



Searches for

Jet pT = 104 GeV

at the LHC

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MET = 269 GeV



Why do we keep searching for SUSY?

W

SM is an effective theory. We would like to understand physics in a more generic framework which completes the missing pieces. SUSY unifies gauge couplings at the GUT scale because contributions from new particles modify running of the gauge couplings.

SM does not incorporate gravity. SUSY could.

Fine tuning in the corrections to the Higgs mass can be resolved by adding new particles with different spin. SUSY contributions to Higgs mass cancel SM contributions and stabilize the EW scale. A symmetry called R-parity forces the lightest supersymmetric particle (LSP) to be stable. When LSP is heavy, neutral and stable, it is a good dark matter candidate.

Naturalness drives most SUSY searches

Hierarchy problem: Measured Higgs mass is 125 GeV despite the divergent corrections from the top loop. The divergencies can be cancelled and EW scale can be stabilized by contributions from SUSY particles – but this imposes requirements on the SUSY mass spectrum:

- Leading contribution to the Higgs mass comes from Higgsinos → ≤ few hundred GeV
- Stops contribute to Higgs mass via 1-loop corrections → ≤ few hundred GeV
- Sbottom left can be tied to stop left $\rightarrow \leq$ few hundred GeV.
- Gluinos contribute to Higgs mass via 2-loop corrections → ≤ few TeV
- Rest of the spectrum can be decoupled / heavy.
- But *strained* by the current mass limits.





R.Barbieri & D.Pappadopulo JHEP 0910:061,2009



SUSY(-like) models for interpretation

- Use mainly simplified models. Occasional use of full models like phenomenological MSSM (generic interpretation at the end of a Run).
- A simplified model spectrum (SMS) is defined by a set of hypothetical particles and a sequence of their production and decays.
- Mainly production of 2 particles.
- Each particle decays directly or via a cascade to particles X + a neutral, undetected particle (i.e., neutralino lighterst SUSY particle), or to SM particles (for RPV).
- For each SMS point, experimental acceptance times efficiency (A X e) is calculated.
- From this information, a 95% confidence level upper limit on the product of σ x BR is derived as a function of the particle mass.





A typical SUSY analysis at ATLAS and CM



Run1: 2008-2013, 7TeV, ~5fb⁻¹ + 8 TeV, ~20fb⁻¹ Run2: 2015-2018 13TeV, ~36 fb⁻¹ (data taking continues)

~1000 papers from ATLAS and CMS.

ATLAS and CMS are generic purpose LHC detectors designed to optimally make standard model measurements and new physics searches.





SUSY has O(100) free parameters. It can be realized in diverse ways: different sparticle masses, cross sections, branching ratios...

- In R parity conserving SUSY, sparticles are produced in pairs. Decay chain ends with lightest SUSY particle (LSP).
- Masses up to ~2.1 TeV are accessible at the 13 TeV LHC with ~36 fb⁻¹.



SUSY) [pb] LHC 8TeV 10^{-1}) o(pp-20 10⁻² 10^2 10 # ff NLO(-NLL 10-3 $\widetilde{\chi}^{\dagger}\widetilde{\chi}$ 10^{-4} 1200 1400 1600 200 400 600 800 1000 SUSY sparticle mass [GeV] Heavy sparticles decay to lighter sparticles and SM particles. Can see final states with large number and diversity of particles: multijets, multi-b-jets,-t-quarks, multi-leptons. R parity conserving SUSY: neutralino or gravitino LSP \rightarrow high missing transverse energy (E_T^{miss}). • R parity violating SUSY: decay to SM particles. —> reconstruct resonances, search excesses in SM. Suppress and estimate SM backgrounds with orders of

LPCC SUSY o WG

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dati

magnitude larger cross sections.

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A typical SUSY analysis at ATLAS and CM

- Identify a target signal and a final state of objects
 - target signal: gluinos, stops, EW gauginos, etc.
 - final state: multijets + E_T^{miss}, 2 leptons, 1 lepton + jets, boosted W + jets, etc.
- Design triggers to most efficiently collect data with the target characteristics
- Reconstruct and identify objects (jets, electrons, muons, photons, E^{miss}, ...)
- Find object and event variables to discriminate signal from backgrounds
 - many variables especially for SUSY-like final states: E^{miss}, H_T, m_{eff}, M_{T2}, etc.
- Estimate the irreducible backgrounds in the signal regions
 - use data control regions; MC simulation; functional fits, etc.
- Compute systematic uncertainties
- Results: Compare the estimated background with observed data
- Interpret the result using simplified or full physics models

Results in summary

No

yet

But lots of searches for lots of processes in lots of final states using lots of variables and methods, resulting in lots of limits.



- Inclusive searches are sensitive to:
 - mainly gluinos due to their high production cross sections and diverse decay modes, simple decays or cascades.
 - also squarks and EW gauginos.
- Used multiple objects (jets, b jets, leptons, photons, E_T^{miss}).
- Final states: fully hadronic, single or dileptons, single or diphotons.
- Used kinematic variables designed to discriminate final states with heavy particles and/or E_T^{miss}:
- Used multiple search regions defined by different object multiplicities and kinematic variables.
- Current integrated luminosity allows us to work with disjoint search regions.
 Gluino examples



Kinematic variables for inclusive searches

Missing transverse energy:	$E_T^{miss} = \left -\sum_i^{nvisible} \vec{p}_T^{visible} \right $	for neutral LSP: neutralinos, gravitinos	
Missing hadronic transverse momentum:	$\label{eq:mass_state} {I}_T = H_T^{miss} = -\sum_i^{njets} \vec{p}_T^{jet_i}$		
Hadronic transverse energy	$H_T = \sum_{i}^{n jet} p_T^{jet}$		
Effective mass:	$m_{\rm eff} = E_T^{\rm miss} + H_T$		
Stransverse mass: For \tilde{q} $m_{T2}(m_{\tilde{\chi}}) = \max_{\vec{p}_T^{\tilde{\chi}_1} + \vec{p}_T^{\tilde{\chi}_2}}$	$ \rightarrow X \tilde{\chi}_1^0 $ $ \inf_{\substack{2=\vec{p}_T^{miss}}} \left[\max\left(m_T(\vec{p}_T^{j_1}, \vec{p}_T^{\tilde{\chi}_1}), m_T(\vec{p}_T^{j_2}, \vec{p}_T^{\tilde{\chi}_2}) \right) \right] \le m_{\tilde{q}}^2 $	for high transverse activity	
where the transverse ma	iss is:		
m	$m_{T,W}^2 = m_\ell^2 + m_\nu^2 + 2(p_T^\ell p_T^\nu - \vec{p}_T^\ell \vec{p}_T^\nu)$		

$$(m_\ell, m_\nu \sim 0 \rightarrow) \simeq 2p_T^\ell p_T^\nu (1 - \cos \Delta \phi(\ell, \nu))$$

inclusive

10

For a system with 2 massive particles decaying to 2 massless visible and 2 massive invisible particles. Kinematic variables built by first building a dijet system.

alphaT: For a dijet system with

$$\alpha_T = E_T^{J2} / M_T(jj)$$

where $m_{T}(jj)$ is the dijet transverse mass system.

