

PHOTON 2017

CERN

May 26th, 2017

Isolated photon, photon+jet and diphoton results in ATLAS

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On behalf of the ATLAS Collaboration

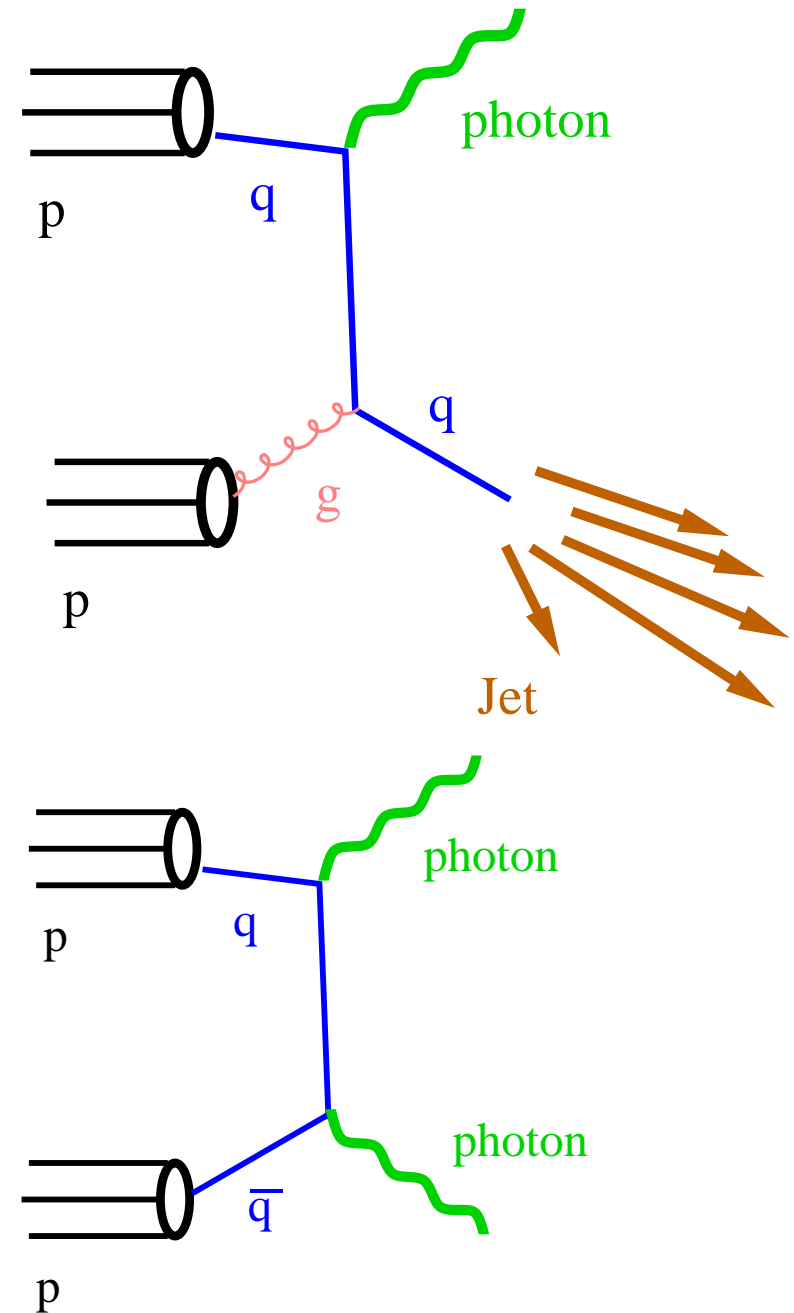
● **Outline**

- **Physics with photons**
- **Inclusive photon production at 13 TeV**
- **Photon + jet(s) production at 8 TeV**
- **Photon pair production at 8 TeV**
- **Summary**



Photon production in pp collisions at LHC

- Photon production in pp collisions
 - allows tests of perturbative QCD predictions
 - provides information on the proton PDFs
- Possibilities to study inclusive production of photons or in association with jets
- Prompt photons represent a cleaner probe of the hard interaction than jet production
- Prompt-photon measurements aid searches involving photons or $E_T^{\text{miss}} + \text{jets}$ (through ratios $Z + \text{jets} / \gamma + \text{jets}$)
- Diphoton production is of special interest as the major background to $H \rightarrow \gamma\gamma$

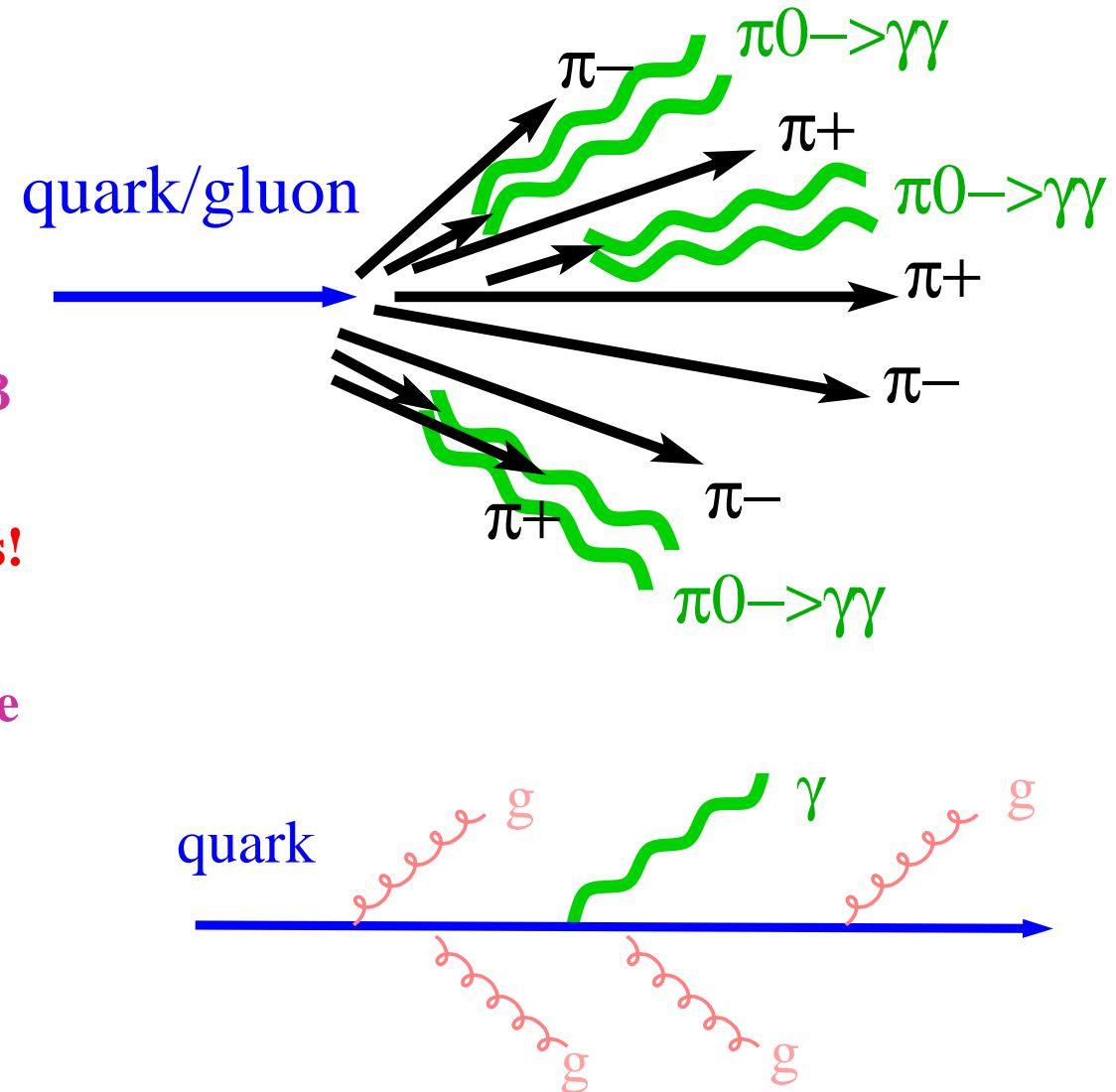


Other sources of photons

- Quarks and gluons are sources of photons

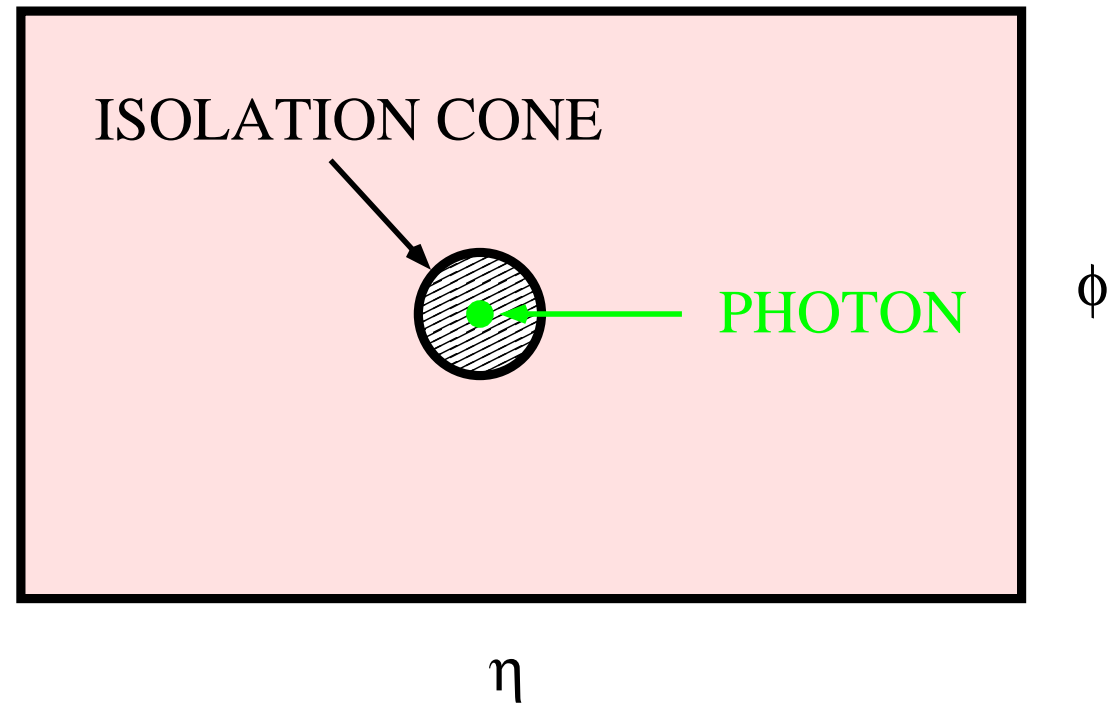
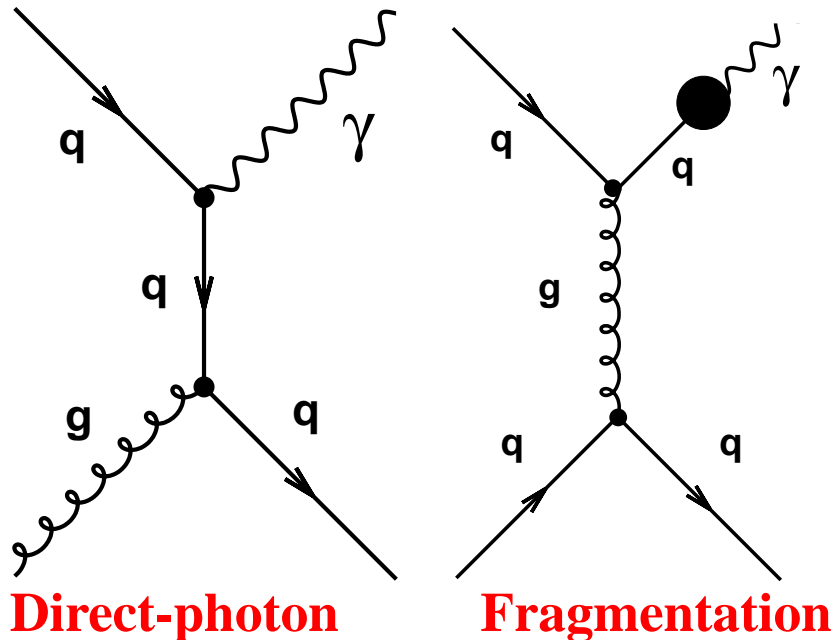
→ Quarks and gluons fragment mostly into pions and, by isospin symmetry, 1/3 are π^0 's, which decay into two photons
 ⇒ γ 's are produced copiously inside jets!

→ Quarks have electric charge and radiate photons
 ⇒ fragmentation function $D_{q/g}^\gamma(z, \mu_f)$



⇒ Distinct feature: these photons are inside jets, i.e. not isolated!

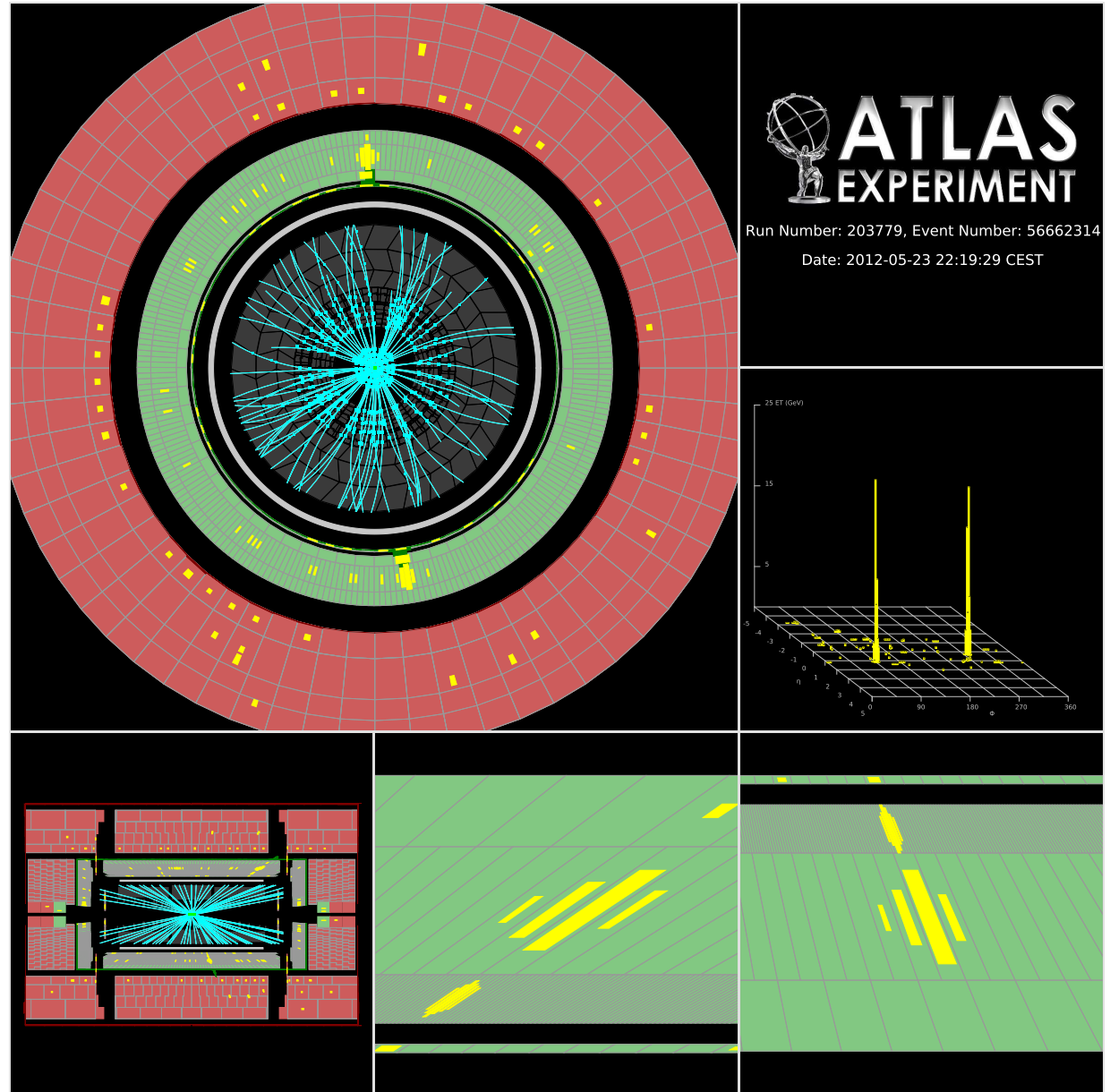
Photon isolation



- It is essential to require the photon to be isolated. It is achieved by requiring $E_T^{iso} \equiv \sum_i E_T^i < E_T^{max}$ with the sum over the particles (except the photon!) inside a cone of radius $R = 0.4$ centered on the photon in the $\eta - \phi$ plane
- The isolation requirement suppresses the contribution of photons inside jets: π^0 (as well as other neutral mesons) decays and the fragmentation contribution

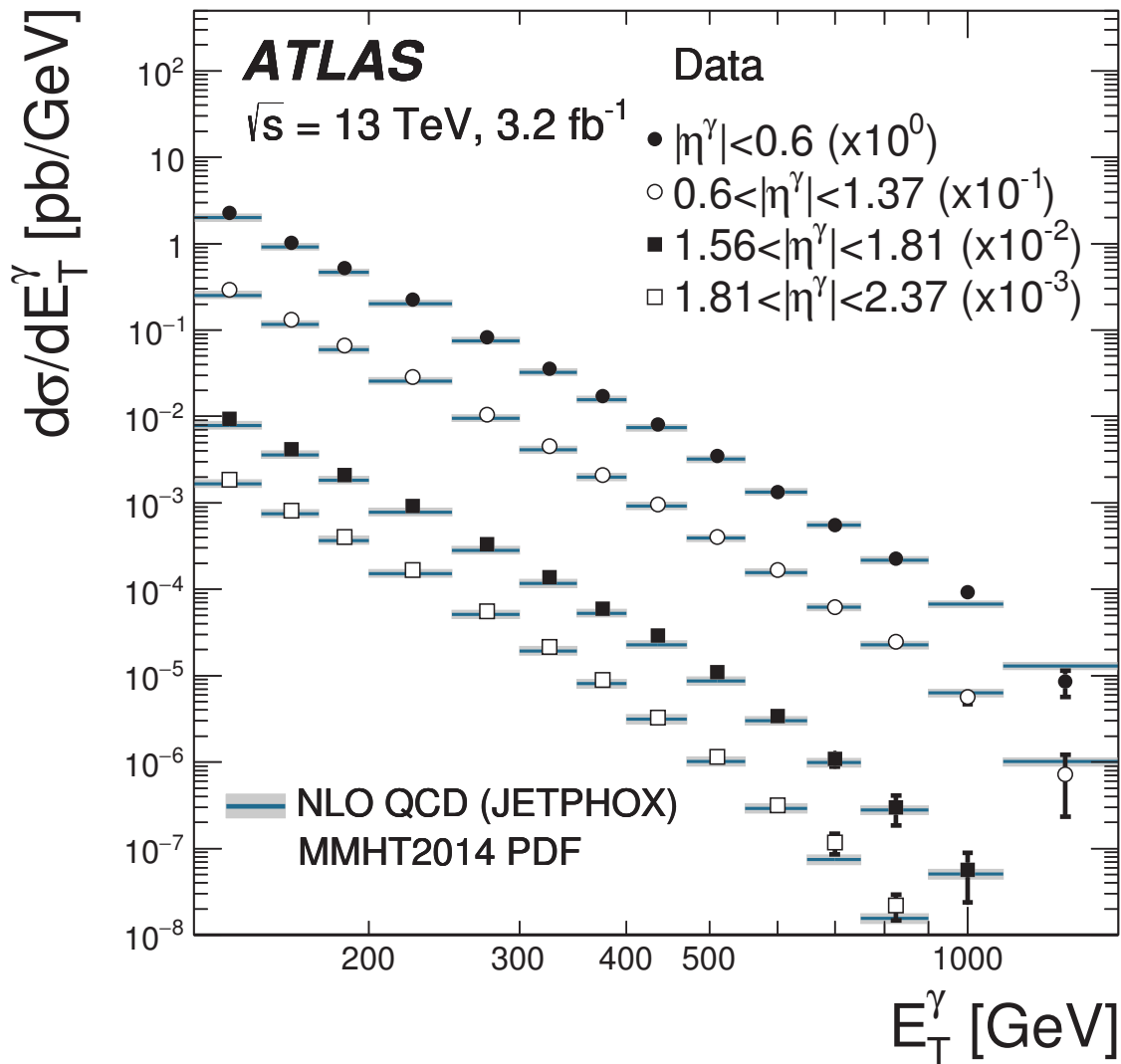
Photon isolation in ATLAS

- E_T^{iso} ($R = 0.4$) computed using clusters of calorimeter cells (EM and HAD) in a cone $R = 0.4$, excluding the contribution from the photon
 - Subtraction of the leakage of the photon energy into that cone (few %)
 - The underlying event and pileup (overlapping pp interactions in the same/neighbouring bunch crossings) contribute to E_T^{iso} !
- Subtracted on event-by-event basis using the jet-area method of M. Cacciari et al
- After isolation requirement, residual background still expected



