

R. Kersevan, M. Ady, CERN-TE-VSC-VSM

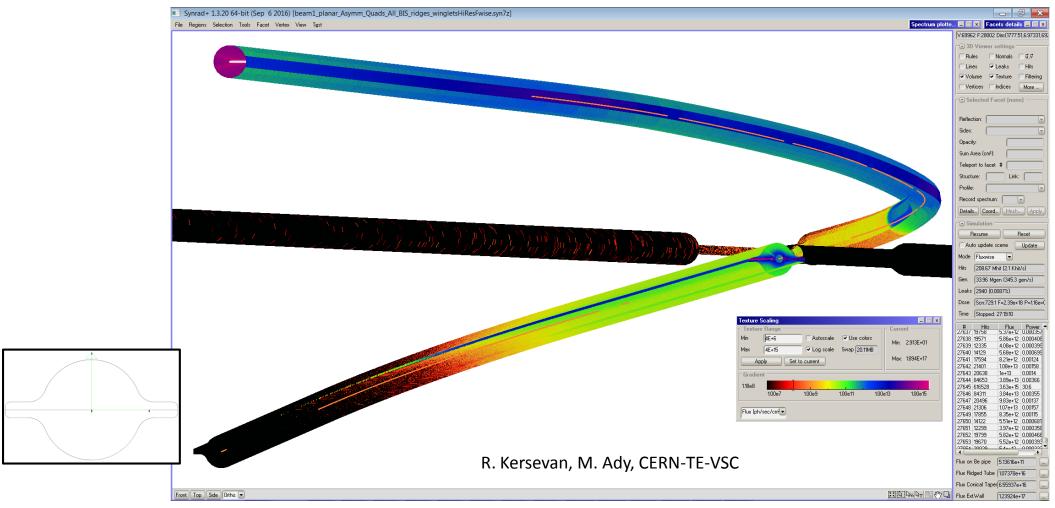


Ray Tracing with SYNRAD+ in FCC-ee IR

R. Kersevan, M. Ady, CERN-TE-VSC-VSM

- The geometry for the ~680 m around the IP for the latest T-pole lattice (as per K. Oide's file on AFS) has been created;
- One beam only has been modelled (assuming the other one is symmetric): it includes 9 dipoles and 9 quadrupoles;
- Two version of the geometry have been created: one with a symmetric opening of the IP quadrupole focusing doublets (20 mm ID) and one with an "exit" doublet chamber twice as big in radius (40 mm ID), following the prescription to let the trapped HOM in the IP "escape" and dissipate their power elsewhere (see previous meetings, this series);
- A third variant of the geometry has included a "winged" geometry, i.e. a chamber cross-section "à-la-SUPERKEKB" which allow the positioning of short localized absorber (to reduce the photon scattering and improve the shielding of high-energy photons, see end of my presentation at FCC Week in Rome);
- ➤ Neither the photon absorbers nor the bellows/contact fingers have been included yet; <</p>
- → This work does NOT aim to take the place of the analysis already done via GEANT4 and/or other calculations (H. Burkhardt et al., M. Sullivan et al.) but simply wants to show the potential of a different code, SYNRAD+, as far as the calculation of photon flux and power on sensitive equipment is concerned (like the Be pipe at the IP);
- The main purpose of this analysis is to prepare a geometry for simulating the pressure profiles (aiming at presenting this as a poster at FCC Week in Berlin, if accepted).

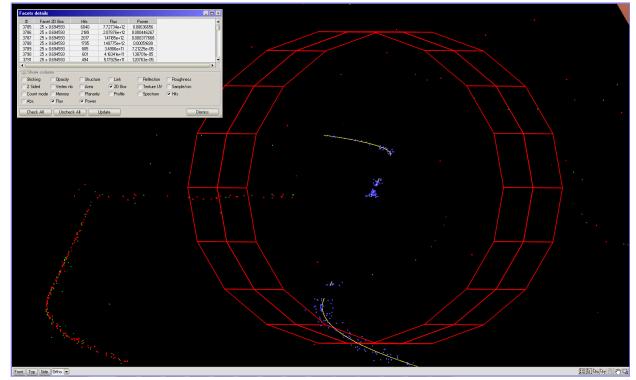
SYNRAD+: SR flux along one ~680 m-long arm of the interaction region of FCC-ee (175 GeV T-pole machine; 6.632 mA; Z-pole study to be done)

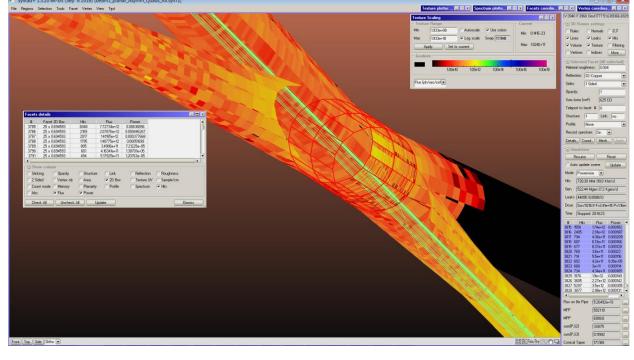


One arm of the IP chambers: \sim -347 m to \sim +337 m;