 $R \equiv M_T^R / M_R$

Razor variables: Estimate

$$|\bar{p}^{q}| = (m_{G}^{2} - m_{\chi}^{2})/2m_{G} = m_{\Delta}/2$$

using longitudional lab fr. observables:

inclusive

$$M_R = \sqrt{\frac{(\vec{p}_z^{q_1} E^{q_2} - \vec{p}_z^{q_2} E^{q_1})^2}{(\vec{p}_z^{q_1} - \vec{p}_z^{q_2})^2 - (E^{q_1} - E^{q_2})^2}} \approx m_\Delta$$

using transverse lab fr. observables:

$$\begin{split} M_R &= \sqrt{\frac{(\vec{p}_z^{q_1} E^{q_2} - \vec{p}_z^{q_2} E^{q_1})^2}{(\vec{p}_z^{q_1} - \vec{p}_z^{q_2})^2 - (E^{q_1} - E^{q_2})^2}} \approx m_\Delta \\ M_T^R &= \sqrt{\frac{E_T^{miss}}{2} (p_T^{q_1} + p_T^{q_2}) - \frac{1}{2} \vec{E}_T^{miss} \cdot (\vec{p}_T^{q_1} + \vec{p}_T^{q_2})} < m_\Delta \end{split}$$





Variables in multijet searches





Multijet search results

CMS SUS-16-036, m_{T2} analysis.





ATLAS-CONF-2017-022







CMS SUS-16-037 $1 lep + jets + E_T^{miss}$ $using M_J = \Sigma_i m(j_i)$. M_J distribution for a signal region.

> ATLAS arXiv:1708.08232 1lep + jets + E^{Tmiss} m^T distribution for the 2jet b veto signal region.





CMS SUS-16-042 $1 lep + jets + E_T^{miss}$. $using m_{T2}$. Data - BG estimate comparisons in 39 search regions defined by n_{jets} , n_b , H_T , $L_T = p_T^{lepton} + E_T^{miss}$ CMS SUS-16-034: Opposite sign dileptons + jets + E_T^{miss}.

Uses dilepton invariant mass $m_{\rm II}$ to search for

inclusive

- a resonance-like excess compatible with the Z boson mass (targets both strongly and electroweakly produced new physics.
- or a kinematic edge (targets strong production)

Kinematic fit to m_{\parallel} to find an edge in the non-resonant search. Signal: triangle convoluted with a Gaussian.



CMS SUS-16-035

Same sign dileptons + jets + E_T^{miss}.

Dominant backgrounds from non-prompt leptons.

Signal regions based on H_T , E_T^{miss} , $m_T^{min}(I_1, I_2)$, n_j , n_b .



Inclusive searches: multileptons

inclusive

Signal region	$N_{\rm leptons}^{\rm signal}$	N _{b-jets}	N _{jets}	$p_{\mathrm{T}}^{\mathrm{jet}}$	$E_{\mathrm{T}}^{\mathrm{miss}}$	m _{eff}	$E_{\rm T}^{\rm miss}/m_{\rm eff}$	Other	
				[GeV]	[GeV]	[GeV]			ATLAS arXiv:
Rpc2L2bS	$\geq 2SS$	≥ 2	≥ 6	> 25	> 200	> 600	> 0.25	_	1706.03731: same
Rpc2L2bH	$\geq 2SS$	≥ 2	≥ 6	> 25	_	> 1800	> 0.15	_	sign dilepton / 3
Rpc2Lsoft1b	$\geq 2SS$	≥ 1	≥ 6	> 25	> 100	_	> 0.3	$20,10 < p_{\rm T}^{\ell_1}, p_{\rm T}^{\ell_2} < 100 {\rm GeV}$	$Ienton + E_{\pm}^{miss}$
Rpc2Lsoft2b	$\geq 2SS$	≥ 2	≥ 6	> 25	> 200	> 600	> 0.25	$20,10 < p_{\rm T}^{\ell_1}, p_{\rm T}^{\ell_2} < 100 {\rm GeV}$	
Rpc2L0bS	$\geq 2SS$	= 0	≥ 6	> 25	> 150	—	> 0.25	-	Gluinos decaying
Rpc2L0bH	$\geq 2SS$	= 0	≥6	> 40	> 250	> 900	_	_	via neutralinos/
Rpc3L0bS	≥ 3	= 0	≥ 4	> 40	> 200	> 600	—	_	charginos
Rpc3L0bH	≥ 3	= 0	≥ 4	> 40	> 200	> 1600	_	_	charginos.
Rpc3L1bS	≥ 3	≥ 1	≥ 4	> 40	> 200	> 600	_	_	Sbottoms.
Rpc3L1bH	≥ 3	≥ 1	≥ 4	> 40	> 200	> 1600	_	_	Signal regions
Rpc2L1bS	$\geq 2SS$	≥ 1	≥6	> 25	> 150	> 600	> 0.25	_	
Rpc2L1bH	$\geq 2SS$	≥ 1	≥ 6	> 25	> 250	_	> 0.2	_	Uses many
Rpc3LSS1b	$\geq \ell^{\pm} \ell^{\pm} \ell^{\pm}$	≥ 1	_		_	_	_	veto $81 < m_{e^{\pm}e^{\pm}} < 101 \text{ GeV}$	discriminating
Rpv2L1bH	$\geq 2SS$	≥ 1	≥ 6	> 50	_	> 2200	_	_	variables.
Rpv2L0b	= 2SS	= 0	≥ 6	> 40	-	> 1800	_	veto $81 < m_{e^{\pm}e^{\pm}} < 101 \text{ GeV}$	
Rpv2L2bH	$\geq 2SS$	≥ 2	≥ 6	> 40	_	> 2000	_	veto $81 < m_{e^{\pm}e^{\pm}} < 101 \text{ GeV}$	No excess.
Rpv2L2bS	$\geq \ell^- \ell^-$	≥ 2	≥ 3	> 50	_	> 1200	_	_	
Rpv2L1bS	$\geq \ell^- \ell^-$	≥ 1	≥ 4	> 50	_	> 1200	_	_	
Rpv2L1bM	$\geq \ell^- \ell^-$	≥ 1	≥ 4	> 50	_	> 1800	_	_	