Inclusive photon production at 13 TeV

Inclusive isolated-photon production in pp collisions at $\sqrt{s} = 13$ TeV



- Measurement of $d\sigma/dE_T^\gamma$ in different ranges in η^γ for $125 < E_T^\gamma < 1500$ GeV using $\mathcal{L} = 3.2 \text{ fb}^{-1}$ of pp collision data at $\sqrt{s} = 13$ TeV

- Isolation: $E_T^{\text{iso}} < 4.2 \cdot 10^{-3} \cdot E_T^\gamma + 4.8$ GeV

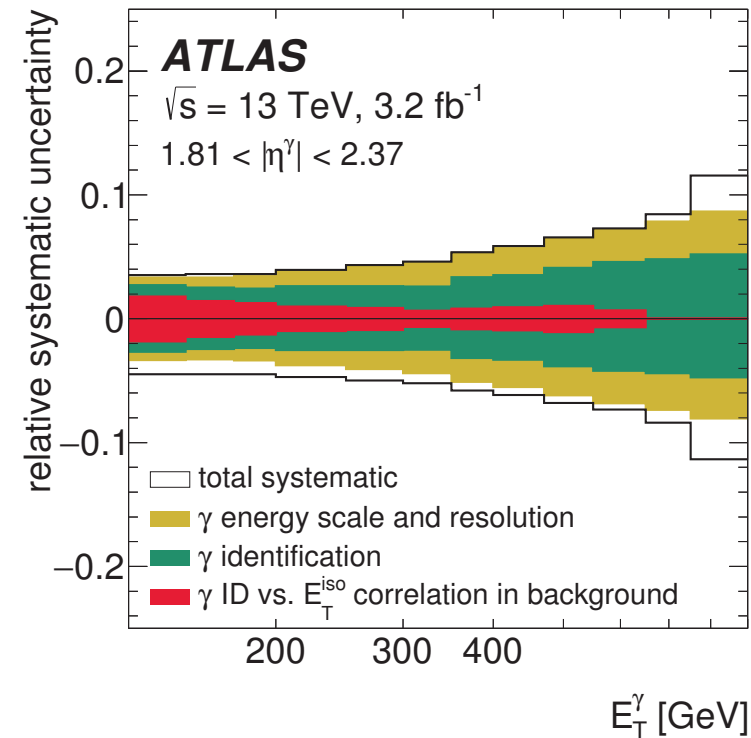
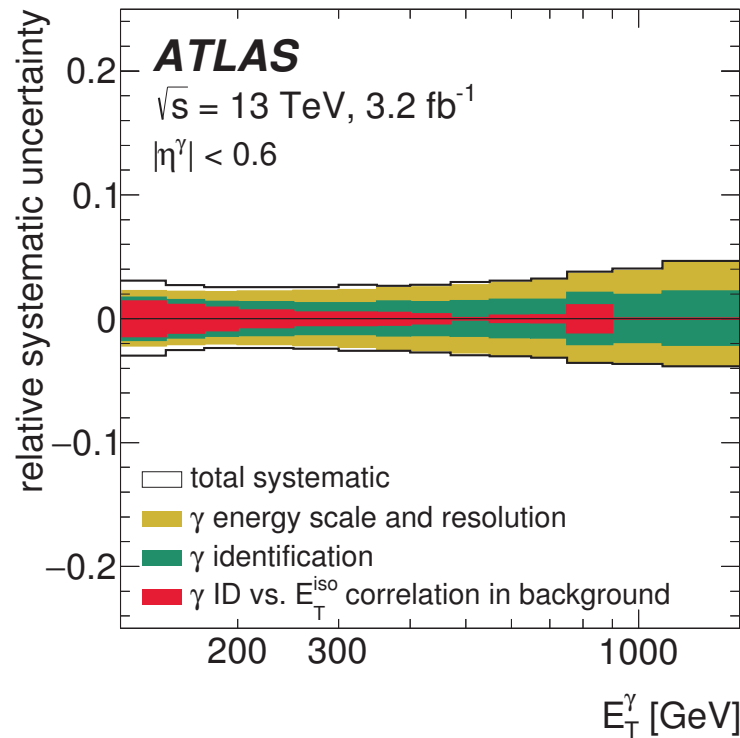
- The measurement covers more than five orders of magnitude in cross section

- $d\sigma/dE_T^\gamma$ increases by a factor 2 (10) at $E_T^\gamma = 125$ (1000) GeV with respect to at $\sqrt{s} = 8$ TeV

- Comparison to NLO QCD predictions computed with JetPhox using the MMHT2014 PDFs

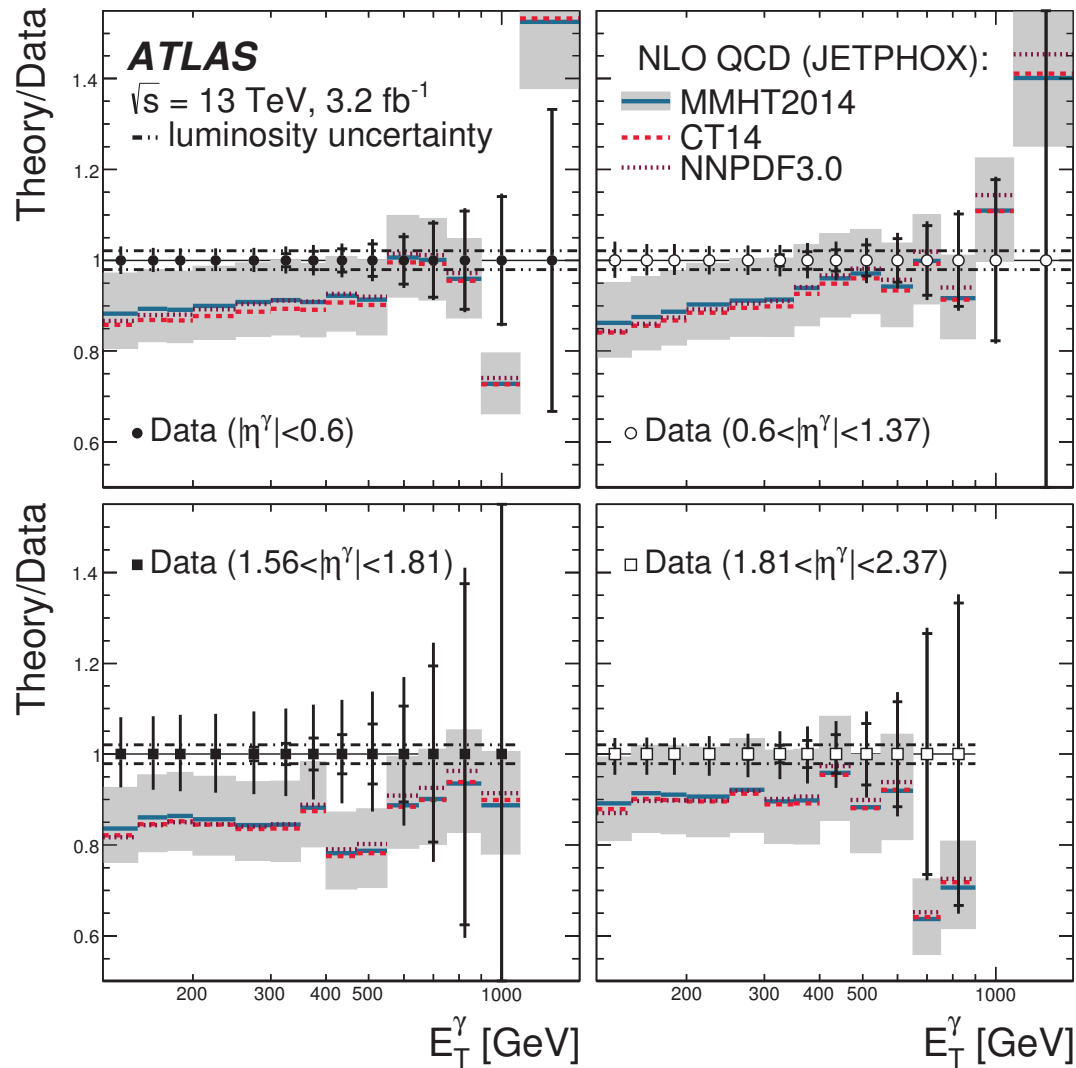
ATLAS Coll., arXiv:1701.06882, accepted PLB

Major experimental uncertainties



- The uncertainty on the photon energy scale dominates at high E_T^γ : 2–5% except for $1.56 < |\eta^\gamma| < 1.81$, where it is 7–18% (on the cross section)
- The uncertainty in the photon identification represents a significant contribution at low E_T^γ : it increases from 1–2% at 125 GeV to 2–6% at ~ 1 TeV (on the cross section)
- The uncertainty in the correlation between the photon ID variables and the isolation is a significant contribution at low E_T^γ : typically smaller than 2% (on the cross section)

Inclusive isolated-photon cross sections vs NLO QCD



- **NLO QCD predictions underestimate data by up to $\approx 10\text{-}15\%$**

- **Theoretical uncertainty 10-15% much larger than experimental uncertainties**

- **For $E_T^\gamma \lesssim 600 \text{ GeV}$ the measurements are systematically limited**

- **NLO QCD provides an adequate description of the data within uncertainties**

- **First measurement of inclusive photon production in the new kinematic regime opened by the LHC at $\sqrt{s} = 13 \text{ TeV}$**

- **Ready for the comparison to NNLO QCD predictions (Campbell, Ellis, Williams arXiv:1612.04333)**

Photon+jet(s) production at 8 TeV

Dynamics of $\gamma + \text{jet}$ production in pp collisions at $\sqrt{s} = 8 \text{ TeV}$

- Study of the $\gamma + \text{jet}$ dynamics by measuring the differential cross sections as functions of

→ Photon: E_T^γ

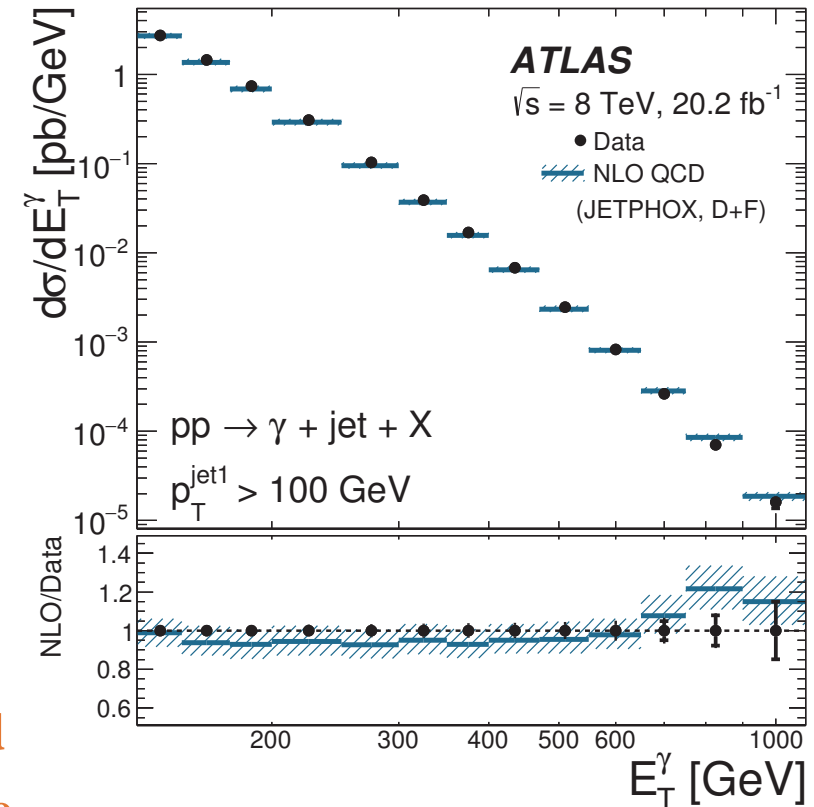
→ Leading jet: $p_T^{\text{jet}1}$

→ Photon+Leading jet: $m^{\gamma\text{-jet}1}, \cos \theta^*$

where $\cos \theta^* = \tanh \frac{1}{2} (y^{\text{jet}1} - \eta^\gamma)$

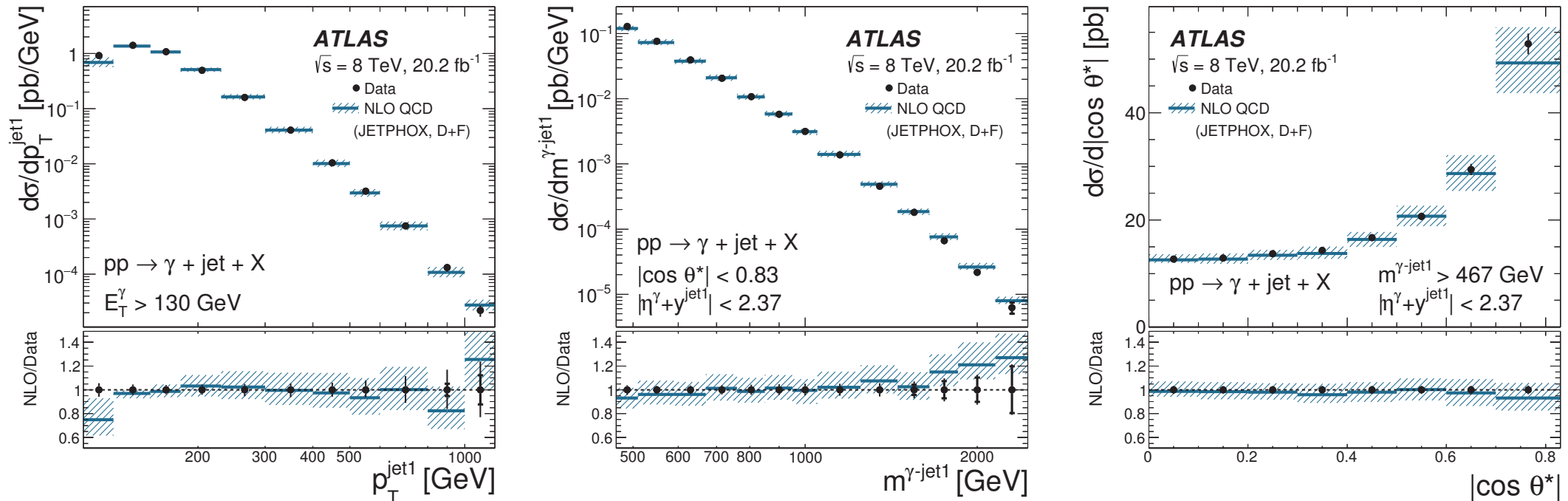
θ^* = scattering angle in centre-of-mass frame for $2 \rightarrow 2$ hard collinear scattering

- Measurements in the phase-space region defined by: $E_T^\gamma > 130 \text{ GeV}$, $|\eta^\gamma| < 2.37$ (excluding the region $1.37 < |\eta^\gamma| < 1.56$), $p_T^{\text{jet}1} > 100 \text{ GeV}$, $|y^{\text{jet}1}| < 4.4$ (anti- k_t algorithm with $R = 0.6$), $E_T^{\text{iso}} < 10 \text{ GeV}$ and $\Delta R_{\gamma j} > 1$
- Comparison to NLO QCD calculation (JETPHOX) corrected for non-perturbative effects
- Good description of the measured $d\sigma/dE_T^\gamma$ by the NLO QCD calculations
- Looking forward to comparison with NNLO QCD calculations (Campbell, Ellis, Williams arXiv:1703.10109)



Dynamics of $\gamma + \text{jet}$ production in pp collisions

ATLAS Coll., NPB918 (2017) 257

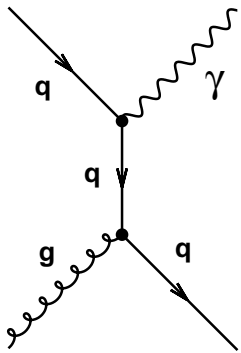


- **Additional requirements for $d\sigma/dm^{\gamma\text{-jet}1}$ and $d\sigma/d|\cos \theta^*|$ to remove biases due to the cuts on the p_T and rapidity of the leading photon and jet:**
 $|\eta^\gamma + y^{\text{jet}1}| < 2.37$, $|\cos \theta^*| < 0.83$, $m^{\gamma\text{-jet}1} > 467 \text{ GeV}$
- **In the selected (unbiased) region the angular distribution increases as $|\cos \theta^*|$ increases**
- **Good description of the data by the NLO QCD calculations within the (small) experimental and theoretical uncertainties \Rightarrow validation of the description of the dynamics of $\gamma + \text{jet}$ production in pp collisions at $\mathcal{O}(\alpha_{em}\alpha_s^2)$**

Dynamics of $\gamma + \text{jet}$ production in pp collisions

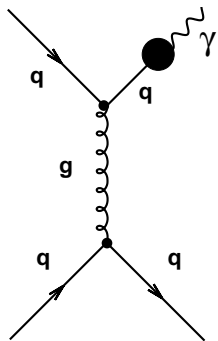
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- Angular distribution $d\sigma/d|\cos\theta^*|$ sensitive to the spin of the exchanged (virtual) particle: quark(1/2) vs gluon(1)



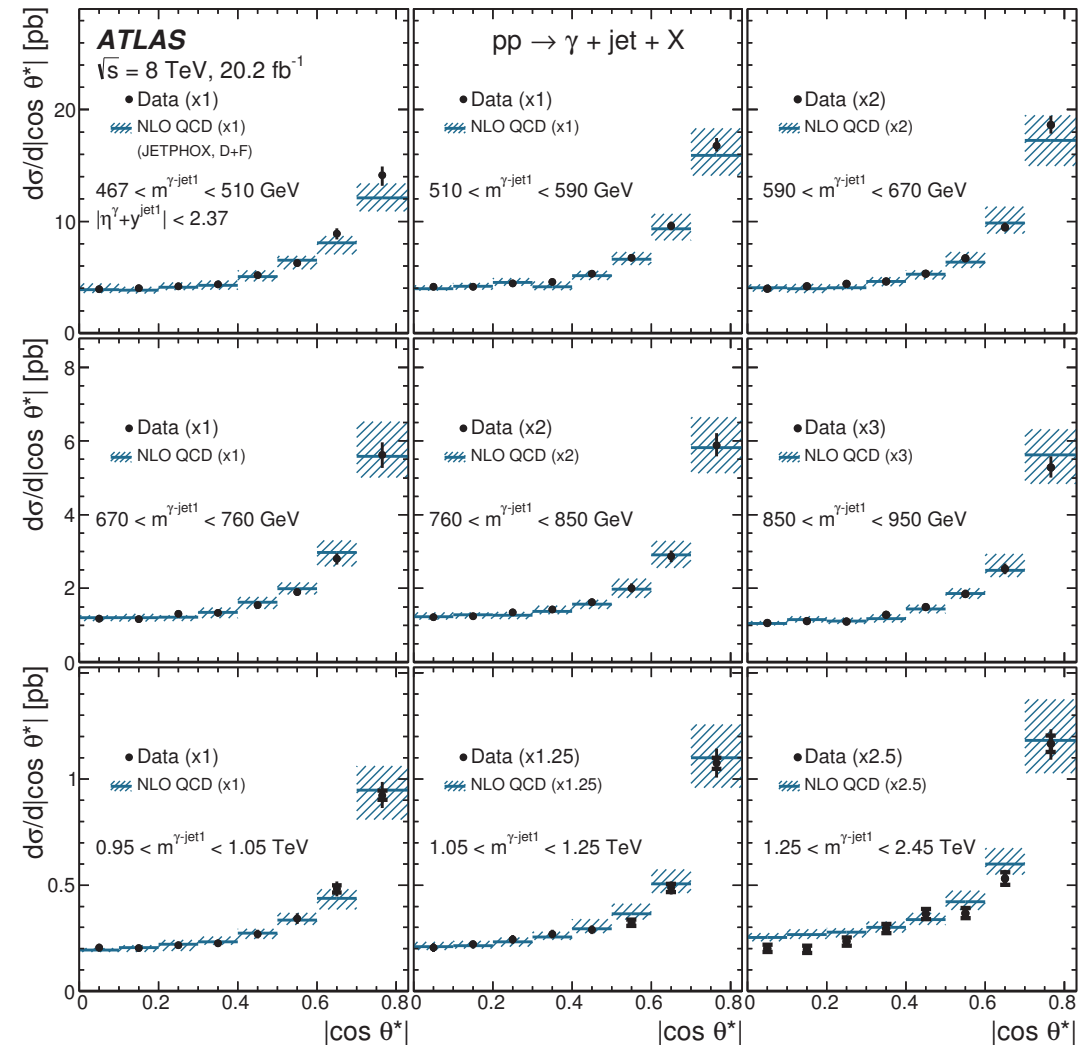
direct-photon process

$$d\sigma/d|\cos\theta^*| \sim (1 - |\cos\theta^*|)^{-1}$$



fragmentation process

$$d\sigma/d|\cos\theta^*| \sim (1 - |\cos\theta^*|)^{-2}$$

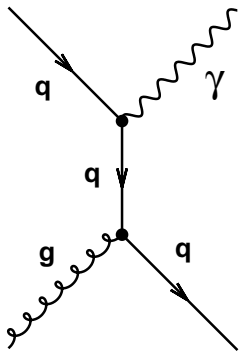


- Measured angular distribution in regions of photon-jet invariant mass \Rightarrow good description of the data by NLO QCD in shape and normalisation

