



New Dark Matter Search Strategies at DUNE

Jason Kumar

University of Hawaii

collaborators

- Seongjin In
- Carsten Rott
- David Yaylali

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dark matter and monoenergetic neutrinos

- can search for **dark matter** using **neutrino detectors**
 - dark matter **scatters** off solar nuclei and collects in the **core of the Sun**
 - **annihilates** to Standard Model products
 - **neutrinos** get out and reach **detector** on earth
- focus is typically on a **smooth** distribution of **high-energy** events above background
- I'll focus on a **different possibility**
 - models in which dark matter can produce **monoenergetic sub-GeV neutrinos**
 - detectors and strategies which can **resolve a line signal**
 - obtaining **direction information** about neutrino
- **DUNE** is an ideal setting for this type of search



standard lore

- expect to get a **continuum** signal
 - dark matter annihilates to **intermediate particles**
 - **decays** give a continuum neutrino spectrum
- look for **high energy** neutrinos
 - **larger cross section** with detector
 - **smaller background** from atmospheric neutrinos
- use **directionality**, but only for **high energy neutrinos**
 - try to identify neutrinos arriving from the **direction of the Sun**
 - looking for **charged lepton** produced by charged-current interaction
 - **points away from source**, but only for $E > \text{GeV}$
 - for lower energies, charged lepton is roughly isotropic



basic points

- theory
 - u, d, s final state quarks produce plenty of K^+
 - light hadrons stop before they decay (producing more K^+)
 - decay produces 236 MeV monoenergetic neutrino
- experiment
 - DUNE will do very well at total energy reconstruction for a charged-current interaction
 - sensitive to a line signal
 - DUNE can also get the direction of the neutrino from the nucleon recoil
 - new type of directionality search
 - great for reducing systematic uncertainty

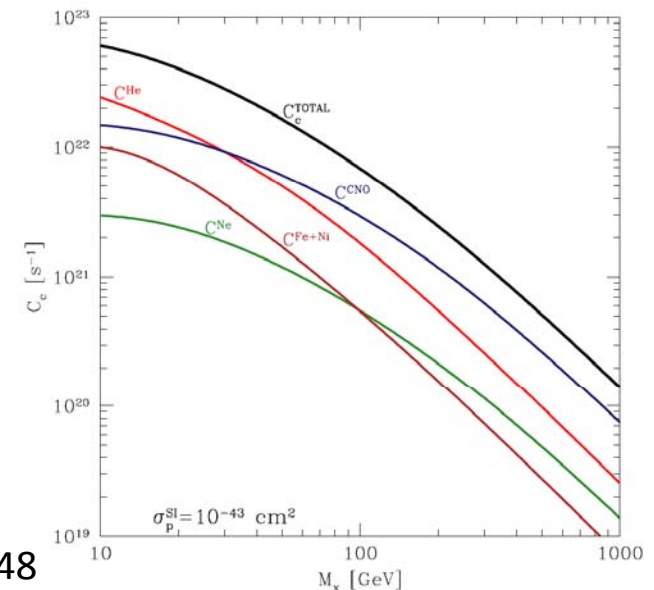
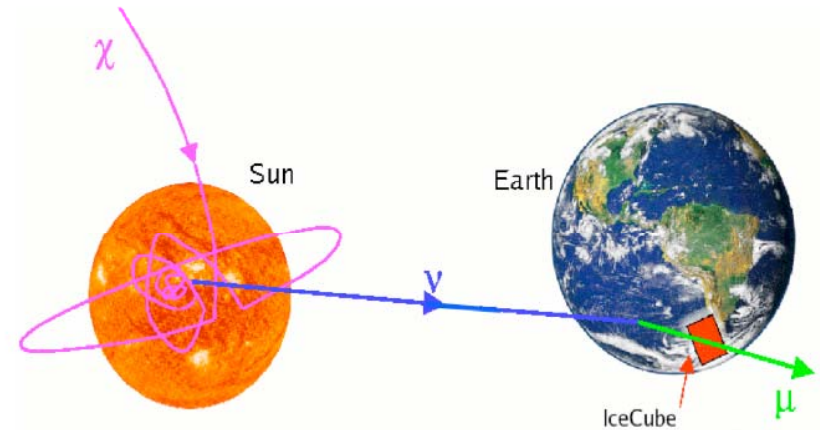


neutrinos from the Sun

(see Carsten Rott's talk....)

