## FIRST SEARCH FOR BOOSTER HIGGS→BB WITH GMS BOOSTED GGF H(BB)

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LHC HIGGS XS WG WORKSHOP CERN GENEVA, SWITZERLAND GENEVA, 2017

Javier Duarte  $\blacktriangleright$ Fermilab



## OUTLINE

 $q/g$ 

- **Motivation**
- Experimental techniques
	- Jet substructure and grooming
	- Double-b-tagging  $\mathbf{g}$
- Event selection
- Data-driven QCD estimation
- Higgs  $p_T$  modeling
- **Results**
- **BKG: QCD** *[PUPPI'ED INPUTS]*  • Summary and outlook







#### MOTIVATION  $\frac{1}{\epsilon}, \mathbf{t}, \mathbf{0}$  $t, \widetilde{t}, X?$  $\overline{V}$ involves in the contract of th twoBreak thestage of clustering. $\blacktriangle$  $\overline{\phantom{a}}$  $\overline{ }$  $\overline{\mathcal{A}}$  $\begin{bmatrix} \phantom{-} \end{bmatrix}$  $\overline{1}$ ATI ≃  $\Delta$  $\lambda$ to $\overline{\bigvee}$  $\overline{\phantom{a}}$  $\overline{\phantom{a}}$  of jet j  $\overline{)}$  $\overline{\phantom{a}}$  ordering $\overline{\phantom{a}}$  $\bigcap$  $\frac{1}{2}$  $\Lambda$ µ  $\mathsf{N}$ still trigger 0 67).) eliminates $(1)$   $(1)$  $\mathsf{I} \vee \mathsf{H}$

