approx \sum_{k} \left[\max \left\{ X(S^{0}) - X(S_{k}^{+}), X(S^{0}) - X(S_{k}^{-}), 0 \right\} \right]^{2}$$

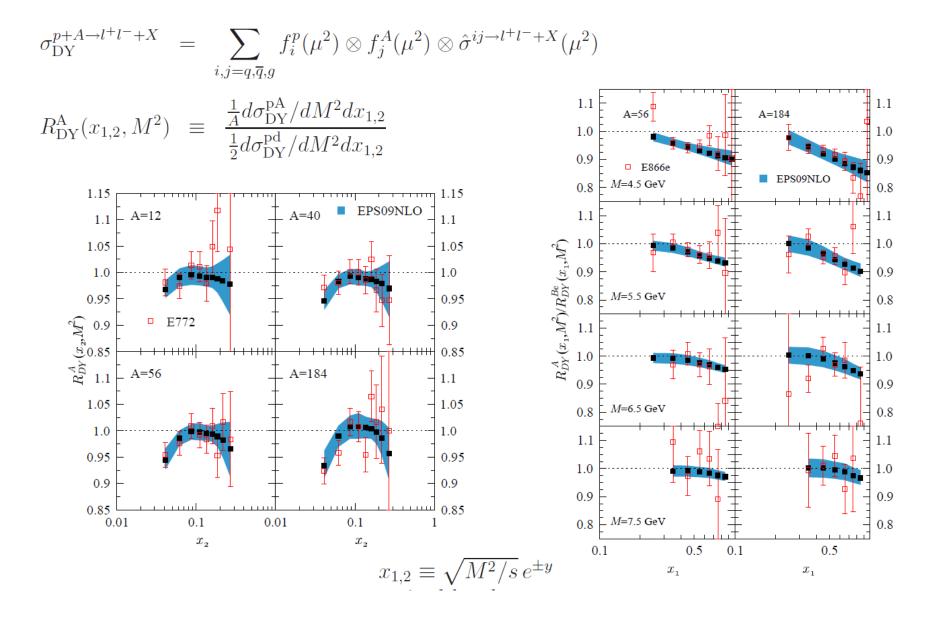


2. Comparison with data: NLO DIS + uncertainties



Good fits & Propagation of the exp. uncertainties into the nPDFs is OK !

3. Comparison with data: NLO DY + uncertainties



4. Comparison with data: DIS vs. NLO DGLAP evolution

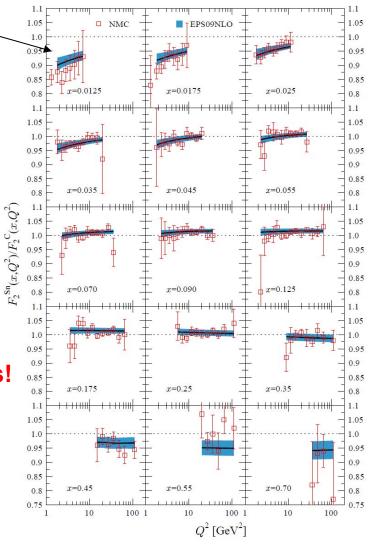
 $F_2^{\mathrm{Sn}}/F_2^{\mathrm{C}}$

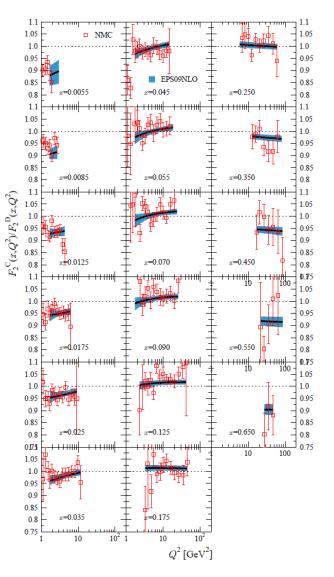
 $F_2^{\rm C}/F_2^{\rm D}$

Due to this panel, gluons at x~0.03 are ~well constrained

but...

smaller-x data would be needed to reduce the large uncertainties in small-x gluons!

