



Report on the VH subgroup activities (Theory part)

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VH subgroup conveners:

John Campbell, Carlo Pandini, Luca Perrozzini and FT

Outline

- NNLOPS
- NLOPS QCD+EW corrections
- VH(bb)@NNLO
- ggZH@NLO

NNLO+PS MATCHING: ASSOCIATED HIGGS PRODUCTION

*Wojciech Bizoń
(University of Oxford)*

*in collaboration with:
William Astill, Emanuele Re, Giulia Zanderighi*

LHC HXSWG – WG1 meeting (29/06/2017)

- Already in YR4 NNLOPS for HW: in a nutshell

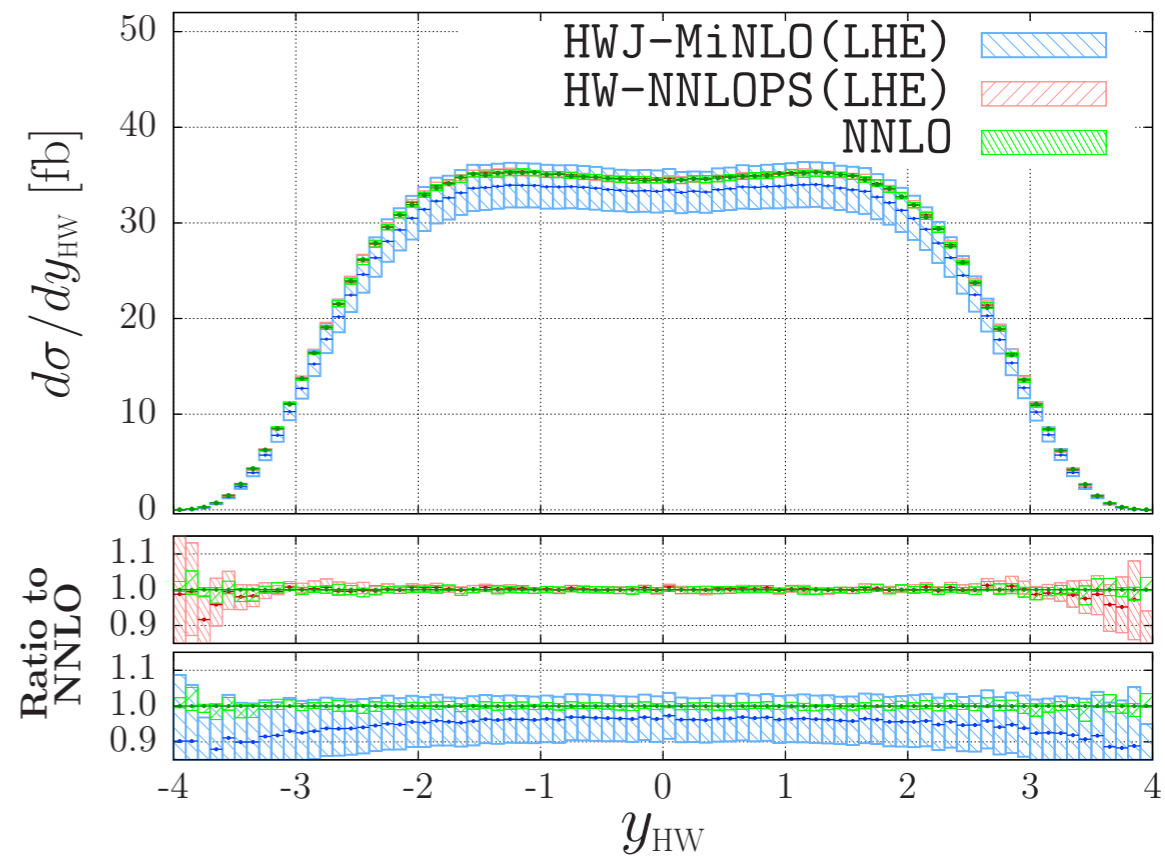
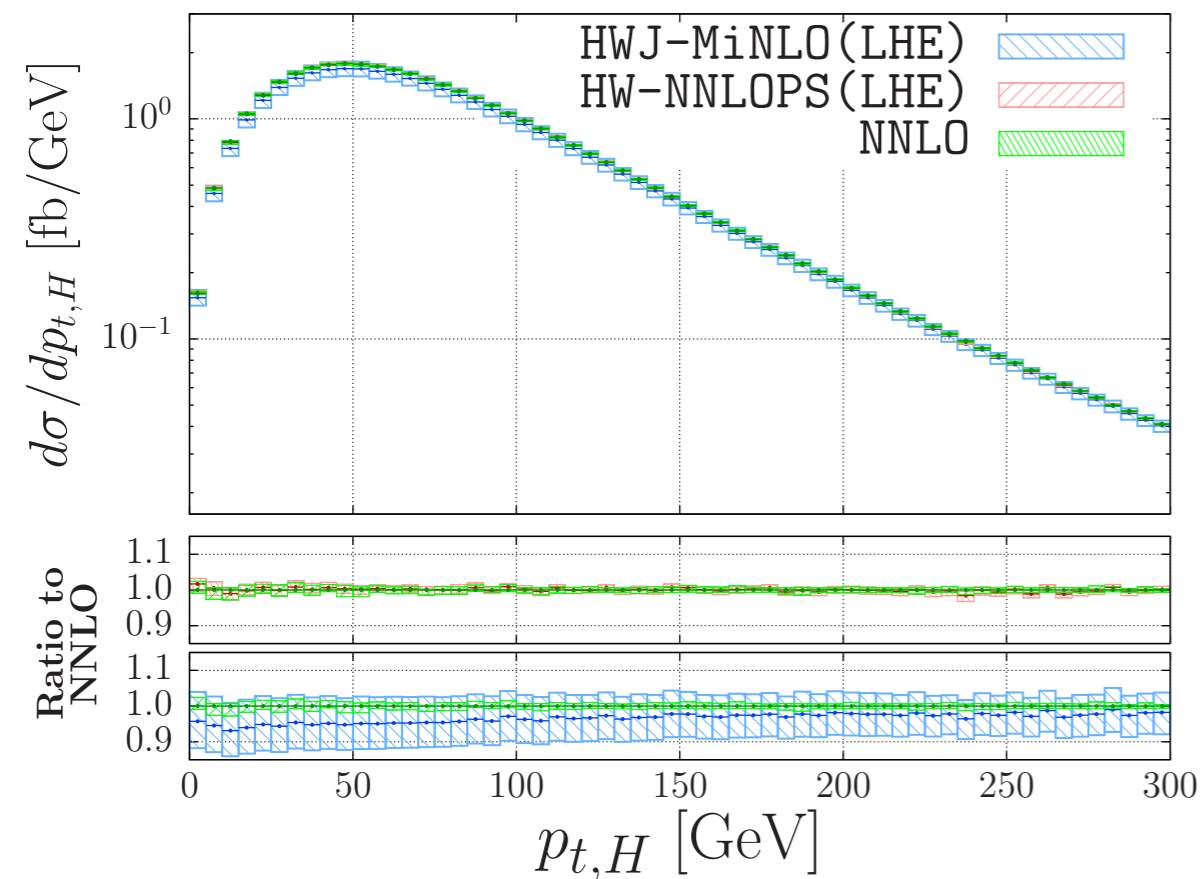
Starting from the VHJ generator:

$$\blacktriangleright \tilde{B}_{\text{MiNLO}} = \alpha_s(q_T) \Delta^2(q_T, \bar{\mu}_R) \left[B \left(1 - 2\Delta^{(1)}(q_T, \bar{\mu}_R) \right) + \alpha_s(\bar{\mu}_R) \left(V(\bar{\mu}_R) + \int d\Phi_r R \right) \right]$$

Preserve NLO⁰ accuracy for VH

$$\blacktriangleright W(\Phi_B) = \frac{\left(\frac{d\sigma}{d\Phi_B} \right)_{\text{NNLO}}}{\left(\frac{d\sigma}{d\Phi_B} \right)_{\text{MiNLO}}} = \frac{d\sigma^{(0)} + d\sigma^{(1)}\alpha_s + d\sigma^{(2)}\alpha_s^2}{d\sigma^{(0)} + d\sigma^{(1)}\alpha_s + d\tilde{\sigma}^{(2)}\alpha_s^2} = 1 + \frac{d\sigma^{(2)} - d\tilde{\sigma}^{(2)}}{d\sigma^{(0)}}\alpha_s^2 + \mathcal{O}(\alpha_s^3)$$

Preserve NLO accuracy for VHJ production



GROWING COMPLEXITY

- Easy to imagine: with bigger phase-space (formally simple) procedure becomes computationally involving...

