## Higgs Bosons in the NMSSM

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## Higgs content of the NMSSM:

- Two SU(2) doublets  $H_u$  (couples to up-type quarks) and  $H_d$  (couples to down-type quarks and leptons)  $\rightarrow$  like in the MSSM,
- A singlet S whose vev generates a Dirac mass term for higgsinos (replaces the  $\mu$  term of the MSSM)

## Mass eigenstates (- Goldstone boson):

- Three neutral scalars:
  - $H_{SM} 
    ightarrow M_{H_{SM}} \sim 125~{
    m GeV}$
  - $H_S \rightarrow$  mostly singlet-like
  - $H \rightarrow$  mostly MSSM-like, heavy, see below
- Two neutral pseudoscalars
  - $A_S \rightarrow \text{mostly singlet-like}$
  - $A \rightarrow \text{mostly MSSM-like, heavy, see below}$
- One charged Higgs  $H^{\pm} \rightarrow MSSM$ -like, heavy, see below

## Indirect constraints on the masses of the states beyond $H_{SM}$ :

 $H^{\pm}$  contributes to the  $BR(b \rightarrow s + \gamma)$  which is in agreement with the SM

ightarrow  $M_{H^\pm} \gtrsim$  350 GeV > scale of EW symmetry breaking

 $\rightarrow$  H<sup>±</sup>, H and A form a nearly degenerate SU(2) multiplet with M  $\gtrsim$  350 GeV

BUT: The masses  $M_{H_S}$ ,  $M_{A_S}$  of the mostly singlet-like (pseudo-)scalars depend on unknown parameters and can vary from 0...1000 GeV (and are different);

 $M_{H_S} \sim 60 - 110$  GeV is natural, helps to explain  $M_{H_{SM}} \sim 125$  GeV without inducing a too large  $BR(H_{SM} \rightarrow H_S H_S)$  which could reduce the SM-like branching fractions like  $H_{SM} \rightarrow Z^*Z$  below its measured values

# Couplings of the BSM Higgs states to SM particles (essential for searches!):

Use:  $\tan \beta = \frac{\langle H_u \rangle}{\langle H_d \rangle} = 1 \dots 60;$ 

MSSM (typically): tan  $eta\gtrsim$  10 for  $M_{H_{SM}}\sim$  125 GeV,

NMSSM (typically): tan  $\beta \sim 2-4$  for NMSSM-specific contributions to  $M_{H_{SM}}$ 

• *H*/*A*:

- $g_{Htt}$  suppressed by  $1/\tan\beta$  (still dominant for  $\tan\beta < 60$ )
- $g_{Hbb}, g_{H\tau\tau}$  enhanced by tan  $\beta$
- $g_{HWW}$ ,  $g_{HZZ}$  strongly suppressed for  $M_H \gg M_{H_{SM}}$  ("alignment")
- $H_S$ : Mixes with  $H_{SM}$  and/or H
  - $M_{H_S} < 125$  GeV: most couplings like  $H_{SM}$  reduced by a common factor (mixing angle)  $\rightarrow$  very similar branching fractions, except: Couplings of  $H_S$  to  $\gamma\gamma$  can be considerably enhanced! (Due to a reduced coupling to  $b\bar{b}$ )
  - $M_{H_S} \gg 125$  GeV: most couplings like H reduced by a common factor  $\rightarrow$  very similar branching fractions

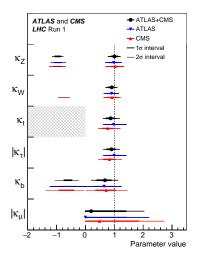
Note: Mixing between  $H_S$  and  $H_{SM}$  reduces the couplings of " $H_{SM}$ " to ZZ, WW below their SM values (corresponding to  $\kappa = 1$ );

Run I ATLAS and CMS combination:

From  $\kappa^2_{H_{SM}ZZ} + \kappa^2_{H_5ZZ} \lesssim 1$  and  $\kappa_{H_{SM}ZZ} > 0.84$  (at  $2\sigma$ ):

