

Interpretation of $e^+e^- \rightarrow b\bar{b}$

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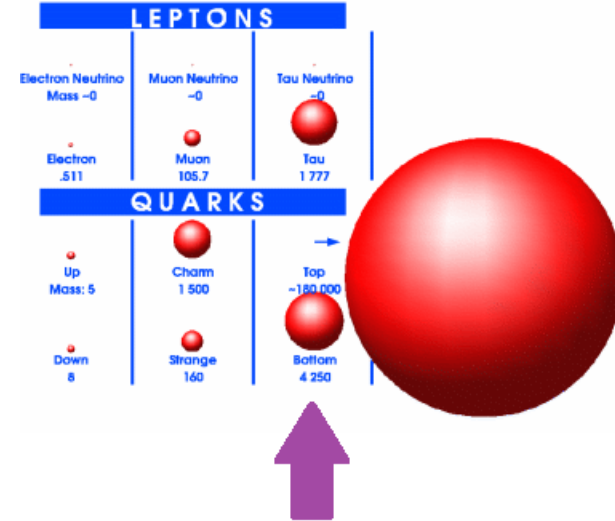
in collaboration with S. Bilokin and R. Poeschl

[Workshop on top physics at the LC 2017](#)

CERN 7-9 juin 2017

Introduction

- The main attention has been up to now focused on the top sector, but e^+e^- also offers a unique opportunity to study weak couplings of **b quarks** to Z and γ
- While Z couplings can, as for LEP1, be measured at the Z resonance, outside the resonance **beam polarisation** is needed to separate the Z and γ components since, in contrast to the top, b decays are dominated by **scalar meson channels** and offer no observables to disentangle the two components
- Several **composite models** predict a high degree of compositeness both for the t and b quarks. In this respect the well-known deviation observed at LEP1 should be verified/discarded in the future
- Can this be done at **ILC**? To which accuracy ?



Purpose of this talk

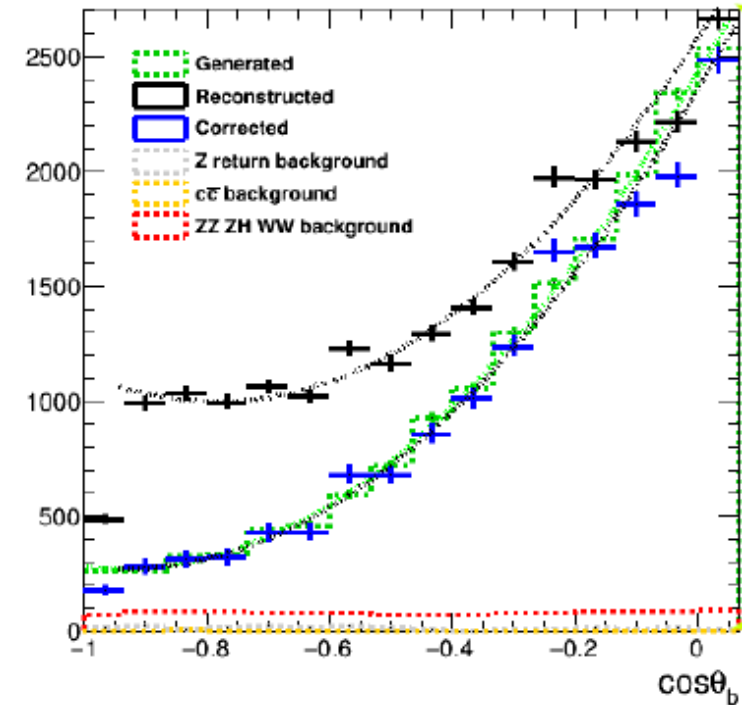
- This presentation intends to discuss a very recent analysis of the channel $ee \rightarrow bb$ with the ILD set up, collecting 500 fb^{-1} at 250 GeV with **polarized e^- (80%) and e^+ (30%)**
- I will briefly recall (see talk of **S. Bilokin**) the main features of this analysis relying on b charge determination using the **microvertex information** and the charged **kaon identification** from the TPC of ILD
- I will then derive the accuracies achievable from these measurements on the vector and axial couplings
- I will finally go through the interpretation of these results, recalling the **LEP1 puzzle** and putting this measurement in context with top physics