(a) Higgs production (ggH): 1-dim 1 variable (1D histogram = 25 bins)

(b) Drell-Yan production: 3-dim 3 variables (3D histogram = 15 625 bins)

(c) VH production: 6-dim 6 variables (6D histogram = ??? [244M bins])

- phase-space parametrisation:

1	2	3	4	5	6
y_{VH}	$p_{t,H}$	Δy	θ^*	ϕ^*	$m_{\ell\bar{\ell}'}$

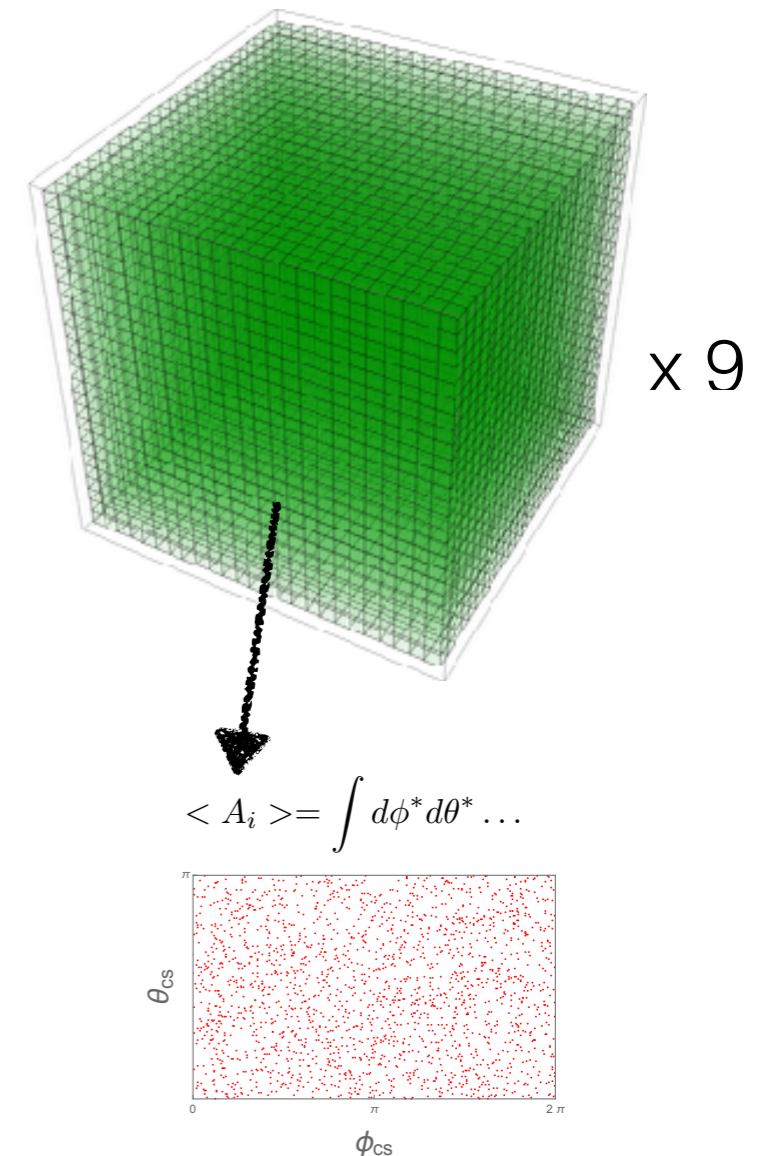
- cross-section in terms of Collins-Soper angles:

$$\frac{d\sigma}{d(\cos\theta^*)d\phi^*} = \frac{3\sigma}{16\pi} \left[(1 + \cos^2\theta^*) + A_0 \frac{1}{2} (1 - 3\cos^2\theta^*) + A_1 \sin 2\theta^* \cos\phi^* \right. \\ \left. + A_2 \frac{1}{2} \sin^2\theta^* \cos 2\phi^* + A_3 \sin\theta^* \cos\phi^* + A_4 \cos\theta^* \right. \\ \left. + A_5 \sin\theta^* \sin\phi^* + A_6 \sin 2\theta^* \sin\phi^* + A_7 \sin^2\theta^* \sin 2\phi^* \right]$$

- neglect dependence on $m_{\ell\bar{\ell}'}$ (validated)

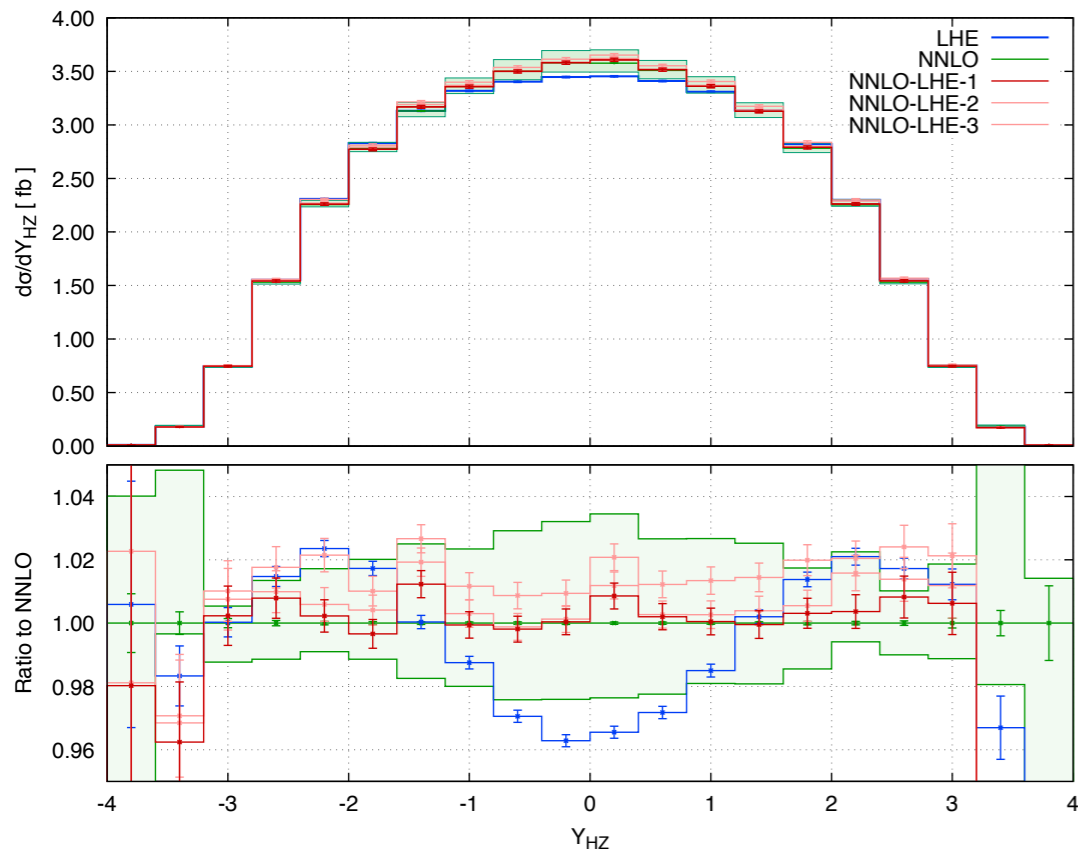
FINALLY:

- one 3D histogram for each A-coefficient (8+1 tables)
- still numerically challenging as each bin is an integral over 2-dim phase-space

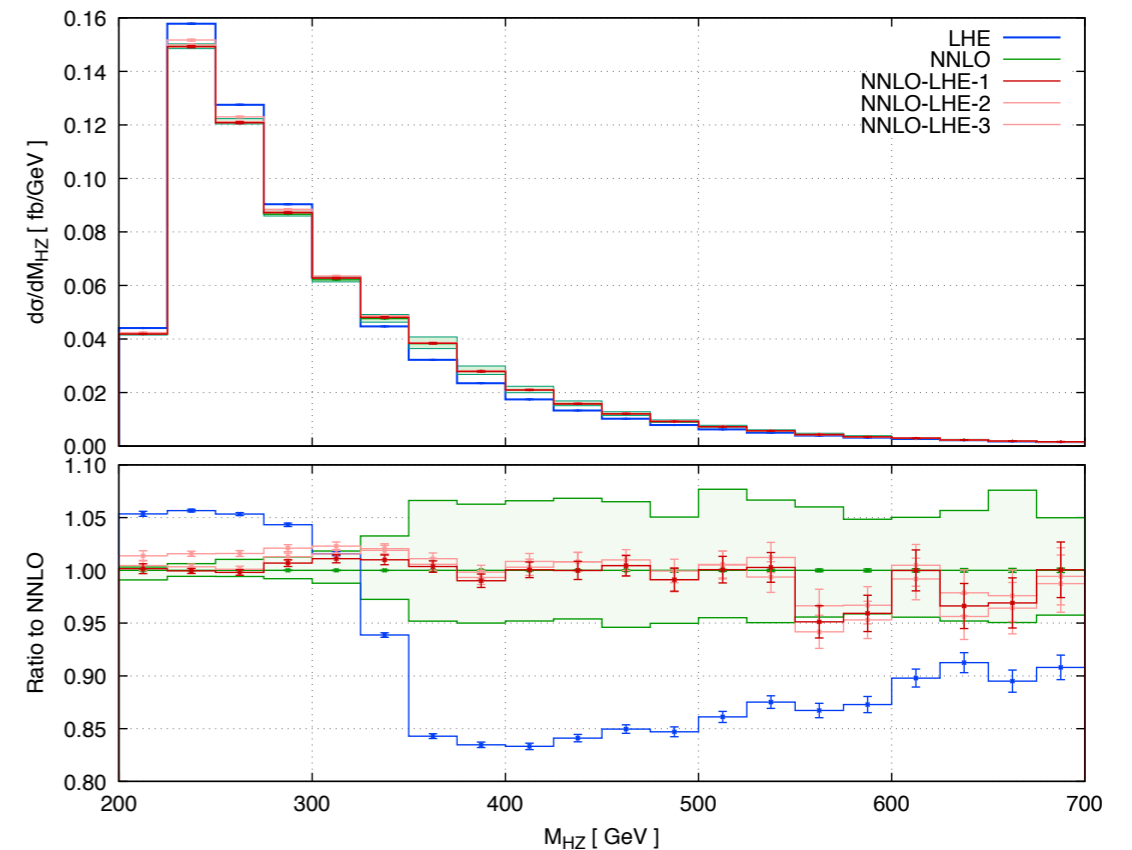


REWEIGHTING UNCERTAINTY (HZ)

- ▶ large phase-space ==> computationally heavy task
 - ▶ HW@NNLOPS: “smooth” enough distributions required very long runs (~ 1 month x 300cores)
 - ▶ Is it essential to have that long runs?
 - results below (HZ) were prepared with NNLO runs (~ 2 days x 2000cores) and 12.5M HZJ events
 - we have used various setups:
 - (a) reweight only with three basic variables (neglect Collins-Soper angles: $A(i)=0$)
 - (b) neglect $A(i)$ coefficients with large uncertainties (stat.err > 200%, stat.err > 50%)
- => use less precise histograms for reweighting but assign an error associated with this procedure