Dawn Williams

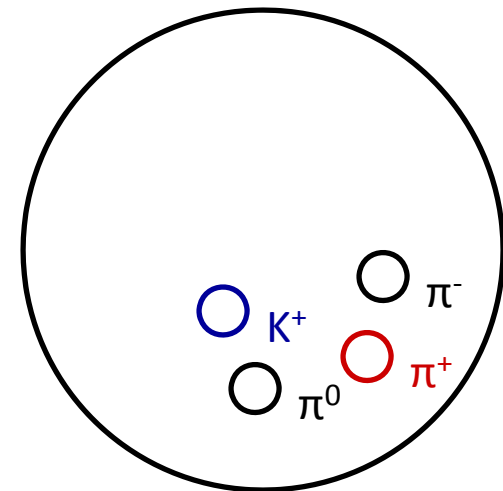
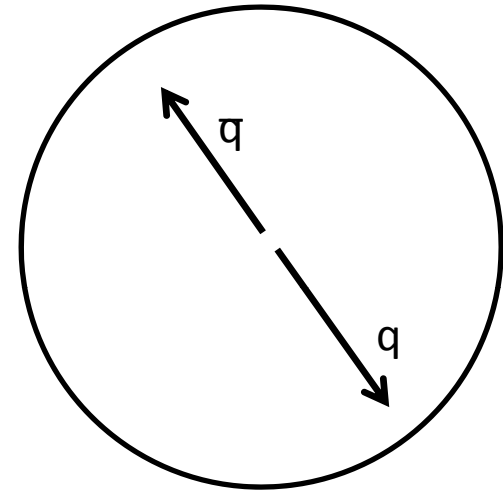
- basic idea
 - DM scatters off solar nuclei, loses energy through **elastic scattering**
 - falls below $v_{\text{esc}} \rightarrow$ **captured**
 - orbits, eventually collects in core
 - rate depends on mass, σ
 - DM **annihilates** to SM matter
 - SM decay yields **neutrinos** \rightarrow seen at detector
 - DM in **equilibrium** $\rightarrow \Gamma_C = 2 \Gamma_A$
 - so neutrino event rate probes DM capture rate ($\propto \sigma_{\text{SI}}, \sigma_{\text{SD}}$)
- usually ignore light q final state
 - why?





dark matter annihilation to light quarks

- u, d, s final states \rightarrow hard!
 - $u, d, s \rightarrow$ **light hadrons** which **stop in the Sun before decay**
 - resulting ν spectrum is very **soft**
 - **large background, small detector effective area**
- but the stopping process produces a large number of π^+, K^+
 - trade a hard spectrum for a softer one, but with **larger flux**
[Beacom, [Rott](#), Siegal-Gaskins (1208.0827); Bernal, Martin-Albo, Palomares-Ruiz (1208.0834)]





spectrum

- care about π^+ and K^+
 - $\pi^0 \rightarrow \gamma\gamma$
 - π^- Coulomb-captured by nuclei, and absorbed (not a lot of neutrinos)
- main relevant decay is $\pi^+, K^+ \rightarrow \nu_\mu \mu^+$
 - monoenergetic ν with $E = 29.8 \text{ MeV}$ (π^+ - 100%) or 235.5 MeV (K^+ - 64%)
 - line signal
 - include oscillation effects
- just need the fraction of DM energy which goes into stopped π^+, K^+
 - determine with Pythia/GEANT
 - use Pythia to simulate showering and hadronization; output the spectrum of long-lived hadrons
 - GEANT deals with interaction in dense solar medium



K^+

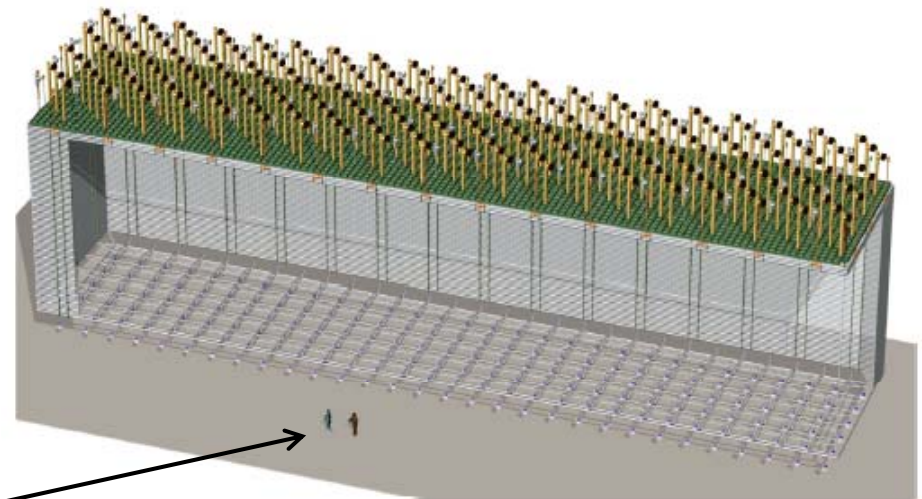
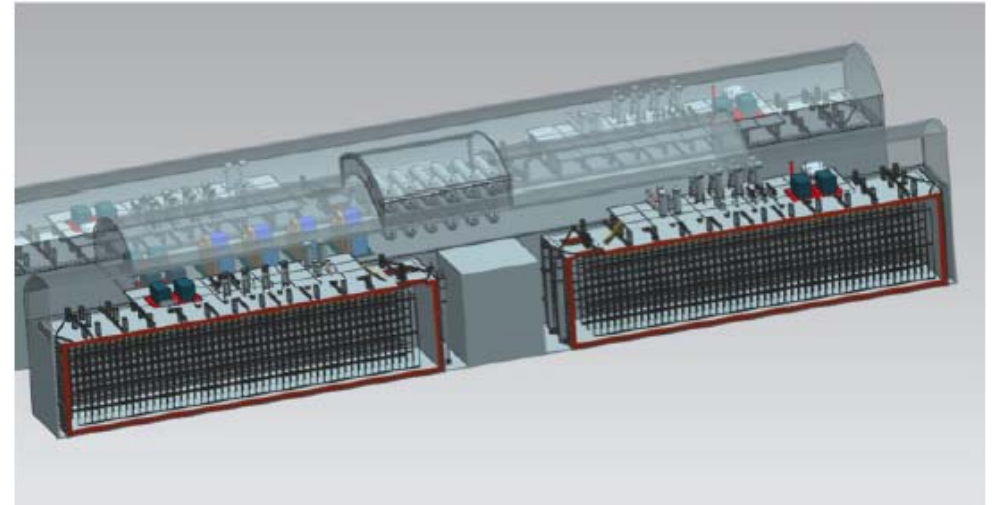
- I'll focus on the **236 MeV neutrino** arising from stopped K^+ decay
- **much larger cross section** with detector target
 - more than offsets smaller number of kaons per annihilation
- now have all the pieces
 - given **dark matter mass, scattering cross section, and annihilation channel**, can get the **flux of 236 MeV neutrinos** from the Sun
 - with the **energy resolution**, can get the flux from **atmospheric neutrino background**
 - gives us the **signal-to-background ratio**
 - with the **neutrino-nucleus scattering cross section** (numerical) and **exposure**, can get **signal significance**



