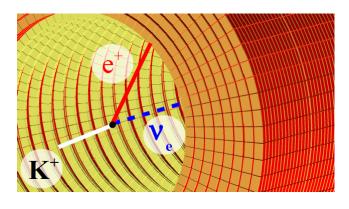


Enhanced neutrino beams from kaon tagging (ENUBET): opportunities and challenges

- Enable the technology for a new generation of neutrino beams with superior control of the ν flux and flavor composition
- The ERC ENUBET Project (2016-2021):

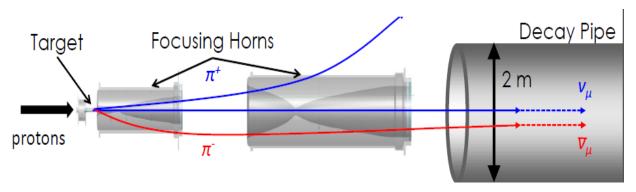
A pure ν_e source monitored by tagging large angle positrons in the decay tunnel originating from kaon decay $K^+ \to \pi^0 \; e^+ \; \nu_e$



- Opportunities for a new generation of short baseline neutrino experiments
- Technical challenges:
 - Proton extraction scheme
 - Focusing system
 - Transfer line and beam dumps
 - Instrumented decay tunnel
- Conclusions



Accelerator neutrino beams



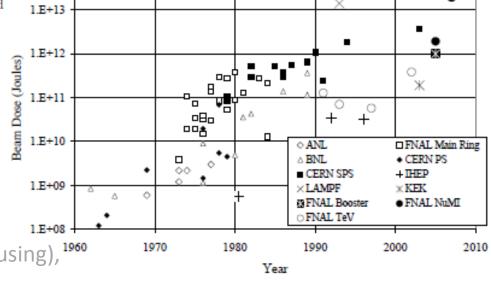
1.E+14

The "van Der Meer's paradigm" (A. Bross)

CERN Informal Workshop on Neutrino Physics, CERN-63-37, (1963).

S. van der Meer, "A directive device for charged particles and its use in an enhanced neutrino beam," CERN-61-07 (1961).

- Fast proton extraction (<10 ms)
- Focus of the largest amount of mesons in the forward direction
- Sign and momentum selection performed only by the horns (*)
- Decay tunnel just after the horns (**)



Notable exceptions: (*) NuTeV (static focusing), 1960 (**) Narrow Band beams (transfer line)

It has been the workhorse of major discoveries in electroweak and neutrino physics because it is scalable to very high beam power (up to MW, right now)

Limitations

In general, accelerator neutrino beams do not employ any of the beam manipulation techniques (FODO lattices, charged particle monitoring etc.) that are common in accelerator physics.

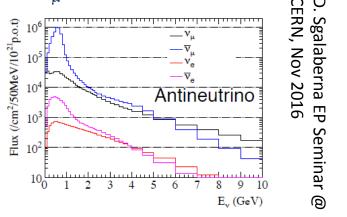
Meson selection performed only by the horns

Decay tunnel just in front of the target

No instrumentation between target and dump

[notable exception: K2K]

Large energy tails and "wrong flavor" (e.g. v_u in anti- v_u beam) contamination



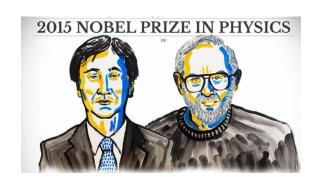
Radiation environment too hostile to directly measure the flux of decayed mesons.

Flux and flavor content estimates must resort to an ab-initio simulation of of the beamline. They are thus sensitive to secondary production yields, proton beam, target horn stability, detailed material budget etc. A tremendous task! (see e.g. L. Aliaga, Phys.Rev. D94 (2016) 092005)

An accelerator neutrino source at the GeV scale is mostly a ν_{μ} source whose flux is known with a precision of >8%

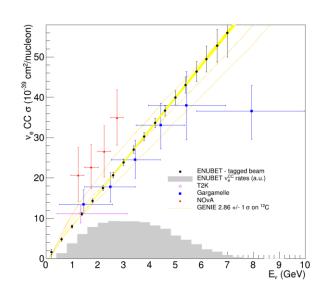
Who cares? No-one... until 2012 ©

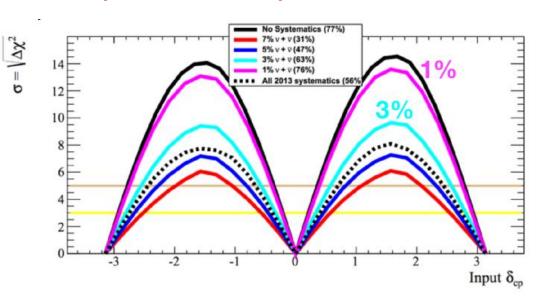
For oscillation studies this can be mitigated by comparison of events at <u>near and far detectors</u> and <u>hadro-production experiments</u>. That was enough in the discovery era of neutrino oscillation physics



The discovery of a large mixing between ν_{μ} and ν_{e} at the atmospheric neutrino scale (large θ_{13}) opened up the precision era of neutrino oscillation physics and boosted a global billion\$ programme in US (DUNE, SBL), Japan (HK, T2K-II), China (JUNO) and Europe (CERN Neutrino Platform).