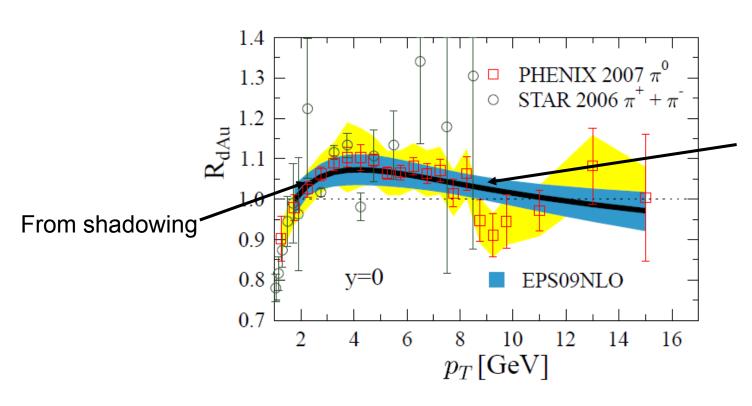




5. Comparison with data: pi0 (y=0, pT>1.7 GeV) in d+Au at RHIC

$$\sigma^{A+B\to\pi+X} = \sum_{i,j,k=q,\overline{q},g} f_i^A(\mu^2) \otimes f_j^B(\mu^2) \otimes \hat{\sigma}^{ij\to k+X}(\mu^2) \otimes D_{k\to\pi}(\mu^2)$$

$$R_{\rm dAu}^{\pi} \equiv \frac{1}{\langle N_{\rm coll} \rangle} \frac{d^2 N_{\pi}^{\rm dAu}/dp_T dy}{d^2 N_{\pi}^{\rm pp}/dp_T dy} \stackrel{\rm min.bias}{=} \frac{\frac{1}{2A} d^2 \sigma_{\pi}^{\rm dAu}/dp_T dy}{d^2 \sigma_{\pi}^{\rm pp}/dp_T dy}$$



Surprisingly good, fairly tensionless fit

Suggests an EMC-depletion for gluons – the treatment of overall normalization errors important!

EPS09 – summary

- Excellent agreement between NLO pQCD and the hard-process nuclear data for DIS, DY, and π^o production in the kinematical range 0.005 < x < 1, 1.69 GeV² < Q² < 150 GeV² --- chi2/N = 0.79
- No significant tension between the data sets from different processes

→ Factorization seems to work well in describing high-E inclusive nuclear hard processes

• The **EPS09** nPDF package

central set (best fit) + 30 error sets both in NLO and LO, https://www.jyu.fi/fysiikka/en/research/highenergy/urhic/nPDFs

→ Estimation of the nPDF-uncertainty propagation into hard nuclear XSs is now, finally, possible for <u>anybody</u>!

- further tests of factorization and nPDFs provided by
 - RHIC data for

direct photons in d+Au (also Au+Au) heavy quarks in d+Au pion production at fwd-y

p+A at the LHC – but are such runs foreseen?

future facilities eRHIC,

- on our near-future work list
 - + include also neutrino DIS data (NuTeV, ...)
 - \rightarrow further constraints for quarks?
 - + extend the analysis to general-mass framework