DUNE

1601.02984

- Deep **U**nderground **N**eutrino **E**xperiment
- perfect for this type of search
 - large exposure
 - good total energy resolution
 - can identify outgoing particle tracks with good energy and angular resolution
- our benchmarks
 - angular resolution $\sim 5^\circ$
 - total energy res. – $\epsilon \sim 10\%$



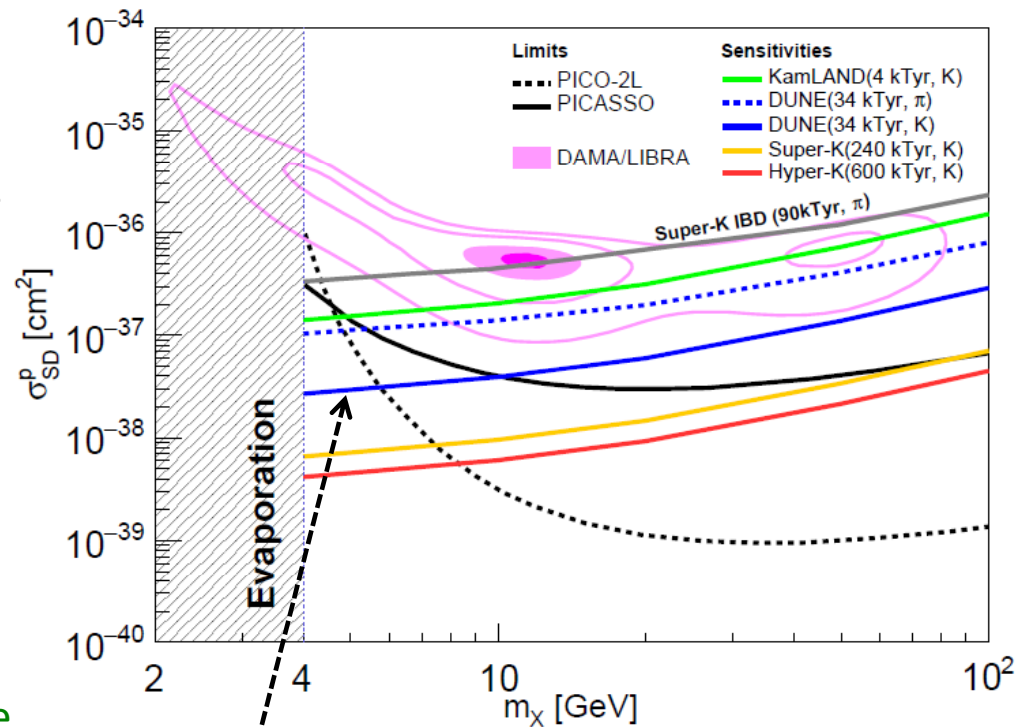
a theorist, for scale



sensitivity for non-directional search

- assume 34 kT yr exposure
 - electron channel
 - ~ 50 bgd. events
 - 90%CL exclusion, assuming observation consistent with bgd.
 - sig. signif. $\propto (\text{exposure} / \epsilon)^{1/2}$
- competitive with direct detection at ~ 4-5 GeV (but PICO-60 wins above this)
- SK, HK \rightarrow win with exposure
 - WC detectors \rightarrow size advantage
- other neutrino searches not sensitive (focused on high-energy neutrinos)