How to measure $e^+e^- \rightarrow b\bar{b}$

- Secondary vertices allow to unambiguously identify **charged B mesons** and therefore the b quark charge
- Unfortunately this can only be done with limited purity, $\sim 80\%$, given that the secondary track efficiency is limited at $\sim 95\%$ (half of this inefficiency being due to an absence of significant offset) recalling that, on average, one needs to reconstruct 5 charged tracks.
- To cope with this, one requires the presence of two charged B mesons, which reduces the efficiency
- This efficiency is increased to 13% by selecting **charged kaons** which give a comparable purity

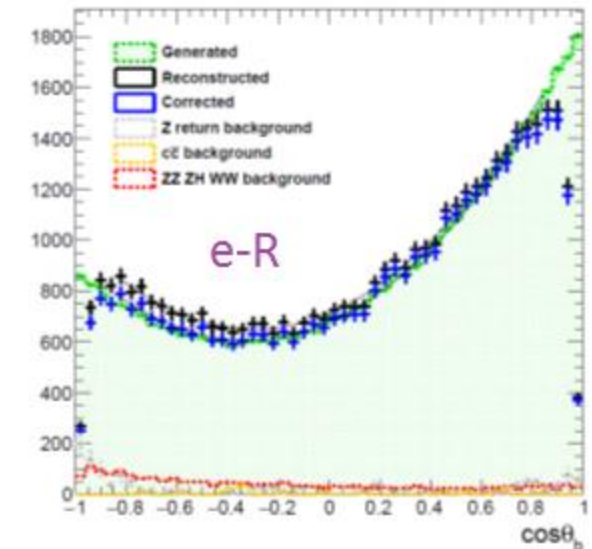
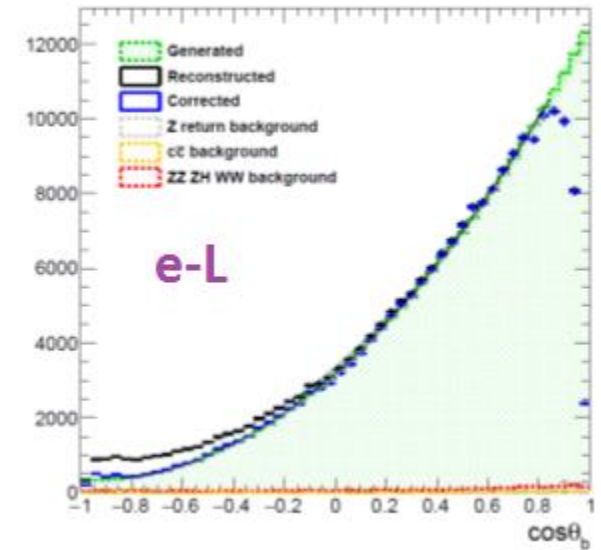
Limitation

- Selecting two charged B mesons is however insufficient, in particular to measure $e^-_L e^- \rightarrow b\bar{b}$ which is very forward backward asymmetric
- A method was derived, using the amount of wrong sign events B^+B^+ and B^-B^- to correct for this effect
- Migration due to track losses can be **perfectly corrected** from the data themselves
- If BSM physics populate the backward hemisphere, it will be detected with very **high sensitivity**