Round pipe (70 mm ID) everywhere except along the incoming beam, which has winglets;

SYNRAD+: SR flux along one ~640 m-long arm of the interaction region of FCC-ee (175 GeV T-pole machine; 6.632 mA)



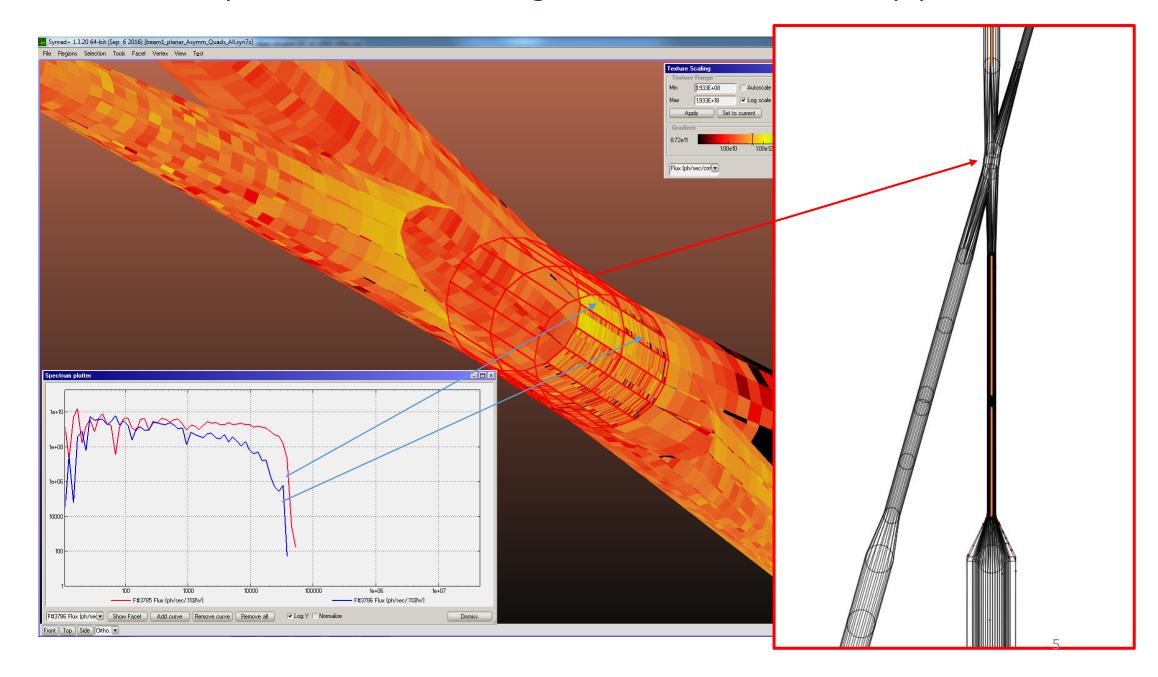


View of the source points (blue) and centroid trajectory (yellow); Red and Green points are locations of absorbed/reflected photons;

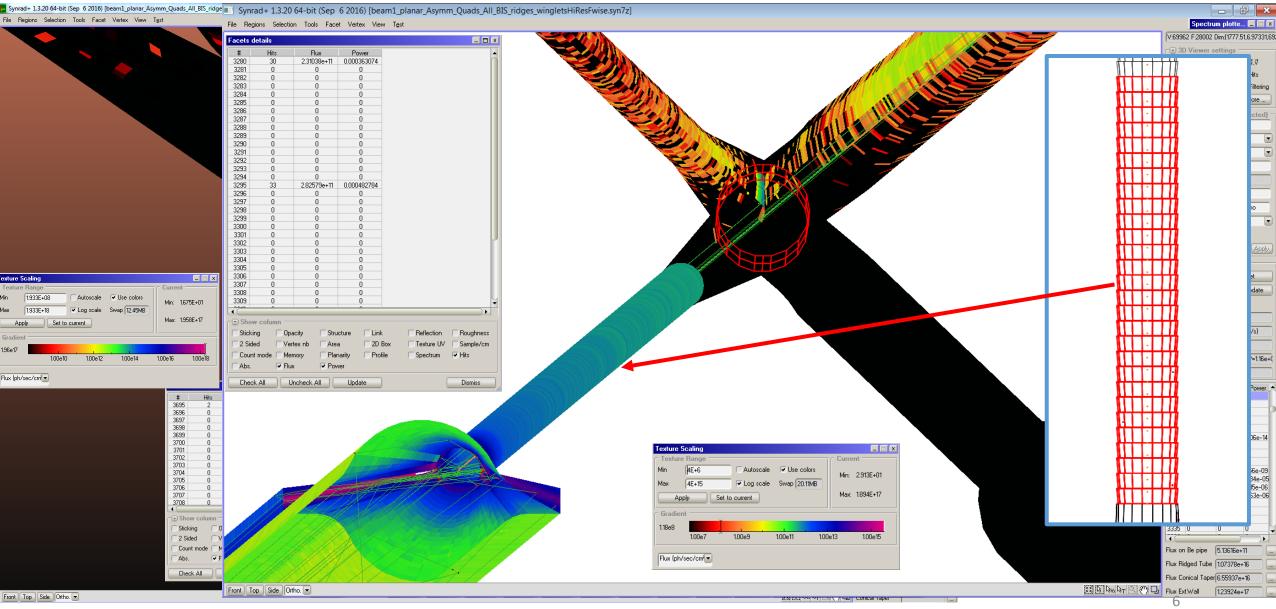
Red lines represent the 50 cm-long Be pipe