Dynamics of $\gamma + \text{jet}$ production in pp collisions

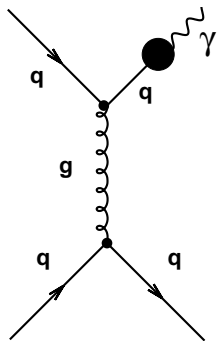
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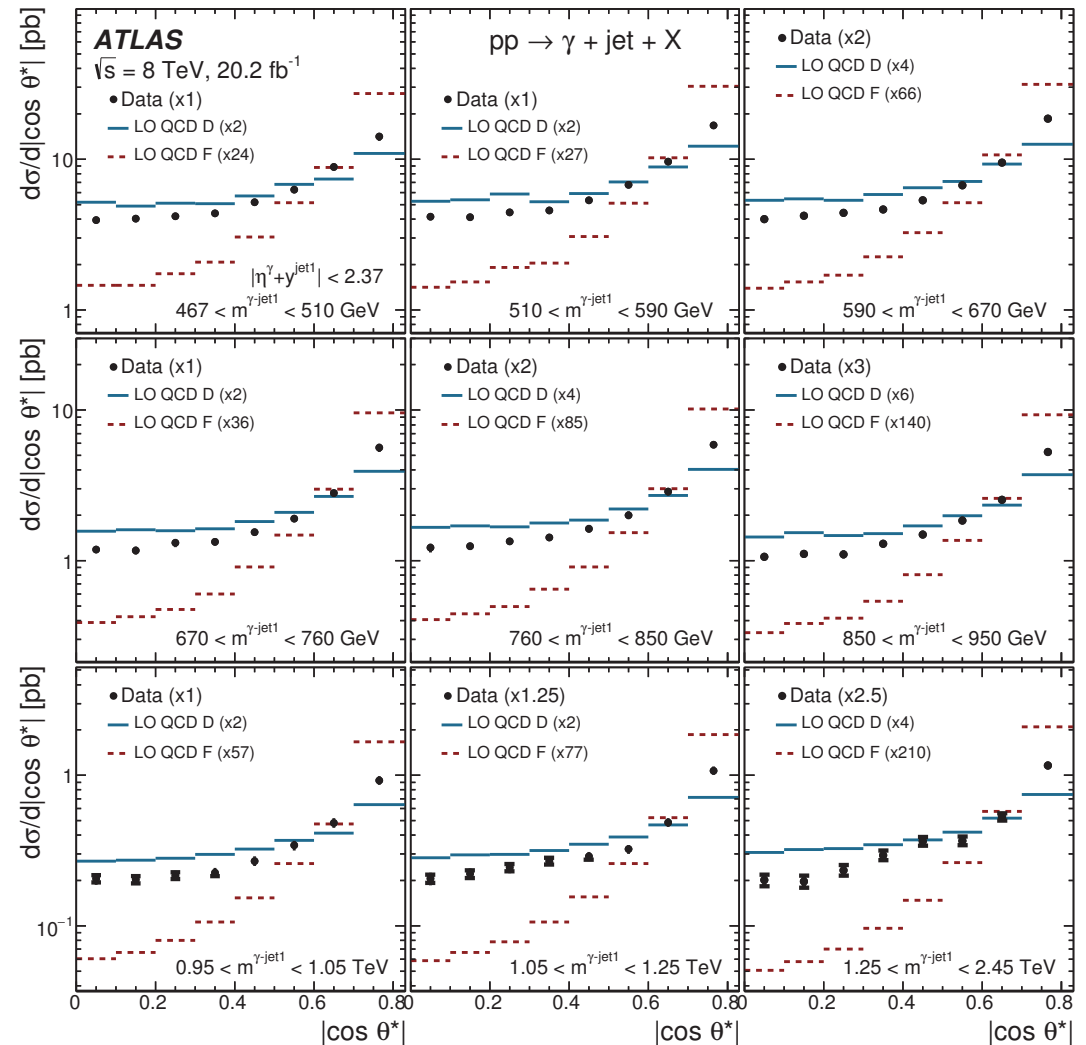
direct-photon process

$$d\sigma/d|\cos\theta^*| \sim (1 - |\cos\theta^*|)^{-1}$$



fragmentation process

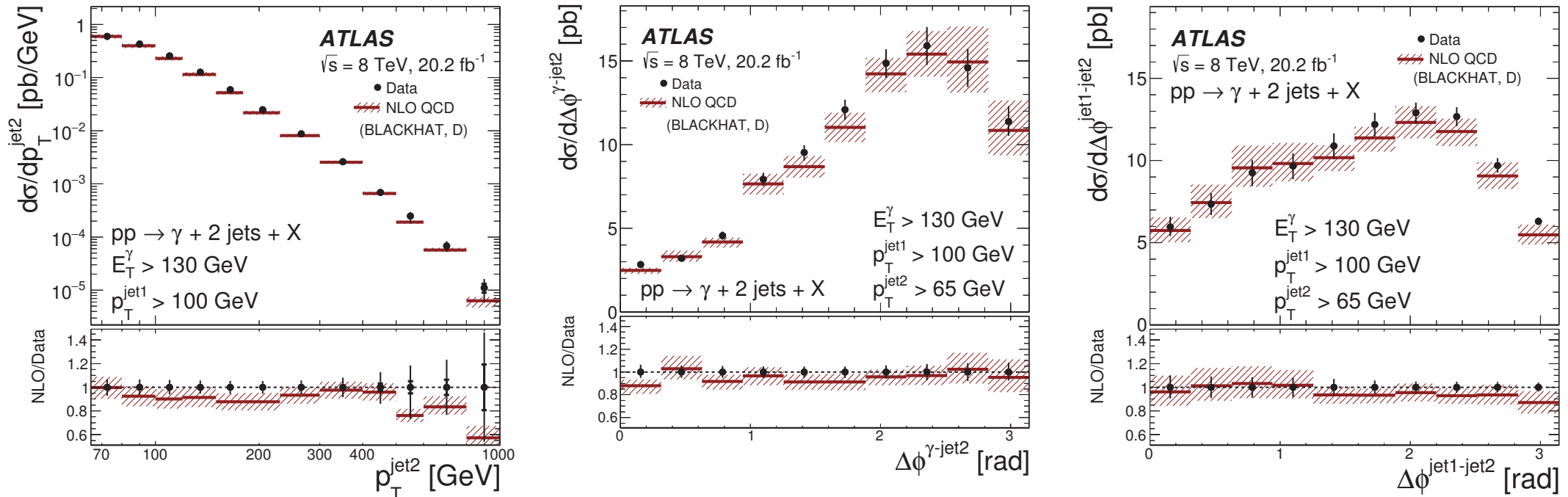
$$d\sigma/d|\cos\theta^*| \sim (1 - |\cos\theta^*|)^{-2}$$



- Measured angular distribution closer to that of direct-photon processes than fragm. \Rightarrow consistent with the dominance of processes in which a virtual quark is exchanged

Dynamics of $\gamma + 2\text{jet}$ production in pp collisions

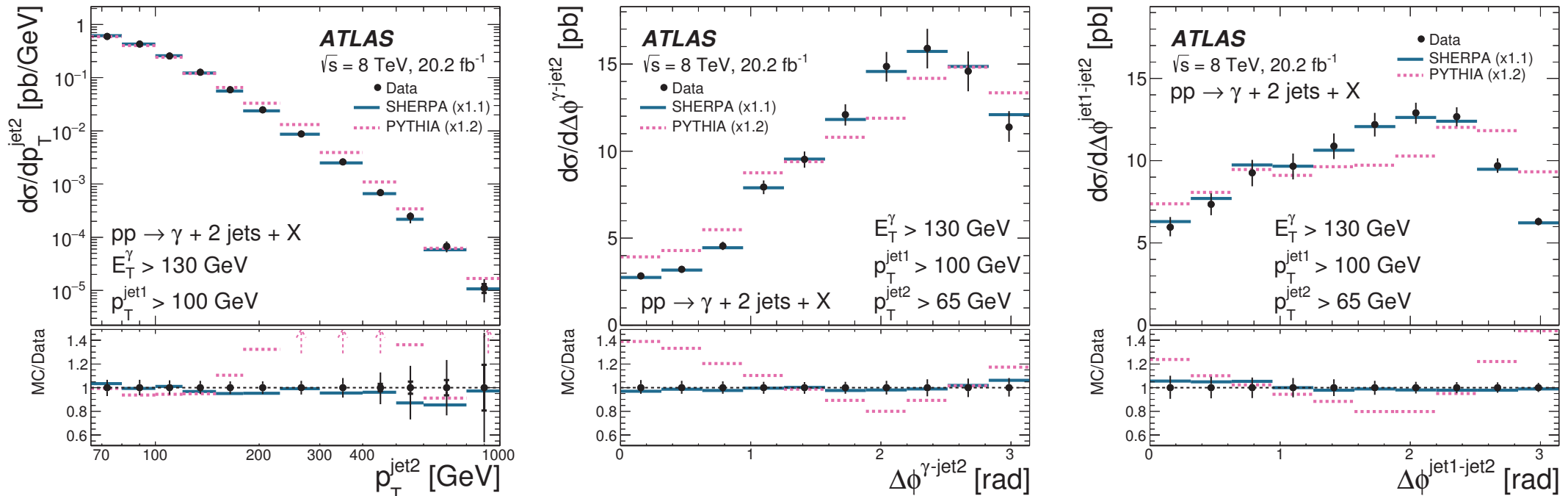
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- **First measurement of $\gamma + 2\text{jet}$ production in pp collisions at $\sqrt{s} = 8 \text{ TeV}$:**
 $E_T^\gamma > 130 \text{ GeV}, p_T^{\text{jet}1} > 100 \text{ GeV}$ and $p_T^{\text{jet}2} > 65 \text{ GeV}$
- **Measurement of $d\sigma/dp_T^{\text{jet}2}$ and angular correlations between the photon and the jets**
 → $\Delta\phi$ between the photon and subleading jet ($\Delta\phi^{\gamma\text{-jet}2}$)
 → $\Delta\phi$ between the leading and subleading jets ($\Delta\phi^{\text{jet}1\text{-jet}2}$)
- **Good description of the data both in shape and normalisation by the NLO QCD predictions computed with Blackhat**

Dynamics of $\gamma + 2\text{jet}$ production in pp collisions

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- **Comparison to the predictions of Monte Carlo generators:**

- **PYTHIA: 2** → 2 matrix elements plus parton showers

- **SHERPA: 2** → n ($n = 2, \dots, 5$) matrix elements plus parton showers

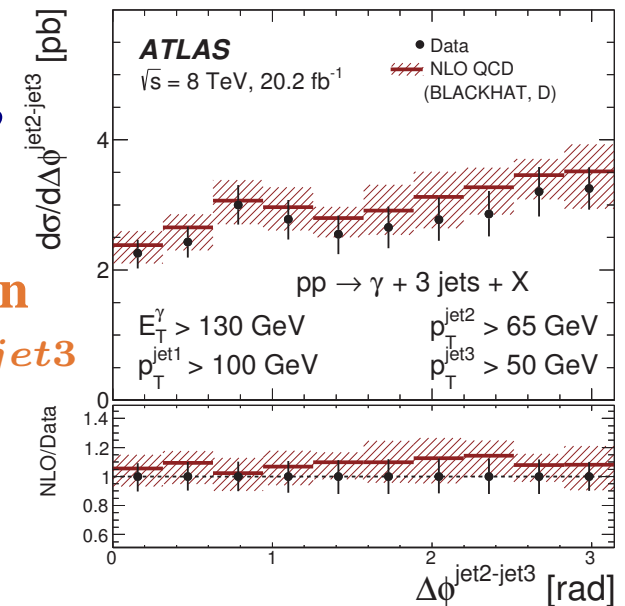
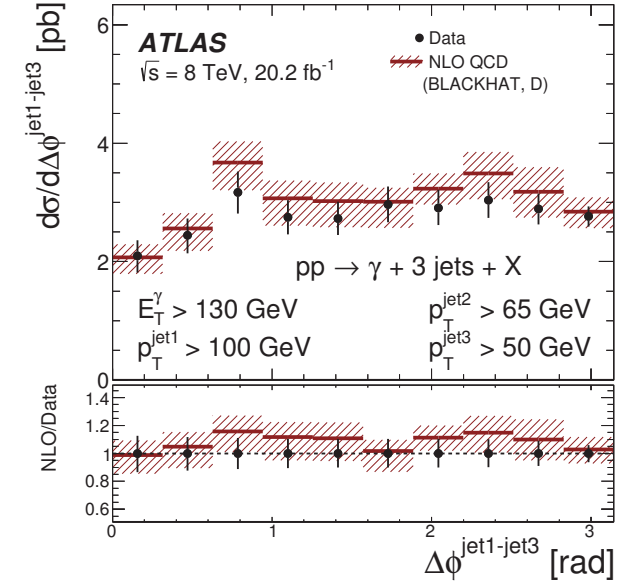
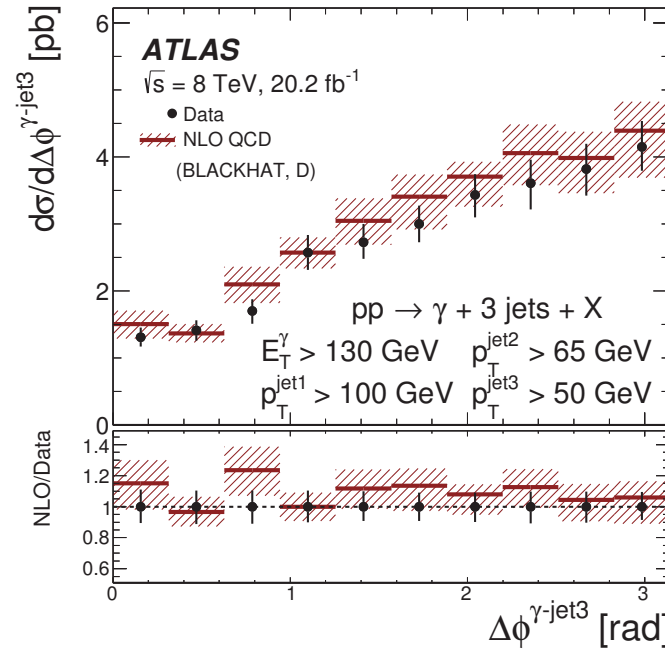
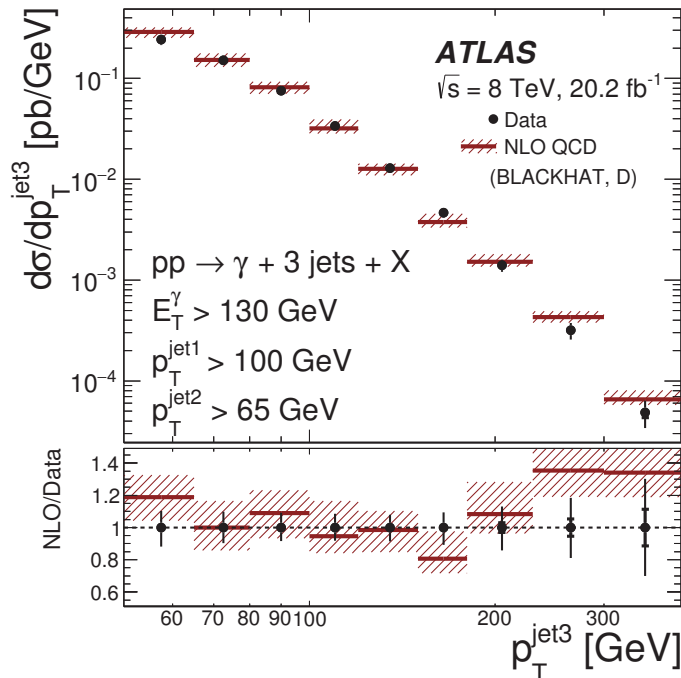
- **MC predictions normalised to data: shape comparison only**

- **Good description of the data by the SHERPA predictions while PYTHIA fails to describe the distribution in p_T^{jet2} and the angular correlations**

⇒ **Inclusion of higher-order tree-level ME in SHERPA improves description of data significantly**

Dynamics of $\gamma + 3\text{jet}$ production in pp collisions

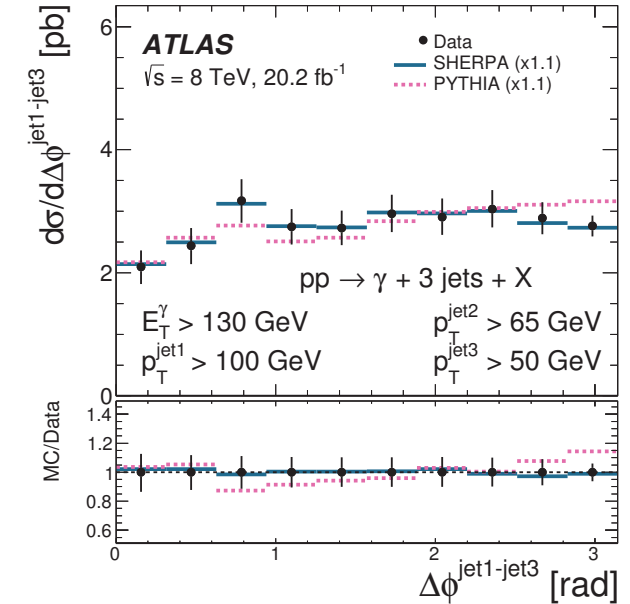
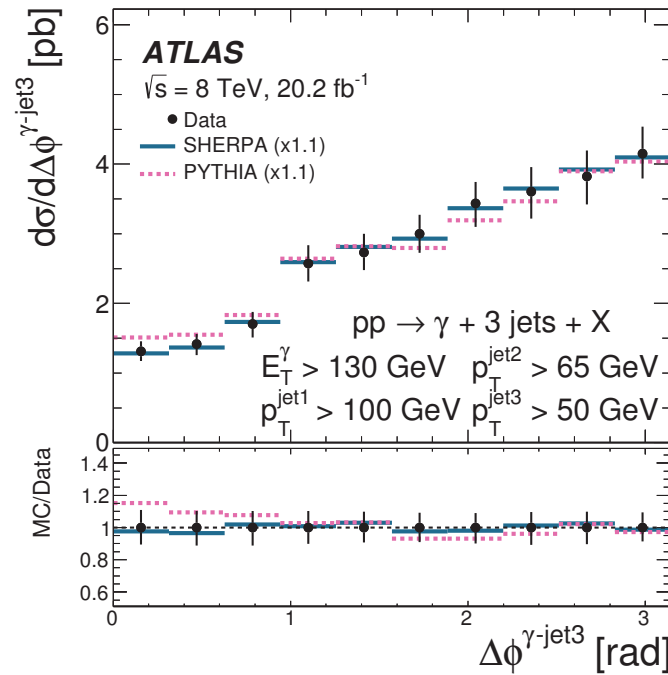
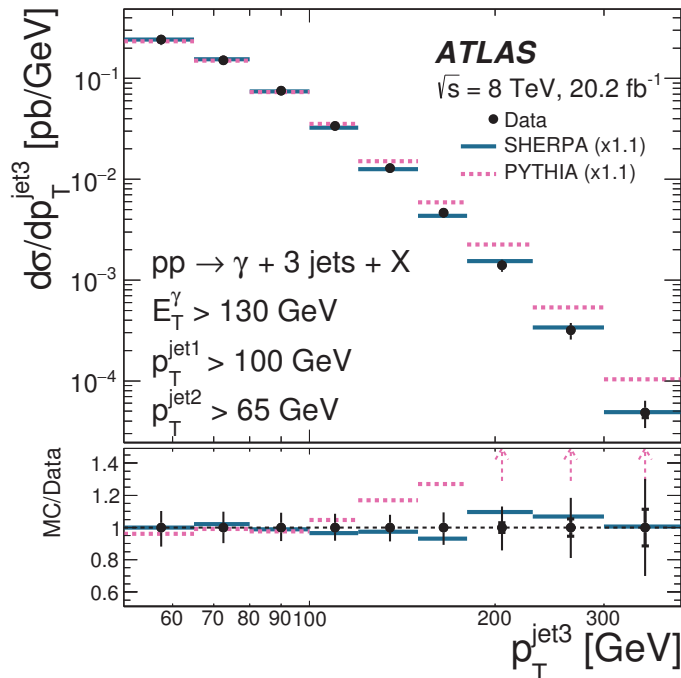
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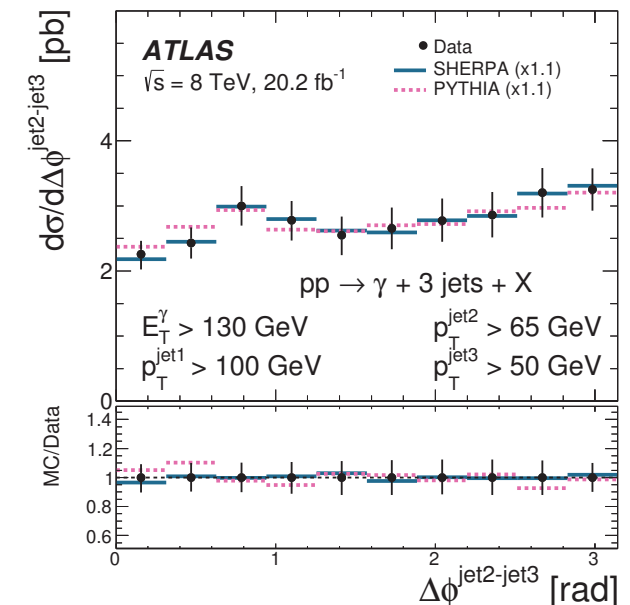
- First measurement of $\gamma + 3\text{jet}$ at the LHC: $E_T^\gamma > 130 \text{ GeV}$, $p_T^{\text{jet}1} > 100 \text{ GeV}$, $p_T^{\text{jet}2} > 65 \text{ GeV}$ and $p_T^{\text{jet}3} > 50 \text{ GeV}$
- Measurement of $d\sigma/dp_T^{\text{jet}3}$ and angular correlations between the photon and the jets: $\Delta\phi^{\gamma\text{-jet}3}$, $\Delta\phi^{\text{jet}1\text{-jet}3}$, $\Delta\phi^{\text{jet}2\text{-jet}3}$
- Adequate description of the data by the NLO QCD predictions computed with Blackhat