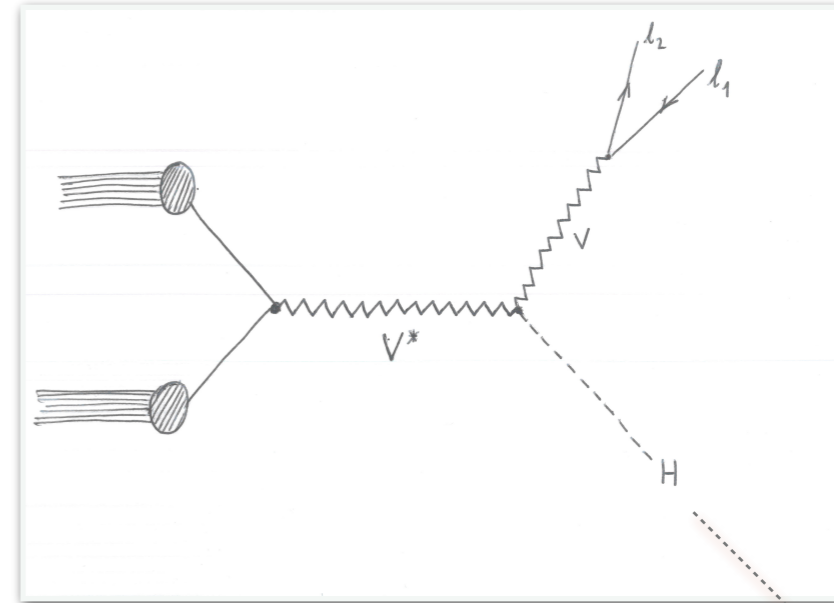
PRELIMINARY!



PRELIMINARY!

INCORPORATING HBB DECAY AT NLO

- **Hbb** – largest SM Branching Ratio ($\sim 60\%$)
- Allows for “precision” measurements in non-primary H-production channels



REWEIGHTING: TREATMENT OF THE DECAY

- in narrow width approximation phase-space split into production/decay:

$$d\Phi_{Vb\bar{b}} = d\Phi_{VH} \times d\Phi_{(H\rightarrow b\bar{b})}$$

- NNLO reweighting performed using Born kinematics hence we can use the same setup as without Higgs decay (we are actually changing setup but purely for practical purposes).
- This approach secures NNLO accuracy in production stage.
- NNLO-LHE: Hbb decay is treated at NLO within POWHEG (i.e. virtual corrections + some events contain real emission from bb-pair) which enables probing decay observables at NLO.

USING THE CODE

- VH Reweighting requires two sources of input
 - (1) HWJ / HZJ @POWHEG+MiNLO
 - (2) HW / HZ @NNLO
- our code contains:
 - patches (analysis, identical physical parameters, ...) to produce compatible results
 - hv_minnlo: program for reweighting event files using multidimensional histograms
- for HW(NNLOPS) we have used HVNNLO code:
[1107.1164; G.Ferrera, M.Grazzini, F.Tramontano], [1407.4747; G.Ferrera, M.Grazzini, F.Tramontano]
- for HZ(NNLOPS) we are using MCFM-8.0 for NNLO distributions:
[1601.00658; J.Campbell, R.K.Ellis, C.Williams]
- we are planning to release the full code with detailed manual shortly after HZ publication
- for the time being, we are able to provide multidimensional HW distributions used in first paper [1603.01620], disadvantage: fixed settings

Should other NNLO codes become available in the meantime,
we can help interested users to interface it with our NNLO-reweighter!

NLO QCD+EW CORRECTIONS FOR HV AND HV +JET IN THE POWHEG BOX RES

Carlo Oleari

Università di Milano-Bicocca, Milan

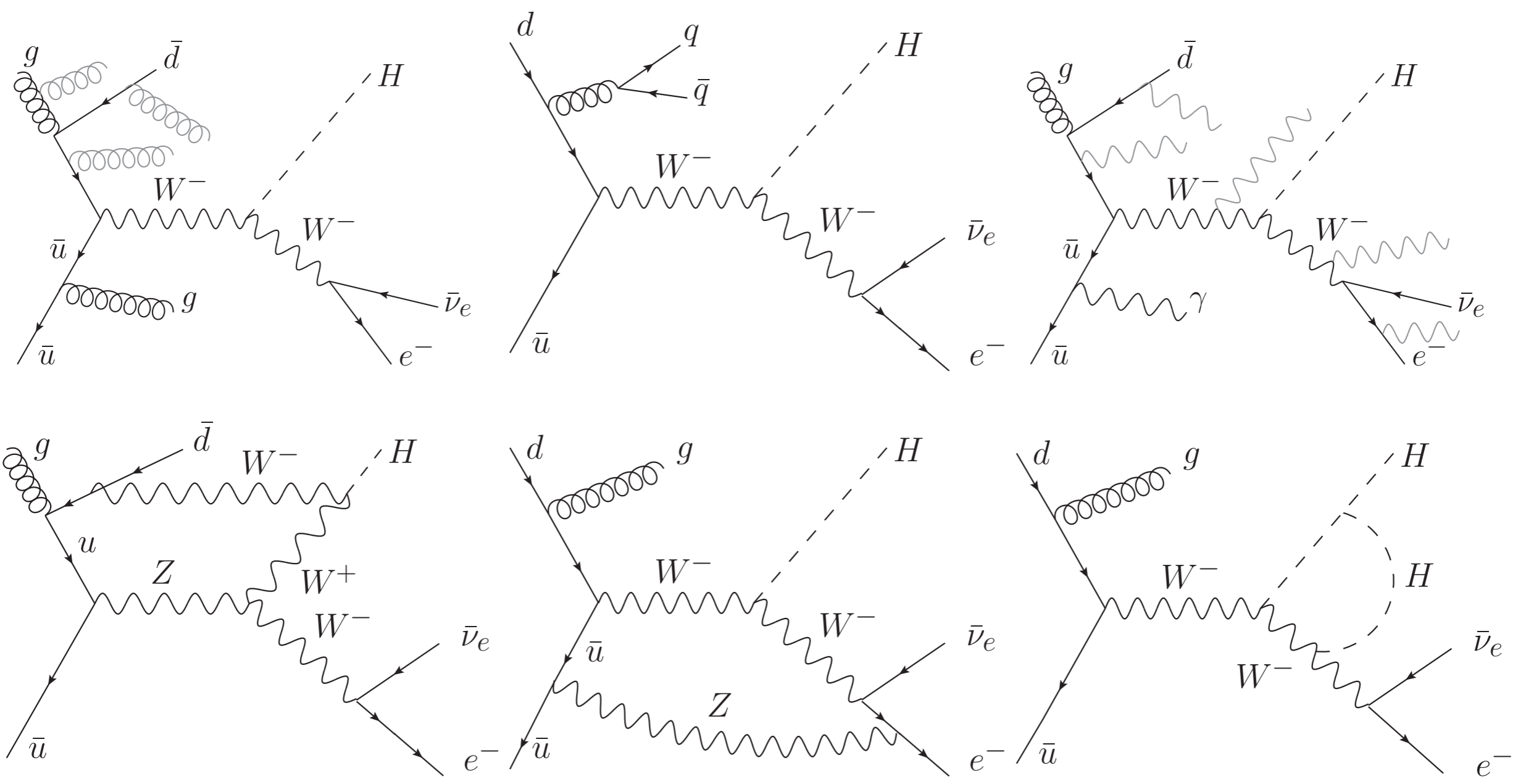
29 June 2017

In collaboration with: **F. Granata, J. Lindert and S. Pozzorini**

arXiv:1706.03522

- ✗ NLO QCD+EW HV and HVj production
- ✗ Resonances and the POWHEG BOX RES code
- ✗ $HVj+$ MiNLO
- ✗ A few results
- ✗ Conclusions

QCD+EW corrections to HVj



Born: $\mathcal{O}(\alpha_S \alpha_{EM}^3)$

QCD real+virtual: $\mathcal{O}(\alpha_S^2 \alpha_{EM}^3)$

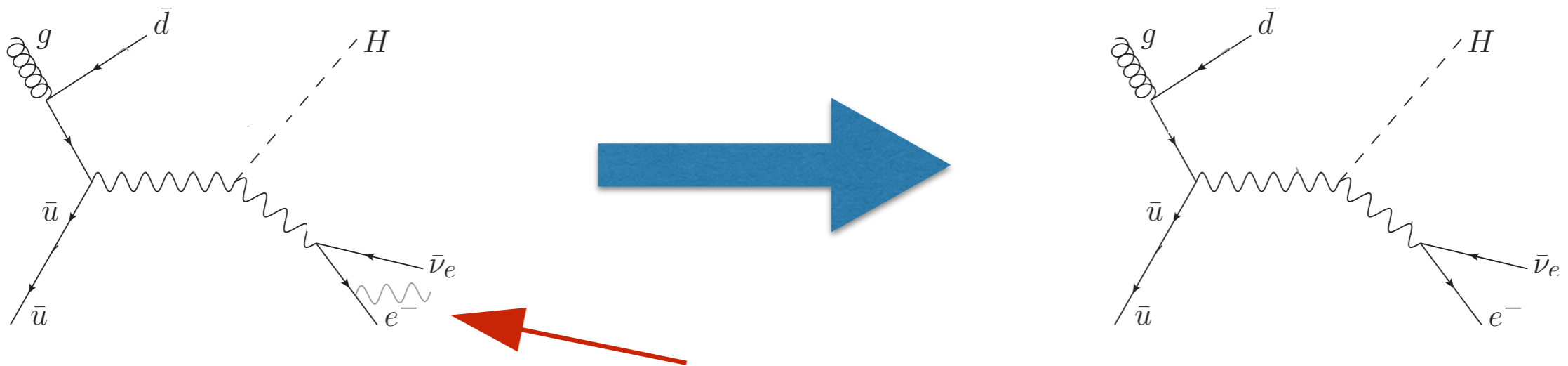
EW real+virtual: $\mathcal{O}(\alpha_S \alpha_{EM}^4)$

Sensitive to the **trilinear** Higgs boson coupling.