In, Kumar, Rott, Yaylali - 1510.00170



90% CL
annihilation to u or d
spin-dep. scattering



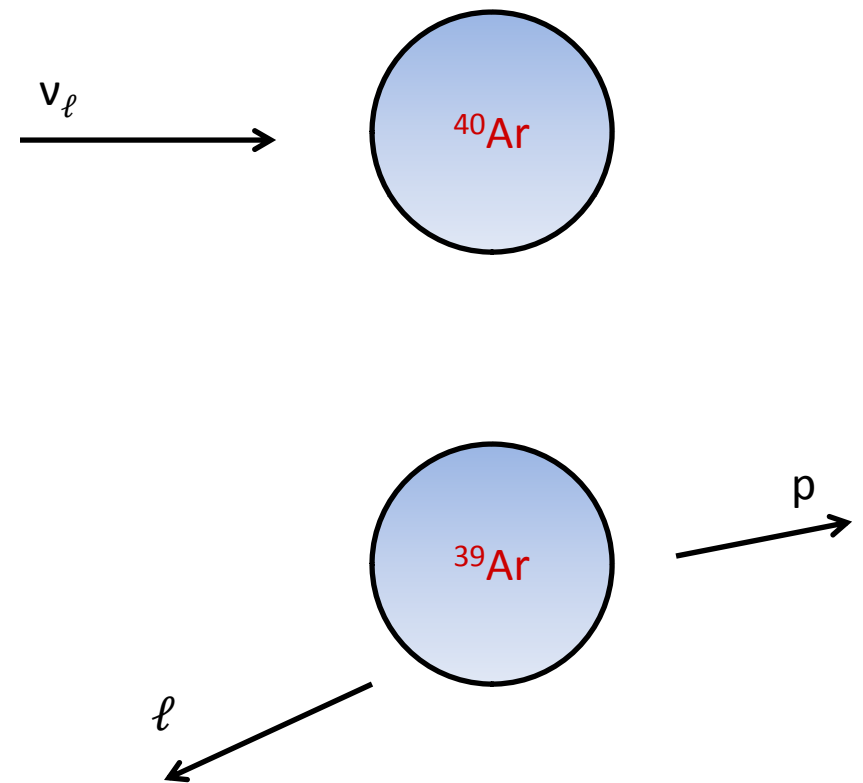
directionality

- for 34 kT yr exposure, DUNE atm. ν background is significant
- would be great to get a **directionality** cut
 - preferentially select events where ν arrives from the direction of the Sun
- reduces **systematic** uncertainties in background by comparing **on-axis** to **off-axis** event rates
 - want $S / B > \delta B_{\text{sys.}} / B \rightarrow$ excess not just a systematic error
 - can measure B by going off axis (**reduces $\delta B_{\text{sys.}} / B$**)
 - **increases S / B** by picking events from the direction of the Sun
- can **improve statistical significance**
- most searches for neutrinos arising from dark matter annihilation utilize directionality...
- ... but usually when looking for a very **energetic neutrino**
 - CC-interaction produces a forward-peaked **charged lepton**



directionality for sub-GeV ν s

- for **sub-GeV** ν , the **charged lepton** produced is mostly **isotropic**
- but the **hadronic recoil is not!**
- at this energy, get a lot of events where a single **proton** is ejected
 - $\nu_\ell + {}^{40}\text{Ar} \rightarrow \ell + p + {}^{39}\text{Ar}$
- ejected in the **forward** direction
 - **cut on proton direction**
- but analytic approximations to the **cross sections** and **distributions** are lacking
- rely on **numerical** techniques
- NuWro





NuWro

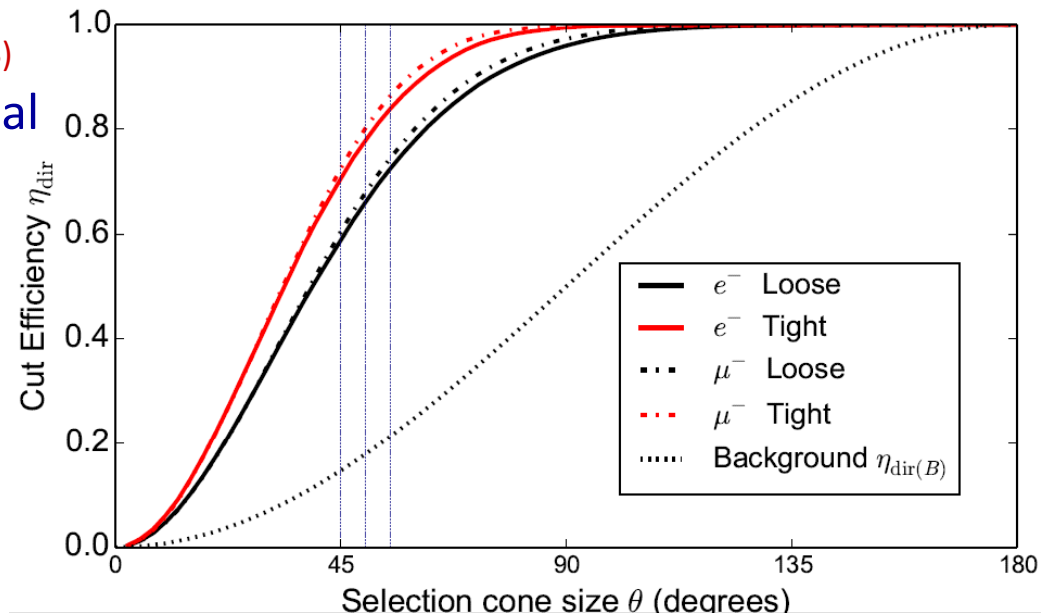
- generate 10^5 events per flavor (ν_e or ν_μ)
- select events with...
 - one charged lepton track identified
 - one ejected proton track identified (kills $\bar{\nu}$ bgd.)
 - cuts at event generation level (no attempt to model detector)
 - just need particles generated above a threshold
- lepton threshold \rightarrow 30 MeV
- proton threshold \rightarrow 50 MeV (according to DUNE CDR...)
 - “tight”
- we’ll also consider a more optimistic proton cut \rightarrow 20 MeV
 - “loose”
- determine efficiency for signal events (η_S) and bgd events (η_B) to satisfy event selection and angular cuts



