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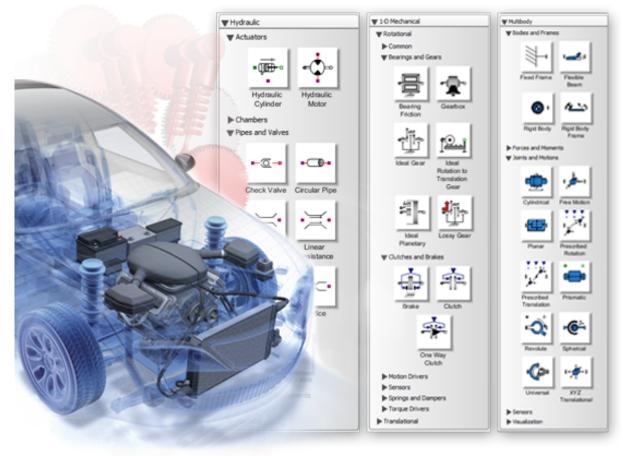
VHiggsTools & Maplesoft

- Giulia Gonella interactive equation solving
- Agnieszka Ilnicka generalizing bivariate limit code & extending Feynman diagram code
- Raquel Gomez evaluating Magma library for matrix algebra on GPUs (like LAPACK)
- Hjalte Frellesvig implement the generalized polylogarithm in Maple

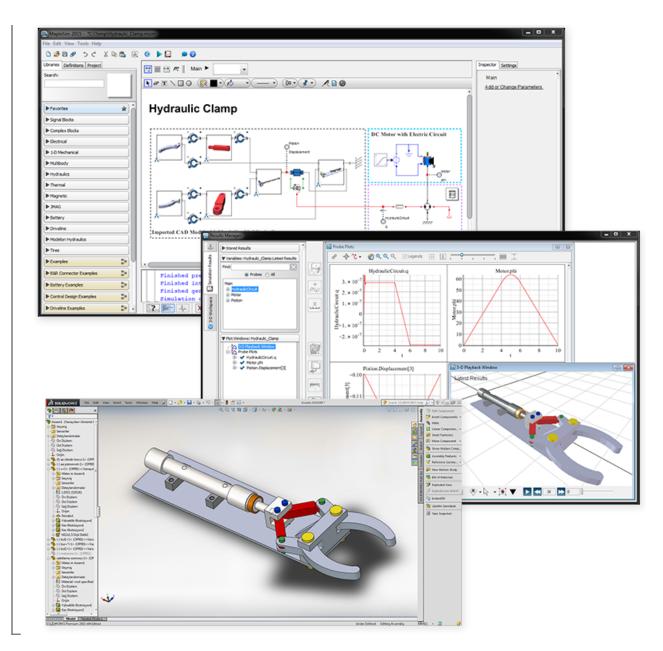
What's new at Maplesoft

MapleSim

• Multi-Domain Physical Modeling and Simulation



• Fast simulation and visualization



Möbius

- Online courseware environment that focuses on science, technology, engineering, and mathematics
- Automated assessment
- Gradebook and analytics

Maple

• At the heart of it all

What's new with Maple

Kernel

- Multithreaded garbage collection, mostly in a thread separate from the "main" thread
- Updated builtin libraries for polynomial arithmetic

GUI

• Construct collections of user interface elements programmatically - see this example

• Workbook: collection of worksheets and data files

Import("this://IRIS.csv");

[1	2	3	4	5	6	7	8		
1	5.1	0.222222222	3.5	0.625	1.4	0.06779661	0.2	0.041666667		
2	4.9	0.166666667	3	0.416666667	1.4	0.06779661	0.2	0.041666667		
3	4.7	0.111111111	3.2	0.5	1.3	0.050847458	0.2	0.041666667		
4	4.6	0.083333333	3.1	0.458333333	1.5	0.084745763	0.2	0.041666667		(2.2.1)
5	5	0.19444444	3.6	0.666666667	1.4	0.06779661	0.2	0.041666667		(3.2.1)
6	5.4	0.305555556	3.9	0.791666667	1.7	0.118644068	0.4	0.125		
7	4.6	0.083333333	3.4	0.583333333	1.4	0.06779661	0.3	0.083333333		
8	5	0.19444444	3.4	0.583333333	1.5	0.084745763	0.2	0.041666667		