All EW amplitudes computed with OpenLoops that recently achieved automation also for EW corrections

Resonances

When dealing with **resonances** whose decay products can radiate, we have two technical problems to tackle. Consider for example $e^- \bar{\nu}_e \mu^+ \nu_\mu b \bar{b}$



1. mismatch of resonance virtuality among real and subtractions in the NLO computation
2. more seriously this mismatch affect the R/B in POWHEG event generation

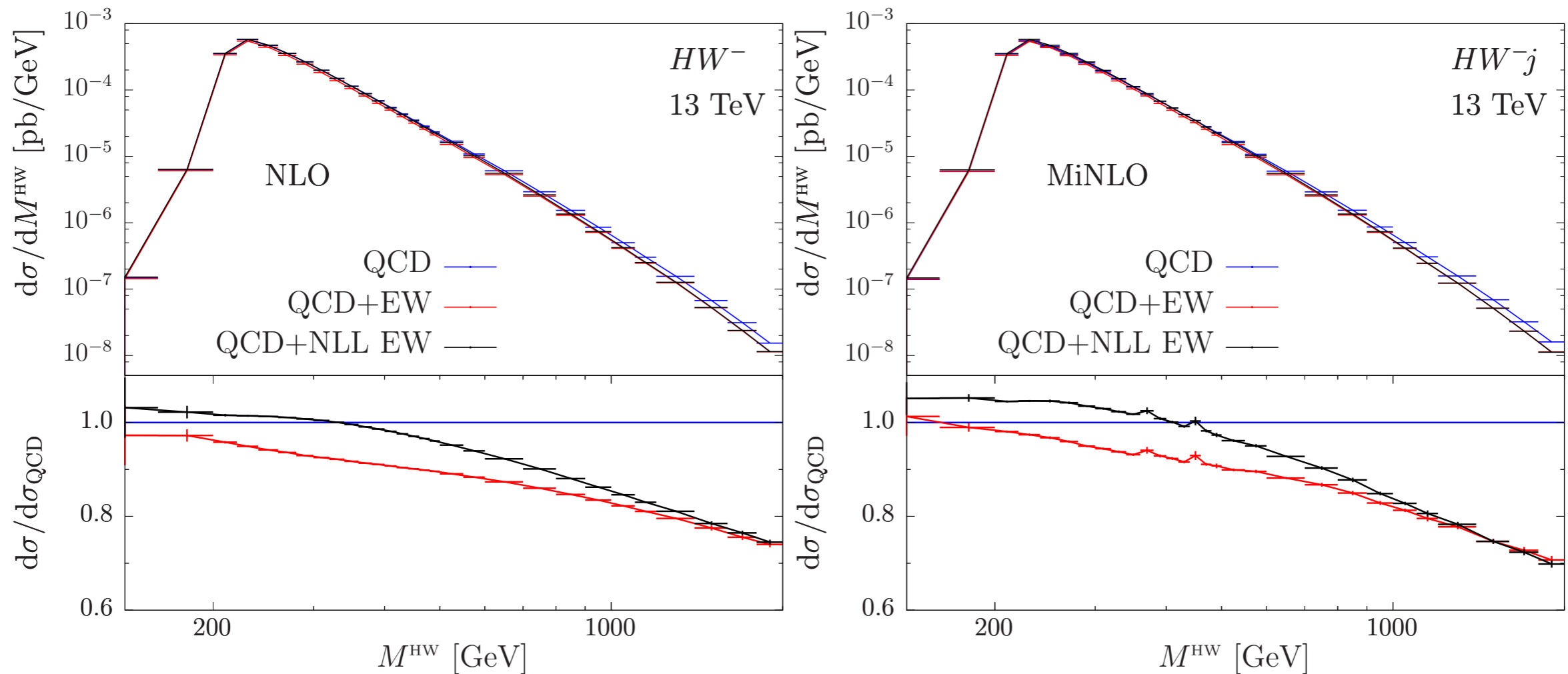
The POWHEG BOX RES

The solutions have been discussed in [Jezo, Nason, arXiv:1509.09071](#). The output of this has been a **major revision** of the POWHEG BOX V2 code: the **POWHEG BOX RES**.

- For each flavour structure, the code automatically finds all the possible **resonance histories** compatible with the partonic process at hand and keeps track of them, while generating radiation from each resonance, **preserving the virtuality** of the resonances.

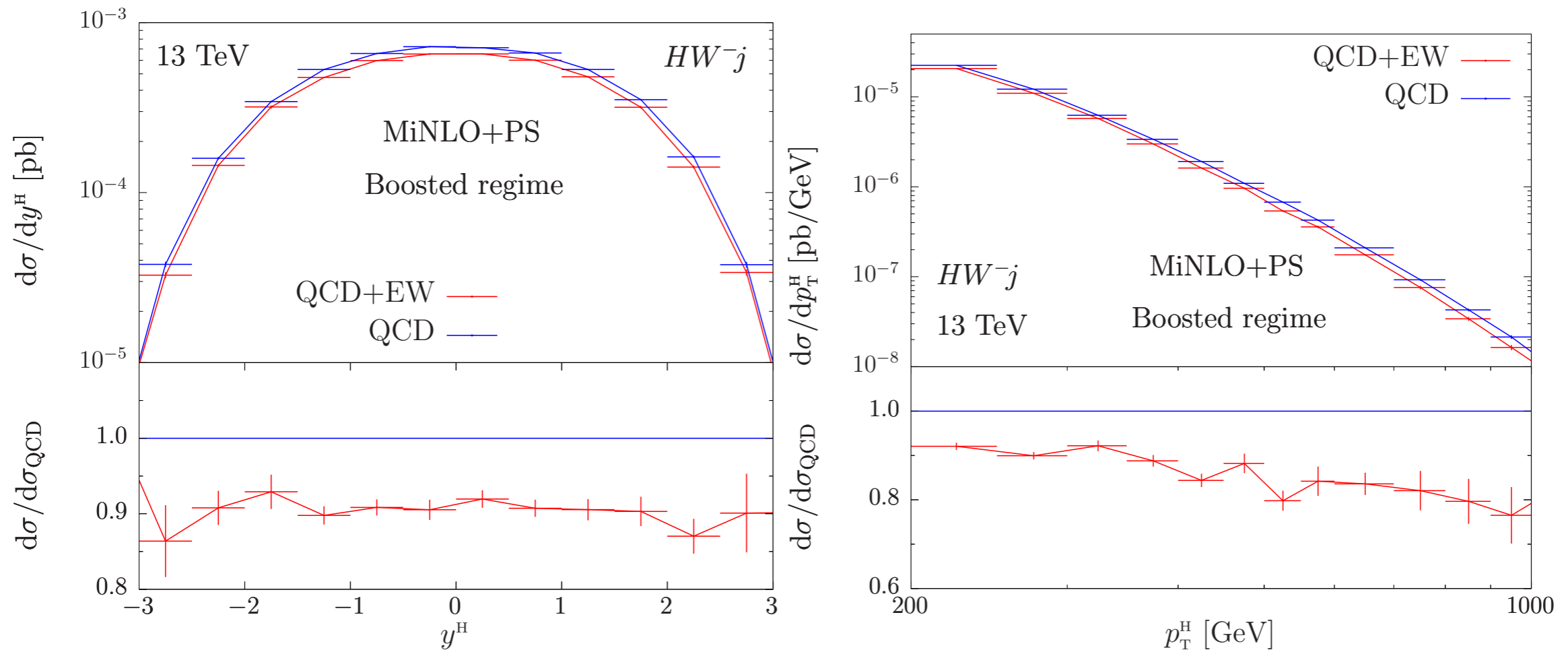
Applied now to HV and HVj production, where the **virtuality** of the **V boson** is preserved when **photon radiation** is produced.