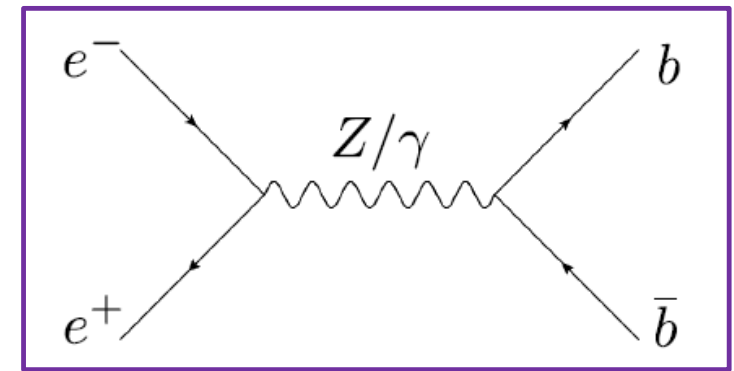
Results

- The two curves are obtained with the two beams **fully polarized** and an integrated luminosity of 250 fb^{-1} for each polarisation
- 3% background coming from ZZ, ZH and radiative return to Z
- In practice 2/3 Lumi. taken with e_L and 1/3 with e_R [1506.07830](#)
- One can obtain these two distributions L and R from combining data obtained with partial polarisation
- $d\sigma_{-,+}/d\cos\theta = 0.58L + 0.035R$ and $d\sigma_{+,-}/d\cos\theta = 0.58R + 0.035L$
 $L = d\sigma_L/d\cos\theta$ and $R = d\sigma_R/d\cos\theta$
- L and R are fitted by $d\sigma/d\cos\theta \sim S(1 + \cos^2\theta) + A\cos\theta$
- $S \sim F1V^2 + F1A^2$ and $A \sim F1VF1A$ terms in $1/\gamma^2$ can be neglected
- For e_L S and A **strongly correlated** ($\rho \sim 0.8$) which reduces the errors
- The fit uses the region where $|\cos\theta| < 0.8$

