

Summary of Works and ideas in collaboration with

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based on the model µvSSM hep-ph/0508297, PRL (2006)

The fact that the Higgs is:

- -likely an elementary scalar
- -with a mass of ~125 GeV

puts support on the idea of SUSY...

scalar particles exist,..., they produce the hierarchy problem, ..., SUSY predicts the Higgs with a mass ≤ 140 GeV

- The simplest interpretation of SUSY is to assume R-parity conservation:
 - Particles and sparticles are not mixed
 - the LSP is stable producing missing energy at colliders
 - However, novel signals with missing energy have not been discovered yet
- If R-parity is violated (RPV), SUSY particles can decay to standard model particles, and the bounds become significantly weaker

Supersymmetry with right-handed neutrinos

Rigth-handed neutrinos are likely to exist in order to generate neutrino $Y_{ij}^{\nu} H_u L_i \nu_i^c$ masses

$$L_i = \begin{pmatrix} \nu_i \\ e_i \end{pmatrix}, \qquad \stackrel{e_i^c}{\underset{i}{v_i^c}}, \quad Q_i = \begin{pmatrix} u_i \\ d_i \end{pmatrix}, \quad \stackrel{d_i^c}{\underset{i}{u_i^c}}$$
 $H_d = \begin{pmatrix} H_d^0 \\ H_d^- \end{pmatrix}, \quad H_u = \begin{pmatrix} H_u^+ \\ H_u^0 \end{pmatrix}$

R-parity is naturally violated:

W =
$$Y_{ij}^e H_d L_i e_j^c + Y_{ij}^d H_d Q_i d_j^c - Y_{ij}^u H_u Q_i u_j^c - Y_{ij}^v H_u L_i v_j^c$$

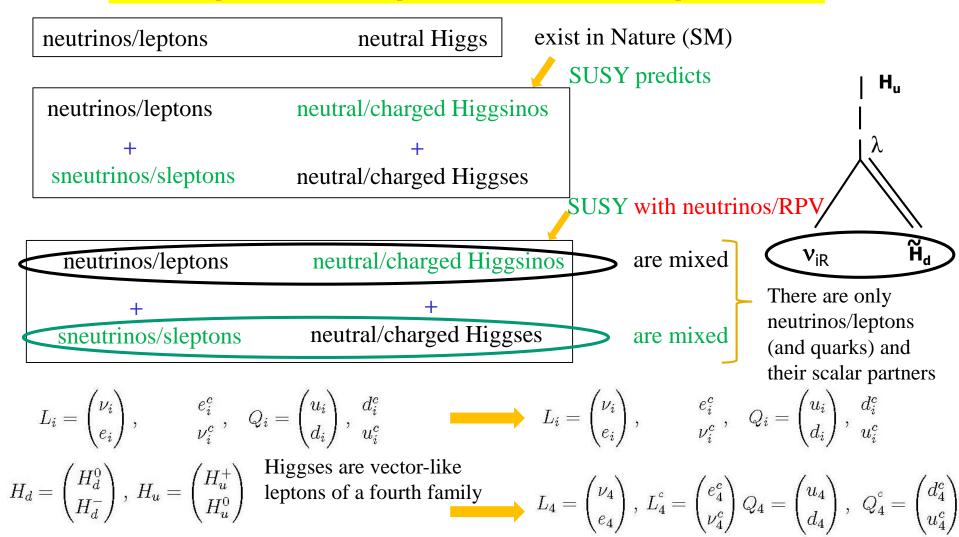
+ $\lambda_{ijk} L_i L_j e_k^c + \lambda'_{ijk} L_i Q_j d_k^c + (\kappa_{ijk} v_i^c) v_j^c v_k^c + \lambda_i H_u H_d v_i^c$.
Only one scale in the model: the soft SUSY-breaking scale ~ TeV $v_{ip}(+1)$ H_d (-1)

- * Only one scale in the model: the soft SUSY-breaking scale ~ TeV
- * The VEVs of the right sneutrinos generate:
 - an effective μ-term solving the μ problem
 - effective Majorana masses for neutrinos solving the v problem (how to accommodate neutrino data) electroweak-scale seesaw with $Y_{\rm V} \lesssim 10^{-6}$ like the electron Yukawa
 - an effective bilinear term

μνSSM Lopez-Fogliani, C. M., PRL 2006

 $v_{iR}(+1)$

A simple re-interpretation of the spectrum



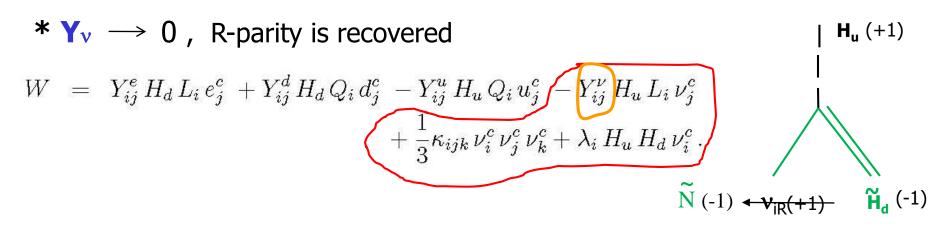
Heavier quarks and leptons will decay to the lightest ones