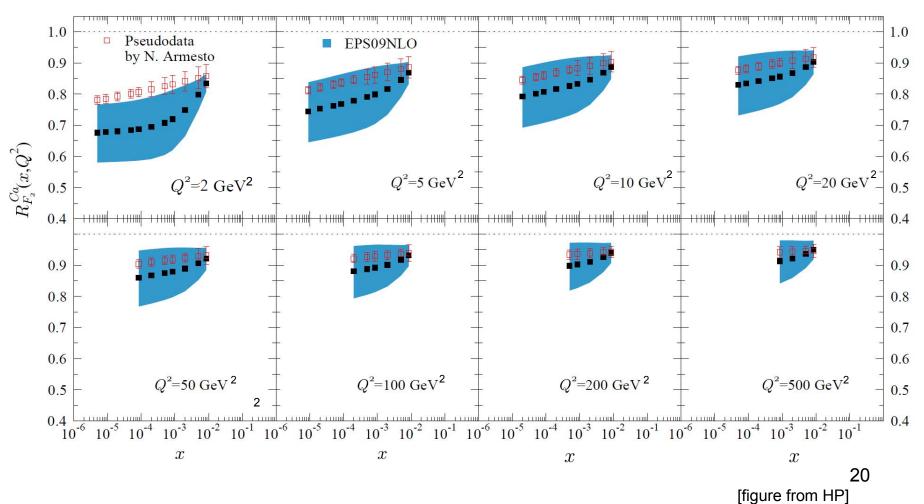
a longer-term goal

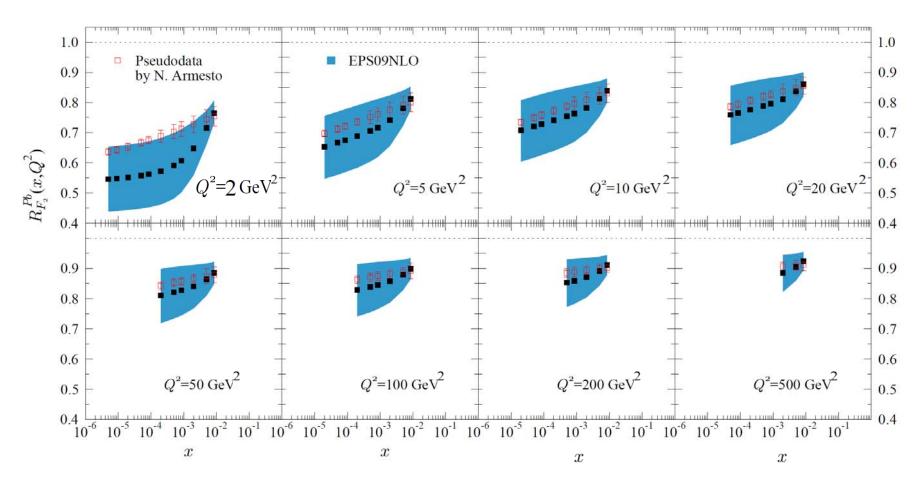
a master global analysis which combines PDFs and nPDFs

Impact of the LHeC for the global analysis of nPDFs?

1. Plot the LHeC pseudodata [generated by N. Armesto] against EPS09

Calsium, A=40





Inclusion of such LHeC data in the global analysis should reduce the nPDF uncertainties at small-x significantly

2. Global NLO fit with LHeC pseudodata [from N. Armesto] included [performed by Hannu Paukkunen for this workshop]

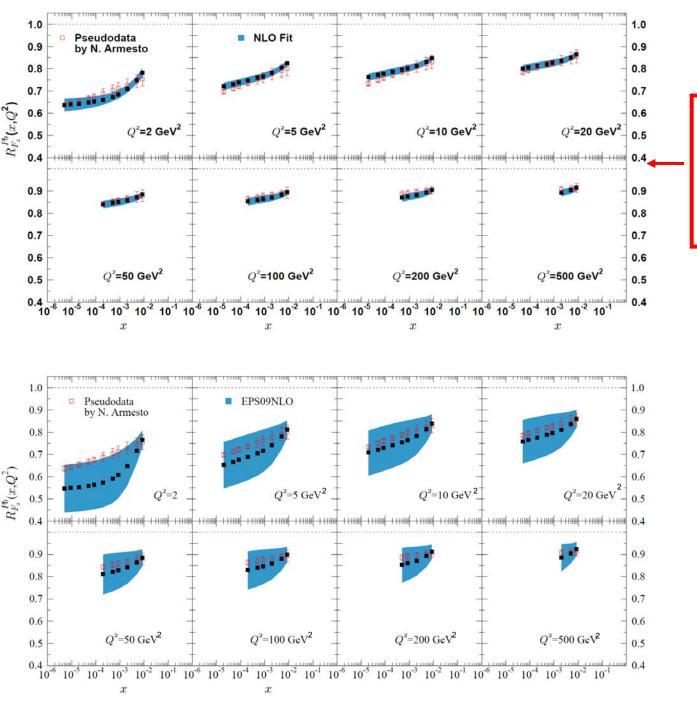
- keep the EPS09 set-up (fit functions, khi2 definition, weighting, error analysis, etc)
- keep the same $\Delta \chi^2 = 50$ in the error analysis
- decrease the PHENIX data weight (now more data on gluons, can now do this)
- with the LHeC data, we can release one more gluon parameter (x_a)
- $\rightarrow \qquad \underline{\text{Quantify}} \text{ how much the nPDF uncertainties can be expected to reduce}$