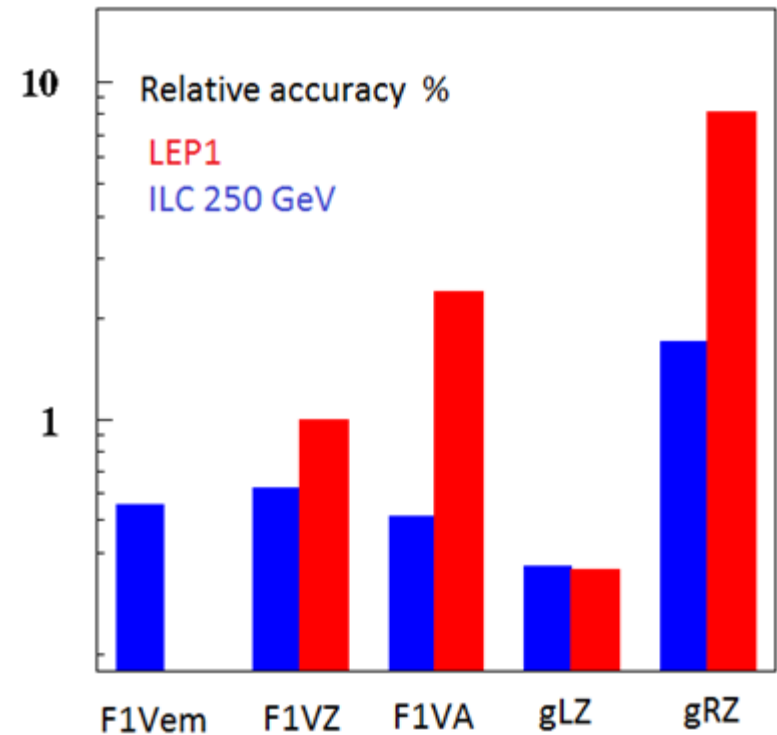


Separating γ from Z exchange

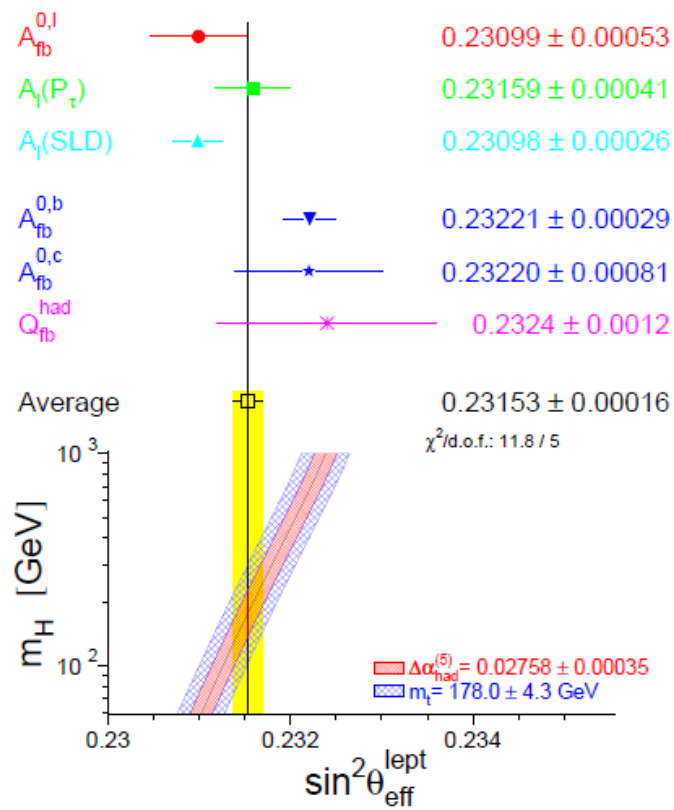
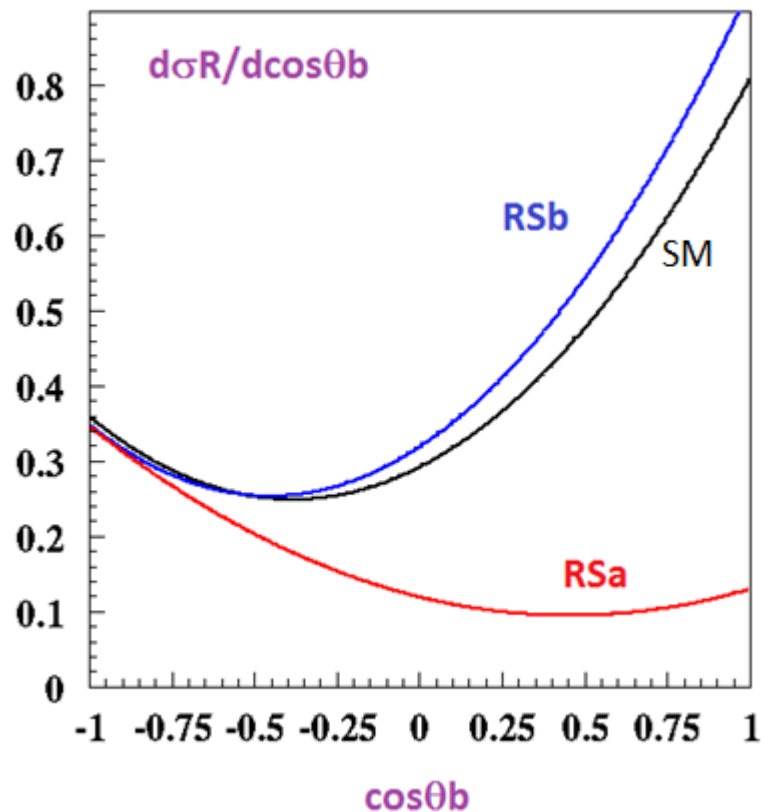
- With the four measurement of S and A for two polarisations one can **fully disentangle** F1VZ, F1AZ (or g_{RZ} and g_{LZ}) and F1Vem
- Precision is not limited **by systematics** (P, Lumi., ϵ_b , back. $\sim 0.4\%$) and errors can be reduced with higher luminosity (2000 fb^{-1} is planned at ILC)
- $\delta g_{RZ}/g_{RZ} \sim 2\%$ sufficient to **confirm at $>5\sigma$** or to **discard** the LEP1 effect which is at the **25%** level
- Recall the sign uncertainty on LEP1 solutions $dg_{RZ}/g_{RZ} = 25\%$ or $dg_{RZ}/g_{RZ} = -225\%$
- Not a problem at 250 GeV to make the right choice for the sign