It comes to no surprise that the physics reach of these facilities is completely dominated by source-related systematics







The ENUBET Project

- ✓ Enable a new technology for accelerator neutrino beams: monitor neutrino production in the decay tunnel at the single particle level
- ✓ In particular, monitor e^+ production in a beam where ν_e are only originated by K^+ → $\pi^0~e^+~\nu_e~\Rightarrow$ a pure ν_e source whose flux is known at the 1% level

ENUBET is a project approved by the European Research Council (ERC) for a 5 year duration (Jun 2016 – May 2021) with an overall budget of 2 Meuro.

Grant: ERC Consolidator Grant, 2015 (PE2)

Principal Investigator: Andrea Longhin

Host Institution: INFN

Collaboration: 41 physicists from 10 institutions





Expression of Interest

Enabling precise measurements of flux in accelerator neutrino beams: the ENUBET project

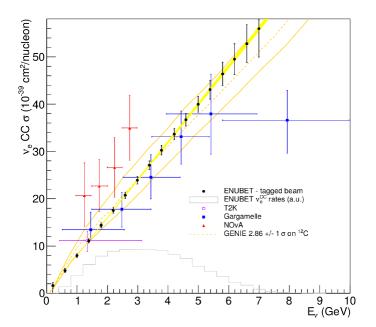
A. Berra^{a,b}, M. Bonesini^b, C. Brizzolari^{a,b}, M. Calviani^m, M.G. Catanesi^l, S. Cecchini^c, F. Cindolo^c, G. Collazuol^{k,j}, E. Conti^j, F. Dal Corso^j, G. De Rosa^{p,q}, A. Gola^o, R.A. Intonti^l, C. Jollet^d, M. Laveder^{k,j}, A. Longhin^{j(*)}, P.F. Loverre^{n,f}, L. Ludovici^f, L. Magaletti^l, G. Mandrioli^c, A. Margotti^c, N. Mauri^c, A. Meregaglia^d, M. Mezzetto^j, M. Nessi^m, A. Paoloni^e, L. Pasqualini^{c,g}, G. Paternoster^o, L. Patrizii^c C. Piemonte^o, M. Pozzato^c, M. Prest^{a,b}, F. Pupilli^e, E. Radicioni^l, C. Riccio^{p,q}, A.C. Ruggeri^p, G. Sirri^c, F. Terranova^{b,h}, E. Vallazzaⁱ, L. Votano^e, E. Wildner^m

See CERN-SPSC-2016-036: SPSC-EOI-014 for details

A very timely R&D

This is **just the right time** to go beyond the standard paradigm (**ENUBET**, **NuSTORM**, **NuPIL** etc.). The ENUBET technology is well suited for short baseline experiment where the intensity requirement are less stringent. There are **three major applications ENUBET can enable**:

- <u>a new generation of cross section experiments</u> operating with a neutrino source that is controlled at the <1% level. This is the main aim of ENUBET as funded by ERC.
- a phase II sterile neutrino search, especially in case of positive signal from the Fermilab short baseline program
- The first step toward a <u>tagged neutrino beam</u> where the $v_{\rm e}$ CC interaction at the detector is time-correlated with the observation of the lepton in the decay tunnel



Impact on v_e cross section measurement assuming the parameters of EPJ C75 (2015) 115 (see below)

A pure v_e source from K decays

Enhanced NeUtrino BEams from kaon Tagging

K⁺ decays

μ⁺ decays in flight

Parent Momentum (GeV/c)

0.025

0.02

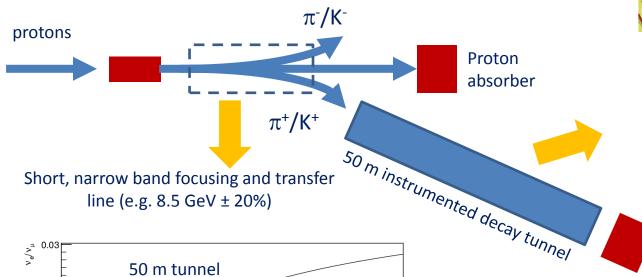
0.015

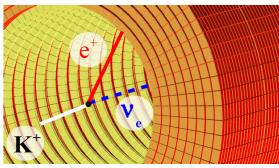
0.01

0.005

100 m tunnel

Inspired by the "tagged neutrino beam" concept [Hand, 1969; Pontecorvo, 1979, Denisov, 1981; Bernstein, 1989; Ludovici, Zucchelli, 1999; Ludovici, Terranova, 2010] and by the design of nuSTORM and nuPIL