Similar CMS search: CMS-SUS-16-041, using H_T , E_T^{miss} , m_T , n_b , and invariant mass of the opposite charge same flavor dilepton pairs. No excess observed.

Kinematic variables for inclusive searches

inclusive



- CMS probed gluinos up to mass ~2TeV in the t/b decay channels.
 ATLAS has similar results.
- Having different searches provides sensitivity to different final states and kinematical regions.
- More sensitivity to gluino decays via b quarks due to better measured objects.

Kinematic variables for inclusive searches



inclusive

- ATLAS gluino mass limits for different decay channels.
- Highest gluino mass limits are for light quark decay channels.
- Highest neutralino mass limits are for 3rd gen decay channels.



- Probed 1st/2nd generation squarks up to ~1.6TeV in different decay channels. ATLAS and CMS similar.
- Probed by signal regions with low jet multiplicities, or by reducing the event to a dijet-like state.

Physics on the stop-neutralino plane



3rd gen Hadronic stop searches



 \geq 4 jets + 1b + E_T^{miss}.

discriminating variable.

m_T^{b,min}(b_{1,2},E_T^{miss}) is an effective

CMS-SUS-16-050, Stage Top tagging + ET^{miss}. (13 TeV) CMS 0.9 Combined Simulation Preliminary Monojet 0.8╞ Dijet Top quark tagger efficiency Tagging both boosted tops (with Trijet **0.7**₽ measured in T2tt(850,100) 0.6 05 n-subjettiness) and $\frac{3}{2}$ resolved tops (with 0.4 Top tagging 0.3 efficiencies random forest $\stackrel{\circ}{\vdash}$ 0.2 0 1 discriminator) 100 200 300 500 600 700 800 900 400 p_{τ}^{gen} [GeV] Uses M_{T2}. 35.9 fb⁻¹ (13 TeV) Events 10⁸ ĊMS Rare Observed Z(νν) QCD multijet 107 ---- T2tt(1000,1) Lost lepton Low Δm ---- T2ttC(500, 450 T2tt(700,400) $N_{\rm b} = 1, m_{\rm T}^{\rm b} < 175$ XXX Bkg. uncertaintv 10⁵ $N_{SV} = 0,300 < p_{T}^{ISR} < 500$ $N_{ev} = 0, p$ 10⁴ $p^{b} < 40$ $40 \le p_{-}^{b} < 70$ $p_{-}^{b} < 40$ $40 \le p^{b} < 70$ 10³ 10² 10

 $[300, [400, 500], \hat{6}600 [300, [400, 500], \hat{6}000 [450, 550], \hat{6}500, \hat{5}750 [450, 550], \hat{6}500, \hat{5}750]$

10⁻¹

10⁻²

2

0

N_{obs}/N_{exp}

CMS-SUS-16-049: Top tagging + ET^{miss}. Tagging both resolved and merged tops using an alternative method. Signal regions defined by n_j, n_t, n_W, n_b, n_{resolved top}, M_T^b, ET^{miss}.

[300. 1400.

p_miss [GeV]

1 lepton stop searches



zra ger

ATLAS-CONF-2017-037 1lep + jets + E_T^{miss}. Wide range of stop scenarios with various mass splittings. Used a boosted decision tree as a powerful signal discriminator.

Data / pred

Events

 10^{3}

10²

10



3rd gen 2 leptons stop searches

CMS-SUS-17-001: 2lep + jets + b jet + E_T^{miss} . Longer stop cascades with charginos/sleptons.



ATLAS arXiv:1708.03247

Compressed stop; stop to Z, h



Bra gen Sensitivity to stops



- Huge diversity of searches probing different regions in the stop-neutralino parameter space.
- Stop branching ratios are 100%.
- Dedicated stop searches perform better compared to inclusive.
- Very similar results in CMS.
- Most difficult region lies between m_{stop} - m_{neutralino} = m_W and m_{stop} -m_{neutralino} = m_{top}.
- Exploring top quark properties which would discriminate stop pair production from top pair production further could help to close the gaps.
- Increased mass limits suggest a tension with Naturalness.

3rd gen Direct sbottom searches



3rd gen Sensitivity to sbottoms



- Probed sbottoms up to ~1.2TeV.
- Interplay of direct sbottom searches with the inclusive searches help to obtain the best sensitivity.
 - direct sbottom search reaches higher m_{sbottom}.
 - inclusive M_{T2} search accesses lower sbottom neutralino mass difference.



Searches with taus



ATLAS-CONF-2016-048: τ_{had} + 1lep + jets + E_T^{miss} . stop —> stau —> gravitino.

Variable	SR requirement					
$N_{b ext{-jet}}$	≥ 1					
$E_{\mathrm{T}}^{\mathrm{miss}}$	> GeV180					
$p_{\mathrm{T}}(au)$	> GeV70					
$m_{\mathrm{T2}}(\ell, au)$	> GeV 100					





CMS-SUS-17-003 $2\tau_{had} + E_T^{miss}$.

Discriminate using $M_{T2}(\tau_1, \tau_2)$ and ΣM_T

> Interpret for stau pair production with different stau helicities.



