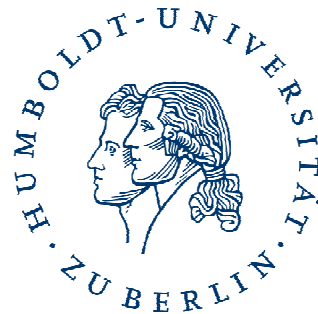


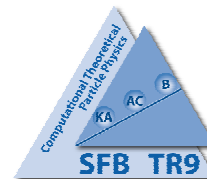
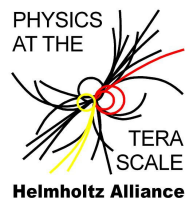
Methodology of Computations in Theoretical Physics

--- Summary ---

Peter Uwer



Tools and Precision Calculations for Physics Discoveries at Colliders





- $5 + 4 + 6 = 15$ presentations
- 450 min = 7.5 h
- In total 367 transparencies, 1.2 min / slide
- Average number of transparencies: 24.4 / talk
- Extreme values: min: 10, max: 56



Where the speakers came from



Europe: 8, Japan: 2, Russia: 2, US: 2



Main topics

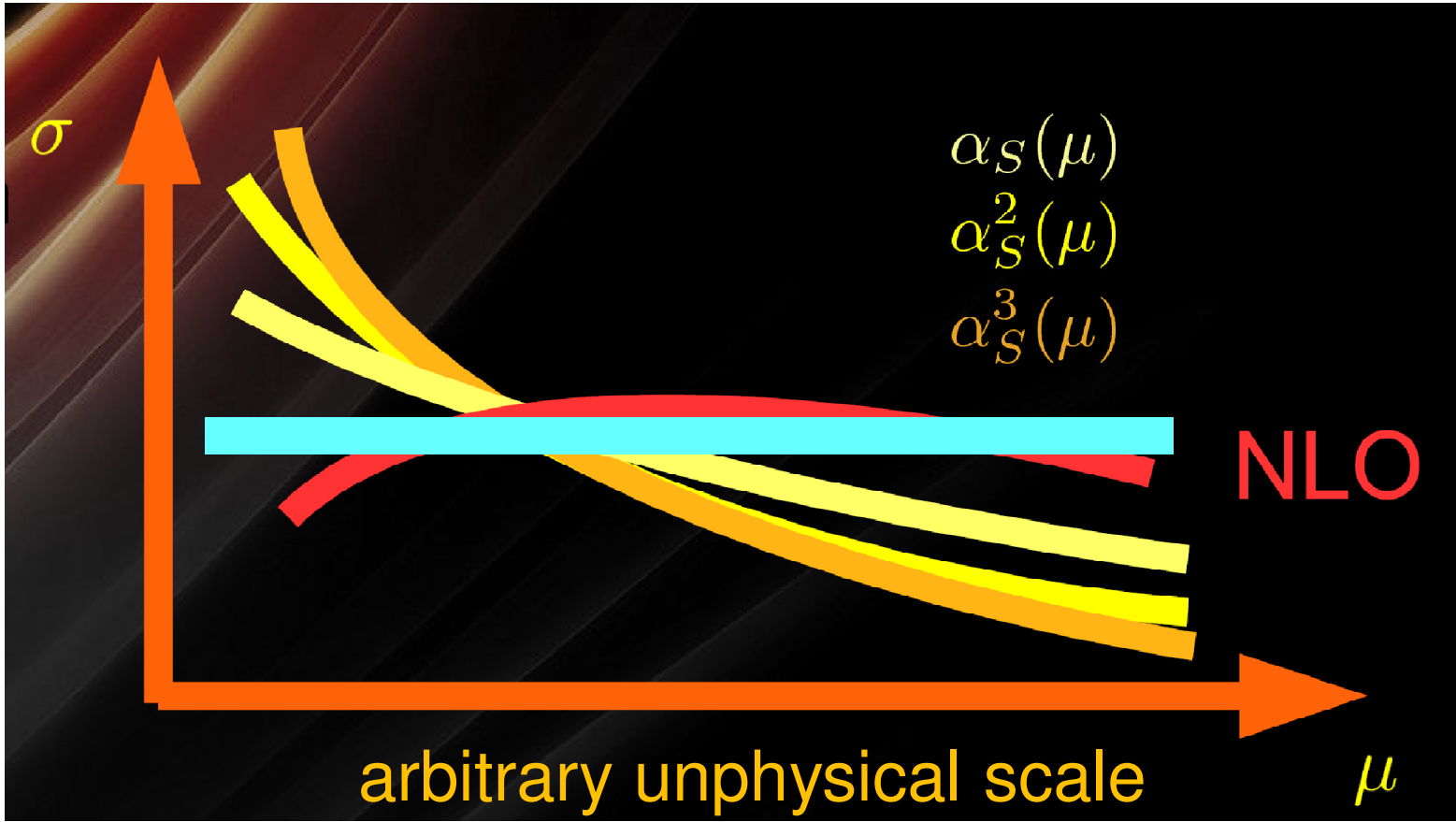


- Automation of higher order corrections
 - Techniques for loop integrals 6
 - Computational aspects 2
 - Real corrections and subtractions 2
- Computer Algebra 2
- Various topics 3



Automation of higher order corrections

- What is the basic problem ?



