

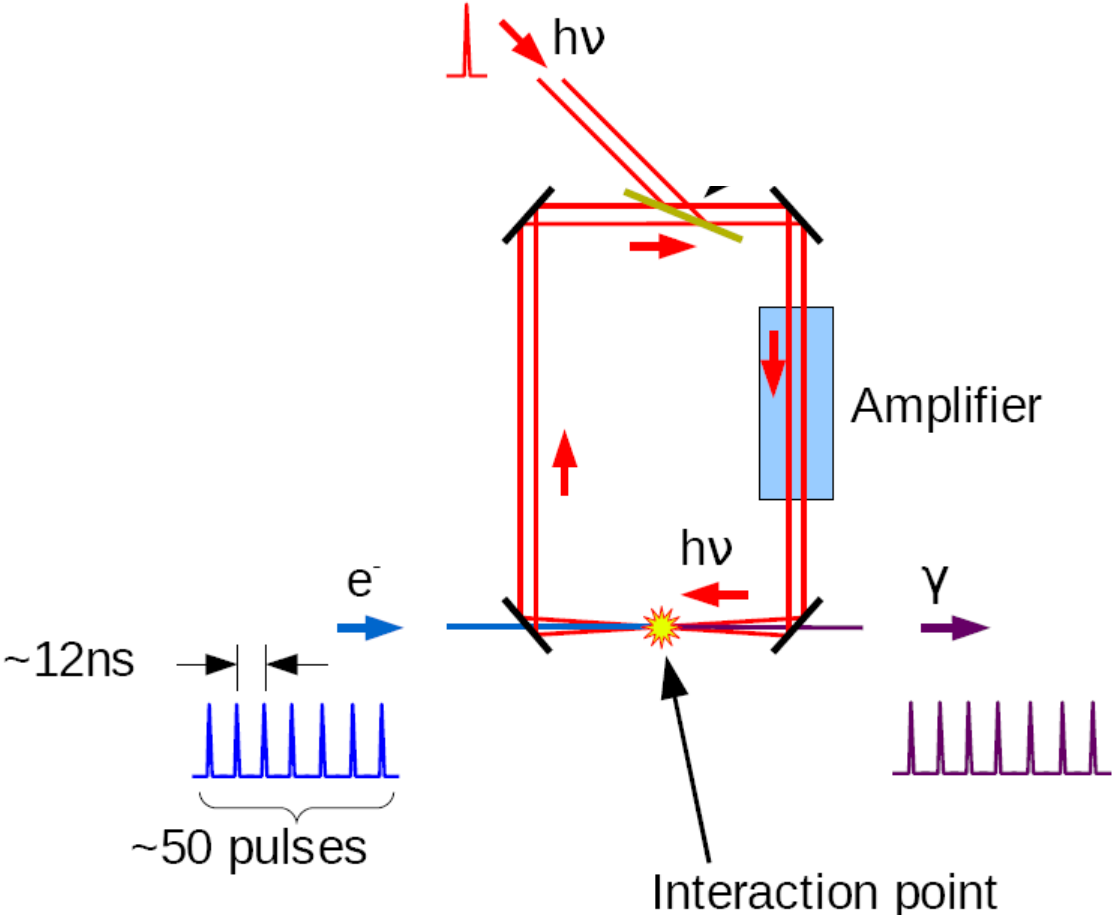
High intensity gamma source for ... positrons (LeHC)