NLO results at fixed order for HW^- and HW^-j production



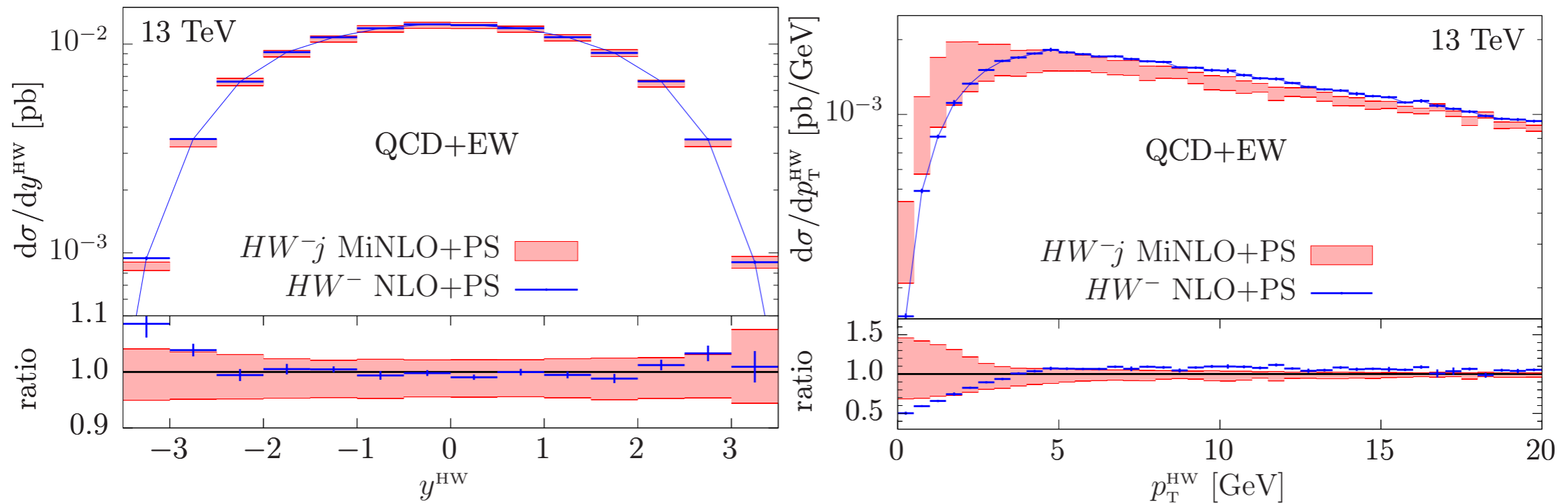
- **EW corrections** can largely exceed the ten percent level in the **high-energy** regions, where **Sudakov logarithms** become **dominant**.
- An example is the invariant mass of the HV pair in HV and HVj production, where the EW corrections reach **-30%** around 2 TeV.

MiNLO + Parton Shower results for HW^-j production



- These results **closely agree** with the corresponding ones for HW^- production.
- This supports the fact that the **MiNLO** predictions for HVj should preserve **NLO QCD+EW** accuracy for **inclusive** (with respect to the jet) quantities.

HV vs. HVj generators



- **Scale variation** bands ([details in arXiv:1706.03522](https://arxiv.org/abs/1706.03522))
- With **MiNLO**, the y^{HW} and p_T^{HW} distributions computed with the HWj generator are **finite** and agree with the results for HW .
- y^{HW} has **NLO** accuracy both in HV and with HVj .
 p_T^{HW} has **LO** accuracy for HV and **NLO** accuracy for HVj .

Possible recipe for QCD@NNLOPS+EW@NLOPS (C. Oleari)

In principle one could get distributions with the highest achievable accuracy combining 3 event samples as follows:

- 1) event sample with QCD @ NNLOPS
- 2) event sample with EW @ NLOPS
- 3) event sample with LO PS

$$\text{QCD NNLO} + \text{EW NLO} + \text{PS} = 1 + 2 - 3$$

VH production with $H \rightarrow b\bar{b}$ decay in full NNLO QCD

Giancarlo Ferrera

Università di Milano & INFN Milano



**LHC Higgs Cross Section Working Group
VH subgroup – June 29th 2017**

In collaboration with: M. Grazzini, G. Somogyi & F. Tramontano

Associated VH production with $H \rightarrow b\bar{b}$ decay

$$h_1(p_1) + h_2(p_2) \rightarrow V + H + X \rightarrow \ell_1 \ell_2 + b\bar{b} + X$$

where $V = Z^0, W^\pm$ and $\ell_1 \ell_2 = \ell^+ \ell^-, \ell \nu_\ell$

QCD factorization formula

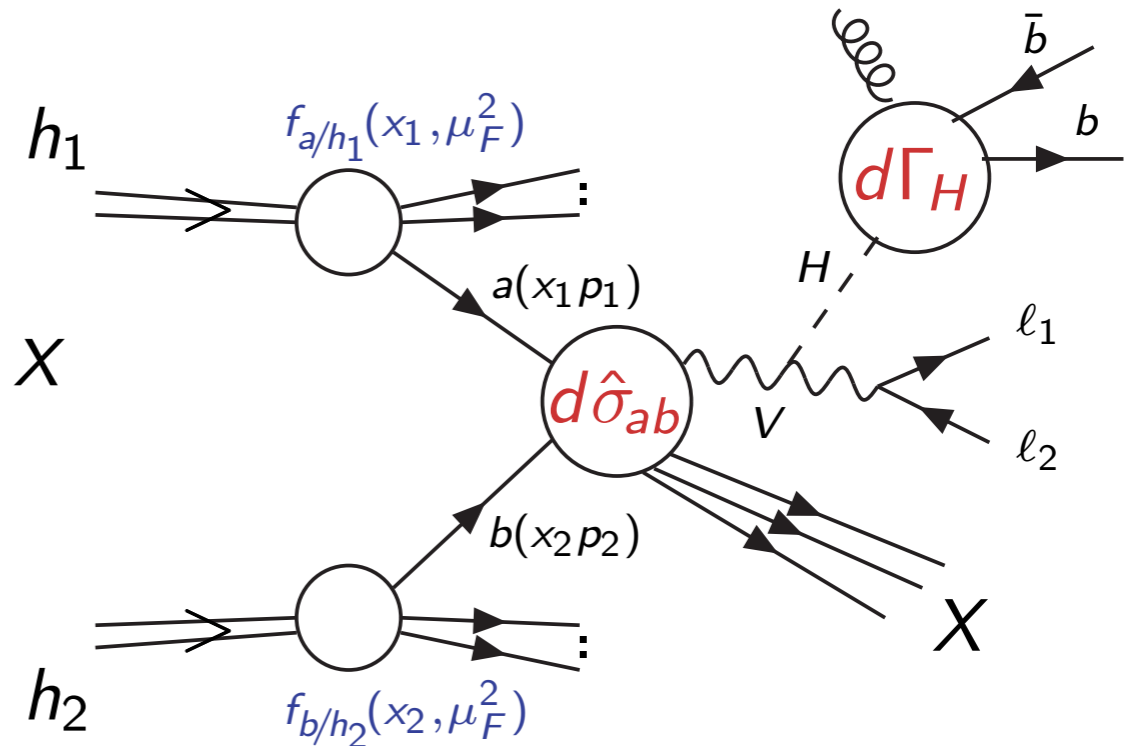
$$d\sigma = \sum_{a,b} \int_0^1 dx_1 \int_0^1 dx_2 f_{a/h_1}(x_1, \mu_F^2) f_{b/h_2}(x_2, \mu_F^2) d\hat{\sigma}_{ab}(x_1 p_1, x_2 p_2; \mu_F^2) + \mathcal{O}\left(\frac{\Lambda_{\text{QCD}}}{Q}\right)^p$$

- By using the zero width approximation ($\Gamma_H \ll m_H$)

$$d\sigma_{VH \rightarrow Vb\bar{b}} = d\sigma_{VH} \times \frac{d\Gamma_{H \rightarrow b\bar{b}}}{\Gamma_H} = d\sigma_{VH} \times \frac{d\Gamma_{H \rightarrow b\bar{b}}}{\Gamma_{H \rightarrow b\bar{b}}} \times \text{Br}(H \rightarrow b\bar{b}),$$

- Perturbative expansion gives

$$d\sigma_{VH \rightarrow Vb\bar{b}}^{\text{NNLO+nlo}} = \left[d\sigma_{VH}^{(0)} \times \frac{d\Gamma_{H \rightarrow b\bar{b}}^{(0)} + d\Gamma_{H \rightarrow b\bar{b}}^{(1)}}{\Gamma_{H \rightarrow b\bar{b}}^{(0)} + \Gamma_{H \rightarrow b\bar{b}}^{(1)}} + \left(d\sigma_{VH}^{(1)} + d\sigma_{VH}^{(2)} \right) \times \frac{d\Gamma_{H \rightarrow b\bar{b}}^{(0)}}{\Gamma_{H \rightarrow b\bar{b}}^{(0)}} \right] \times \text{Br}(H \rightarrow b\bar{b}),$$



Associated VH production with $H \rightarrow b\bar{b}$ decay

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QCD factorization formula

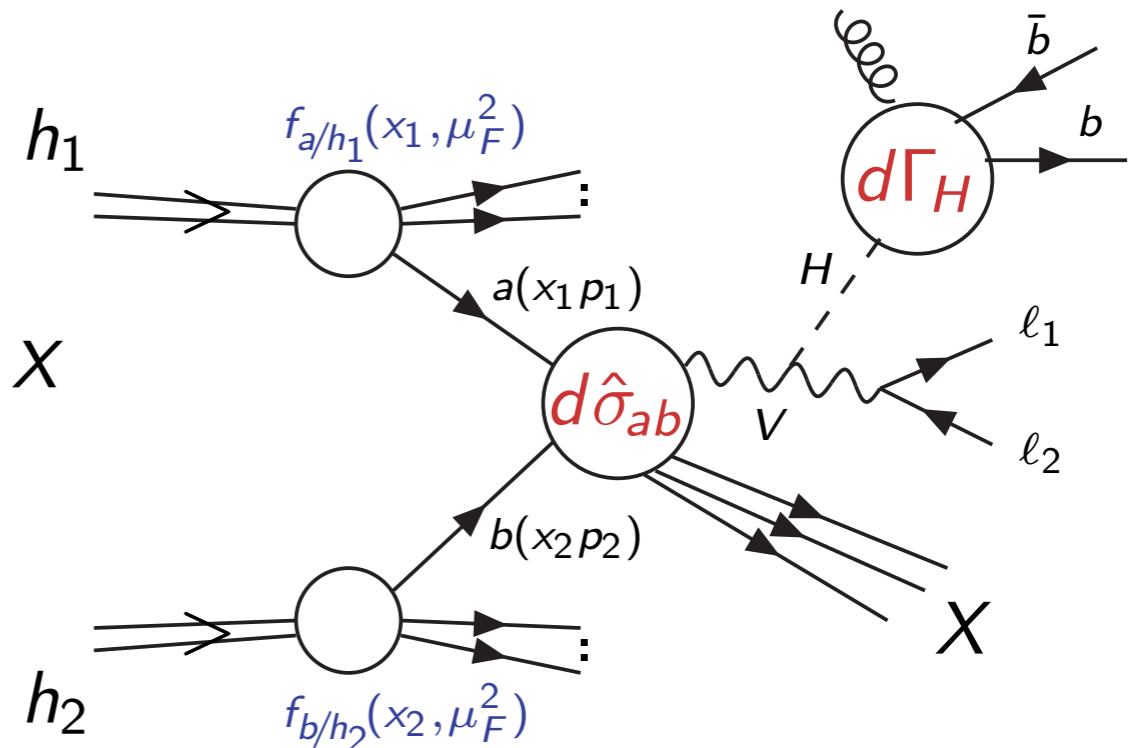
$$d\sigma = \sum_{a,b} \int_0^1 dx_1 \int_0^1 dx_2 f_{a/h_1}(x_1, \mu_F^2) f_{b/h_2}(x_2, \mu_F^2) d\hat{\sigma}_{ab}(x_1 p_1, x_2 p_2; \mu_F^2) + \mathcal{O}\left(\frac{\Lambda_{\text{QCD}}}{Q}\right)^p$$