Zoom into the IP region: 50 cm-long Be pipe and local photon flux density distribution; A total of 5.26E+13 ph/s hit the Be pipe;

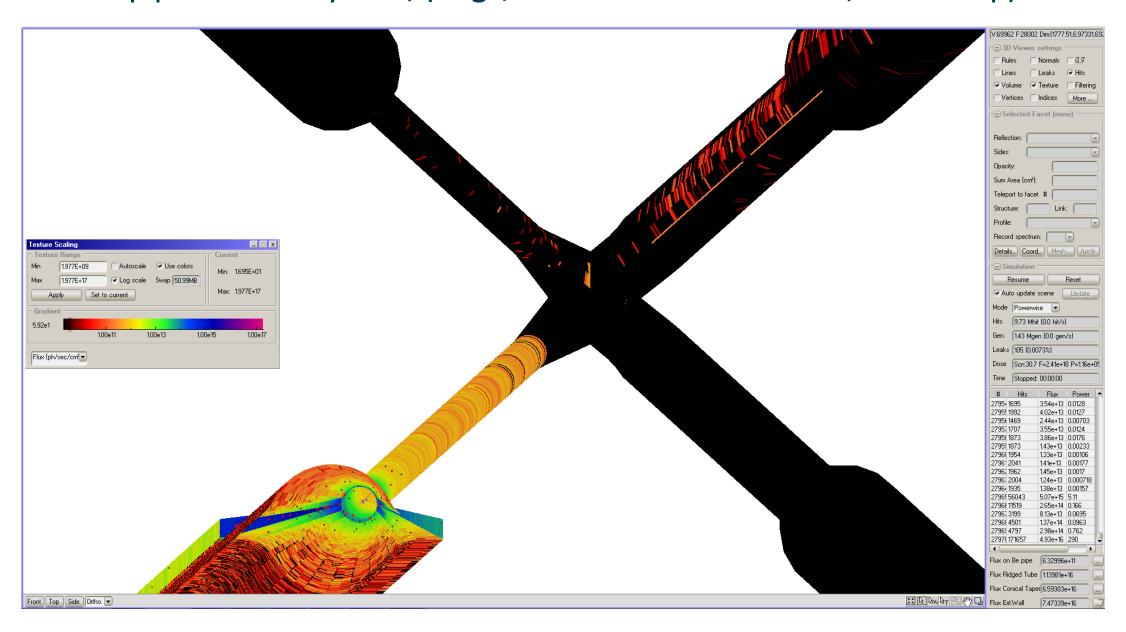
Photon flux spectrum on the two highest-flux facets of the Be pipe