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BNL

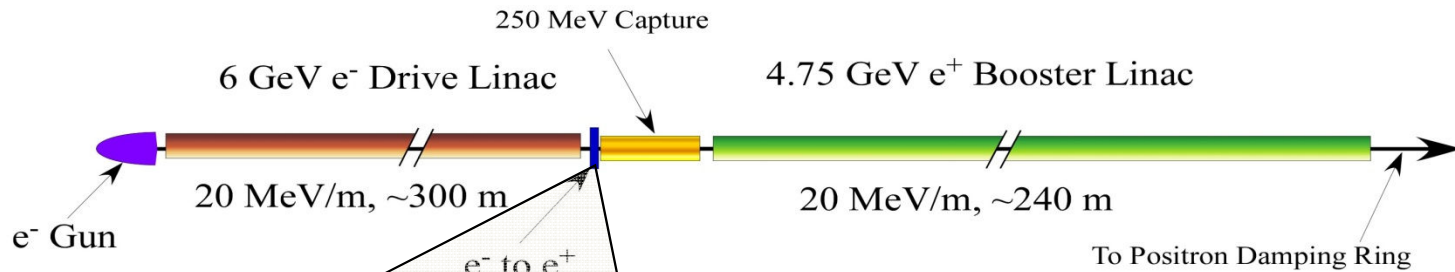
Divonne, September 2, 2008

Concept of the Multi-kHz gamma source

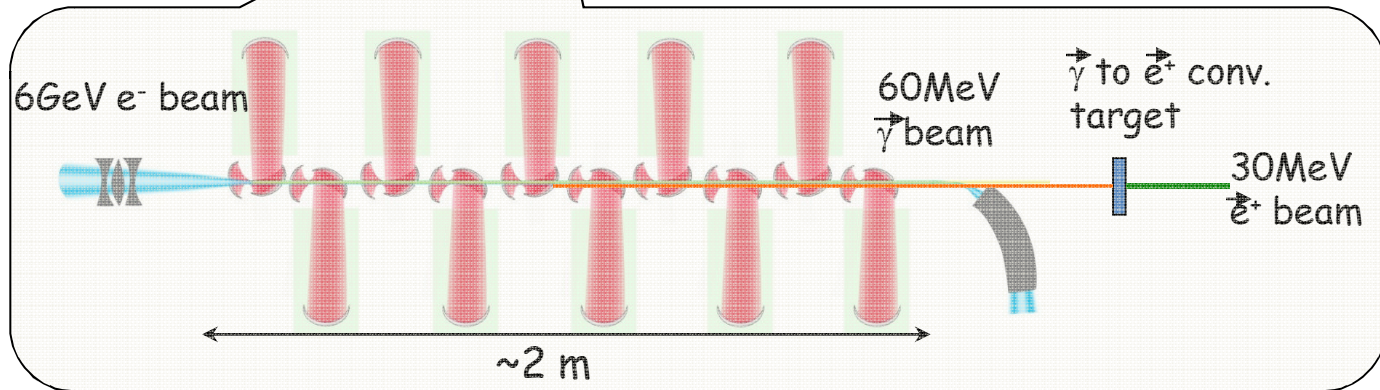


Polarized Positrons Source for ILC, CLIC, Super B

Conventional Non-Polarized Positrons:

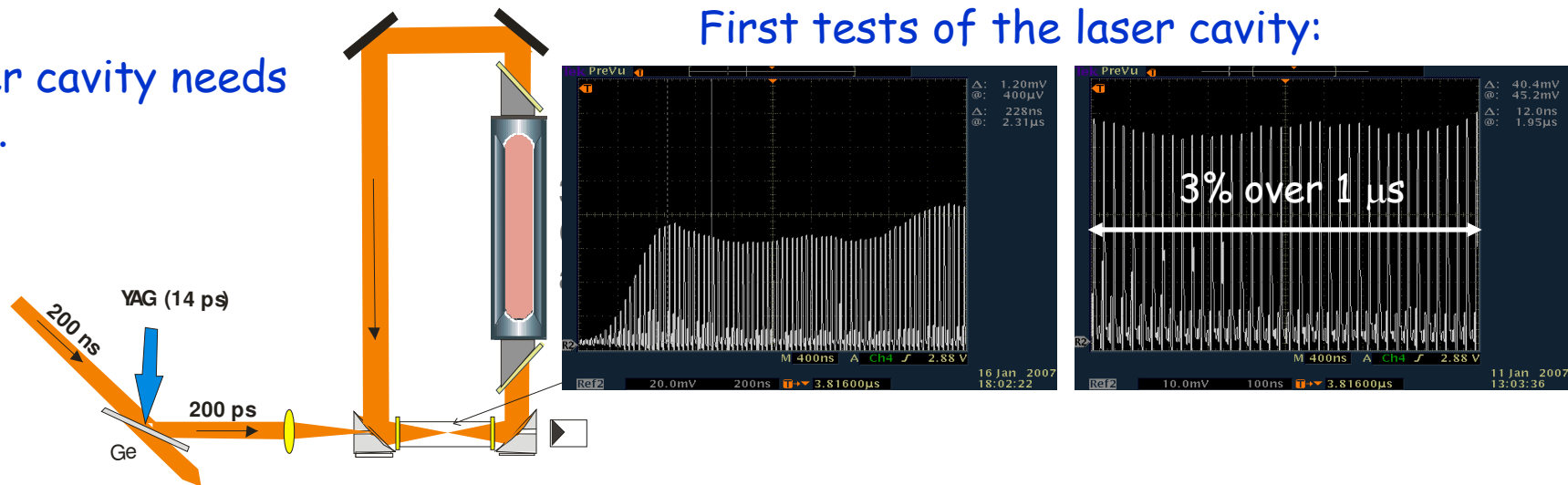


Polarized γ -ray beam is generated in the Compton back scattering inside optical cavity of CO2 laser beam and 6 GeV e-beam produced by linac.



Laser cavity needs R&D.

First tests of the laser cavity:



Simulations

Input: 5 ps, 10 μ J

Active medium: 30 cm

Gas:

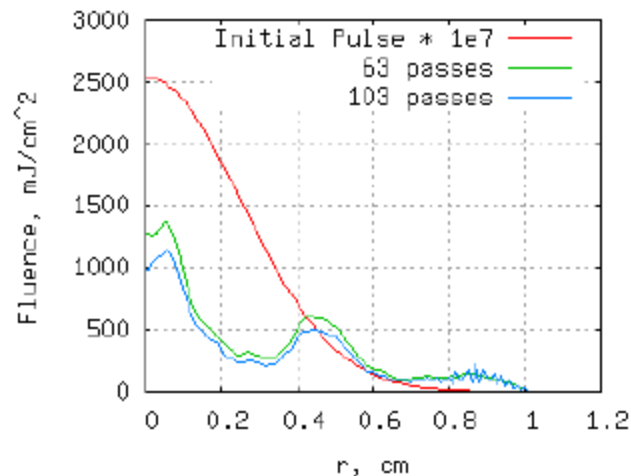
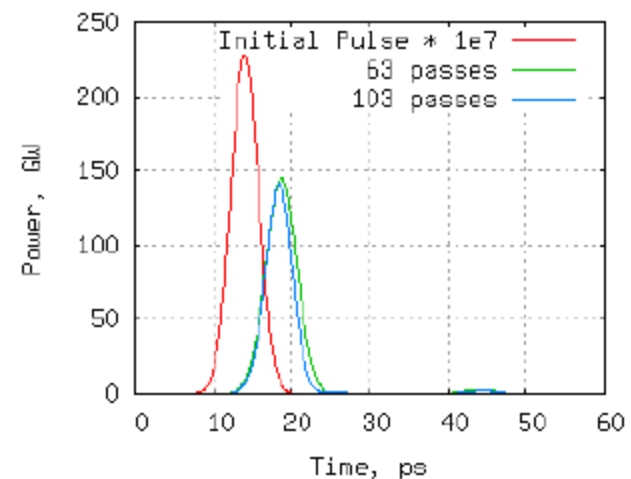
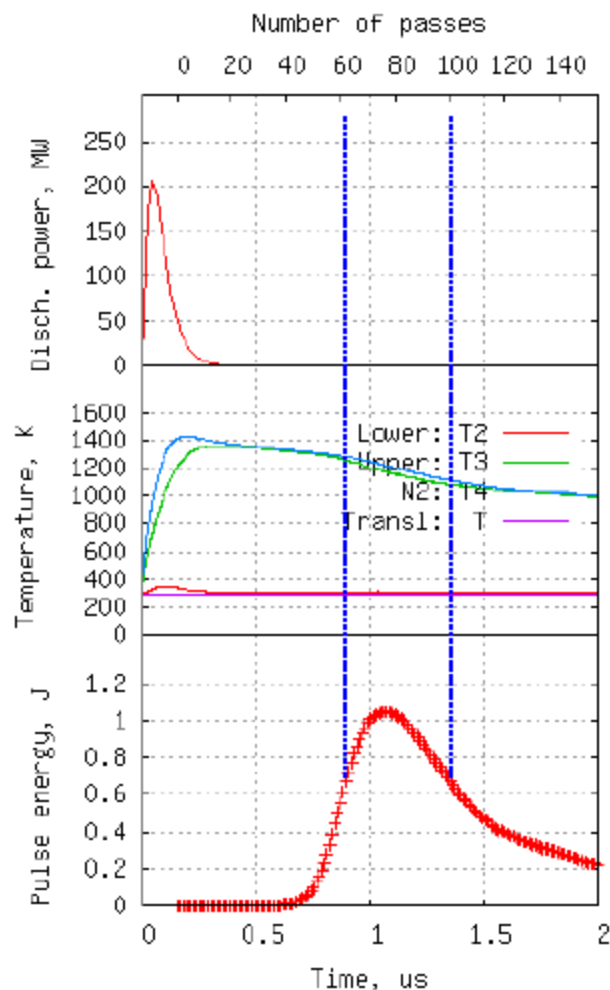
0.6:0.6:8.8 atm
(CO₂:N₂:He)