- Search for  $gg \rightarrow H \rightarrow bb$ historically thought impossible asymmetric, due to overwhelming and difficult to predict QCD background background Fig. 2: Sample of the leading-order Feynman diagrams, in the Full Theory with finite top mass equal top mass equal Theory with finite top mass equal top mass equal to the Full Theory with finite top mass equal to the contr FIG. 2: Sample of the leading-order Feynman diagrams, in the Full Theory with finite top mass e $\frac{1}{\sqrt{2}}$ to $\epsilon$  the  $\overline{\mathsf{b}}$ ).1  $\overline{c}$ candidate Higgs identify  $h$  $\check{ }$  $\overline{a}$  the  $\rightarrow$  1 nd<br>  $\overline{\phantom{0}}$  $\frac{1}{2}$  $\bigcup$ of our ar is $CL$ in<br>Jar<br>D م<br>د ا ししょう かんしゅう かんじょう かんしゃく かんしゃく しゅっとう かんしゅう かんしゅう かんしゅう かんしゅう かんしゅう かんしゅう けんしゅう かんしゅう かんしゅう かんしゅう  $\Omega$ C np  $\overline{\phantom{0}}$ l',  $\overline{\mathsf{C}}$ ing<br>a  $\frac{1}{2}$ t im<br>ina find  $\overline{\phantom{a}}$ e<br>... Weict l (  $\overline{\Gamma}$  $\overline{\phantom{a}}$  $\mathbf{r}$  $\overline{\mathsf{m}}$  $e$ ed<sup>-</sup> i i<br>Ight effective<br>Morti  $\rightarrow \vdash$ ore in ind<br>- $\overline{\phantom{a}}$ and modern m<br>Single pp<br>P  $\overline{a}$  $O($  $\Gamma$ pu<br>D ี<br>;hoเ  $\overline{\phantom{0}}$  $\blacktriangledown$ OL g ·<br>。。 el to r<sub>C</sub> gg<sup>r</sup>  $+1$  $\sqrt{2}$  $\overline{ }$ compared to those⊥⊔cui<br>ackgi ی ر<br>ال same  $\bigcirc$ alt <sup>.</sup><br>.  $\overline{a}$ mass.cult on.<br>allر  $\overline{\phantom{a}}$  butor g ca to<br>. diffic  $f<$ stor<br>Je to ric  $chfc$
- We can access this process in the boosted dijet topology of the boosted dijet topology  $\begin{array}{|c|c|c|c|c|}\n\hline\n\textbf{1} & \textbf{1} & \text$  $\blacksquare$  $\boldsymbol{J}$ ir y<br>y  $\mathsf{ss}\ \mathsf{ir}$  $\mathcal{C}$  $\frac{1}{2}$ mised and mistage<br>in<br>the mistage ces not $\overline{\mathcal{L}}$  $\overline{C}$  $\overline{\phantom{a}}$  $\overline{C}$  $p$  $\sim$  $\overline{\phantom{a}}$ is prc pi<br>nn is p<br>top:  $i$ et  $\ddot{\cdot}$  $\prod$ hi:  $\overline{\mathbf{u}}$  $\sinh \theta$  $\overline{\phantom{a}}$ s!  $\tilde{\mathbf{c}}$  $\overline{\phantom{a}}$
- Probing Higgs couplings at high momentum transfer (Q) accesses large new physics and the loop induced **g** *a g g g g g g g g g g g g g g g g g g g***<sub></sub> <b>***g***</del> <b>***g g g g g g g*  $\text{energy scale }(\Lambda)$  and the additional  $\frac{0.85}{50}$  for  $\frac{100}{100}$  and  $\frac{200}{250}$  and  $\frac{300}{350}$ shown in Figure 2. Note that in the production of the mediators in channels with associated *b* or *t* quarks is largely shown in Figure 2. Note that in the production of the mediators in channels with associated *b* or *t* quarks is largely  $\frac{10^{13}}{\frac{1}{2}}$  the tree-level terms, the in  $\Lambda$ filorie in a single coupling value, as in Higgs physics. A similar diagram induces couplings to photons. At leading-order, the  $\begin{bmatrix} 1 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 \end{bmatrix}$ momentu If the external particles in the loop induced *g g* (*A*) interaction are on-shell, then it can be exactly calculated accesses large new physics. At leading-order, the on-shell Lagrangian coupling-order on-shell lagrangians for o  $\sim$   $\sim$   $\sim$   $\sim$   $\sim$ ↵*s* ight<br>  $\frac{1}{2}$  $\mathbf{I}$ hig n.  $\mathsf{t}$ g<br>Since יטי<br>ר  $\overline{a}$ optimal

*gv*

*v*

8⇡

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*LS,*loop =

**v**<br>and the set of  $\alpha$  and  $\alpha$ 

↵*<sup>S</sup>*

dl





### SO HOW CAN WE DO IT?

- Inspiration from boosted Z'+jet search? V Use ISR jet to get you above the trigger threshold
	- Requires one boosted fat jet ✔
	- Substructure and jet grooming to ✔ enhance S/B
	- Data-driven background estimate ✔
- Inspiration from machine learning and b-tagging?
	- Double b-tagger selects fat jets ✔containing two b-quarks









BOOSTED GGF H(BB)



## SUBSTRUCTURE AND TAGGING

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## [arXiv:1307.0007](https://arxiv.org/abs/1307.0007)

#### arXiv:1307.0007<br>arXiv:1402.2657<br> distributions, the study show their could be understanding the study of the stu Stop when

- Provides good separation between W/Z/H-jets from q/g jets e Provides apod separation between  $M/Z/H$  iets from a/a iets of  $\mathcal{S}$  is substructure, the authors of Ref.  $\mathcal{S}$  developed the modified mass drop tagger (mMDT) substructure. naration between W/7/H-jets from a/a je para  $\sim$ satisfy software  $J/J$ criterion
- Grooming removes soft and wide-angle radiation (soft drop / modified mass soft drop) **BDRS, MMDT, Soft Drop, JHU top tagger, CMSTT** both subjects satisfy condition. The condition of the condition of the condition of the condition of the condi<br>In the condition. The condition of the con