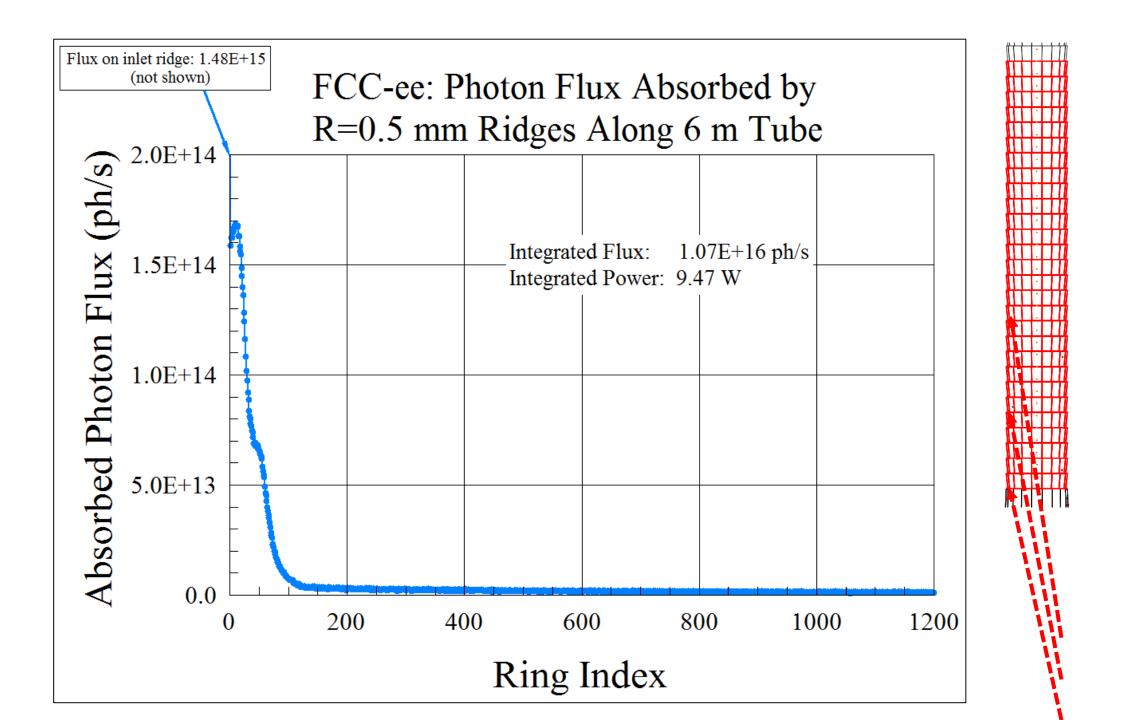


Adding a "ridged" (sawtooth) profile to the 20 mm ID doublet quad pipe: reduces the flux onto Be pipe to virtually zero; (ridge/sawtooth: 0.5 mm radial, 5 mm step)