A=208

LHeC pseudodata included in the fit; uncertainties much smaller

 $\Delta \chi^2$ = 50 seems OK here, too

EPS09 – LHeC pseudodata not included



LHeC pseudodata included in the fit; uncertainties much smaller

A=40

1.0

0.9 0.8

0.7

0.6

0.5

0.4

0.9

0.8

0.7

0.6

0.5

0.4

10.00

 $Q^2=20 \text{ GeV}^2$

2 X X

 $Q^2 = 500 \text{ GeV}^2$

x

10⁻⁵ 10⁻⁴ 10⁻³

10⁻² 10⁻¹

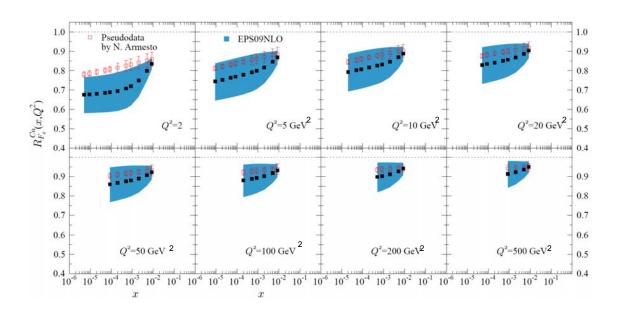
 $O^2 = 10 \text{ GeV}^2$

 $Q^2 = 200 \text{ GeV}^2$

10⁻³ 10⁻²

x

10-1 10-6



1.0

0.9

0.8 0.7

0.5

0.4

0.9

0.8

0.7