 $\kappa_{H_SZZ} (= \kappa_{H_SWW}) \lesssim 0.5$ 

The (relative)  $H_S - ZZ/WW$  couplings squared can be at most  $\sim 0.25$ 

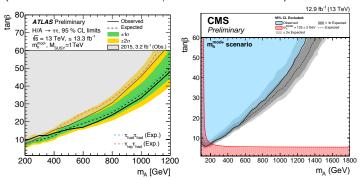


#### Possible direct production of the BSM Higgs states:

• *H*/*A*:

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- ggF: dominant, but large QCD background to  $H \rightarrow tt$ ,  $H \rightarrow bb$
- Ass. production with *b* quarks: enhanced for large tan  $\beta$  where also the BRs into *bb*,  $\tau\tau$  become larger; *bb* + resonant  $\tau\tau$  final state: background under control, but: no excess  $\rightarrow$  constraints in the plane tan  $\beta - M_A$  ( $\simeq M_H, M_{H^{\pm}}$ ) (from ATLAS-CONF-2016-085/CMS-PAS-HIG-16-037):



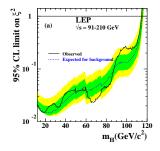
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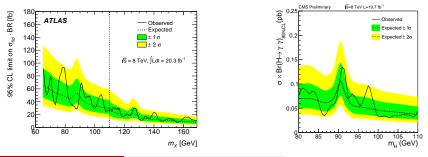
• *H<sub>S</sub>*/*A<sub>S</sub>*:

LEP search for a light scalar with reduced coupling  $\xi^2$  to ZZ (recall:  $\xi^2 \lesssim 0.25$  from  $H_{SM}ZZ$  coupling):

The region in the  $\xi^2 - m_H$  plane below the black line is allowed

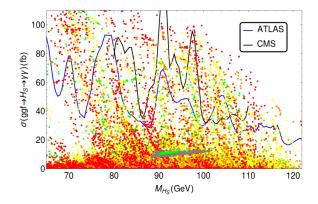


ATLAS/CMS searches for  $ggF \rightarrow H_S \rightarrow \gamma\gamma$  at 8 TeV:



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Do the ATLAS/CMS searches touch possible values for  $\sigma(ggF \rightarrow H_S \rightarrow \gamma\gamma)$  within the LEP-allowed NMSSM parameter space? (M. R.-Vázquez, U.E.):



YES, but far from exclusion... even light  $H_S/A_S$  states may have too small direct production cross sections for discovery, even at 13 TeV

Possible indirect production of the BSM Higgs states:

Potentially (relatively) large NMSSM-specific trilinear Higgs couplings:

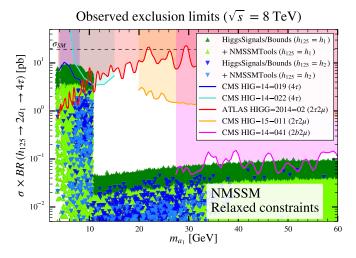
 $g_{H_{SM}H_SH_S}$ ,  $g_{H_{SM}A_SA_S}$ ,  $g_{HH_SH_{SM}}$ 

 $\rightarrow$  allow for decays

- $H_{SM} \rightarrow H_S H_S$ ,  $H_{SM} \rightarrow A_S A_S$ (if kinematically allowed, i.e.  $M_{H_S,A_S} \lesssim 60$  GeV)
- $H \rightarrow H_{SM}H_S$

Exotic  $H_{SM}$  decays reduce its SM-like branching ratios, and are limited by its SM-like signal rates (see above); still ...

## Run I searches for $H_{125} \rightarrow A_S A_S (H_S H_S) \rightarrow 4$ leptons (From R. Aggleton et al., JHEP 1702 (2017) 035, arXiv:1609.06089)



Light green/blue points: viable in the NMSSM after LEP/other LHC constraints

 $\rightarrow$  These searches for  $A_S/H_S$  have only scratched the NMSSM parameter space ...

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## ... and are limited to $M_{H_s,A_s} \lesssim 60$ GeV; how to search for heavier $H_s/A_s$ ?