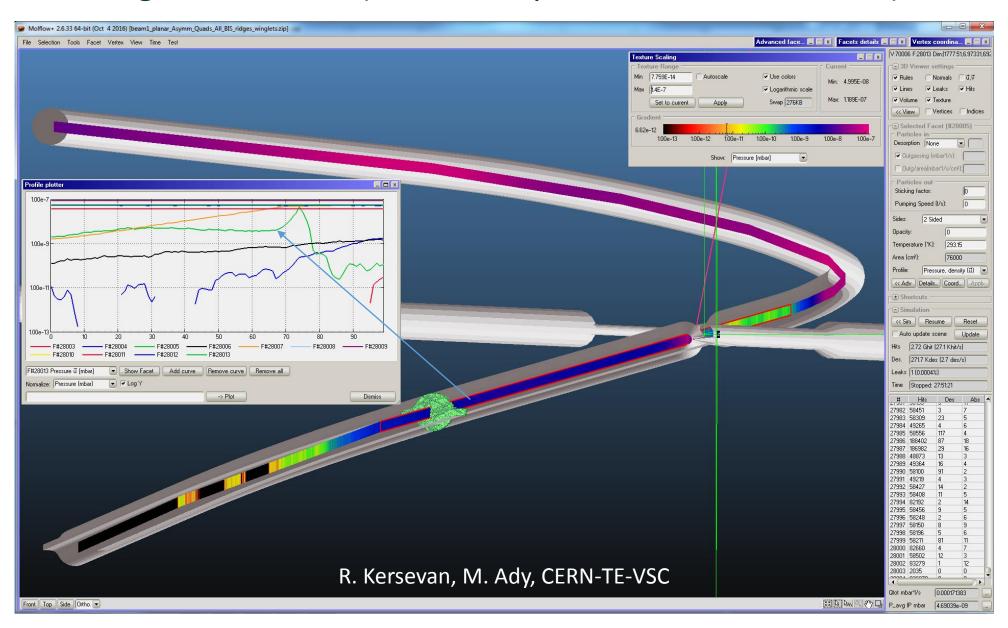


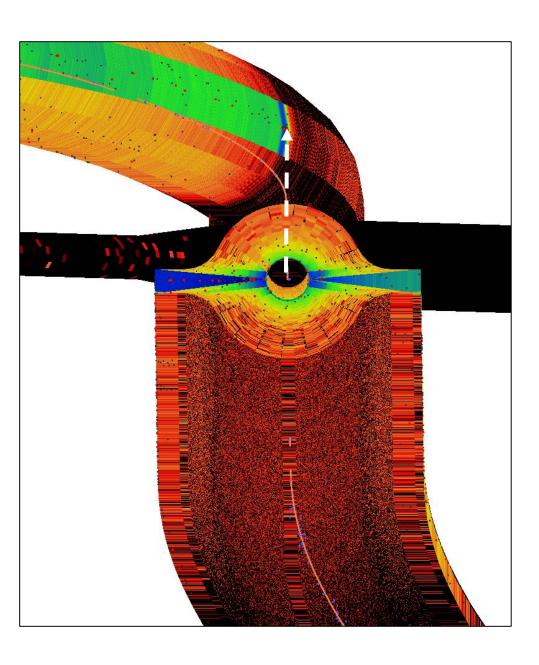
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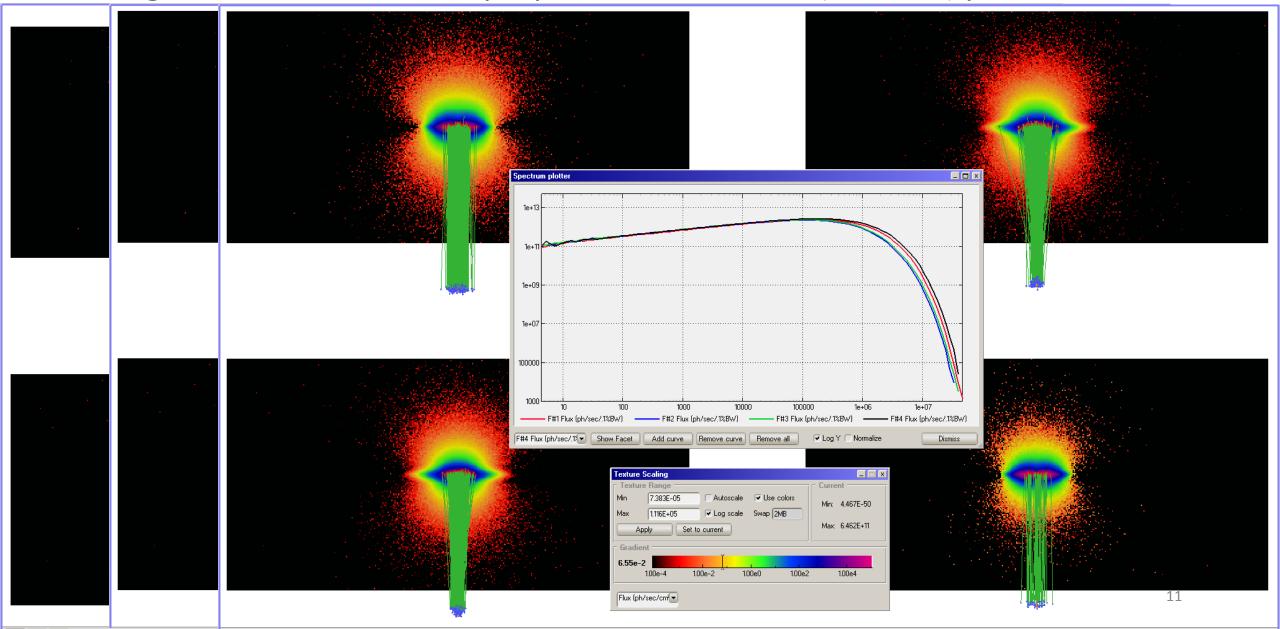
Molflow+: Pressure profile along one ~680 m-long arm of the interaction region of FCC-ee (175 GeV T-pole machine; 6.632 mA)





