



CKM Fits

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- Phase invariant parameterisation conserving the CKM matrix unitarity at any order in λ

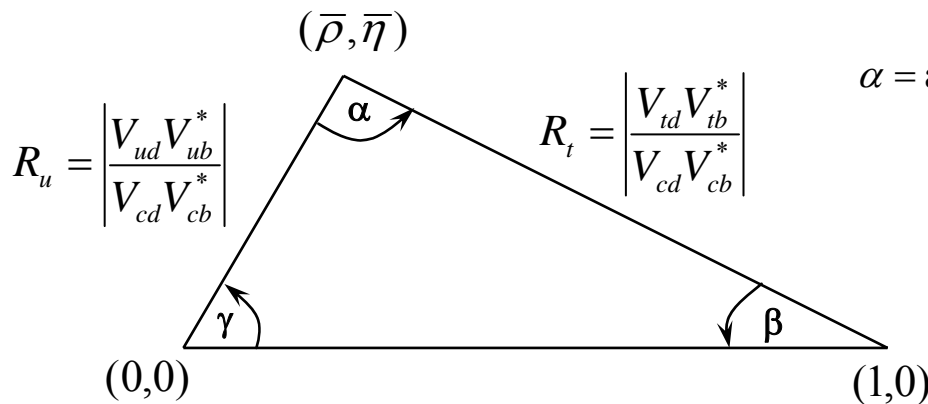
Wolfenstein parameterisation with Jarlskog like phase invariants as in EPJ C41,1-131 (2005)

4 free parameters, $A, \lambda, \bar{\rho}$ and $\bar{\eta}$, taken such that:

$$\lambda = \frac{|V_{us}|}{\sqrt{|V_{ud}|^2 + |V_{us}|^2}}, \quad A\lambda^2 = \frac{|V_{cb}|}{\sqrt{|V_{ud}|^2 + |V_{us}|^2}} \quad \text{and} \quad \bar{\rho} + i\bar{\eta} = -\frac{V_{ud}V_{ub}^*}{V_{cd}V_{cb}^*}, \quad \text{with} \quad V_{\text{CKM}} = \begin{bmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{bmatrix}$$

- λ is measured from $|V_{ud}|$ and $|V_{us}|$ in superallowed nuclear β -decays and (semi)leptonic K decays, resp.
- A is determined from $|V_{cb}|$ and λ .
- $\bar{\rho} + i\bar{\eta}$ is to be determined from angles and sides measurements of the B_d unitarity triangle.

B_d Unitarity Triangle (UT)



$$\alpha = \arg \left[-\frac{V_{td}V_{tb}^*}{V_{ud}V_{ub}^*} \right], \quad \beta = \arg \left[-\frac{V_{cd}V_{cb}^*}{V_{td}V_{tb}^*} \right], \quad \gamma = \arg \left[-\frac{V_{ud}V_{ub}^*}{V_{cd}V_{cb}^*} \right]$$

$$\beta_s = -\arg \left[-\frac{V_{cs}V_{cb}^*}{V_{ts}V_{tb}^*} \right]$$

in B_s

■ CKM matrix fit

- + Use **Frequentist Hypothesis testing** to build statistical **significance** (p-value) functions from which estimates and **confidence intervals** are obtained; test statistic = Maximum Likelihood Ratio = $\Delta\chi^2$.
- + Dedicated **RFit** scheme for the treatment of theoretical systematics.

- **data = weak** \otimes **QCD** \Rightarrow need for hadronic inputs; **often LQCD**: Our Own Average (**OOA**)

Observable category	Experimental sources	Theory methods & inputs
$ V_{ud} $	Superallowed β decays	Towner & Hardy, PRC77 025501 (2008)
$ V_{us} $	K_{l3} , WA FlaviaNet: KLOE	$f_+^{K\pi}(0)=0.964(5)$ (most precise: RBC-UKQCD)
$ V_{cb} $	HFAG incl.+excl. $B \rightarrow X_c l \nu$	$40.59(38)(58) \times 10^{-3}$
$ V_{ub} _{SL}$	HFAG incl.+excl. $B \rightarrow X_u l \nu$	OOA (spec. uncert. budget.): $3.87(9)(46) \times 10^{-3}$
$B[B \rightarrow \tau \nu] (\rightarrow V_{ub})$	2008 WA: BaBar & Belle	OOA $f_{B_s}/f_{B_d}=1.196(8)(23)$ & $f_{B_s}=228(3)(17)$
$\Delta m_d (\rightarrow V_{td})$	HFAG WA \bar{B}_d - B_d mixing	OOA $B_{B_s}/B_{B_d}=1.05(2)(5)$ + f_{B_s}/f_{B_d}
$\Delta m_s (\rightarrow V_{ts})$	CDF B_s - \bar{B}_s mixing	OOA $B_{B_s}=1.23(3)(5)$ + f_{B_s}
$ \varepsilon_K $	K^0 - \bar{K}^0 , PDG 2008: KLOE, NA48, KTeV	PDG param. (Buchalla et al. '96) + OOA $B_K=0.721(5)(40)$
α / ϕ_2	WA $\pi\pi$, $\rho\pi$, pp New!	isospin SU(2), Gronau & London, PRL65, 3381 (1990)
β_{cc} / ϕ_1	WA HFAG charmonium	0.671(23)
γ / ϕ_3	WA HFAG $B^- \rightarrow D^{(*)}K^{(*)-}$	GLW / ADS / GGSZ

CP

A, λ

R_u, R_t

Modulus
and sides
from rates

CP

$\bar{\rho}, \bar{\eta}$

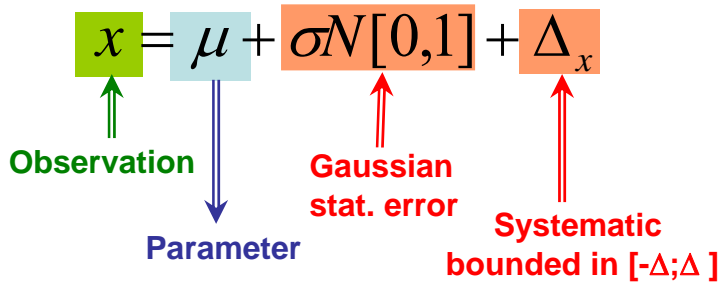
Angles from
phases in
interferences

RFit scheme

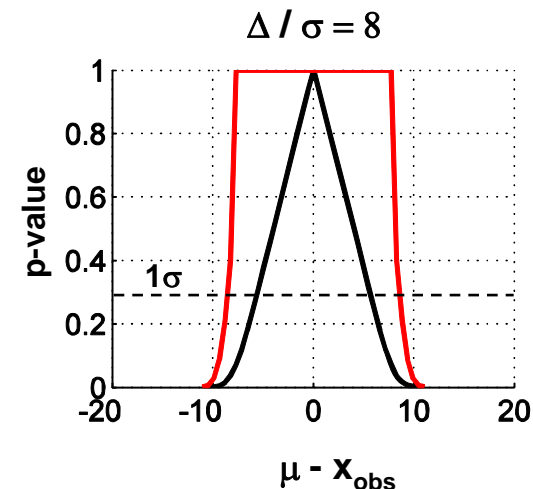
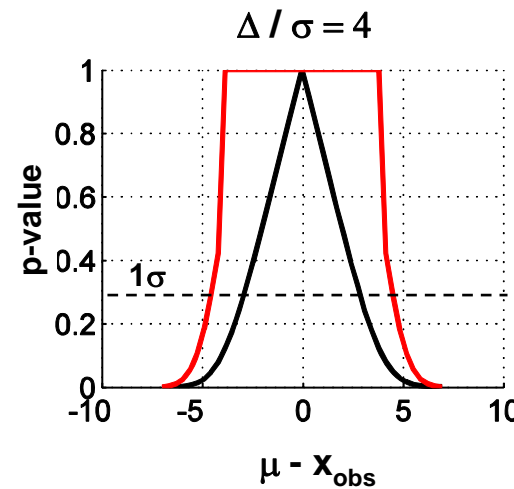
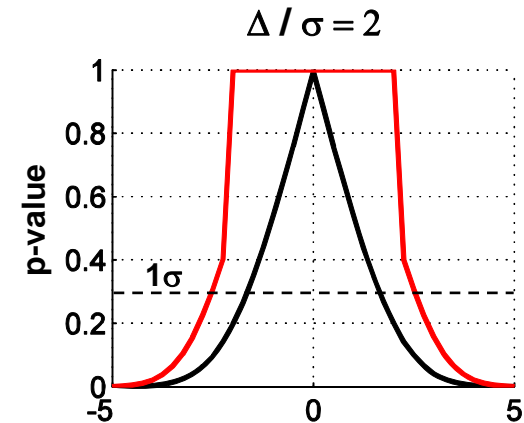
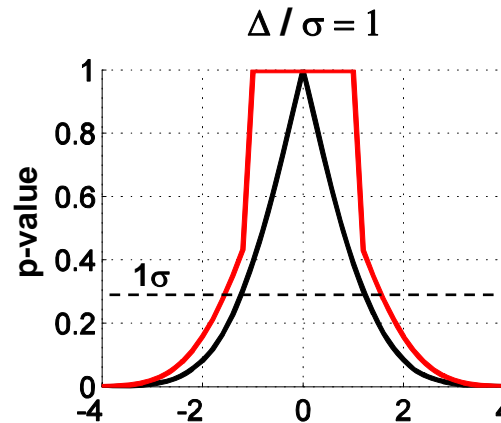
Theoretical systematics are considered as additional **nuisance parameters** bounded over a confident enough range. Considering the worst case –supremum–, on the latter interval the significance is flat.