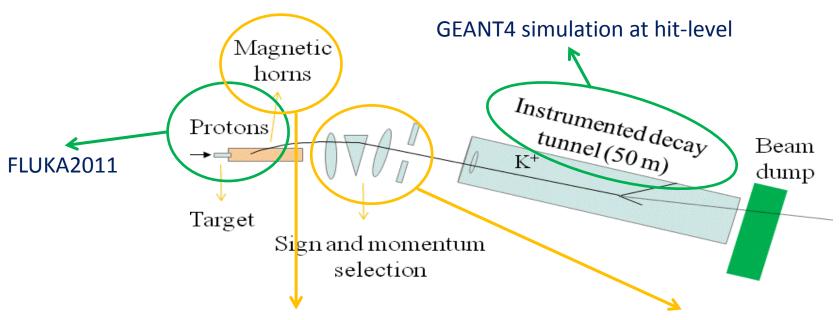
 π^+ and μ decay at small angles and reach the dump without crossing the wall of the tunnel. K decay products cross the instrumented walls and are detected.

Hadron dump

Since the decay tunnel is short and the secondary momentum is 8.5 GeV, 97% of the $\nu_{\rm e}$ are from K decay.

the e^+ rate is a direct measurement of the ν_e flux

Before ENUBET: the SCENTT R&D and the Reference Design

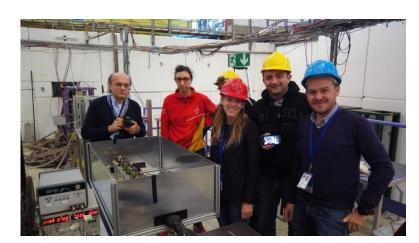


Assuming 85% efficiency for secondaries inside the ellipse $\varepsilon_{xx'}=\varepsilon_{yy'}=0.15$ mm rad in the (x,x',y,y') phase space

Assuming 20% momentum bite at 8.5 GeV and flux reduction due to decay (15 m).

Tunnel instrumentation mostly based on calorimetric techniques (SCENTT: longitudinally segmented shashlik calorimeters)

A. Longhin, L. Ludovici, F. Terranova, EPJ C75 (2015) 155, NIM A824 (2016) 693; A. Berra et al., NIM A830 (2016) 345, A. Berra et al., NuFact15 Proceedings, arXiv:1512.08202 etc.



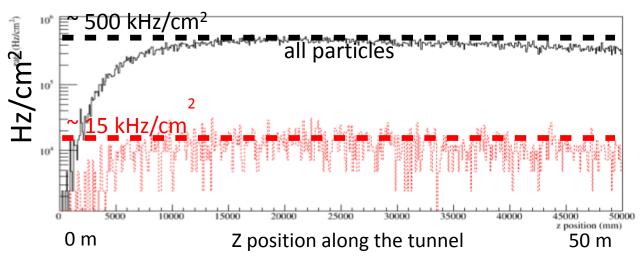
Reference parameters: **500 ton neutrino detector**(*) at 100 m from the entrance of the tunnel. How many protons-on-target are needed to observe 10^4 v_e CC events in the detector (1% statistical uncertainty on cross section)?

(*) e.g. ICARUS@Fermilab, Protodune SP/DP @CERN

	E (GeV)	π^+/PoT	K^+/PoT	PoT for a $10^{10} \pi^+$	PoT for $10^4 \nu_e$ CC
		(10^{-3})	(10^{-3})	spill (10^{12})	(10^{20})
JPARC	30	4.0	0.39	2.5	5.0
Protvino	50	9.0	0.84	1.1	2.4
FIOLVIIIO	60	10.6	0.97	0.94	2.0
	70	12.0	1.10	0.83	1.76
Fermilab	120	16.6	1.69	0.60	1.16
CERN-SPS	450	33.5	3.73	0.30	0.52

For 10^{10} π^+ in a 2 ms spill at the entrance of the tunnel rates are well below 1 MHz/cm²

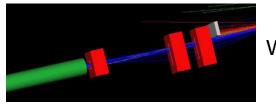
Positron identification based on calorimetric techniques S/N ratio >6