0.6

0.5

0.4

10-5

10-4 10-3

 $R_{F_2}^{Ca}(x,Q^2)$ 0.6 Pseudodata

by N. Armesto

 $O^2 = 2 \text{ GeV}^2$

 $Q^2 = 50 \text{ GeV}^2$

x

10-2

10-1

NLO Fit

 $Q^2 = 100 \text{ GeV}^2$

10⁻¹

10⁻⁶ 10⁻⁵

104

10⁻⁴ 10⁻³ 10⁻²

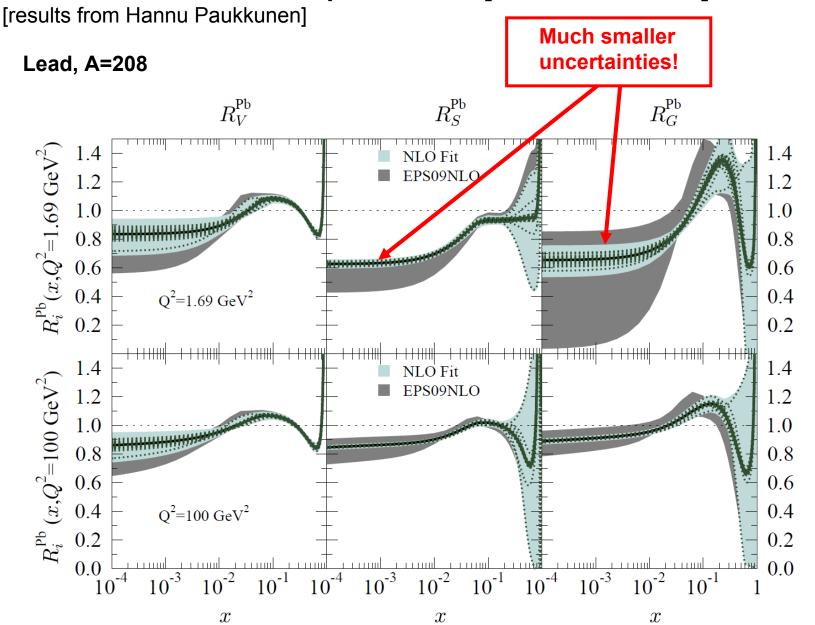
x

10-5

10-6

 $Q^2 = 5 \text{ GeV}^2$

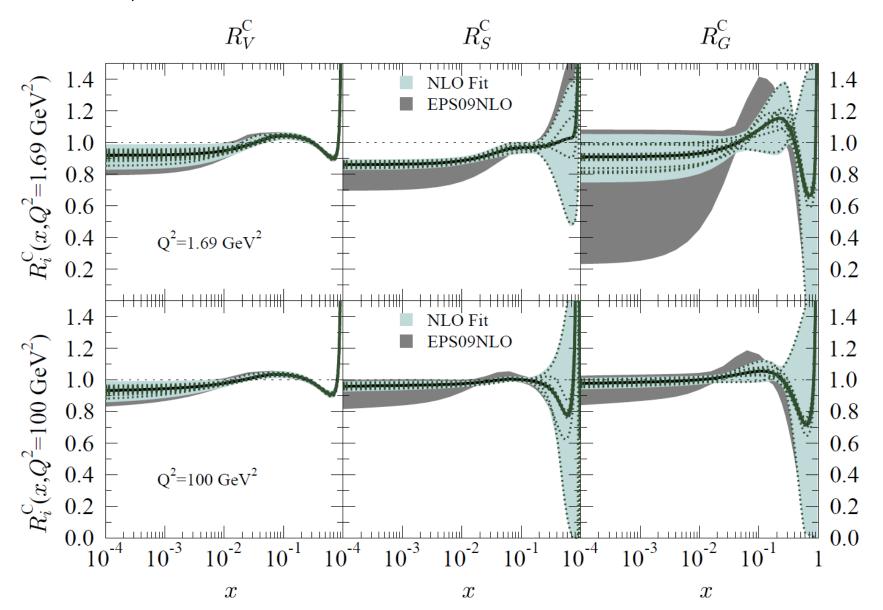
EPS09; LHeC pseudodata not included



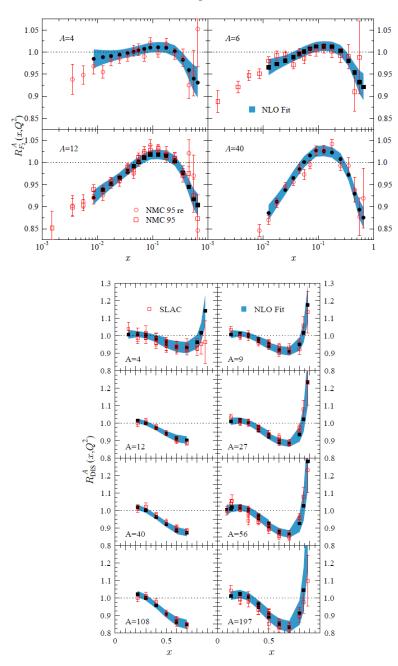
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Global NLO fit with LHeC pseudodata [from N. Armesto] included