[Daniel Le Maitre]

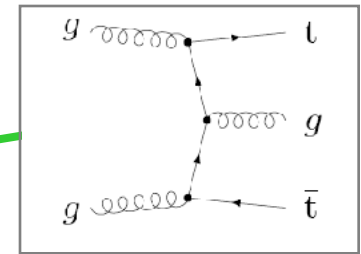
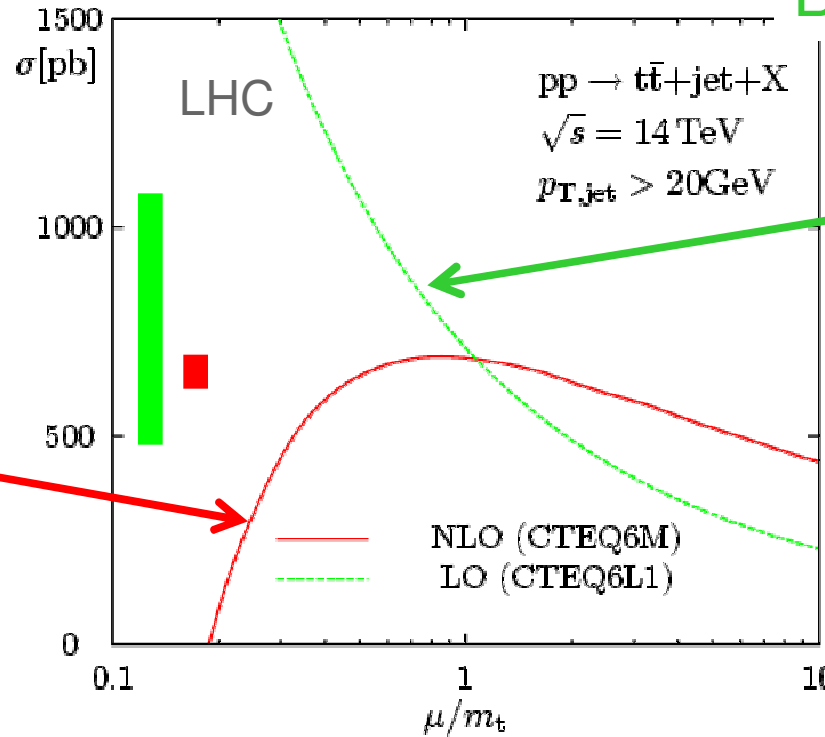


Automation of higher order corrections



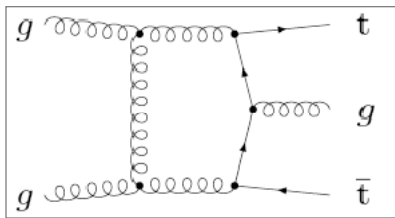
Example: Top-quark production + 1 Jet

Born approximation



O(10) diagrams

Next-to-leading order (NLO)



350 diagrams

→ Born approximation is not reliable, need to go beyond leading-order



Generic one-loop calculation

$$\begin{aligned}
\sigma_{ij} = & \int \left| \text{[n-legs diagram]} \right|^2 && \text{Leading-order, Born approximation} \\
& \underbrace{\int 2\text{Re} \left[\text{[n-legs diagram]} \times \text{[virtual corrections diagram]} \right]}_{\text{IR divergent}} + \int \left| \text{[(n+1)-legs, real corrections diagram]} \right|^2 && \text{Next-to-leading order (NLO)} \\
& \underbrace{\hspace{15em}}_{\text{IR divergent}}
\end{aligned}$$



Bottleneck in one-loop calculation:

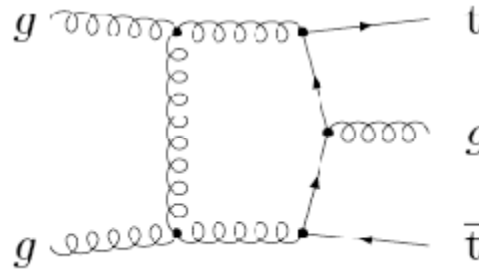
- Calculation of the virtual corrections
 - Many diagrams
 - Each with complicated analytical structure
 - Numerical stability and speed
- Combination of virtual corrections with real ones
 - Cancellation of IR singularities, conceptually solved, but cumbersome if done by hand

Need for new methods and automation

Algorithms crucial for automation



A typical one-loop diagram



$$E_{\mu\nu\alpha\dots} = \int d^d \ell \frac{l_\mu l_\nu l_\alpha l_\beta}{(\ell^2 + i\epsilon)((\ell + p_1)^2 + i\epsilon)((\ell + p_1 - p_t)^2 - m_t^2 + i\epsilon)} \times \frac{1}{((\ell - p_1 + p_{\bar{t}})^2 - m_t^2 + i\epsilon)((\ell - p_2)^2 + i\epsilon)}$$

→ complicated function of many variables

How do we calculate this efficiently?



Different approaches:

- Refinement of mixed approaches
- Improved integration methods
- New algorithms



Techniques for loop integrals



[Tord Riemann]

Basic idea:

Make use of the fact that all the scalar one-loop integrals are known analytically

→ Derive reduction avoiding leading Gram determinants in the denominator

Explicit reduction formulae are implemented in Computer code

Exceptional configuration with vanishing Gram determinants are handled by special reduction (extrapolation)



Techniques for loop integrals



[Tord Riemann]

Using LoopTools call and our math numerics (preliminary):

```
x      D111

 7 :  0.007106204244698895      +0.0046539807850273325 T  D0i [dd111]
     -3.15345811639208  -10      -3.318373348243635  -10 T  Z4d30,Z4d20,T4id20

-6 : -3.2313079078584034-06      -2.8963160014947846-06 I  D0i [dd111]
     3.1479286753545824  10      3.318332145498356  10 T  Z4d30,Z4d20,T4id20

 5 :  5.5231182028025025  09      +3.4832284324178667  09 T  D0i [dd111]
     3.0926394107374516  10      3.3179201270079527  10 T  Z4d30,Z4d20,T4id20

x<  4:      LoopTools dies out.

-4 : -3.1544928789869657-10      -3.33218368329059  -10 I  D0i [dd111]
     -3.0798250216856066-10      -3.3447698103297804-10 I  f1ei

x < -3:      loss of accuracy

-3 : -3.153742175665908  -10      -3.31639655233478  -10 I  D0i [dd111]
     -3.1537481925176414-10      -3.3164147721227693-10 I  f1ei

-2 : -3.1500799889469005-10      -3.29915924109457  -10 I  D0i [dd111]
     -3.150080001830792  -10      -3.2991592067243136-10 I  f1ei

-1 : -3.112267506942415  -10      -3.135823319774082  -10 I  D0i [dd111]
     -3.1122675069507063-10      -3.1358233197649007-10 I  f1ei
```

x allows to test numerical stability



Techniques for loop integrals



[Giovanni Ossola]

