

Circular polarisation as a new probe of BSM physics in space

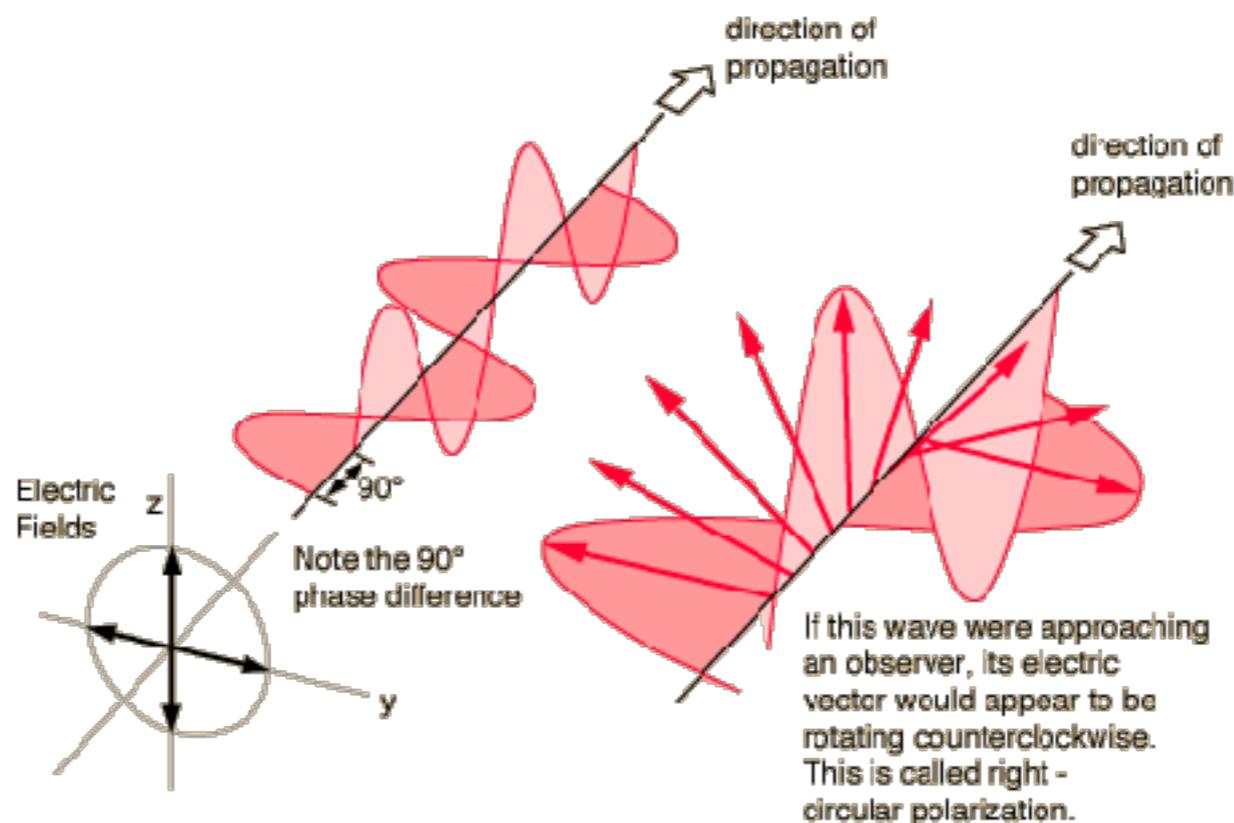
Céline Boehm

Céline Degrande

Olivier Mattelaer

Aaron C. Vincent

arXiv:1701.02754



*Courtesy
Hyper physics*

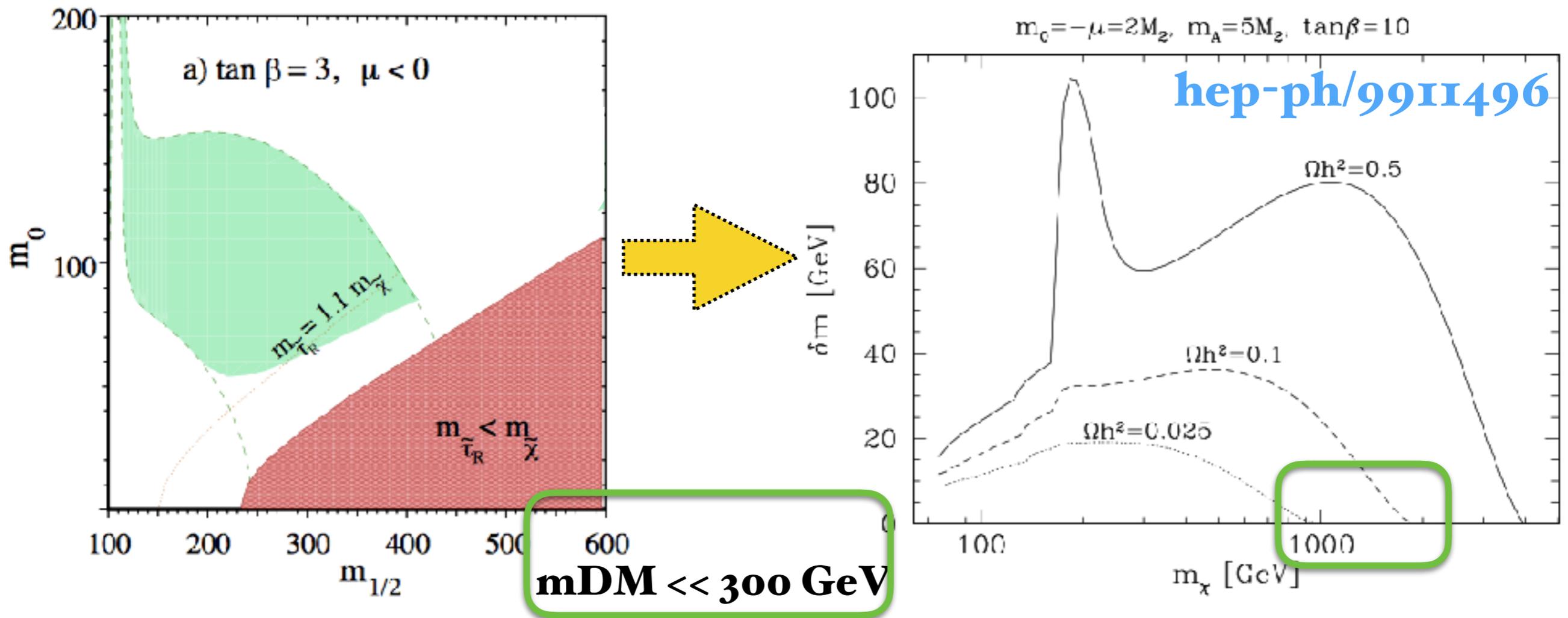
Quick recap of the DM world

“Intensity” & energy frontiers in the **DM** world

Energy is critical as DM may be heavy

Supersymmetry Extra-Dimensions ...

Relic density criterion alone (before WMAP!)



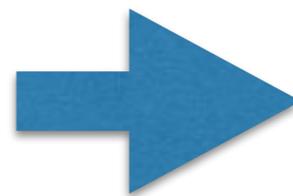
Larger DM masses need higher energy colliders!

“Intensity” & energy frontiers in the **DM** world

Intensity is critical! **not only for new heavy mediators but also for Dark (5th) force**

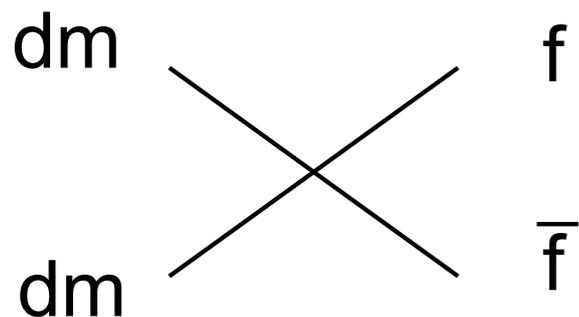
Hut, Lee&Weinberg 77

$$\frac{dn}{dt} = -3Hn - \sigma v (n^2 - n_0^2)$$

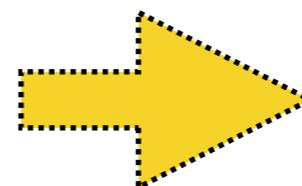


$$\Omega h^2 \simeq \frac{3 \times 10^{-27} \text{ cm}^3 / \text{s}}{\langle \sigma v \rangle}$$

$$\sigma v \sim 3 \times 10^{-26} \text{ cm}^3 / \text{s}$$



$$\sigma v \propto \frac{m_{\text{dm}}^2}{m_{\text{w}}^4}$$



$$\Omega_{\text{DM}} h^2 \propto m_{\text{DM}}^{-2}$$

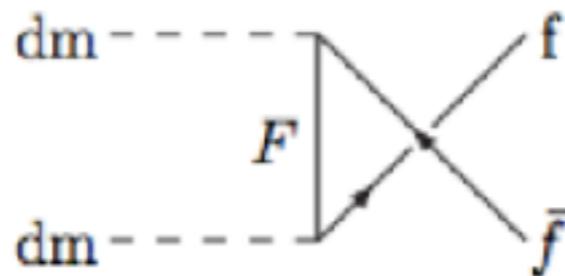
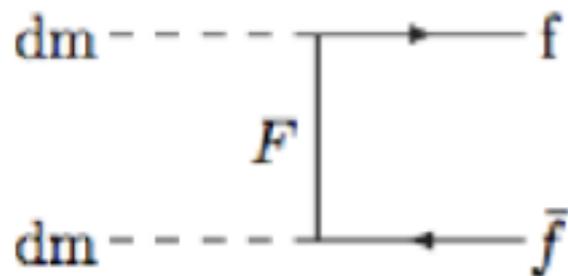
Dark Matter heavier than a proton unless new types of interactions

“Intensity” & energy frontiers in the **DM** world

Intensity is critical! [astro-ph/0208458v3](#) [hep-ph/0305261](#)

$$\Omega_{\text{DM}} h^2 \propto \cancel{m_{\text{DM}}^{-2}} \quad \text{or...}$$

Dark 5th force and vector-like fermions

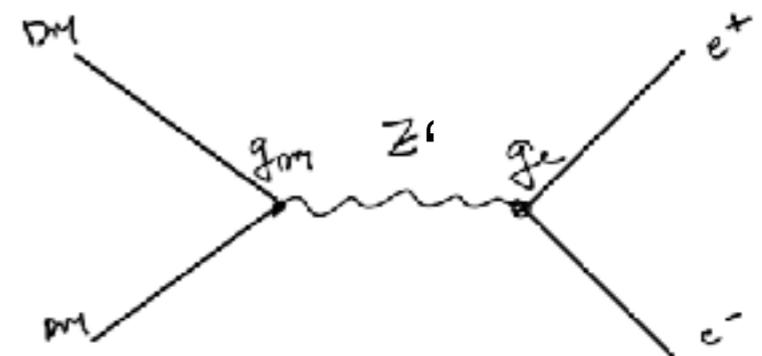


$$\sigma v \propto \frac{1}{m_F^4} \left((C_l^2 + C_r^2) m_f + 2C_l C_r m_F \right)^2$$

$$\sigma v \propto \frac{1}{m_F^2}$$

mirror fermions

**Such heavy mediators need
fb, ab luminosity!**



$$\sigma v \propto v^2 \frac{m_{\text{DM}}^2}{m_{Z'}^4} g_{\text{DM}}^2 g_e^2$$

$$m_{\text{DM}} \simeq m_{Z'}$$

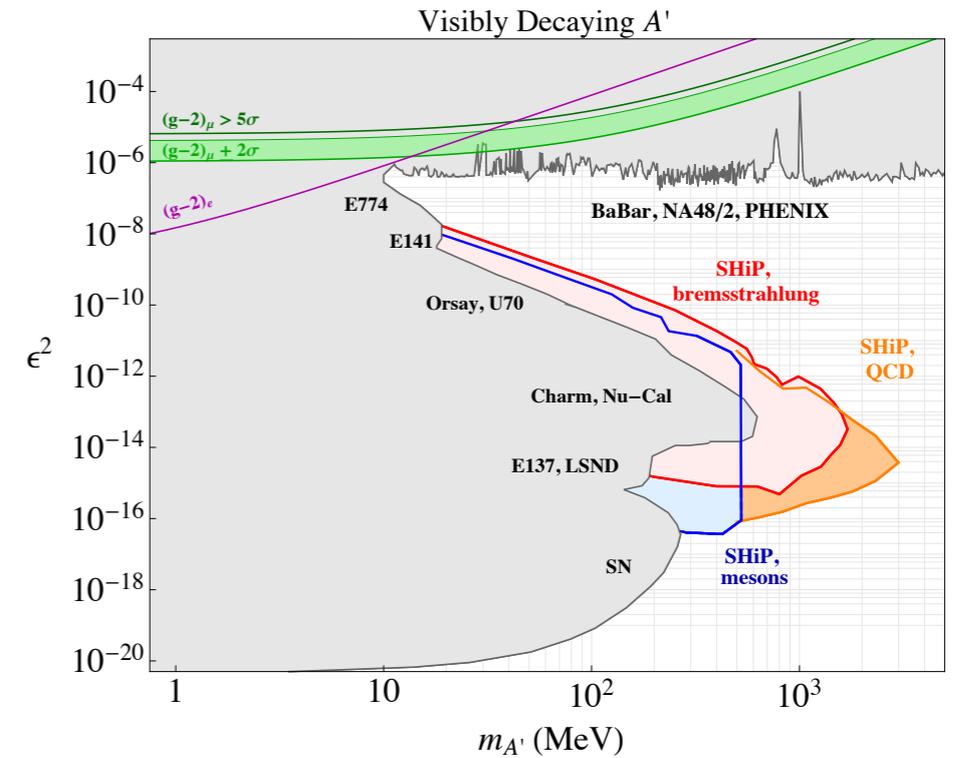
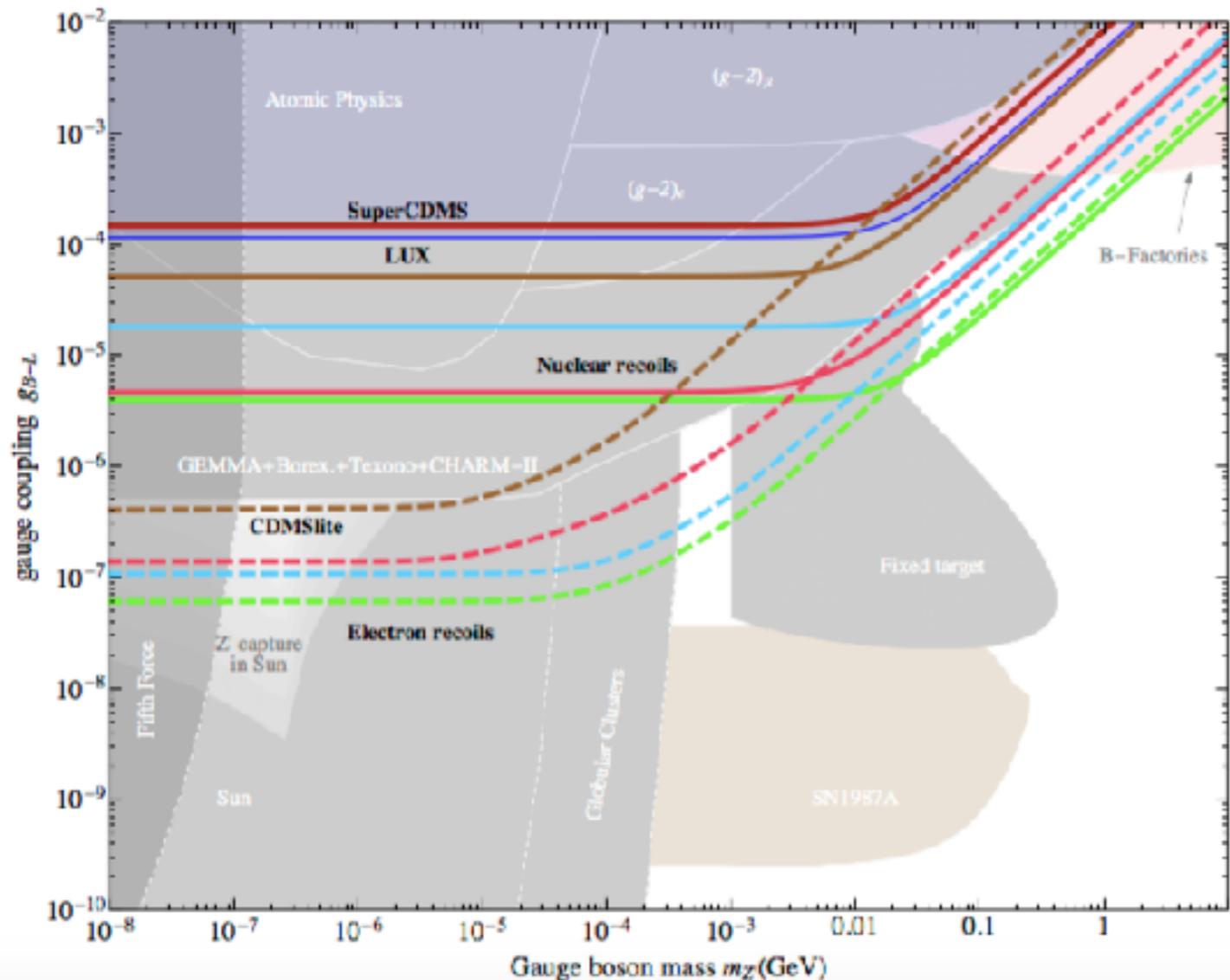
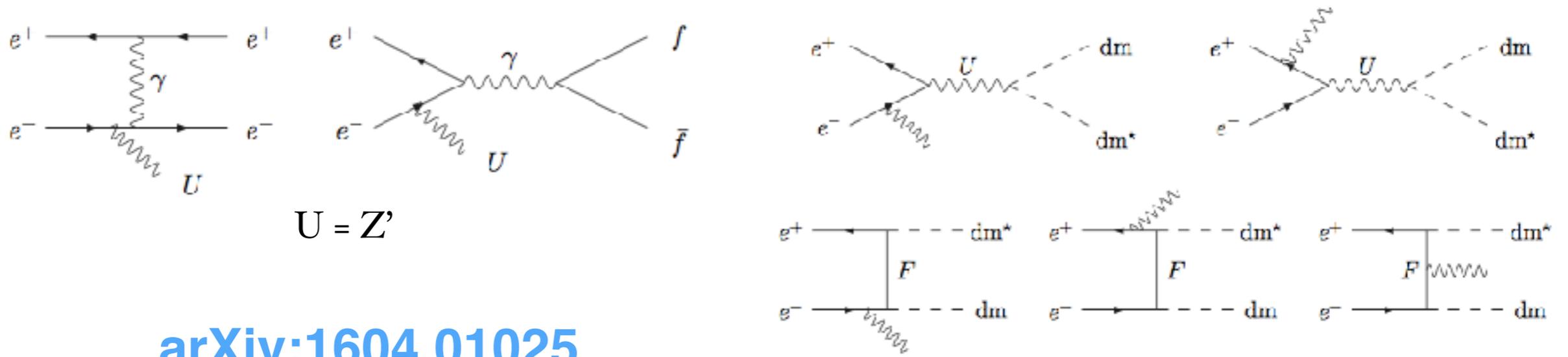
dark photons/dark Z'