Dynamics of $\gamma + 3\text{jet}$ production in pp collisions

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- Comparison to the predictions of Monte Carlo generators of PYTHIA (2 → 2 ME+PS) and SHERPA (2 → n ME+PS) normalised to data (shape comparison)
 - Good description of the data by the SHERPA predictions while PYTHIA describes poorly the distribution in p_T^{jet3} at large values
- ⇒ Inclusion of higher-order tree-level ME in SHERPA improves description of data significantly



Photon pair production at 8 TeV

Isolated-photon pair production in pp collisions

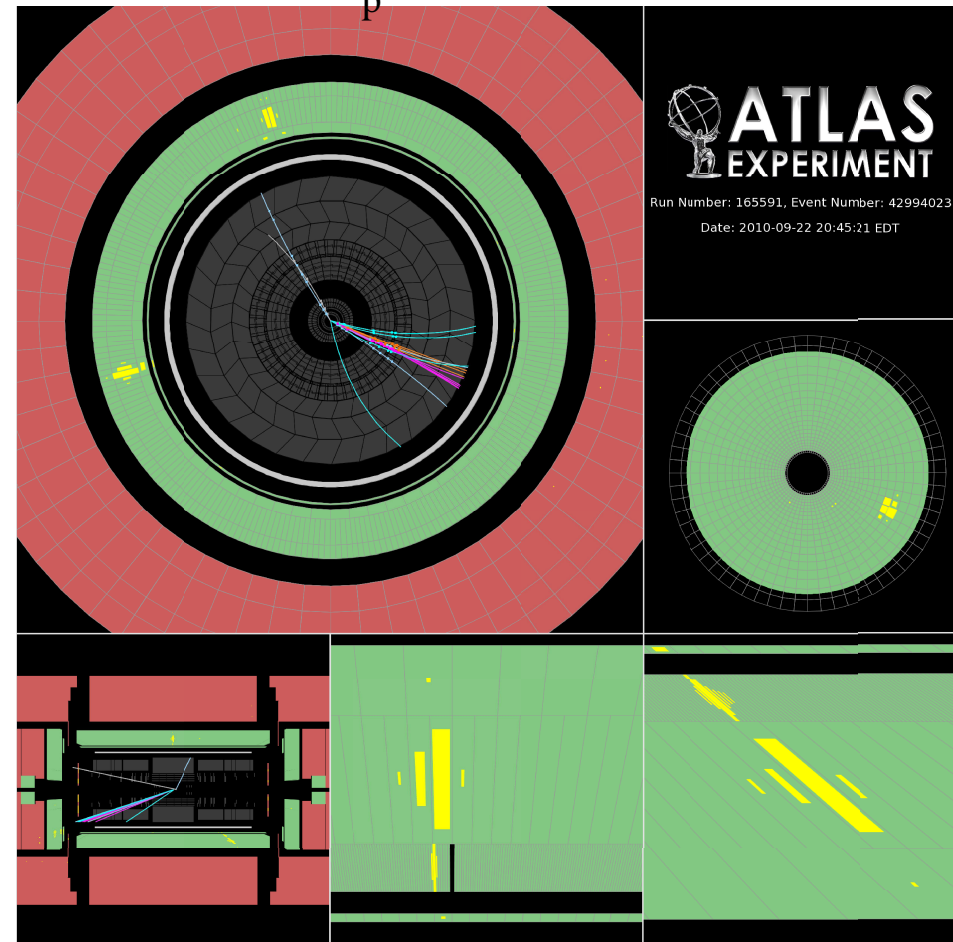
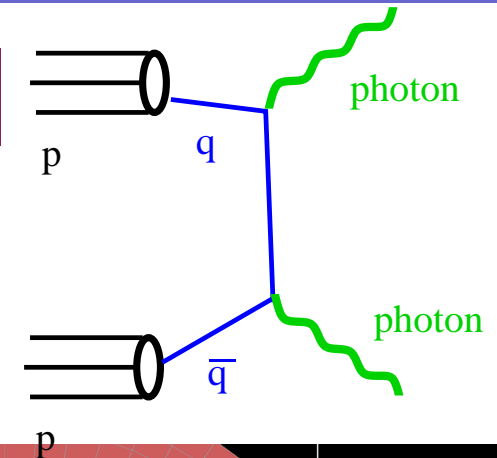
- Measurements of the process $pp \rightarrow \gamma\gamma + X$ with the aim of testing pQCD and understanding the irreducible background to new physics processes involving photons or $H \rightarrow \gamma\gamma$

- Measurement of differential cross sections

- diphoton invariant mass, $m_{\gamma\gamma}$
- diphoton transverse momentum, $p_{T,\gamma\gamma}$
- azimuthal separation in LAB frame, $\Delta\phi_{\gamma\gamma}$
- $\cos\theta_\eta^*$ → $\phi_\eta^* \equiv \tan\left(\frac{\pi - \Delta\phi_{\gamma\gamma}}{2}\right) \sin\theta_\eta^*$
- transverse component of $\vec{p}_{T,\gamma\gamma}$ with respect to thrust axis (a_T)

in the phase-space region defined by:

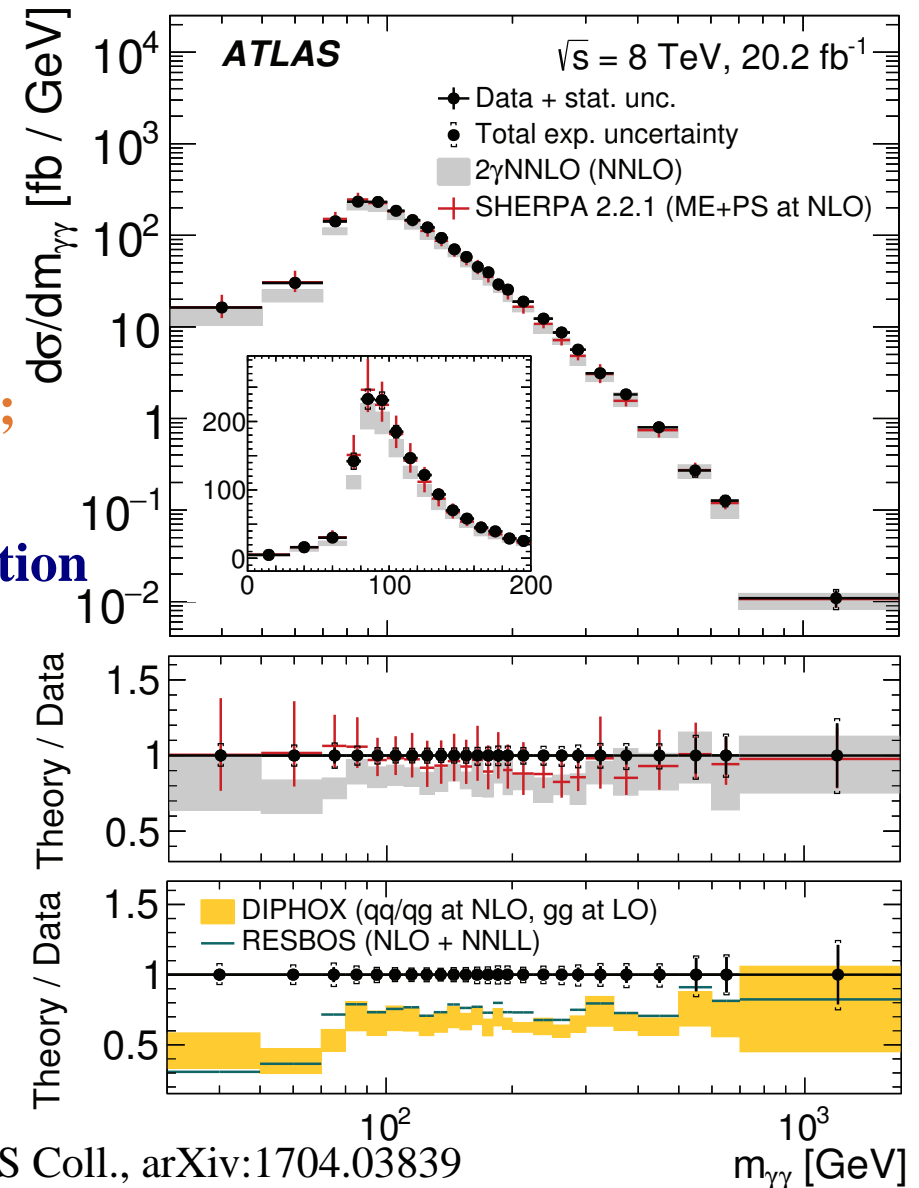
$E_T^{\gamma 1,2} > 40(30) \text{ GeV}$, $|\eta^\gamma| < 2.37$ (excluding the region $1.37 < |\eta^\gamma| < 1.56$), $\Delta R_{\gamma\gamma} > 0.4$ and $E_T^{iso} < 11 \text{ GeV}$



Isolated-photon pair production in pp collisions at $\sqrt{s} = 8$ TeV

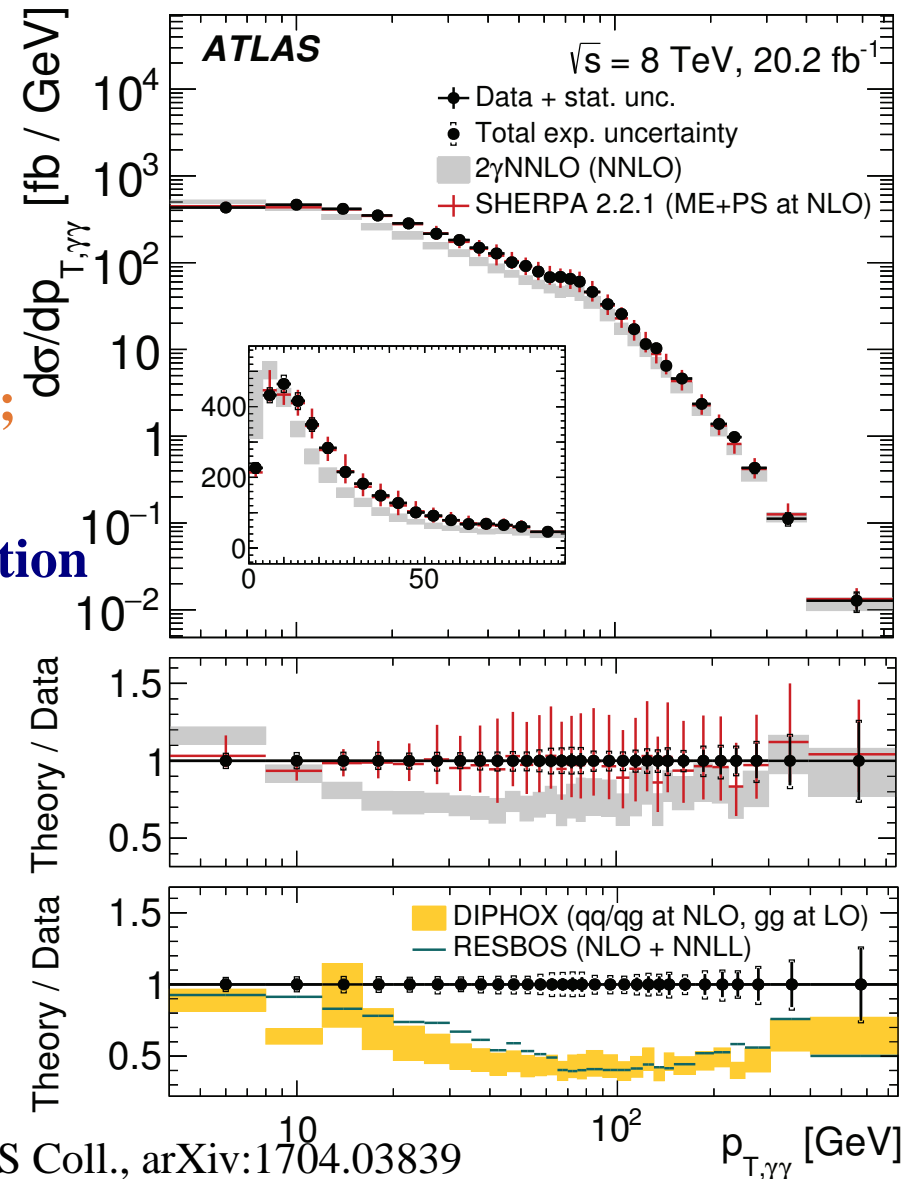
- Comparison to theoretical calculations
- Fixed-order QCD calculations (NP corrected)
 - 2γ NNLO program; NNLO calculation of direct-photon contribution (no fragm.)
 - DIPHOX program; NLO calculation of direct-photon and fragmentation contributions; box diagram $gg \rightarrow \gamma\gamma$ included
 - RESBOS program; NLO plus NNLL resummation
- New SHERPA (v2.2.1) calculation combining
 - $\gamma\gamma$ and $\gamma\gamma + 1p$ at NLO
 - $\gamma\gamma + 2p$ and $\gamma\gamma + 3p$ at LO
 - parton showers
- The small contribution from $H \rightarrow \gamma\gamma$ is neglected

⇒ SHERPA prediction in agreement with data

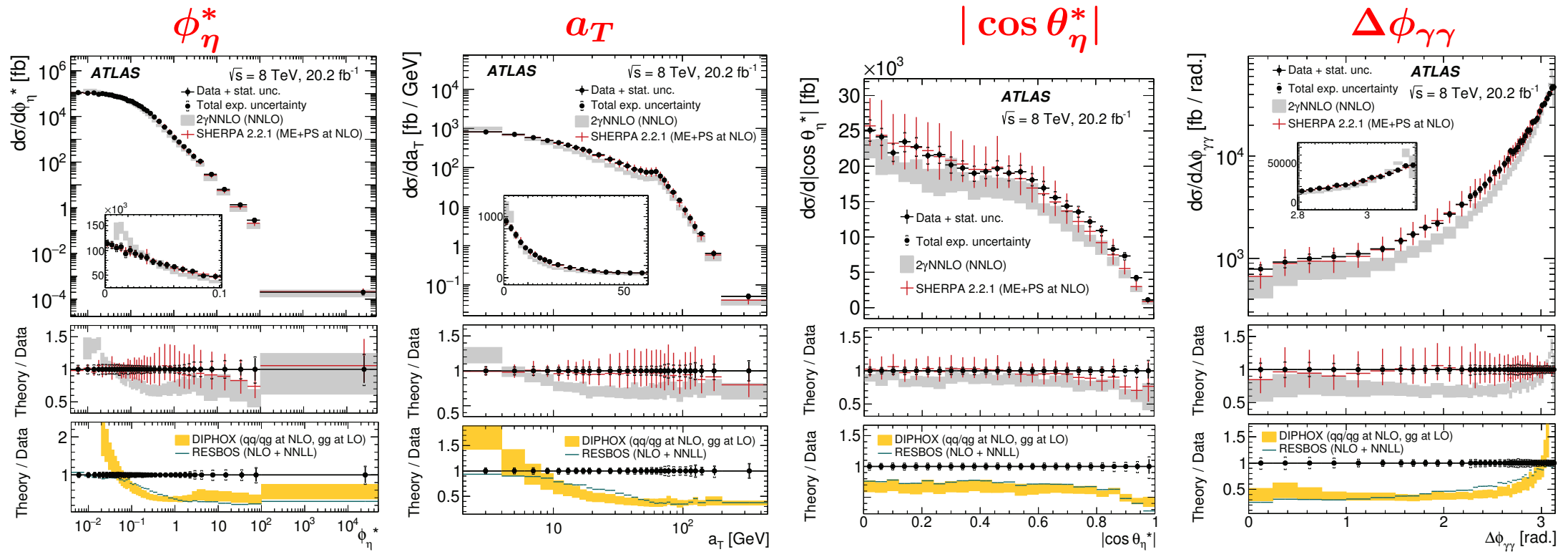


Isolated-photon pair production in pp collisions at $\sqrt{s} = 8$ TeV

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 - Fixed-order QCD calculations (NP corrected)
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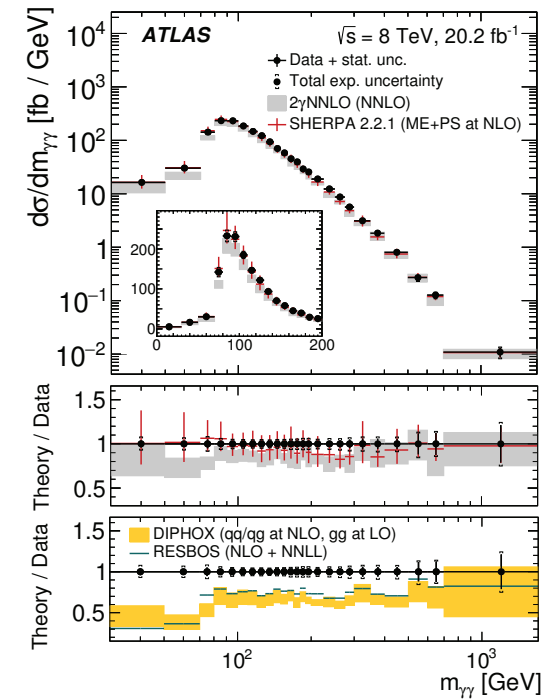
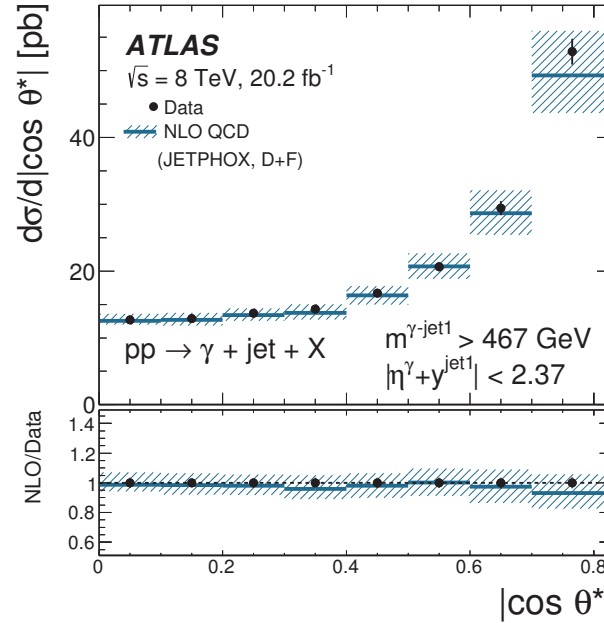
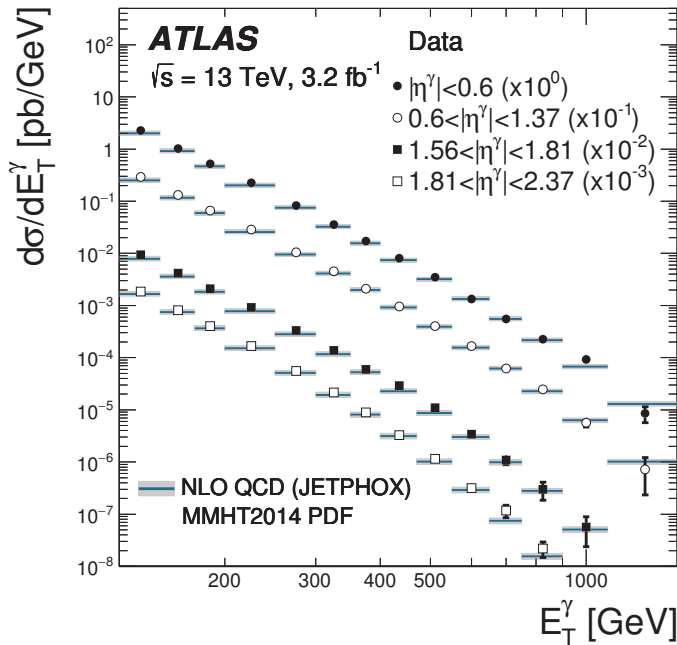
Isolated-photon pair production in pp collisions at $\sqrt{s} = 8$ TeV