Note that this result is very different from what one would get from a statistical modelling of the systematic (ex: uniform over the range)

Simple illustrative example allowing analytical resolution:



$$p\text{-value} = \begin{cases} 1 & \text{if } x - \mu \in [-\Delta; \Delta] \\ \frac{1}{2} \left(\operatorname{erfc}\left[\frac{|x - \mu| + \Delta}{\sqrt{2}\sigma}\right] + \operatorname{erfc}\left[\frac{|x - \mu| - \Delta}{\sqrt{2}\sigma}\right] \right) & \text{elsewhere} \end{cases}$$



— Gaussian pdf + uniform pdf for systematic
— Gaussian pdf + parametric systematic

- More and more **accurate theoretical predictions** (ex: $f_{B_s}/f_{B_d} \sim 2-3\%$) but various methods, results and error estimates depending on collaborations. Need to **combine these results**; several methods also ☹️

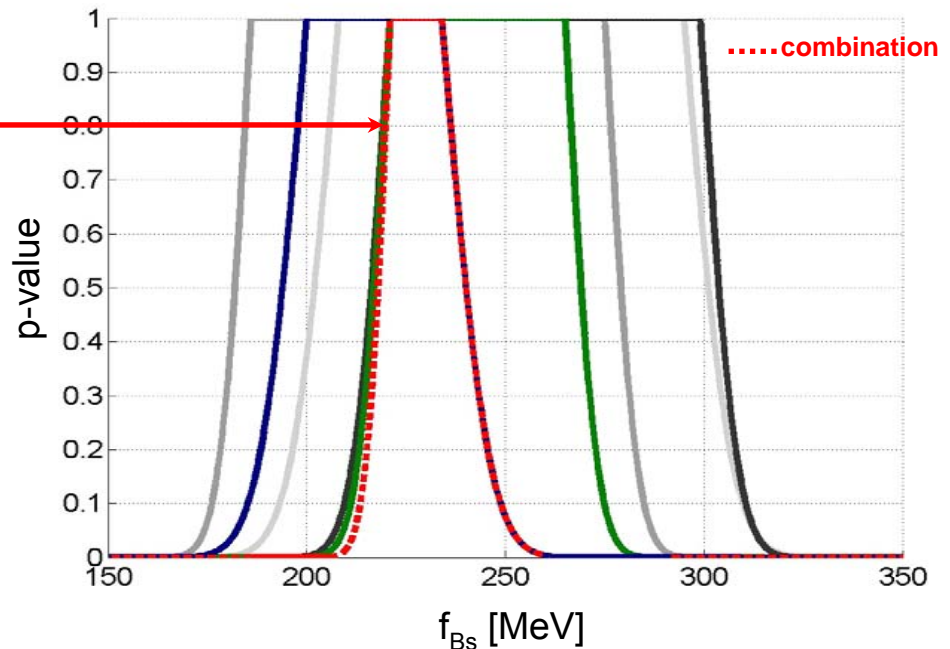
⇒ For now we perform **our own average** using an **algorithmic procedure**

Our Own Average: Educated *RFit* scheme illustrated here with f_{B_s}

1) From selected LQCD results estimate f_{B_s} **central value in the *RFit* scheme**, distinguishing statistic and systematic contributions to uncertainties.

2) Perform and **educated combination of uncertainties**; Not more nor less accurate than the most precise individual LQCD prediction.

$$\Rightarrow f_{B_s} = 228 \pm 3 \pm 17 \text{ MeV}$$



For more details:

- + V. Tisserand (CKMfitter Group), Moriond EW 2009 proceedings [arXiv:0905.1572];
- + http://ckmfitter.in2p3.fr/plots_Moriond09/ckmEval_results_Moriond09.pdf;

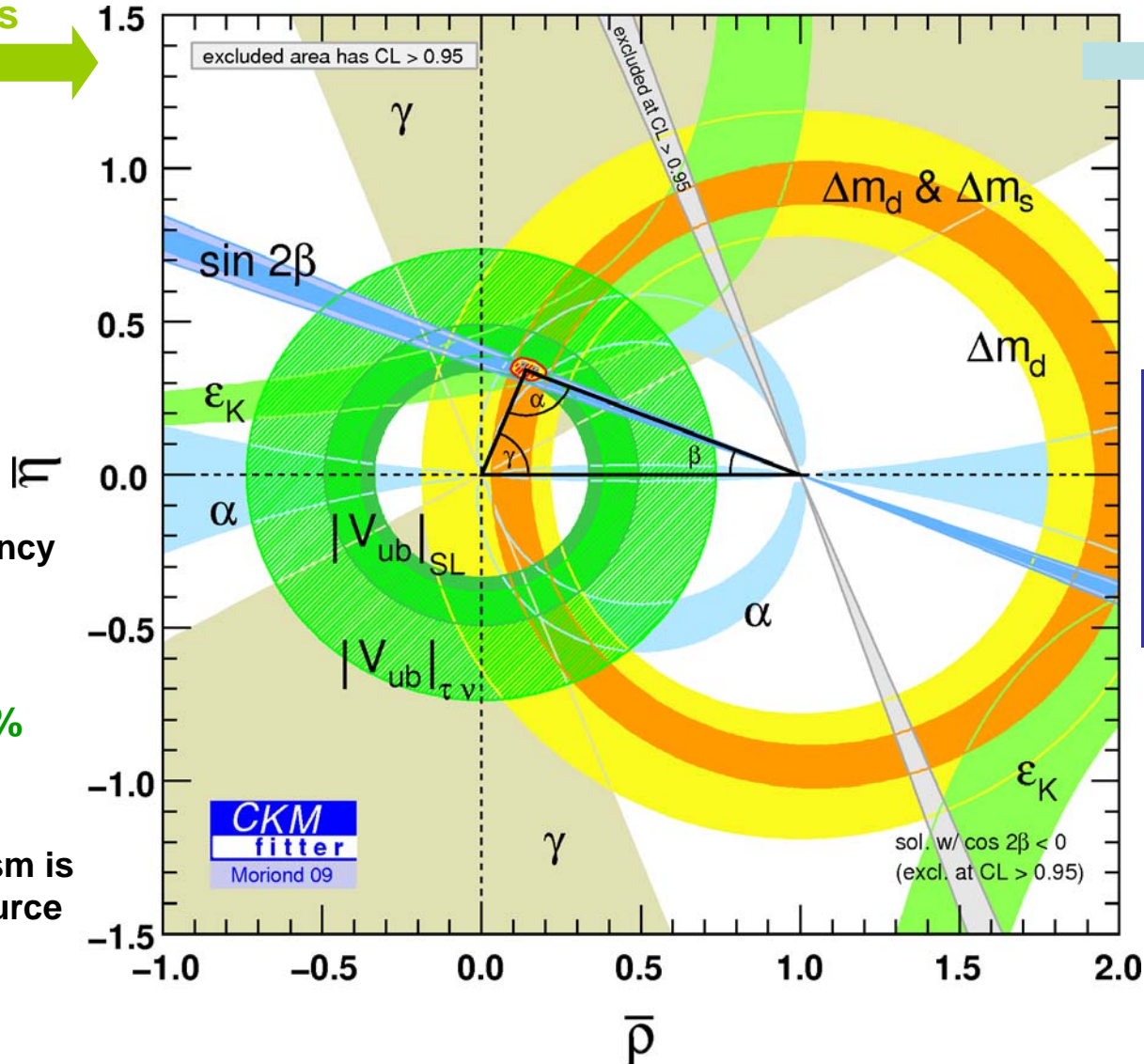
Observables

- $|V_{ud}|, |V_{us}|$
- $|V_{cb}|, |V_{ub}|$
- $B[B \rightarrow \tau \nu]$
- $\Delta m_d, \Delta m_s$
- $|\varepsilon_K|$
- $\alpha, \sin(2\beta), \gamma$

Overall consistency
at 95% CL

Global Fit
p-value $\approx 45\%$
(0.8 σ)

The KM mechanism is
the dominant source
of \mathcal{CP} in B's



$$\bar{\rho} = 0.139^{+0.025}_{-0.027}$$

$$\bar{\eta} = 0.341^{+0.016}_{-0.015}$$

(1 σ interval)

Fit dominated by
 $\sin(2\beta)$, $\Delta m_d/\Delta m_s$ and
 α/ϕ_2 (!). Excellent
agreement between
these 3 inputs.
p-value $\approx 95\%$ (0.1 σ)

■ A overview of numerical results

Wolfenstein parameters and Jarlskog invariant:

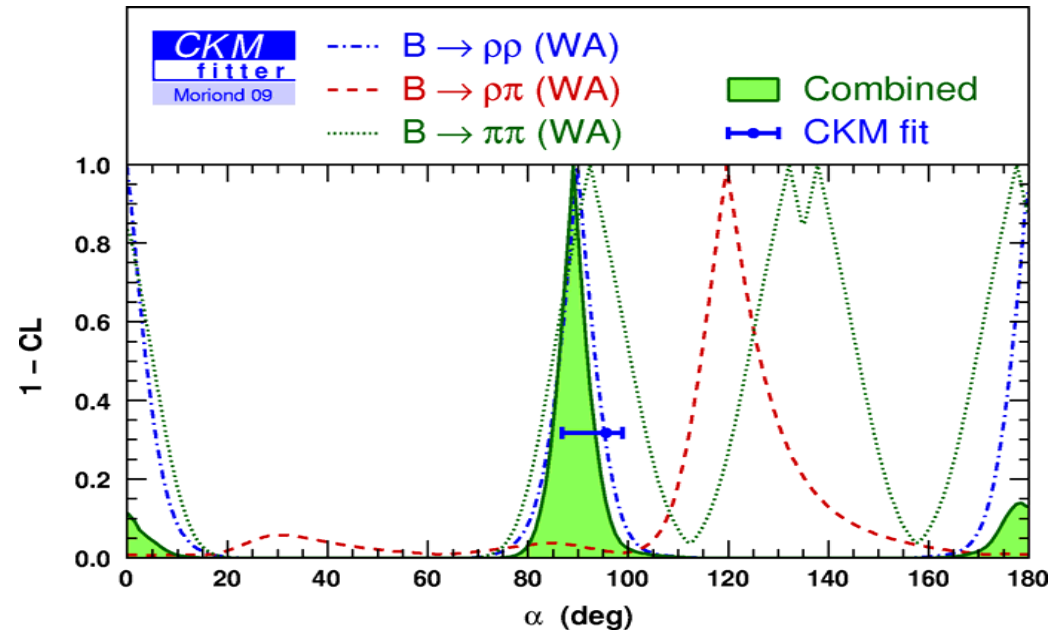
Observable	Central $\pm 1 \sigma$	$\pm 2 \sigma$	$\pm 3 \sigma$
A	0.8116 [+0.0097 -0.0241]	0.812 [+0.019 -0.036]	0.812 [+0.029 -0.045]
λ	0.22521 [+0.00082 -0.00082]	0.2252 [+0.0016 -0.0016]	0.2252 [+0.0025 -0.0025]
ρ bar	0.139 [+0.025 -0.027]	0.139 [+0.053 -0.040]	0.139 [+0.073 -0.052]
η bar	0.341 [+0.016 -0.015]	0.341 [+0.032 -0.025]	0.341 [+0.048 -0.034]
J [10^{-5}]	2.92 [+0.15 -0.15]	2.92 [+0.30 -0.19]	2.92 [+0.45 -0.23]

UT angles and sides:

Observable	Central $\pm 1 \sigma$	$\pm 2 \sigma$	$\pm 3 \sigma$
$\sin 2\alpha$	-0.02 [+0.15 -0.13]	-0.02 [+0.22 -0.26]	-0.02 [+0.28 -0.35]
$\sin 2\alpha$ (meas. not in the fit)	-0.20 [+0.30 -0.11]	-0.20 [+0.41 -0.17]	-0.20 [+0.48 -0.23]
$\sin 2\beta$	0.684 [+0.023 -0.021]	0.684 [+0.046 -0.035]	0.684 [+0.068 -0.049]
$\sin 2\beta$ (meas. not in the fit)	0.817 [+0.026 -0.040]	0.817 [+0.039 -0.114]	0.817 [+0.052 -0.171]
$ \sin (2\beta+\gamma) $	0.934 [+0.023 -0.030]	0.934 [+0.039 -0.051]	0.934 [+0.049 -0.071]
α [deg]	90.6 [+3.8 -4.2]	90.6 [+7.5 -6.3]	90.6 [+10.2 -8.2]
α [deg] (meas. not in the fit)	95.6 [+3.3 -8.8]	95.6 [+5.2 -11.8]	95.6 [+6.8 -13.9]
α [deg] (dir. meas.)	89.0 [+4.4 -4.2]	89.0 [+9.1 -8.3] 178.3 [+1.7 -5.6] 0 [+5 -0]	89 [+21 -13] 178.3 [+1.7 -13.8] 0 [+12 -0]
β [deg]	21.58 [+0.91 -0.81]	21.6 [+1.8 -1.4]	21.6 [+2.8 -1.9]

Many more results and plots are available at: http://ckmfitter.in2p3.fr/plots_Moriond09

α/ϕ_2 from $B \rightarrow \pi\pi, \rho\pi, \rho\rho$

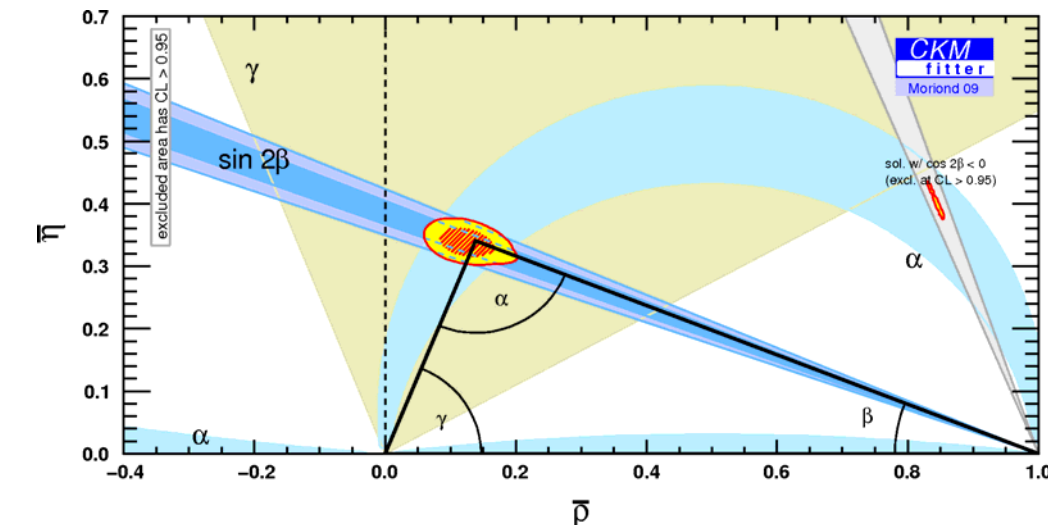


■ Get α From a global fit to rates and CP asymmetries further **assuming SU(2) isospin triangular relations** between amplitudes: $A^{+-} + \sqrt{2}A^{00} = \sqrt{2}A^{+0}$, to simultaneously extract the additional strong phase; [Gronau & London PRL65, 3381 (1990)]

$$\alpha = (89.0^{+4.4}_{-4.2})^\circ \quad \text{p-value} \approx 86\% \text{ (0.2 } \sigma)$$

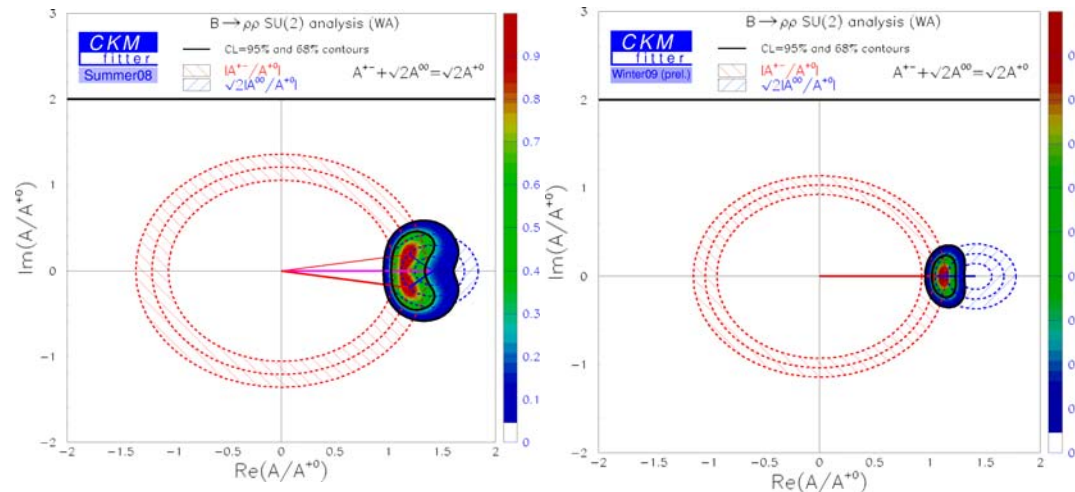
The fit is **dominated by $B \rightarrow \rho\rho$** analysis

■ **New measurement of $B[B^+ \rightarrow \rho^+\rho^0]$ from BaBar** (PRL102, 141802 (2009)) which strongly constrains the isospin triangles and hence α extraction.



α has become a precise measurement @ **5%** comparable to β @ **4.2%**
 $\sin(2\beta_{cc}) = 0.671(23)$ [HFAG]

$\rho^+\rho^0$ BaBar update: normalisation A^{+0} increased



previously 2 folded solution now degenerated
 \Rightarrow increased accuracy by a factor ~ 2 on α from $B \rightarrow \rho\rho$

α from $B \rightarrow \rho\rho$
 Old summer 08

$$\alpha = (90.9^{+6.7}_{-14.9})^\circ$$

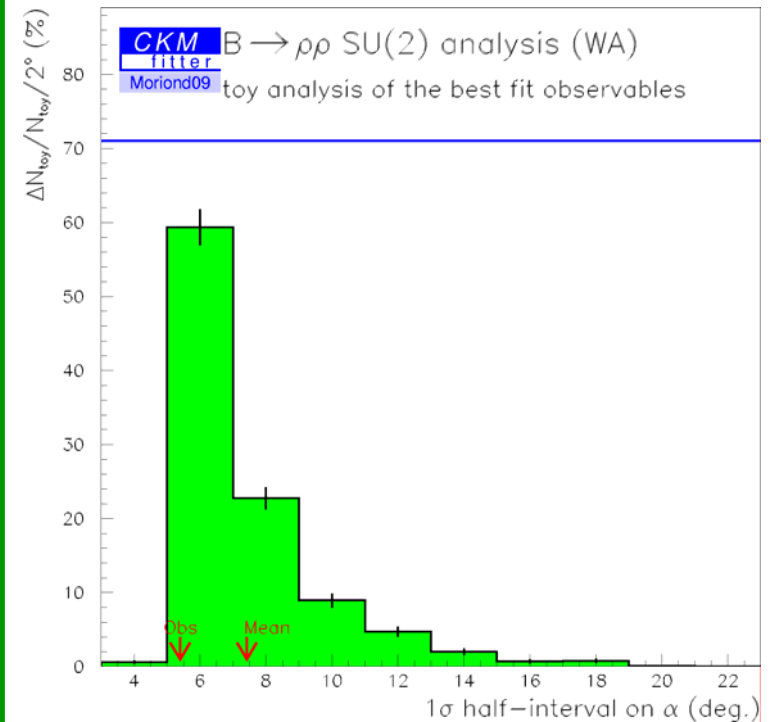
$$\Delta\alpha = (0.5^{+12.6}_{-5.5})^\circ$$

α from $B \rightarrow \rho\rho$
 Winter 09

$$\alpha = (89.9 \pm 5.4)^\circ$$

$$\Delta\alpha = (1.4 \pm 3.7)^\circ$$

How lucky are we?



Study from **toy Monte-Carlo** using fit best guess parameter values (plugin):

\Rightarrow **Only 34% of the toy MC triangles close**

\Rightarrow Average toy error is 7.5° to be compared to observed value of 5.4° . **68% of toys have larger error than the data.**

- Various possible sources: QCD ($\mu \neq md$), QED, ... amount to $1-3^\circ$ [J. Zupan Nucl Phys Proc Suppl 170 33 2007]

Largest effect presumably comes from **finite width: $\Gamma_\rho \neq 0$ allows $\Delta I=1$ transitions**. Breaking as $\propto O\left(\frac{\Gamma_\rho^2}{m_\rho^2}\right) \approx 4\%$