Math library

• Most of the stuff that will interest you most

What's new with the math library

ThermophysicalData

with(ThermophysicalData) :
Property(temperature_triple, Water, useunits);
273.1600000 K

(4.1.1)

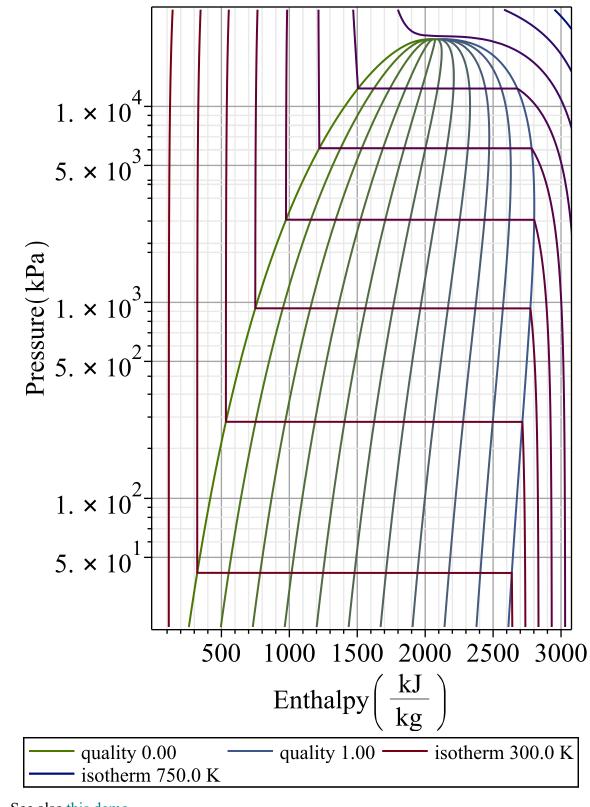
(4.1.2)

Property (*temperature_dew_point*, *HumidAir*, *temperature_dry_bulb* = 300, *pressure* = 1 atm,

humidity =
$$\frac{1}{2}$$

288.7139414 K

PHTChart(Water)



See also <u>this demo</u>.

Bivariate limits $expr := \frac{x^3 \cdot y}{x^6 + y^2}$

$$expr := \frac{x^{3}y}{x^{6} + y^{2}}$$
(4.2.1)

$$expr_line := expr|_{y=ax}$$

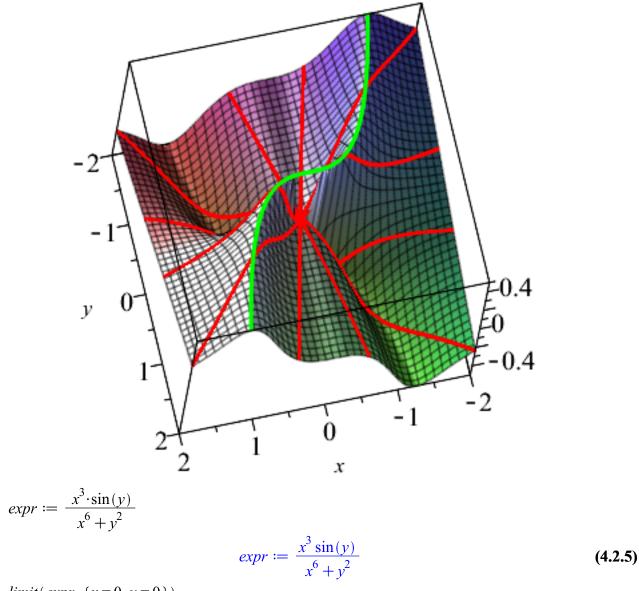
$$expr_line := \frac{x^{4}a}{x^{6} + a^{2}x^{2}}$$
(4.2.2)

$$\lim_{x \to 0} expr_line$$
0
(4.2.3)