sensitivity and systematics

- two efficiencies
 - event selection (η_{sel})
 - common to S and B
 - directional (η_{dir})
 - fraction of events in forward cone from the Sun
 - better for S than for B
- total efficiency $\eta_{S,B} = \eta_{sel} \times \eta_{dir(S,B)}$
- care about improvement to signal significance, and to S-to-B ratio
- we'll choose cuts to maximize improvement for signal significance for fixed exposure
 - other choices possible...

cut	proton threshold	selection efficiency (η_{sel})
tight:electron	$E_{kin} > 50$ MeV	0.43
tight:muon	$E_{kin} > 50$ MeV	0.28
loose:electron	$E_{kin} > 20$ MeV	0.83
loose:muon	$E_{kin} > 20$ MeV	0.75





cuts and efficiencies

$$\frac{S}{B} \rightarrow \left(\frac{\eta_S}{\eta_B} \right) \times \frac{S}{B}$$

$$\frac{S}{\sqrt{B}} \rightarrow \left(\frac{\eta_S}{\eta_B} \sqrt{\eta_B} \right) \times \frac{S}{\sqrt{B}}$$

cut	S/B enhancement	sensitivity enhancement
tight:electron	4.8	1.2
tight: muon	4.5	1.0
loose:electron	3.4	1.4
loose:muon	3.5	1.4

tight → win on S/B (up to S/B ~0.4)