General structure of one-loop amplitude:

$$\mathcal{M} = \int d^n \bar{q} A(\bar{q}) = \int d^n \bar{q} \frac{N(\bar{q})}{\bar{D}_0 \bar{D}_1 \dots \bar{D}_{m-1}}$$

$$\begin{aligned} \mathcal{M} &= \sum_i d_i \text{Box}_i + \sum_i c_i \text{Triangle}_i \\ &+ \sum_i b_i \text{Bubble}_i + \sum_i a_i \text{Tadpole}_i + \mathbf{R}, \end{aligned}$$



Automation of higher order corrections

Reduction at the integrand level

[Giovanni Ossola]

$$\begin{aligned} N(q) &= \sum_{i_0 < i_1 < i_2 < i_3}^{m-1} \left[d(i_0 i_1 i_2 i_3) + \tilde{d}(q; i_0 i_1 i_2 i_3) \right] \prod_{i \neq i_0, i_1, i_2, i_3}^{m-1} D_i \\ &+ \sum_{i_0 < i_1 < i_2}^{m-1} \left[c(i_0 i_1 i_2) + \tilde{c}(q; i_0 i_1 i_2) \right] \prod_{i \neq i_0, i_1, i_2}^{m-1} D_i \\ &+ \sum_{i_0 < i_1}^{m-1} \left[b(i_0 i_1) + \tilde{b}(q; i_0 i_1) \right] \prod_{i \neq i_0, i_1}^{m-1} D_i \\ &+ \sum_{i_0}^{m-1} \left[a(i_0) + \tilde{a}(q; i_0) \right] \prod_{i \neq i_0}^{m-1} D_i \end{aligned}$$

- Structures in red vanish after integration and their form is known \rightarrow finite number of terms

\rightarrow Determine coefficients by solving linear system of equations

OPP method (Ossola, Pittau, Papadopoulos)



[Giovanni Ossola]

- OPP method very powerfull
- Available as Fortran program CutTools
- Can be combined with automated amplitude generation (combination with HELAC already done)
- Many new results recently ($pp \rightarrow Wjjj$, $pp \rightarrow ttbb$, $pp \rightarrow ttjj$)

One-loop amplitudes solved ?



Techniques for loop integrals



[Elise de Doncker]

Basic idea:

Recursive (deterministic) integration
over Feynman parameter, combined
with extrapolation

$$\int_0^1 dx \int_0^1 dy \frac{2\alpha y}{(x+y-1)^2 + \alpha^2}, \quad \alpha = 10^{-p}$$

p	DQAGE × DQAGE		DCUHRE	
	ABS. ERR.	# EVAL.	ABS. ERR.	# EVAL.
1	0.00e+00	21255	2.06e-12	144165
2	2.40e-13	93135	5.96e-12	1998675
3	3.49e-13	208035	1.37e-12	21040551
4	1.58e-13	388125	8.04e-12	99999963
5	4.49e-13	561585	4.40e-07	99999963
6	1.69e-09	527205	3.38e-02	99999963
7	1.42e-10	686745	1.99e+00	99999963
8	3.94e-10	902145	3.04e+00	99999963
9	3.20e-08	106965	3.13e+00	99999963
10	4.32e-09	1964385	3.14e+00	99999963
11	1.87e-01	58651365	3.14e+00	99999963

- use DQAGE from QUADPACK recursively



Techniques for loop integrals



[Elise de Doncker]

Six-point scalar integrals are reduced algebraically to
3- and 4-point scalar integrals

3- and 4-point integrals are then evaluated numerically

Technique also applicable to tensor integrals

**Recursive integration might be
interesting also for other fields**



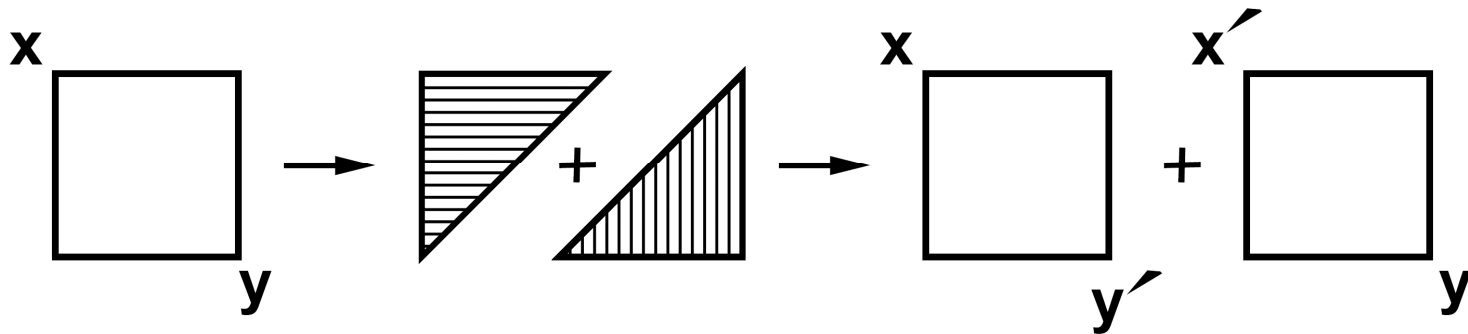
Techniques for loop-integrals



Sector decomposition to isolate singularities [Mikhail Tentyukov]

$$\int_0^1 \frac{dx dy f(x, y)}{(x + y)^{2-\varepsilon}}$$

with $\varepsilon \rightarrow 0$ problems in both $x = 0$ and $y = 0$ *simultaneously*.