CO₂ isotopes:

0.15:0.3:0.15 atm
(CO₂¹⁶:CO¹⁶O¹⁸:CO₂¹⁸)

Injection: 1% reflector

Optical losses: 5% / pass

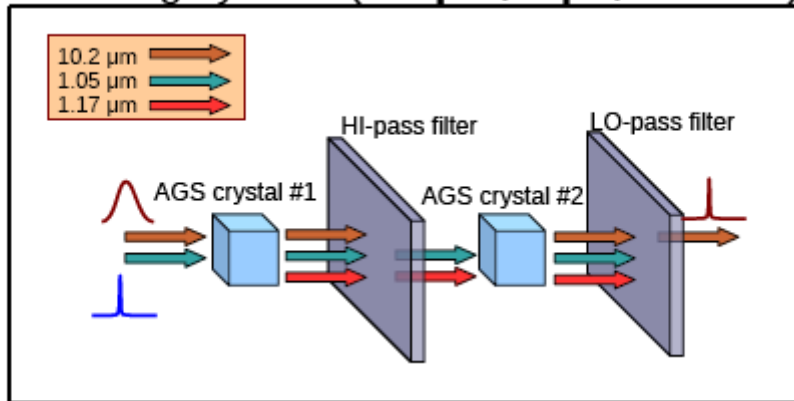


Laser system

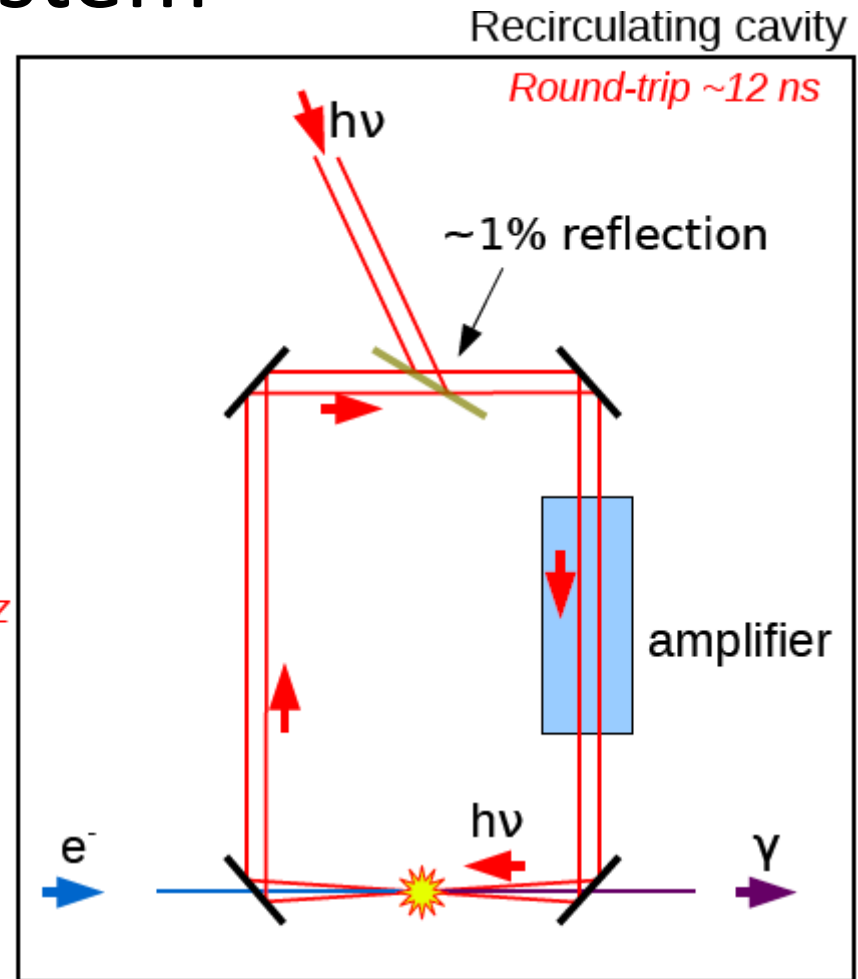
CO₂ laser system
(10 μm, 80 ns, 120 Hz)