**Light DM and mediator
but intensity is crucial!**

Low E. intensity frontier ...

hep-ph/0305261

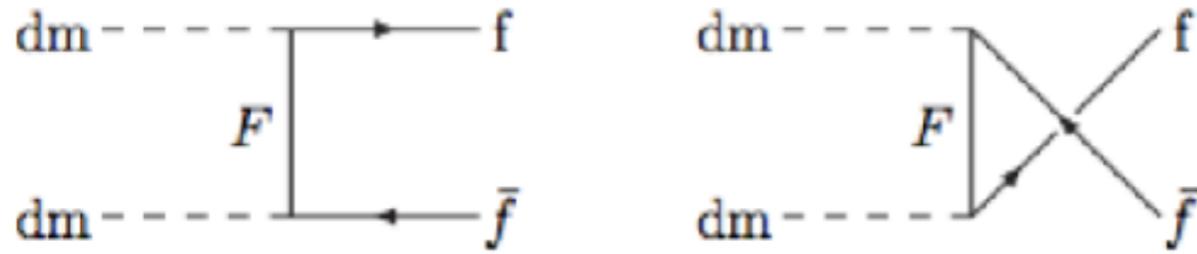
U-strahlung.



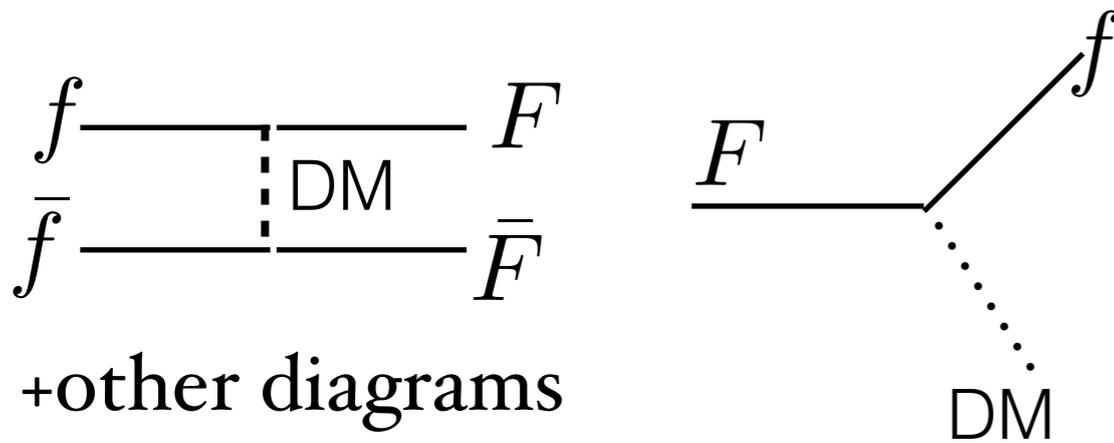
- 1406.3028, 1412.8378
- 1606.08849 1407.0993
- 1604.08206 1607.01789
- 1504.04855

High E. intensity frontier ...

09I2.5373



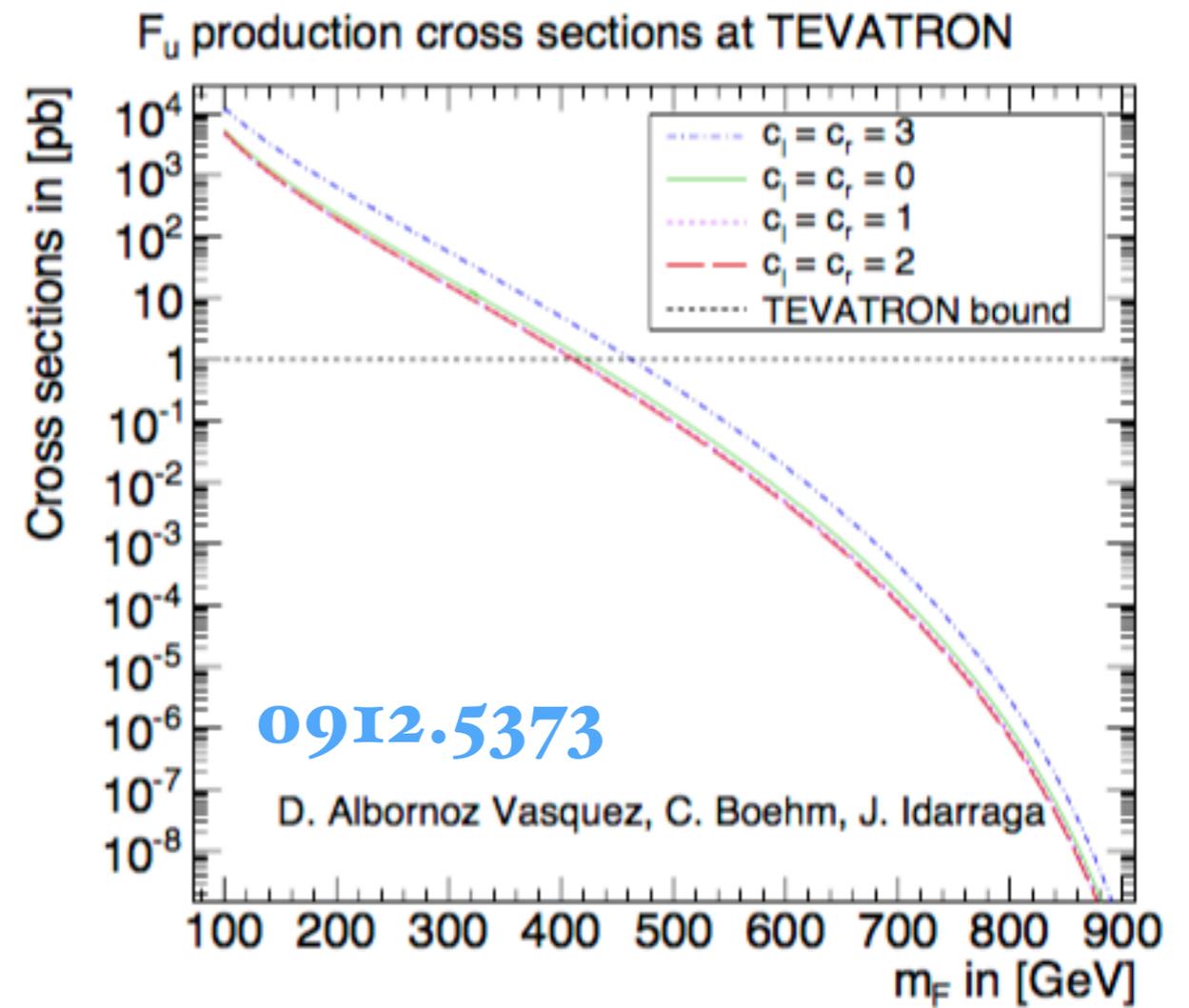
t-channel mediators



+other diagrams

The mediator is produced through the DM exchange

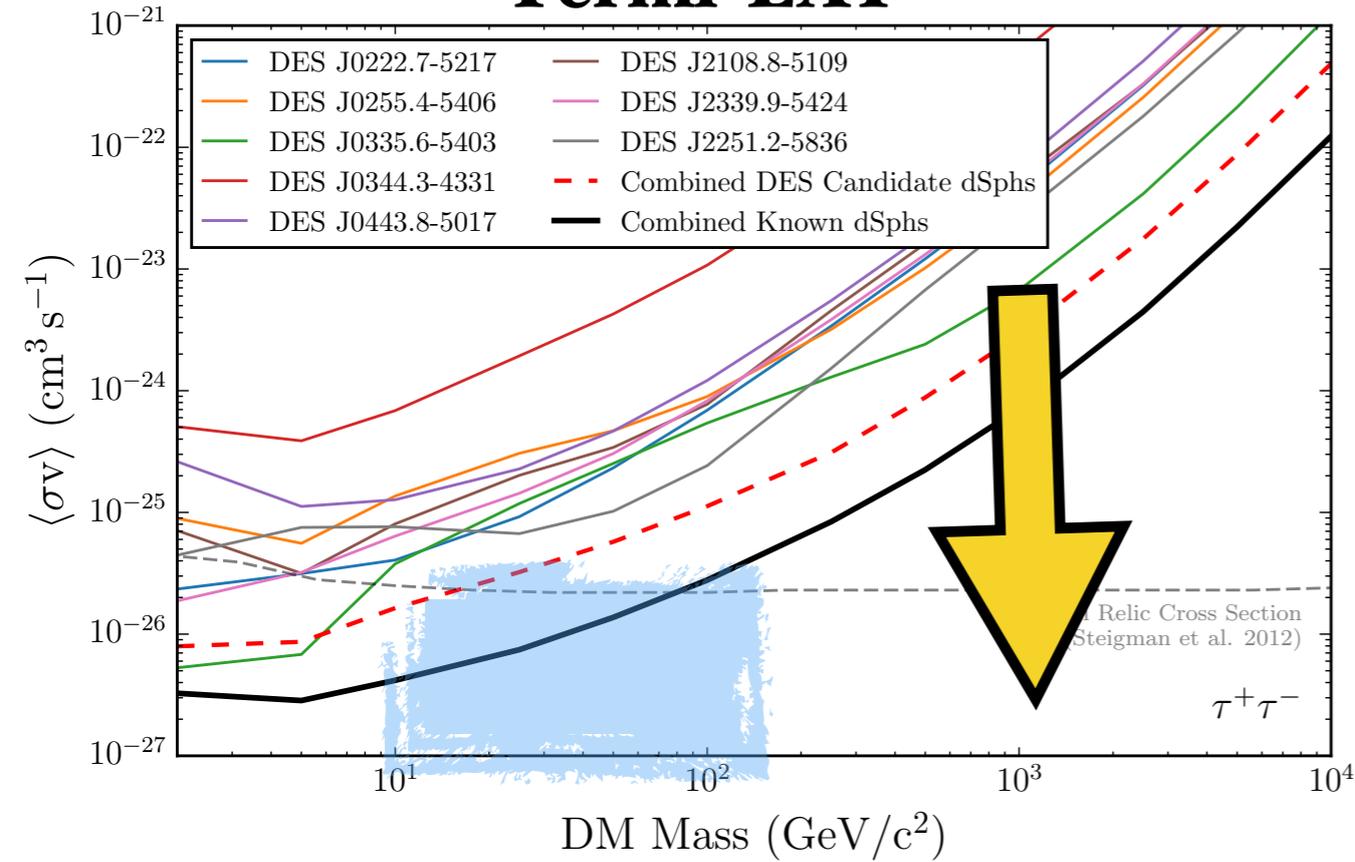
first example of DM searches
in colliders
using simplified models



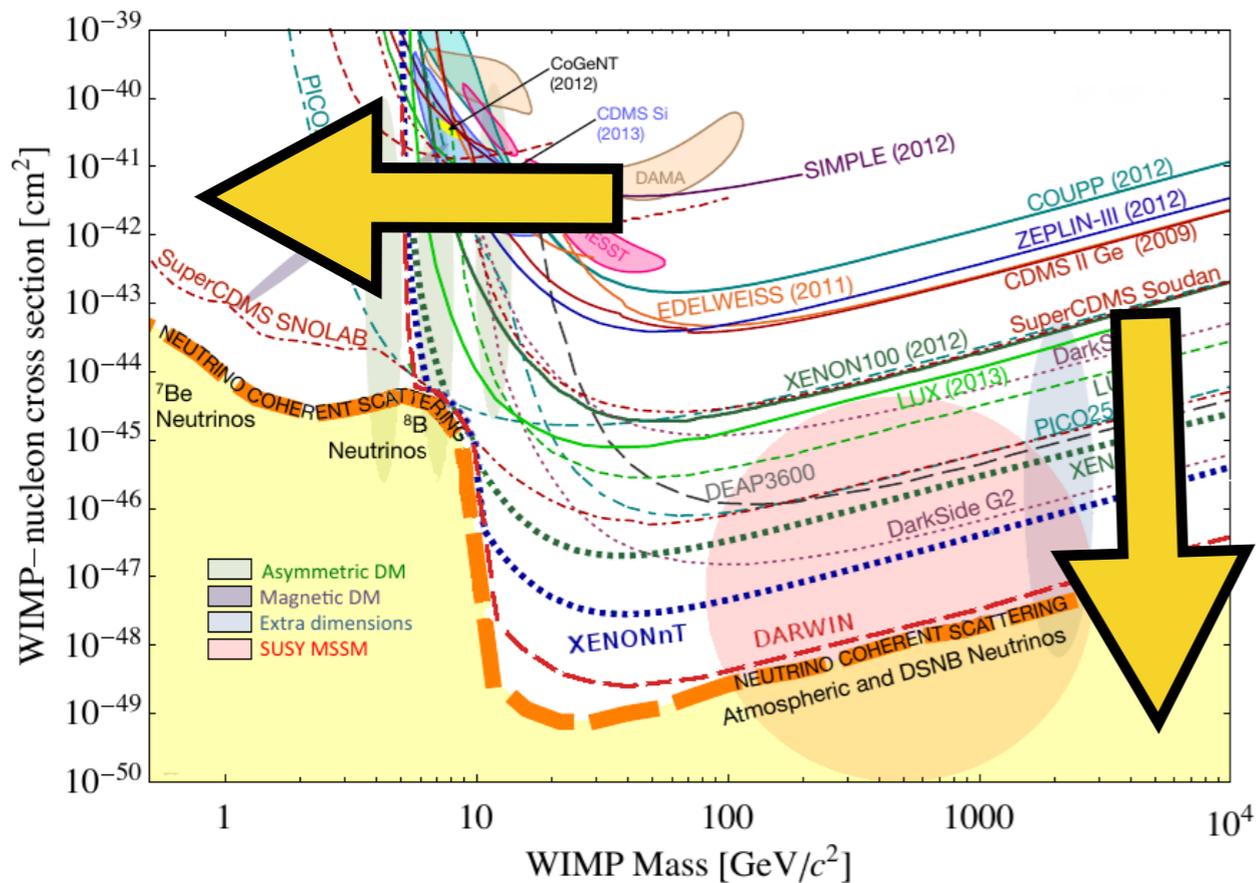
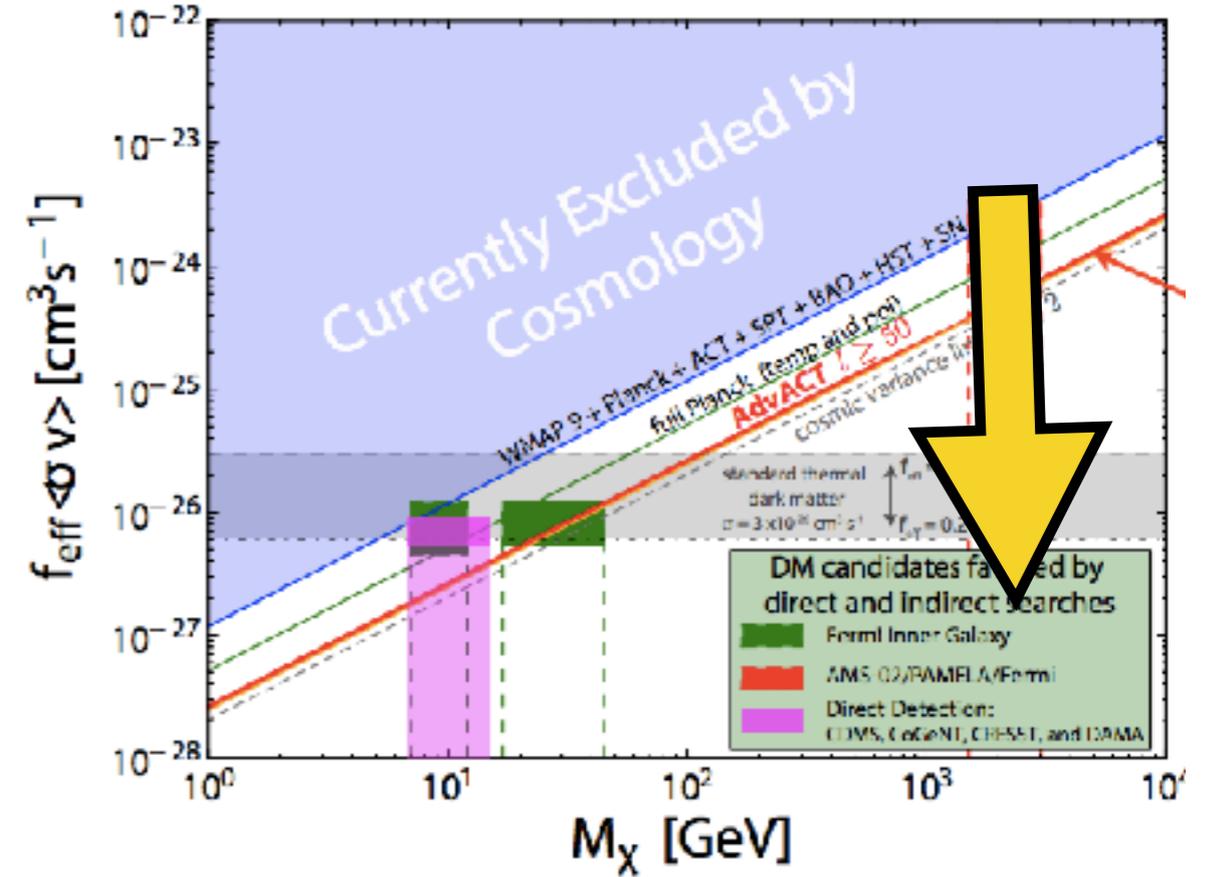
**Astrophysical/Astroparticle
observations
are going the same direction...**