Soft Drop Condition:  $\frac{m}{f}$ 

Soft Drop Condition: 
$$
\frac{\min(p_{T1}, p_{T2})}{p_{T1} + p_{T2}} > z_{\text{cut}} \left(\frac{\Delta R_{12}}{R_0}\right)^{\beta}
$$

 $\bigcap A C \cdot \mathbb{R}^n \cap \bigcap A \cap \mathbb{R}^n \cap \bigcap A \cap \mathbb{R}^n$ is the rapidity-azimuth plane, *z*cut is the soft of the soft drop threshold, and  $\sim$  $CMS: z<sub>cut</sub> = 0.1, β = 0$ 



### [arXiv:1609.07483](https://arxiv.org/abs/1609.07483) JET SUBSTRUCTURE

- How many "prongs" are in the jet? ow many "prongs" are in the jet?
- Generalized energy correlation functions are sensitive to N-point **26 2** Correlations within a jet
- A two-pronged jet has  $_2e_3 < (1e_2)^2$ • A two-pronged jet has <sub>2</sub>e<sub>3</sub> < (1e<sub>2</sub>)<sup>2</sup> 2-point 3-point 2-point 3-point
	- $\sim$  $\mathbf{J}^{-1}$  . • Taking a ratio gives  $N^1{}_2$  $\Delta$ 3 astudies we conclude that  $\Delta$ <sup>424</sup> power and shows similar discrimination power as *t*<sup>21</sup>  $\frac{1}{2}$   $\frac{1}{2}$

$$
1e_2^{\beta} = \sum_{1 \leq i < j \leq n_j} z_i z_j \Delta R_{ij}^{\beta}
$$

$$
2e_3^{\beta} = \sum_{1 \le i < j < k \le n_J} z_i z_j z_k \min\{\Delta R_{ij}^{\beta} \Delta R_{ik}^{\beta}, \Delta R_{ij}^{\beta} \Delta R_{jk}^{\beta}, \Delta R_{ik}^{\beta} \Delta R_{jk}^{\beta}\}\
$$

$$
\beta = 1
$$

*ik*D*R<sup>b</sup>*

*ik*D*R<sup>b</sup>*

*ij*D*R<sup>b</sup>*



*b*

*b*



) )

*zs*

) ✓*cc*  $\mathbf{r}$ 

C-Soft

) )

*zs*

 $\mathbf{r}$ 

*zcs*

)

*zcs*

)

ment itself allows for a powerful understanding of the jet's energy and angular structure.

C-Soft

*jk}* (8)

*zs*



ment itself allows for a powerful understanding of the jet's energy and angular structure.

*ij* (7)

*jk* , (10)

## $\frac{EXO-17-001}{2}$  $\frac{EXO-17-001}{2}$  $\frac{EXO-17-001}{2}$  N<sup>1</sup><sub>2</sub> IN PRACTICE

- Here's what the boosted Z'+jet analysis looks like after kinematic selection
- Difficult to use the QCD Monte Carlo to predict the background in this phase space
- Fitting this mass distribution directly requires high order polynomial → large background uncertainties
- Can we try a data-driven sideband prediction?







### SIDEBAND QCD PREDICTION

- Core idea: predict QCD jet mass distribution from region failing the tagger
- Possible problem: does tagger sculpt jet mass distribution?







SIDEBAND QCD PREDICTION

• Solution: define new substructure variable intended to be decorrelated from jet mass







### Z' RESULTS

• Jet mass distribution is fit down to 40 GeV





[EXO-17-001](http://cds.cern.ch/record/2264843?ln=en)



### SIDEBAND QCD PREDICTION (REDUX)

- Can we use the same QCD prediction when using a double-b tagger?
	- Yes if it's sufficiently decorrelated from jet mass and  $p_T$







### BOOSTED H(BB)

- With large boost, both b quarks merge into a single large radius jets **H(bb̄ )**
	- How can we best exploit the presence of the b-quarks in the jet in a tagger?  $\frac{1}{200}$