Pulse forming system (10 μm, 3 ps, 120 Hz)



Solid state laser system
(1 μm, 3 ps, 120 Hz)



Realistic goal:

Trains: 25~50 pulses @ 120 Hz
(effective repetition rate: 3~6 kHz)

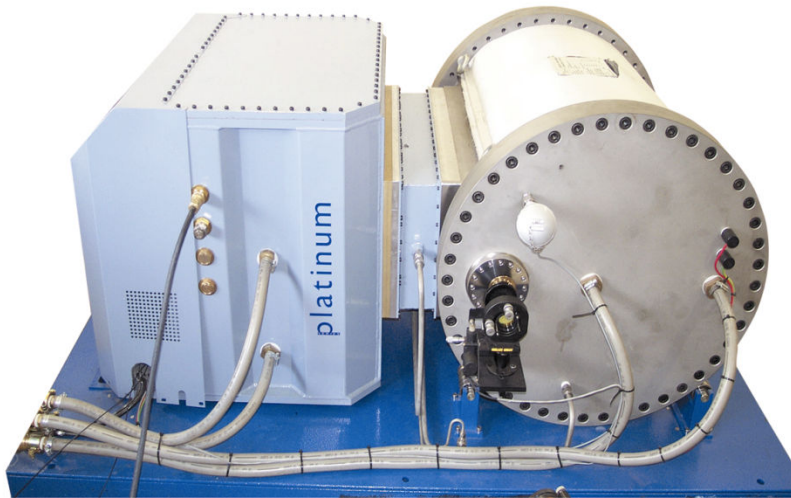
Pulse energy: 0.5~1.0 J

Pulse duration: 3~5 ps (fwhm)