 $\Sigma M_T = M_T(\tau_1, E_T^{miss}) + M_T(\tau_2, E_T^{miss}).$

\mathcal{EWK} EW production of sparticles

- SUSY can be produced via EW interactions, leading to direct production of charginos, neutralinos and sleptons.
 - Cross sections depend on the EW state composition (bino, wino, Higgsino) for charginos / neutralinos, and on chirality for sleptons.
 - EW sector could be the only accessible sector at the LHC if the colored sparticles are above 3-4 TeV.
- Complex decay structures:
 - directly via leptons
 - indirectly via lighter EW gauginos, sleptons/sneutrinos, W/Z/h.
 - Investigated cases with both neutralino and gravitino LSP.
- Searches mainly exploit the multilepton nature of the final states.



\mathcal{EWK} EW production in multilepton searches

CMS-SUS-16-039:

Search for direct EW production of charginos and neutralinos.

Final states with:

- 2 same charge light leptons + E_T^{miss}.
- ≥3 leptons including up to 2 hadronically decaying τ leptons + E_T^{miss}.

Using M_T , $p_T(II)$, M_{II} , E_T^{miss} , $M_{T2}(I_1, I_2)$ as discriminating variables.

No excess.

Extended the probed mass space considerably compared to the previous searches.



$\mathcal{WK} \ \tilde{\chi}^{\pm_1} \ \tilde{\chi}^{0_2}$ in Wh and WZ channels

CMS-SUS-16-043: Search for chargino-neutralino2 production with decays to W(->lv)h(->bb).

 $1 \text{ lep} + 2b + E_T^{\text{miss}}$.

Look for an excess in m_{bb} around the Higgs mass.

CMS-SUS-16-048:

2 opposite charge soft leptons + E_T^{miss}.

Nearly degenerate chargino-neutralino2 production, as in natural compressed higgsino models. Exploit decay to W Z(-> soft II).





Events/30 GeV

\mathcal{EWK} Combining searches with W/Z/h

CMS-SUS-17-004: Combined search for EW production of charginos and neutralinos.

Targets decays with W/Z/h.

Additional optimized analysis with ≥ 3 charged leptons, for models with Δm (neutralino2-neutralino1 ~= mZ.

Searches in combination and target signal topologies

	Signal topology							
Search	WZ	WH	ZZ	ZH	HH			
1ℓ 2b		\checkmark						
4b					\checkmark			
2ℓ on-Z	\checkmark		\checkmark	\checkmark				
2ℓ soft	\checkmark							
2ℓ SS, $\geq 3\ell$	\checkmark	\checkmark	\checkmark	\checkmark	\checkmark			
$\mathrm{H}(\gamma\gamma)$		\checkmark		\checkmark	\checkmark			



\mathcal{EWK} Searches with sleptons and staus

$2\ell + ext{jets signal region definitions}$									
SR2-int SR2-high SR2-low-2J SR2-lo									
$n_{\rm non-b-tagged}$ jets	2	2	2	3-5					
$m_{\ell\ell} [{ m GeV}]$	81	-101	81-101	86-96					
$m_{jj} [{ m GeV}]$	70	-100	70-90	70-90					
$E_{\rm T}^{\rm miss}$ [GeV]	>150	> 250	>100	>100					
$p_{\rm T}^{Z} ~[{ m GeV}]$	>	80	> 60	> 40					
p_{T}^W [GeV]	>	100							
$m_{\mathrm{T2}} \; [\mathrm{GeV}]$	>	100							
$\Delta R_{(jj)}$	<1.5			<2.2					
$\Delta R_{(\ell\ell)}$	<	1.8							
$\Delta \phi_{(\vec{E}_{m}^{miss},Z)}$			< 0.8						
$\Delta \phi_{(\vec{E}_{T}^{miss},W)}$	0.5-3.0 > 1.5		< 2.2						
$E_{\mathrm{T}}^{\mathrm{miss}}/p_{\mathrm{T}}^{Z}$			0.6 - 1.6						
$E_{\mathrm{T}}^{\mathrm{miss}}/p_{\mathrm{T}}^{W}$			< 0.8						
$\Delta \phi_{(\vec{E}_{T}^{miss}, ISR)}$				> 2.4					
$\Delta \phi_{(\vec{E}_{T}^{\mathrm{miss}}, \mathrm{jet1})}$				> 2.6					
$E_{\rm T}^{\rm miss}/{\rm ISR}$				0.4-0.8					
$ \eta(Z) $				< 1.6					
$p_{\mathrm{T}}^{\mathrm{jet3}}$ [GeV]				> 30					

ATLAS-CONF-2017-039: EWK gauginos, sleptons at the 2 or 3 leptons + E_T^{miss} channel.



ATLAS arXiv:1708.07875: Search for charginos/ neutrainos decaying to τs . $2\tau s + E_T^{miss}$.

SR-lowMass	SR-highMass				
At least one opp	osite-sign tau pair				
<i>b</i> -jet veto					
Z-v	eto				
At least two medium tau candidates	at least one medium and one tight tau candidates				
	$m(\tau_1, \tau_2) > 110 \text{ GeV}$				
$m_{\mathrm{T2}} > 70 \ \mathrm{GeV}$	$m_{\mathrm{T2}} > 90 \mathrm{GeV}$				
$Di-tau+E_T^{miss}$ trigger	di-tau+ $E_{\rm T}^{\rm miss}$ trigger	asymmetric di-tau trigger			
$E_{\rm T}^{\rm miss} > 150 { m ~GeV}$	$ \begin{array}{ c c c c c c c c c c c c c c c c c c c$				
$p_{\mathrm{T},\tau_1} > 50 \mathrm{GeV}$					
$p_{\mathrm{T},\tau_2} > 40 \mathrm{GeV}$	$p_{\mathrm{T},\tau_2} > 40 \mathrm{~GeV}$	$p_{\mathrm{T},\tau_2} > 65 \mathrm{~GeV}$			



\mathcal{TWK} Sensitivity to EW gauginos



- Diversity of searches and channels give a complementary picture.
- For decays via W/Z/h bosons, sensitivity is up to ~600 GeV.
- For decays via sleptons, sensitivity is up to ~1.2 TeV. Highest sensitivity is reached in the slepton/sneutrino decays.

Photons + E_T^{miss} searches

Target gauge mediated SUSY breaking (GMSB) models with gravitino LSP. Neutralino next-to-LSP decays to gravitino + γ + others.

hotoms

CMS Preliminary







CMS-SUS-16-046: Isolated γ + E_T^{miss}. **Studied scenarios** where neutralino has bino or wino components. Use E_T^{miss} , S_T^{γ} , $m_T(\gamma)$, Er^{miss})



High H_T , decays with $h \rightarrow bb$





Anomalous h —> $\gamma\gamma$

CMS-SUS-16-045:

Anomalous Higgs production in the $\gamma\gamma$ decay channel in association with >= 1 jet.