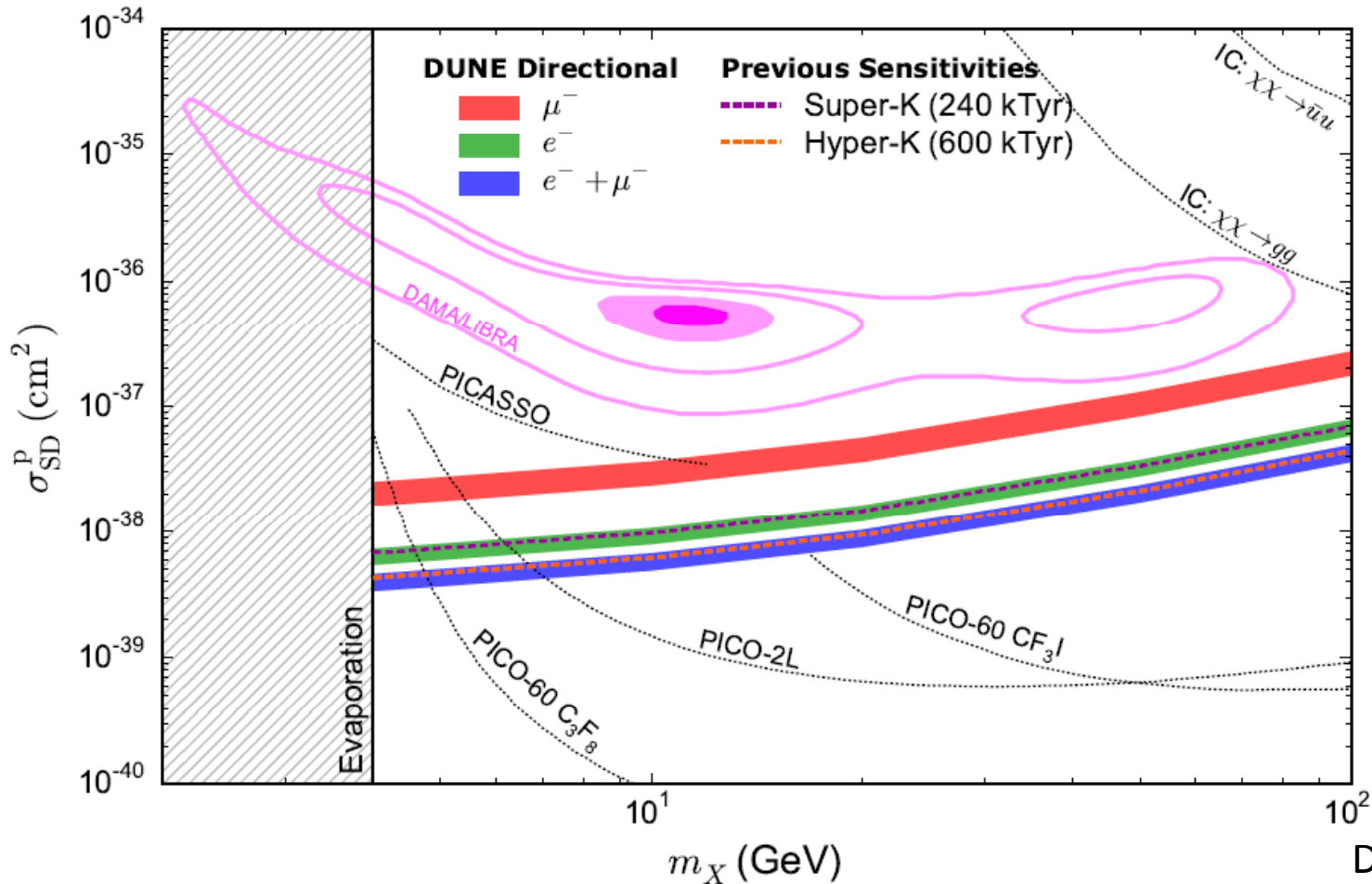
loose → win on sensitivity

- cuts: cone half-angle (\gg ang. res)
- tight: muon → 45°
- tight: electron → 50°
- loose: electron → 55°
- loose: muon → 55°
- S/B can improve by up to $\times 5$
 - very good for on-/off-axis
- but signal significance only sees a modest improvement
 - big hit from small selection efficiencies
- win more on systematics than statistics



results

90%CL, q=u,d



assume 340 kT yr ... need large exposure to offset selection efficiencies
dozens of background events
need a long run-time just to catch up to Super-K and PICO-60



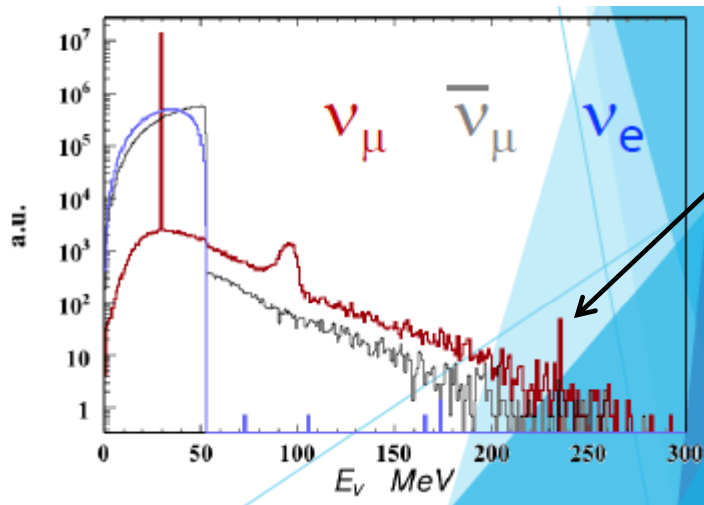
what's the point of doing this at DUNE?

- for signal significance, **WC detectors** will always win because of exposure
- except for a small mass range, **PICO** is already winning
- but there are **good reasons** to search at DUNE
- **directionality** gives a new handle on systematic uncertainties and bgd.
 - **no such directionality possible** with WC detectors
 - PICO sensitivity is **degrading** rapidly < 10 GeV
 - different **astrophysics uncertainties** than direct detection
- if a **signal** is seen in the future, can get a handle on **annihilation channel**
 - is it **asymmetric dark matter**?
 - a 236 MeV line signal at DUNE from the Sun would be **striking evidence** of dark matter annihilation producing **light quarks**
 - cross section could be $\ll 1$ pb, with Sun still in equilibrium
 - especially for low mass DM, **hard to see this any other way**
- important as a **complementary** search strategy



resolving uncertainties

- a lot of uncertainty in the **neutrino-nucleus scattering cross section**, etc.
 - really a **proof-of-principle**
- can “calibrate” by comparing rates **on-axis** vs. background off-axis
- but can also calibrate directly with a **stopped kaon experiment**
- a **stopped pion experiment** is also a stopped kaon experiment
- stopped pion proposals like **DAEδALUS** are under consideration for DUNE
- can also put an LArTPC at a stopped pion experiment
 - **CAPTAIN** at **SNS** (see Lisa Whitehead Koerner’s parallel talk....)



conclusion

- dark matter annihilation in the Sun can produce monoenergetic 236 MeV neutrinos
 - produce numerous stopped K^+
- LArTPC ν -detectors can reconstruct energy and direction of products
 - can detect a neutrino line with good total energy resolution
 - can get directionality from ejected proton
- reduced backgrounds and systematic uncertainties
- sub-GeV ν directionality is a unique capability of DUNE
- stopped kaon experiment would help with calibration
- above all, need lots of exposure

Mahalo!



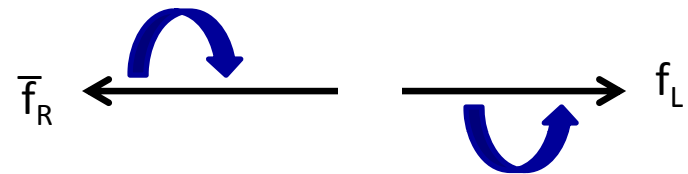
Back-up slides



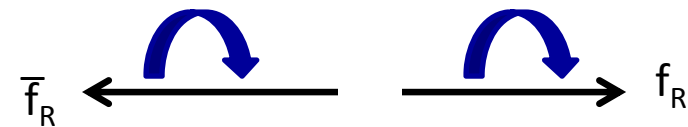
why (not) $\chi\chi \rightarrow \bar{f}f$ ($f=u,d,s$)?