### MULTIPLE APPROACHES



- Based on standard b-tagging algorithm
- Not designed for two b's in the same jet



- Defines sub-jets
- Standard b-tagging applied to each subject
- Identifies two b hadron decay chains in the same fat jet
- Does not define subjects, but uses N-subjettiness axes





## DOUBLE B-TAGGER

- Combines tracking and vertexing information in a multivariate classifier with 27 observables **d** Combines tracking  $\sum_{i=1}^{n}$  multiparticle approach  $\sum_{i=1}^{n}$
- Targets the bb signal with additional aims: er additional aims<sup>.</sup>
	- $\bullet$  jet mass and  $p_T$  independent
- cover a very wide  $p_T$  range • **training strategy** is designed to cover a very wide pt range of the state of th
- inputs are chosen to avoid  $p_T$  $|$  *correlation* 
	- e.g. no ΔR-like variables, no substructure info







### CORRELATIONS?



• No strong correlations in double-b tagger versus  $m_{SD}$  or  $p_T$ in QCD background









# EVENT SELECT

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#### [HIG-17-010](http://cms.cern.ch/iCMS/analysisadmin/cadilines?line=HIG-17-010)

## EVENT SELECTION

- Online selection asks for a high  $p_T$  single jet or large hadronic activities
	- $p_T > 360$  GeV (m  $> 30$ ) or  $\Sigma$  p<sub>T</sub> > 900 GeV
- Offline: Highest  $p_T$  jet
	- $p_T > 450$  GeV,  $|\eta| < 2.5$
	- jet mass  $m_{SD} > 40$  GeV
- lepton veto, p<sub>T</sub><sup>miss</sup> veto
- $-6.0 < p = log(m_{SD}^2/p_T^2) < -2.1$







[HIG-17-010](http://cms.cern.ch/iCMS/analysisadmin/cadilines?line=HIG-17-010)

## EVENT SELECTION

**Substructure**: two prong discrimination, ~50% sig. efficiency, 26% bkg. efficiency

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**Double-b tagger:** ~30% sig. efficiency, 1% bkg. efficiency (tight working point)



#### SIGNAL COMPOSITION [HIG-17-010](http://cms.cern.ch/iCMS/analysisadmin/cadilines?line=HIG-17-010)

- Analysis is inclusive in Higgs production mode
- Dominant contribution is ggF (74%)
	- 12% VBF
	- 8% VH
	- $\bullet$  6% ttH







### BACKGROUND ESTIMATION BOOSTED GGF H(BB)





### BACKGROUND STRATEGY

- Backgrounds estimated from data
	- QCD (90%): from failing double b-tag x transfer factor
	- tt+jets (3%): from 1μ control region
- Backgrounds estimated from MC including NLO QCD + EWK corrections and jet mass, resolution, and substructure tagging scale factors
	- $W/Z + jets (5%)$
	- single-t,  $VV$  (<1%)

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#### QCD TRANSFER FACTOR D TRANK

- If the double b-tagger were completely uncorrelated from jet mass and  $p_T$ , the transfer factor would be flat
- Taylor expand as a polynomial in  $\rho$  and  $p_T$  to parameterize any small correlations **many** small correlations  $\overline{\Omega}$
- F-test determined 2<sup>nd</sup> order in  $\rho$  and 1<sup>st</sup> order in  $p_T$  is sufficient to fit the ratio

$$
N_{\text{pass}}^{\text{QCD}}(m_{\text{SD}}, p_{\text{T}}) = R_{\text{p/f}}(\rho, p_{\text{T}}) \cdot N_{\text{fail}}^{\text{QCD}}(m_{\text{SD}}, p_{\text{T}})
$$

$$
N_{\text{pass}}^{\text{QCD}}(m_{\text{SD}i}, p_{\text{T}j}) = \left( \sum_{k,\ell} a_{k\ell} \rho_{ij}^{k} p_{\text{T}j}^{\ell} \right) \cdot N_{\text{fail}}^{\text{QCD}}(m_{\text{SD}i}, p_{\text{T}j})
$$