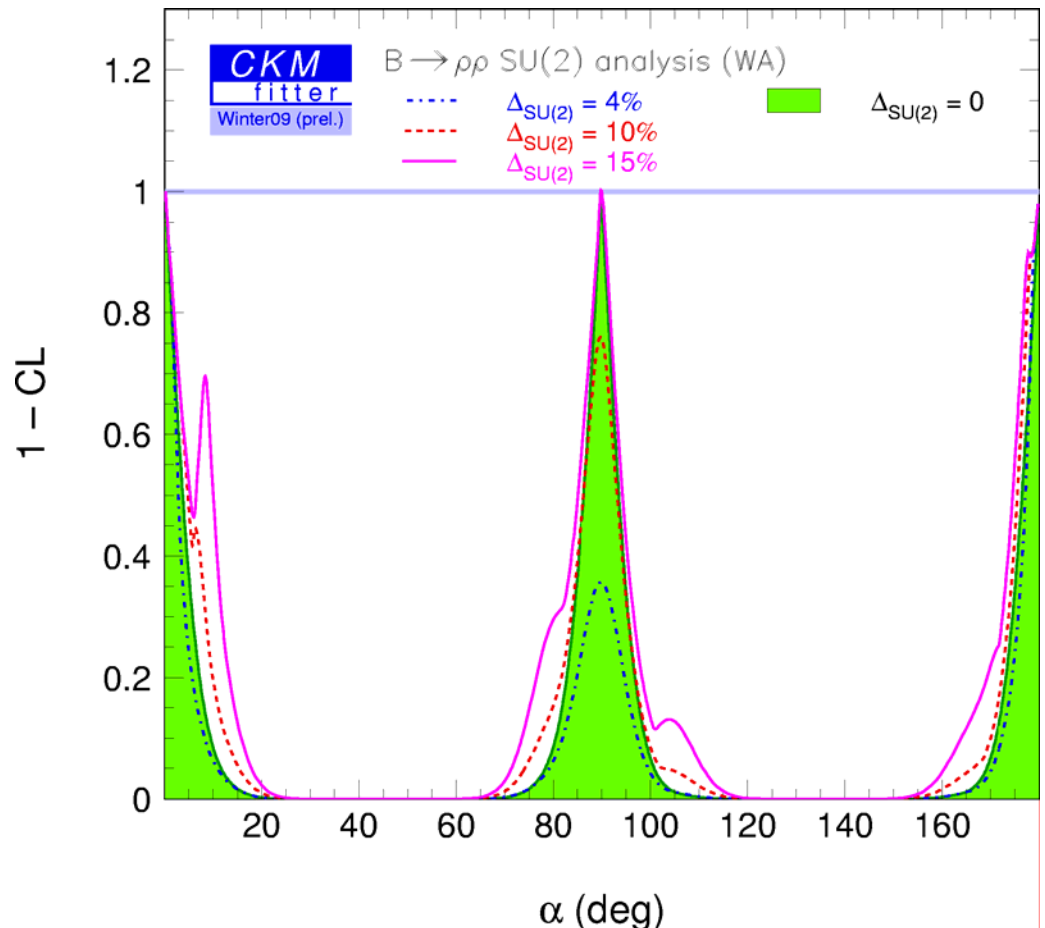
■ Phenomenological study

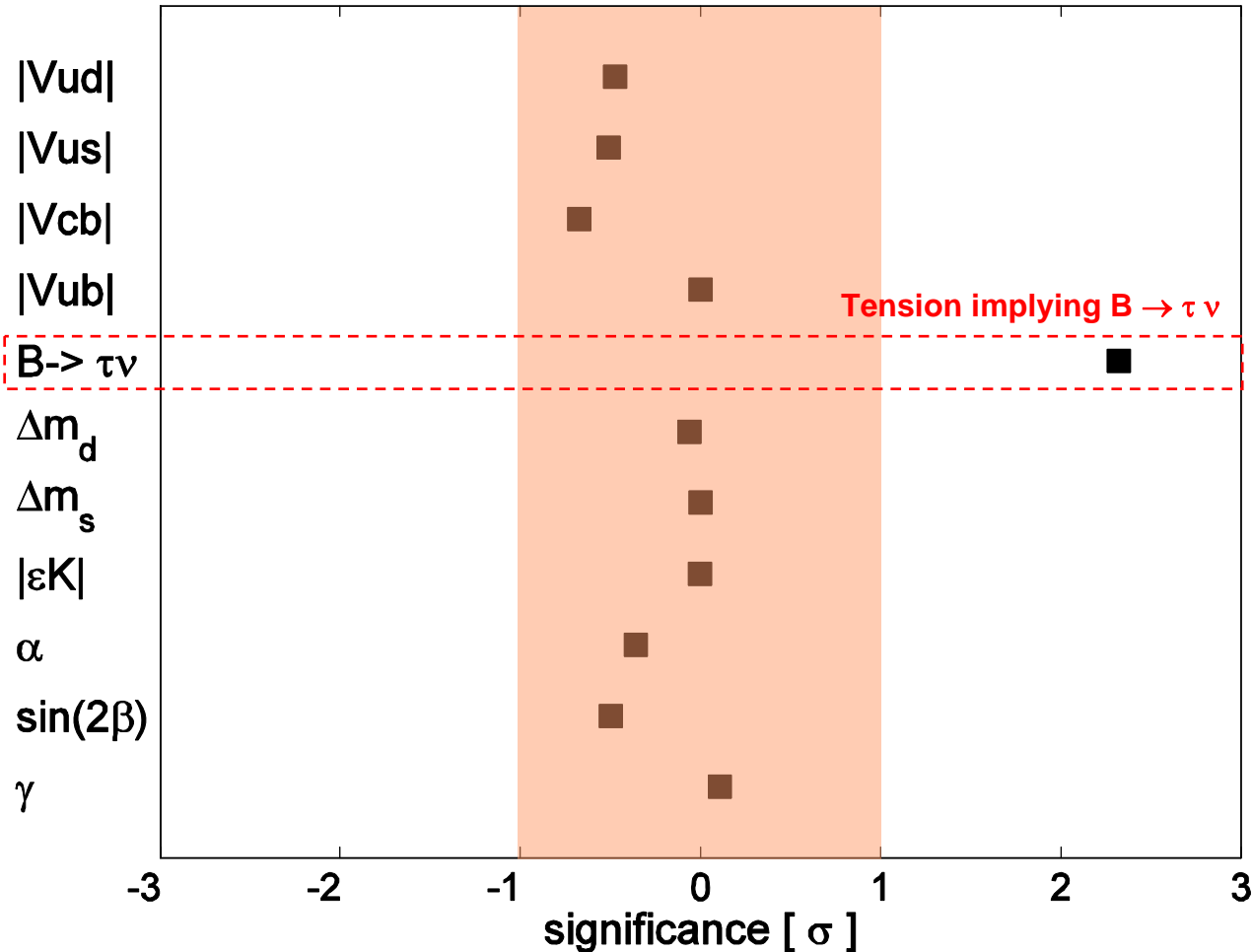
Test **4%**, 10% and 15% violation of the triangular relation with arbitrary additional amplitude:

$$\Rightarrow A^{+0} \rightarrow A^{+0} + \Delta A^{+0}$$

Small values of isospin breaking do not change the pattern.

Small impact on $\pi\pi$, $\rho\pi$, $\rho\rho$ WA combination; only visible at 95% CL
A nice constraint in any case





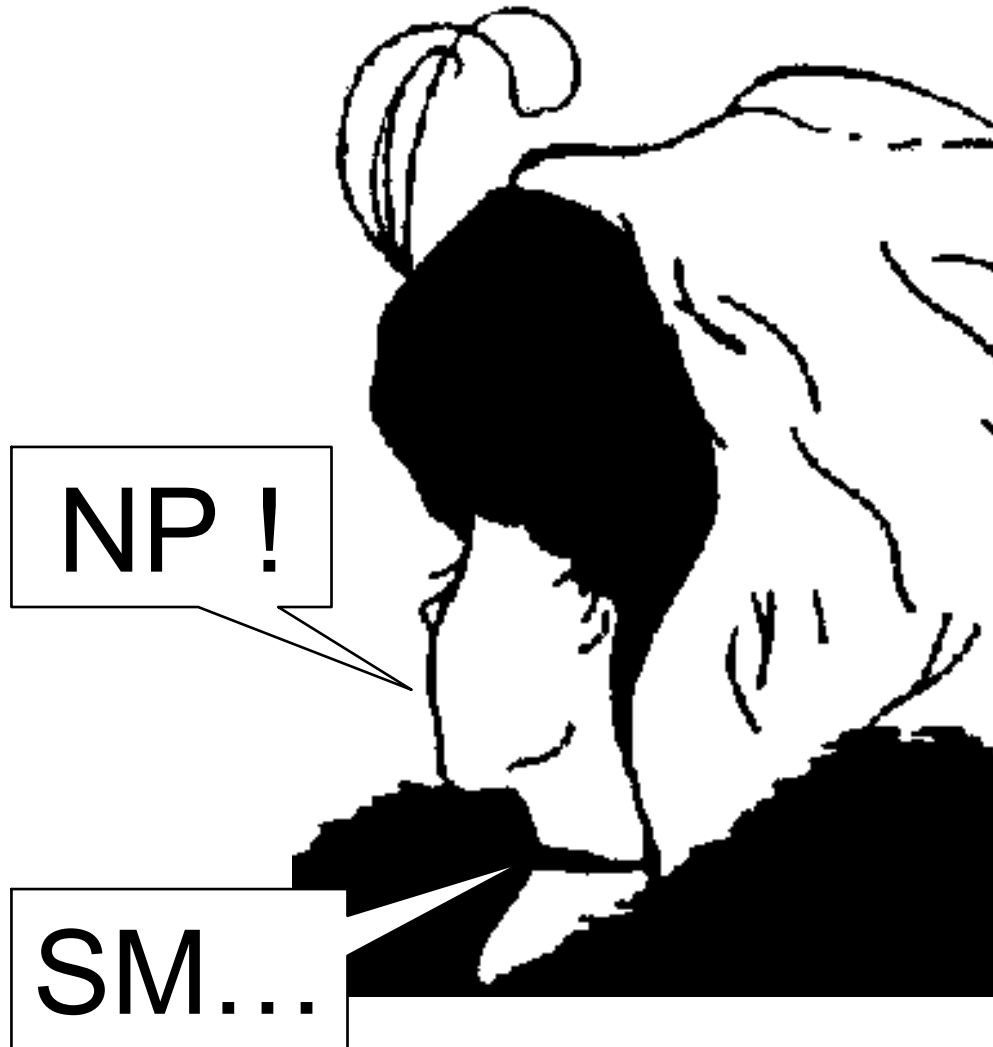
H₀: the observable central value is given by the global fit best guess. The measurement is polluted by a known Gaussian error σ and a systematic lying in a known range $\pm\Delta$.

What is the significance of the observed value?

Caution: In the *RFit* scheme, significance close to 0σ means that the measurement lies within the systematics range.

Conservative approach used here. Most measurements can be accommodated within the range of hadronic uncertainties

Any Tensions ?



- The *significant* difference between $|V_{ub}|$ derived from inclusive and exclusive $B \rightarrow X_u l \nu$ semileptonic measurements drops when treating systematics within the *Educated RFit* scheme