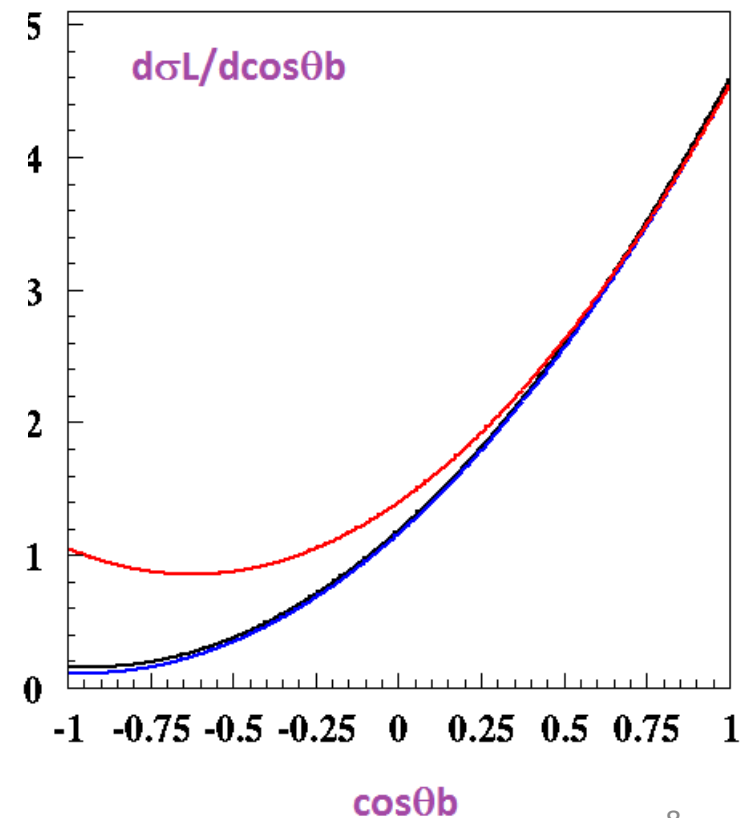
Preliminary



LEP1 effect



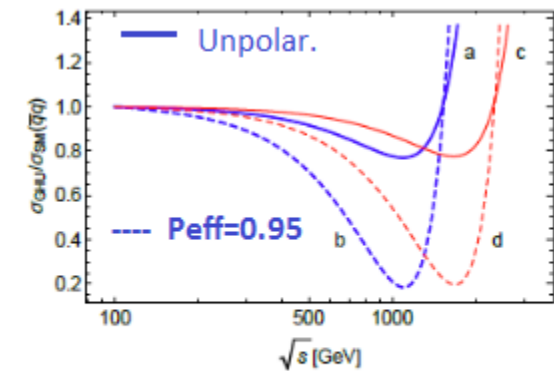
hep-ph/0610173



Interpretations

- Either use **specific models** like RS which has many versions or use **EFT**
- RS alone: $M_{\text{KK}} > 10$ TeV to pass S,T constraints
- RS with **custodial symmetries** to pass S,T constraints: $M_{\text{KK}} > 5$ TeV
- RS with special warping: $M_{\text{KK}} > 1$ TeV **1011.2205 LHCb anomalies**
- Will be explored by t and b measurements as shown in **1403.2893**
- RS with **gauge-Higgs unification**: $M_{\text{KK}} 5\text{-}10$ TeV **1705.05282**
- All flavour affected and will be tested not only with t,b but also will light quarks (**c to be investigated**) and leptons

f	$g_{Z^{(1)}f}^L$	$g_{Z^{(1)}f}^R$	$g_{Z_\mu^{(1)}f}^L$	$g_{Z_\mu^{(1)}f}^R$	$g_{\gamma^{(1)}f}^L$	$g_{\gamma^{(1)}f}^R$
ν_e	-0.2225	0	0	0	0	0
ν_μ	-0.2225	0	0	0	0	0
ν_τ	-0.2224	0	0	0	0	0
e	0.1196	0.9981	0	-1.3762	0.1880	-1.8165
μ	0.1196	0.9369	0	-1.3029	0.1880	-1.7051
τ	0.1195	0.8847	0	-1.2401	0.1879	-1.6102
u	-0.1539	-0.6536	0	0.9034	-0.1253	1.1896
c	-0.1539	-0.6041	0	0.8439	-0.1253	1.0994
t	0.6888	-0.3431	1.3208	0.5253	0.5616	0.6258
d	0.1882	0.3268	0	-0.4517	0.1303	-1.2369
s	0.1882	0.3021	0	-0.4220	0.1303	-1.1431
b	-0.8470	0.1720	1.3189	-0.2625	-0.5840	-0.6506