Proposal of new (vector-like) quarks

In this framework, the first scalar particle discovered at the LHC is a sneutrino belonging to a 4th-family vector-like doublet representation

Phenomenology of displaced vertices at the LHC

It is related to neutrino physics:



* But $Y_{\nu} \lesssim 10^{-6}$, implying that there is a relevant area of the parameter space producing displaced vertices

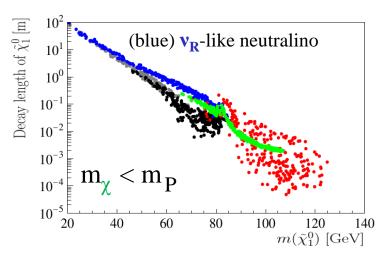
Since the 'LSP' is no longer stable, all particles are potential LSPs: neutralino, sneutrino, stau, squark, chargino

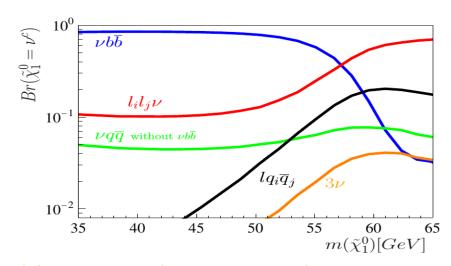
Neutralino LSP

$$Y_{\rm V} \sim 10^{-6}$$

 $m_{\gamma} \sim 20 - 120 \,\text{GeV}$ have decay lengths $\sim m - 1 \,\text{mm}$

Bartl, Hirsch, Vicente, Liebler, Porod, JHEP 2009





Example: Displaced neutralino from SM-like Higgs

Fidalgo, Lopez-Fogliani, C.M., Ruiz de Austri, JHEP 2011 Ghosh, Lopez-Fogliani, Mitsou, C.M., Ruiz de Austri, PRD 2013

$$h_{4} \rightarrow \tilde{\chi}^{0} \tilde{\chi}^{0} \rightarrow 2P2\nu \rightarrow 2\tau^{+}2\tau^{-}2\nu$$

$$h_{4} \rightarrow \tilde{\chi}^{0} \tilde{\chi}^{0} \rightarrow 2P2\nu \rightarrow 2\mathbf{b} \ 2\mathbf{\bar{b}} \ 2\nu$$

$$\uparrow^{P_{i}} \qquad \tau^{+} (\mathbf{b})$$

$$\uparrow^{A_{4}} \qquad \tau^{-} (\mathbf{\bar{b}})$$

$$\uparrow^{A_{5}} \qquad \tau^{+} (\mathbf{b})$$

$$\uparrow^{A_{7}} \qquad \tau^{+} (\mathbf{b})$$

$$\uparrow^{A_{7}} \qquad \tau^{-} (\mathbf{\bar{b}})$$

$$\uparrow^{A_{7}} \qquad \tau^{-} (\mathbf{\bar{b}})$$

$$\uparrow^{A_{7}} \qquad \uparrow^{A_{7}} \qquad \uparrow^{A_{7}} \qquad \uparrow^{A_{7}} \qquad \uparrow^{A_{7}} \qquad \uparrow^{A_{7}} \qquad \downarrow^{A_{7}} \qquad \downarrow^{A_{7}}$$

 $(2m_{\tau} < m_p < 2m_b)$ multileptons $(2m_b < m_p)$ multijets

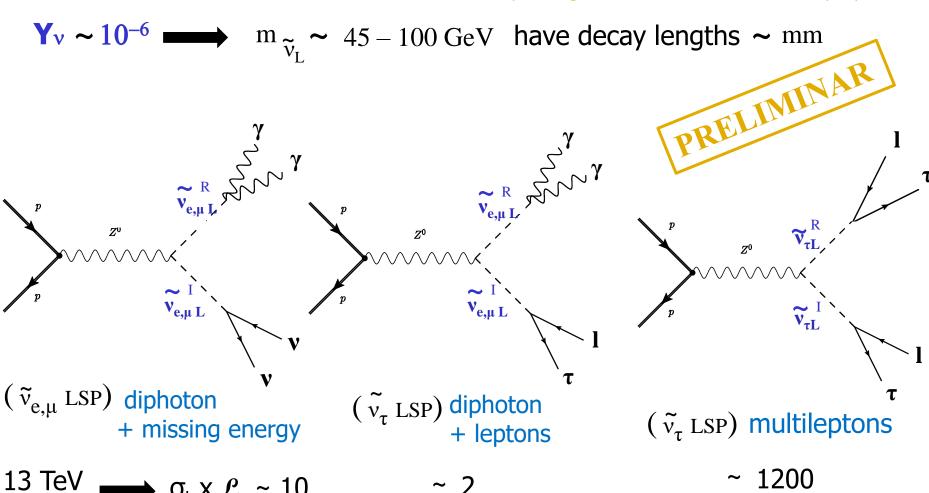
$$au^{+}(\mathbf{b})$$
 $Br(\tilde{\chi}_{4}^{0} \to \sum_{i=1}^{3} v_{i} \tau^{+} \tau^{-}) \approx 99\%$ $Br(h_{4} \to \tilde{\chi}_{4}^{0} \tilde{\chi}_{4}^{0}) \approx 1\%$