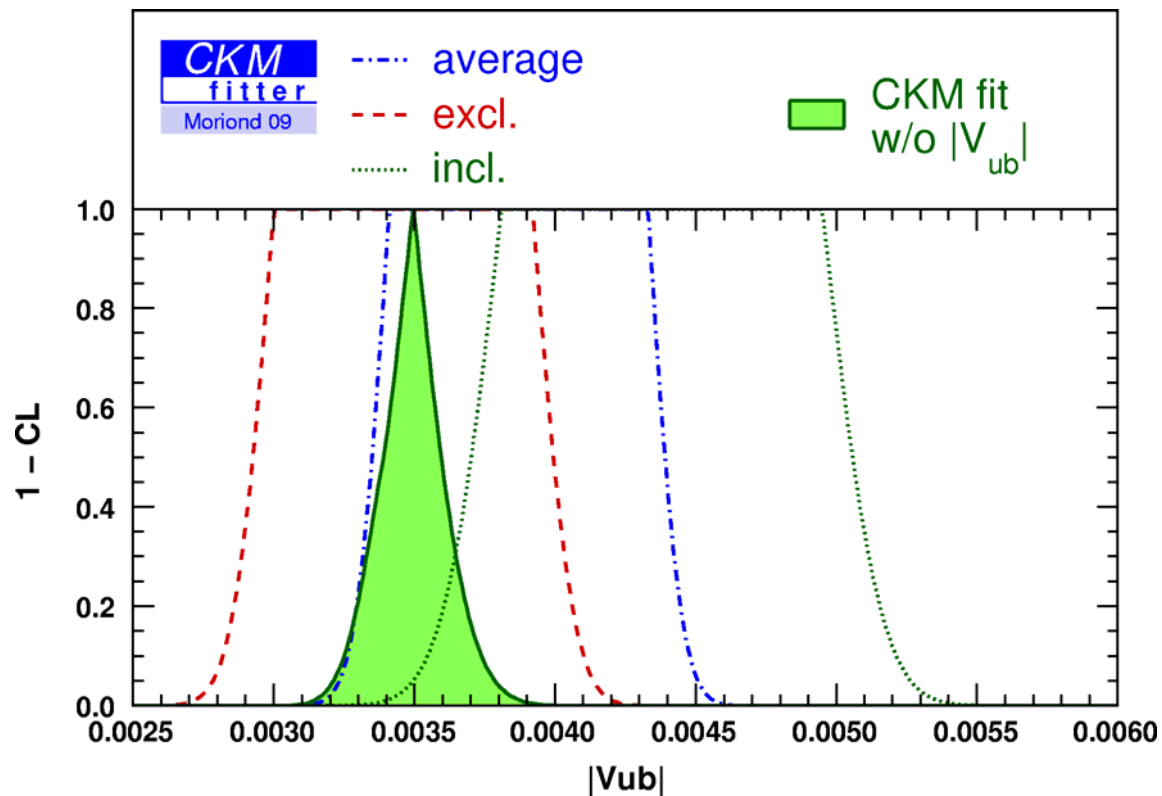
OOA adapted from HFAG summer 08

$$|V_{ub}|_{incl.} = 4.38(16)(57) 10^{-3}$$

$$|V_{ub}|_{excl.} = 3.46(11)(46) 10^{-3}$$

$$|V_{ub}| = 3.87(9)(46) 10^{-3}$$

$\Delta|V_{ub}| = 0.92$, consistent with **OOA** error budget



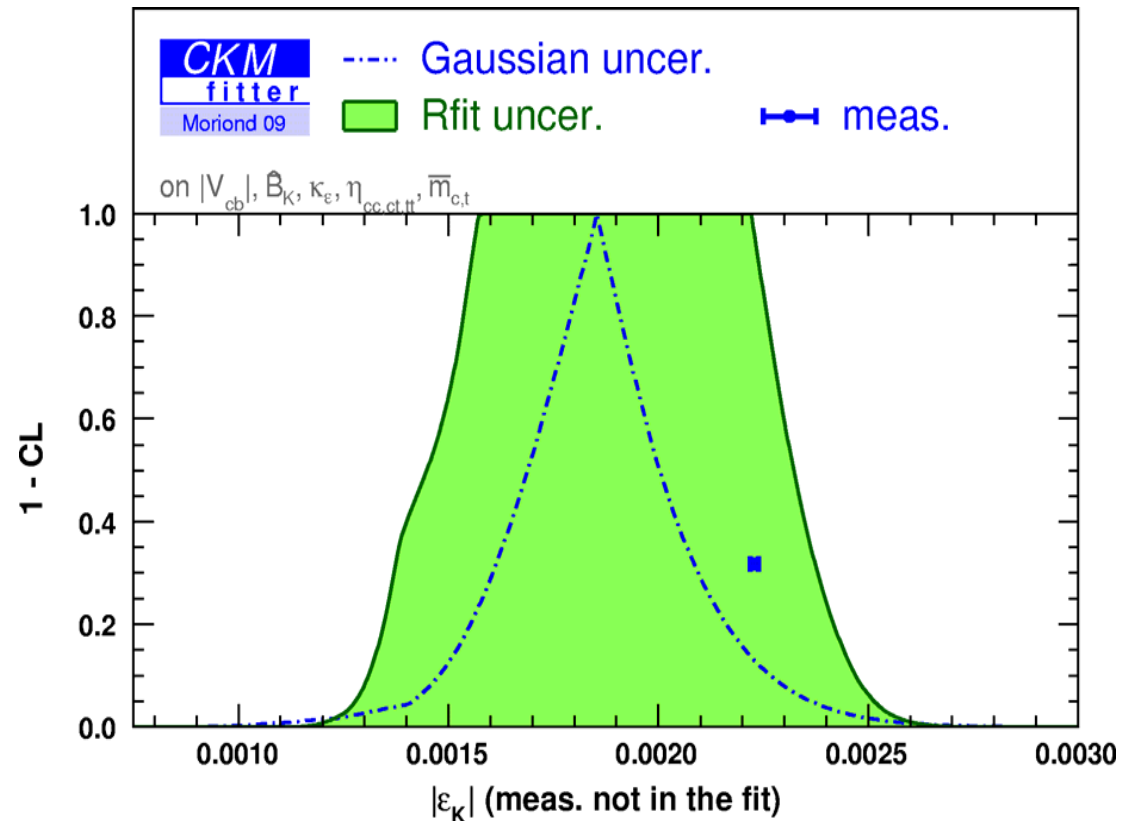
- Reminder from Buras & Guadagnoli (Phys. Rev. D78, 033005 (2008)): there is an additional suppression factor, κ_ε , to $|\varepsilon_K|$, estimated to be $\kappa_\varepsilon = 0.92(2)$.

⇒ Note that this factor has not yet been accounted for in the Global Fit. Its in the line for next update

- Any tension between direct measurement of $|\varepsilon_K|$ and indirect measurement from the global fit (through $\sin(2\beta_{cc})$)?

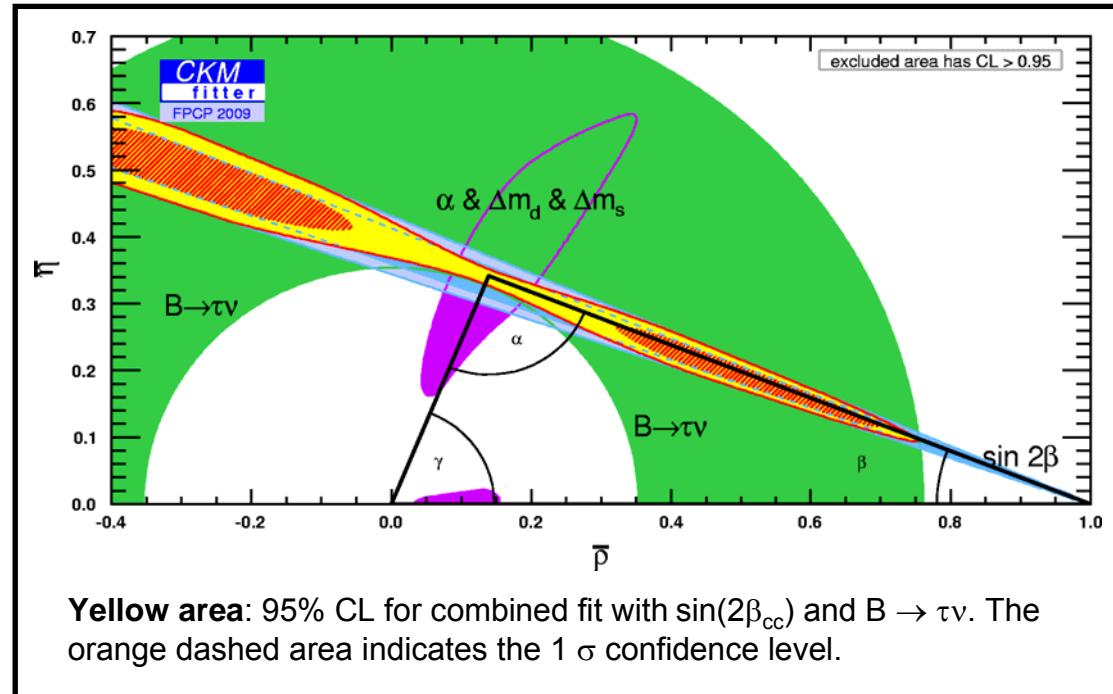
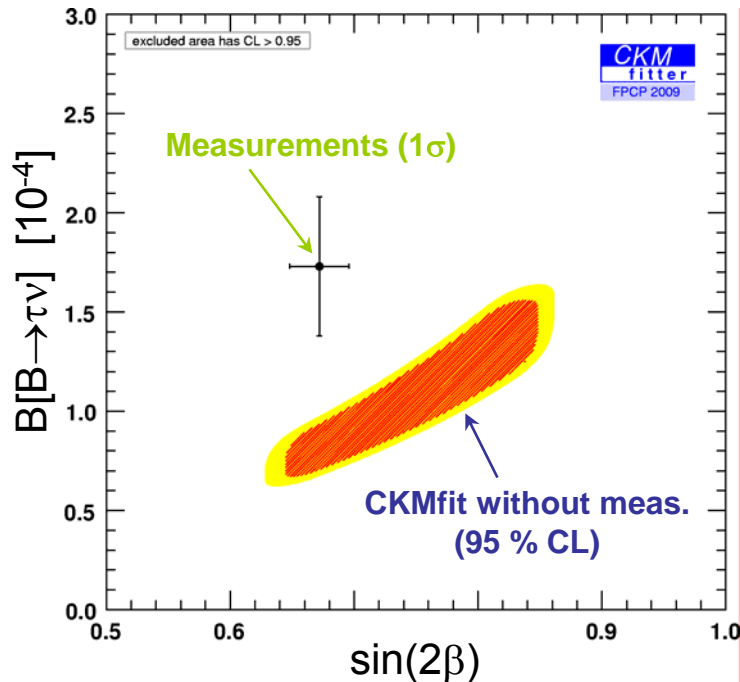
⇒ Using **Gaussian distributions** for systematic uncertainties and including the factor $\kappa_\varepsilon = 0.92(2)$ we get **1.5 σ** deviation

⇒ With our **Educated RFit** treatment of systematics **no deviation is seen**. The measurement is compatible with our fit best guess considering **uncertainties** on CKM parameters (through $|V_{cb}|^4 \sim 7\%$) mainly and hadronic uncertainties from B_K (**$\sim 5\%$**).