The deliverables of ENUBET

Conceptual design of the beamline

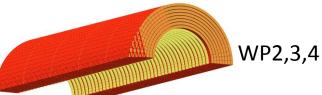




WP1,5

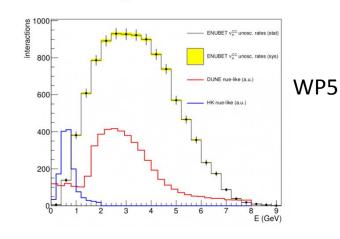
Construction of a 3 m section of the instrumented decay tunnel (~10% of the actual instrumentation)





Assessment of physics potential (systematics reduction)





Tests at PS-T9, LNL

Tests at LNL, LNF

Tests at PS-T9, EHN1

Jun 2016

Prototyping (WG2,3,4) SPS-based Design (WG1,5) **Dec 2018**

Review of the Design LHC LS2

Jun 2021

Demonstrator (WG2,2,4)

Final Design (WG1-5)

Work Packages (WP)

PI A. Longhin



WP1 Conceptual design of the beamline see below

WP2 Design and prototyping of the positron taggers

WP coordinator: M. Pozzato

WP3 SiPM and front-end electronics for the instrumented decay tunnel

WP coordinator: V. Mascagna

WP4 Design and prototyping of the photon veto (e/ γ separation)

WP coordinator: G. Sirri

WP5 Simulation and assessment of the systematics

WP coordinator: A. Meregaglia

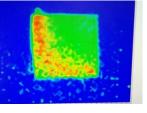




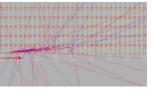




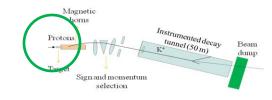








Proton extraction schemes



At SPS, 5 10^{19} pot must be accumulated to perform a v_e cross section measurement at 1% (about 1 year of CNGS) but the extraction scheme must be compatible with the maximum sustainable rates in the tunnel:

Fast extractions (few µs)

[e.g. CNGS]



Not compatible

(rates at GHz/cm²)

Interesting for options like NuPIL(*) (BCT in the transfer line) but, currently, not in the ENUBET scope

WANF-like extractions (few ms)

CERN - ECP / 95-14 31 July 95 **Compatible** with both rates and the use of horns

Baseline option for ENUBET

Slow extractions (>100 ms)

[e.g. SHIP]

Excellent for rates but **not compatible** with horns

Very interesting option for ENUBET if a **static focusing system** can be devised

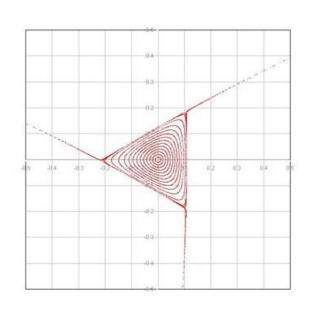
An interesting possibility An interesting possibility Instrumented decay Sign and momentum selection Sign and momentum selection

Deplete the SPS using "slow" (2-10 ms) resonant extractions repeated at multi-Hz rate during the flat top.

Example: 6 ms (WANF) extractions (9 10^{11} pot/6 ms) at 10 Hz provide about 1.8 10^{13} per supercycle for a 2s flat top. It completely depletes the SPS (4 10^{13} pot per supercycle) for a 4.4 s flat top.

Interesting (not only for ENUBET) but non conventional:

- ✓ Can we extract these protons employing a slow (thousand turns) resonant extraction?
- ✓ Can we cross the 3rd order resonance at the sought-for rate (multi-Hz)
- ✓ Are losses tolerable?

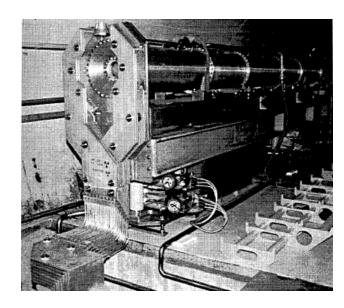


Horn

The main advantage of ms extraction is the possibility to employ magnetic signand momentum selection horns in the beamline as envisaged in projects that employ sign selection and transfer lines similar to the one of ENUBET (NuSTORM and NuPIL).