“Intensity” & energy frontiers in the DM world

Fermi-LAT



Planck



Madhavacheril, NS, Slatyer 2014, PRD, (1310.3815)

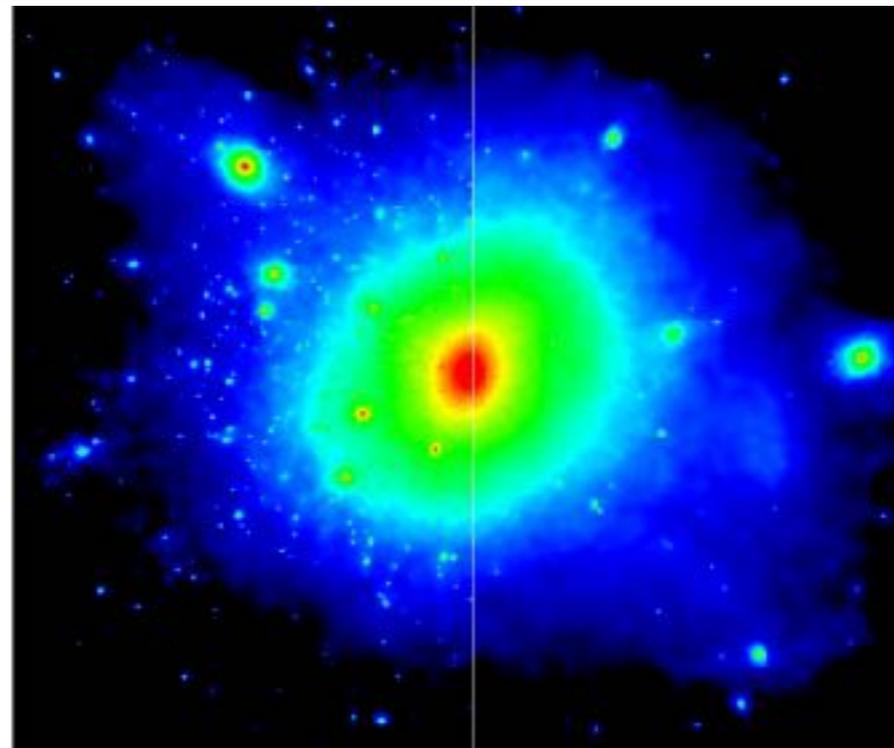
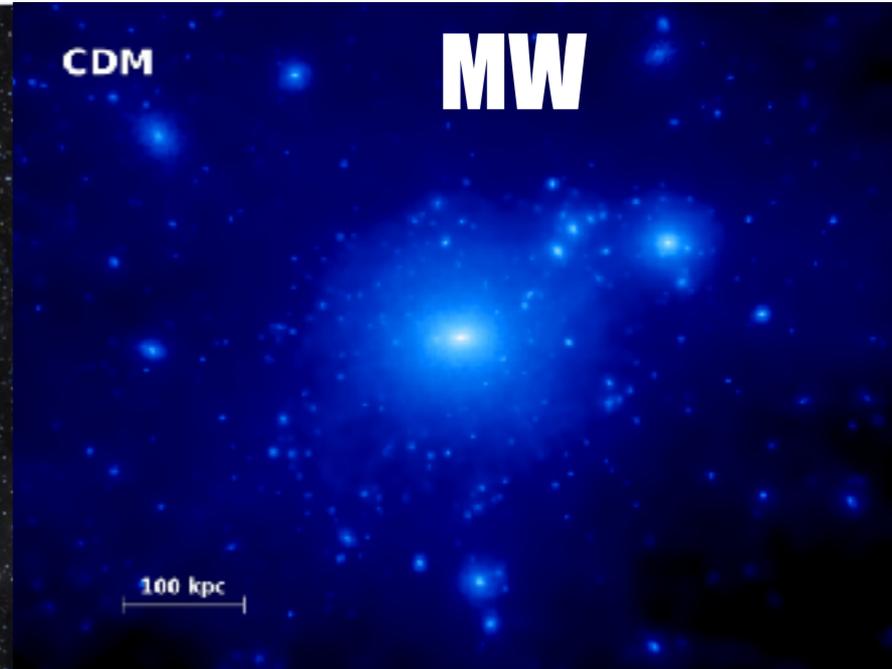
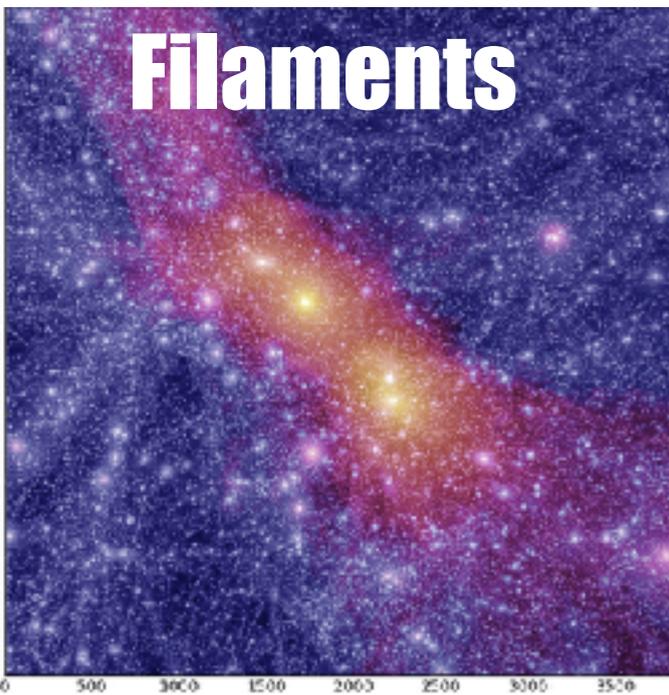
**Going lower/higher DM masses
Dig deeper**

And yet

DM is everywhere!

The landscape for new searches

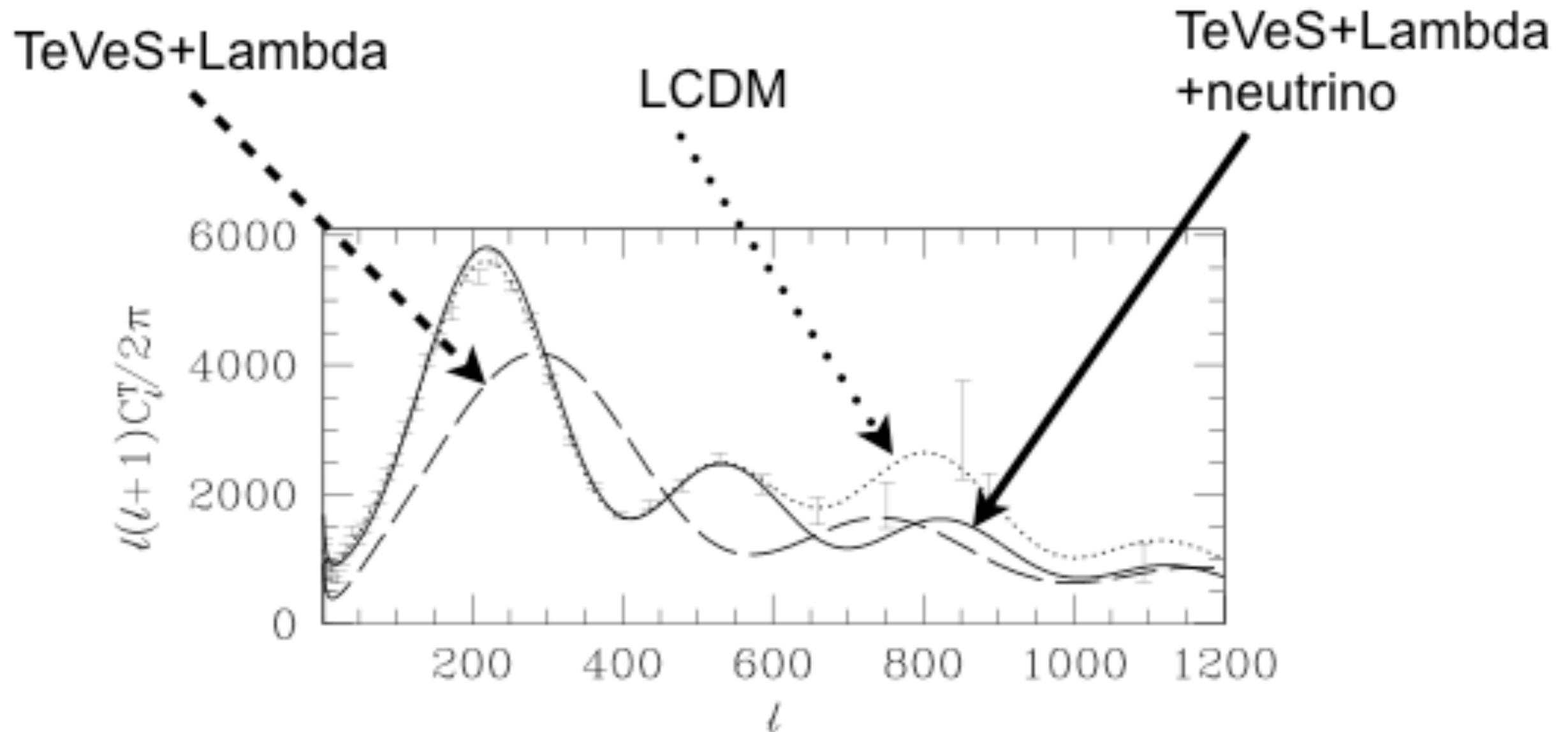
Dark Matter is (more or less) everywhere/at all scales



Once upon a time in a **DM-free** world

Mond/Bekenstein Bekenstein [astro-ph/0403694](#)

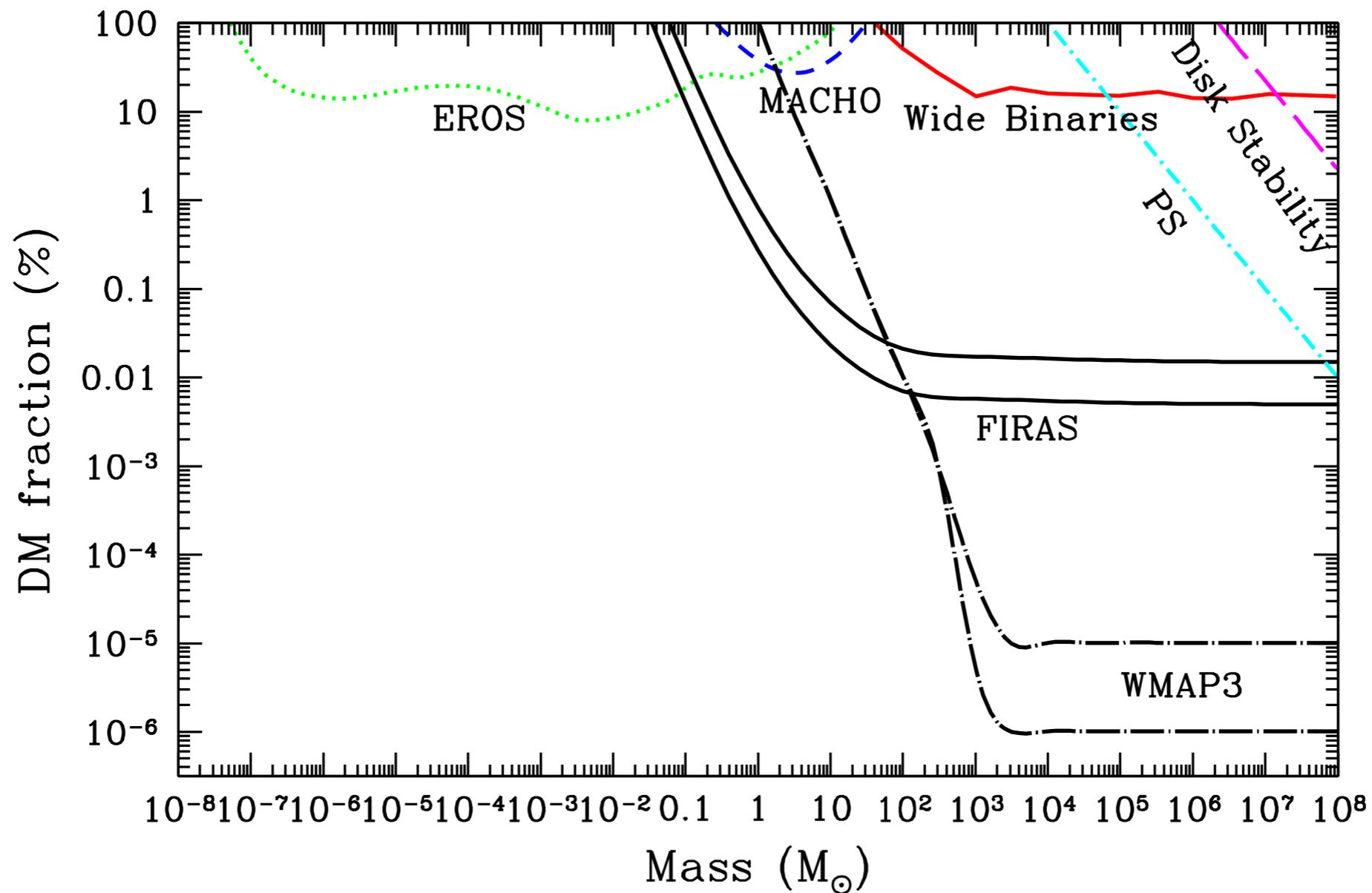
C. Skordis, D. Mota, P. Ferreira, C.Boehm : [astro-ph/0505519](#)



Impossible to explain Planck 2015!

Once upon a time in a **DM-free** world

Black Holes 0709.0524v1



A fraction of PBH is possible but one still needs a collisionless fluid.

Back to basics

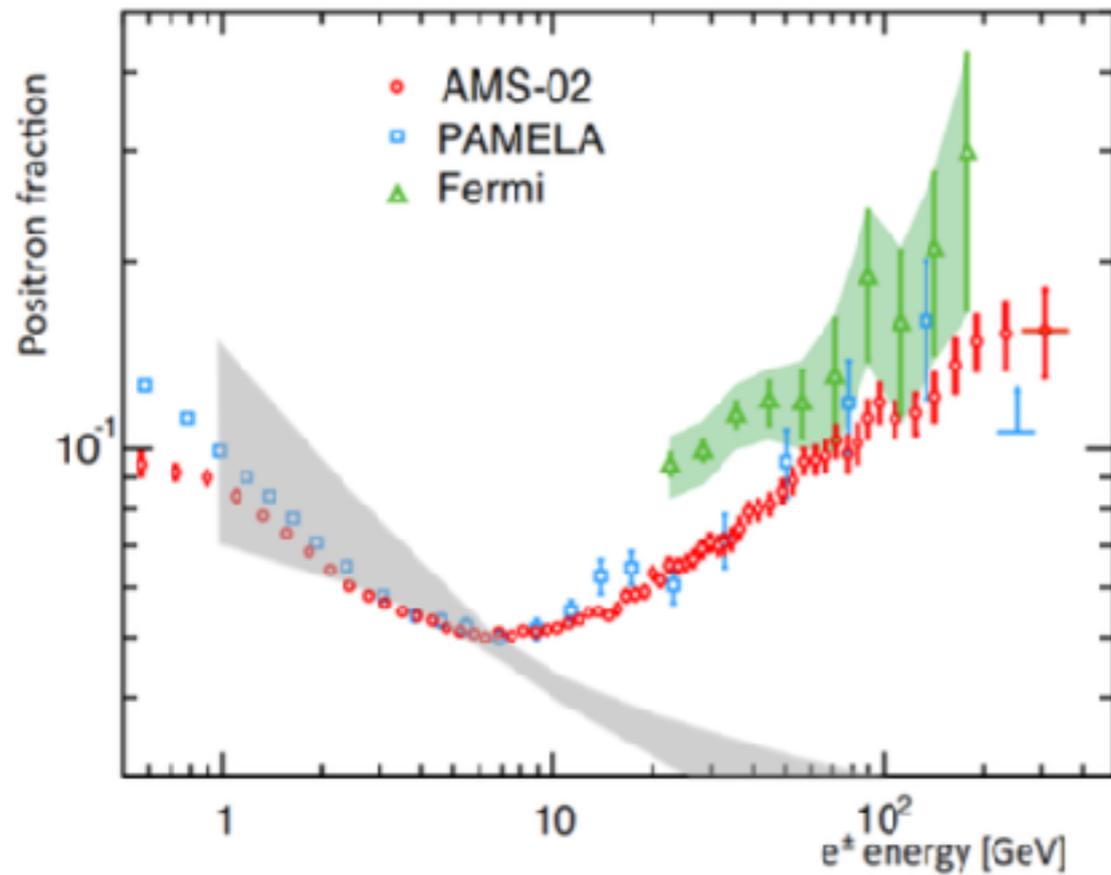
**The real proof of DM comes from
astrophysics/cosmology**

natural colliders: AGNs > 10 TeV

but many anomalies and DM is always a good fit!