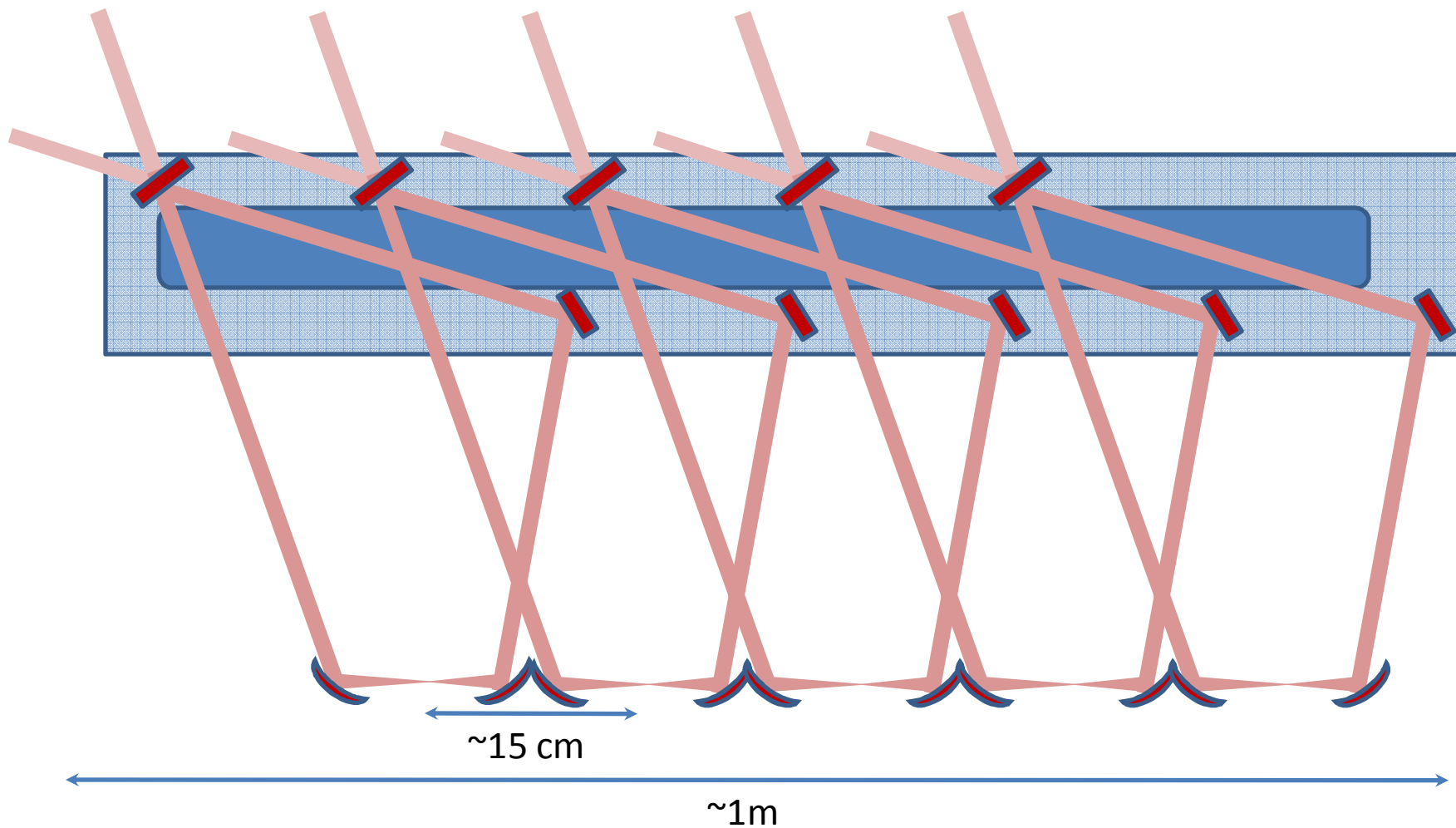
Lasers from SDI

<http://www.lightmachinery.com/SDI-CO2-lasers.html>

Wavelength	9 - 11 μ m, Line Tunable
Continuous	20 Hz 100 Hz 350 Hz 500 Hz
Repetition Rate	
Pulse Energy	1.5 J
Mode Type	Multimode
Optional:	TEM ₀₀ , custom beam shapes, SLM
Beam Size	13 x 13 mm ²
Average Power	30 W 150 W 525 W 750 W
Power Stability	< 7 %



Possible configuration with 5 IPs and 1 laser amplifier



Choice of parameters

- Train of 50 bunches is generated at 300Hz and will form ~ 3000 bunches x 5Hz of ILC beam.
- ~ 6 -24 ns bunch spacing can be selected.
- High power/10 atm. (picoseconds beam) CO₂ lasers are commercially available at up to 500Hz.
- ~ 30 -40 μm laser focus is set by practical considerations of electron and laser beams focusing and requires ~ 5 ps long pulses (hour glass effect).
- Nonlinear effects in Compton back scattering limit laser energy at $\sim 1\text{J}$
- Train of ~ 5 nC electron bunches are required to produce 1 nC of polarized positrons per bunch.
- ~ 1 γ -ray per 1 electron per laser IP, 10 IPs (each electron emits 10 γ)
- Conversion efficiency of gammas into captured polarized positrons is simulated at $\sim 2\%$.