ATLAS Coll., arXiv:1704.03839

- $\Delta\phi_{\gamma\gamma} \sim \pi$ or at low values of $p_{T,\gamma\gamma}$, a_T and ϕ_η^* (soft gluon resummation important): **RESBOS** and **SHERPA** do well
- NLO QCD calculations without higher order terms (DIPHOX, RESBOS) are insufficient
- NNLO corrections (2 γ NNLO) improve the description, but still insufficient
- **SHERPA** predictions agree with the data

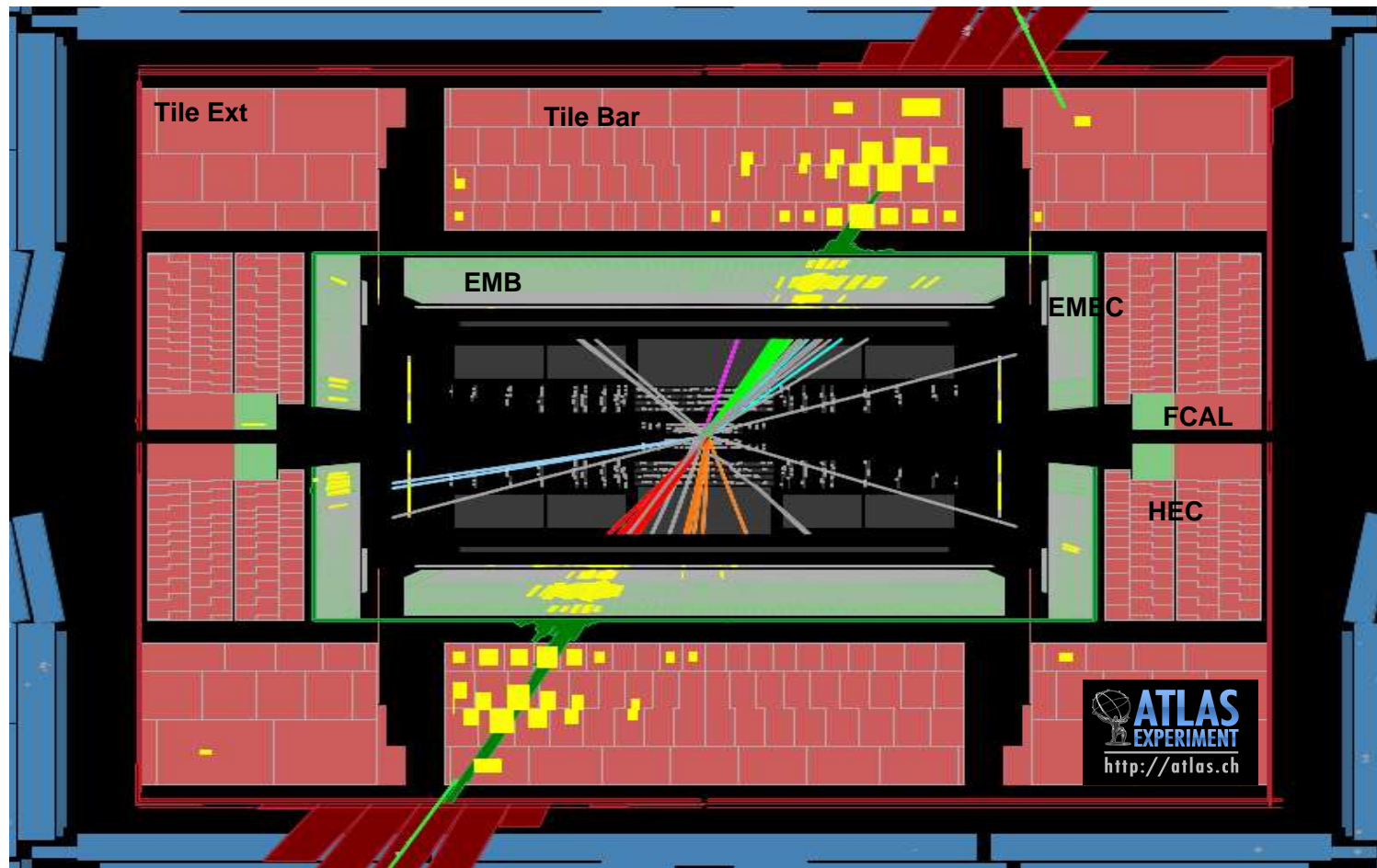
Summary



- Exploration of isolated photon production in pp collisions up to $E_T^\gamma \sim 1 \text{ TeV}$
- Additional experimental information on the gluon density in the proton
- Measurement of the dynamics of photon+jet(s) and diphoton production
- Understanding (in pQCD) the background to Higgs into $\gamma\gamma$
- Overall, perturbative QCD succeeds in describing the data!

Backup

The ATLAS detector



- Inner detector (ID): tracking and particle identification in $|\eta| < 2.5$
- Calorimeters: electromagnetic (LAr) \rightarrow barrel $|\eta| < 1.475$, endcap $1.375 < |\eta| < 3.2$, forward $3.1 < |\eta| < 4.9$; hadronic (scintillator/steel, LAr/Cu, LAr/W) \rightarrow barrel $|\eta| < 0.7$ extended barrel $0.8 < |\eta| < 1.7$, endcap $1.5 < |\eta| < 3.2$ and forward $3.1 < |\eta| < 4.9$

Photon reconstruction in the ATLAS LAr Calorimeter

● Layout of the ATLAS electromagnetic calorimeter (Lead-liquid Argon)

- barrel section, $|\eta| < 1.475$
- two end-cap sections, $1.375 < |\eta| < 3.2$
- three longitudinal layers

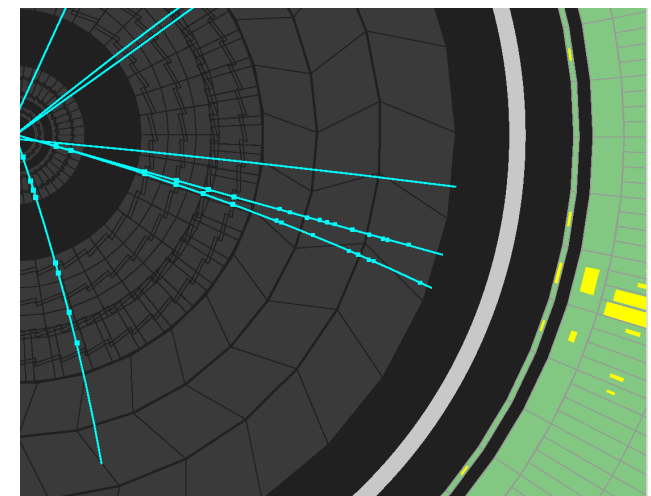
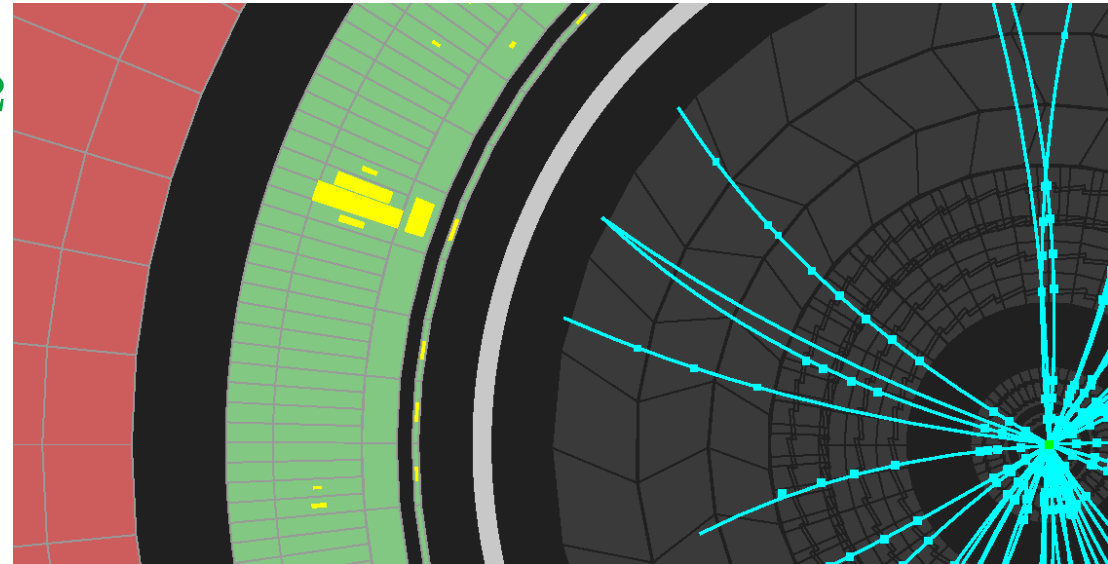
- First layer: high granularity in η direction, width 0.003-0.006 (except for $1.4 < |\eta| < 1.5$ and $|\eta| > 2.4$)
- Second layer: collects most of the energy, granularity 0.025×0.025 in $\eta \times \phi$
- Third layer: used to correct for leakage

● Cluster of EM cells without matching track:

- “unconverted” photon candidate

● Cluster of EM cells matched to pairs of tracks (from reconstructed conversion vertices in the inner detector) or matched to a single track consistent with originating from a photon conversion

- “converted” photon candidate



Photon identification in the ATLAS LAr Calorimeter

● **To discriminate signal vs background:** shape variables from the lateral and longitudinal energy profiles of the shower in the calorimeters; “loose” and “tight” identification criteria.

● **“Loose” identification criteria:**

→ leakage $R_{had} = E_T^{had} / E_T$ (1st layer hadronic calorimeter)

→ $R_\eta = E_{3 \times 7}^{S2} / E_{7 \times 7}^{S2}$; S2=second layer of EM calorimeter

→ RMS width of the shower in η direction in S2

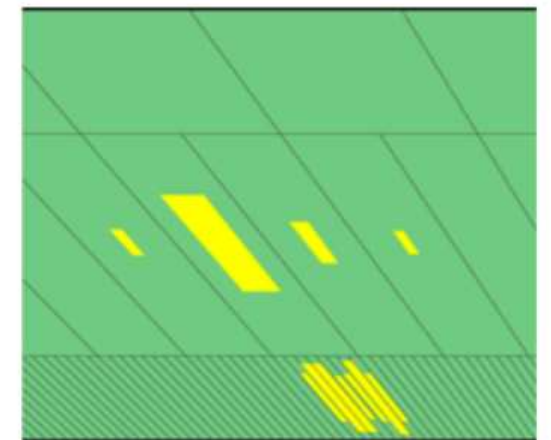
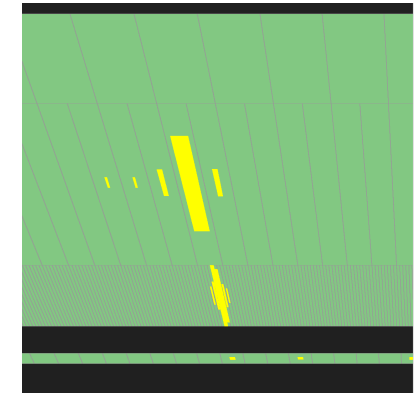
● **“Tight” identification criteria:**

→ the requirements applied in “Loose” are tightened

→ $R_\phi = E_{3 \times 3}^{S2} / E_{3 \times 7}^{S2}$

and shower shapes in the first layer (to discriminate single-photon showers from overlapping nearby showers, such as $\pi^0 \rightarrow \gamma\gamma$)

→ e.g. asymmetry between the 1st and 2nd maxima in the energy profile along η (S1)



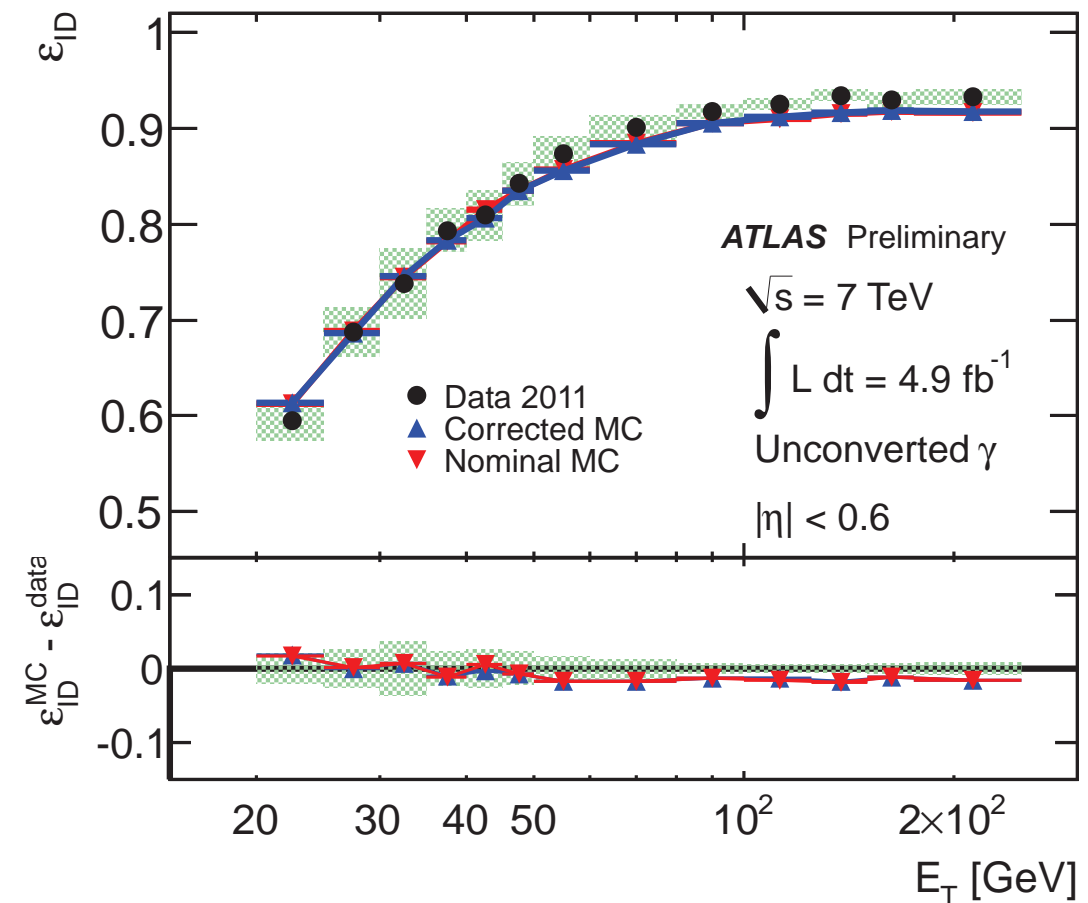
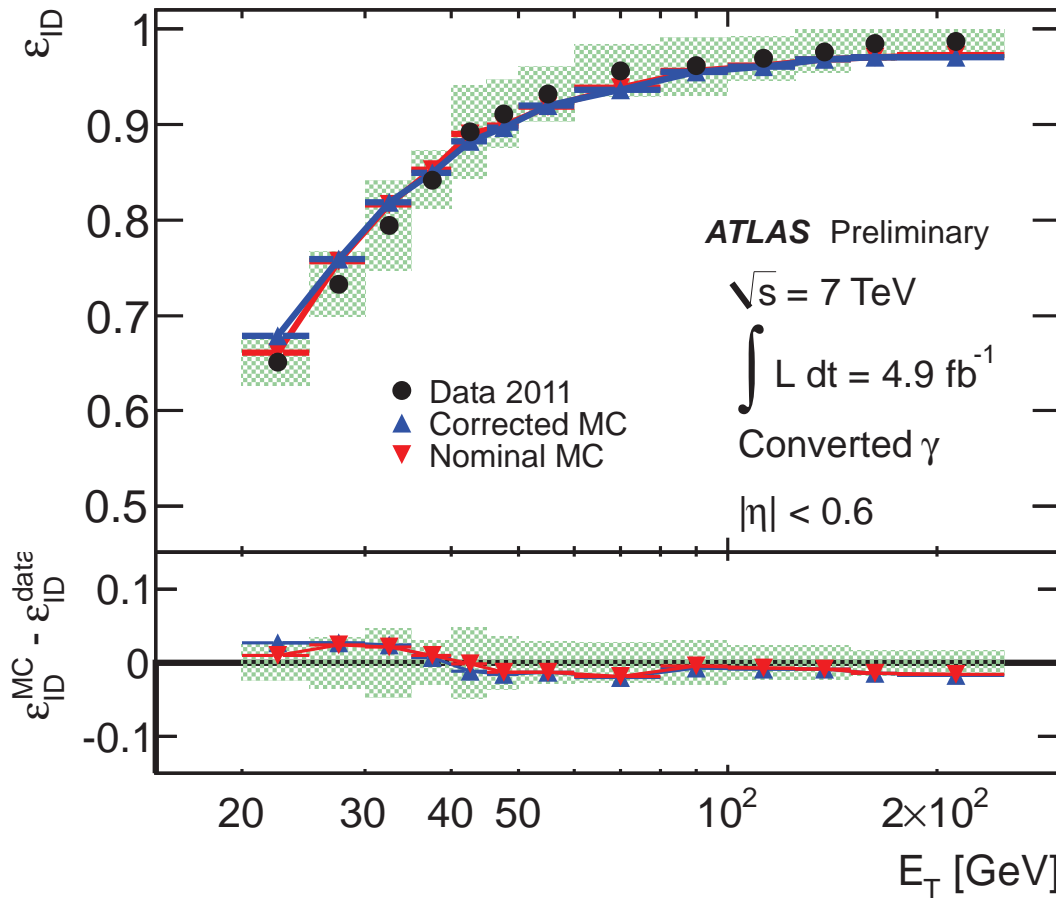
$\pi^0 \rightarrow \gamma\gamma$