AMS 02

perhaps not so DM after all...



1702.08436

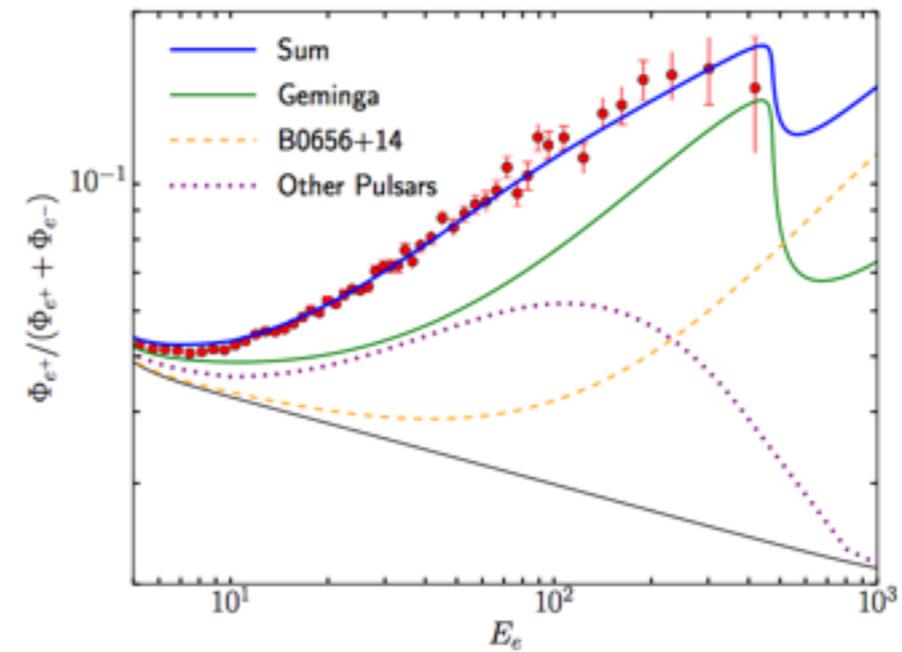
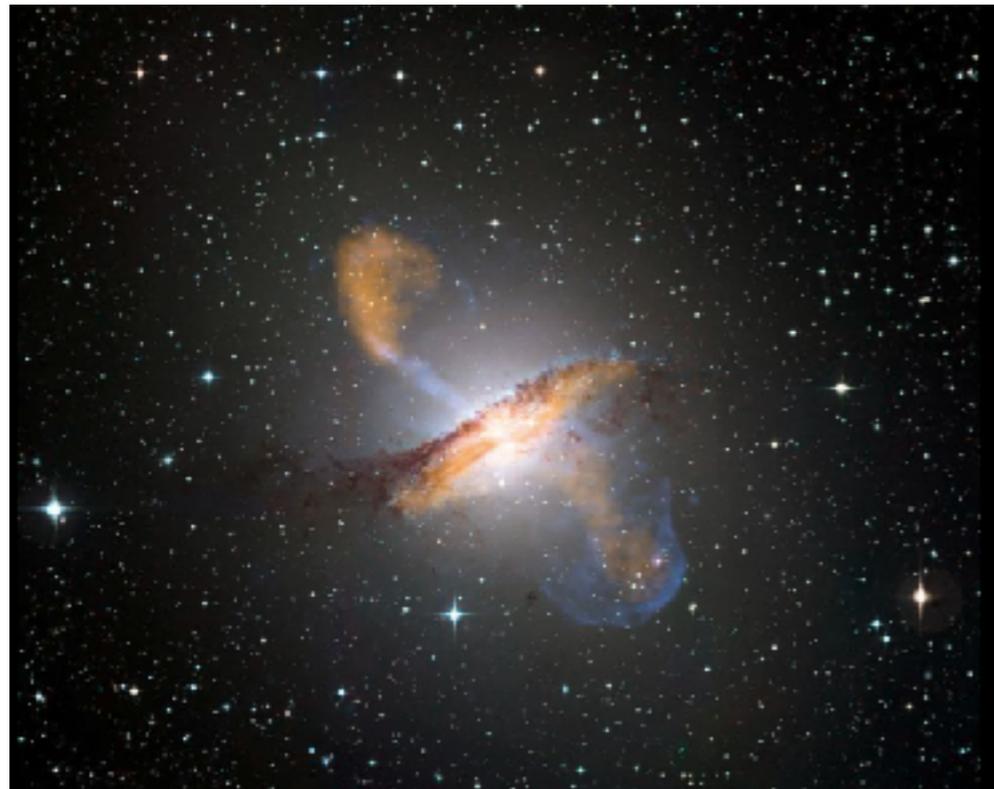


Figure 6. As in Fig. 5, but showing contributions from Geminga, B0656+14, and from all pulsars more than 0.5 kpc from the Solar System. For each source, we adopted $\alpha = 1.9$, $E_c = 49$ TeV, $v_c = 554$ km/s $\times (r_{\text{region}}/10 \text{ pc})$, and normalized their contributions with $\tau = 4.3 \times 10^3$ years, adopting a total birth rate of two pulsars per century in the Milky Way. While we expect many of these parameters to vary from pulsar-to-pulsar, making a detailed prediction of this kind difficult and possibly unreliable, this calculation provides significant support for the conclusion that a sizable fraction of the observed positron excess originates from pulsars.

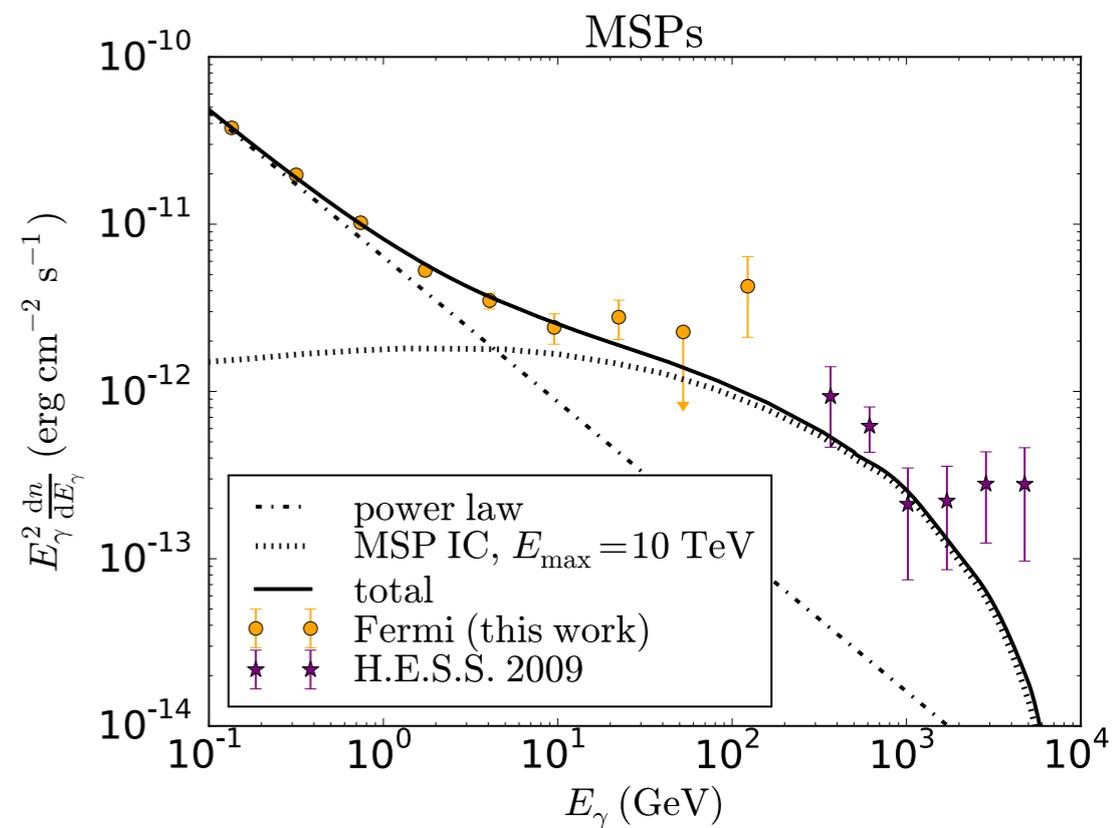
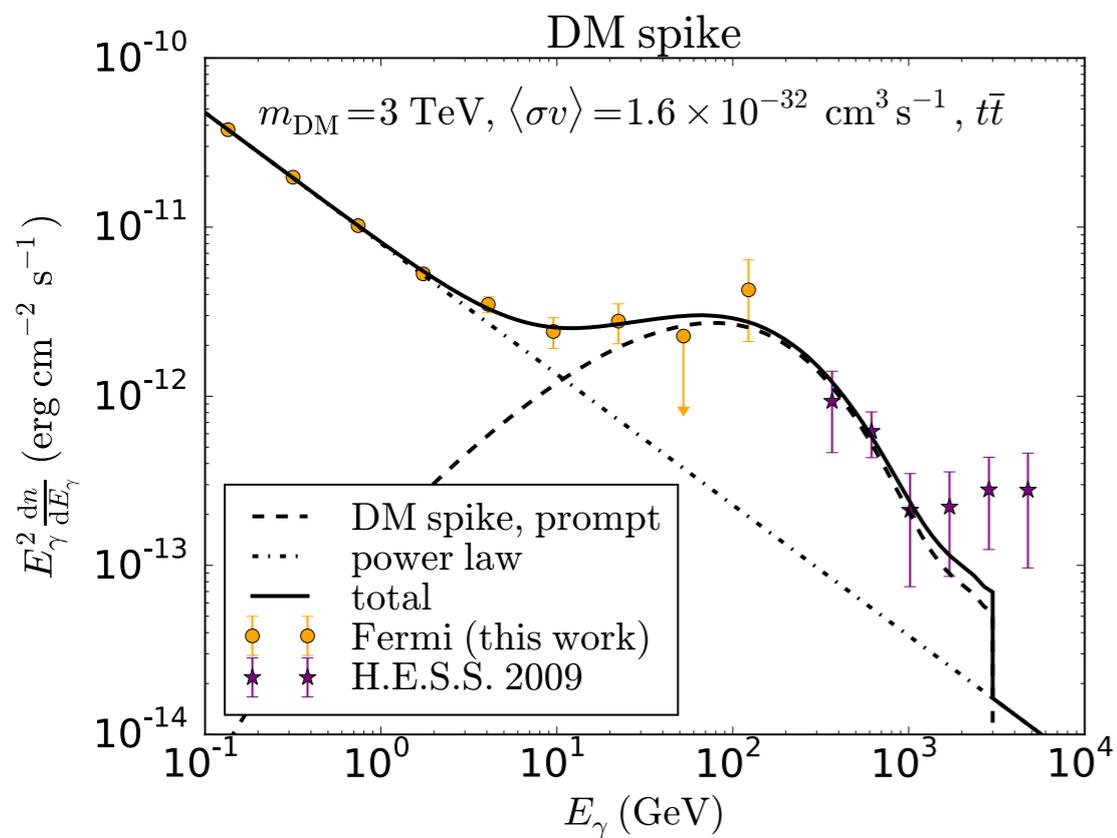
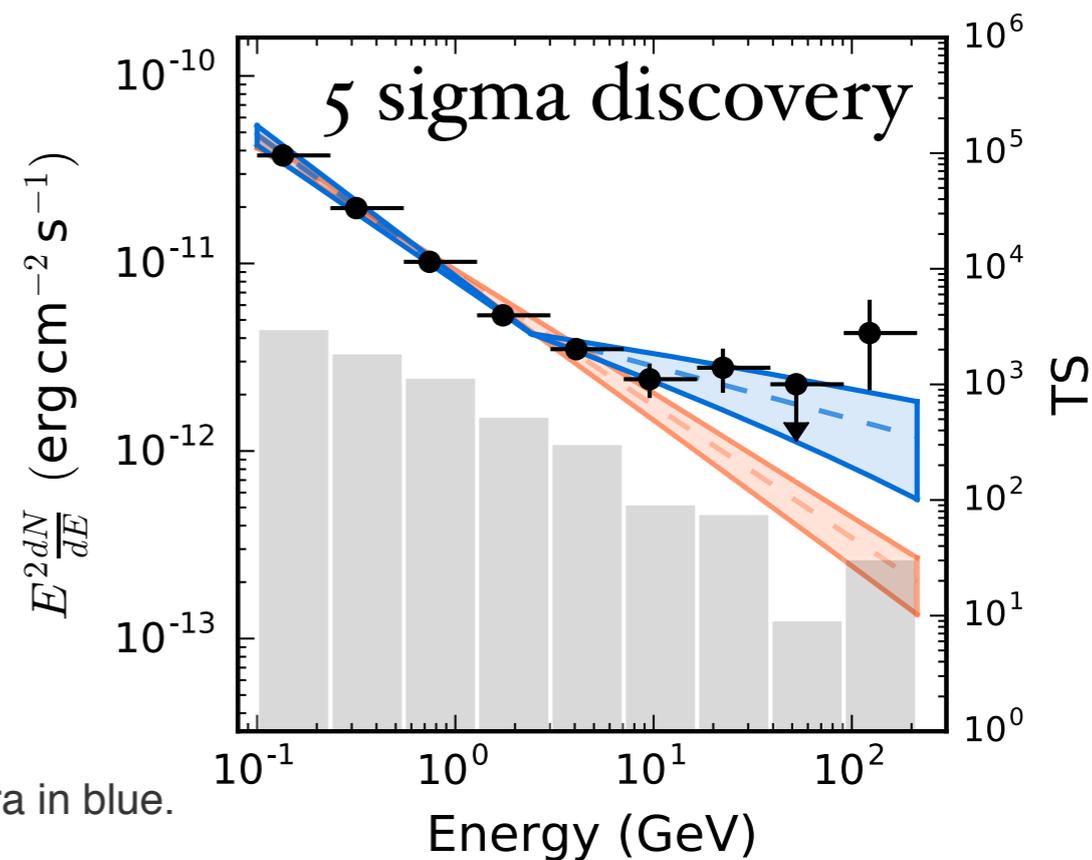
3-5 Mpc away

Centaurus A

1603.05469



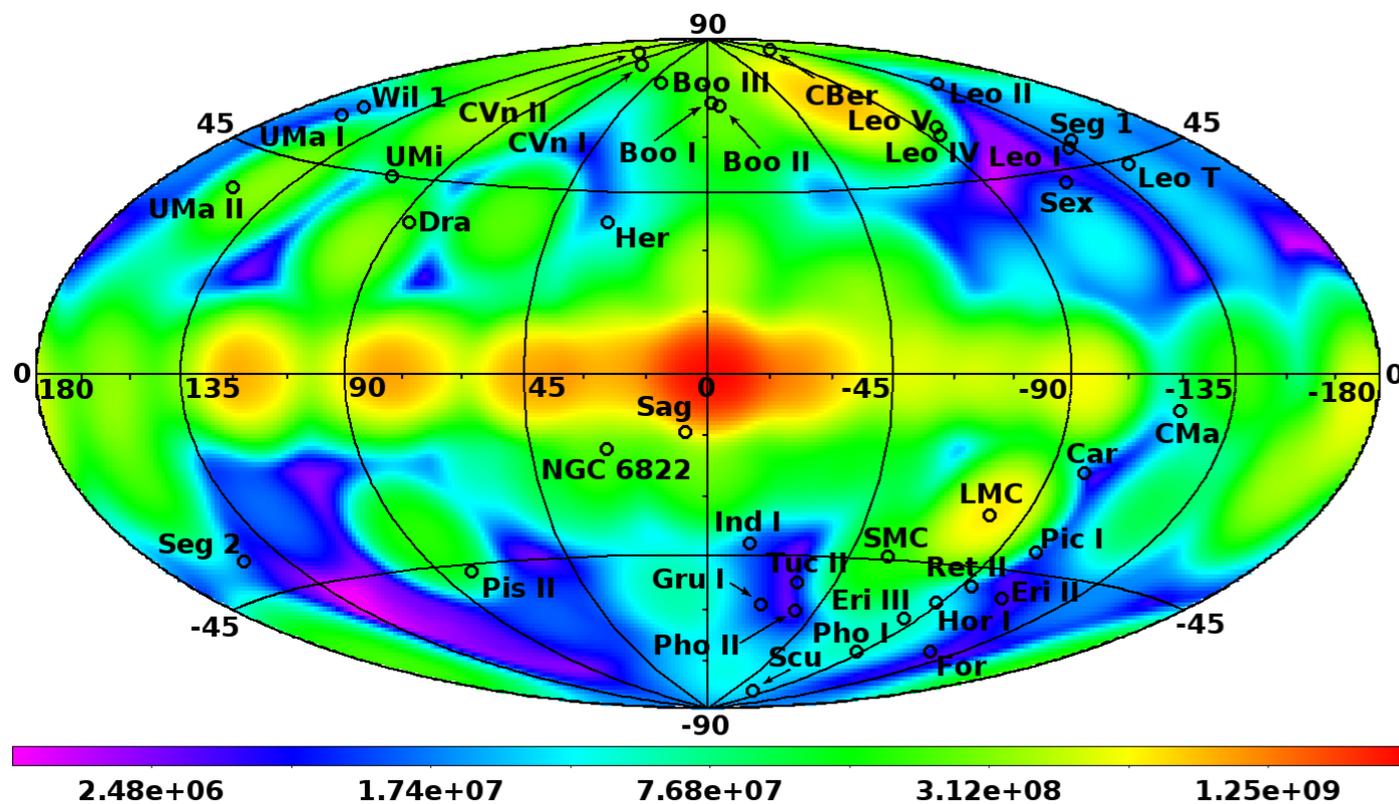
870-micron data from LABOCA on APEX in orange. X-ray data from Chandra in blue. Visible light data from the Wide Field Imager (WFI) in "true colour".