Again, there are specific issues that have to be investigated:

- ✓ What it is optimal pulse length and rep. rate at the flat top accounting both for the engineering constraints (mostly Joule heating) and the maximum tolerable rate at the instrumented decay tunnel?
- ✓ Compatibility between the proton extraction scheme and the thermo-mechanical properties of the horns

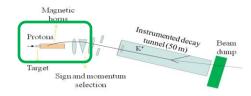


WANF horn (max 0.4 Hz, 120 kA, 6 ms proton pulse)



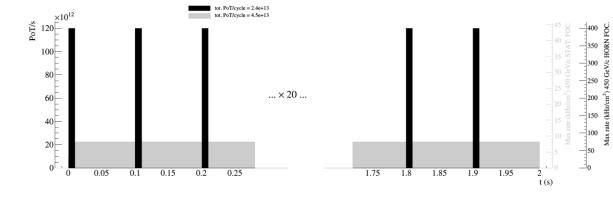
MiniBooNE horn (max 15 Hz, 170 kA, 0.2ms proton pulse) $_{14}$

Slow extraction schemes and static focusing systems

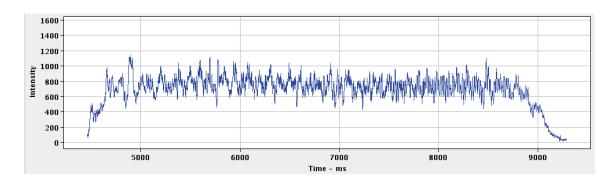


Quad-based focusing systems come to the expense of a strong flux reduction, especially for low energy, high intensity beams. For short-baseline experiments (ENUBET) this is still tolerable (up to a factor 2-3) and bring major advantages:

Decrease by one order of magnitude the particle rate in the instrumented tunnel

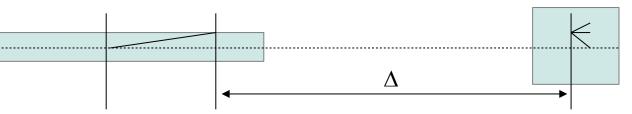


Allows for the implementation of more conventional slow resonant extraction schemes (1-4 s)



Tagged neutrino beams

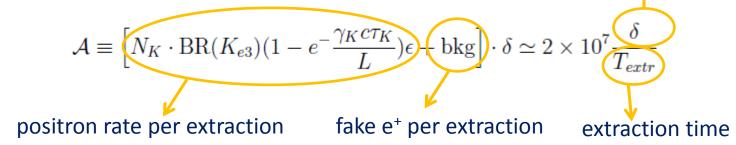
Time coincidence between the v_e CC and the positron: $|\delta t-\Delta/c| < \delta$



 δ t is the difference between the e⁺ and the ν_e CC time (\approx 100 ns).

 δ is the linear sum of the timing resolutions of the e⁺ tagger and neutrino detector

Time correlation between can work if we can beat the number of accidentals:



If $T_{\text{extr}} = 2$ ms we have one positron every 70 ps. δ is unrealistically small. As a consequence, a double tag facility must be operated with **long extraction (1 s)** as it was the "Tagged Neutrino facility" developed in USSR in the 80's.

For
$$T_{extr}$$
 = 1s (1 observed e⁺ every 30 ns) and δ = 1 ns \rightarrow A = 2%

The possibility of using tagged-neutrino beams in high-energy experiments must have occurred to many people. In tagged-neutrino experiments it should be required that the observed event due to the interaction of the neutrino in the neutrino detector would properly coincide in time with the act of neutrino creation $(\pi \to \mu\nu, K \to \mu\nu,$

Tagged beams: the holy grail of v physics

- ✓ Know the flavor of the observed neutrino before it changes due to oscillation.
- ✓ Reconstruct the energy of every single neutrino by the kinematics reconstruction in the decay tunnel

The proton extraction time must be ~1s

Cannot use any more the horns. Must rely on static systems \rightarrow reduction of acceptance

The time resolution of the tag must be <1 ns

OK

The time resolution of the neutrino detector must be ~1 ns

Feasible and strongly synergic with current R&D's

The cosmic background increases by x10 [i.e. by $A \times (1s)/(2 ms)$]

Comparable to what have to be dealt by ICARUS@Fermilab

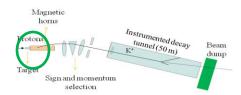
The momentum bite of the K+ must be small enough not to limit the $\nu_{\rm e}$ energy reconstruction

A 20% energy resolution for v on event by event basis would be a major breakthrough

Time synchronization between the tagger and the detector ≪1 ns

OK [direct optical link at short baselines]

Target



Formally, the ENUBET beamline is a O(100 kW) mean power conventional narrow-band neutrino beam.