Emittances $E_e = 6\text{GeV}$ $L_{target} \approx 0.4X_0$

- Emittance of the positron beam at the target exit $\varepsilon_N \approx \gamma_{e+} \sigma \sigma'$;

- Acquired angular spread $\sigma'_{e+} \approx \frac{1}{\sqrt{2}} \frac{14\text{MeV}}{E_{e+}} \sqrt{\frac{L_{target}}{X_0}} \approx 12 / \gamma_{e+} \approx 0.14$

- Beam size

- Gamma beam divergence $\sigma_{\gamma d} \approx \frac{1}{2\gamma_{e-}} \frac{L_{IR}}{\sqrt{2}} \approx \frac{1}{2 \cdot 12000} \frac{2m}{\sqrt{2}} \approx 60\mu m$

- Scattering in the target $\sigma_{e+sc} \approx \frac{\sqrt{2}}{3} \sigma'_{e+} L_{target} \approx \frac{\sqrt{2}}{3} 0.14 1.2\text{mm} \approx 80\mu m$

- e^- beam size on target $\sigma_{\gamma e-} \approx \sqrt{\frac{\varepsilon_{Ne-}}{\gamma_{e-}} \beta_{e-}} \approx \sqrt{\frac{10\mu m}{12000}} 10\text{cm} \approx 10\mu m$

$$\varepsilon_{Nsc} \sim 1\text{mm} \quad \text{and} \quad \varepsilon_{N\gamma d} \sim 0.7\text{mm}$$

- Longitudinal emittance

$$\varepsilon_{\square N} \approx \Delta\gamma_{e+} \sigma_{\tau e-} = \frac{120 - 60}{4} 30\mu m = 450\mu m$$

Emittances $E_e=12\text{GeV}$ $L_{\text{target}} \square 0.1\text{X}0$

- Emittance of the positron beam at the target exit $\varepsilon_N \square \gamma_{e+} \sigma \sigma'$;

- Acquired angular spread $\sigma'_{e+} \square \frac{1}{\sqrt{2}} \frac{14\text{MeV}}{E_{e+}} \sqrt{\frac{L_{\text{target}}}{X_0}} \approx 6/\gamma_{e+} \approx 0.017$

- Beam size

- Gamma beam divergence $\sigma_{\gamma d} \approx \frac{1}{2\gamma_{e-}} \frac{L_{IR}}{\sqrt{2}} \approx \frac{1}{2 \cdot 24000} \frac{2\text{m}}{\sqrt{2}} \approx 28\mu\text{m}$

- Scattering in the target $\sigma_{e+sc} \approx \frac{\sqrt{2}}{3} \sigma'_{e+} L_{\text{target}} \approx \frac{\sqrt{2}}{3} 0.017 \cdot 0.3\text{mm} \approx 2\mu\text{m}$

- e^- beam size on target $\sigma_{\gamma e-} \approx \sqrt{\frac{\varepsilon_{Ne-}}{\gamma_{e-}} \beta_{e-}} \approx \sqrt{\frac{10\mu\text{m}}{24000}} 10\text{cm} \approx 10\mu\text{m}$

$$\varepsilon_{Nsc} \sim 12\mu\text{m}, \quad \varepsilon_{N\gamma d} \sim 160\mu\text{m} \quad \text{and} \quad \varepsilon_{Ne-} \sim 40\mu\text{m}$$

- Longitudinal emittance

$$\varepsilon_{\square N} \approx \Delta\gamma_{e+} \sigma_{\tau e-} = \frac{480 - 240}{4} 20\mu\text{m} \approx 1.5\mu\text{m}$$

Conclusion

- Linac-Compton source based on CO₂ laser cavity can be direct upgrade of the non-polarized source. Would require drive linac RF power increase and ~2-3 meters long Compton interaction region with ~10 laser IPs.
- Laser cavity is the only uncertain point of the scheme. Demonstrated at ~10% power on existing hardware. Purchase of adequate amplifier is not funded.
- Pulsed drive linac, conversion and capture components 300Hz for 1 μ s makes it more efficient and very close to conventional technologies.
- High efficiency of the Linac-Compton source (high energy of gammas and clean spectrum) allows for a much simpler design of the target
- Transverse emittance is dominated by scattering in the target. Improvement possible by using shorter target and higher drive beam energy to improve conversion efficiency and reduce scattering.