$$curves := plots[display](seq[plots[spacecurve]([x, a x, expr_line], x = max(-2, -\frac{2}{|a|})..min(2, \frac{2}{|a|}), color = red, thickness = 3), a = [-4, -1, -\frac{1}{4}, \frac{1}{4}, 1, 4]))$$
0,4
0,4
0,4
0,2
0,-1
0,2
-1
0,-1
0,-1
1
limit(expr, {x=0, y=0})
undefined
(4.2.4)

plots[display] (plot3d(expr, x = -2..2, y = -2..2, grid = [200, 200]), curves, plots:-

spacecurve
$$\left(\left[x, x^3, eval(expr, y = x^3) \right], x = -2^{\frac{1}{3}} \dots 2^{\frac{1}{3}}, color = green, thickness = 3 \right) \right)$$



limit(*expr*, $\{x = 0, y = 0\}$)

undefined (4.2.6)

This will be computed correctly in Maple 2017!

Partial differential equations

$$pde := \frac{\partial}{\partial x} f(x, y, z) + \frac{\partial}{\partial y} f(x, y, z) + \frac{\partial}{\partial z} f(x, y, z) = f(x, y, z)$$
$$pde := \frac{\partial}{\partial x} f(x, y, z) + \frac{\partial}{\partial y} f(x, y, z) + \frac{\partial}{\partial z} f(x, y, z) = f(x, y, z)$$
(4.3.1)

pdsolve(pde)

$$f(x, y, z) = FI(-x + y, -x + z) e^{x}$$
(4.3.2)

$$bc := f(\alpha + \beta, \alpha - \beta, 1) = \alpha \cdot \beta$$

$$bc := f(\alpha + \beta, \alpha - \beta, 1) = \alpha \beta$$
(4.3.3)

pdsolve([bc, pde])

$$f(x, y, z) = \frac{(x - 2z + 2 + y) (x - y) e^{z - 1}}{4}$$
(4.3.4)

How does this work?

$$eval((4.3.2), [x = alpha + beta, y = alpha - beta, z = 1])$$

$$f(\alpha + \beta, \alpha - \beta, 1) = FI(-2\beta, -\alpha - \beta + 1) e^{\alpha + \beta}$$

$$(4.3.1.1)$$

eval(%, *bc*)

$$\alpha \beta = F1(-2\beta, -\alpha - \beta + 1) e^{\alpha + \beta}$$
(4.3.1.2)

 $isolate(\%, op(indets(\%, specfunc(_F1)))))$

$$FI(-2\beta, -\alpha - \beta + 1) = \frac{\alpha\beta}{e^{\alpha + \beta}}$$
(4.3.1.3)

solve([op(lhs((4.3.1.3)))] = [x0, y0], [alpha, beta]);

$$\left[\left[\alpha = 1 - y\theta + \frac{x\theta}{2}, \beta = -\frac{x\theta}{2} \right] \right]$$
(4.3.1.4)

eval((4.3.1.3), %[1])

$$FI(x0, y0) = -\frac{\left(1 - y0 + \frac{x0}{2}\right)x0}{2e^{1-y0}}$$
(4.3.1.5)

eval((4.3.2), F1 = unapply(rhs(%), [x0, y0]))

$$f(x, y, z) = -\frac{\left(\frac{x}{2} - z + 1 + \frac{y}{2}\right)(-x + y)e^{x}}{2e^{x - z + 1}}$$
(4.3.1.6)

simplify(%)

$$f(x, y, z) = \frac{(x - 2z + 2 + y) (x - y) e^{z - 1}}{4}$$
(4.3.1.7)

$$pde := 3 (u(x, y) - y)^{2} \left(\frac{\partial}{\partial x} u(x, y)\right) - \left(\frac{\partial}{\partial y} u(x, y)\right) = 0$$

$$pde := 3 (u(x, y) - y)^{2} \left(\frac{\partial}{\partial x} u(x, y)\right) - \left(\frac{\partial}{\partial y} u(x, y)\right) = 0$$
(4.3.5)

pdsolve(*pde*)

$$u(x, y) = RootOf(-y^{3} + 3y^{2} Z - 3y Z^{2} + Z^{3} - FI(Z) - x)$$

$$DETools[remove_RootOf](\%)$$
(4.3.6)

$$-y^{3} + 3y^{2}u(x, y) - 3yu(x, y)^{2} + u(x, y)^{3} - FI(u(x, y)) - x = 0$$
(4.3.7)