Novel signals from novel R-parity violation

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Left Sneutrino LSP

Ghosh, Lara, Lopez-Fogliani, C.M., Ruiz de Austri, in preparation



 \mathcal{L} =20 fb⁻¹

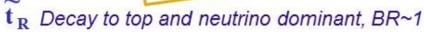
 $\sigma_t \times \mathcal{L} \sim 10$

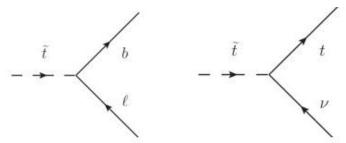
$Y_{\rm V} \sim 10^{-6}$

Stop LSP



Two dominant decay channels

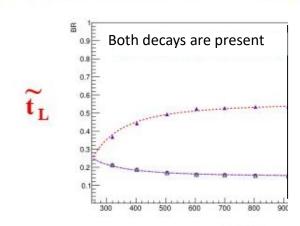


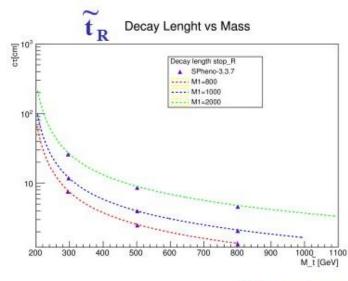


$$\begin{split} \Gamma_{\widetilde{t} \to \nu_i t} \sim \frac{|M_{\widetilde{t}}^2 - M_t^2|^2}{16\pi M_{\widetilde{t}}^3} \left[|\frac{2g_1^2}{\sqrt{3}} \frac{v_{\nu_i}}{M_1} Z_{1,6}^U|^2 + |\frac{1}{2} Z_{1,3}^U (\frac{g_1^2 v_{\nu_i}}{\sqrt{3} M_2} - \frac{\sqrt{3} g_2^2 v_{\nu_i}}{M_1})|^2 \right] \\ \Gamma_{\widetilde{t}} = \frac{M_{\widetilde{t}}}{M_1} |V_{t_i} - Z_{t_i}^U \frac{Y_{\nu_{jj}} v_{\nu_{j+1}^2}}{V_{\nu_{jj}} v_{\nu_{j+1}^2}} + |Z_{t_i}^U \frac{Y_{e} v_{\nu_{j+1}^2}}{V_{e} v_{\nu_{j+1}^2}} \end{split}$$

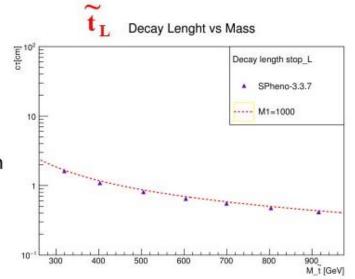
$$\Gamma_{\widetilde{t} \to \nu_i t} \sim \frac{1}{16\pi M_{\widetilde{t}}^3} \left[\left| \frac{2J_1}{\sqrt{3}} \frac{\delta \nu_i}{M_1} Z_{1,6}^C \right|^2 + \left| \frac{2}{2} Z_{1,3}^C \left(\frac{J_1 - \nu_i}{\sqrt{3} M_2} - \frac{V_2 J_2 \nu_i}{M_1} \right) \right|^2 \right]$$

$$\Gamma_{\widetilde{t} \to \ell_j b} \sim \frac{M_{\widetilde{t}}}{16\pi} \left[|Y_{d,33} Z_{1,3}^U \frac{Y_{\nu_{jj}} v_{\nu_j^c}}{\lambda_i v_{\nu_i^c}} |^2 + |Z_{1,6}^U \frac{Y_e v_{\nu_j}}{\lambda_i v_{\nu_i^c}} |^2 \right]$$





Decay length ~ cm



Cross section similar to MSSM

Conclusions

SUSY with right-handed neutrinos naturally produces R-parity violation

- Interesting theoretical advantages:
 - solves the μ problem
 - solves the v problem
 - reinterpretion of the Higgs(es) as a "4th family" of lepton superfields
- Interesting LHC phenomenology:

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    Novel signals with displaced vertices,
multilepton final states,
multijets
diphoton + leptons
diphoton + missing energy
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