### FITTING TRANSFER FACTOR

• Pre-fit both regions have the same predicted shape  $n$  ie same predicted shap<br>*a* 







#### $FITTING TRANSFER FACTOR$ fail (*m*SD, *p*T) (2)

• Post-fit signal region has slightly different shape with the ratio given by the polynomial transfer factor *j* |
|-<br>|-*· <sup>N</sup>*QCD







#### FINAL TRANSFER FACTOR  $\mathbf{h} \mathbf{h} \cap \mathbf{h}$  $\sqrt{ }$

• Two views of the same transfer factor function  $\mathcal{O}$ 











# HIGGS PT MODE

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- iu<br>. • Other CMS Higgs results use Powheg: 1 jet + m $_{\rm t}$  = ∞, [arXiv:1111.2854](http://arxiv.org/abs/1111.2854)
- or bo  $\ddot{\mathbf{e}}$  $\sf{corrections}$  and  $\sf{finite}$  top mass LL <sup>2</sup> LO+N <sup>3</sup> MRT, N • We want to account for both effects of **higher order** 
	- .<br>1. 1.  $\bullet\,$  No real NLO + finite top mass calculation available in the literature





• LO H+0-2jet, finite mt

•

•



- $\overline{\phantom{0}}$ gs results use Powheg: 1 jet + m *H*<sub>1</sub> *M*<sub>1</sub> ວ<sub>(</sub>  $\frac{1}{\sqrt{2}}$ • Other CMS Higgs results use Powheg: 1 jet + m<sub>t</sub> = ∞, <u>[arXiv:1111.2854](http://arxiv.org/abs/1111.2854)</u>
- corrections and finite top mass • We want to account for effects of higher order
	- $-9$ • No real NLO + finite top mass calculation available in the literature
- Adopt a factorized approach:

•

•

acceptance cuts <u>arXiv:1410.5806</u> → We build on this 0.8 1 • LO H+0-2jet, finite  $\mathsf{m}_{\mathsf{t}'}^{}$   $\mathsf{p}_{\mathsf{t}}^{}$ H up to 600 GeV, including WW



*<sup>B</sup>* (*t*)

GF H(NNLO + 
$$
m_t
$$
) = (1 jet  $m_t \to \infty$ ) ×  $\frac{MG\ LO\ 0-2\ \text{jet }m_t}{(1\ \text{jet }m_t \to \infty)}$ 

#### CNNW merged CKKW merged





- Other CMS Higgs results use Powheg: 1 jet +  $m_t = \infty$ , [arXiv:1111.2854](http://arxiv.org/abs/1111.2854)
- We want to account for both effects of higher order corrections and finite top mass
	- No real NLO + finite top mass calculation available in the literature
- Adopt a factorized approach: H

•

- LO H+0-2jet, finite  $m_{t'}^{\phantom{\dag}}$  p<sub>t</sub> up to 600 GeV, including WW acceptance cuts  $arXiv:1410.5806 \rightarrow$  $arXiv:1410.5806 \rightarrow$  We build on this 4
- NLO H+1jet finite  $m_{_t}$  up to 1/ $m_{_t}$  expansion: <u>arXiv:</u> 1609.00367



$$
GF H(NNLO + m_t) = (1 \text{ jet } m_t \to \infty) \times \frac{MG\text{ LO } 0 - 2 \text{ jet } m_t}{(1 \text{ jet } m_t \to \infty)} \times \frac{NLO\text{ 1 jet } m_t}{LO\text{ 1 jet } m_t}
$$
  

$$
CKKW\text{ merged} \qquad \text{factor of } 2
$$





#### HIGGS PT SPECTRUM  $\blacksquare$  to the Highest Eft order to the Highest Eft of the Highest Eft  $\begin{array}{c} \n\cdot \quad \text{L} \quad \text{L$