Razor variables M_R , R^2 to suppress SM.

Categorize events using M_R , R^2 ; momentum and mass resolution of the $\gamma\gamma$ system.



Extract signal from non-resonant QCD BG through a fit to the yy mass distribution.

Interpret in terms of SMS sbottom —> h b neutralino.



ng- Searches for long lived particles

In many cases, particles may decay non-promptly, outside the beamspot.

- A wide diversity of searches for non-prompt particles complement the prompt searches. Improving and exploring further in Run2.
- Not much SM background outside the beamspot. Mostly detector noise, cosmic rays, reconstruction failure, etc. estimated from data.





ATLAS-CONF-2017-017: Short track + jets + E_T^{miss}.

Long-lived chargino decay, common to wino LSP scenario - ~1/3 of pMSSM phase space.

 $\tilde{\chi}^{\pm}$ NLSP almost degenerate with $\tilde{\chi}^{0}$ LSP $-> \tilde{\chi}^{\pm} -> \tilde{\chi}^{0}\pi^{\pm}$ (soft) $-> \pi^{\pm}$ not reconstructed -> disappearing track.



Disappearing condition: Tracking algorithm with shorter tracks than standard tracks (tracklets). Looking for tracklets with hits only in pixeldetector (pixel tracklets):





Improved sensitivity in the gluino-chargino mass plane compared to inclusive searches, specifically for low gluino-chargino mass difference.

Searches for long lived particles



- Dedicated searches for long-lived gluinos enhance sensitivity significantly as lifetime gets longer.
- For long-lived charginos, sensitivity drops for shorter lifetimes. Using additional SUSY search variables, tracker only triggers or short tracks to probe such signatures.

\mathcal{RPV} Gluinos



\mathcal{RPV} Stops with 2 body decays; neutralinos

ATLAS-CONF-2017-036: RPV stop —> b l.

Upper limits on RPV stops between 600-1100 GeV.





ATLAS-CONF-2016-075: Charginos decaying to RPV neutralinos. 4 leptons final state. Excluded charginos up to mass ~1.1 GeV.