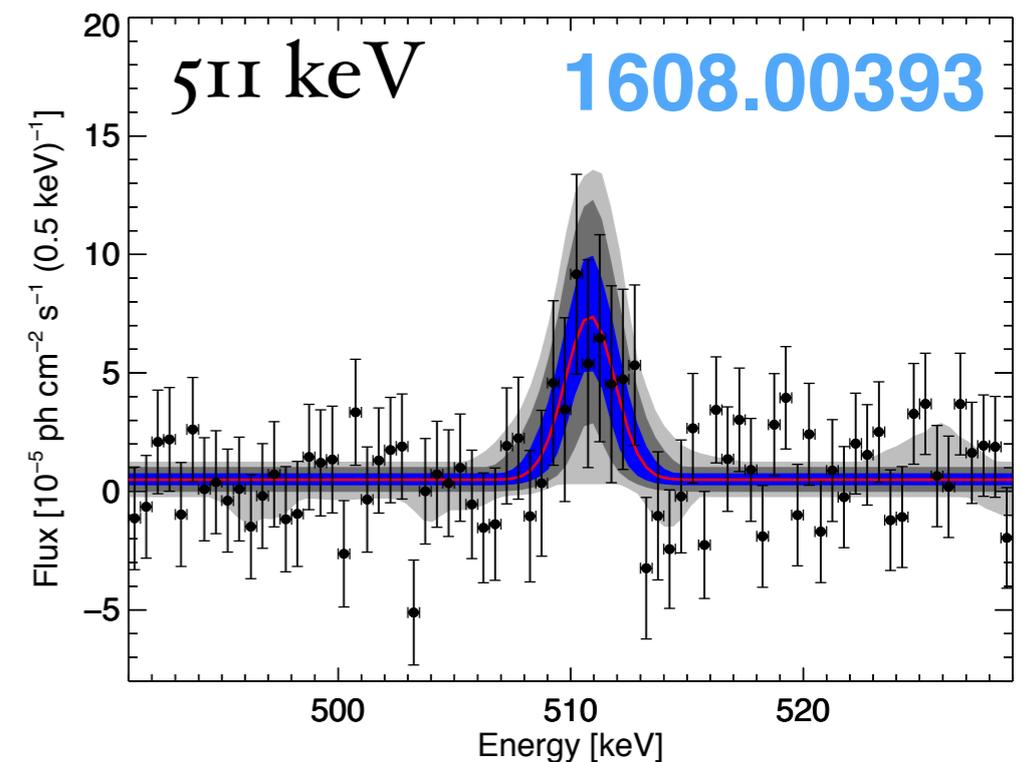
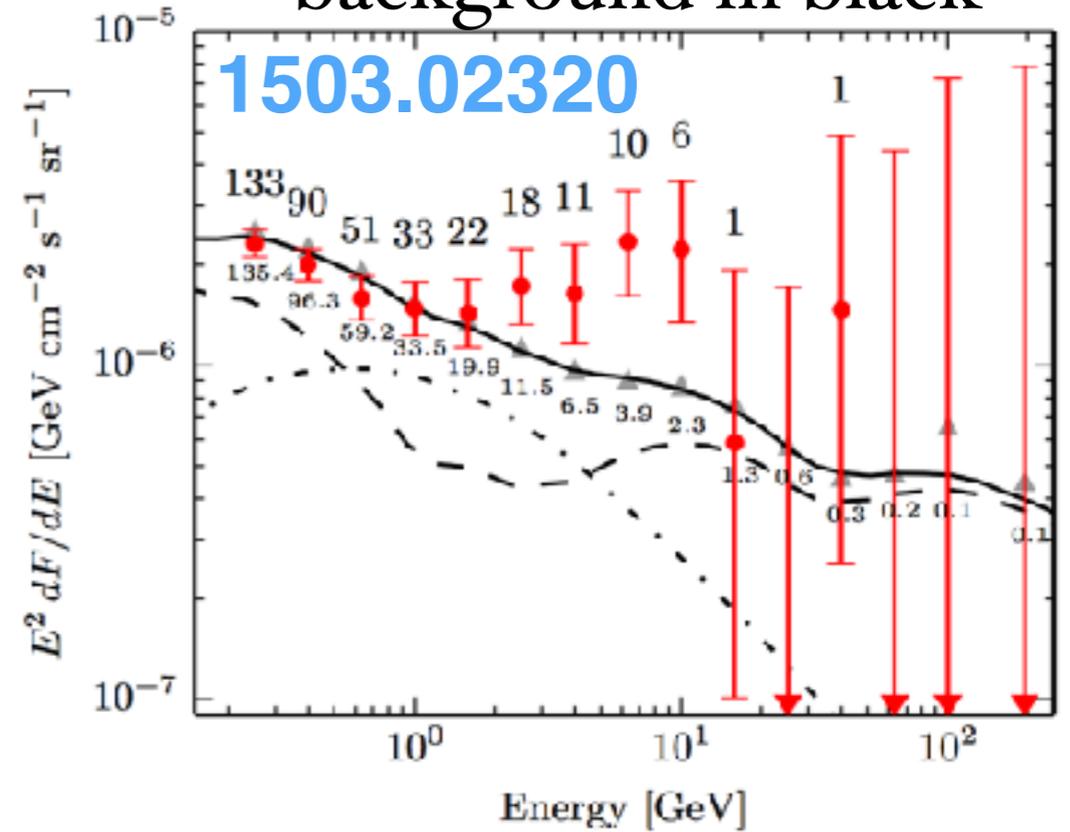
Reticulum II (?)

Excess of gamma-ray
in the GeV range but also
at 511 keV...

Integral/SPI exposure map
Courtesy: T. Siegert



background in black



**Can we discriminate
DM from astro sources?**

Circular polarisation

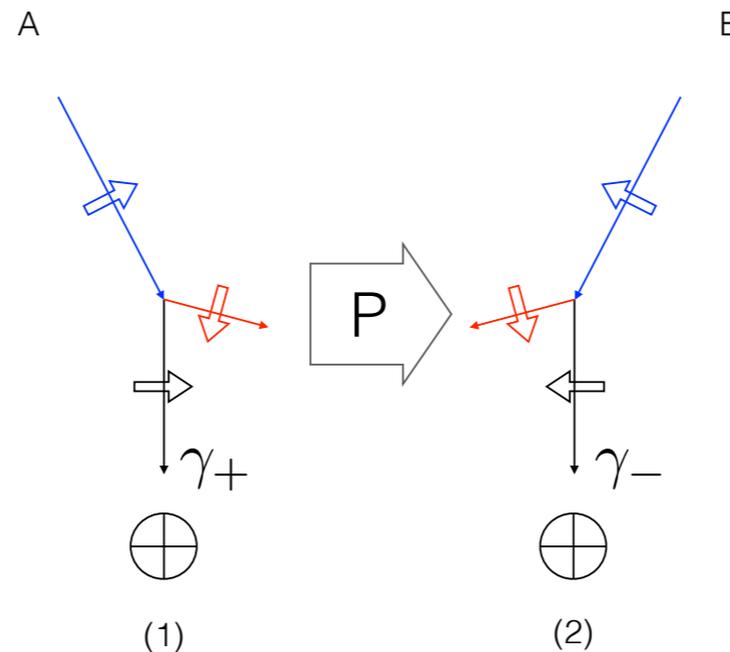
$$\epsilon_1^\mu(k) = (0, 1, 0, 0)$$

$$\epsilon_2^\mu(k) = (0, 0, \frac{k_z}{|k_z|}, 0).$$

$$\epsilon_\pm^\mu(k) = \frac{1}{\sqrt{2}} (\mp \epsilon_1^\mu - i \epsilon_2^\mu)$$

At process level: needs an excess of + over - or vice versa

[arXiv:1701.02754](https://arxiv.org/abs/1701.02754)



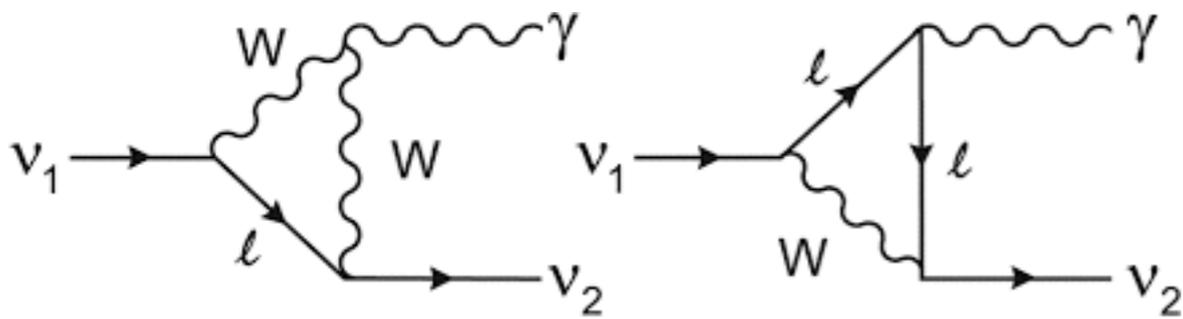
**If P violation,
then one state may
be dominant.**

**Processes with
gamma_5 are
promising**

In the Universe: needs P and CP violation

Asymmetry in number densities

Check with neutrino decay



$$\mathcal{A}(\nu_i \rightarrow \nu_j \gamma) = -i\sigma^{\mu\nu} \epsilon_{\mu} q_{\nu} (\mu_{ij} + i\varepsilon_{ij} \gamma_5)$$

$$\frac{\mu_{ij}}{i\varepsilon_{ij}} = -\frac{3eG_f}{32\sqrt{2}\pi^2} (m_{\nu_i} \pm m_{\nu_j}) \sum_{l=e,\mu,\tau} U_{il}^{\dagger} U_{lj} \left(-2 + \frac{m_l^2}{m_W^2} \right).$$

$$\Gamma(\nu_i \rightarrow \nu_j \gamma) = \frac{1}{8\pi} \left(\frac{m_{\nu_i}^2 - m_{\nu_j}^2}{m_{\nu_i}} \right)^3 \left(|\mu_{ij}|^2 + |\varepsilon_{ij}|^2 \right)$$

$$\frac{\Gamma(\nu_i \rightarrow \nu_j \gamma_+)}{\Gamma(\nu_i \rightarrow \nu_j \gamma)} = \frac{|\mu_{ij} - i\varepsilon_{ij}|^2}{2 \left(|\mu_{ij}|^2 + |\varepsilon_{ij}|^2 \right)} = \frac{m_{\nu_i}^2}{m_{\nu_j}^2 + m_{\nu_i}^2}$$

$$\frac{\Gamma(\nu_i \rightarrow \nu_j \gamma_-)}{\Gamma(\nu_i \rightarrow \nu_j \gamma)} = \frac{|\mu_{ij} + i\varepsilon_{ij}|^2}{2 \left(|\mu_{ij}|^2 + |\varepsilon_{ij}|^2 \right)} = \frac{m_{\nu_j}^2}{m_{\nu_j}^2 + m_{\nu_i}^2}$$

$$\begin{aligned} m_{\nu_1} &= 8.5 \times 10^{-4} \text{ eV} \\ m_{\nu_2} &= 8.7 \times 10^{-3} \text{ eV} \\ m_{\nu_3} &= 5.016 \times 10^{-2} \text{ eV} \end{aligned}$$



$$\Gamma(\nu_3 \rightarrow \nu_1 \gamma_+) / \Gamma(\nu_3 \rightarrow \nu_1 \gamma) = 0.99$$

$$\Gamma(\nu_3 \rightarrow \nu_2 \gamma_+) / \Gamma(\nu_3 \rightarrow \nu_2 \gamma) = 0.96$$

circular polarisation! Almost 100%

SM Neutrino decay



HST composite image of a supernova explosion SN 2014J in M82, at a distance of approximately 11.5 million light-years from Earth. Using a ground-based telescope, the explosion was discovered on Jan 21, 2014. The photograph was taken on January 31, as the SN approached its peak brightness

$$\phi_{\nu_i \rightarrow \nu_j \gamma_{\pm}} = \frac{\dot{n}_{\nu,0}}{4\pi D^2} f_i \left(1 - e^{-\frac{m_{\nu_i}}{E_{\nu}} \frac{D}{c} \Gamma_{\nu_i \rightarrow \nu_j \gamma_{\pm}}} \right)$$

$$\phi_{\nu_i \rightarrow \nu_j \gamma_{\pm}} \simeq \phi_{\nu,0} \frac{m_{\nu_i}}{E_{\nu}} \frac{D}{c} f_i \Gamma_{\nu_i \rightarrow \nu_j \gamma_{\pm}}$$

f_i = fraction of neutrino "i"

circular polarisation!

Net circular polarisation

$$\phi_{\gamma, pol}^{\nu_i} = (\xi - \bar{\xi}) (\phi_{\nu_i \rightarrow \nu_j \gamma_+} - \phi_{\nu_i \rightarrow \nu_j \gamma_-})$$

$$\phi_{\gamma, pol}^{\nu_i} = \phi_{\nu,0} \frac{m_{\nu_i}}{E_{\nu}} \frac{D}{c} f_i \Delta_{CP} \Gamma_{\nu_i, pol}$$

$$\phi_{\gamma, pol}^{\nu_i} \simeq 10^{16} f_i \Delta_{CP} \left(\frac{m_{\nu_i}}{\text{eV}} \right) \left(\frac{\Gamma_{\nu_i, pol}}{\text{s}^{-1}} \right) \text{ph cm}^{-2} \text{s}^{-1}.$$

$$\Delta_{CP} = (\xi - \bar{\xi})$$

$$\xi = n_{\nu_i} / (n_{\nu_i} + n_{\bar{\nu}_i})$$

$$\bar{\xi} = n_{\bar{\nu}_i} / (n_{\nu_i} + n_{\bar{\nu}_i})$$

SM & BSM neutrino decay

$$\phi_{\gamma,pol}^{\nu_i} \simeq 10^{16} f_i \Delta_{CP} \left(\frac{m_{\nu_i}}{\text{eV}} \right) \left(\frac{\Gamma_{\nu_i,pol}}{\text{s}^{-1}} \right) \text{ph cm}^{-2} \text{s}^{-1}.$$

for a 10 kpc SN emitting a MeV neutrino flux $\phi_{\nu,0} \sim 10^{10} \text{ cm}^{-2} \text{ s}^{-1}$

$$\Gamma_{\nu_i,pol} = (\Gamma_{\nu_i \rightarrow \nu_j \gamma_+} - \Gamma_{\nu_i \rightarrow \nu_j \gamma_-}) \quad \longrightarrow \quad \Gamma_{\nu_i,pol} = \Gamma_{\nu_i \rightarrow \nu_j \gamma_+}$$

“Standard Model” (massive) neutrinos

$$\Gamma(\nu_3 \rightarrow \nu_2 \gamma) = 1.27 \times 10^{-51} \text{ s}^{-1}$$

So the corresponding flux is invisible ...
unless the source produces a huge amount of neutrinos

But if we allow for new Physics, we can expect fluxes of

If we request the decay rate to be $< 10^{-17}/\text{s}$

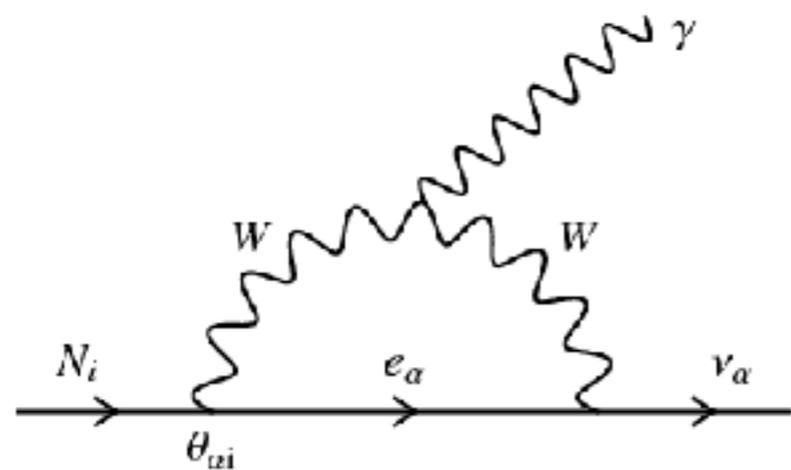
$$\phi_{\gamma,pol}^{\nu_i} < 2 \times 10^{-2} \left(\frac{m_{\nu_i}}{\text{eV}} \right) \text{ph cm}^{-2} \text{s}^{-1}$$