$$\int_0^1 dx \int_0^x dy + \int_0^1 dy \int_0^y dx \quad \left(\begin{array}{l} y \rightarrow xy' \\ x \rightarrow x'y \end{array} \right) = \int_0^1 \frac{dx dy' f(x, xy')}{x^{1-\varepsilon} (1+y')^{2-\varepsilon}} + \int_0^1 \frac{dx' dy f(x', y)}{y^{1-\varepsilon} (1+x')^{2-\varepsilon}}$$

Singularities are factorized, can be integrated analitcally

to find the decomposition of a complicated integrand highly non-trivial



Techniques for loop-integrals



FIESTA

[Mikhail Tentyukov]

Feynman I ntegral E valuation by a S ector decomposi T ion A pproach

*Computer algebra part in Mathematica combined with
numerical integration routine*

Important:

- Publicly available
- Different Algorithms for sector decomposition
- Applicable to multi-loop integrals
- Important new 4-loop results
- Circumvent memory problem in Mathematica
- Own interpreter to process formulae of TeraByte length



Techniques for loop-integrals



Alternative algorithm for sector decomposition [Toshiaki Kaneko]

Sector decomposition based on computational geometry

Diagram	A	B	C	S	X	H	This method	Exponential S.D.
Bubble	2	2	2	2*	2		2	2
Triangle	3	3	3	3*	3		3	3
Box	12	12	12	12	12		12	8
Tbubble	58	48	48	48*	48		48	36
Double box, $p_i^2 = 0$	775	586	586	362	293	282	266	106
Double box, $p_4^2 \neq 0$	543*	245*	245*	230*	192*	197	186	100
Double box, $p_i^2 = 0$ nonplanar	1138	698	698	441*	395		360	120
D420	8898	564	564	180	F		168	100
3 loop vertex (A8)	4617*	1196*	1196*	871*	750*	684	684	240
Triple box	M	114256	114256	22657	10155		6568	856

- “H” : G. Heinrich, Int. J. Mod. Phys. A23 (2008) 1457.
- “A”, “B”, “C”, “S”, “X” (without “*”) is cited from Borger & Weinzierl, and Smirnov & Tentyukov.

→ implementation underway



Methods rely on

- Increased computational power
- Increased main memory
- Parallelization is used frequently



[Theodoros Diakonidis]

Apply reduction scheme
presented by Tord Riemann
to $gg \rightarrow ttgg$ @ 1-loop