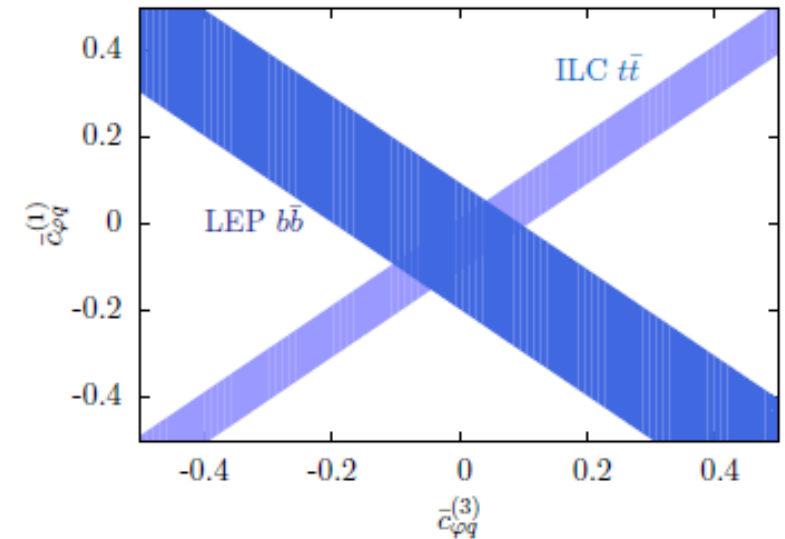


at ILC $P_{\text{eff}}=0.89$

EFT

- Recently have compared and e+e- and LHC potential using the Effective Field Theory approach [1704.01782](#)
- It clearly states that b measurement are a key feature to **narrow down uncertainties**
- This reference uses a coarse analysis which can be improved using ILC results on top [1505.06020](#)
- Reach $\Lambda \sim 10 \text{ TeV}$ is achievable

$$C_{1V} = \frac{v^2}{\Lambda^2} \Re \left[c_{\varphi q}^{(3)} - c_{\varphi q}^{(1)} - c_{\varphi u} \right]^{33}$$
$$C_{1A} = \frac{v^2}{\Lambda^2} \Re \left[c_{\varphi q}^{(3)} - c_{\varphi q}^{(1)} + c_{\varphi u} \right]^{33}$$
$$C_{2V} = \sqrt{2} \frac{v^2}{\Lambda^2} \Re \left[\cos \theta_W c_{uW} - \sin \theta_W c_{uB} \right]^{33}$$
$$C_{2A} = \sqrt{2} \frac{v^2}{\Lambda^2} \Im \left[\cos \theta_W c_{uW} + \sin \theta_W c_{uB} \right]^{33}$$