- can understand just from **angular momentum**
- for Majorana fermion, wavefunction is **anti-symmetric**
 - $L=0, S=0$ or $L=1, S=1$
- if outgoing fermions on z-axis
 - $L_z=0$ ($Y_{lm}(\theta=0,\phi) \neq 0$ only if $m=0$)
 - $S_z = J_z$
- if $S_z=0$ need f, \bar{f} with **same helicity**
 - **not** CP-conjugate
 - need Weyl spinor **mixing**
 - in **MFV**, mixing scales with **mass**
- if $S_z=\pm 1$ need f, \bar{f} with **opp. helicity**
 - **no mixing** needed

$$J=0, L_z=0 \rightarrow S_z=0$$



$$J_z=1, L_z=0 \rightarrow S_z=1$$





monoenergetic neutrinos

- this argument underlies the **theoretical prejudice** towards searches for the $\bar{b} b$, $\bar{\tau} \tau$ and $W^+ W^-$ channels
- but the **chirality suppression** arises from the assumption of **Majorana fermion dark matter** and **minimal flavor violation**
 - certainly true for the **CMSSM**, but **need not be true in general**
 - WIMPs need not be Majorana, and **MFV can fail even in the general MSSM**
- if dark matter is a **Dirac fermion**, then the initial state can be $L=0, S=1, J=1$, so **s-wave annihilation**, but **no mixing needed**
- if we **drop minimal flavor violation**, then **mixing need not scale with quark mass**
- either way, $\chi\chi \rightarrow \bar{q}q$ ($q = u, d, s$) **branching fraction could be $\mathcal{O}(1)$**
- worth studying these annihilation channels



signal limited – K^+

- compare to π^+ channel
- larger cross section
 - larger effective area ($> \times 100$)
 - need smaller exposure to get signal (or bgd.) events
- fewer K^+ per annihilation
 - backgrounds similar
 - smaller flux, larger bin
 - factor 5-10 smaller S/B
- upshot
 - better sensitivity with small exposure
 - leaves linear regime first
 - ultimate sensitivity comparable

$$N_{\text{bgd.}, 236 \text{ MeV}}^{\text{DUNE}} \approx (486)\epsilon \left(\frac{M_{\text{Ar}}}{34 \text{ kT}} \right) \left(\frac{\text{T}}{\text{yr}} \right)$$

$$\frac{S}{B} \approx \frac{0.83}{\epsilon} \frac{\sigma_{\text{SD}}^p}{\text{pb}}$$

$$E_{\text{opt}}^{\text{Ar}} \approx \frac{0.07}{\epsilon}$$



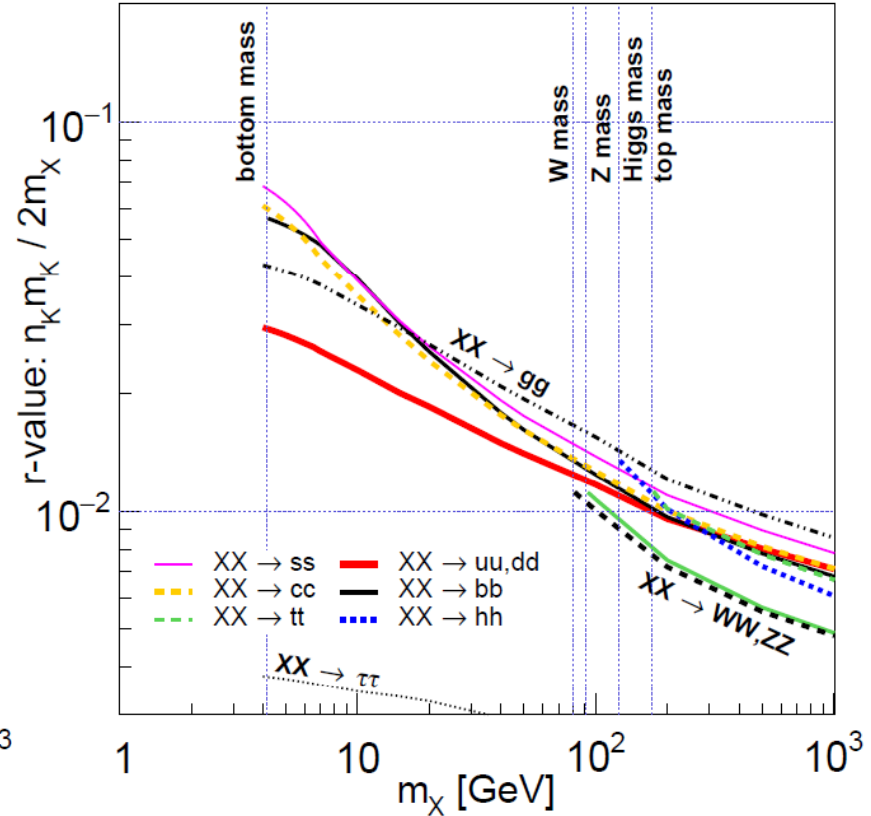
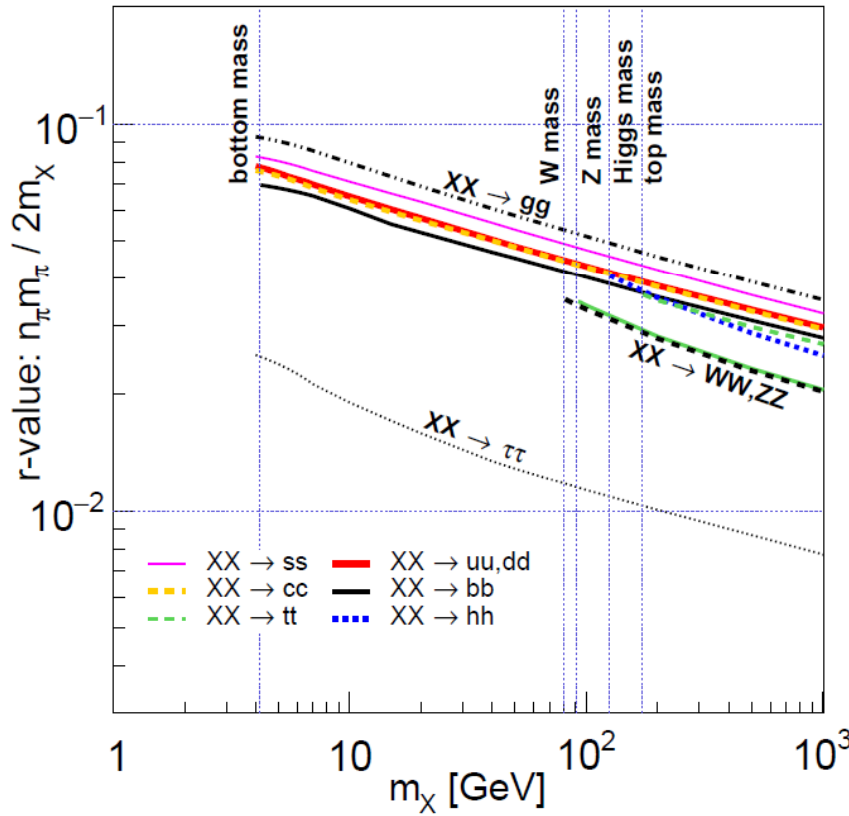
issues with the cross section at $\mathcal{O}(100)$ MeV (a novice's view)