- $\rightarrow$  Recall:  $g_{HH_SH_{SM}}$  can be large (in contrast to  $g_{HH_{SM}H_{SM}}$ ,  $g_{HH_SH_S}$ )
- → The  $BR(H \rightarrow H_S H_{SM})$  can be large (~ 30%, competing only with  $H \rightarrow t\bar{t}$ , reducing *BR* for the search into  $\tau\tau$ )
- $\rightarrow$  Looks like resonant Higgs pair production, but with one SM Higgs replaced by  $H_S$  with unknown mass in the range up to  $M_H 125$  GeV
- $\rightarrow$  Look for  $b\bar{b}b\bar{b}$  (4 *b*-tagged jets) with
  - one *bb* pair:  $M_{bar{b}} \sim 125$  GeV,
  - another  $b\bar{b}$  pair:  $M_{b\bar{b}} \sim M_{Hs}$  (unknown),
  - $M_{b\bar{b}b\bar{b}} \sim M_H$  (unknown)

 $(b\bar{b}\tau\tau$  final states are slightly less promising;  $b\bar{b}\gamma\gamma$  final states possibly promising for  $M_H \lesssim 500$  GeV, under study)

## Best Strategy (M. R.-Vázquez)

(Borrowed from ATLAS/CMS searches for resonant SM Higgs pair production)

Use "test" mass  $M_{b\bar{b}} \sim M_{H_S}$ ; for given  $M_{H_S}$ : Optimise the pairing of 4 *b*-tagged jets into 2 × 2 *b*-tagged jets, cut on *bb* masses arond 115 GeV and  $M_{H_S} - 10$  GeV (allow for "losses" outside the R = 0.4 - jets)

Study the distribution of  $M_{b\bar{b}b\bar{b}}$  from the 4 *b*-tagged jets (after correcting  $M_{b\bar{b}}$  near 115 GeV to 125 GeV)

 $\rightarrow$  The (by far) dominant 4b QCD background is a smoothly decreasing function of  $M_{b\bar{b}b\bar{b}}$ , obtained from SHERPA (during a secondment in Durham)

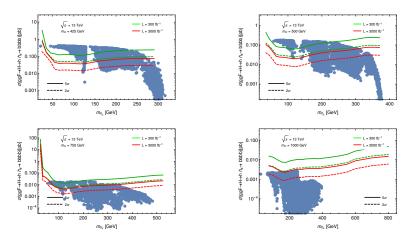
→ If  $M_{H_S}$  was chosen correctly, one observes a "bump" in  $M_{b\bar{b}b\bar{b}}$  near  $M_H$  whose significance can be computed as function of  $M_{H_S}$ ,  $M_H$  and notably the  $\sigma(ggF \rightarrow H \rightarrow H_{SM} + H_S \rightarrow b\bar{b}b\bar{b})$ 

 $\rightarrow$  Expected 2  $\sigma$  exclusion limits and 5  $\sigma$  discovery limits can be obtained

→ These are model independent (assuming just a width not larger than a few GeV), but can be compared to possible values for  $\sigma(ggF \rightarrow H \rightarrow H_{SM} + H_S \rightarrow b\bar{b}b\bar{b})$  in the NMSSM:

## Expected sensitivities to $\sigma(ggF \rightarrow H \rightarrow H_{SM} + H_S \rightarrow b\bar{b}b\bar{b})$ as function of $M_{HS}$ :

Upper left: assuming  $M_H = 425$  GeV, upper right: assuming  $M_H = 500$  GeV Lower left: assuming  $M_H = 750$  GeV, lower right: assuming  $M_H = 1000$  GeV Full/dotted green:  $5/2\sigma$  excess for  $L = 300 fb^{-1}$ Full/dotted red:  $5/2\sigma$  excess for  $L = 3000 fb^{-1}$ ; Blue: NMSSM points



# Conclusions

- The NMSSM contains a rich BSM Higgs sector which is hardly tested
- Mostly singlet-like scalars  $H_S$ /pseudoscalars  $A_S$  (light < 125 GeV, or heavy > 125 GeV) are far from being excluded; their direct production is detectable only in some regions of the NMSSM parameter space
- If lighter than 60 GeV they may be detectable via  $H_{SM}$  decays (many ongoing ATLAS/CMS studies on exotic Higgs decays)
- Otherwise they may be detectable via  $H \rightarrow Z + A_S$  (ongoing ATLAS/CMS studies) or  $H \rightarrow H_{SM} + H_S$  decays (see above, no LHC studies yet)

(These decays, if taken into account, reduce the sensitivities to MSSM-like H/A states in the standard  $b\bar{b}\tau\tau$  channel!)