- The Global fit χ^2_{\min} drops by $\sim 2.4 \sigma$ if taking out $\sin(2\beta_{cc})$ or $B \rightarrow \tau \nu$.
- The combination $\sin(2\beta_{cc})$ and $B \rightarrow \tau \nu$ favours 2 solutions in contradictions with other inputs. Therefore **one can not accommodate both inputs simultaneously in the global fit.**

\Rightarrow The current fit **best guess favours $\sin(2\beta_{cc})$** better measured (4% vs 20%) while reflecting most errors on $B \rightarrow \tau \nu$



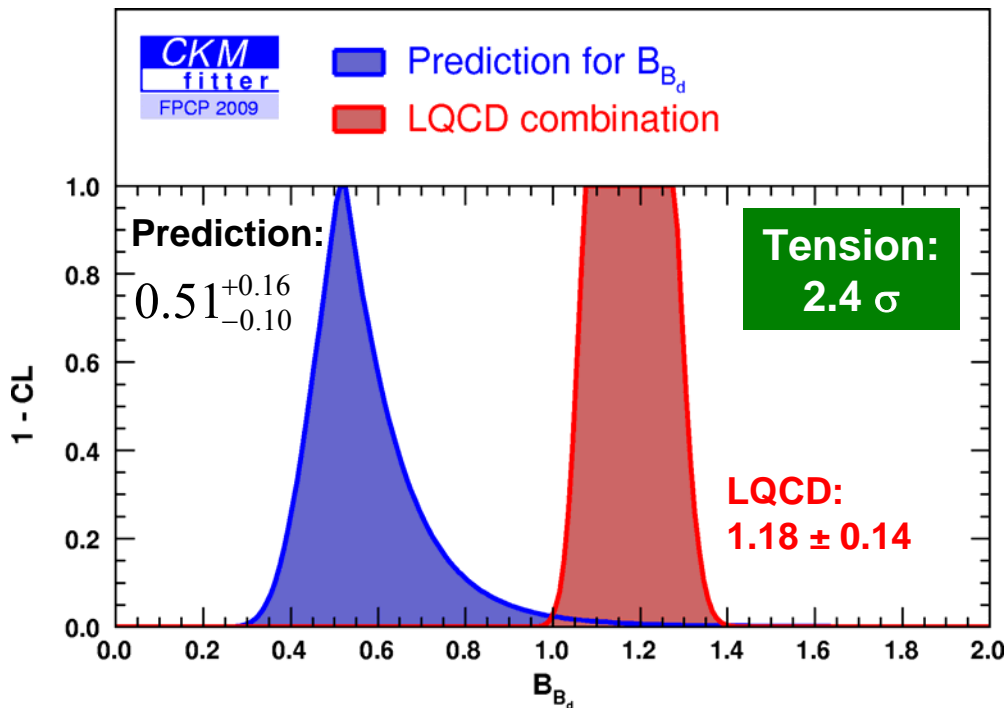
Closer look at 3 scenarios

- 1) Underestimated hadronic uncertainties
- 2) Measurements, fluctuations
- 3) Indication of New Physics

- The bag parameter B_{B_d} can be measured from the ratio of $B \rightarrow \tau \nu$ to Δm_d eliminating the dependency to f_{B_d} , as:

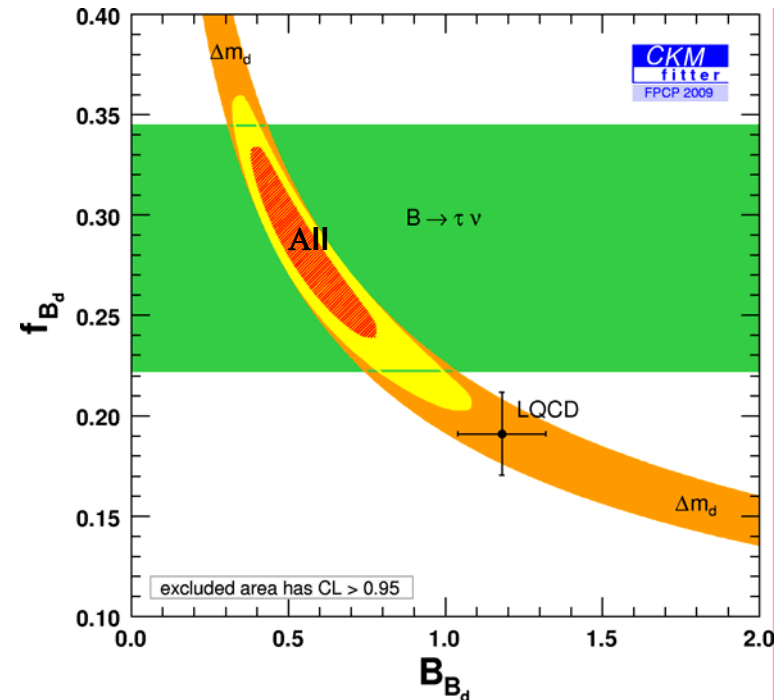
$$\frac{B[B \rightarrow \tau \nu]}{\Delta m_d} = \frac{3\pi}{4} \frac{m_\tau^2 \tau_B}{m_W^2 \eta_B S[x_\tau]} \left(1 - \frac{m_\tau^2}{m_B^2}\right)^2 \frac{\sin^2(\beta)}{\sin^2(\alpha + \beta)} \frac{1}{|V_{ud}|^2 \hat{B}_{B_d}}$$

The tension is still there at $\sim 2.4\sigma$! But a factor of 2 off on B_{B_d} while keeping f_{B_d} wouldn't work in the global fit ...



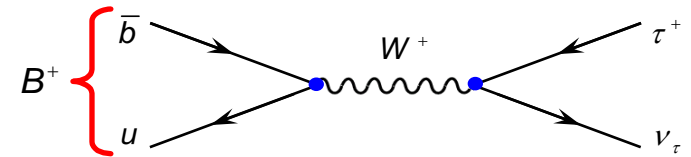
- Let's let f_{B_d} and B_{B_d} be completely free and fit them from all observables. What do we get?

\Rightarrow No more tension / no more constraints
 \Rightarrow The global fit is accommodated keeping $f_{B_d}^2 \times B_{B_d} \approx \text{const}$ to fit Δm_d while increasing f_{B_d} to fit $B \rightarrow \tau \nu$



Would require significantly out of range LQCD errors

$$\text{BR}(B^+ \rightarrow \tau^+ \nu) = \frac{G_F^2 m_B \tau_B}{8\pi} m_\tau^2 \left(1 - \frac{m_\tau^2}{m_B^2}\right)^2 f_{B_d}^2 |V_{ub}|^2$$



- Helicity-suppressed annihilation decay sensitive to $(f_{B_d} \times |V_{ub}|)^2$

Experimental measurements

	$B[B \rightarrow \tau \nu] \times 10^4$
Belle (hadronic)	1.79 ± 0.71 [2006]
Belle (semi-leptonic)	1.65 ± 0.52 [ICHEP08]
Belle	1.70 ± 0.42
BABAR (hadronic)	1.80 ± 1.00 [2007]
BABAR (semi-leptonic)	1.80 ± 0.81 [CKM08]
BABAR	1.80 ± 0.63
World Average	1.73 ± 0.35

CKMfit prediction: $(0.796^{+0.154}_{-0.093}) \times 10^{-4}$ (1σ , without meas.)

The various measurements for $B \rightarrow \tau \nu$ look consistent; we combine them using a weighted mean and assume Gaussian distributions. The p-value for this hypothesis is 11% (1.6σ).

$\sin(2\beta)$ from HFAG charmonium WA

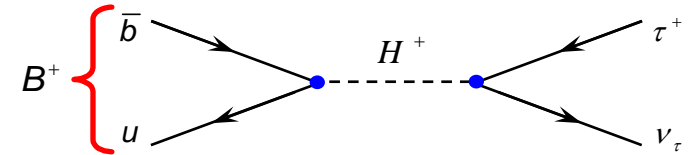
$$\sin(2\beta_{cc}) = 0.671(23)$$

No obvious tension

There is an overall experimental agreement that $B[B \rightarrow \tau \nu]$ is too high or $\sin(2\beta_{cc})$ too low

Charged Higgs contributions: increase $B[B \rightarrow \tau \nu]$ prediction

$$\text{BR}(B^+ \rightarrow \tau^+ \nu) = \frac{G_F^2 m_B \tau_B}{8\pi} m_\tau^2 \left(1 - \frac{m_\tau^2}{m_B^2}\right)^2 f_{B_d}^2 |V_{ub}|^2 \times (1 + r_H^B)^2$$



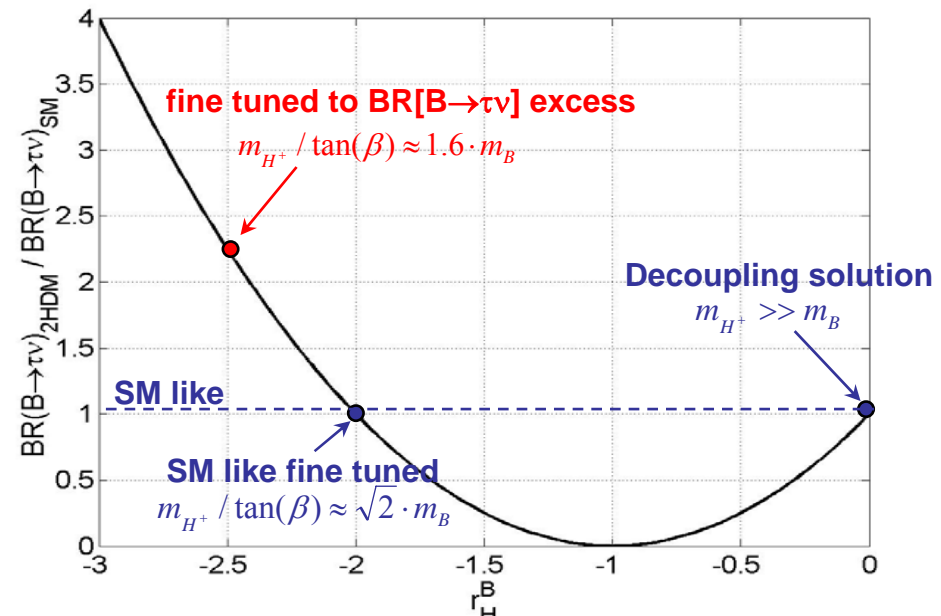
Charged Higgs contribution can modify $B[B \rightarrow \tau \nu]$ as a multiplicative term: $r_H^B \approx -\tan^2(\beta) m_B^2 / m_{H^\pm}^2$ in **2HDM Type II** model. Note that one would need $r_H^B \approx -2.5$ to fit $B[B \rightarrow \tau \nu]$ (fine tuned solution).

\Rightarrow **Requires a global analysis with other observables to check implications.**

Agreement with the SM can be recovered 2 ways:

- $r_H^B \rightarrow 0 \Rightarrow m_{H^\pm} / m_B \rightarrow \infty$ irrespective $\tan(\beta)$. This is the **decoupling solution**
- $r_H^B = -2 \Rightarrow m_{H^\pm} / \tan(\beta) \approx \sqrt{2} \cdot m_B$; requires a **fine tuning of $m_{H^\pm} / \tan(\beta)$ to the meson mass**.