Photon identification efficiency

ATLAS Coll., ATLAS-CONF-2012-123

CONVERTED PHOTONS

UNCONVERTED PHOTONS



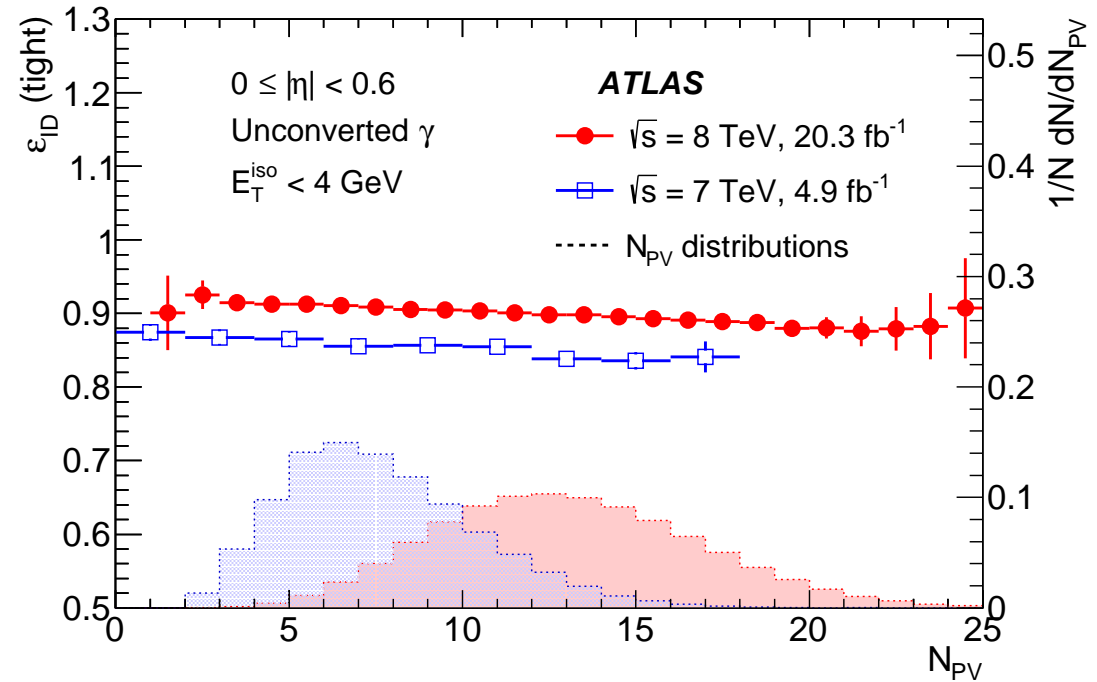
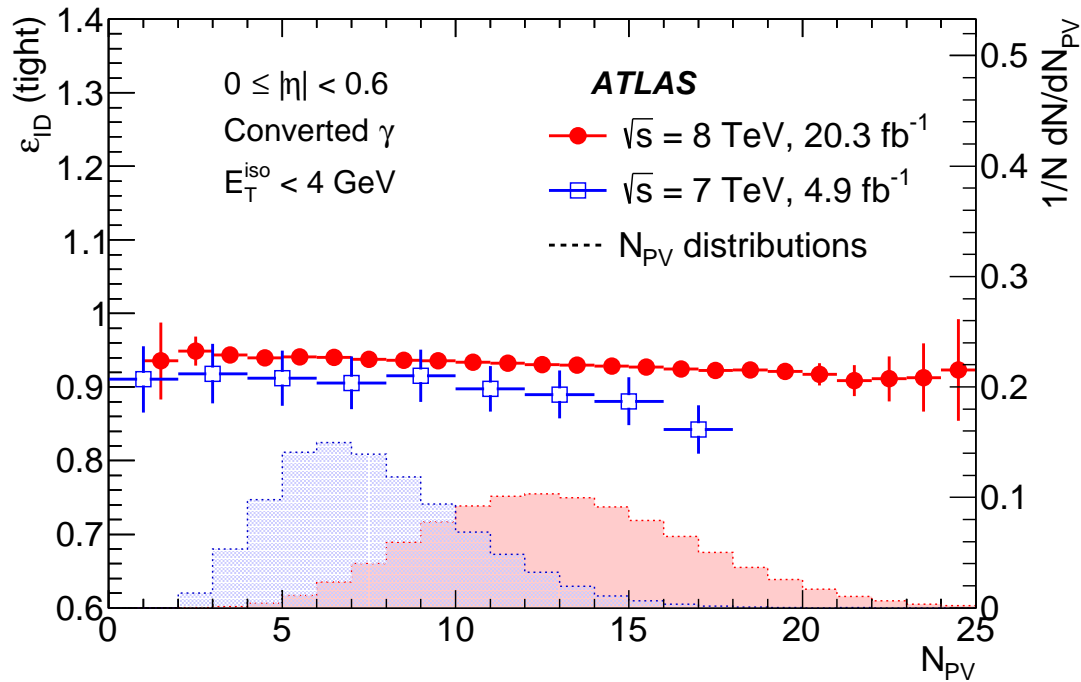
- Data-driven measurements of photon identification efficiency for converted and unconverted photons (radiative Z decays, extrapolation from e^\pm and matrix method) compared to estimations based on Monte Carlo simulations

Photon identification efficiency vs pile-up

ATLAS Coll., EPJC 76 (2016) 666

CONVERTED PHOTONS

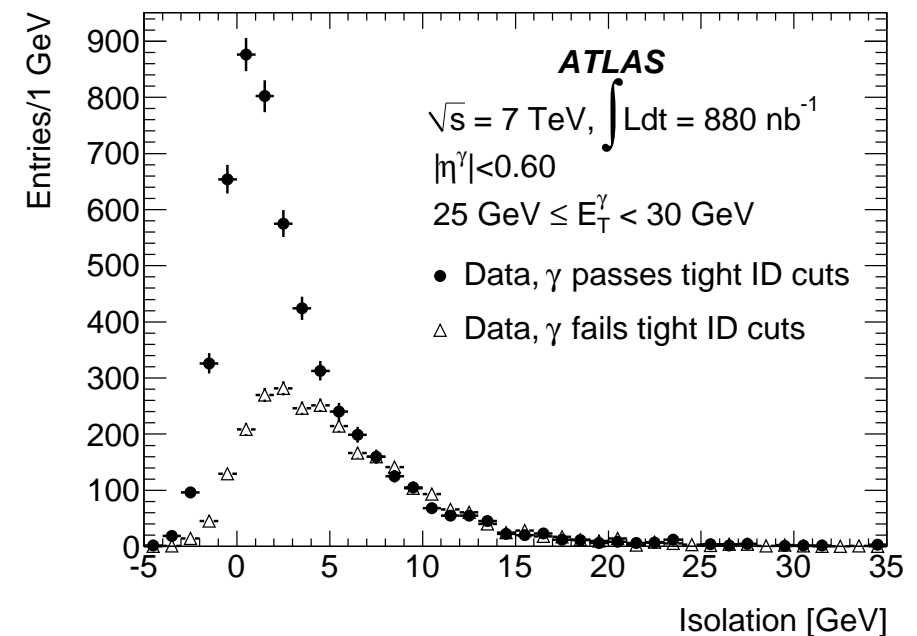
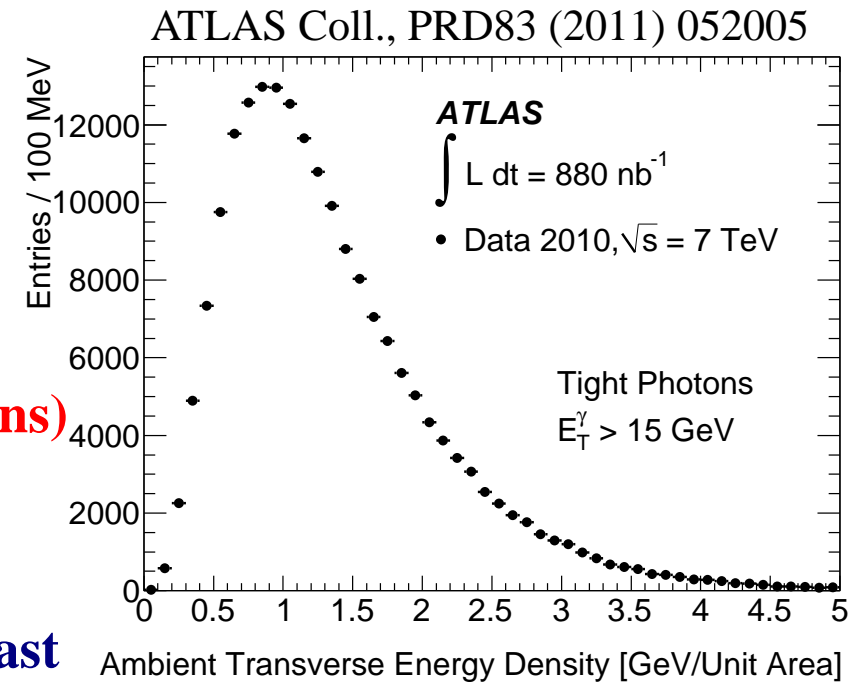
UNCONVERTED PHOTONS



- Comparison of data-driven efficiency measurements for converted and unconverted photons performed with the 2011 and 2012 datasets as a function of the number of reconstructed primary vertex candidates (N_{PV}), for $|\eta^\gamma| < 0.6$. The 2011 measurements are performed with the matrix method for photons with $E_T^\gamma > 20 \text{ GeV}$ and the 2012 measurements with the electron extrapolation method for photons with $E_T^\gamma > 30 \text{ GeV}$.**

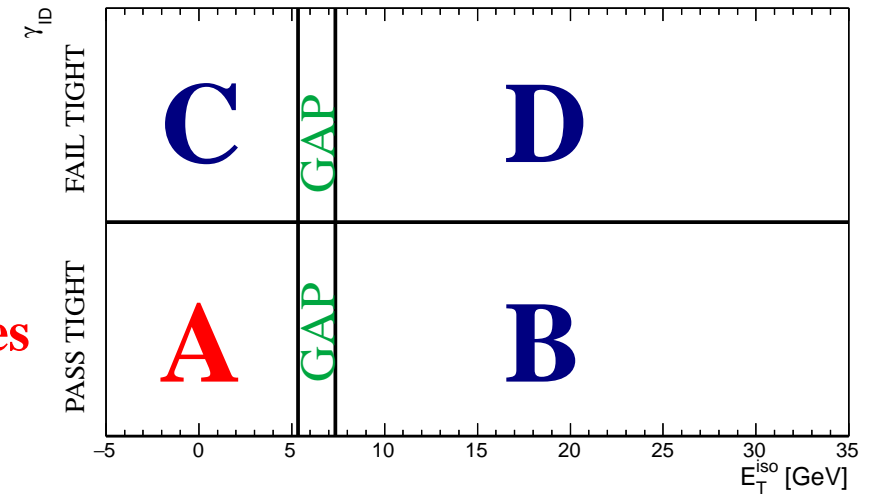
Photon isolation in ATLAS

- E_T^{iso} is corrected by subtracting the estimated contributions from the underlying event and pileup; the correction is computed on an event-by-event basis (to avoid the large fluctuations) using the jet-area method (M. Cacciari et al.)
 \Rightarrow ambient transverse-energy density
 540 MeV (in $R = 0.4$ cone) for events with at least one photon candidate with $E_T > 15$ GeV and exactly one PV (+170 MeV for each extra PV)
- After the correction the E_T^{iso} distribution is centered at zero with a width of 1.5 GeV in simulated signal events



Background subtraction

- Residual background still expected even after the tight identification and isolation requirements
- A data-driven method necessary to avoid relying on detailed simulations of the background processes
- The two-dimensional sideband method:
→ photon identification γ_{ID} vs E_T^{iso} plane



- It is assumed that for background events there is no correlation between γ_{ID} and E_T^{iso}

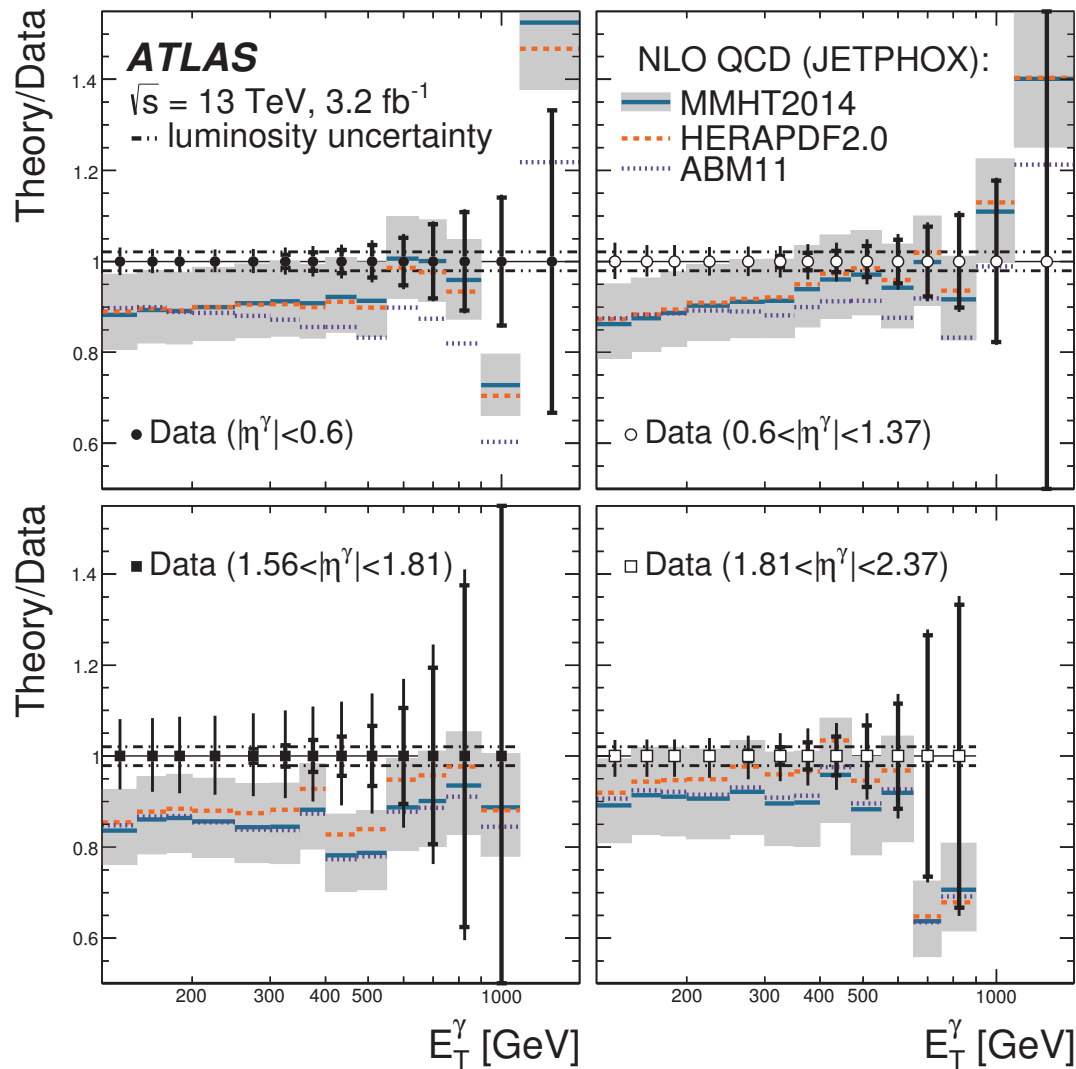
$$\frac{N_A^{bkg}}{N_B^{bkg}} = \frac{N_C^{bkg}}{N_D^{bkg}} \Rightarrow R_{bkg} \equiv \frac{N_A^{bkg} \cdot N_D^{bkg}}{N_B^{bkg} \cdot N_C^{bkg}} = 1$$

and the effects of the small signal contaminations can be accounted for by using

$$\frac{N_A - N_A^{sig}}{N_B - \epsilon_B N_A^{sig}} = \frac{N_C - \epsilon_C N_A^{sig}}{N_D - \epsilon_D N_A^{sig}} \quad \text{to extract the signal yield } N_A^{sig}$$

the leakage fractions ($\epsilon_K \equiv N_K^{sig} / N_A^{sig}$, $K = B, C, D$) are estimated using MC samples of signal \Rightarrow purity rises from 60% ($E_T^\gamma \sim 25$ GeV) to 100% ($E_T^\gamma \sim 300$ GeV)

Inclusive isolated-photon cross sections vs NLO QCD



- **NLO QCD predictions underestimate data by up to $\approx 10\text{-}15\%$**

- **Theoretical uncertainty 10-15% much larger than experimental uncertainties**

- **For $E_T^\gamma \lesssim 600 \text{ GeV}$ the measurements are systematically limited**

- **NLO QCD provides an adequate description of the data within uncertainties**

- **First measurement of inclusive photon production in the new kinematic regime opened by the LHC at $\sqrt{s} = 13 \text{ TeV}$**

- **Ready for the comparison to NNLO QCD predictions (Campbell, Ellis, Williams arXiv:1612.04333)**

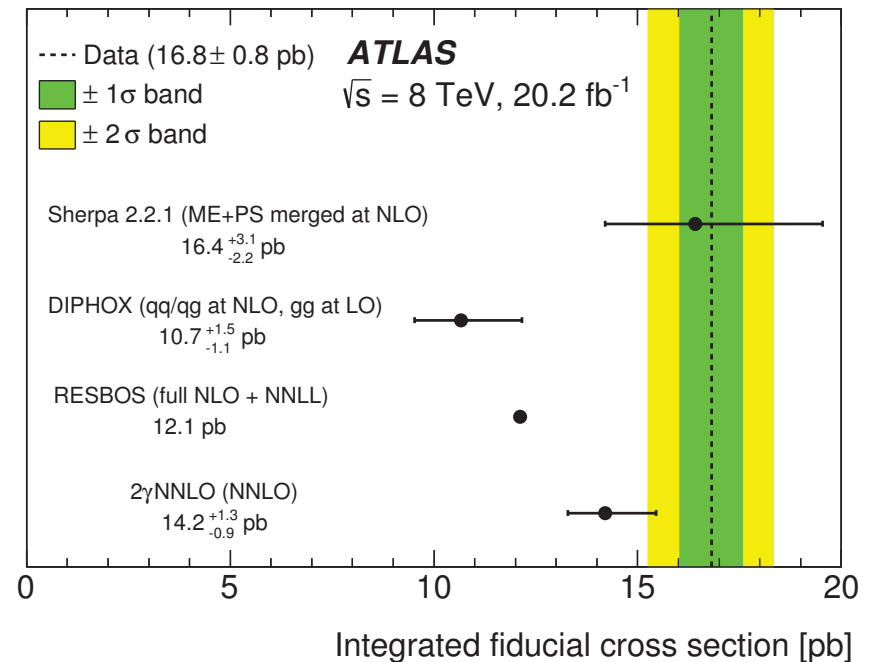
ATLAS Coll., arXiv:1701.06882, accepted PLB

Diphoton: sample composition and experimental uncertainties

Process	Event fraction [%]	
	Two-dimensional template fit	Matrix method
$\gamma\gamma$	75.3 ± 0.3 (stat) $^{+2.6}_{-2.8}$ (syst)	73.9 ± 0.3 (stat) $^{+3.1}_{-2.7}$ (syst)
γj	14.5 ± 0.2 (stat) $^{+2.7}_{-2.8}$ (syst)	14.4 ± 0.2 (stat) $^{+2.0}_{-2.4}$ (syst)
$j\gamma$	6.0 ± 0.2 (stat) $^{+1.4}_{-1.5}$ (syst)	5.8 ± 0.1 (stat) ± 0.6 (syst)
jj	1.6 ± 0.2 (stat) $^{+0.9}_{-0.4}$ (syst)	2.4 ± 0.1 (stat) $^{+0.6}_{-0.5}$ (syst)
ee	2.6 ± 0.2 (stat) $^{+0.9}_{-0.4}$ (syst)	3.5 ± 0.1 (stat) ± 0.4 (syst)