Back to DM: sterile neutrino decay

(was the object of 1610.04532 though no estimates)

Sterile neutrinos

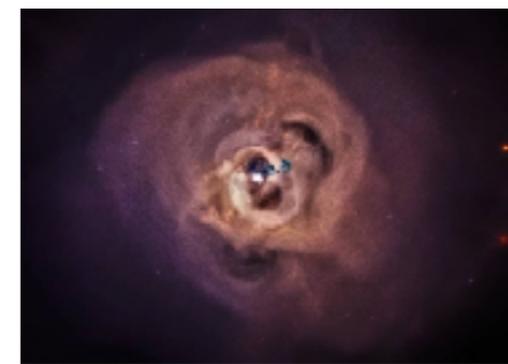
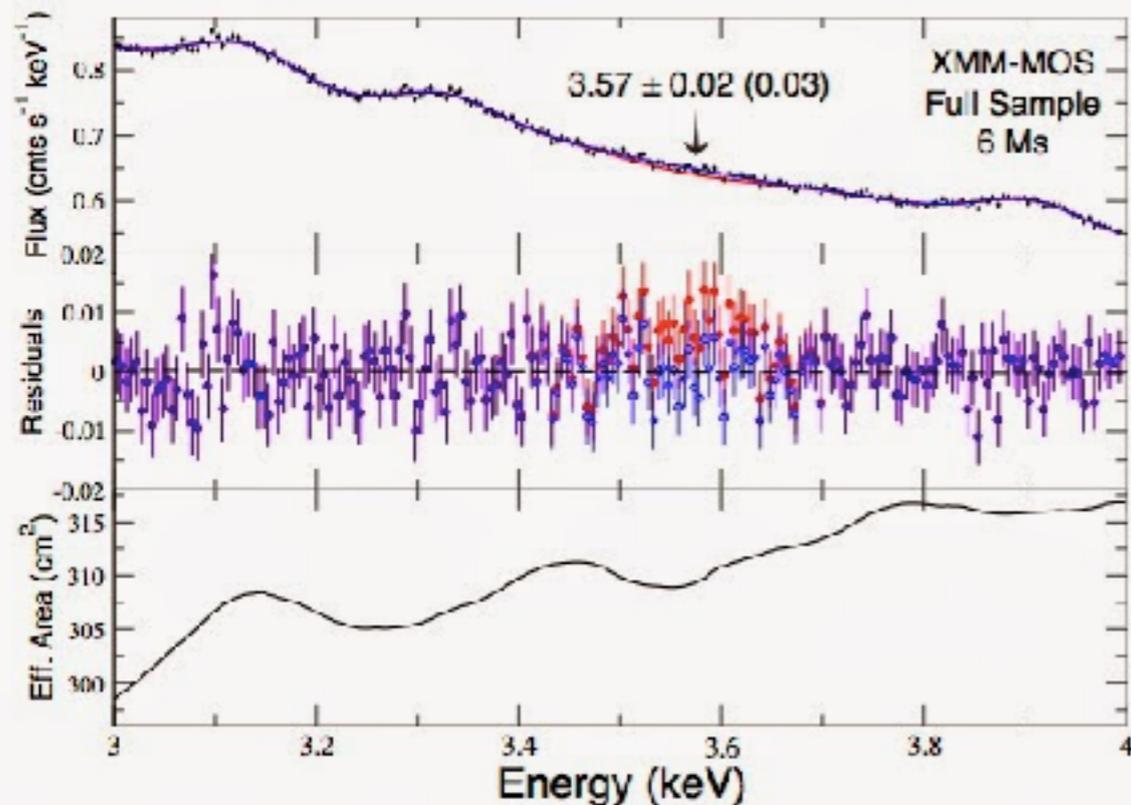
(Gravitinos etc)



$$E_\gamma = \frac{m_{sterile} \nu}{2}$$

7 keV line

seen in Perseus cluster (+ in few others)
but Hitomi did not see the line

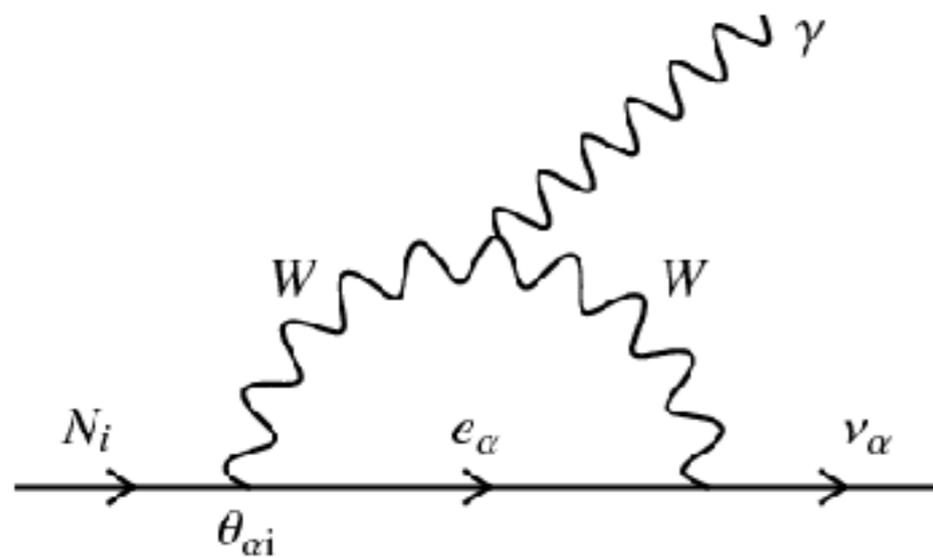


3.5 keV ~ Potassium (Bananas ...)

arxiv.org/pdf/1409.4143v1
arxiv.org/abs/1408.1699v1

Back to DM: sterile neutrino decay

(was the object of 1610.04532 though no estimates)



$$\Gamma_{\nu_s} = 1.38 \times 10^{-30} \left(\frac{\sin^2(2\theta)}{10^{-8}} \right) \left(\frac{m_s}{\text{keV}} \right)^5 \text{ s}^{-1}$$

$$\phi_{pol} = \Delta_{CP}^{\nu_s} \Gamma_{\nu_s} \int_{FoV} d\Omega \int ds(r) \frac{\rho_{DM}(r)}{m_{DM}} (1+z)^{-1}$$

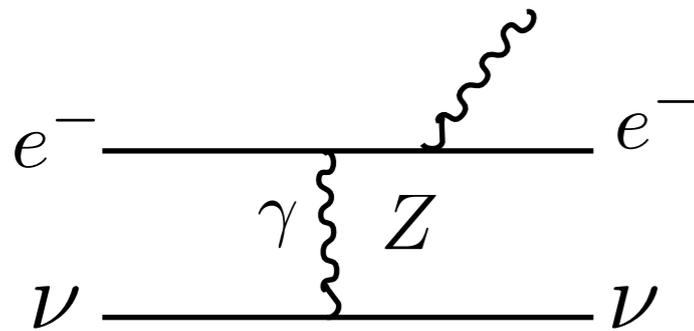
100% circular polarisation

For Perseus and a 7 keV sterile neutrino $z=0.0179$, $r_s = 164/h$ kpc, $d=70$ Mpc

$$\Gamma_{\nu_s} \simeq 1.63 \cdot 10^{-28} \text{ s}^{-1} \quad \sin^2(2\theta) \sim 7 \times 10^{-11}$$

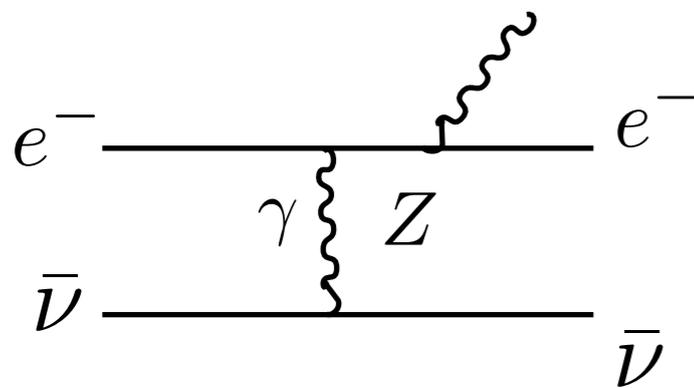
$$\phi_{pol, 3.5keV}^{Perseus} \sim 10^{-6} \Delta_{CP}^{\nu_s} \text{ ph cm}^{-2} \text{ s}^{-1}$$

Application 3: neutrinos scattering off electrons



same with e+ and anti-nu
(muon neutrino here)

\sqrt{S}	cross-section	γ_+ (%)	γ_- (%)	net pol (%)	$E_{\min}(\gamma)$
1 TeV	1.0 pb	43	57	-14.1 ± 1.2	100 GeV
1 TeV	3.1 pb	47	53	-6.3 ± 1.2	10 GeV
1 TeV	5.3 pb	48	52	-3.2 ± 1.2	1 GeV
100 GeV	0.35 pb	36	64	-28.5 ± 1.2	10 GeV
100 GeV	1.19 pb	44	56	-11.2 ± 1.2	1 GeV
10 GeV	0.0052 pb	34	66	-32.7 ± 1.2	1 GeV
10 GeV	0.0185 pb	44	56	-12.6 ± 1.2	0.1 GeV

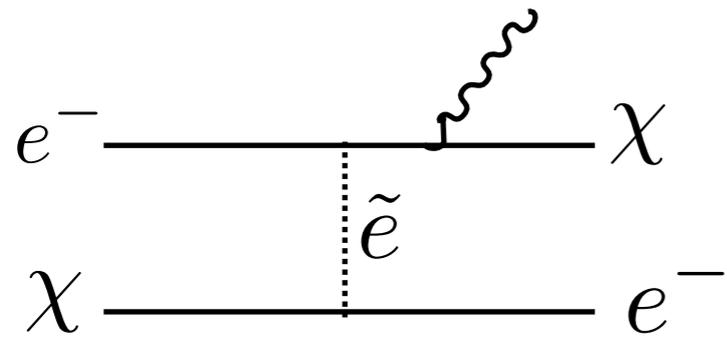


same with e+ and nu

\sqrt{S}	Cross section	γ_+ (%)	γ_- (%)	net pol (%)	$E_{\min}(\gamma)$
1 TeV	1.0 pb	54	46	8.8 ± 1.2	100 GeV
1 TeV	3.0 pb	52	48	4.2 ± 1.2	10 GeV
1 TeV	5.1 pb	51	49	2.6 ± 1.2	1 GeV
100 GeV	0.30 pb	55	45	9.8 ± 1.2	10 GeV
100 GeV	1.00 pb	52	48	4.2 ± 1.2	1 GeV
10 GeV	0.0042 pb	59	41	17.2 ± 1.2	1 GeV
10 GeV	0.015 pb	53	47	6.4 ± 1.2	0.1 GeV

we could probe the distribution of neutrinos near astro sources

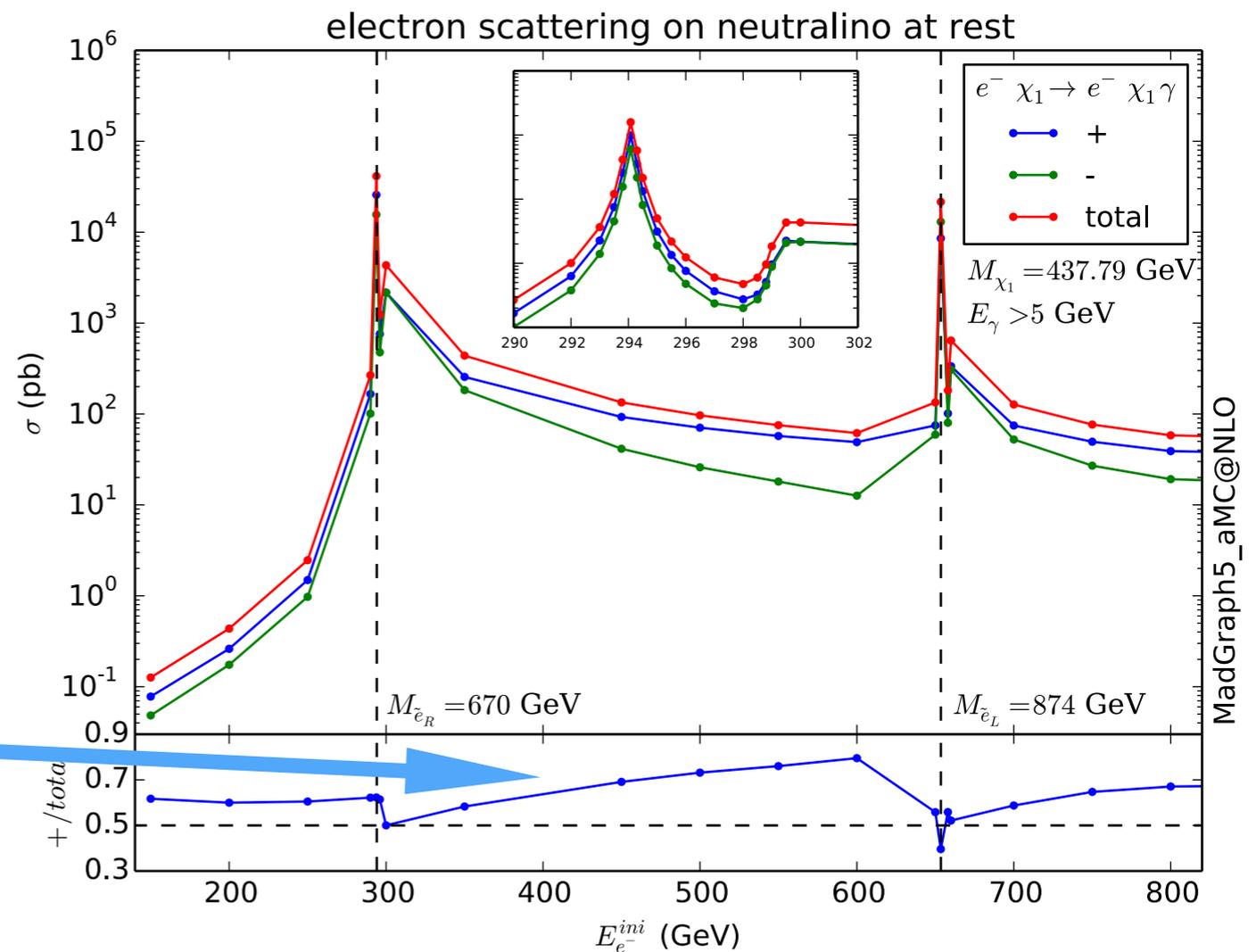
Application 4: dark matter scattering off electrons



selectrons are scalar particles but their couplings are sensitive to electron chirality

One also expects the generation of a circular polarisation

$$\begin{aligned} \text{sign}(\mu) &> 0 \\ m_0 &= 550 \text{ GeV} \\ m_{1/2} &= 1020 \text{ GeV} \\ \tan \beta &= 19.16 \\ A_0 &= -2878 \text{ GeV} \end{aligned}$$