Here fine tuning to adjust $\text{BR}[B \rightarrow \tau \nu]$ excess

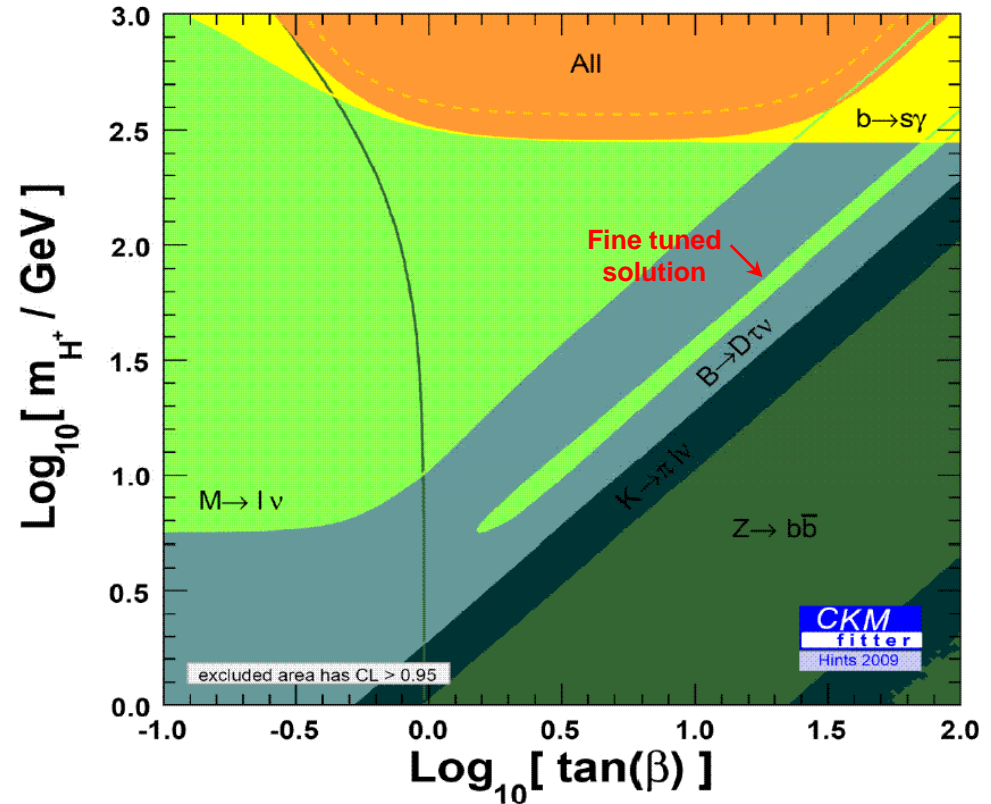
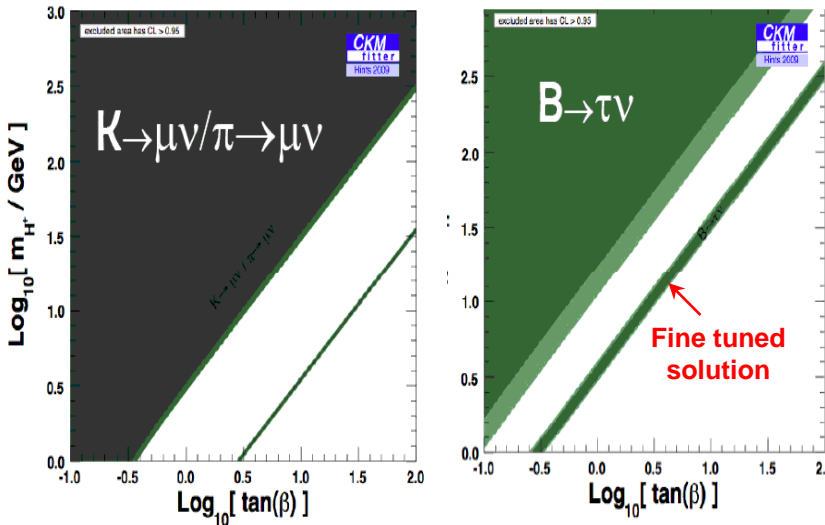


- **Combined 2HDM(II) analysis** within CKMfitter including modified constraints from mesons leptonic and semileptonic tree decays and loop radiative $b \rightarrow s\gamma$ decays and $Z \rightarrow b\bar{b}$ partial width.

Observables

$B[B \rightarrow \tau\nu], B[D \rightarrow \mu\nu], B[D_s \rightarrow \mu\nu], B[D_s \rightarrow \tau\nu]$
 $B[K \rightarrow \mu\nu]/B[\pi \rightarrow \mu\nu], B[B \rightarrow D\tau\nu], B[K \rightarrow \pi l\nu]$
 $B[b \rightarrow s\gamma], \Gamma[Z \rightarrow b\bar{b}]/\Gamma[Z \rightarrow \text{hadrons}]$

Individual constraints from some leptonic decays



Fine tuned solution ruled out at 95% CL, mostly from $B[K \rightarrow \mu\nu]/B[\pi \rightarrow \mu\nu]$ and $B[B \rightarrow D\tau\nu]$ constraints.

- **Only marginal improvement of the χ^2_{\min}** when going from SM to 2HDM(II), $\Delta\chi^2_{\min} = 1.4$ which corresponds to a p-value of 31%, **1.0 σ effect**, from a toy Monte-Carlo study.

⇒ **We see no particular indication for a charged Higgs effect in a 2HDM(II) scheme**

■ New Physics in $B_{q=d,s}$ mixing: decrease $\sin(2\beta_{cc})$ prediction

Assume that **NP only** affects shorts distance Physics in $\Delta B = 2$. Model independent parameterisation. Use Cartesian coordinates (Lenz&Nierste 2006):

$$\langle B_q | \mathbf{H}_{\Delta B=2}^{\text{SM+NP}} | \bar{B}_q \rangle = \langle B_q | \mathbf{H}_{\Delta B=2}^{\text{SM}} | \bar{B}_q \rangle \times (\text{Re}[\Delta_q] + i \text{Im}[\Delta_q])$$

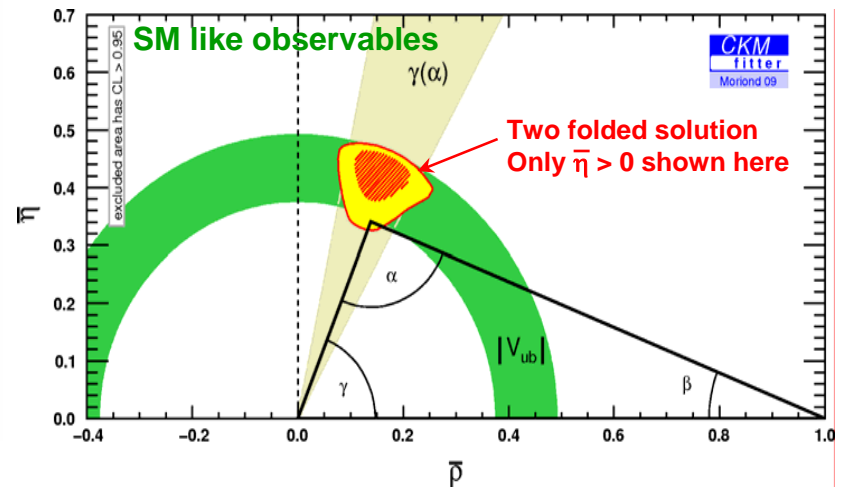
Note that:

$$\Delta q = |\Delta q| e^{2i\Phi_q^{\text{NP}}}$$

parameter	prediction in the presence of NP
Δm_q	$ \Delta_q^{\text{NP}} \times \Delta m_q^{\text{SM}}$
2β	$2\beta^{\text{SM}} + \Phi_d^{\text{NP}}$
$2\beta_s$	$2\beta_s^{\text{SM}} - \Phi_s^{\text{NP}}$
2α	$2(\pi - \beta^{\text{SM}} - \gamma) - \Phi_d^{\text{NP}}$
$\Phi_{12,q} = \text{Arg}\left[-\frac{M_{12,q}}{\Gamma_{12,q}}\right]$	$\Phi_{12,q}^{\text{SM}} + \Phi_q^{\text{NP}}$
A_{SL}^q	$\frac{\Gamma_{12,q}}{M_{12,q}^{\text{SM}}} \times \frac{\sin(\Phi_{12,q}^{\text{SM}} + \Phi_q^{\text{NP}})}{ \Delta_q^{\text{NP}} }$
$\Delta\Gamma_q$	$2 \Gamma_{12,q} \times \cos(\Phi_{12,q}^{\text{SM}} + \Phi_q^{\text{NP}})$

phases

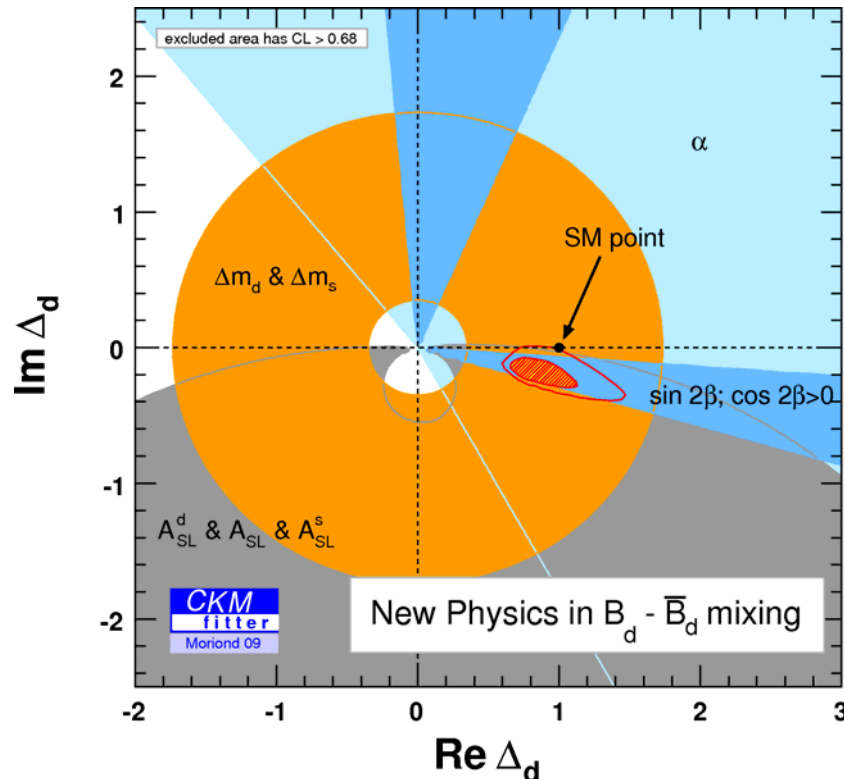
$$|V_{ud}|, |V_{us}|, |V_{cb}|, |V_{ub}|_{SL+B \rightarrow \tau\nu}, \gamma, \gamma(\alpha) = \pi - \alpha - \beta$$



Inputs:

Δm_d
 Δm_s
 $\sin(2\beta)$
 α
 $\Delta\Gamma_d$
 $A_{SL}^{B_d}$
 $A_{SL}^{B_s}$

Warning : only 68% CL regions are shown



The Cartesian parameterisation $\Delta_d = \text{Re}(\Delta_d) + i\text{Im}(\Delta_d)$ allows for a simple geometrical interpretation of each individual constraints.