- basic idea \rightarrow **impulse approximation** (IA)
 - neutrino interacts with a **single struck nucleon**
 - **subsequent interactions** between struck nucleon and rest of nucleus
- can model the nucleus state as...
 - **Fermi gas**
 - using a more detailed **spectral function** obtained from theory and electron scattering experiments
- **spectral function is a better model...**
- ...but analysis still based on IA
- IA becomes **less valid** an approximation for $E_\nu < 100$ MeV
 - Ankowski, Soczyk -- 0709.2139
 - no good tool for going beyond IA, though
 - best to just calibrate



r-factors

1510.00170



r decreases with m_X

about $\times 10$ more 30 MeV vs than 236 MeV vs per annihilation for u and d channels



the pieces we need....

- we have the **neutrino fluxes from the Sun arising from DM....**
- we have **estimates of the $\nu_{e,\mu}$ background** at $E \sim 236$ MeV (atm. ν)
- **charged current neutrino-nucleus scattering cross section** ($\nu_\ell + n \rightarrow \ell + p$)
 - for $E \sim 236$ MeV, theory complicated
 - dominant contribution is **quasi-elastic**
 - not very well understood
 - rely on numerical packages
 - **NuWro**

Battistoni, Ferrari, Montaruli, Sala

$$\frac{d^2\Phi_B^e}{d\Omega dE} \approx 1.2 \text{ m}^{-2}\text{s}^{-1}\text{sr}^{-1}\text{MeV}^{-1}$$

$$\frac{d^2\Phi_B^\mu}{d\Omega dE} \approx 2.3 \text{ m}^{-2}\text{s}^{-1}\text{sr}^{-1}\text{MeV}^{-1}$$

($\bar{\nu}$ similar)

$$\sigma_{CC}^e(236 \text{ MeV}) \approx 4.2 \times 10^{-38} \text{ cm}^2$$

$$\sigma_{CC}^\mu(236 \text{ MeV}) \approx 2.7 \times 10^{-38} \text{ cm}^2$$

NuWro



90% CL numbers

non-directional search, electron channel

experiment	status	exposure	N_B^π	N_{obs}^π	f_S^π	N_S^π	N_B^K	N_{obs}^K	f_S^K	N_S^K
KamLAND	current	4 kT yr	—	—	—	—	5.1	6	0.68	5.5
DUNE	future	34 kT yr	0.2	0	1	2.3	50	50	0.68	10.3
Super-K	current	240 kT yr	—	—	—	—	305	305	0.68	28.7
Hyper-K	future	600 kT yr	—	—	—	—	762.5	763	0.68	45.4

directional search, 340 kT yr exposure

cuts	expected N_B	assumed N_{obs}	expected N_S for exclusion
tight: electron	14.8	15	6.5
tight: muon	14.9	15	6.4
loose: electron	41.6	42	10.0
loose: muon	47.5	48	10.7