Average distance travelled by SR photons generated in the 2 doublets is ~ 63 m

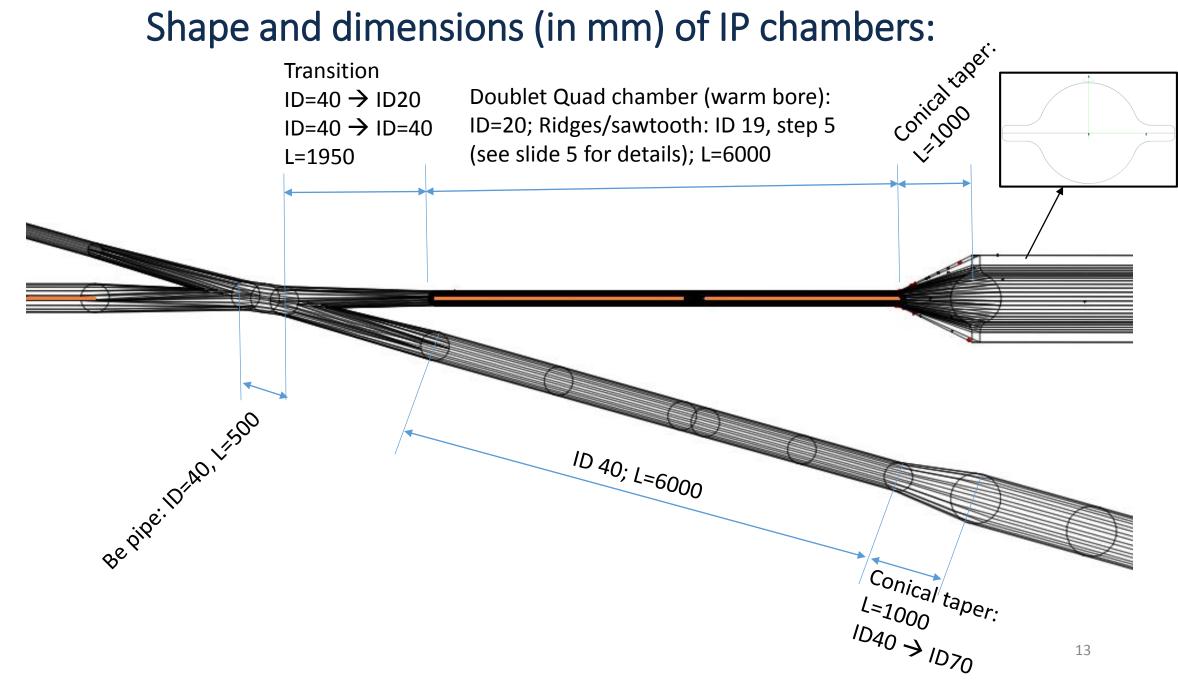
Synrad+: angular distribution of the SR generated along the 4 SC doublet magnets, viewed on a flat perpendicular screens (4x2 cm²) placed at 63 m