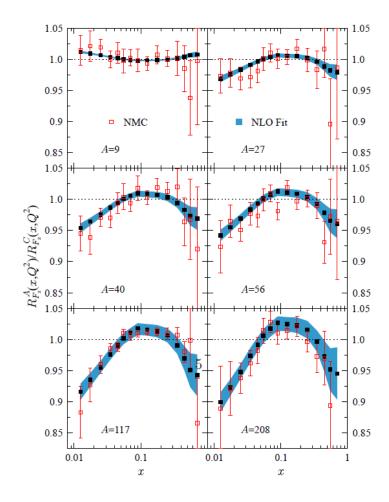
Carbon, A=12



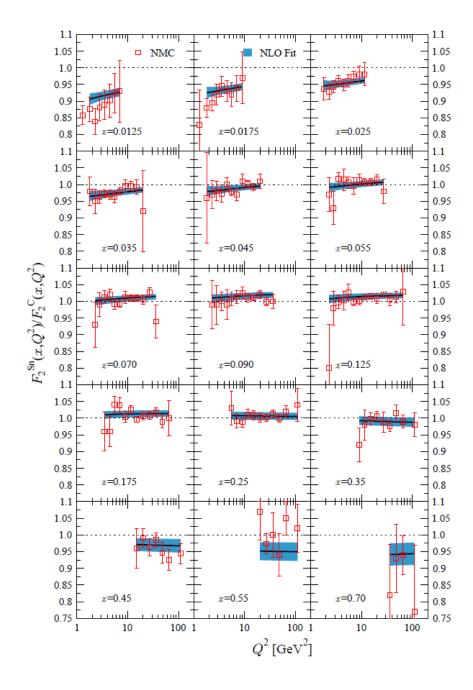
Also with the LHeC pseudodata included, the global NLO fit remains very good...

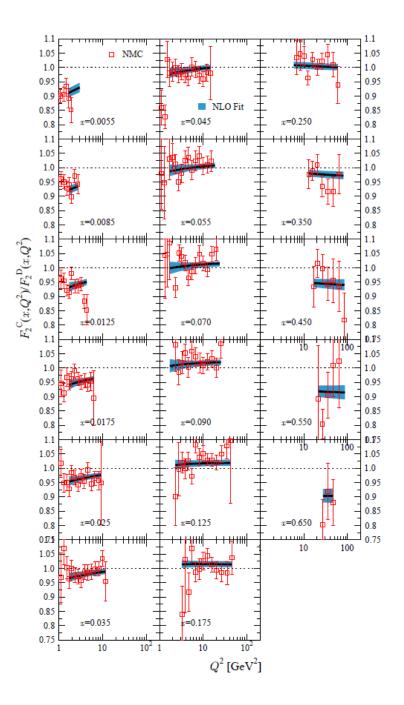


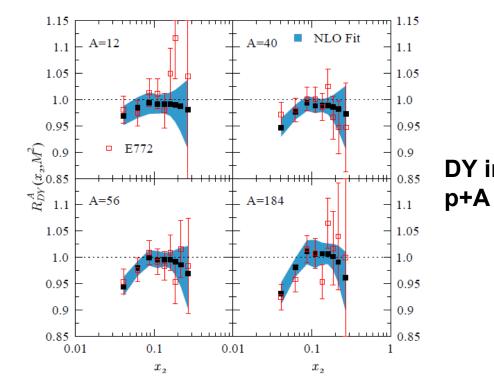




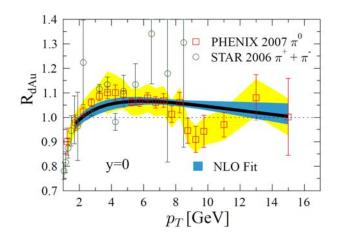
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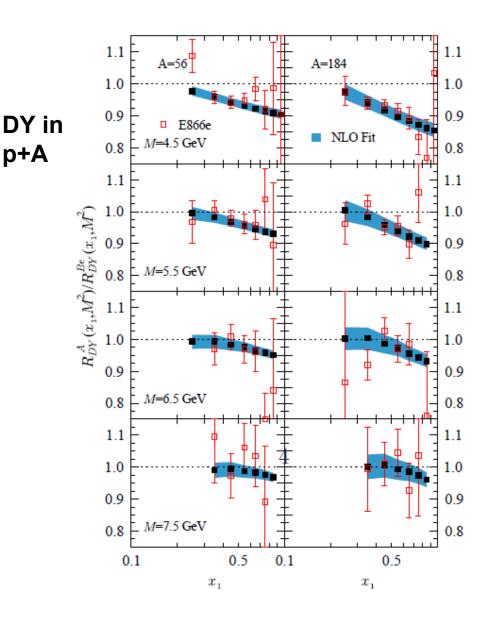