LHC impact on SUSY

ATLAS SUSY Searches* - 95% CL Lower Limits

similar for CMS



ATLAS Preliminary

Μ	ay 2017 Model	e, μ, τ, γ	Jets	E_{T}^{miss}	∫ <i>L dt</i> [fb	⁻¹] Mass limit	$\sqrt{s}=7,$	$\frac{1}{\sqrt{s}} = 13 \text{ TeV}$	$\sqrt{s} = 7, 8, 13 \text{ TeV}$ Reference
Inclusive Searches	$\begin{array}{l} MSUGRA/CMSSM\\ \tilde{q}\tilde{q}, \tilde{q} \rightarrow q\tilde{\chi}_{1}^{0} \\ \tilde{q}\tilde{q}, \tilde{q} \rightarrow q\tilde{\chi}_{1}^{0} \ (\text{compressed}) \\ \tilde{g}\tilde{g}, \tilde{g} \rightarrow q\bar{q}\tilde{\chi}_{1}^{0} \\ \tilde{g}\tilde{g}, \tilde{g} \rightarrow q\bar{q}\tilde{\chi}_{1}^{0} \rightarrow qqW^{\pm}\tilde{\chi}_{1}^{0} \\ \tilde{g}\tilde{g}, \tilde{g} \rightarrow qq\tilde{\chi}_{1}^{\pm} \rightarrow qqW^{\pm}\tilde{\chi}_{1}^{0} \\ \tilde{g}\tilde{g}, \tilde{g} \rightarrow qqQ\tilde{\chi}_{1}^{0} \\ \tilde{g}\tilde{g}, \tilde{g} \rightarrow qqWZ\tilde{\chi}_{1}^{0} \\ GMSB (\tilde{\ell} \ NLSP) \\ GGM \ (\text{bino \ NLSP}) \\ GGM \ (\text{higgsino-bino \ NLSP}) \\ GGM \ (\text{higgsino-bino \ NLSP}) \\ GGM \ (\text{higgsino \ NLSP}) \\ GGM \ (\text{higgsino \ NLSP}) \\ Gravitino \ LSP \end{array}$	$\begin{array}{c} 0-3 \ e, \mu/1-2 \ \tau \\ 0 \\ mono-jet \\ 0 \\ 0 \\ 3 \ e, \mu \\ 0 \\ 1-2 \ \tau + 0-1 \ \ell \\ 2 \ \gamma \\ \gamma \\ 2 \ e, \mu \ (Z) \\ 0 \end{array}$	2-10 jets/3 b 2-6 jets 1-3 jets 2-6 jets 2-6 jets 4 jets 7-11 jets 0-2 jets 1 b 2 jets 2 jets mono-jet	y Yes Yes Yes Yes Yes Yes Yes Yes Yes Yes	20.3 36.1 3.2 36.1 36.1 36.1 36.1 36.1 36.1 3.2 3.2 20.3 13.3 20.3 20.3	q q q q 608 GeV ğ 608 GeV ğ g	1.85 TeV 1.57 TeV 2.02 TeV 2.01 TeV 1.825 TeV 1.8 TeV 2.0 TeV 1.65 TeV 1.37 TeV 1.8 TeV	$\begin{split} & m(\tilde{q}) = m(\tilde{g}) \\ & m(\tilde{\chi}_{1}^{0}) < 200 \text{ GeV}, \ m(1^{\mathrm{st}} \mathrm{gen.} \tilde{q}) = m(2^{\mathrm{nd}} \mathrm{gen.} \tilde{q}) \\ & m(\tilde{q}) - m(\tilde{\chi}_{1}^{0}) < 5 \text{ GeV} \\ & m(\tilde{\chi}_{1}^{0}) < 200 \text{ GeV}, \ m(\tilde{\chi}^{\pm}) = 0.5(m(\tilde{\chi}_{1}^{0}) + m(\tilde{g})) \\ & m(\tilde{\chi}_{1}^{0}) < 200 \text{ GeV}, \ m(\tilde{\chi}^{\pm}) = 0.5(m(\tilde{\chi}_{1}^{0}) + m(\tilde{g})) \\ & m(\tilde{\chi}_{1}^{0}) < 400 \text{ GeV} \\ & c\tau(NLSP) < 0.1 \text{ mm} \\ & m(\tilde{\chi}_{1}^{0}) < 950 \text{ GeV}, \ c\tau(NLSP) < 0.1 \text{ mm}, \ \mu < 0 \\ & m(\tilde{\chi}_{1}^{0}) > 680 \text{ GeV}, \ c\tau(NLSP) < 0.1 \text{ mm}, \ \mu > 0 \\ & m(NLSP) > 430 \text{ GeV} \\ & m(\tilde{G}) > 1.8 \times 10^{-4} \text{ eV}, \ m(\tilde{g}) = m(\tilde{g}) = m(\tilde{g}) = 1.5 \text{ TeV} \end{split}$	1507.05525 ATLAS-CONF-2017-022 1604.07773 ATLAS-CONF-2017-022 ATLAS-CONF-2017-022 ATLAS-CONF-2017-030 ATLAS-CONF-2017-033 1607.05979 1606.09150 1507.05493 ATLAS-CONF-2016-066 1503.03290 1502.01518
§ med.	$egin{array}{l} ilde{g}ar{g}, ilde{g} ightarrow bar{b} ilde{\chi}_1^0 \ ilde{g}ar{g}, ilde{g} ightarrow tar{\lambda}_1^0 \ ilde{g}ar{g}, ilde{g} ightarrow btar{\lambda}_1^0 \ ilde{g}ar{g}, ilde{g} ightarrow btar{\lambda}_1^0 \end{array}$	0 0-1 <i>e</i> ,μ 0-1 <i>e</i> ,μ	3 b 3 b 3 b	Yes Yes Yes	36.1 36.1 20.1	ē ē ē	1.92 TeV 1.97 TeV 1.37 TeV	$\begin{array}{l} m(\tilde{\chi}^0_1) \! < \! 600 \mathrm{GeV} \\ m(\tilde{\chi}^0_1) \! < \! 200 \mathrm{GeV} \\ m(\tilde{\chi}^0_1) \! < \! 300 \mathrm{GeV} \end{array}$	ATLAS-CONF-2017-021 ATLAS-CONF-2017-021 1407.0600
or gen. squarks direct production	$ \begin{split} \tilde{b}_{1}\tilde{b}_{1}, \tilde{b}_{1} \to b\tilde{\chi}_{1}^{0} \\ \tilde{b}_{1}\tilde{b}_{1}, \tilde{b}_{1} \to t\tilde{\chi}_{1}^{\pm} \\ \tilde{i}_{1}\tilde{t}_{1}, \tilde{t}_{1} \to b\tilde{\chi}_{1}^{\pm} \\ \tilde{i}_{1}\tilde{t}_{1}, \tilde{t}_{1} \to b\tilde{\chi}_{1}^{\pm} \\ \tilde{i}_{1}\tilde{t}_{1}, \tilde{t}_{1} \to C\tilde{\chi}_{1}^{0} \\ \tilde{i}_{1}\tilde{t}_{1}, \tilde{t}_{1} \to c\tilde{\chi}_{1}^{0} \\ \tilde{t}_{1}\tilde{t}_{1} \text{ (natural GMSB)} \\ \tilde{t}_{2}\tilde{t}_{2}, \tilde{t}_{2} \to \tilde{t}_{1} + Z \\ \tilde{t}_{2}\tilde{t}_{2}, \tilde{t}_{2} \to \tilde{t}_{1} + h \end{split} $	$\begin{array}{c} 0 \\ 2 \ e, \mu \ (\text{SS}) \\ 0 - 2 \ e, \mu \\ 0 - 2 \ e, \mu \end{array} \\ 0 \\ 2 \ e, \mu \ (Z) \\ 3 \ e, \mu \ (Z) \\ 1 - 2 \ e, \mu \end{array}$	2 b 1 b 1-2 b 0-2 jets/1-2 b