Preliminary conclusions:

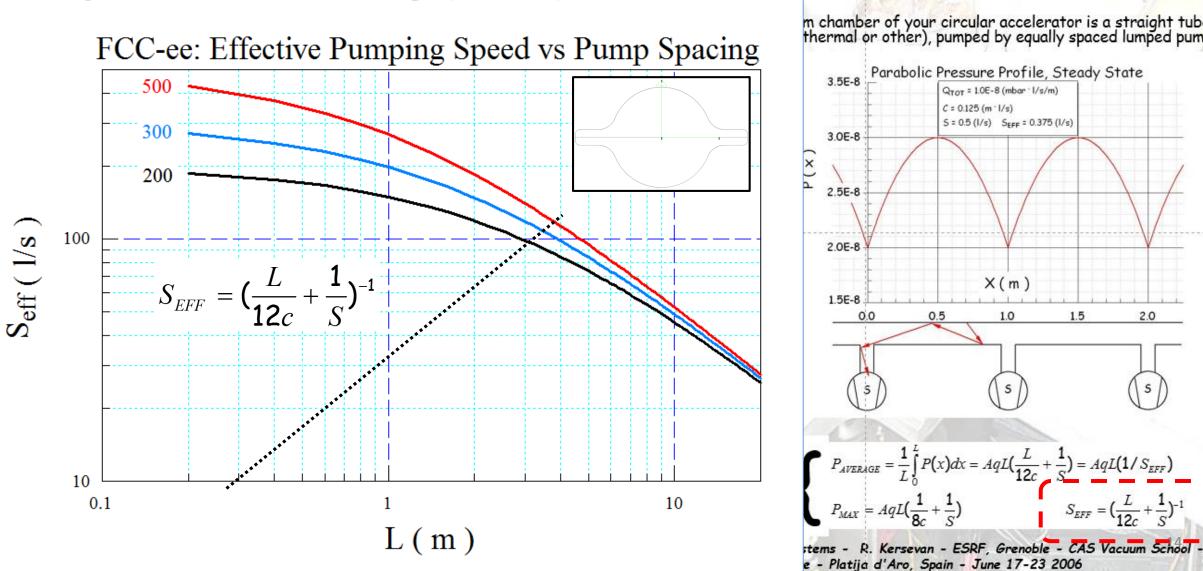
- > The ray-tracing montecarlo code SYNRAD+ has been applied to the FCC-ee IP region;
- A model of approximately 680 m length around the IP has been made: for the time being it doesn't have details about many important vacuum components, which could change the way low-energy photons are scattered (low-energy='those photons with energies below the Compton threshold', ~100 keV);
- It is evident that without a proper masking of the Be pipe, the pipe will get a non-negligible photon flux with photons up to several 10's keV: is this a problem for the detector's hardware and electronics?
- It is also evident that a rather simple to implement ridged (sawtooth-ed) geometry somehow machined on the internal part of the cold bore focusing doublet helps in reducing a lot (virtually to zero) the photon flux on the Be pipe; it needs to be coupled to a larger-bore 'exit' tube (which would also be beneficial for avoiding trapped modes in the Be pipe area (see E. Belli's presentation);
- The 4 quadrupole magnets of the doublet generate a rather large and extremely hard photon flux, with photons reaching the energy range of several TENS MeV: they mostly land on a small spot on the exit side of the beam, about ~63 m downstream, past the detector: careful shielding of that area must be envisaged;

Bonus slides (following yesterday's discussions):



Conductance limitation severely affects the effective pumping speed, and therefore the final pressure

Specific conductance of 70 mm ID, winged (SUPERKEK) cross-section: 48.65 1*m/s



m chamber of your circular accelerator is a straight tube, with thermal or other), pumped by equally spaced lumped pumps (L=1m): Parabolic Pressure Profile, Steady State $C = 0.125 (m \cdot 1/s)$ S = 0.5 (I/s) SEFF = 0.375 (I/s) 3.0E-8 2.5E-8 2.0E-8 X(m) $P_{AVERAGE} = \frac{1}{L} \int_{0}^{L} P(x) dx = AqL(\frac{L}{12c} + \frac{1}{S}) = \underline{AqL(1/S_{EFF})}$ $P_{MAX} = AqL(\frac{1}{8c} + \frac{1}{S})$ $S_{EFF} = (\frac{L}{12c} + \frac{1}{S})^{-1}$

Example: What average effective pumping speed would be necessary for the Z-pole machine in order to have an average pressure lower than 1.0E-9 mbar?