- Other CMS Higgs results use Powheg: 1 jet +  $m_t = \infty$ , [arXiv:1111.2854](http://arxiv.org/abs/1111.2854)
- We want to account for both effects of higher order corrections and finite top mass
	- No real NLO + finite top mass calculation available in the literature
- Adopt a factorized approach: H
- LO H+0-2jet, finite  $m_{t'}^{\phantom{\dag}}$  p<sub>t</sub> up to 600 GeV, including WW acceptance cuts  $arXiv:1410.5806 \rightarrow$  $arXiv:1410.5806 \rightarrow$  We build on this 4
- NLO H+1jet finite  $m_{_t}$  up to 1/ $m_{_t}$  expansion: <u>arXiv:</u> 1609.00367
- NNLO H+1jet,  $m_t = \infty$ ,  $p_T$  up to ~200 GeV, arXiv: H 1408.5325, arXiv:1302.6216, arXiv:1504.07922, arXiv: 1505.03893, [arXiv:1508.02684](http://arxiv.org/abs/1508.02684)



$$
GF H(NNLO + m_t) = (1 jet m_t \rightarrow \infty) \times \frac{MG LO 0 - 2 jet m_t}{(1 jet m_t \rightarrow \infty)} \times \frac{NLO 1 jet m_t}{LO 1 jet m_t} \times \frac{NNLO 1 jet m_t \rightarrow \infty}{NLO 1 jet m_t \rightarrow \infty}
$$
  

$$
CKKW merged \quad factor of 2 \quad factor of 1.25
$$





- Pythia version of CKKW-L merged 0,1,2jet LO finite top mass
- ME generation in aMC@NLO  $(pt > 20)$  with xqcut = 30 GeV
- CKKW shower is extended down to a merging scale of  $TMS = 20 GeV$
- Two factorized systematic uncertainties:
	- 30% overall normalization
	- 30% linear change in slope (no effect on overall norm.)











# RESULTS

BOOSTED GGF H(BB)

#### FIT RESULTS [HIG-17-010](http://cms.cern.ch/iCMS/analysisadmin/cadilines?line=HIG-17-010) • Simultaneous fit for Z(bb) and H(bb)



5.1 $\sigma$ ,  $\mu$ <sub>z</sub> = 0.78<sup>+0.23</sup><sub>-0.19</sub>



constraint of H(bb) signal systematics



#### FIT RESULTS [HIG-17-010](http://cms.cern.ch/iCMS/analysisadmin/cadilines?line=HIG-17-010) • Simultaneous fit for Z(bb) and H(bb)



observed H(bb) significance:

**1.5** $\sigma$ **,**  $\mu_{\rm H} = 2.3^{+1.8}$ **<sub>-1.6</sub>** 



#### [HIG-17-010](http://cms.cern.ch/iCMS/analysisadmin/cadilines?line=HIG-17-010)

#### FIT RESULTS passing and failing regions. Contributions from W and Z boson production are clearly visible

· Two dimensional likelihood scan in the single-julian the substructure and b-tagging strategy functions and b-tagging strategy for the substruct Higgs boson search in the same topology. The measured cross section of the Z+jets process







#### [HIG-17-010](http://cms.cern.ch/iCMS/analysisadmin/cadilines?line=HIG-17-010)

### P<sub>T</sub> CATEGORIES



![](_page_36_Picture_3.jpeg)

![](_page_36_Picture_4.jpeg)

### SUMMARY AND OUTLOOK

- First LHC search for  $gg \rightarrow H \rightarrow bb$  in boosted topology
	- First observation of Z(bb) in single-jet topology, 5.1σ observed (5.8σ expected)
	- Observed significance of H(bb) is 1.5σ
- Measured cross sections agree with SM
- Search probes **previously unexplored** regions of Higgs phase space
- New and generic strategy to search for boosted hadronic Higgs decays
	- Future prospects are bright
	- Means we need help from LHC H XS WG for best possible theory prediction in boosted Higgs regime —  $\mathsf{p}_\mathsf{T}$ H up to 1 TeV and beyond…