mono-jet 1 b 1 b 4 b	Yes Yes Yes Yes Yes Yes Yes Yes	36.1 36.1 9.7/13.3 20.3/36.1 3.2 20.3 36.1 36.1 36.1	$ \begin{array}{c ccccccccccccccccccccccccccccccccccc$	V	$\begin{array}{l} m(\tilde{\chi}_{1}^{0}){<}420GeV \\ m(\tilde{\chi}_{1}^{0}){<}200GeV,m(\tilde{\chi}_{1}^{\pm}){=}m(\tilde{\chi}_{1}^{0}){+}100GeV \\ m(\tilde{\chi}_{1}^{\pm}){=}2m(\tilde{\chi}_{1}^{0}),m(\tilde{\chi}_{1}^{0}){=}55GeV \\ m(\tilde{\chi}_{1}^{0}){=}1GeV \\ m(\tilde{\chi}_{1}^{0}){=}15GeV \\ m(\tilde{\chi}_{1}^{0}){=}150GeV \\ m(\tilde{\chi}_{1}^{0}){=}0GeV \\ m(\tilde{\chi}_{1}^{0}){=}0GeV \\ m(\tilde{\chi}_{1}^{0}){=}0GeV \end{array}$	ATLAS-CONF-2017-038 ATLAS-CONF-2017-030 1209.2102, ATLAS-CONF-2016-077 1506.08616, ATLAS-CONF-2017-020 1604.07773 1403.5222 ATLAS-CONF-2017-019 ATLAS-CONF-2017-019
E W direct	$ \begin{array}{l} \tilde{\ell}_{L,R} \tilde{\ell}_{L,R}, \tilde{\ell} \rightarrow \ell \tilde{\chi}_{1}^{0} \\ \tilde{\chi}_{1}^{+} \tilde{\chi}_{1}^{-}, \tilde{\chi}_{1}^{+} \rightarrow \tilde{\ell} \nu(\ell \tilde{\nu}) \\ \tilde{\chi}_{1}^{+} \tilde{\chi}_{1}^{+} \tilde{\chi}_{2}^{0}, \tilde{\chi}_{1}^{+} \rightarrow \tilde{\tau} \nu(\tau \tilde{\nu}), \tilde{\chi}_{2}^{0} \rightarrow \tilde{\tau} \tau(\nu \tilde{\nu}) \\ \tilde{\chi}_{1}^{+} \tilde{\chi}_{2}^{0} \rightarrow \tilde{\ell}_{L} \nu \tilde{\ell}_{L} \ell(\tilde{\nu}\nu), \ell \tilde{\nu} \tilde{\ell}_{L} \ell(\tilde{\nu}\nu) \\ \tilde{\chi}_{1}^{+} \tilde{\chi}_{2}^{0} \rightarrow W \tilde{\chi}_{1}^{0} h \tilde{\chi}_{1}^{0}, h \rightarrow b \bar{b} / W W / \tau \tau / \gamma \gamma \\ \tilde{\chi}_{2}^{0} \tilde{\chi}_{3}^{0}, \tilde{\chi}_{2,3}^{0} \rightarrow \tilde{\ell}_{R} \ell \\ \text{GGM (wino NLSP) weak prod., } \tilde{\chi}_{1}^{0} \rightarrow \end{array} $	$2 e, \mu$ $2 e, \mu$ 2τ $3 e, \mu$ $2-3 e, \mu$ e, μ, γ $4 e, \mu$ $\gamma \tilde{G} 1 e, \mu + \gamma$ $\gamma \tilde{G} 2 \gamma$	0 0 	Yes Yes Yes Yes Yes Yes Yes Yes	36.1 36.1 36.1 36.1 20.3 20.3 20.3 20.3	$\begin{array}{cccccccccccccccccccccccccccccccccccc$.16 TeV $m(\tilde{\chi}_{1}^{\pm}) = m(\tilde{\chi}_{2}^{0}) =$	$\begin{split} & m(\tilde{\chi}_{1}^{0}) {=} 0 \\ & m(\tilde{\chi}_{1}^{0}) {=} 0, m(\tilde{\ell}, \tilde{\nu}) {=} 0.5(m(\tilde{\chi}_{1}^{\pm}) {+} m(\tilde{\chi}_{1}^{0})) \\ & m(\tilde{\chi}_{1}^{0}) {=} 0, m(\tilde{\tau}, \tilde{\nu}) {=} 0.5(m(\tilde{\chi}_{1}^{\pm}) {+} m(\tilde{\chi}_{1}^{0})) \\ & m(\tilde{\chi}_{2}^{0}), m(\tilde{\chi}_{1}^{0}) {=} 0, m(\tilde{\ell}, \tilde{\nu}) {=} 0.5(m(\tilde{\chi}_{1}^{\pm}) {+} m(\tilde{\chi}_{1}^{0})) \\ & m(\tilde{\chi}_{1}^{\pm}) {=} m(\tilde{\chi}_{2}^{0}), m(\tilde{\chi}_{1}^{0}) {=} 0, \tilde{\ell} \text{ decoupled} \\ & m(\tilde{\chi}_{1}^{\pm}) {=} m(\tilde{\chi}_{2}^{0}), m(\tilde{\chi}_{1}^{0}) {=} 0, \tilde{\ell} \text{ decoupled} \\ & m(\tilde{\chi}_{3}^{0}), m(\tilde{\chi}_{1}^{0}) {=} 0, m(\tilde{\ell}, \tilde{\nu}) {=} 0.5(m(\tilde{\chi}_{2}^{0}) {+} m(\tilde{\chi}_{1}^{0})) \\ & c\tau {<} 1 mm \\ c\tau {<} 1 mm \end{split}$	ATLAS-CONF-2017-039 ATLAS-CONF-2017-039 ATLAS-CONF-2017-035 ATLAS-CONF-2017-039 ATLAS-CONF-2017-039 1501.07110 1405.5086 1507.05493 1507.05493
particles	Direct $\tilde{\chi}_{1}^{\dagger}\tilde{\chi}_{1}^{-}$ prod., long-lived $\tilde{\chi}_{1}^{\pm}$ Direct $\tilde{\chi}_{1}^{\dagger}\tilde{\chi}_{1}^{-}$ prod., long-lived $\tilde{\chi}_{1}^{\pm}$ Stable, stopped \tilde{g} R-hadron Stable \tilde{g} R-hadron Metastable \tilde{g} R-hadron GMSB, stable $\tilde{\tau}, \tilde{\chi}_{1}^{0} \rightarrow \tilde{\tau}(\tilde{e}, \tilde{\mu}) + \tau(e, \mu)$ GMSB, $\tilde{\chi}_{1}^{0} \rightarrow \gamma \tilde{G}$, long-lived $\tilde{\chi}_{1}^{0}$ $\tilde{g}\tilde{g}, \tilde{\chi}_{1}^{0} \rightarrow eev/e\mu\nu/\mu\mu\nu$ GGM $\tilde{g}\tilde{g}, \tilde{\chi}_{1}^{0} \rightarrow Z\tilde{G}$	Disapp. trk dE/dx trk 0 trk dE/dx trk $1-2 \mu$ 2γ displ. $ee/e\mu/\mu$ displ. vtx + jet	1 jet - 1-5 jets - - - μ - ts -	Yes Yes - - Yes - Yes -	36.1 18.4 27.9 3.2 3.2 19.1 20.3 20.3 20.3	$\begin{array}{c ccccccccccccccccccccccccccccccccccc$	1.58 TeV 1.57 TeV	$\begin{split} & m(\tilde{\chi}_{1}^{\pm})\text{-}m(\tilde{\chi}_{1}^{0})\text{-}160 \; MeV, \; \tau(\tilde{\chi}_{1}^{\pm})\text{=}0.2 \; ns \\ & m(\tilde{\chi}_{1}^{\pm})\text{-}m(\tilde{\chi}_{1}^{0})\text{-}160 \; MeV, \; \tau(\tilde{\chi}_{1}^{\pm})\text{<}15 \; ns \\ & m(\tilde{\chi}_{1}^{0})\text{=}100 \; GeV, \; 10 \; \mu\text{s}\text{<}\tau(\tilde{g})\text{<}1000 \; s \\ & m(\tilde{\chi}_{1}^{0})\text{=}100 \; GeV, \; \tau\text{>}10 \; ns \\ & 10\text{<}tan\beta\text{<}50 \\ & 1\text{<}\tau(\tilde{\chi}_{1}^{0})\text{<}3 \; ns, \; SPS8 \; model \\ & 7 \; < \! c\tau(\tilde{\chi}_{1}^{0})\text{<}740 \; mm, \; m(\tilde{g})\text{=}1.3 \; TeV \\ & 6 \; < \! c\tau(\tilde{\chi}_{1}^{0})\text{<}480 \; mm, \; m(\tilde{g})\text{=}1.1 \; TeV \end{split}$	ATLAS-CONF-2017-017 1506.05332 1310.6584 1606.05129 1604.04520 1411.6795 1409.5542 1504.05162 1504.05162
RPV	$ \begin{array}{l} LFV \ pp \rightarrow \tilde{v}_{\tau} + X, \tilde{v}_{\tau} \rightarrow e\mu/e\tau/\mu\tau \\ Bilinear \ RPV \ CMSSM \\ \tilde{\chi}_{1}^{+}\tilde{\chi}_{1}^{-}, \tilde{\chi}_{1}^{+} \rightarrow W\tilde{\chi}_{1}^{0}, \tilde{\chi}_{1}^{0} \rightarrow eev, e\mu v, \mu\mu v \\ \tilde{\chi}_{1}^{+}\tilde{\chi}_{1}^{-}, \tilde{\chi}_{1}^{+} \rightarrow W\tilde{\chi}_{1}^{0}, \tilde{\chi}_{1}^{0} \rightarrow \tau\tau v_{e}, e\tau v_{\tau} \\ \tilde{g}\tilde{g}, \tilde{g} \rightarrow qqq \\ \tilde{g}\tilde{g}, \tilde{g} \rightarrow qq\tilde{\chi}_{1}^{0}, \tilde{\chi}_{1}^{0} \rightarrow qqq \\ \tilde{g}\tilde{g}, \tilde{g} \rightarrow t\tilde{\chi}_{1}^{0}, \tilde{\chi}_{1}^{0} \rightarrow qqq \\ \tilde{g}\tilde{g}, \tilde{g} \rightarrow \tilde{t}_{1}t, \tilde{t}_{1} \rightarrow bs \\ \tilde{t}_{1}\tilde{t}_{1}, \tilde{t}_{1} \rightarrow b\ell \end{array} $	$e\mu,e\tau,\mu\tau$ 2 e, μ (SS) 4 e, μ 3 e, μ + τ 0 4- 0 4- 1 e, μ 8 1 e, μ 8 0 2 e, μ	- 0-3 <i>b</i> - 5 large- <i>R</i> je -5 large- <i>R</i> je -10 jets/0-4 -10 jets/0-4 2 jets + 2 <i>b</i> 2 <i>b</i>	- Yes Yes tts - tts - b - b - -	3.2 20.3 13.3 20.3 14.8 14.8 36.1 36.1 15.4 36.1	$ \begin{array}{c} \tilde{v}_{\tau} \\ \tilde{q}, \tilde{g} \\ \tilde{\chi}_{1}^{\pm} \\ \tilde{\chi}_{1}^{\pm} \\ \tilde{\chi}_{1}^{\pm} \\ \tilde{g} \\ \tilde{g} \\ \tilde{g} \\ \tilde{g} \\ \tilde{g} \\ \tilde{g} \\ \tilde{i}_{1} \\ \tilde{i}_{1} \\ \end{array} $	1.9 TeV 1.45 TeV 14 TeV 8 TeV 1.55 TeV 2.1 TeV 1.65 TeV 0.4-1.45 TeV	$\begin{split} \lambda'_{311} = & 0.11, \ \lambda_{132/133/233} = & 0.07 \\ m(\tilde{q}) = m(\tilde{g}), \ c\tau_{LSP} < 1 \ mm \\ m(\tilde{\chi}_1^0) > & 400 \text{GeV}, \ \lambda_{12k} \neq 0 \ (k = 1, 2) \\ m(\tilde{\chi}_1^0) > & 0.2 \times m(\tilde{\chi}_1^{\pm}), \ \lambda_{133} \neq 0 \\ \text{BR}(t) = & \text{BR}(b) = & \text{BR}(c) = & 0\% \\ m(\tilde{\chi}_1^0) = & 800 \ \text{GeV} \\ \forall \ m(\tilde{\chi}_1^0) = & 1 \ \text{TeV}, \ \lambda_{112} \neq 0 \\ m(\tilde{\tau}_1) = & 1 \ \text{TeV}, \ \lambda_{323} \neq 0 \\ \text{BR}(\tilde{\tau}_1 \rightarrow be/\mu) > & 20\% \end{split}$	1607.08079 1404.2500 ATLAS-CONF-2016-075 1405.5086 ATLAS-CONF-2016-057 ATLAS-CONF-2016-057 ATLAS-CONF-2017-013 ATLAS-CONF-2017-013 ATLAS-CONF-2017-013 ATLAS-CONF-2017-036
ther	Scalar charm, $\tilde{c} \rightarrow c \tilde{\chi}_1^0$	0	2 c	Yes	20.3	č 510 GeV		$m(\tilde{\chi}_1^0)$ <200 GeV	1501.01325
Only only	a selection of the available ma	iss limits on r limits are bay	new states	s or	1) ⁻¹	1	Mass scale [TeV]	

phenomena is shown. Many of the limits are based on simplified models, c.f. refs. for the assumptions made.

Mass scale [TeV]



To conclude....

We found no significant hint for SUSY after 7 years at the LHC...

- ~36 fb⁻¹ of 13 TeV data
- ~60 searches at 13 TeV in ATLAS and CMS in a diversity of final states involving jets, b-jets, leptons, photons, ET^{miss}. Special searches for longlived particles and R parity violation.
- Not all results were included in this talk. Full list of results here:
 - ATLAS: https://twiki.cern.ch/twiki/bin/view/AtlasPublic/ SupersymmetryPublicResults
 - CMS: https://twiki.cern.ch/twiki/bin/view/CMSPublic/PhysicsResultsSUS
- Further constraints from direct SUSY Higgs sector searches.
- Even further constraints from 125GeV Higgs, w/Z, top measurements.
- Many other searches being performed with upcoming data
- Designing refined analysis methods and tools to improve sensitivity for probing SUSY most efficiently.

Perhaps in the next Runs? The journey looks as delightful as ever :)