pi0 in d+A at RHIC





29

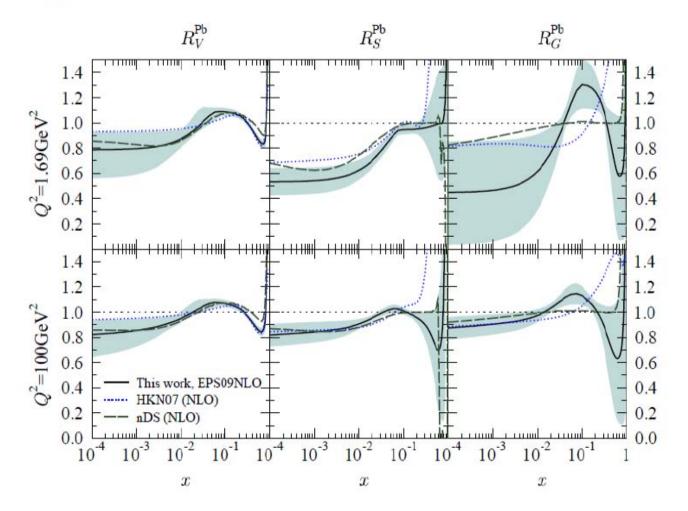
Conclusion

Performing a global analysis of nPDFs with a set of simulated LHeC pseudodata, we have demonstrated the effect of the expected LHeC data on the nPDF uncertainties.

The high-precision LHeC data would play a crucial role in pinning down the nPDFs in the small-x region at perturbative scales. Extra slides

NLO	Param.	Valence R_V^A	Sea R_S^A	Gluon R_G^A
1	y_0	Baryon sum rule	0.785	Momentum sum rule
2	p_{y_0}		-0.136	
3	x_a	$6.56 imes10^{-2}$	$8.74 imes10^{-2}$	0.1 fixed
4	p_{x_a}	0, fixed	0, fixed	0, fixed
5	x_e	0.688	as valence	as valence
6	p_{x_e}	0, fixed	0, fixed	0, fixed
7	y_a	1.05	0.970	1.207
8	p_{y_a}	$1.47 imes10^{-2}$	-8.12 $ imes 10^{-3}$	$2.72 imes10^{-2}$
9	y_e	0.901	1.076	0.625
10	p_{y_e}	-2.81 $ imes 10^{-2}$	as valence	as valence
11	\dot{eta}	1.31	1.3, fixed	1.3, fixed
12	p_{eta}	4.63×10^{-2}	0, fixed	0, fixed

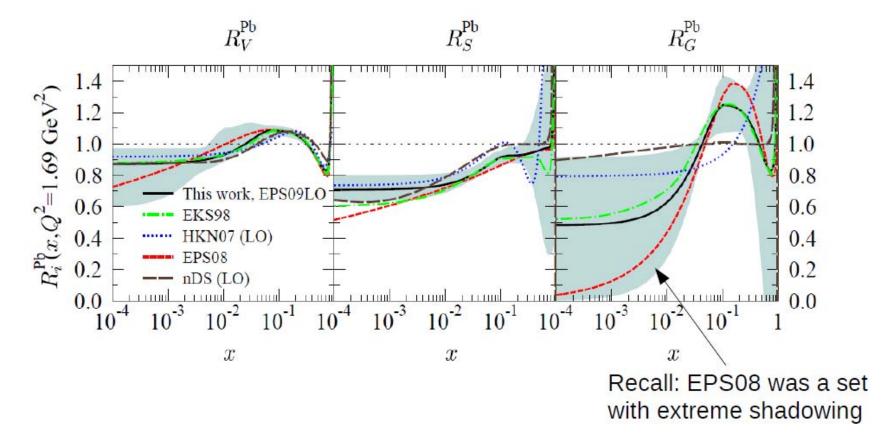
Comparison with earlier works: NLO



 Difference with EPS09NLO & HKN07 gluons results mainly from the more restricted form of HKN07 fit function and inclusion of pion data