Photon flux:
$$F(ph/s) = 8.08E+17 \cdot E(GeV) \cdot I(mA) = 5.34E+22(ph/s)$$

If k=2.47E+19 (mol/mbar/liter), then the gas load Q is:

Q (mbar·l/s) = F (ph/s) / k(mol/mbar/l)
$$\cdot \eta$$
(mol/ph)

with η = photodesorption yield.

 \rightarrow Typically a machine is considered vacuum conditioned when η < 1.0E-6 (mol/ph):

```
Q' (mbar·l/s/m) = Q/ 2\pi\rho = 3.44E-8 (mbar·l/s/m) (in the arcs, \rho ~10 km) P_{avg} (mbar) = Q' (mbar·l/s/m) / S_{eff} (l/s) · L(m)
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Explicitating for S_{eff} , one gets: S_{eff} (I/s) = 34.42 · L(m) (dotted line at 45 deg on previous plot)

- This means that the Z-pole machine, due to its huge photon flux, would condition to P_{avg} <1.0E-9 only if 200 ÷ 500 l/s pumps are placed at a distance < 3 ÷ 4 m from each other;
- This, in turn, would mean that in the arcs one would need 19000 ÷ 25000 pumps/beam;
- Each lumped pump would in turn need pumping slots machined on the vacuum chamber: impedance budget?... and compare this to impedance generated by NEG-coating.

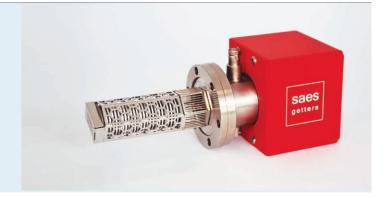
NEG pumps: the NEXTorr family (SAES Getters, Milan, Italy)

Integrated NEG pump with "small" noble ion-pump (~10 l/s)



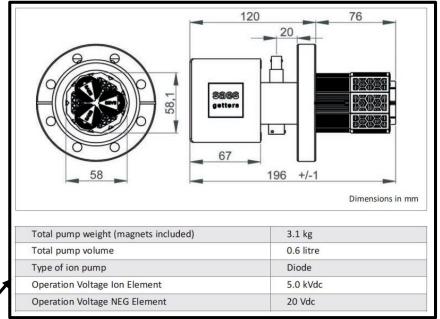


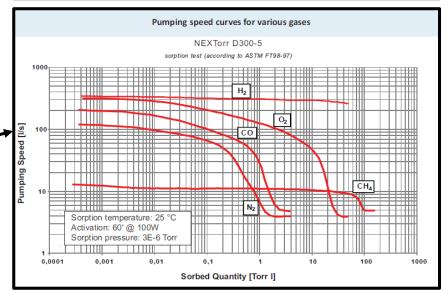
NEXTorr® D 200-5



NEXTorr® D 300-5







Twin-aperture magnets for FCC-ee and ante-chamber design compatibity:

PHYSICAL REVIEW ACCELERATORS AND BEAMS 19, 112401 (2016)

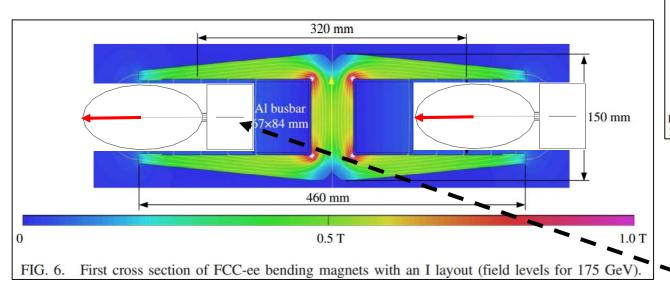
Efficient twin aperture magnets for the future circular e^+/e^- collider

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CERN-The European Organization for Nuclear Research, CH-1211 Geneva, Switzerland (Received 29 June 2016; published 2 November 2016)

We report preliminary designs for the arc dipoles and quadrupoles of the FCC-ee double-ring collider. After recalling cross sections and parameters of warm magnets used in previous large accelerators, we focus on twin aperture layouts, with a magnetic coupling between the gaps, which minimizes construction cost and reduces the electrical power required for operation. We also indicate how the designs presented may be further optimized so as to optimally address any further constraints related to beam physics, vacuum system, and electric power consumption.

DOI: 10.1103/PhysRevAccelBeams.19.112401



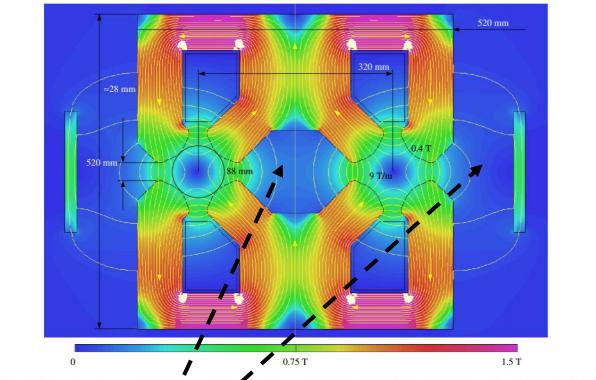


FIG. 7. First cross section of FCC-ee twin quadrupoles; the geometry is to scale with the twin dipole of Fig. 6 (field levels for 175 GeV).

Space for antechamber???? Yes for the quads...

... <u>difficult</u> for one of the dipoles!

17

Beam ionization and its electron density (follow up to E. Belli's presentation)

EUROPEAN LA	BORATORY FOR PARTICLE PHYSICS	
CERN - LHC DIVISION		
LHC-VAC/AGM	Vacuum Technical Note 96-01 January 1996	
Beam-Gas Ionis	sation Cross Sections at 7.0 TeV	
A.G. A	Mathewson and S. Zhang	

At energies greater than 100 keV the ionisation cross section σ for a gas by a PHYSIC **TOBER 1972** particle of charge Ze is given by: Ioniza ectrons and $\sigma = 4\pi Z^2 \left(\frac{\hbar}{mc}\right)^2 \frac{1}{\beta^2} \left(M^2 x + C\right)$ (2) σ is the ionisation cross section for the gas molecule (m²) where c is the speed of light=2.998 108 (m s⁻¹) v is the speed of the ionising particle (m s⁻¹) $\beta = v/c = 0.999999998203$ for 7.0 TeV protons m is the mass of the electron=9.109 10-31 (kg) M² and C are constants depending on the molecule and the function x is given by: $x = \ln \left(\frac{\beta^2}{1 - \beta^2} \right) - \beta^2$ (3)since β is ~1 this expression is more useful in the following form: $x = 2 \ln(\gamma) - \beta^2$ (4) where γ is the ratio of the energy of the proton relative to its rest mass and the rest mass of the proton =0.9383 GeV the rest mass of the electron =5.11 MeV and It is interesting that expression (2) depends only on the charge on the ionising particle and is independent of the mass i.e. in the high energy limit > 100keV, protons and electrons of the same velocity, β (or γ), have the same 18 ionising cross section. For example 7.0 TeV protons have the same ionising effect

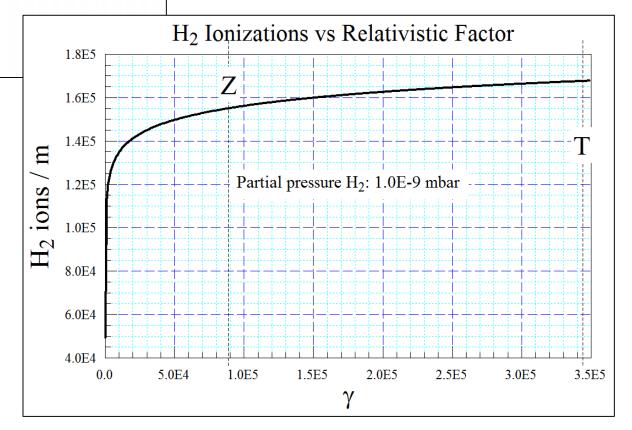
Beam ionization and its electron density

Putting in values for the fundamental constants the expression (2) for σ can be written:

$$\sigma = 1.874 \ 10^{-24} \frac{Z^2}{\beta^2} (M^2 x + C)$$
 m²

Table 2			
Gas	M ²	С	
H ₂	0.695	8.115	
Не	0.752	7.571	
CH4	4.23	41.85	
CO	3.70	35.14	
CO ₂	5.75	55.92	

Example: H_2 , 1.0E-9 mbar 1.0E-9 mbar $\rightarrow \rho$ = 2.47E+13 mol/m³



If the primary electrons traverse a path L in a gas of concentration ρ then the average number of primary ionizations z per primary electron is: $Z = \sigma \cdot L \cdot \rho$ Substituting for $n_{e-}(Z_{1450\,mA})=9.05E+13/m \rightarrow$ the average number of ionizations per meter length:

~ 1.55E+5 vs the 1.0E+8 assumed by E. Belli. ????

Negligible!