![](_page_37_Picture_10.jpeg)

### BOOSTED GGF H(BB)

### BACKUP

![](_page_38_Picture_2.jpeg)

![](_page_38_Picture_3.jpeg)

#### $N<sup>1</sup>$ <sub>2</sub> DDT

• Cut value map used to transform  $N^1{}_2$ 

![](_page_39_Figure_2.jpeg)

![](_page_39_Picture_3.jpeg)

![](_page_39_Picture_4.jpeg)

![](_page_40_Picture_0.jpeg)

### EVENT SELECTION

 $-6.0 < \rho < -2.1$ 

![](_page_40_Figure_3.jpeg)

![](_page_40_Picture_4.jpeg)

![](_page_40_Picture_5.jpeg)

#### EFFICIENCY AND MIS-TAG **Efficiency vs. Mistage rates of the second value of the secon**  $\Box$

![](_page_41_Figure_1.jpeg)

• Mis-tag is reduced by more than 40% at 30% signal efficiency for a tight working point **D** IVIIS-TAC *Mistag is reduced by more than 40% at 30% signal efficiency (~ tight working point)*

![](_page_41_Picture_3.jpeg)

![](_page_41_Picture_4.jpeg)

EFFICIENCY IN DATA **Efficiency measurement in data**

![](_page_42_Figure_1.jpeg)

- Using g(bb) jet as proxy in double muon tagged jet sample **24**
	- Associated data/MC uncertainty 3-5%

![](_page_42_Picture_5.jpeg)

![](_page_43_Figure_0.jpeg)

- Signal systematic uncertainties from merged W sample in semi-leptonic ttbar events (external constraint) *W* sample in semi-leptonic tion events in the single-jet topology (further value and b-tagging strategy for the substructure and b-tagging strategy for
- SM candles: presence of W/Z in final jet mass distribution provides additional in-situ constraint The measured H boson street is *µH* = 2.3 and includes the corrections to the corrections mass abundulum provides additional in-situ  $\sim$   $\sim$  summarizes the measured signal strengths and  $\sim$  the Higgs and Z bosonic the Higgs and Z bosonic bosonic bosonic that  $\sim$

processes. In particular, they are also reported for the case the corrections to the Higgs *p*<sup>T</sup> spec-

![](_page_43_Figure_3.jpeg)

**ErfExp Data comp. ErfExp MC comp.**

<sup>40</sup>

200

400

600

800

1000

1200

1400

![](_page_43_Picture_4.jpeg)

### P<sub>T</sub> CATEGORIES

![](_page_44_Figure_1.jpeg)

![](_page_44_Figure_2.jpeg)

![](_page_44_Picture_3.jpeg)

 $m_{SD}$  (GeV)

![](_page_44_Picture_4.jpeg)

#### [HIG-17-010](http://cms.cern.ch/iCMS/analysisadmin/cadilines?line=HIG-17-010)

### P<sub>T</sub> CATEGORIES

![](_page_45_Figure_2.jpeg)

![](_page_45_Picture_3.jpeg)

![](_page_45_Picture_4.jpeg)

## LARGE HADRON COLLIDER

• Proton-proton collisions at

**ENS** 

 $\frac{1}{2}$ 

Javier Duarte

Fermilab

![](_page_46_Picture_1.jpeg)

### **CERN**

**1 0Dy**

**0**

**7RtDO In tHJUDtHdLumLn**

**CMS** 

27 kilometer ring

**5**

**10**

**15**

**20**

**sLt y(**fb

−1 **)**

**25**

**1 Jun**

**1 JuO**

**1 AuJ**

**LHC DHOLVHUH COS 5HC 5HC 5HC 5HC** 

**1 SHS**

**DD** 

**1 2ct**

**1 1Rv**

**1 DHc**

**0**

**10**

**15**

### What is the best Higgs  $p_{\tau}$ :Options

- The key is to identify two different effects
	- Finite top mass effect
	- NNLO differential corrections
- What are the orders known:
	- Differential EFT : NNLO H+1jet production
	- Finite top mass : almost NLO
	- At MC level EFT : NLO H+0/1/2jet
	- At MC level finite top mass : LO 0/1/2