[from HP]

Comparison with earlier works: LO



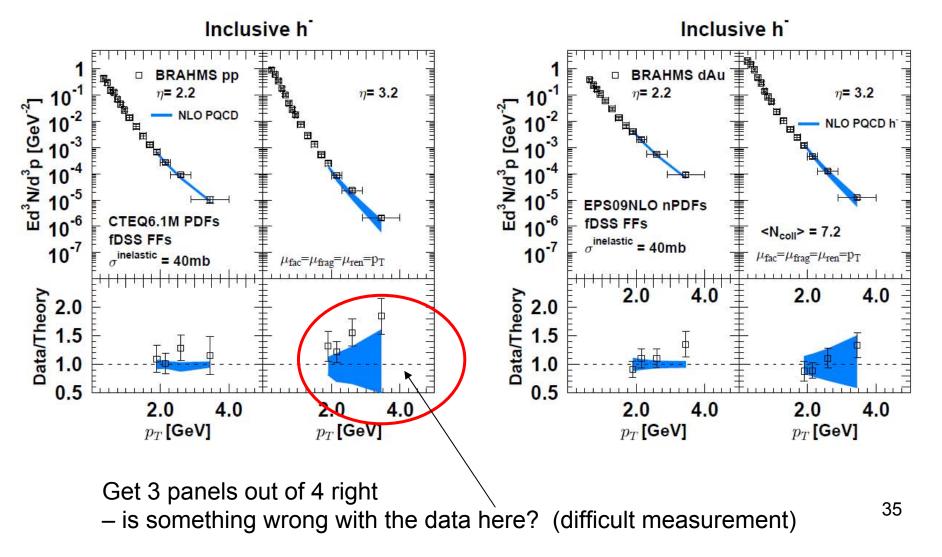
The EPS09LO errorbands are clearly larger than in NLO

[from HP]

4. An Application: fwd-y h⁻ production at RHIC vs. BRAHMS data

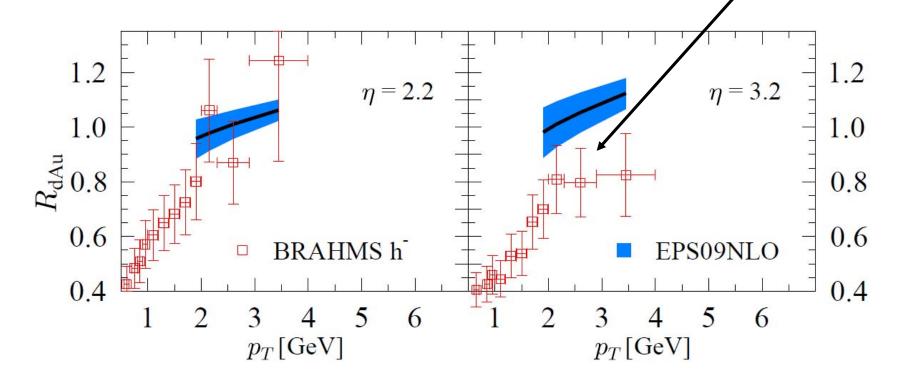
p+p w. CTEQ error sets

d+Au w. CTEQ&EPS09 error sets



Origin of the much-debated suppression in the BRAHMS R_{dAu}??

EPS09 \rightarrow the suppression/saturation observed in the ratio R_{dAu} at pT>2 GeV, y=3.2, is caused by the excessive yield measured in the p+p case, NOT by a suppression in d+Au!



N.B: This data set was <u>not</u> included in EPS09