Maximum conceivable	at SPS
Proton energy:	400 GeV
Nominal flat top:	4.8 s
Repetition rate:	15 s (0.07 Hz)
Protons per cycle:	4.5 10 ¹³
Average beam power	192 kW
Power during flat top	600 kW
Pot/y (200 day livetime)	5.2 10 ¹⁹ pot/y

(physics programme in 1 y of data taking)



CNGS target design average beam power: 750 kW

ENUBET early calculations

Be target, 110 cm length, 3 mm diameter

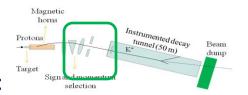
In progress

Graphite target (NuMI-like)

Other options

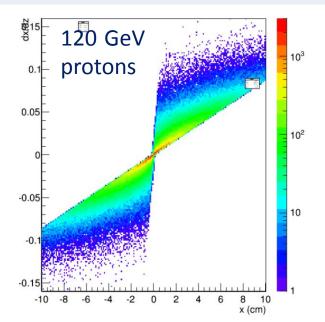
CNGS-like; INCONEL

Transfer line



This is a **key part of the Project** because the transfer line determines:

Source	Expectation	
The K yield at the entrance of the tunnel	20% reduction for 15 m	
The contamination due to neutral and off momentum particles	Bending >5deg	
The quality of the beam at the entrance of the tunnel	10x10 cm ² with angle <3 mrad	
The beam related background	Halo muons (not critical) and debris from collimators	



Acceptance of the transfer line matched with the geometry of the decay tunnel

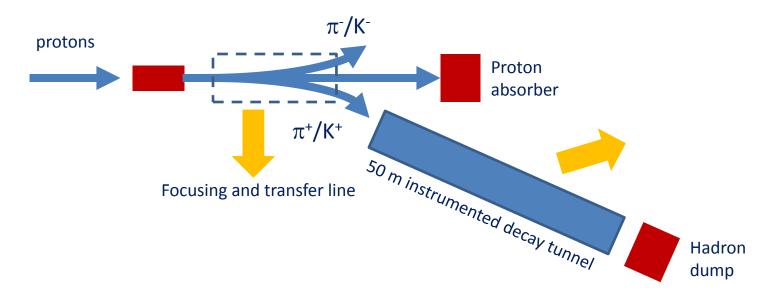
$$A = 4 \times (5 \text{ cm}) \times (3 \text{ mrad}) = 4\epsilon_{xx'} = 4\epsilon_{yy'} = 0.60 \text{ mm rad}$$

$$\epsilon_{xx'} = \epsilon_{yy'} = 0.15 \text{ mm rad}$$

⇒a narrow band beam configuration

Full simulation (G4beamline) is ongoing

Proton absorber and hadron dumps



Both are necessary in a narrow band beam configuration. These are conventional facilities but relevant for ENUBET for the following reasons:

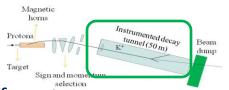
Proton absorber:

Large power (tens of kW)
Shielding and hence, cost assessment of the implementations

Hadron dump:

Low power (<0.5 kW)
Activation and thermal behaviour of the tunnel (relevant for instrumentation)

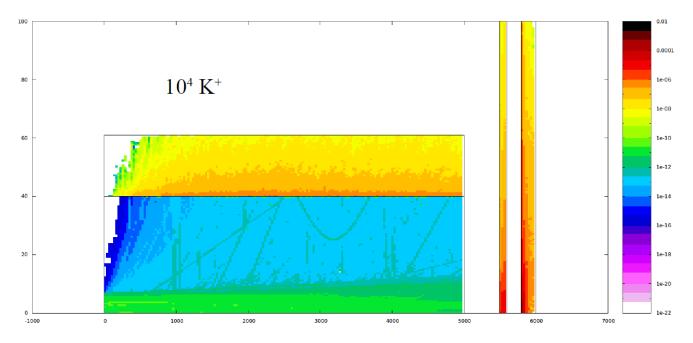
Instrumented decay tunnel



A too harsh environment in conventional wide band beam. Here it is different:

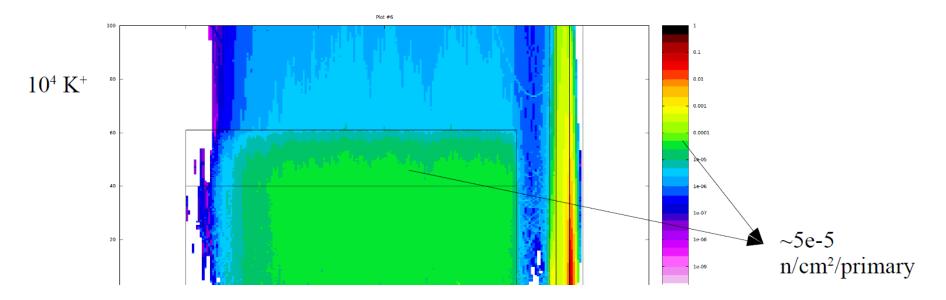
- the tunnel is located after a narrow band transfer line
- the instrumentation is located in the peripheral area of the tunnel

Ionizing and non-ionizin doses evaluated with Fluka2011 assuming instrumentation located at 40 cm from the tunnel axis. Entrance windows 10x10 cm², <1 mbar pressure in the tunnel.