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- Perturbative expansion gives

$$d\sigma_{VH \rightarrow Vb\bar{b}}^{\text{full NNLO}} = \left[d\sigma_{VH}^{(0)} \times \frac{d\Gamma_{H \rightarrow b\bar{b}}^{(0)} + d\Gamma_{H \rightarrow b\bar{b}}^{(1)} + d\Gamma_{H \rightarrow b\bar{b}}^{(2)}}{\Gamma_{H \rightarrow b\bar{b}}^{(0)} + \Gamma_{H \rightarrow b\bar{b}}^{(1)} + \Gamma_{H \rightarrow b\bar{b}}^{(2)}} + d\sigma_{VH}^{(1)} \times \frac{d\Gamma_{H \rightarrow b\bar{b}}^{(0)} + d\Gamma_{H \rightarrow b\bar{b}}^{(1)}}{\Gamma_{H \rightarrow b\bar{b}}^{(0)} + \Gamma_{H \rightarrow b\bar{b}}^{(2)}} + d\sigma_{VH}^{(2)} \times \frac{d\Gamma_{H \rightarrow b\bar{b}}^{(0)}}{\Gamma_{H \rightarrow b\bar{b}}^{(0)}} \right] \times \text{Br}(H \rightarrow b\bar{b}),$$



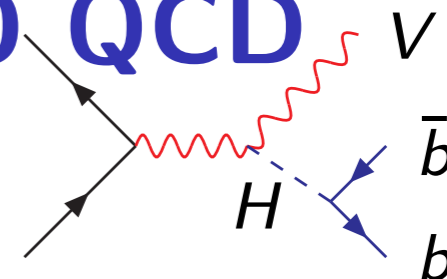
VH production and decay in full NNLO QCD

G.F., Somogyi, Tramontano arXiv:1705.0304

Fully differential NNLO calculation for VH production including $H \rightarrow b\bar{b}$ at NNLO and $V \rightarrow l_1 l_2$ decays with spin correlations.

- NNLO calculation for $h_1 h_2 \rightarrow VH + X$ production calculated in [G.F., Grazzini, Tramontano('11, '15)] within the q_T -subtraction formalism [Catani, Grazzini('07)] requires:
 - Up to $d\sigma_{NLO}^{VH+jets}$.
 - $\mathcal{H}^{VH(1)}$ and $\mathcal{H}^{VH(2)}$ [Catani, Cieri, de Florian, G.F., Grazzini('09, '12)]: contains the finite-part of the one- and two-loops amplitude $c\bar{c} \rightarrow VH$.
 - Up to $d\sigma_{NLO}^{CT}$: depends by the (universal) q_T -resummation coefficients [Bozzi, Catani, de Florian, Grazzini('09, '12)].
- $H \rightarrow b\bar{b}$ decay at NNLO calculated by [Del Duca, Duhr, Somogyi, Tramontano, Trocsanyi('15)] with CoLoRFu1NNLO method [Del Duca, Somogyi, Trocsanyi('07)].
- Fully inclusive QCD effects in the H decay taken into account by normalizing the $Hb\bar{b}$ branching fraction to the LHCHSWG-YR result.

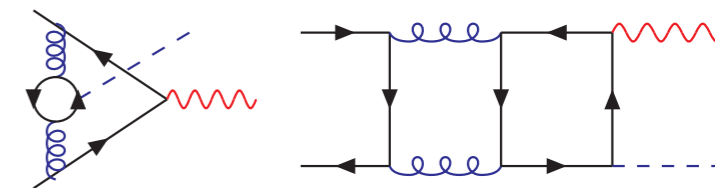
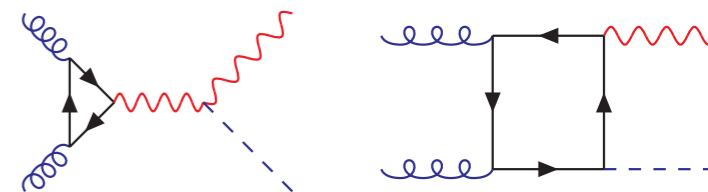
VH production and decay in full NNLO QCD



Our fully differential calculation implemented in the parton level code **HVNNLO**.

For VH prod. we have consistently included:

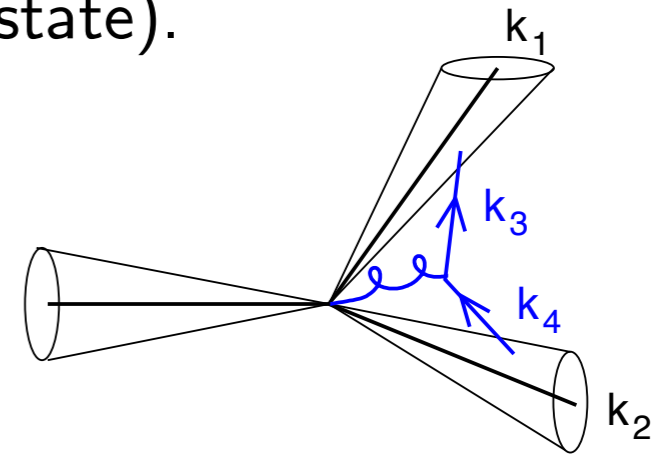
- NNLO DY-like QCD corrections
(bulk of NNLO correction for WH)
[Van Neerven et al. ('91)]
- $gg \rightarrow HZ$ top-loop $\sim g^2 \lambda_t^2 \alpha_S^2$
(non DY-like) corrections [Kniehl ('90)]
(important at the LHC due to large gg luminosity).
- NNLO top-mediated contributions
 $\sim g^3 \lambda_t \alpha_S^2$ to VH
[Brein, Harlander, Wieseemann, Zirke ('11)]
(we included only the terms calculated with the full m_t dependence)



b-quark jets identification

We are interested in the identification of the b -quark jet which originate from the Higgs boson (b -quark treated in massless approximation).

- We consistently include b -quark emissions from initial and final state partons (at NNLO up to four b -quarks in the final state).
- Standard jet alg. not infrared and collinear safe definition of flavoured jets: splitting of a gluon in a soft or collinear (massless) $b\bar{b}$ pair affect the flavour of a jet.



- **Collinear unsafety** removed by defining “b-jet” if contains $N(b) - N(\bar{b}) \neq 0$.
- **Infrared unsafety** removed by using the “flavour- k_T ” algorithm [Banfi, Salam, Zanderighi ('06)]