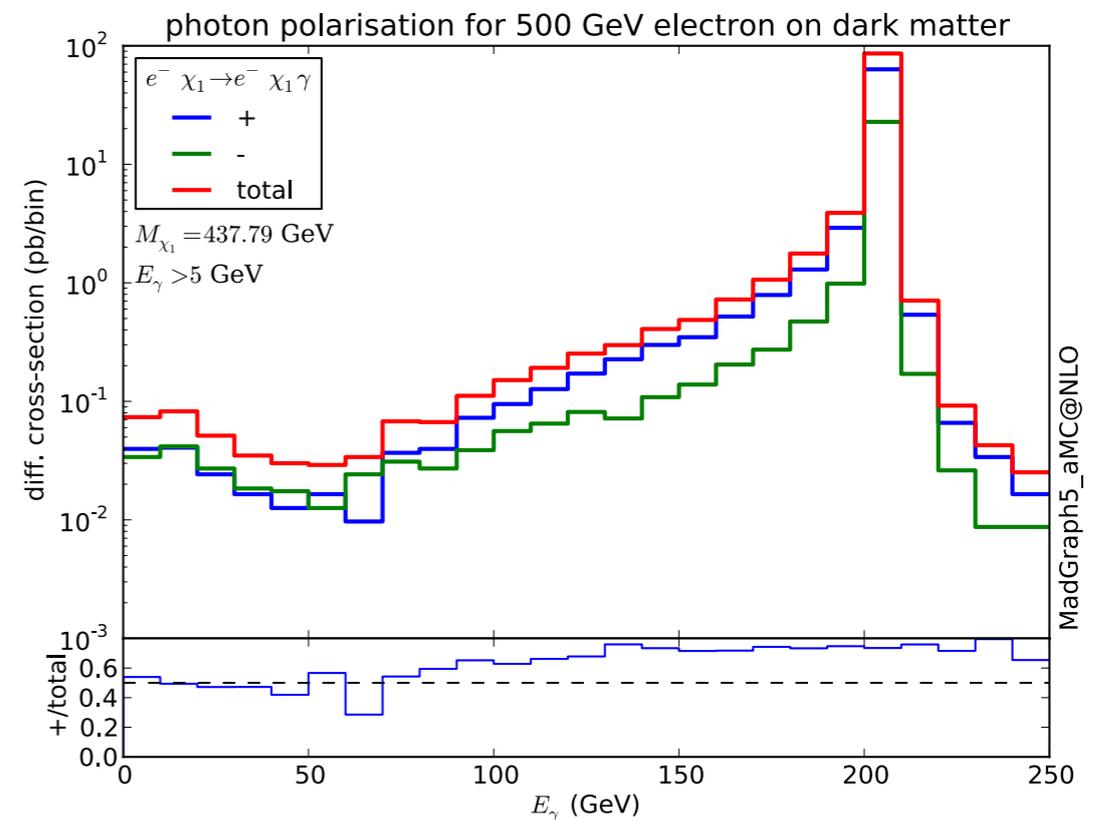
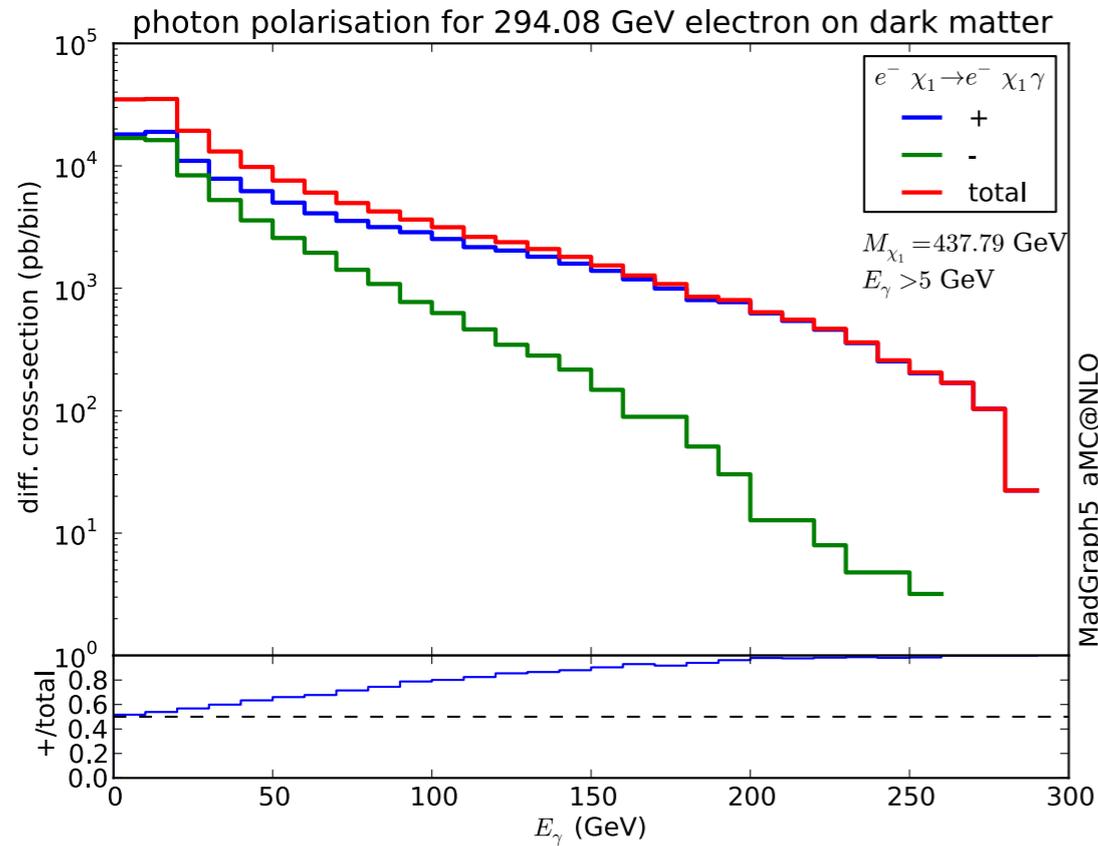


Polarisation signal is there

providing that the number of positrons is not as high as that of electrons!

Application 4: dark matter scattering off electrons

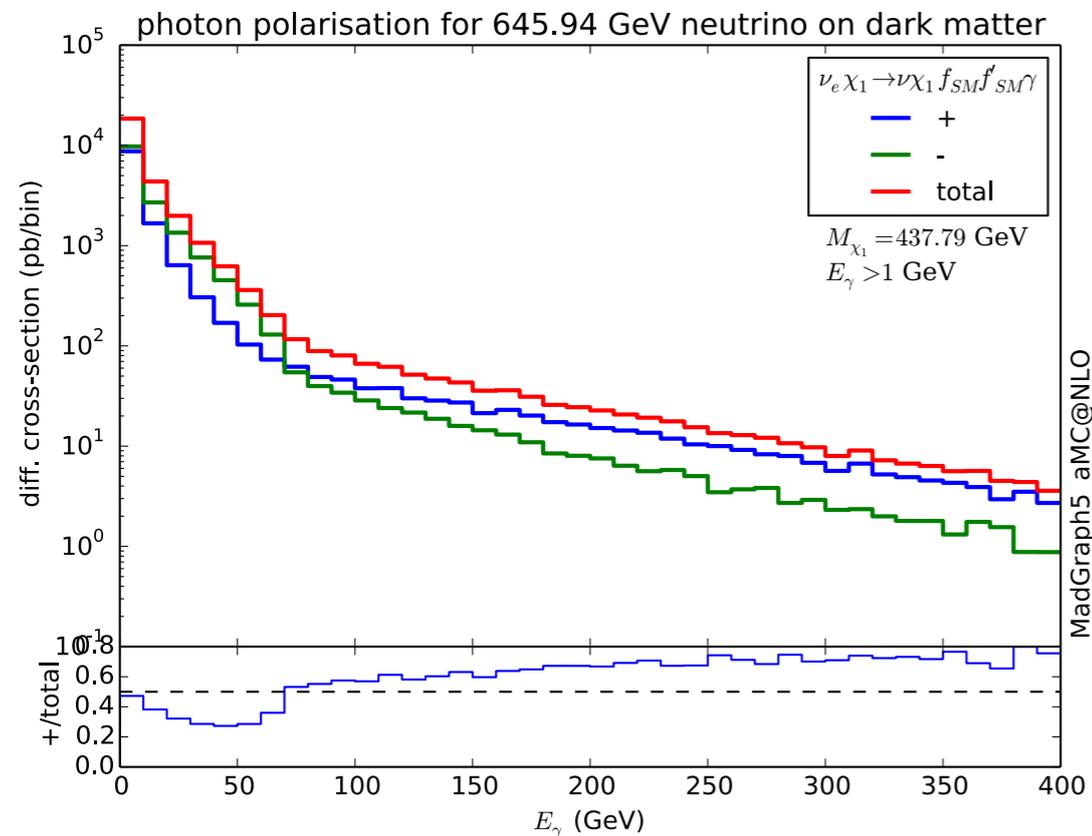
Photon spectrum for a fixed electron energy



Because the photon energy spectra are so different we may be able to reconstruct high energy electrons produced in jets and scattering off the DM halo

Application 5: dark matter scattering off neutrinos

$$\chi_0 \nu_i \rightarrow \chi_0 \nu_j f_{SM,i}^{\pm} f_{SM,j}^{\mp} \gamma,$$



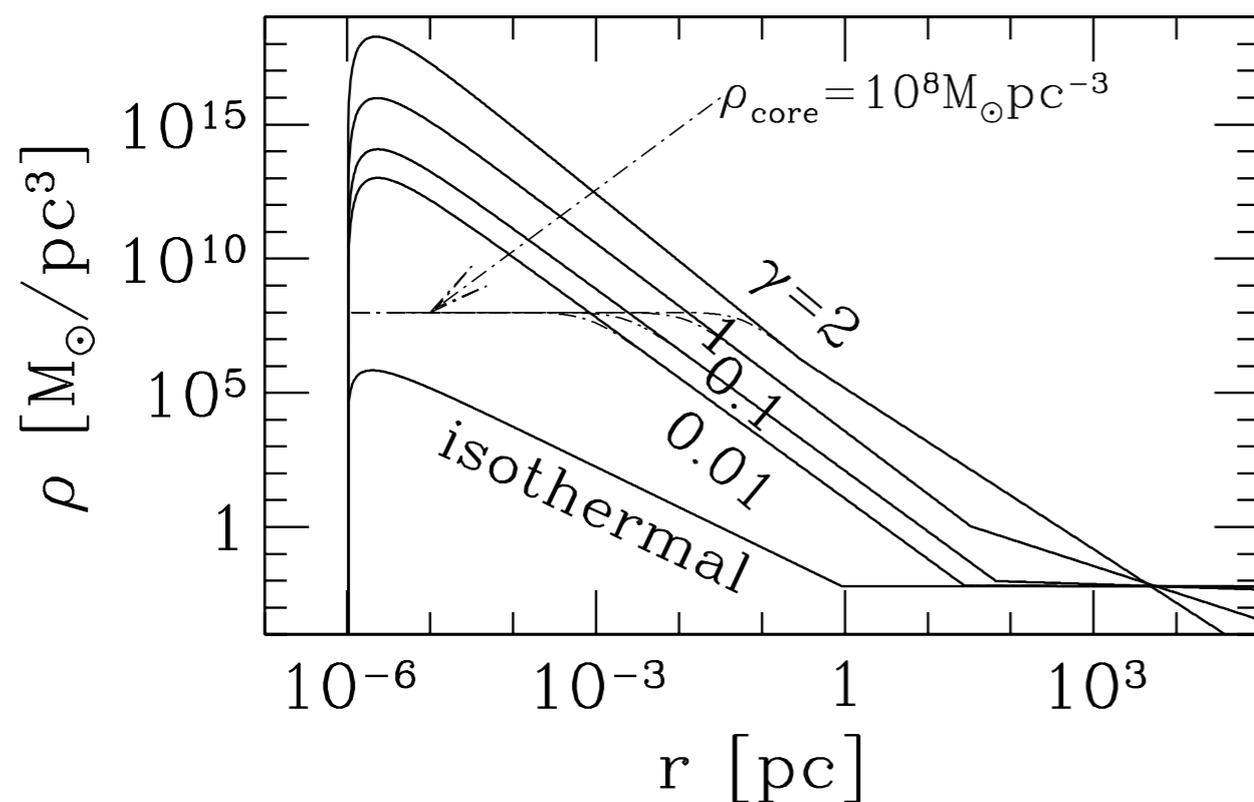
*could be applied also to neutrino-DM interactions
for “complicated” processes*

Many more applications ... Exploration phase

Dark Matter Near Black Holes

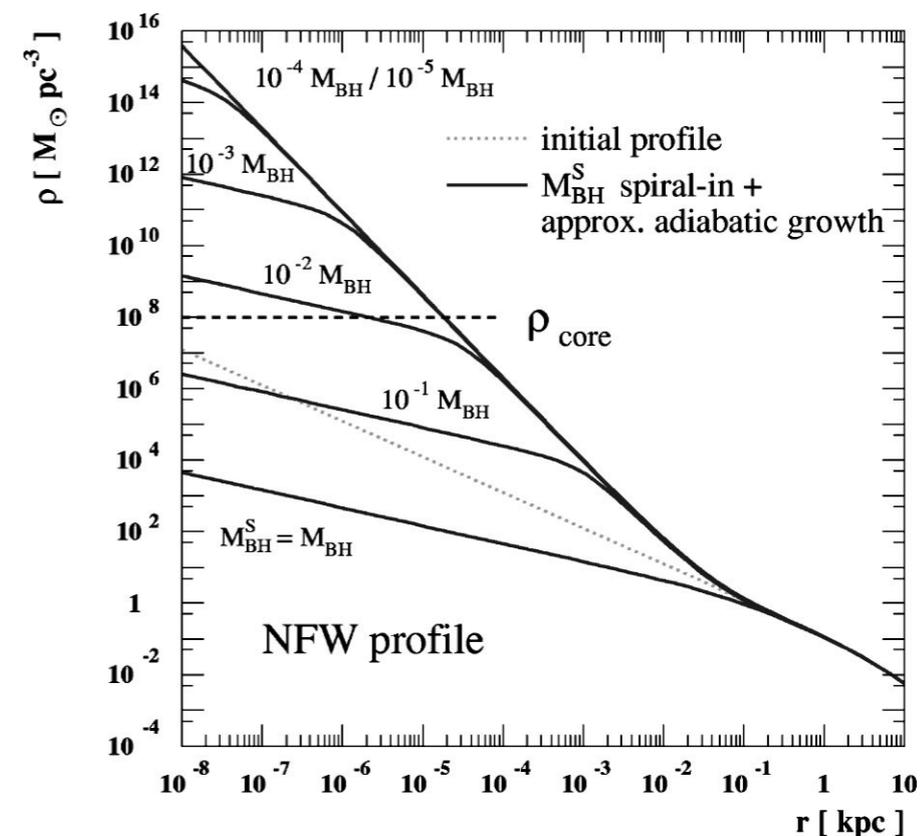
DM accretion near Black Holes

Gondolo & Silk 1999



Strong enhancement of DM profile

Ullio et al 2001



though depends on the BH seeds

Condition of formation of spikes:

- * BHs at the center of galaxies can grow adiabatically
- * Adiabatic growth

Ipsier & Sikivie (1987): isothermal $\rightarrow r^{-3/2}$,

Gondolo & Silk (1999) : NFW $\rightarrow 7/3$

Spikes in the DM density distribution

Conservation of momentum in adiabatic growth from initial to final state

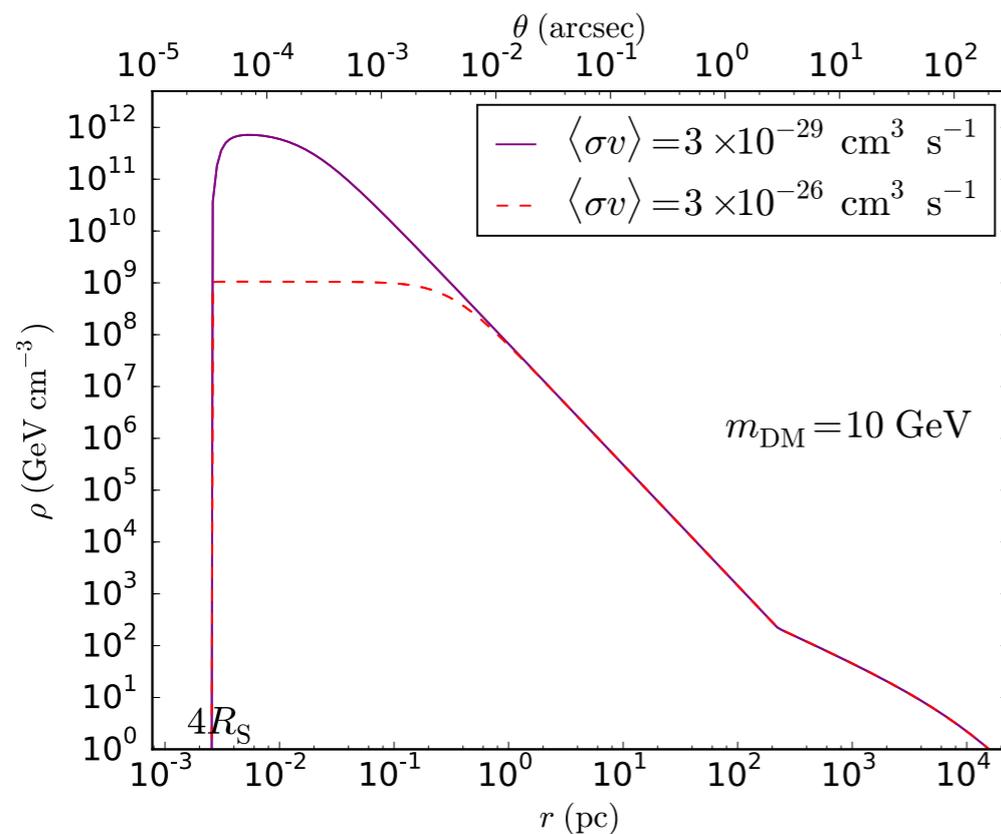
$$r v(r) = cst$$

$$v(r) = (GM(r)/r)^{1/2}$$

$$\int_0^{r_i} \rho_i(r) r^2 dr = \int_0^{r_f} \rho_f(r) r^2 dr$$

$$\gamma_{\text{sp}} = \frac{9 - 2\gamma}{4 - \gamma}$$

NFW: 7/3 inner slope



$$\rho_{\text{sat}} = \frac{m_{\text{DM}}}{\langle \sigma v \rangle t_{\text{BH}}}$$

+ spikes can be destroyed by galaxy dynamics.
 stellar heating (+ mergers, non adiabatic contraction)