Conclusion

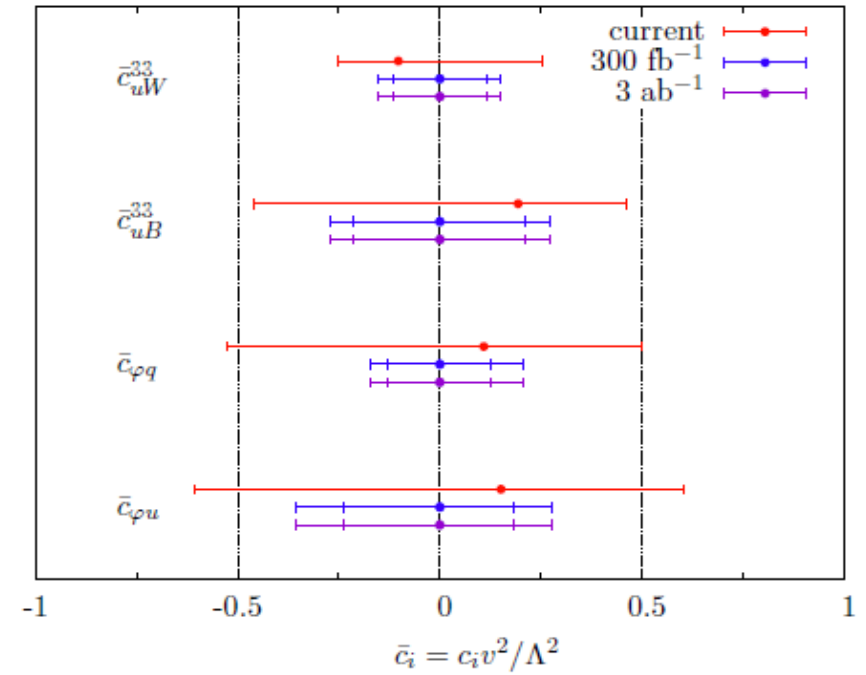
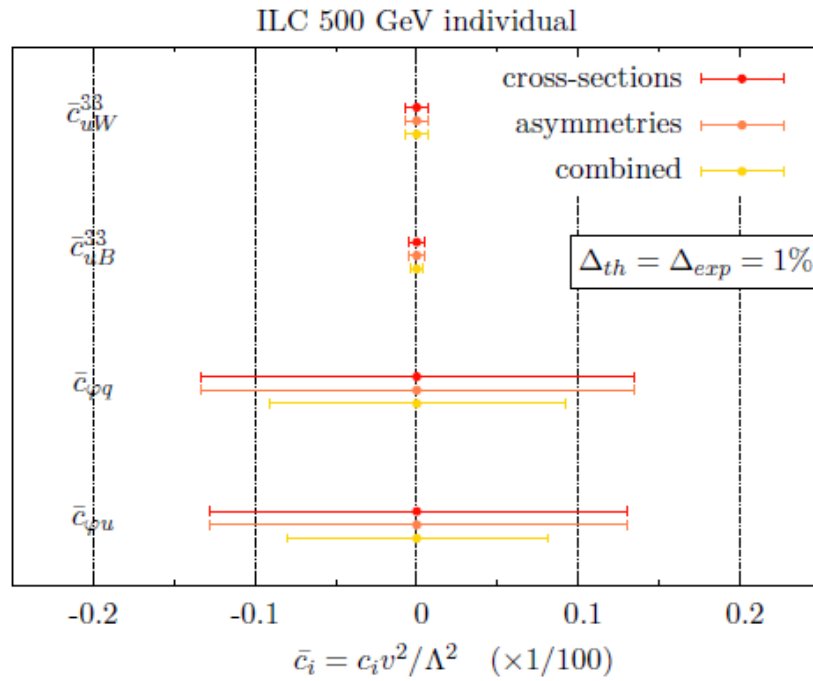
- ee->bb very demanding for the ILD detector (can be used as a **benchmark**)
- **Migration effects can be perfectly corrected** from the data themselves without using external informations
- **Polarisation needed** to separate the γ and Z couplings
- Accuracies surpass LEP1 but **EW corrections are needed**
- The **LEP1 effect** can be conclusively tested at 250 GeV with 500 fb⁻¹
- Reach on composite models, in particular in their RS versions, can be indirectly covered for masses well **beyond the LHC** using top and b EW measurements
- EFT illustrates the strong **complementarity** between **b and t quark** EW measurements and show that a collider reaching 500 GeV can cover effective mass scales up to 10 TeV

BACK UP SLIDES

EFT

1704.01782

LHC



Formulae

- $d\sigma/d\cos\theta = S_0 \{ (1 + \cos^2\theta) [|F_{1V} + F_{2V}|^2 + F_{1A}^2] + (1 - \cos^2\theta) | \gamma^{-1} F_{1V} + \gamma F_{2V} |^2 - 4 \cos\theta F_{1A} (F_{1V} + \Re F_{2V}) \}$

$$F_{1V}^L = -Q_b + e_L BW \frac{-0.25 - Q_b s_W^2}{s_W c_W} \quad F_{1A}^L = e_L BW \frac{0.25}{s_W c_W}$$

- $e_L = (-0.5 + s^2 w) / c w s w$ $BW = s / (s - M_Z^2)$
- $F_{2V} \sim \alpha_s / \gamma^2$ hence $1 - \cos^2\theta$ can be neglected for b quarks
- Keeping interference terms and neglecting $1/\gamma^2$ terms
- $d\sigma/d\cos\theta = S_0 \{ (1 + \cos^2\theta) (F_{1V}^2 + 2F_{1V} \Re F_{2V} + F_{1A}^2) - 4 \cos\theta F_{1A} (F_{1V} + \Re F_{2V}) \}$