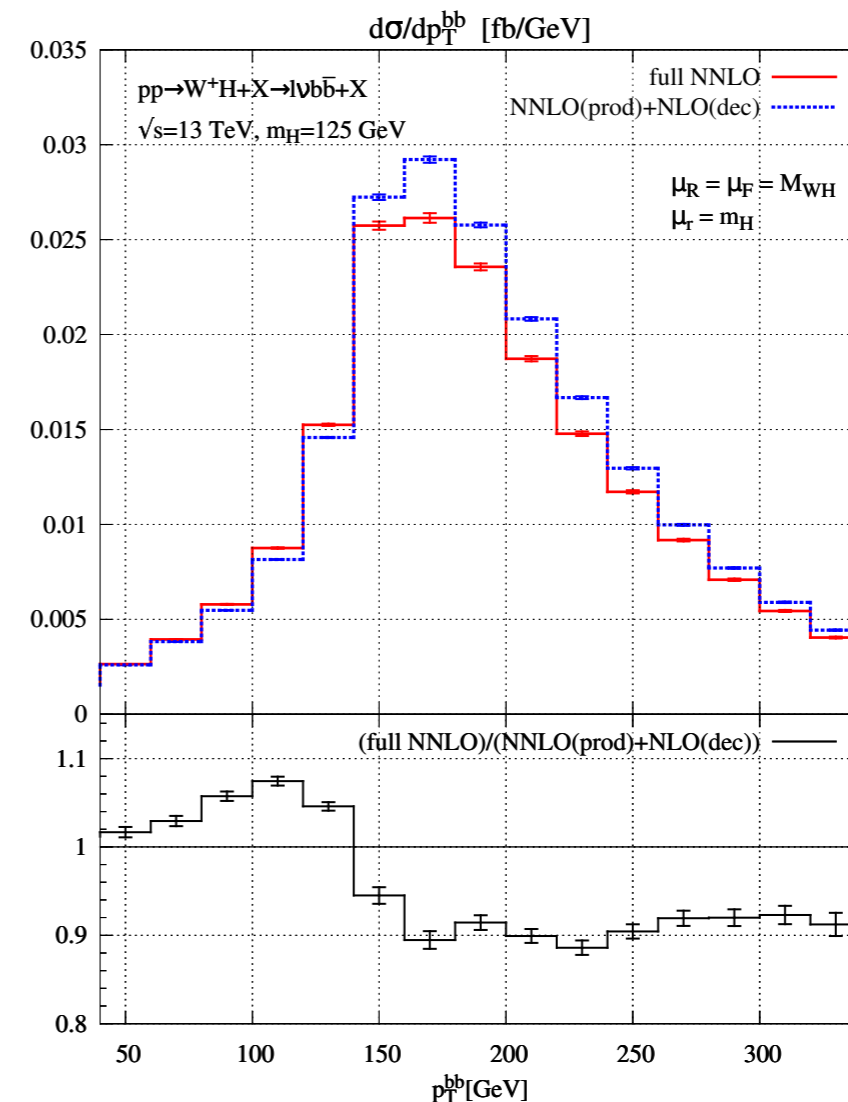
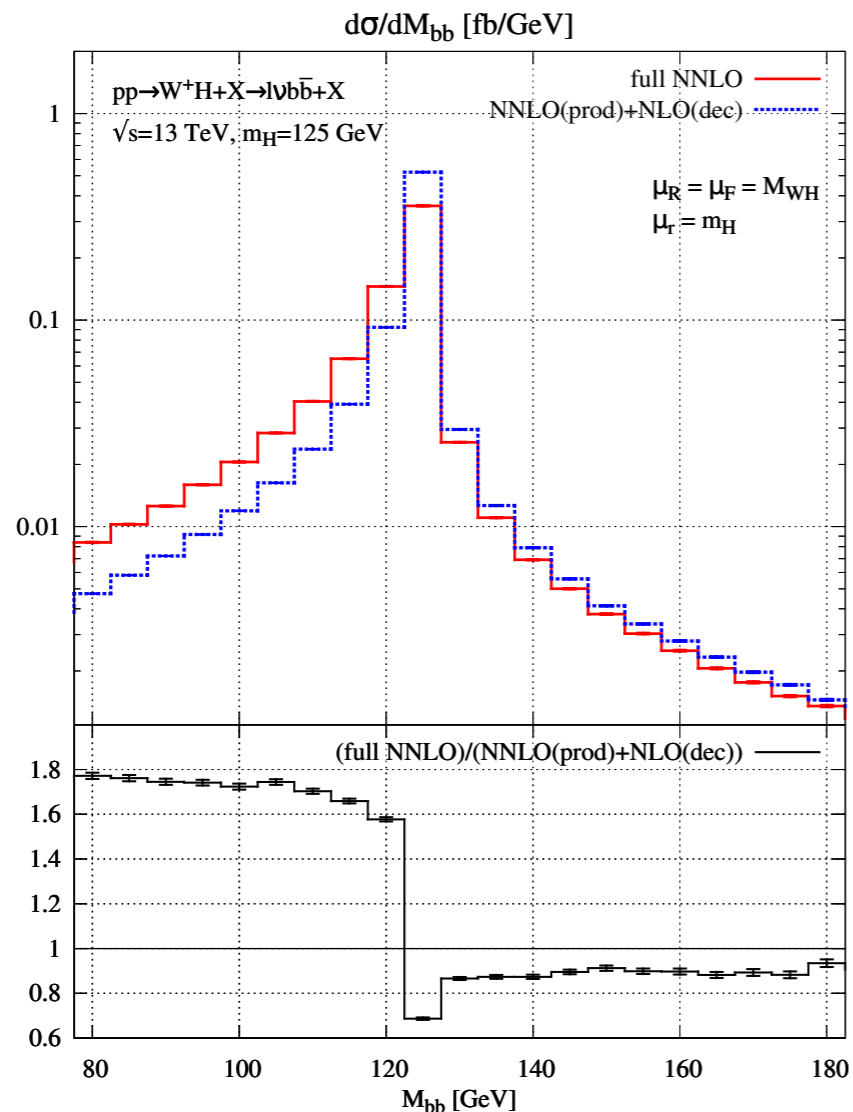
$$d_{ij}^{(F)} = (\Delta\eta_{ij}^2 + \Delta\phi_{ij}^2) \times \begin{cases} \max(k_{ti}^2, k_{tj}^2), & \text{softer of } i, j \text{ is flavoured} \\ \min(k_{ti}^2, k_{tj}^2), & \text{softer of } i, j \text{ is flavourless} \end{cases}$$

(numerical difference with respect to standard alg. small in our case).

NEW: production at full NNLO of WH(bb)

σ (fb)	NNLO(prod)+NLO(dec)	full NNLO
W^+H	4.23 ± 0.02	3.96 ± 0.02

$$K_{fact} \sim -6.5\%$$



LHC13 analysis: $p_T^l > 15$ GeV, $|\eta_l| < 2.5$, $p_T^W > 150$ GeV, 2 b -jets $p_T^b > 25$ GeV, $|\eta_b| < 2.5$, flavour- k_T $R = 0.5$.

Left panel: M_{bb} spectrum of the b -jets pair. Right panel: p_T^{bb} spectrum of the b -jets pair. Lower panels: spectra normalized to the NNLO+nlo results.

Hurdles towards $gg \rightarrow ZH$ @ NLO

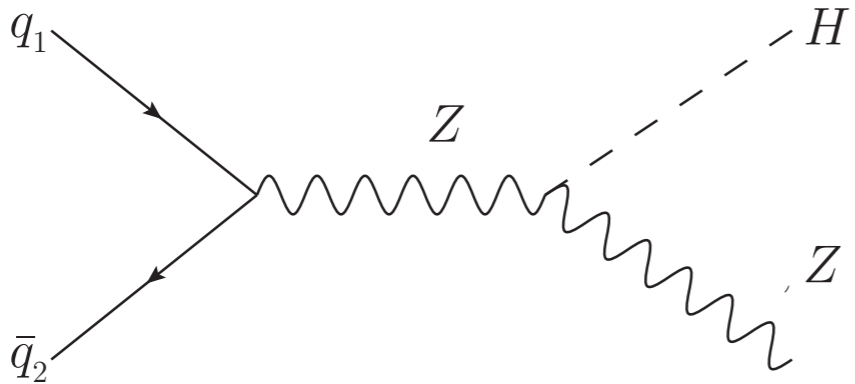
Sophia Borowka (CERN)



HXSWG VH subgroup meeting, 29th June 2017

Why the gg channel is interesting

➔ LO process is Higgs-Strahlung



➔ Drell-Yan component known up to NNLO

Hamberg, Neerven, Matsuura '91, Harlander, Kilgore '02, Brein, Djouadi, Harlander '04

➔ LO gg channel enters at NNLO with $\sim 10\%$ Brein, Harlander, Zirke '12

➔ gluon fusion scale uncertainty large ($\sim 30\%$),
dominates overall $pp \rightarrow ZH$ uncertainty at NNLO

➔ $gg \rightarrow ZH$ @NLO with full top-mass dependence

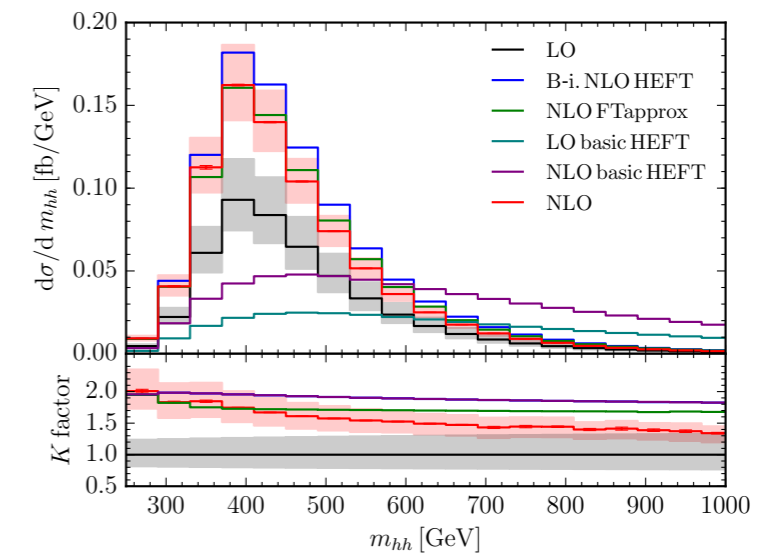
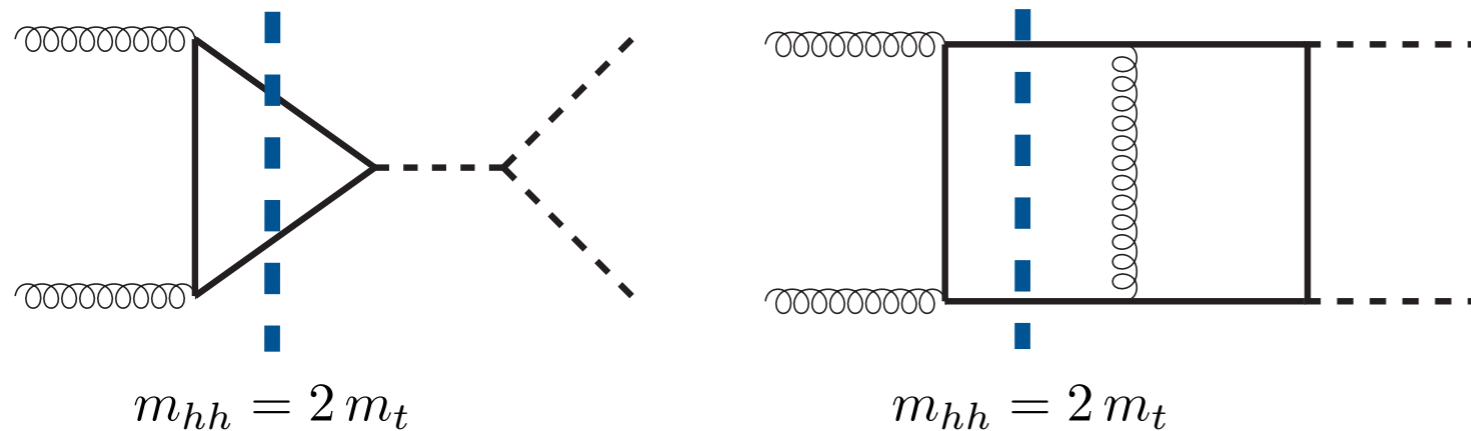
Computation available in the large mass expansion finding
large K-factor for the inclusive cross section and strong hints
for a large k-factor also at differential level