 $eval(\%, _F1 = \sin)$ $-y^3 + 3y^2 u(x, y) - 3y u(x, y)^2 + u(x, y)^3 - \sin(u(x, y)) - x = 0$ (4.3.8)

pdetest(%, pde)

$$bc := u(0, \alpha) = \alpha$$

$$pdsolve([pde, bc])$$

$$u(x, y) = x^{1/3} + y, u(x, y) = -\frac{x^{1/3}}{2} - \frac{1\sqrt{3}x^{1/3}}{2} + y, u(x, y) = -\frac{x^{1/3}}{2}$$

$$+ \frac{1\sqrt{3}x^{1/3}}{2} + y$$

$$(4.3.11)$$

Units

Consider:

 $d \coloneqq 5 \text{ m}$

 $d \coloneqq 5 \text{ m} \tag{4.4.1}$

 $t \coloneqq 5 \text{ s}$

$$t := 5 \text{ s}$$
 (4.4.2)

$$(x + y \cdot d) \cdot (y + t \cdot x)$$

This is a violation of unit consistency: the first factor means that $\frac{x}{y}$ has dimension length,

whereas the second factor implies that $\frac{y}{x}$ has dimension time. That cannot be. Units:-TestDimensions($(x + y \cdot d) \cdot (y + t \cdot x)$); false (4.4.3)

How does this work? Every expression is taken apart and its dimension expressed in terms of the dimensions of its subexpressions; concrete units are expanded in terms of independent base dimensions. Subexpressions that we don't know anything about (such as x, y, f(...)) remain; conceptually, $y \cdot d$ gets turned into $dimension(y) \cdot length$. Inequalities, sums, and equations are recorded: their operands all have the same dimension; we turn that into expressions that must be dimensionless. In the example above:

$$exprs := \left[\frac{dimension(x)}{dimension(y) \cdot length}, \frac{dimension(y)}{dimension(x) \cdot time}\right]: \text{ might get represented as}$$

$$A := \langle 1, -1; -1, 1 \rangle$$

$$A := \left[\begin{array}{c} 1 & -1 \\ -1 & 1 \end{array}\right]$$

$$(4.4.4)$$

$$C := \langle -1, 0; 0, -1 \rangle$$

$$C := \begin{bmatrix} -1 & 0 \\ 0 & -1 \end{bmatrix}$$
(4.4.5)
If we knew the dimensions of x and y, then we could express them in the same way: as an Ansatz

If we knew the dimensions of *x* and *y*, then we could express them in the same way; as an Ansatz, suppose *x* is a velocity and *y* is an area: $B := \langle 1, -1; 2, 0 \rangle$

$$B := \begin{bmatrix} 1 & -1 \\ 2 & 0 \end{bmatrix}$$
(4.4.6)

Now $A \cdot B + C$ give us the dimension of the expressions in *exprs*: $A \cdot B + C$

$$\begin{bmatrix} -2 & -1 \\ 1 & 0 \end{bmatrix}$$
(4.4.7)

Clearly we have failed to make these expressions dimensionless. And we have just shown that $A \cdot B = -C$ has a solution if and only if there is a consistent assignment of dimensions to our atomic expressions. In this case, there is no solution: LinearAlgebra:-LinearSolve(A, -C);

Error, (in LinearAlgebra:-BackwardSubstitute) inconsistent system

For another example, replace one y with z: $(x + y^2 \cdot d) \cdot (z + t \cdot x)$. Now $exprs = \left[\frac{dimension(x)}{dimension(y)^2 \cdot length}, \frac{dimension(z)}{dimension(x) \cdot time}\right],$ $A := \langle 1, -2, 0; -1, 0, 1 \rangle$ $A := \begin{bmatrix} 1 & -2 & 0 \\ -1 & 0 & 1 \end{bmatrix}$ (4.4.8) $C := \langle -1, 0; 0, -1 \rangle$