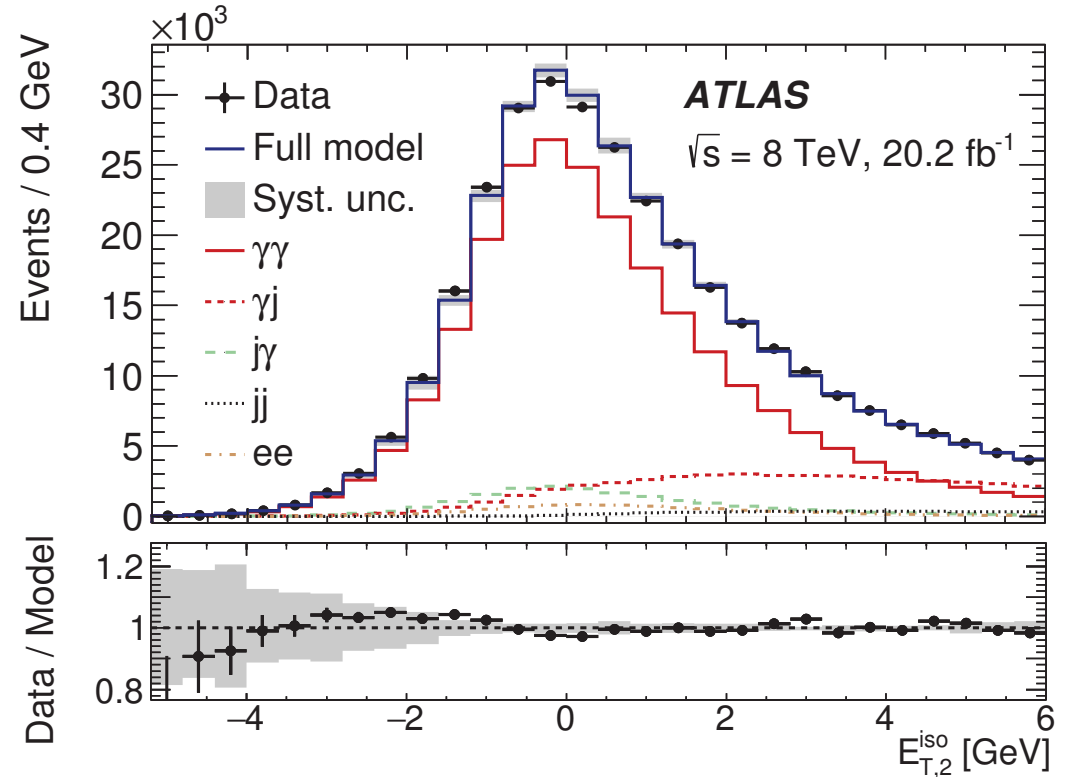
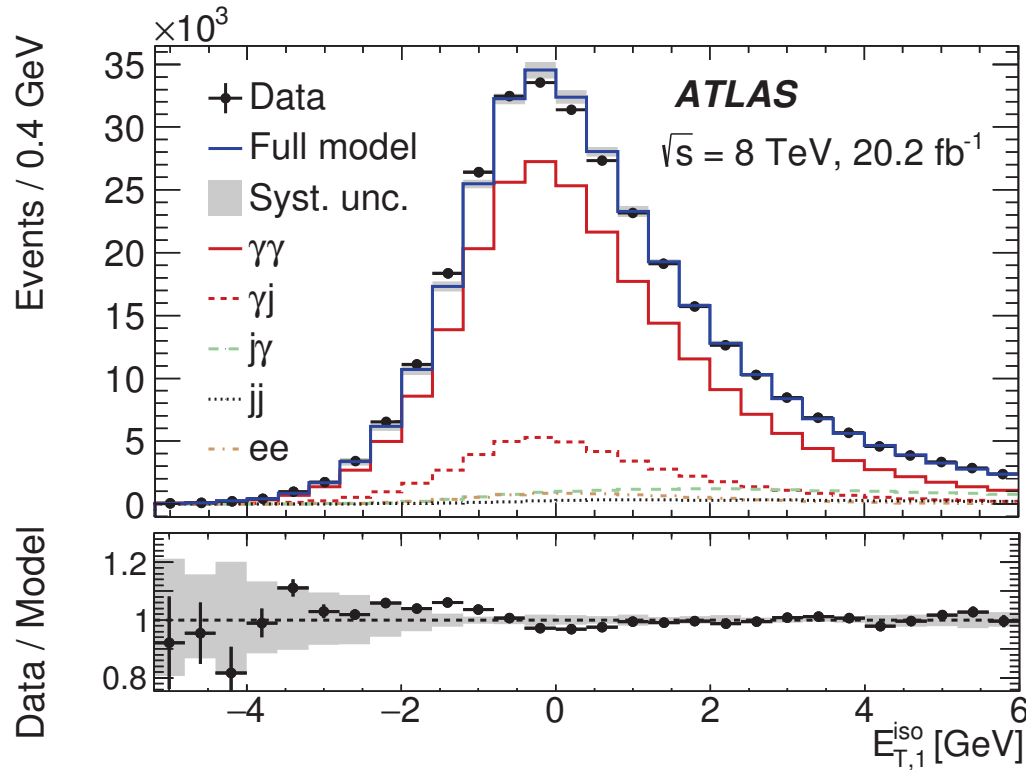
Source of uncertainty	Impact on $\sigma_{\text{tot}}^{\text{fid.}}$ [%]
Photon identification efficiency	± 2.5
Modeling of calorimeter isolation	± 2.0
Luminosity	± 1.9
Control-region definition	$^{+1.5}_{-1.7}$
Track isolation efficiency	± 1.5
Choice of MC event generator	± 1.1
Other sources combined	$^{+0.8}_{-1.3}$
Total	$^{+4.5}_{-4.7}$

Name and type of computation



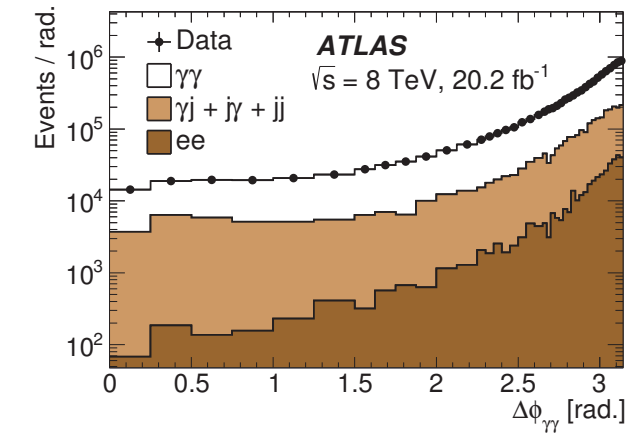
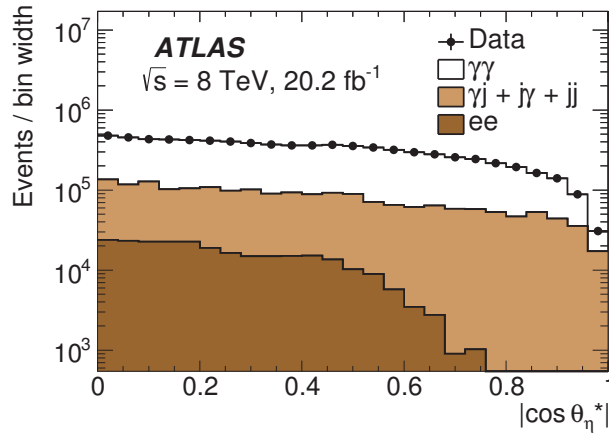
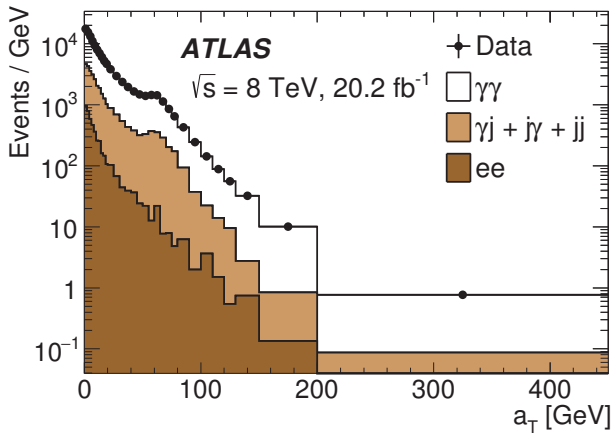
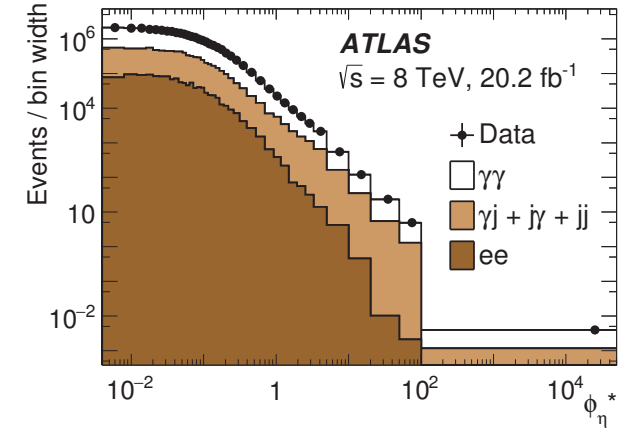
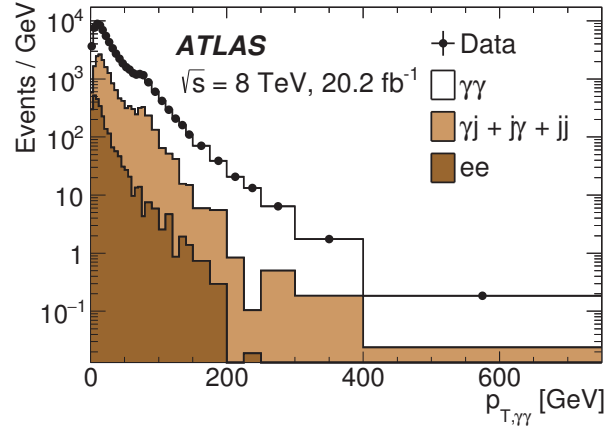
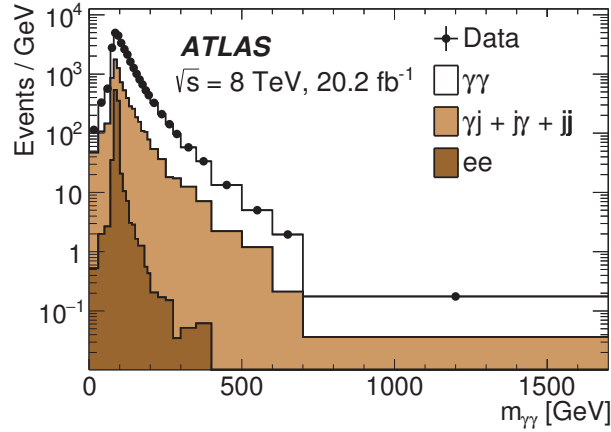
ATLAS Coll., arXiv:1704.03839

Diphoton: E_T^{iso} distributions



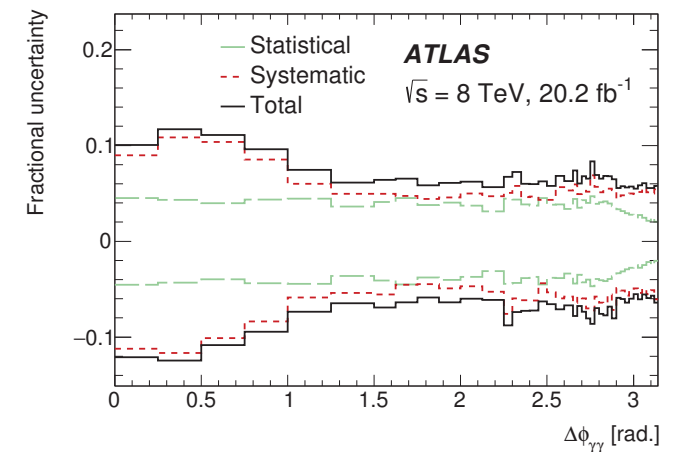
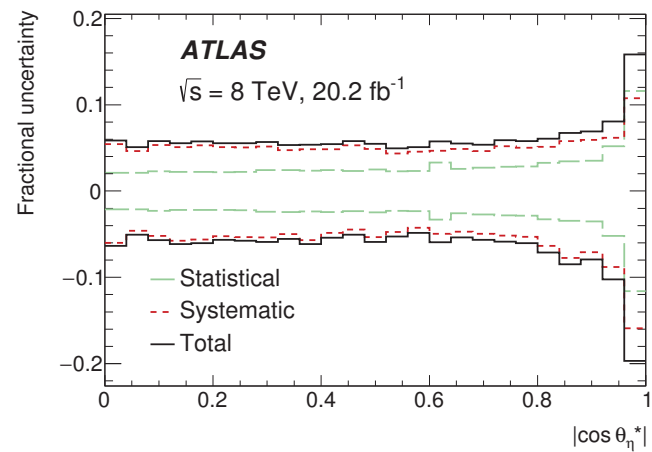
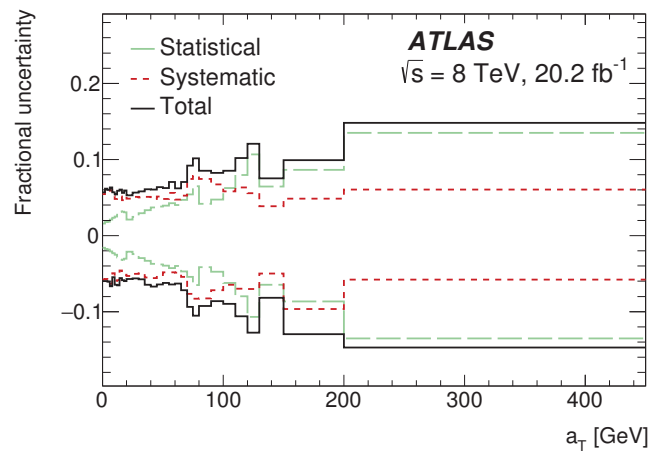
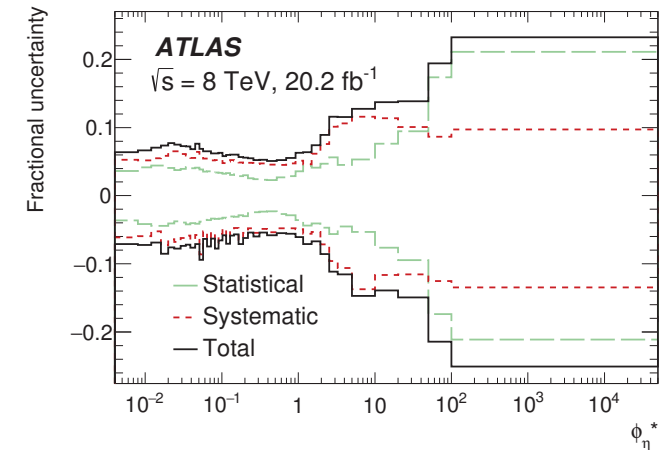
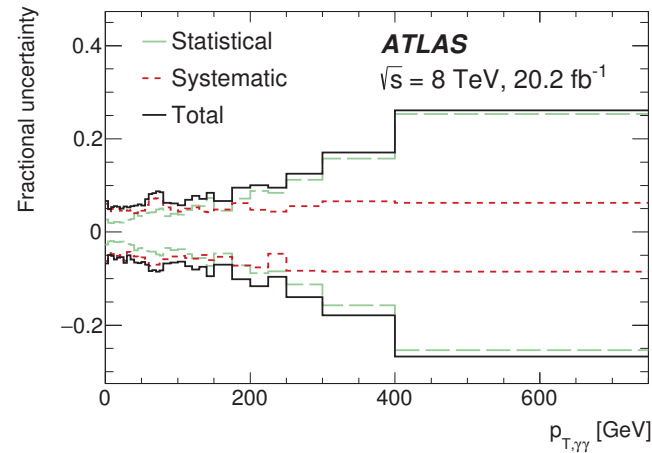
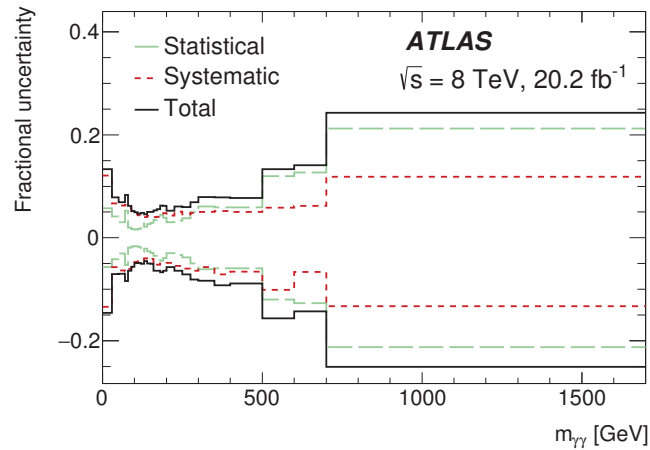
ATLAS Coll., arXiv:1704.03839

Diphoton: sample composition



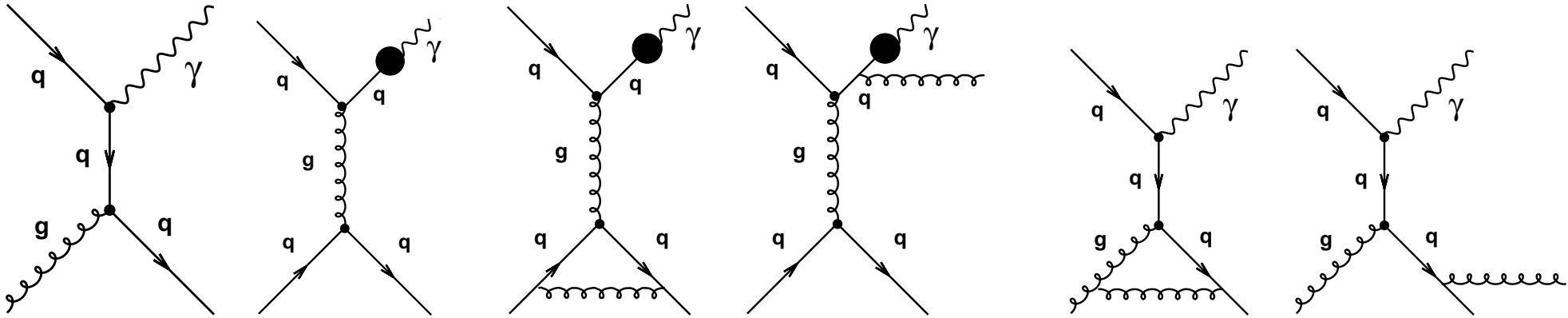
ATLAS Coll., arXiv:1704.03839

Diphoton: experimental uncertainties



ATLAS Coll., arXiv:1704.03839

NLO QCD calculations for inclusive photon production



$$\sigma_{pp \rightarrow \gamma + X} = \sum_{i,j,a} \int_0^1 dx_1 f_{i/p}(x_1, \mu_F^2) \int_0^1 dx_2 f_{j/p}(x_2, \mu_F^2) \hat{\sigma}_{ij \rightarrow \gamma a} +$$

$$\sum_{i,j,a,b} \int_{z_{min}}^1 dz D_a^\gamma(z, \mu_f^2) \int_0^1 dx_1 f_{i/p}(x_1, \mu_F^2) \int_0^1 dx_2 f_{j/p}(x_2, \mu_F^2) \hat{\sigma}_{ij \rightarrow ab}$$

- The calculations includes NLO corrections for both direct-photon and fragmentation contributions; **beware the components are not distinguishable beyond LO**
- The calculations implement the photon isolation requirement at “parton” level: E_T^{iso} calculated with the (few) final-state partons in the perturbative QCD calculation

NLO QCD calculations for inclusive photon production

$$\sigma_{pp \rightarrow \gamma + X} = \sum_{i,j,a} \int_0^1 dx_1 f_{i/p}(x_1, \mu_F^2) \int_0^1 dx_2 f_{j/p}(x_2, \mu_F^2) \hat{\sigma}_{ij \rightarrow \gamma a^+} +$$

$$\sum_{i,j,a,b} \int_{z_{min}}^1 dz D_a^\gamma(z, \mu_f^2) \int_0^1 dx_1 f_{i/p}(x_1, \mu_F^2) \int_0^1 dx_2 f_{j/p}(x_2, \mu_F^2) \hat{\sigma}_{ij \rightarrow ab}$$

- Using the JetPhox program (S. Catani, M. Fontannaz, J. Ph. Guillet and E. Pilon) with

→ $\mu_R = \mu_F = \mu_f = E_T^\gamma$ (nominal)

→ proton PDF set: CT10

→ fragmentation function: BFG set II

→ Corrections for hadronisation and underlying event needed

- Theoretical uncertainties:

→ terms beyond NLO; varying μ_R, μ_F, μ_f by factors 2 and 1/2 (singly or simultaneously)