![](_page_47_Picture_9.jpeg)

![](_page_47_Picture_10.jpeg)

### Going to EFT

• When going to EFT large gain

![](_page_48_Figure_2.jpeg)

![](_page_48_Picture_3.jpeg)

X/Powheg

![](_page_48_Picture_4.jpeg)

### Going to EFT

• Adding the finite top mass merged LO its lower

![](_page_49_Figure_2.jpeg)

![](_page_49_Picture_3.jpeg)

![](_page_49_Picture_4.jpeg)

### $CKKW-L$  | L.Lönnblad, JHEP 05 (2002) 046, L.Lönnblad and S.Prestel,JHEP 03 (2012) 019

Idea: Reduce the dependence to the merging scale MS.

- Start by generate events with  $N_1..N_2$  ME partons, hard and well separated
- Assume an event with n ME partons, reconstruct the possible shower histories, pick one according to the occurence probabilities

![](_page_50_Figure_4.jpeg)

- $\bullet$  Each clustering step *i* is characterized by the emission scale  $\rho_i$ , reweight by the product of  $\alpha_s(\rho_i)/\alpha_s(ME)$
- For  $i=2..n$   $(l=N_2)$ :

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- Generate one emission  $\rho$  with  $\rho_i$  as starting scale.
- **•** If  $\rho > \rho_{i+1} \Rightarrow$  reject the event.
	- **•** This is equivalent to the product of Sudakovs  $\Pi(\rho_i, \rho_{i+1})$ , i=2..n-1.
- if not HME: generate an emission at  $\rho < \rho_n$ , if  $\rho > MS \Rightarrow$  reject the event.
- if HME, accept the event and start shower with  $\rho_n$ .

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Essential parameters for matching/merging:

- ickkw
	- Applies the  $\alpha_s$  reweighting at each QCD vertex in the ME calculation.  $K<sub>T</sub>$ -MLM, Shower- $K<sub>T</sub>$ :1 CKKW-L, UMEPS:0
- xqcut
	- **O** Defines the minimal  $K_T$  between the partons (+beam) at ME level.
- **•** auto\_ptj\_mjj
	- **O** Set to False: leaves the xqcut be the only cut applied to ME partons  $\Rightarrow$  ptj, mmjj=0
- maxjetflavor
	- $\circ$  QCD partons with pdgId $\leq$ maxjetflavor are affected by xqcut ptj,etc... Otherwise, affected by ptb, mmbb, etc... That means that for a n-Flavour prediction, maxjetflavor = n

![](_page_51_Picture_10.jpeg)

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### Practical use: main89.cc

main89ckkwl.cmnd: CKKWL. Essential parameters are

- Merging:TMS = XXX.
	- The merging scale
- Merging:Process = UUU
	- Type of process, e.g. pp>LEPTONS,NEUTRINOS
- Merging:nJetMax = WWW
	- Maximal number of additional jets in the matrix element
- Merging:doPTLundMerging = on
	- $\bullet$  Set the merging scale definition to  $P_{T,evol}$  (cfr definition in the manual)

main89umeps.cmnd: UMEPS. Essential parameters are

- Merging:TMS = XXX.
- Merging:Process = (e.g.) pp>LEPTONS,NEUTRINOS
- Merging:nJetMax = WWW
- Merging:doUMEPSTree = on
	- Reweight events according to the UMEPS prescription for tree-level configurations)

![](_page_52_Picture_16.jpeg)

![](_page_52_Picture_18.jpeg)

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![](_page_52_Picture_19.jpeg)

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