 $C := \begin{bmatrix} -1 & 0 \\ 0 & -1 \end{bmatrix}$ (4.4.9)

B := LinearAlgebra:-LinearSolve(A, -C);

$$B := \begin{bmatrix} -t_{1,1} & -1 + t_{1,2} \\ -\frac{1}{2} + \frac{-t_{1,1}}{2} & -\frac{1}{2} + \frac{-t_{1,2}}{2} \\ -t_{1,1} & -t_{1,2} \end{bmatrix}$$
(4.4.10)

B := eval(B, indets(B) = 1)

$$B := \begin{bmatrix} 1 & 0 \\ 0 & 0 \\ 1 & 1 \end{bmatrix}$$
(4.4.11)

 $A \cdot B + C \begin{bmatrix} 0 & 0 \\ 0 & 0 \end{bmatrix}$ $Units:-TestDimensions((x + y^2 \cdot d) \cdot (z + tx));$ (4.4.12)

$$\frac{true}{true}$$
(4.4.13)

Physics

See the Physics examples help page:

```
> restart;
with(Physics);
Setup(mathematicalnotation = true);
```

['*', `.', Annihilation, AntiCommutator, Antisymmetrize, Assume, Bra, Bracket, Check, Christoffel, Coefficients, Commutator, CompactDisplay, Coordinates, Creation, D, Dagger, Decompose, Define, Dy, Einstein, EnergyMomentum, Expand, ExteriorDerivative, Factor, FeynmanDiagrams, Fundiff, Geodesics, GrassmannParity, Gtaylor, Intc, Inverse, Ket, KillingVectors, KroneckerDelta, LeviCivita, Library, LieBracket, LieDerivative, Normal, Parameters, PerformOnAnticommutativeSystem, Projector, Psigma, Redefine, Ricci, Riemann, Setup, Simplify, SpaceTimeVector, StandardModel, SubstituteTensor, SubstituteTensorIndices, SumOverRepeatedIndices, Symmetrize, TensorArray, Tetrads, ThreePlusOne, ToFieldComponents, ToSuperfields, Trace, *TransformCoordinates*, *Vectors*, *Weyl*, `^`, *dAlembertian*, *d_*, *diff*, *g_*, *gamma_*] [mathematicalnotation = true] (4.5.1)Consider two conjugate observables Q, P, and the corresponding Hermitian operators satisfying $[\underline{Q}, \underline{P}] = I\hbar.$ > macro (h = ` & hbar; `): > Setup (hermitianoperators = {Q,P}, %Commutator(Q,P) = I*h); [algebrarules = { $[Q, P]_= = I\hbar$ }, hermitian operators = { P, Q }] (4.5.2)Suppose now that the system where Q and P act is in some state $|\psi\rangle$ normalized to 1, and set $|\psi\rangle$ as the default state for computing <u>Brackets</u>. > Ket(psi); $|\psi\rangle$ (4.5.3)> Dagger(%) = Bra(psi); $\langle \psi | = \langle \psi |$ (4.5.4)> Bra(psi) . Ket(psi); $\langle \psi | \psi \rangle$ (4.5.5)> Bracket(psi, psi); $\langle \psi | \psi \rangle$ (4.5.6)> Setup(Bracket(psi, psi) = 1, bracketbasis = psi); [*bracketbasis* = ψ , *bracketrules* = { $\langle \psi | \psi \rangle = 1$ }] (4.5.7)We now have: > Ket(psi); $|\psi\rangle$ (4.5.8)> Bra(psi) . %; (4.5.9)The mean values of the operators Q and P in the state $|\psi\rangle$ are then given by: > Qm := Bracket(Q); #Shortcut for Bracket(psi, Q, psi) after having set the bracketbasis to psi $Qm := \langle Q \rangle$ (4.5.10)> Pm := Bracket(P); $Pm := \langle P \rangle$ (4.5.11)Let's introduce another Hermitian operator, Δ , and denote $\Delta(Q)$ and $\Delta(P)$ the operators

representing the observable deviations from these mean values by $\langle Q \rangle$ and $\langle P \rangle$.