Implication of a DM spike for the Milky Way

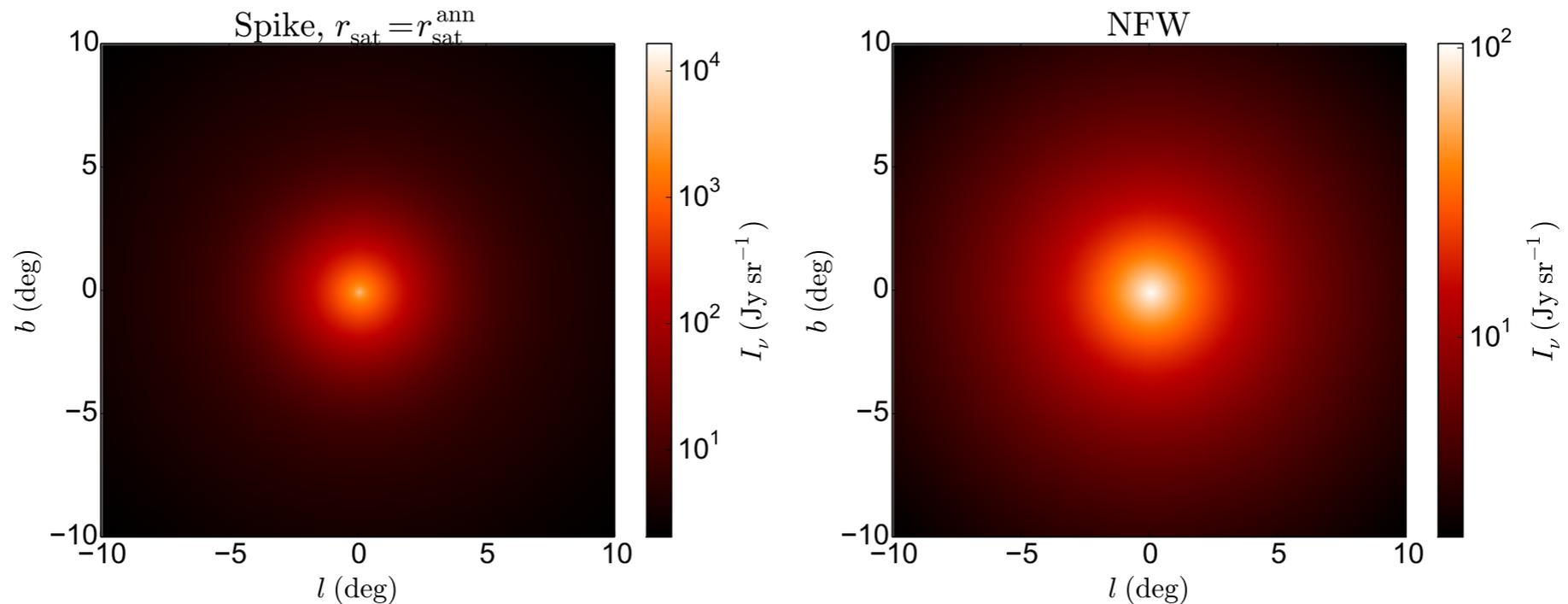
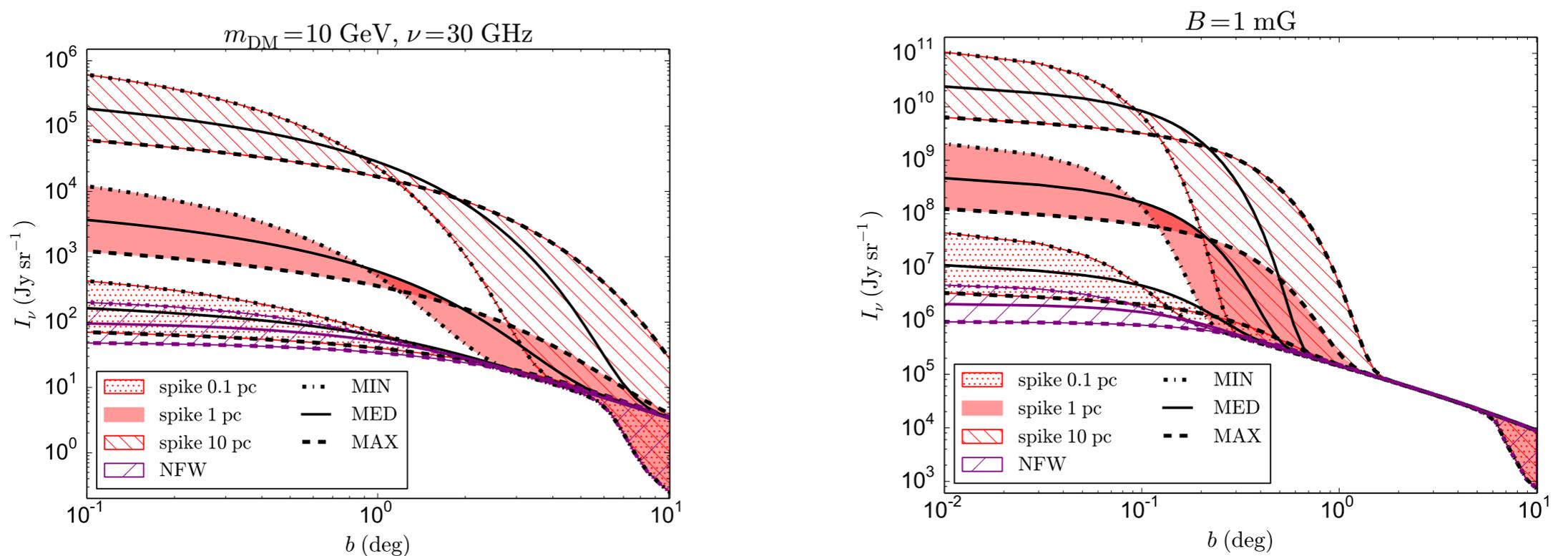


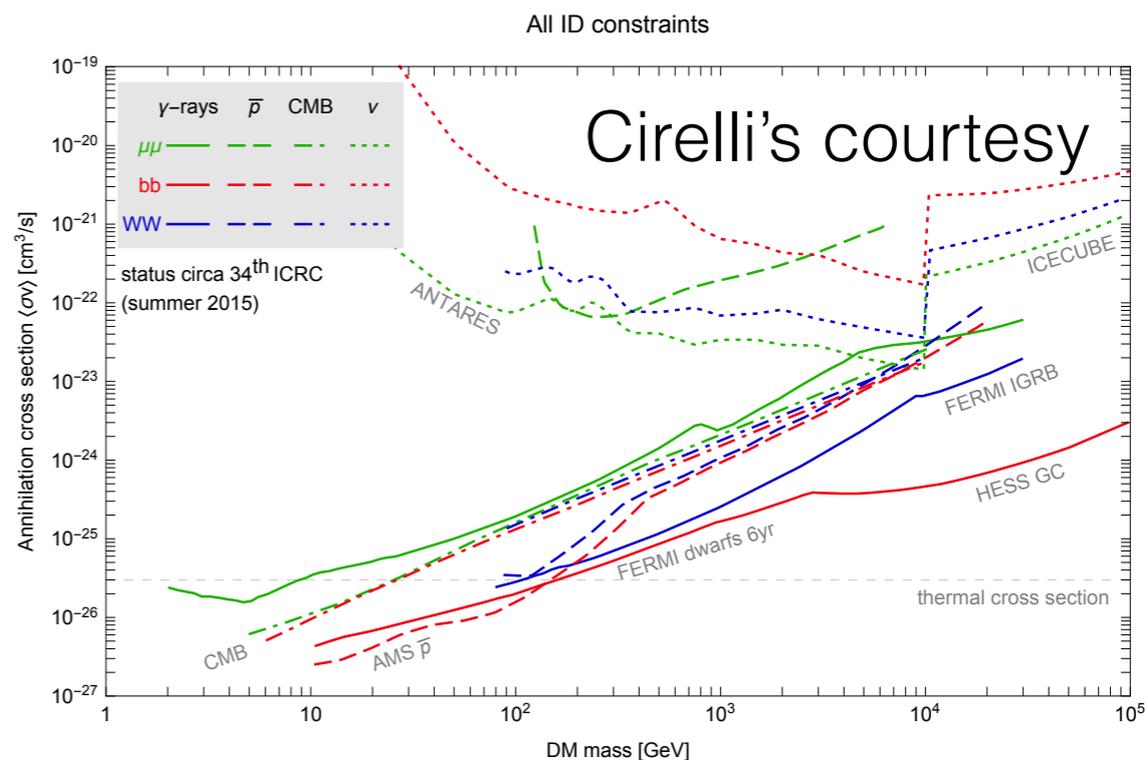
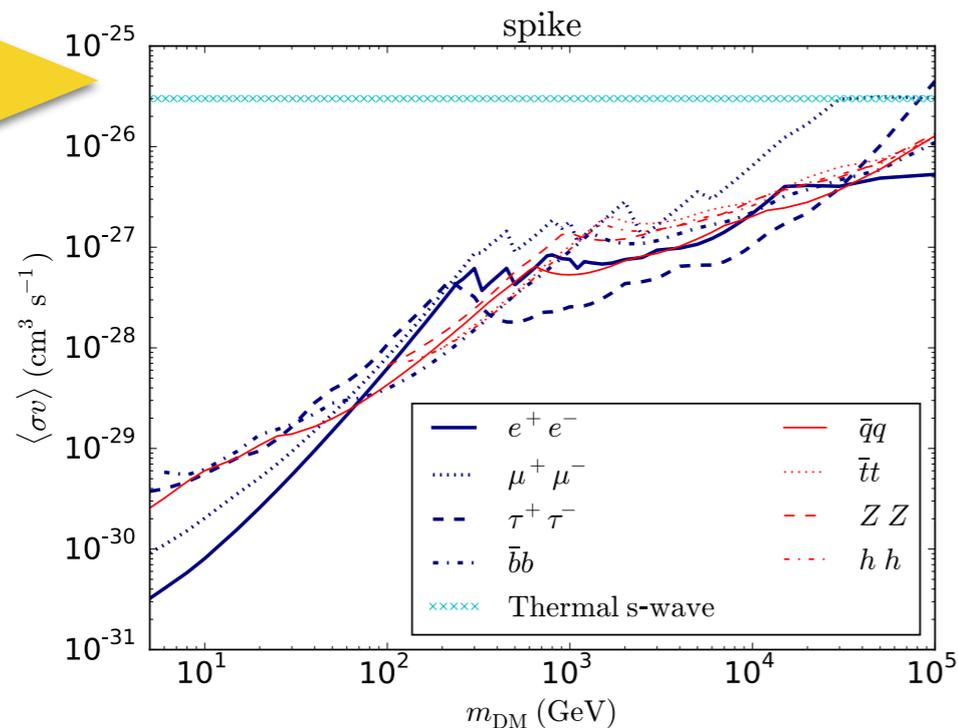
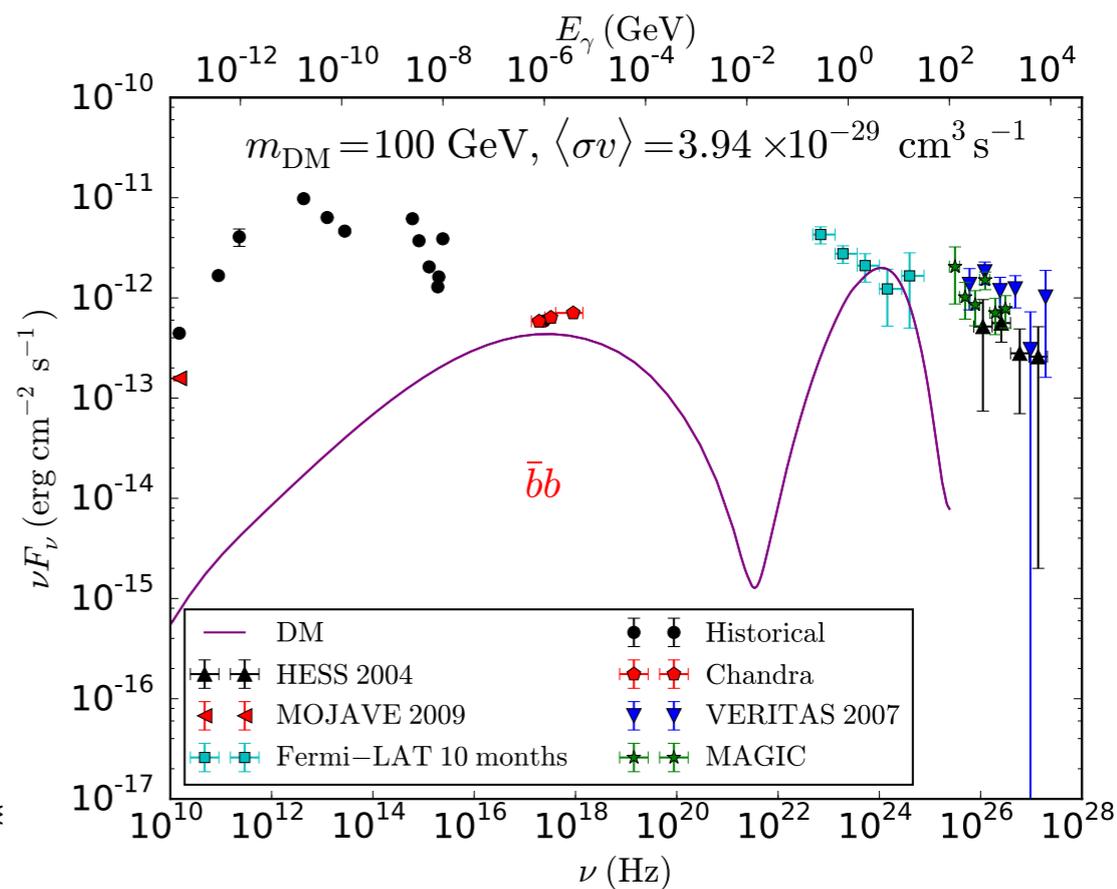
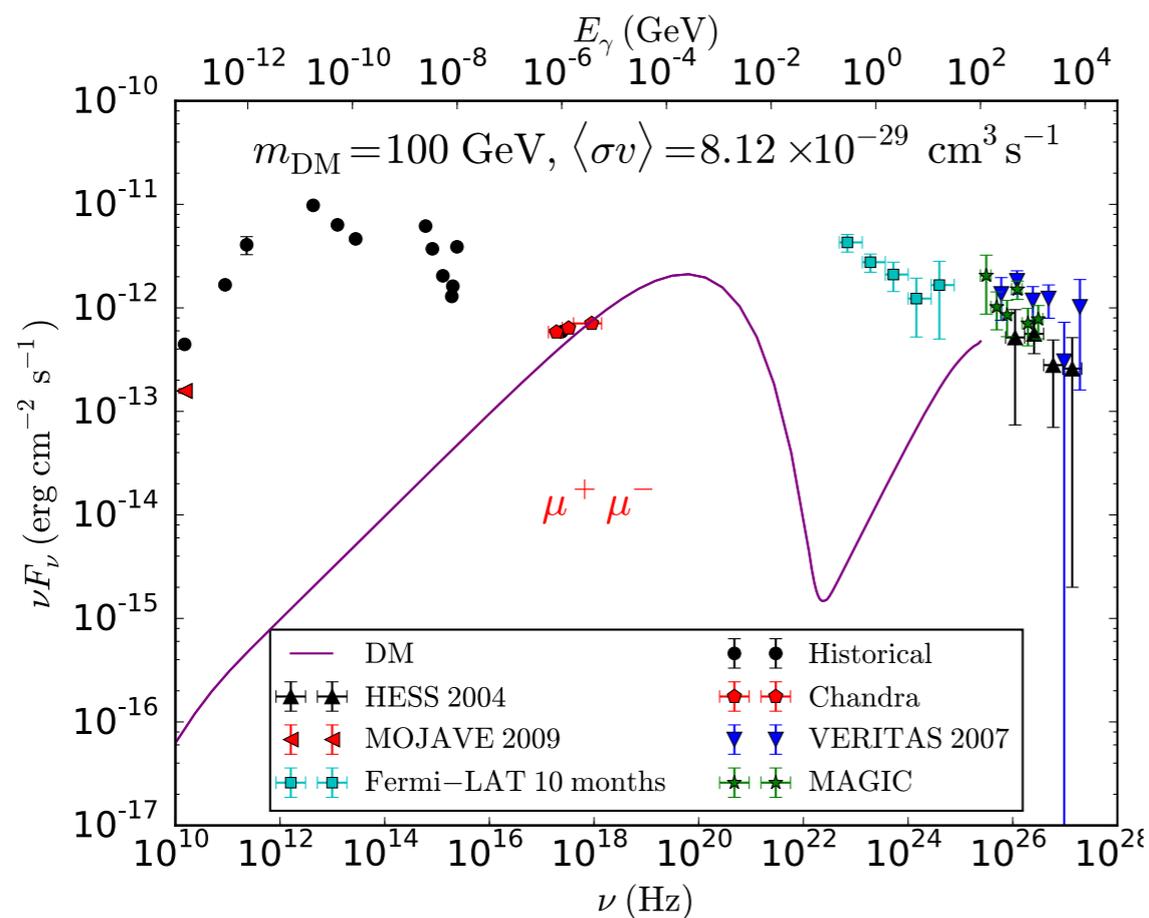
Figure 7.1: 30 GHz maps of the synchrotron intensity induced by 10 GeV DM particles, for $\langle\sigma v\rangle = 3 \times 10^{-26} \text{ cm}^3 \text{ s}^{-1}$, $B = 3 \mu\text{G}$, and the MED set of propagation parameters. The DM profiles used are a spiky profile with $\gamma_{\text{sp}} = 7/3$, $R_{\text{sp}} = 1 \text{ pc}$, with $r_{\text{sat}} = r_{\text{sat}}^{\text{ann}}$ (**left panel**), and the NFW profile (**right panel**).



Implications for AGNs M87

(Virgo, 16 Mpc, ~ 200 x heavier than MW)

adiabatic spike with $\rho \propto r^{-7/3}$



Conclusion

“Standard Model” (massive) neutrinos

$$\Gamma(\nu_3 \rightarrow \nu_2 \gamma) = 1.27 \times 10^{-51} \text{ s}^{-1} \quad \text{Invisible!}$$

Beyond Standard Model (massive) neutrinos

$$\phi_{\gamma, pol}^{\nu_i} < 2 \times 10^{-2} \left(\frac{m_{\nu_i}}{\text{eV}} \right) \text{ ph cm}^{-2} \text{ s}^{-1}$$

Can help to test new neutrino interactions

Sterile neutrino decay

$$\phi_{pol, 3.5keV}^{Perseus} \sim 10^{-6} \Delta_{CP}^{\nu_s} \text{ ph cm}^{-2} \text{ s}^{-1}$$

Neutrino and DM scattering off electrons & neutrinos