**$O(1000)$ Feynman diagrams with complex
structure**

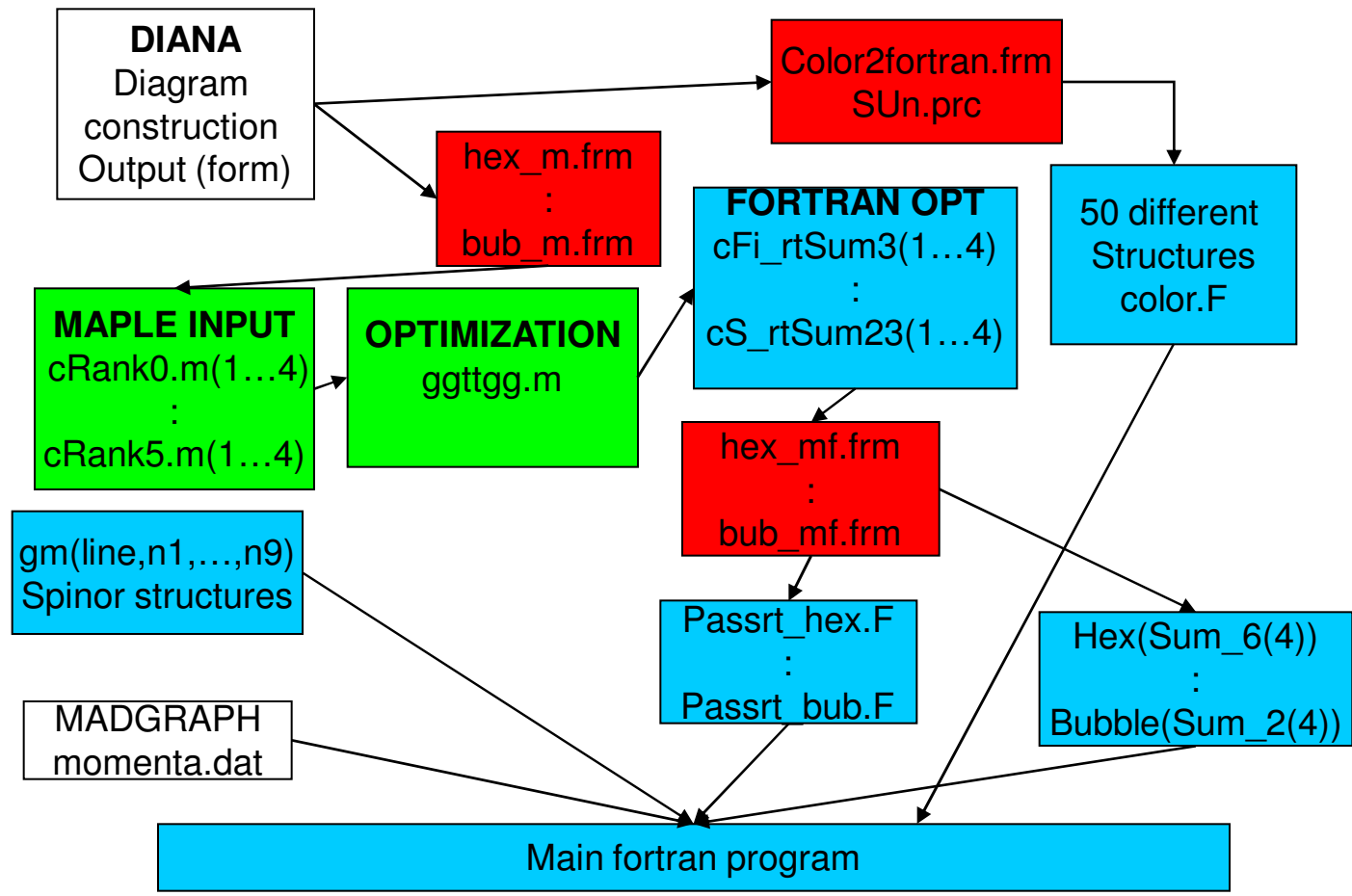
→ Automation needed



Automation of higher order corrections

[Theodoros Diakonidis]

(Qgraf → Form → Maple → Fortran)

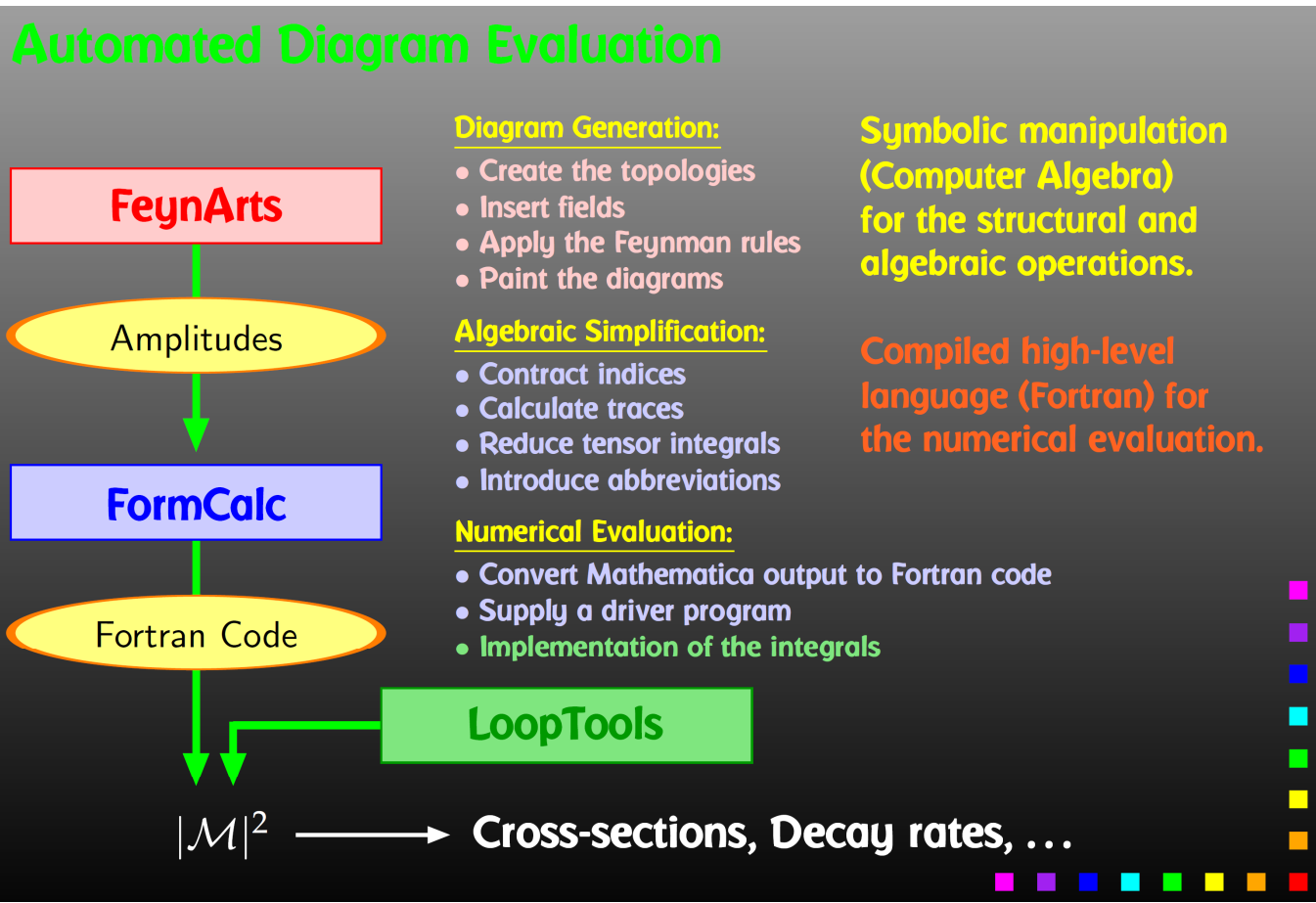


Steering with shell scripts, process specific



Automation of higher order corrections

[Thomas Hahn] Process independent automation based on Feynman diagrams and tensor reduction





[Thomas Hahn]

Many new features in the Feynarts-System

- Tweaking model files
- Diagram selection
- Linear combination of fields (mass vs gauge states)

Computational aspects

- Efficient Fortran code generation, abbreviations to remove common subexpressions
- Parallelization of parameter scans

→ very powerful tool but:

- **On the other hand, you can only go where there are roads. You can't climb a mountain with your car.**



Real corrections and subtractions

- Generic one-loop calculation

$$\begin{aligned}
 \sigma_{ij} = & \int \left| \text{[n-legs diagram]} \right|^2 && \text{Leading-order, Born approximation} \\
 & \underbrace{\int 2\text{Re} \left[\text{[n-legs diagram]} \times \text{[n-legs diagram with loop]} \right]}_{\text{IR divergent}} + \int \left| \text{[(n+1)-legs, real corrections diagram]} \right|^2 && \text{Next-to-leading order (NLO)} \\
 & && \underbrace{\hspace{10em}}_{\text{IR divergent}} \\
 & \text{1/}\epsilon^2, \text{1/}\epsilon \text{ poles} && \text{poles appear after phase} \\
 & \text{appear in scalar loop integrals} && \text{integration in } d \text{ dimensions}
 \end{aligned}$$



Real corrections and subtractions

Problem: Phase space integration cannot be done in d dim.

→ Add and subtract a counterterm which is easy enough to be integrated analytically:

$$\begin{aligned} & \int_0^\alpha dx \frac{1}{x} f(x) x^\epsilon \\ &= \int_0^\alpha dx \frac{1}{x} (f(x) - f(0)) x^\epsilon + \frac{1}{x} f(0) x^\epsilon \\ &= +\frac{1}{\epsilon} \alpha^\epsilon + \int_0^\alpha \frac{1}{x} (f(x) - f(0)) + O(\epsilon) \end{aligned}$$

← Can be done numerically

Construction of subtraction for real corrections more involved,
Fortunately a general solution exists:

→ **Dipole subtraction formalism**



Real corrections and subtractions

$$\sigma_{\text{sub}} = \sum_{\text{dipoles}} \mathcal{D}_{ij,k}(p_i, p_j, p_k)$$

Generic form of individual dipol:

$$\mathcal{D}_{ij,k} = -\frac{1}{(p_i + p_j) - m_{ij}^2} \langle \dots, \tilde{i}j, \dots, \tilde{k}, \dots \left[\frac{\mathbf{T}_a \cdot \mathbf{T}_{ij}}{\mathbf{T}_{ij}} V_{ij,k} \right] \dots, \tilde{i}j, \dots, \tilde{k}, \dots \rangle$$

Leading-order amplitudes
Vector in color space

universal

Color charge operators,
induce color correlation !

Spin dependent part,
induces spin correlation !

Example $gg \rightarrow ttgg$: 36 (singular) dipoles $\mathcal{D}_{g_1 g_3, t, \dots}$

→ Automation required



Automation of NLO subtraction terms

Two different methods:

- Catani-Seymour subtraction
- Frixione-Kunszt-Signer subtraction

→ useful to interface with MC@NLO

→ Fully automated based on Madgraph: MadFKS



[Paolo Bolzoni]

Extension of subtraction method to NNLO

Much more involved due to double unresolved configuration

Analytic integration of subtraction terms highly non trivial

→ Solution using:

- Mellin-Barnes representation
- Special summation algorithms for nested sums (XSummer in Form)



Standard Tool in Theoretical Particle Physics if large expressions are encountered:

Form by Jos Vermaseren et al.

Important features:

- Expression size only limited by disk space (TB)
- Only local operations i.e. no factorization

→ Many ongoing developments

Talks by **Irina Pushkina and Mikhail Tentyukov**



[Irina Pushkina, Mikhail Tentyukov]

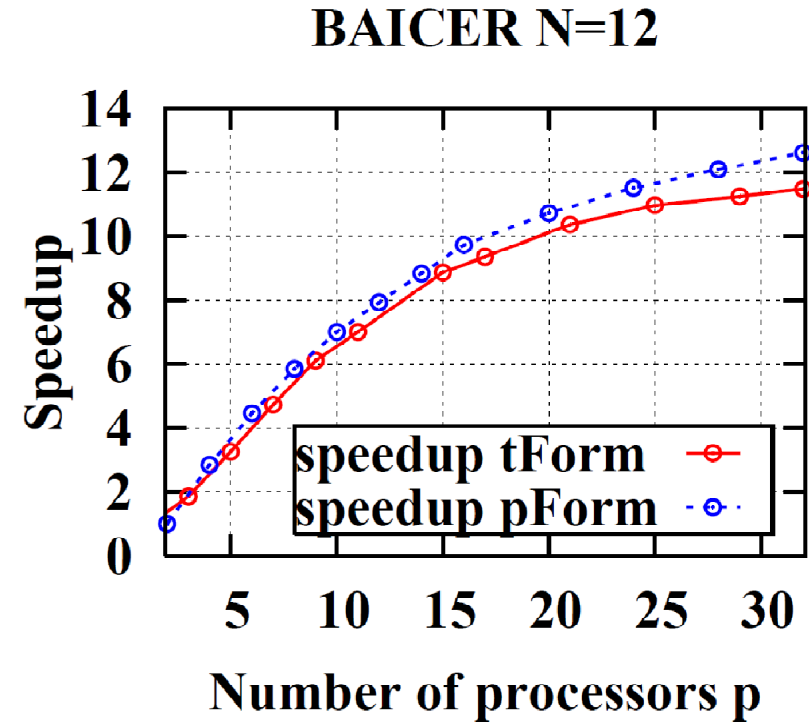
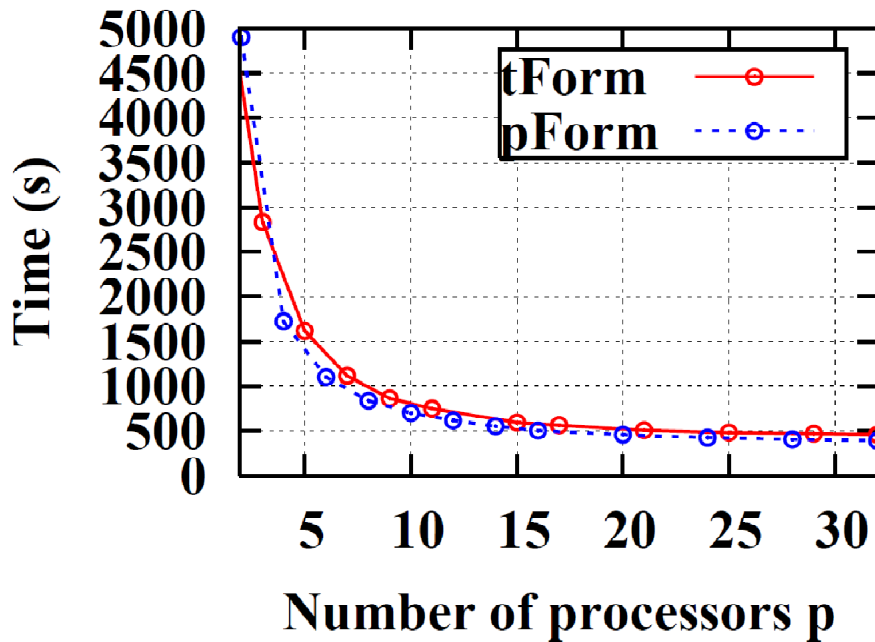
New features:

- Architecture independent file storage (32bit vs 64bit)
- Checkpoints to save intermediate states
- Steps towards open source (summer 2010?)
- Two approaches to parallelisation:
 - Parform based on MPI for cluster
 - Tform based on threads for multi-core machines
- Improved load balancing
- Link to Grace system



Speed up in Form:

[Mikhail Tentyukov]



$$S(p) = \frac{T(1)}{T(p)}$$

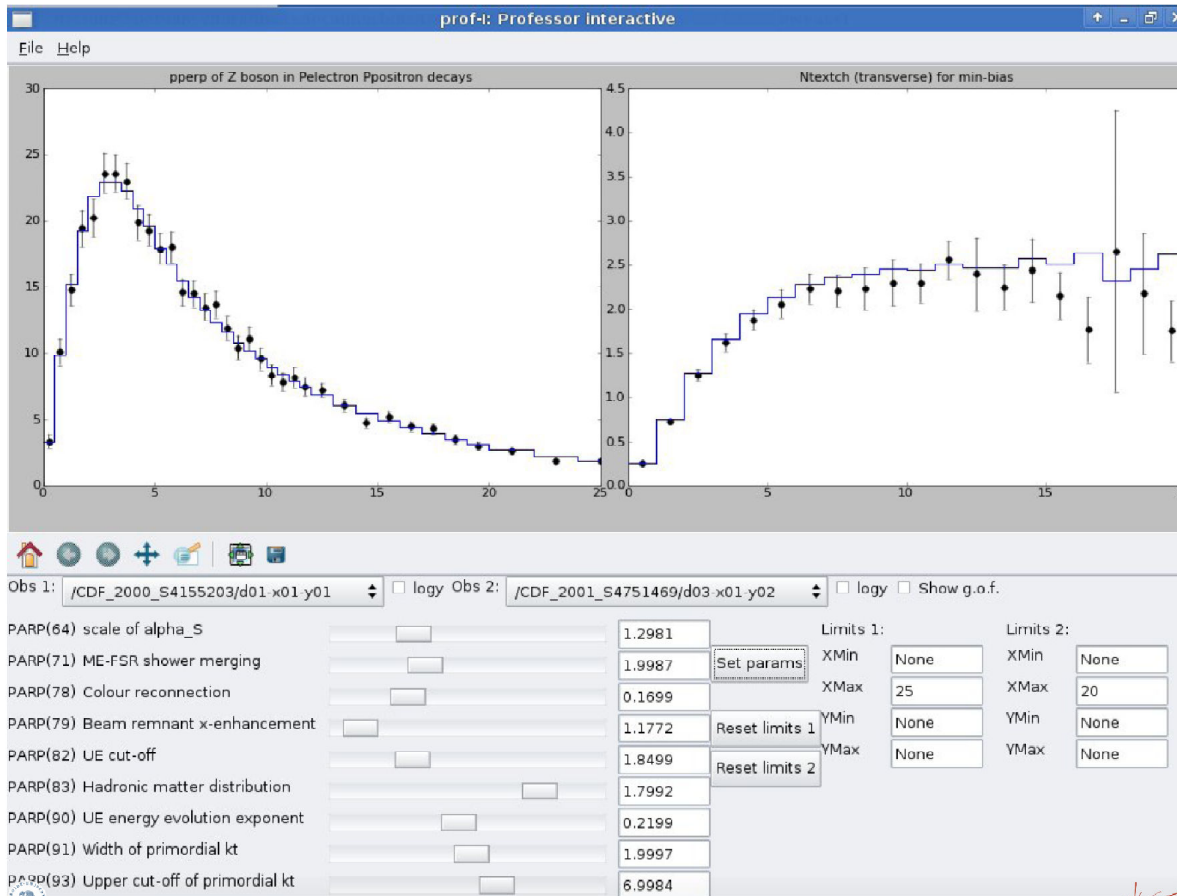


Techniques for Event generator tuning



Tuning framework *Professor*

[James Monk]



Multi-dimensional interpolation

Change Generator Parameters on the fly through interpolation in parameter space



- Main programming language in Theoretical Particle Physics

Fortran

- Powerful mixed approach:

Analytic part is combined with numerical part,
a chain of different tools is connected using scripts