- **Dominant constraint from β and Δm_d .** Semi leptonic asymmetries A_{SL} exclude the symmetric solution with $\eta < 0$. Note that the constraint resulting from **$\sin(2\beta_{cc})$ is shifted from real axis** due to the tension with $B \rightarrow \tau\nu$.

- **The additional phase can accommodate** the discrepancy between $\sin(2\beta_{cc})$ and $B \rightarrow \tau\nu$. For the hypothesis $\Delta_d = 1$ one recovers the **SM**. The latter is **disfavoured at 2.1σ** . At 95 % CL one would have:

$$\Phi_d^{NP} = (-12_{-6}^{+9})^\circ$$

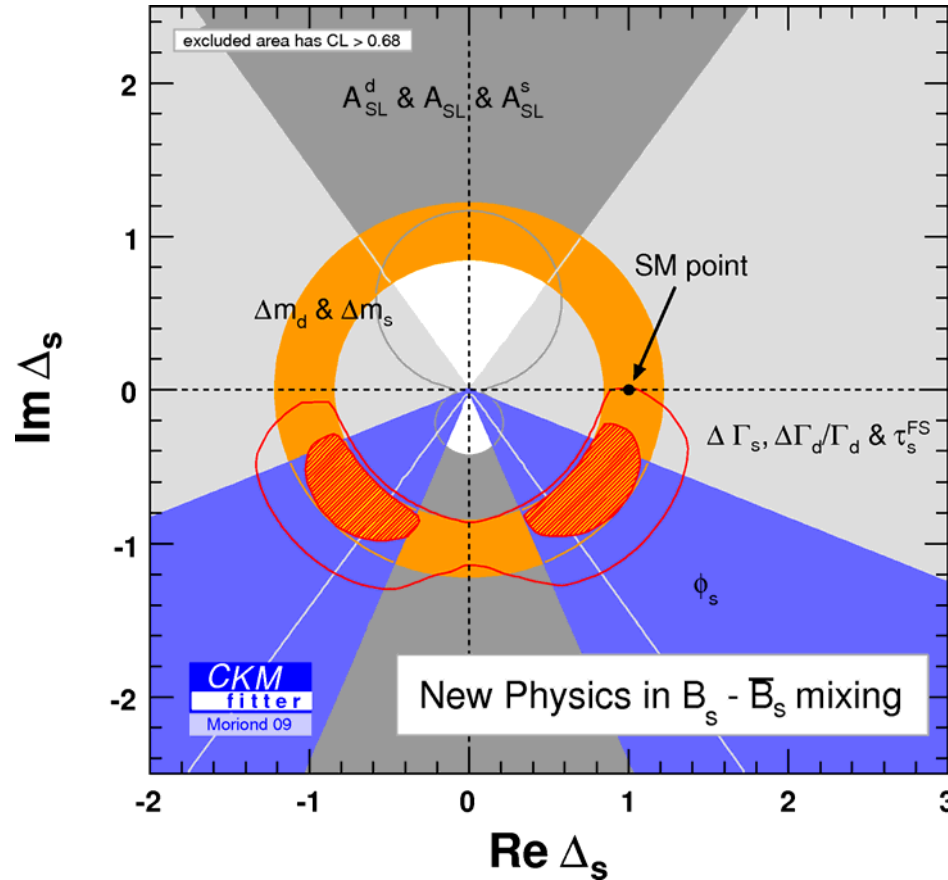
Note that **taking out $B \rightarrow \tau\nu$** one recovers agreement with the SM hypothesis at **0.6σ** .

No striking evidences for NP but it can be seen that sizeable NP contributions ($\phi_d^{NP} \sim 10^\circ$, $|\Delta_d| \sim \{0.5; 1.5\}$) in B_d system are still allowed

Inputs:

- ϕ_s
- Δm_s & Δm_d
- $A_{SL}^{B_d}$ & $A_{SL}^{B_s}$
- $\Delta \Gamma_s$ & τ_s^{FS}
- $B \rightarrow \tau \nu$

Warning : only 68% CL regions are shown



■ **Dominant constraints** from Tevatron direct measurement of $(\phi_s = -2\beta_s, \Delta \Gamma_s)$ in $B_s \rightarrow J/\psi \phi$ and from Δm_s .

ϕ_s D0/CDF (HFAG 08 update, CDF 1.35 fb⁻¹ only) is **2.2 σ away from SM** prediction.

Δm_s agrees with SM which constrains $|\Delta_s|$ to ~ 1 .

■ The **SM** 2D hypothesis $\Delta_s = 1$ is **disfavoured at 1.9 σ** with or without $B \rightarrow \tau \nu$. The tension is almost completely driven by ϕ_s alone.

Eagerly awaiting the Tevatron update. With the expected LHC luminosity in 2010, LHCb is competitive.

■ Standard Model CKM fit

- Very good agreement with the SM for most observables; the **KM phase is the dominant source of CP violation at the EW scale**.
- There is an ongoing **2.4 σ tension** between **$\sin(2\beta_{cc})$** and **$B[B \rightarrow \tau \nu]$** . Caution: in the current SM fit the tension reflects mostly on $B \rightarrow \tau \nu$ which suffers from larger experimental uncertainties.

95% CL interval:

$$\begin{aligned} \bar{\rho} &= 0.139^{+0.053}_{-0.040} \\ \bar{\eta} &= 0.341^{+0.032}_{-0.025} \end{aligned}$$

From all constraints:

$$\gamma = (67.8^{+4.2}_{-3.9})^\circ$$

Agreement with SM:

Hypothesis	Deviation
$\Delta_d = 1$	2.1 σ
$\Delta_s = 1$	1.9 σ
$\Delta_d = \Delta_s = 1$	2.5 σ

■ Origin of the tension?

- Seems hard to explain by LQCD uncertainties; Measurements look consistent on both sides.
- **New Physics**: an additional phase in B_d mixing could accommodate the measurements. Sizeable effects of $\sim 10^\circ$ / $\sim 40\%$ on mixing amplitude are still allowed.

■ Outlooks

- Further taking into account the deviation of ϕ_s , in a generic scenario of New Physics in mixing the **SM is disfavoured at 2.5 σ**
 \Rightarrow Waiting for **new data** to confirm or infirm this.
- 2HDM(II): we see no particular indication for a charged Higgs effect; preferred solution: decoupling (large m_{H^\pm} irrespective $\tan(\beta)$).

■ Winter 2009 SM & NP fit results are available at:

http://ckmfitter.in2p3.fr/plots_Moriond2009/



More details in: V. Tisserand (CKMfitter Group), Moriond EW 2009 proceedings [arXiv:0905.1572]
Or full details at: http://ckmfitter.in2p3.fr/plots_Moriond09/ckmEval_results.html

■ Our Own Average: *Educated RFit* Scheme

- 1) **Preselection**: Use only **unquenched** results with **2 or 2+1** dynamical fermions (sea quarks), also include **staggered fermions**.
⇒ papers & proceedings: RBC, UKQCD, HPQCD, JLQCD, CP-PACS, FNAL Lattice, MILC, ETMC, NPLQCD ...
- 2) For error budget: **split into statistical and systematic uncertainties**.
- 3) Fit the central value of the LQCD observable according to the **RFit** scheme.
- 4) Quadratic average of statistical errors. **Systematic**: add linearly contributions to a single LQCD prediction. Take the systematic of the **most precise method**; not more nor less than the most accurate prediction.

■ Sources of LQD uncertainties

Euclidian, finite, discrete box. Observable = **statistical average over gauge configurations** weighted according to **gauge and fermion actions**.

Statistical: size of the ensemble of gauge configurations; part of errors below when scaling with size of gauge configurations.

Systematics: Fermions action ($N_f=2$, staggered); continuum limit $a \rightarrow 0$; finite volume effects $L \rightarrow \infty$; Quark mass extrapolation (chiral or heavy quark limit).

Input	Central	Stat.	Syst.	Reference
$ \text{Vud} $	0.97418	0.00026		Towner/Hardy [arxiv:0710.3181]
$ \text{Vus} $	0.2246	0.0012		Flavianet Vus (K_{13} only) [arxiv:0801.1817]
$ \text{Vub} $	$3.87 \cdot 10^{-3}$	$0.09 \cdot 10^{-3}$	$0.46 \cdot 10^{-3}$	OOA [arXiv:0905.1572]
$ \text{Vcb} $	$40.59 \cdot 10^{-3}$	$0.38 \cdot 10^{-3}$	$0.58 \cdot 10^{-3}$	OOA [arXiv:0905.1572]
$ \varepsilon_K $	$2.229 \cdot 10^{-3}$	$0.010 \cdot 10^{-3}$		PDG 08* (no κ_K in the fit)
Δm_d	0.507	0.005		PDG 08
Δm_s	17.77	0.12		CDF, PRL 97, 242003 (2006)
α	*	*	*	HFAG (BaBar/Belle)
$\sin 2\beta$	0.671	0.023		HFAG charmonium
$\cos 2\beta$	0.5		0.5	cos2beta>0
γ	*	*	*	
$\mathbf{B}(\mathbf{B} \rightarrow \tau \nu)$	$1.73 \cdot 10^{-4}$	$0.35 \cdot 10^{-4}$		ICHEP08+CKM08(BABAR) WA
\bar{m}_c	1.286	0.013	0.040	
\bar{m}_t	165.017	1.156	0.11	Tevatron Electroweak Working Group [arXiv:0808.1089]
$\alpha_s(m_Z)$	0.1176		0.0020	PDG 08
\mathbf{B}_K	0.721	0.005	0.040	OOA [arXiv:0905.1572]
f_{B_s}	0.228	0.003	0.017	OOA [arXiv:0905.1572]
\mathbf{B}_{B_s}	1.23	0.03	0.05	OOA [arXiv:0905.1572]
f_{B_s}/f_{B_d}	1.196	0.008	0.023	OOA [arXiv:0905.1572]
$\mathbf{B}_{B_s}/\mathbf{B}_{B_d}$	1.05	0.02	0.05	OOA [arXiv:0905.1572]
η_{ct}	0.47		0.04	Herrlich&Nierste [Nucl. Phys. B 419, 292 (1994)]
η_{tt}	0.5765		0.0065	Herrlich&Nierste [private comm.]
δ_1	0		1	Nierste [private comm.]
η_B	0.551		0.007	Buchalla <i>et al.</i> [Rev. Mod. Phys. 68, 1125 (1996)]