Dose at the tagger at the end of the experiment: 7 kGy (0.7 Mrad) It allows even use of Fe-plastic calorimeters (Compass – SHIP)

Neutron fluence and non-ionizing doses



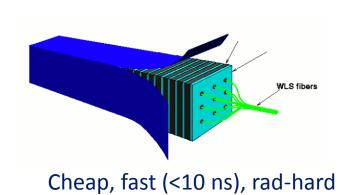
Non ionizing dose at the end of the experiment: $\approx 3 \cdot 10^{12} \cdot 1$ -MeV-equivalent n/cm²

Not critical for sampling calorimeter but significant for solid-state photosensors (e.g. SiPM). Two options:

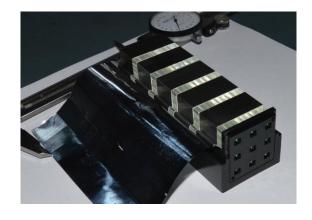
- Scintillation light transported in the proximity of the photosensors with WLS fibers: safe but rather cumbersome
- Use SiPM developed e.g. for CMS hadron calorimeter upgrade (fluence >10¹² n/cm2)

The most elegant solution

Longitudinally segmented shashlik calorimeter with embedded SiPM INFN SCENTT R&D



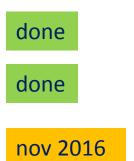


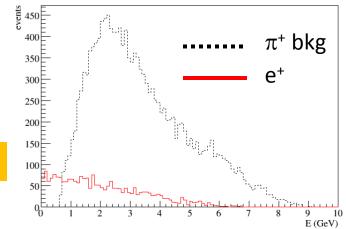


One SiPM for each fiber in the back of each module. Summed signals (9 SiPM per ADC) to reduce cost

Requirements for ENUBET:

- mip sensitivity but no saturation for e.m. showers up to 4 GeV
- energy resolution <25%/E^{1/2}
- recovery time ~10 ns
- validation of MC for e/π separation





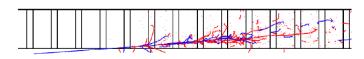
Particle identification in the decay tunnel

Within WP5 activities.

Full GEANT4 simulation of the baseline detector of choice for the instrumented tunnel

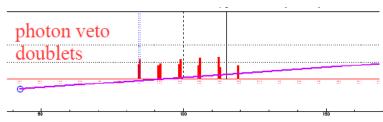
 $e^+/\pi^+/\mu$ separation

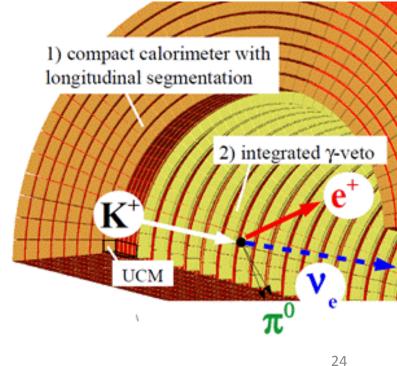
(1) Compact shashlik calorimeter ($3x3x10 \text{ cm}^2 \text{ Fe+scint.}$ modules + energy catcher) with longitudinal (4 X_0) segmentation and SiPM embedded in the bulk of the calorimeter (see below)



 e^+/γ separation

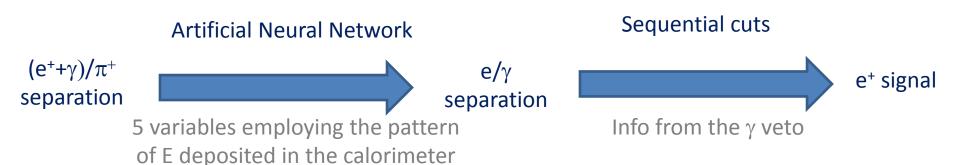
(2) Rings of 3 x 3 cm² pads of plastic scintillator





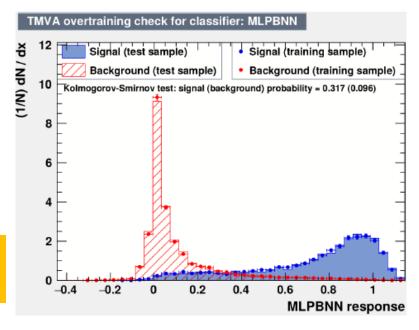
WP5 new results (summer 2016)