- Numerical instabilities

Switch to quadrupel and higher accuracy

→ Very active field, many new and important developments
recently



*Important progress concerning the automation
of one-loop amplitudes*

→ OPP Method



Final Remarks



- Apologies that not everybody could be mentioned
- Many thanks to all the speakers
- Many thanks to the audience of track 3 for their contribution in many lively discussions

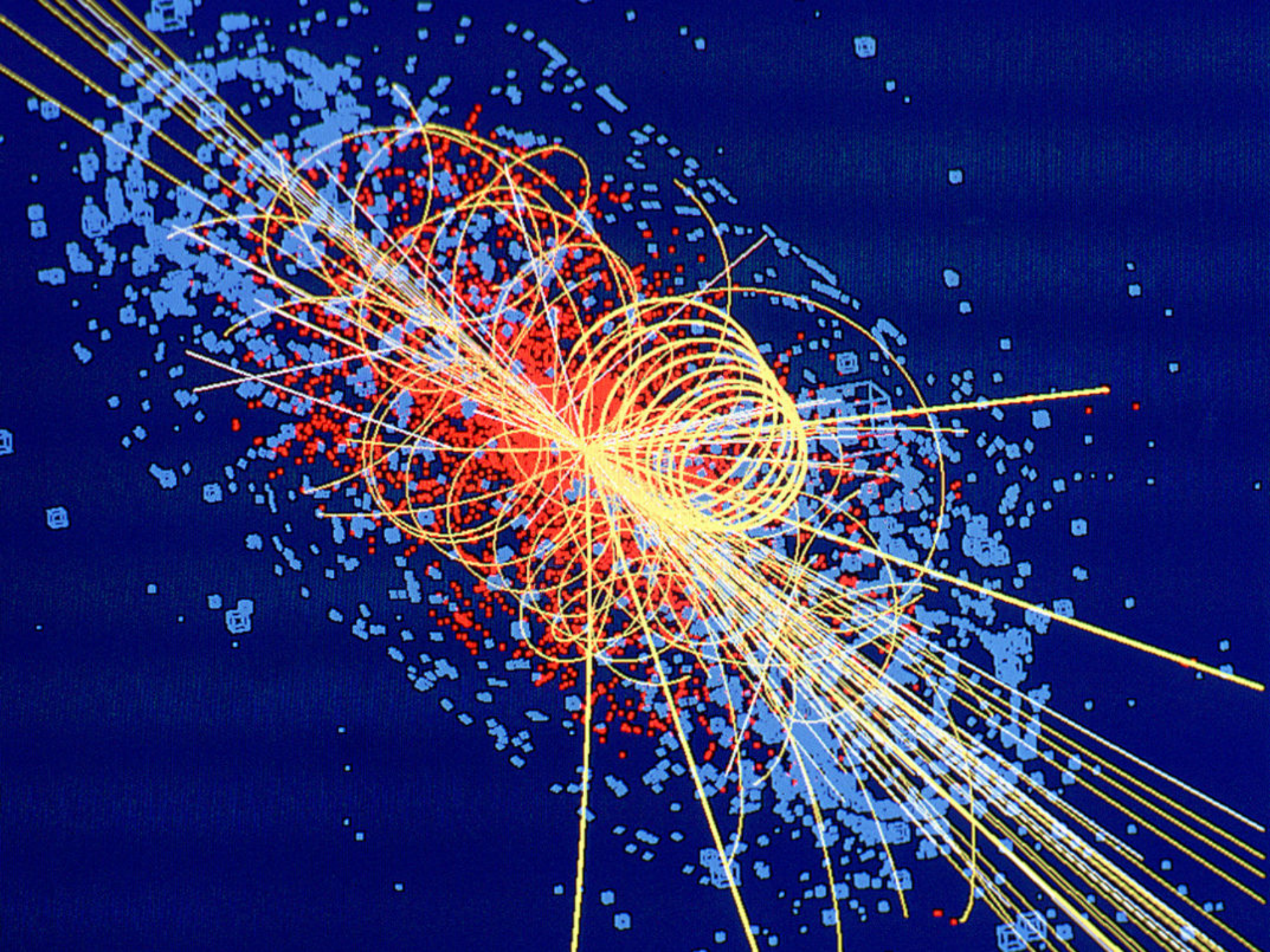
All talks are uploaded, if you want to see the details check Indico



Thank You

Thanks!

- Thanks to IAC for insight
- Sudhir for everlasting help
- All the chairs for creating wonderful workshop atmosphere
- Students etc for flawless technical support
- Speakers for a mosaic of present and future
 - and their help with this summary
- Let's see what the future brings to Brunel!

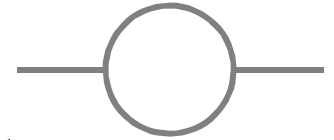




Automation of higher order corrections

Traditional approach to tensor reduction

$$\int d^d \ell \frac{\ell_\mu}{(\ell^2 - m_0^2 + i\epsilon)((\ell + p)^2 - m_1^2 + i\epsilon)} = p_\mu B_1$$



$$B_1 = \frac{1}{2p^2} \left(A(m_0) - A(m_1) + (m_1^2 - m_0^2 - p^2) B_0 \right)$$

$$A(m) = \int d^d \ell \frac{1}{\ell^2 - m^2 + i\epsilon} \quad B_0 = \int d^d \ell \frac{1}{(\ell^2 - m_0^2 + i\epsilon)((\ell + p)^2 - m_1^2 + i\epsilon)}$$

[Passarino, Veltman 78]

Major problem:

(spurious) numerical instabilities for
 exceptional momentum configurations
 vanishing Gram determinants

in the example above: $p^2 \rightarrow 0$ (integral bases degenerates “0/0”)