> Setup(hermitianoperators = Delta); [hermitianoperators = { Δ , P, Q}] (4.5.12) \rightarrow DefDQ := Delta(Q) = Q - Bracket(Q); $DefDQ := \Delta(Q) = Q - \langle Q \rangle$ (4.5.13)> DefDP := Delta(P) = P - Bracket(P); $DefDP := \Delta(P) = P - \langle P \rangle$ (4.5.14)The value of the <u>Commutator</u> between $\Delta(Q)$ and $\Delta(P)$ is a consequence of the value of the Commutator between Q and P, and so it can be computed by rewriting the deviations in terms of Q and P. > %Commutator(Delta(Q), Delta(P)); $\left[\Delta(Q), \Delta(P)\right]$ (4.5.15)> eval(%, {DefDQ, DefDP}); $[Q - \langle Q \rangle, P - \langle P \rangle]_{-}$ (4.5.16)=
> expand(%); OP - PO(4.5.17)> Simplify(%); Ιħ (4.5.18)> eval(Commutator(Delta(Q), Delta(P)), {DefDQ, DefDP}); (4.5.19)Track this result as an algebra rule, so that in what follows we compute directly with $\Delta(Q)$ and $\Delta(P)$. > Setup((4.5.15) = (4.5.19)); $\left[algebrarules = \left\{ \left[Q, P\right]_{-} = I\hbar, \left[\Delta(Q), \Delta(P)\right]_{-} = I\hbar \right\} \right]$ (4.5.20)To show now that $[Q, P]_{-} = I\hbar$ implies $\frac{\hbar^2}{4} \leq \langle \Delta(P)^2 \rangle \langle \Delta(Q)^2 \rangle$, consider the action of these deviation operators $\Delta(Q)$ and $\Delta(P)$ on the state of the system $|\psi\rangle$, and construct with them a new Ket involving a real parameter λ . > Ket(Psi, lambda) := (Delta(Q) + I*lambda*Delta(P)) . Ket(psi) ; $\left| \Psi_{\lambda} \right\rangle := \Delta(Q) \cdot \left| \psi \right\rangle + \mathrm{I}\lambda \left(\Delta(P) \cdot \left| \psi \right\rangle \right)$ (4.5.21)The square of the norm of $|\Psi_{\lambda}\rangle$, for λ real, is > Dagger(%) . % assuming lambda::real; $\langle \Delta(P)^2 \rangle \lambda^2 - \mathrm{I} \langle \Delta(P) \Delta(Q) \rangle \lambda + \mathrm{I} \lambda \langle \Delta(Q) \Delta(P) \rangle + \langle \Delta(Q)^2 \rangle$ (4.5.22)<u>Simplify</u> this norm, taking into account the commutator $[\Delta(Q), \Delta(P)] = I\hbar$, set in (4.5.20) > Simplify(%); $-\hbar \lambda + \langle \Delta(P)^2 \rangle \lambda^2 + \langle \Delta(Q)^2 \rangle$ (4.5.23)This is a polynomial in λ of second degree; its <u>discriminant</u> is negative or zero. > discrim(%, lambda) <= 0;</pre> $\hbar^2 - 4 \left\langle \Delta(P)^2 \right\rangle \left\langle \Delta(Q)^2 \right\rangle \leq 0$ (4.5.24)isolating $\frac{\hbar^2}{4}$, we obtain the lower bound for $\langle \Delta(P)^2 \rangle \langle \Delta(Q)^2 \rangle$.

• isolate (%, h^2) /4;
$$\frac{\hbar^2}{4} \leq \left\langle \Delta(P)^2 \right\rangle \left\langle \Delta(Q)^2 \right\rangle$$
(4.5.25)

Note that this result is a consequence of $[\Delta(Q), \Delta(P)] = I\hbar$, which in turn is a consequence of $[Q, P] = I\hbar$, so that Q and P too satisfy $\frac{\hbar^2}{4} \leq \langle P^2 \rangle \langle Q^2 \rangle$, and in fact the product of *any* two conjugate Hermitian operators, as well as of the root-mean square deviations of them, satisfy this inequality.