The identification algorithms separate positrons from charged and neutral pions combining info from the calorimeter modules and γ veto. Clustering and event building is limited to neighboring modules to avoid pile-up effects and mismatch due to time resolution



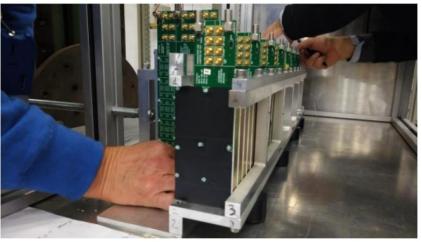
	$\epsilon_{ m geom}$	$\epsilon_{_{\text{sel}}}$
e ⁺	90.7 %	49.0 %
$\pi^{\scriptscriptstyle +}$	85.7 %	2.9 %
π^0	95.1 %	1.2 %

Confirm early results from fast simulation but with a **realistic** and **very cost-effective** setup!

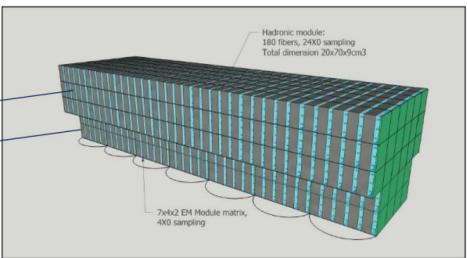


Validation at CERN East Area (T9)

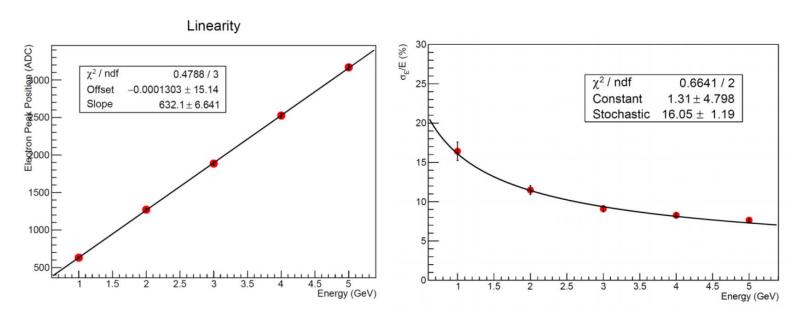






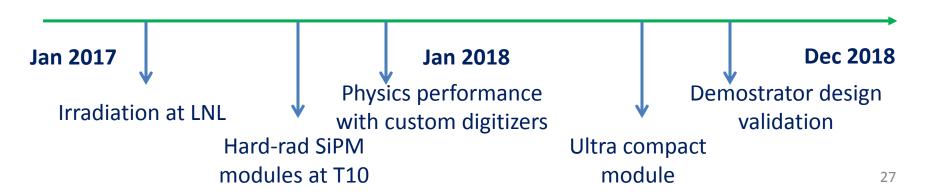


Preliminary results



Data analysis is ongoing (final results in Spring 2017)

Testbeam programme before LS2



Systematics on the flux

Source of uncertainties	Size and mitigation
Statistical error	<1%
kaon production yield	irrelevant (positron tag)
uncertainty on integrated pot	irrelevant (positron tag)
geometrical efficiency	<0.5%
uncertainty on 3-body kinematics and mass	<0.1%
uncertainty on the ν_{e} contam. from μ DIF	<0.5%
secondary beam phase space at the entrance of decay tunnel	can be checked directly with low intensity pion runs
uncertainty on Branching Ratios	irrelevant (positron tag) except for background estimation (<0.1%)
e/π^+ separation and detector stability	can be checked directly at test- beams

The claim of <1% uncertainty is very likely but has to be firmly grounded if ENUBET has to become the standard flux monitoring technique for short baseline neutrino beams.

Conclusions

- The precision era of neutrino oscillation physics requires better control of its artificial sources. At the GeV scale the limited knowledge of the initial flux is **the dominant** contribution to cross section uncertainties
- Such limit can be reduced by one order of magnitude exploiting the K+ \to π^0 e+ ν_e channel (K_{e3})
- In the next 5 years ENUBET will investigate this approach and its application to a new generation of cross section, sterile and time tagged neutrino experiments.
- The technological challenges of this approach are very interesting for accelerator and detector science. For instance:
 - Non conventional extraction schemes (CERN TE-ABT)
 - Focusing, target and beam dump (CERN EN-STI)
 - Transfer Line (CERN EN-EA)
 - Photosensors and instrumentation development (NP3-Plafond)
 - Neutrino physics applications (CERN-EP)
- We hope to enable this technology by 2021 for the next generation of short baseline experiments and – my two cents ☺ - take a step forward beyond the van Der Meer's paradigm.