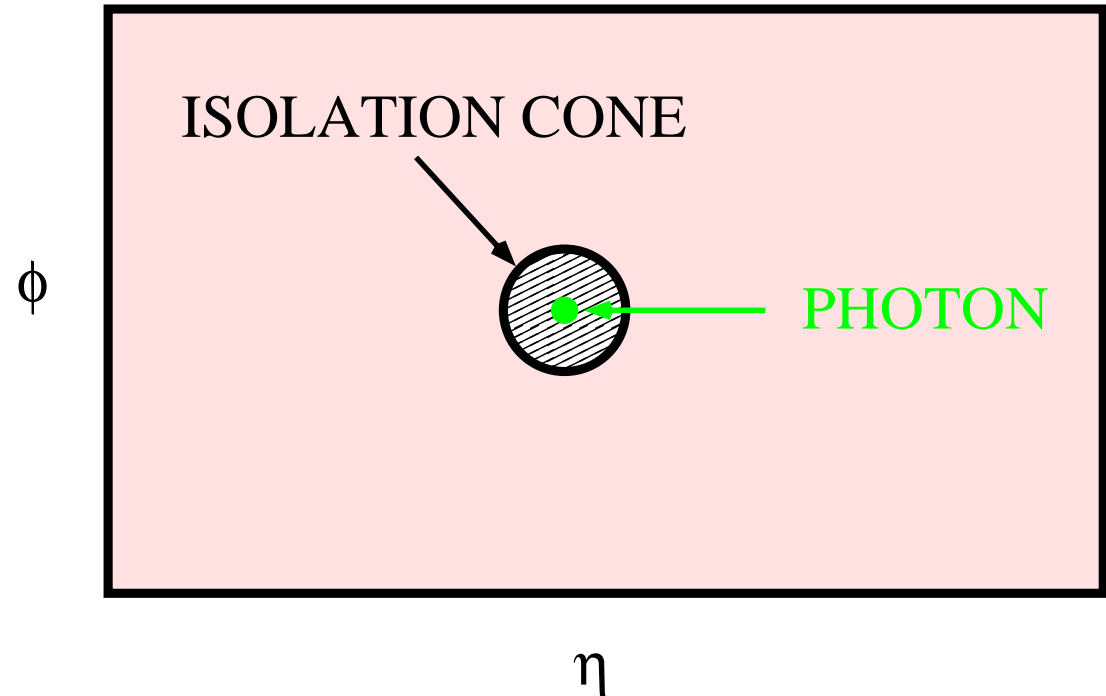
→ PDF-induced uncertainties; estimated using set of PDF eigenvectors

→ uncertainty on α_s ; estimated using PDFs in which different values of α_s are assumed

→ uncertainty on non-perturbative correction; estimated with different MCs

Corrections for non-perturbative effects; photon isolation

- The measurements are corrected for detector effects to the “particle” level
 → to isolated photons, where E_T^{iso} is calculated using all the final-state particles and the jet-area method is also applied
 This is performed using MC simulations



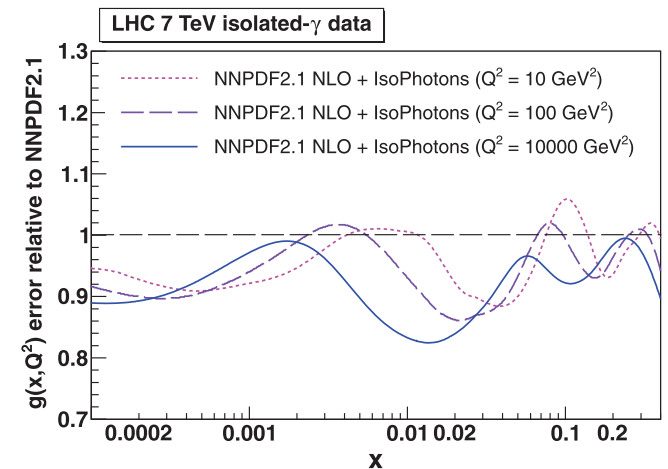
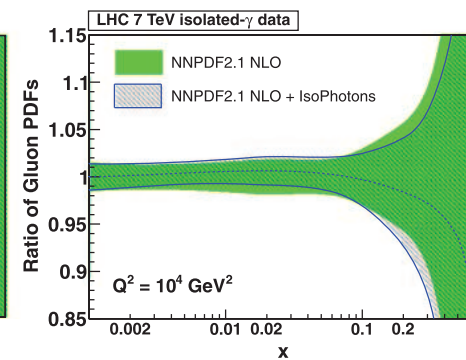
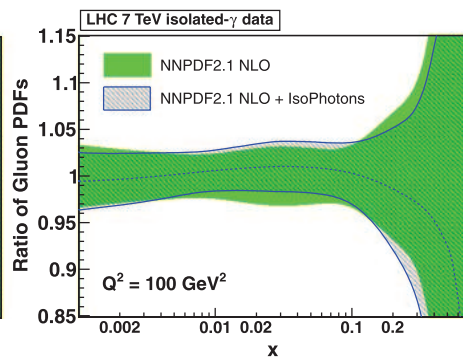
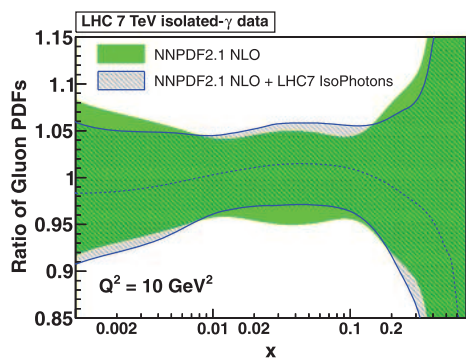
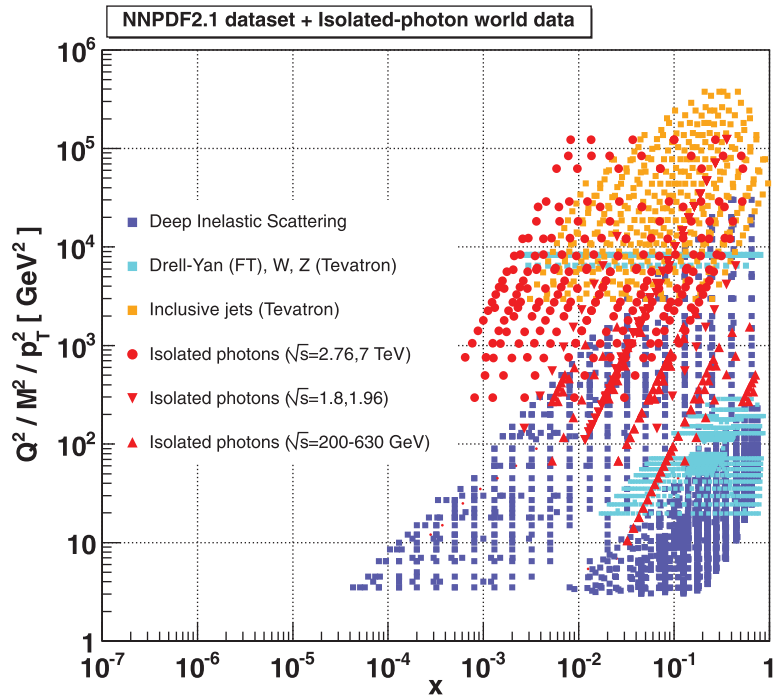
- Corrections for non-perturbative effects (hadronisation and underlying event)

$$C_{NP} = \frac{\sigma_{\gamma+X}(\text{MC, particle - level, UE})}{\sigma_{\gamma+X}(\text{MC, parton - level, no UE})}$$

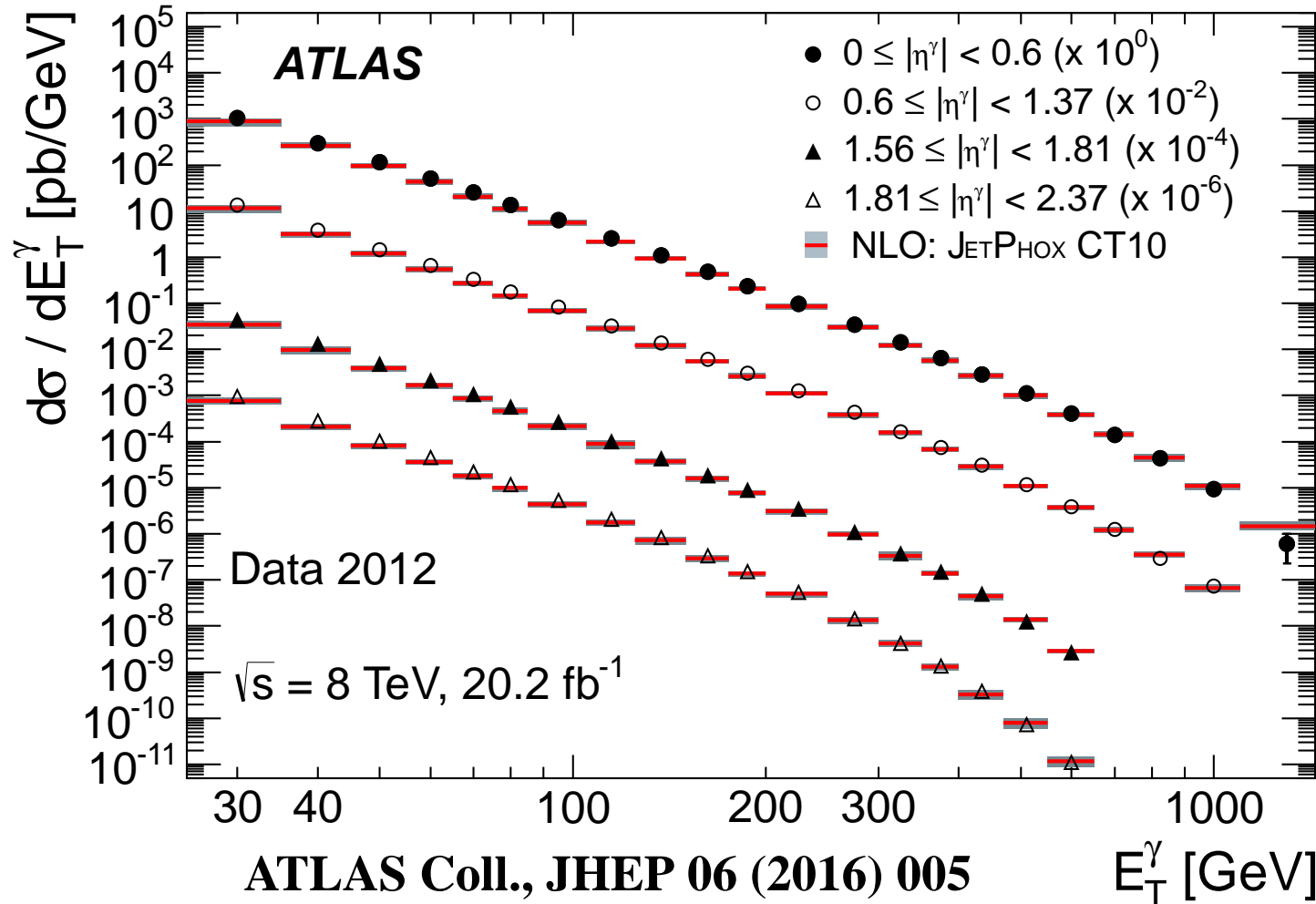
→ Less dependence on the modelling of the final state by having used the jet-area method to subtract the “extra” transverse energy contribution to E_T^{iso}

Impact of inclusive isolated photon measurements at LHC on PDFs

- Analysis by D. d’Enterria and J. Rojo (NPB860,2012,311)
 - Study of the impact on the gluon density of existing isolated-photon measurements from a variety of experiments, from $\sqrt{s} = 200$ GeV up to 7 TeV
 - those at LHC are the more constraining datasets
 - reduction of gluon uncertainty up to 20%
 - localised in the range $x \approx 0.002$ to 0.05
- ⇒ improved predictions for low mass Higgs production in gluon fusion, PDF-induced uncertainty decreased by 20%



Inclusive isolated-photon production in pp collisions at $\sqrt{s} = 8$ TeV



- Measurement of $d\sigma / dE_T^\gamma$ for $25 < E_T^\gamma < 1500$ GeV and different ranges in η^γ using $\mathcal{L} = 20.2 \text{ fb}^{-1}$ of pp collision data at $\sqrt{s} = 8$ TeV

$$E_T^{\text{iso}} < 4.2 \cdot 10^{-3} \cdot E_T^\gamma + 4.8 \text{ GeV}$$

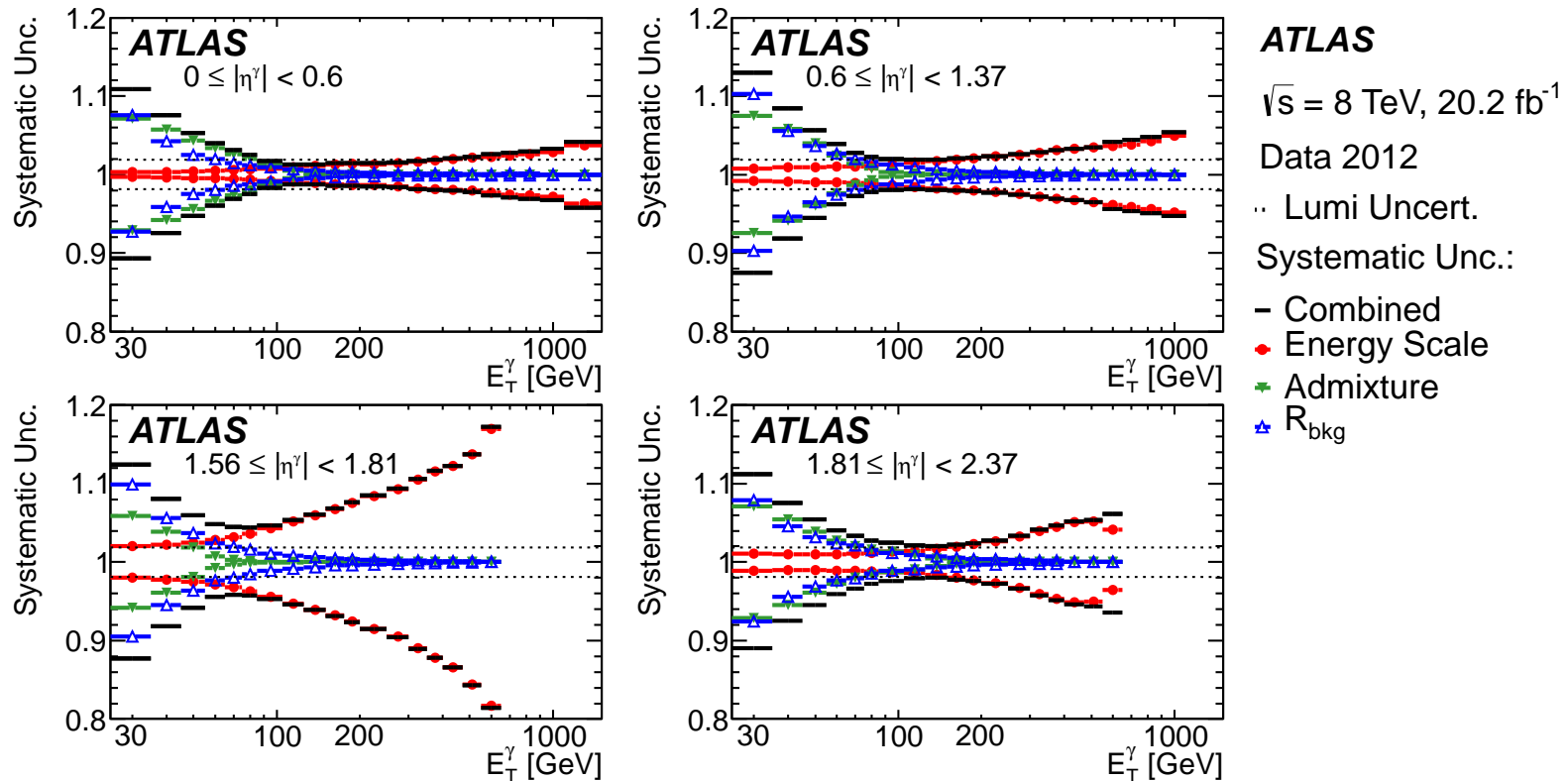
- The measurement covers ten orders of magnitude in cross section

- First measurement of photon production with $E_T^\gamma > 1$ TeV

- Significant improvement in experimental uncertainties over the previous measurements
- Good description (in log scale) of the data by NLO QCD calculations using JetPhox

Major experimental uncertainties

ATLAS Coll., JHEP 06 (2016) 005



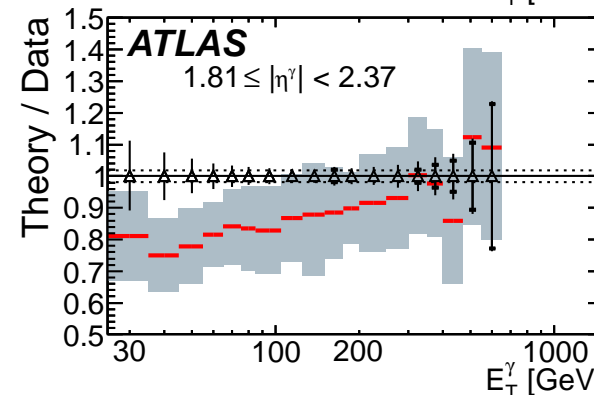
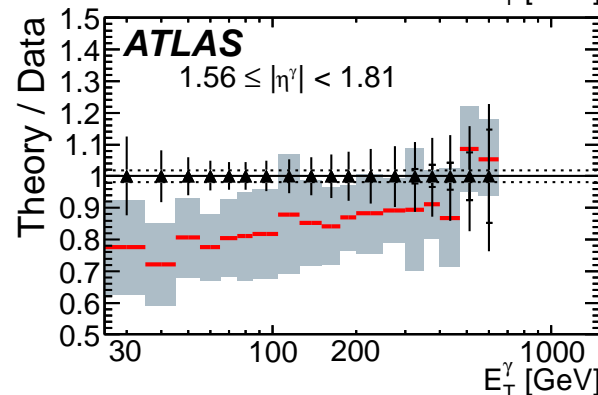
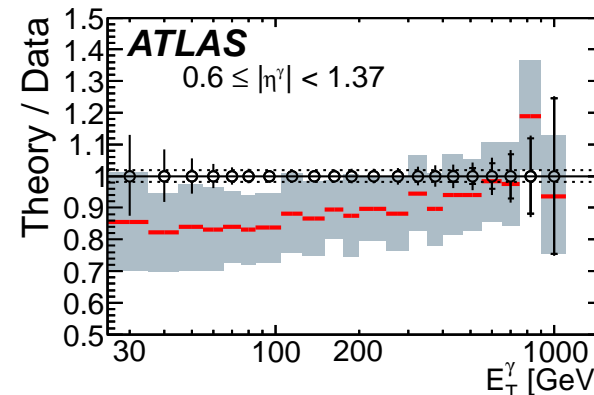
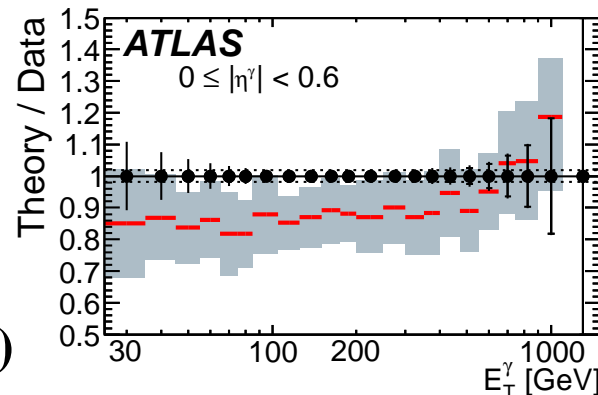
- The uncertainty on the photon energy scale* (about 1% except in the region $1.56 < |\eta^\gamma| < 1.81$) is dominant at high E_T^γ
 - The uncertainty on the correlation in the background ($\pm 10\%$) dominates at low E_T^γ , but negligible at high E_T^γ
 - The uncertainty on the admixture of direct and fragmentation photons increases at low E_T^γ
- * (ATLAS Collaboration, Eur. Phys. J. C74 (2014) 3071)

Inclusive isolated-photon cross sections vs NLO QCD

ATLAS Coll., JHEP 06 (2016) 005

Theoretical
uncertainties
much larger than(!)

experimental
uncertainties



ATLAS

 $\sqrt{s} = 8 \text{ TeV}, 20.2 \text{ fb}^{-1}$

Data 2012

● $0 \leq |\eta^\gamma| < 0.6$
○ $0.6 \leq |\eta^\gamma| < 1.37$
▲ $1.56 \leq |\eta^\gamma| < 1.81$
△ $1.81 \leq |\eta^\gamma| < 2.37$
⋯ Lumi Uncert.

NLO:

■ JETPHOX CT10

- Comparison to NLO QCD calculation using the JetPhox program

- a similar trend is observed at low E_T^γ in all $|\eta^\gamma|$ regions, the NLO QCD predictions underestimate the data by $\approx 20\%$

- the theoretical uncertainty (12-20%) prevents a more precise test of the SM predictions

- Halving the measured uncertainties compared to previous measurements

- ⇒ useful constraint on proton PDFs once included in a global fit