[Altenkamp, Dittmaier, Harlander, Rzehak Zaire 2013]

HH: Differences between SM and HEFT

➔ what happens in the exact Standard Model

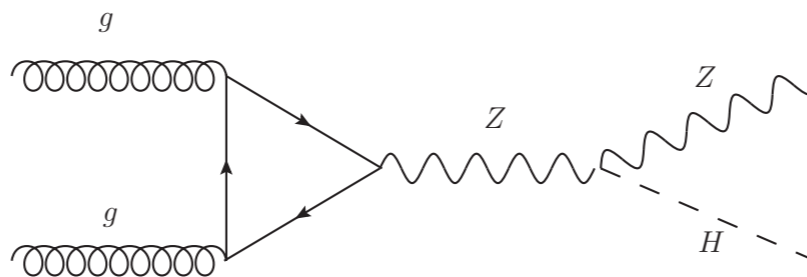
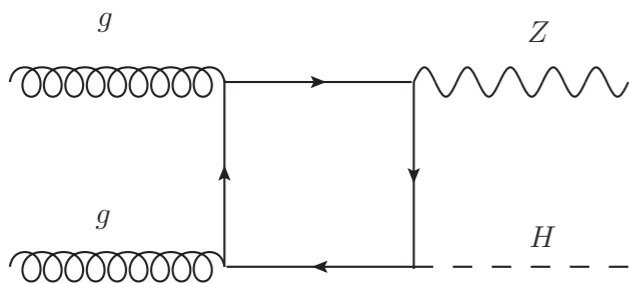
➔ threshold effects can show up



SB, N. Greiner, G. Heinrich, S.P. Jones,
M. Kerner, J. Schlenk, U. Schubert, T. Zirke '16

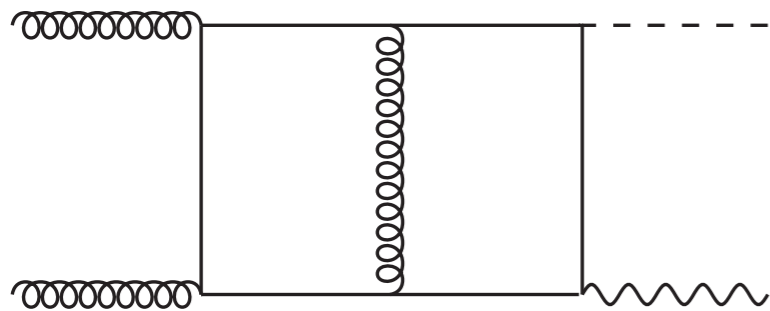
gg->ZH diagrams

Leading Order:

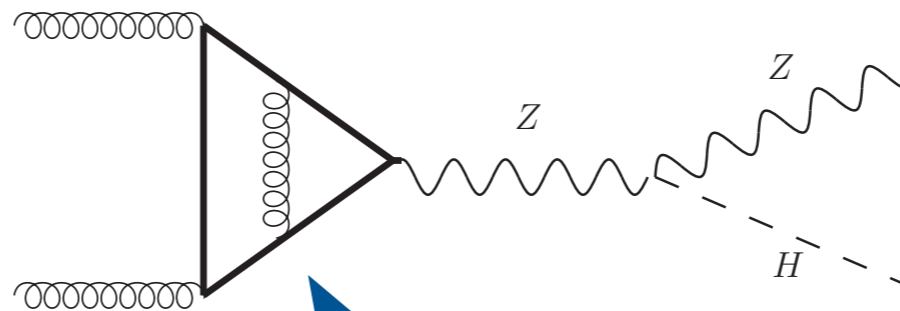


Dicus, Kao '88; Kniehl '90

Exact virtual NLO part:



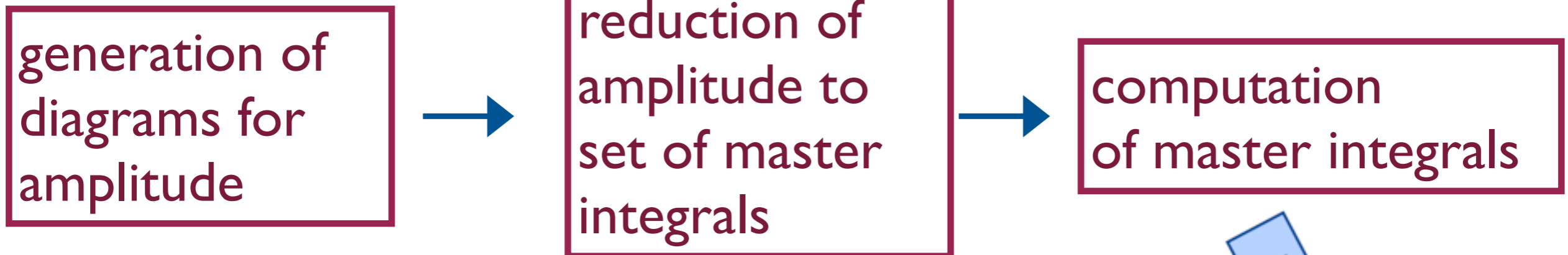
not known yet



master integrals known from
Gehrmann, Huber, Maitre '05

Exact real radiation for NLO by: Hespel, Maltoni, Vryonidou '15

Schematic gg->HH setup (virtual NLO)



- reduction programs: FIRE, KIRA, LiteRed, REDUZE
Smirnov '15; Maierhöfer, Usovitsch, Uwer '17; Lee '13; von Manteuffel, Studerus '12
- REDUZE can generate quasi-finite basis



Important for success:

- + use quasi-finite basis Panzer '14; von Manteuffel, Panzer, Schabinger '14
- + use QMC Dick, Kuo, Sloan '13; Li, Wang, Zan, Zhao '15;
- + only integrate up to necessary accuracy
(2 form factors for HH, 3% for one form factor, $\approx 10\%$ for the other, depending on the ratio of the two)

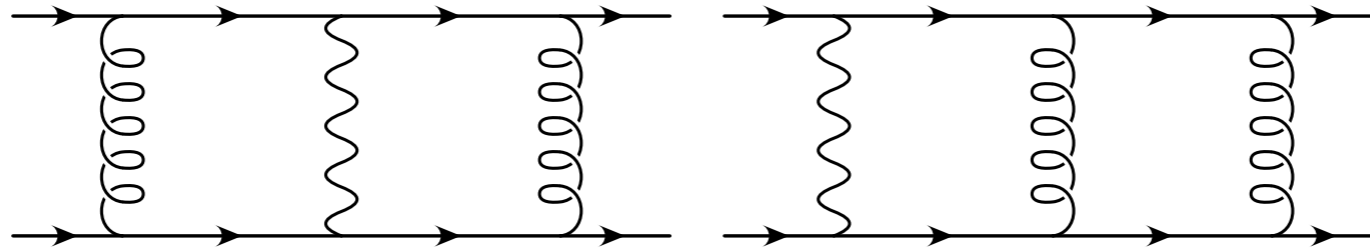
SB, G. Heinrich, S. Jahn, S.P. Jones, M. Kerner, J. Schlenk, T. Zirke '17

Problems that may occur in $gg \rightarrow ZH$

- ➔ additional mass scale makes reduction much more involved
- ➔ if reduction not available no transformation into quasi-finite basis possible
- ➔ if double-box integrals are not finite, numerical convergence significantly worse
- ➔ form factors may be of similar importance (high accuracy also needed for most complicated integrals)
- ➔ numerical convergence in general slower the more scales are involved

A possible recipe that might help in the reduction to master integrals

- The number of scales is the limiting factor for the reduction program to work
- numerics might help to reduce the complexity of the reduction algorithms
 - ▶ Example: t-channel single top at NNLO



[Assadsolimani, Kant, Tausk, Uwer 2014]

- ▶ reduction of double box diagrams successfully achieved exploiting the relation:

$$m_t^2 \approx \frac{14}{3} m_W^2 \quad m_W = 80.385 \pm 0.015 \text{ GeV}/c^2 \quad \longrightarrow \quad m_t \approx 173.65 \text{ GeV}/c^2$$

$$m_t = 173.34 \pm 0.27 \text{ (stat)} \pm 0.71 \text{ (syst)} \text{ GeV}/c^2$$

- for HZ one could use for example:

$$m_Z : m_H : m_t \approx 8 : 11 : 15$$

$$91.1876 : 125 : 173.3 \quad \longrightarrow \quad 91.1876 : 125.4 : 171.0$$

leading to O(1%) error on the correction

Conclusion

- Since the publication of YR4 there has been substantial progress
- Event generation will be available including NNLOPS QCD effects for both WH and ZH
- NLO EW corrections available at the level of NLOPS for both VH and VHJ(-MINLO) processes
- At least at the level of distributions the improvements of the last two items could be combined
- VH(bb)@NNLO computed including higher order corrections for the decay and the product of nlo corrections from production and decay, substantial effects, comparison with MC expectation needed.
- Still no progress on ggZH@NLO, Sophia and collaborators has shown that numerics offers a possible road, but needs of big efforts...