



tW differential cross-sections with ATLAS at 13 TeV (dilepton)

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CoEPP Workshop, Glenelg



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Context

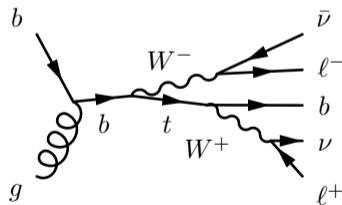
- Follow-up of the total tW cross-section in the dilepton final state ([[arXiv:1612.07231](#) →], submitted to JHEP)
- Particle level **differential** measurement
- Planned for **Moriond EW** with the 2015+2016 ATLAS dataset

Overview

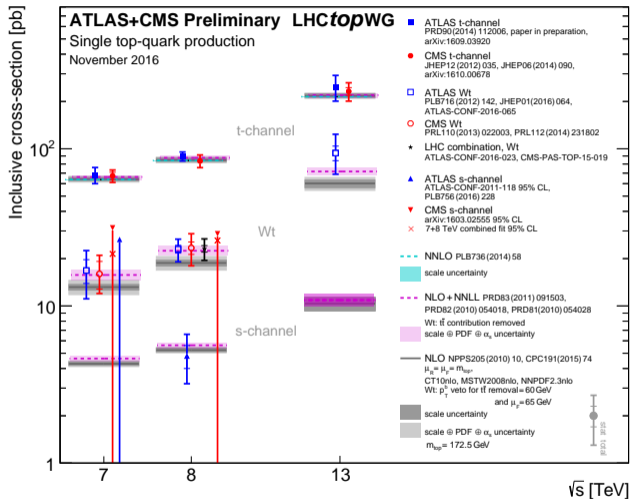
- Dilepton final state
- **BDT** to separate tW from $t\bar{t}$ background
- Unfold with **D'Agostini iterative Bayesian** method
- New plots in this presentation are **simulation only**

Motivation

- The tW channel probes the tWb vertex
- Unlike s - and t -channel it is not affected by the existence of four-fermion operators
- **Practical goal** is to compare Monte Carlo tools and models in this channel

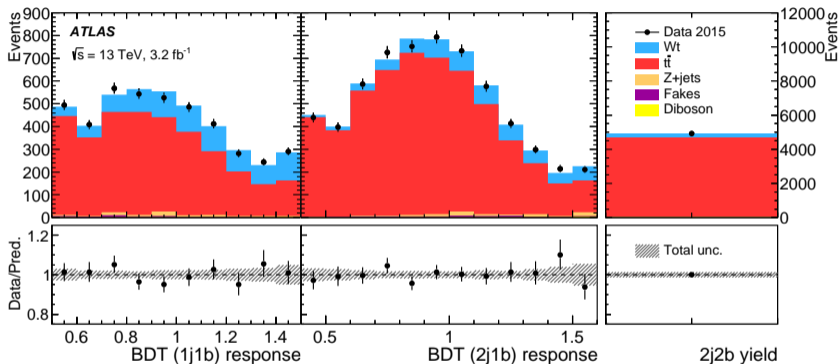


Reminder: Single top cross-sections at the LHC



LHCtopWG preliminary summary, November 2016 [ATLAS →]

Reminder: 13 TeV inclusive tW cross-section



$$\sigma_{w_t} = 94 \pm 10 \text{ (stat.) } {}^{+28}_{-22} \text{ (syst.) } \pm 2 \text{ (lumi.) pb}$$

submitted to JHEP [[arXiv:1612.07231](https://arxiv.org/abs/1612.07231) →]

Selection and definitions

≥ 1 jet with $p_T > 25$ GeV, $|\eta| < 2.5$

2 leptons of opposite charge with $p_T > 20$ GeV,

≥ 1 lepton with $p_T > 27$ GeV, veto third with $p_T > 20$ GeV

lepton matched to the **trigger** object

DF

$E_T^{\text{miss}} > 50$ GeV, if $m_{\ell\ell} < 80$ GeV

$E_T^{\text{miss}} > 20$ GeV, if $m_{\ell\ell} > 80$ GeV

SF

$E_T^{\text{miss}} > 40$ GeV, always

veto, if $m_{\ell\ell} < 40$ GeV

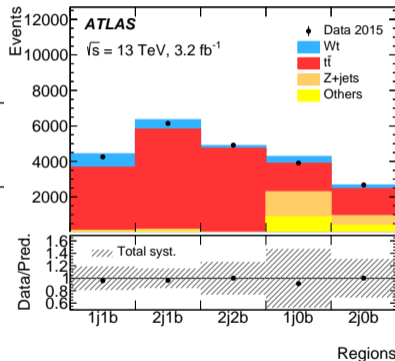
$4E_T^{\text{miss}} > 5m_{\ell\ell}$, if $40 < m_{\ell\ell} < 81$ GeV

veto, if $81 < m_{\ell\ell} < 101$ GeV

$2m_{\ell\ell} + E_T^{\text{miss}} > 300$ GeV, if $m_{\ell\ell} > 101$ GeV


Exactly one jet, which is a *b*-tagged jet (1j1b)

- Summary plot (inclusive; 2015):



- Regions other than 1j1b will be used for validation only

Separating the $t\bar{t}$ background

- Almost entirely $Wt+t\bar{t}$  (1j1b bin)
- We use a **Boosted Decision Tree** (TMVA)
- Inputs are kinematic variables of subsets of objects plus some whole-event variables
- Procedure to narrow down the list:
 1. rank variables by **separation power**
 2. remove highly-correlated variables
 3. find combinations with best separation
 4. tie-break by preferring fewer variables

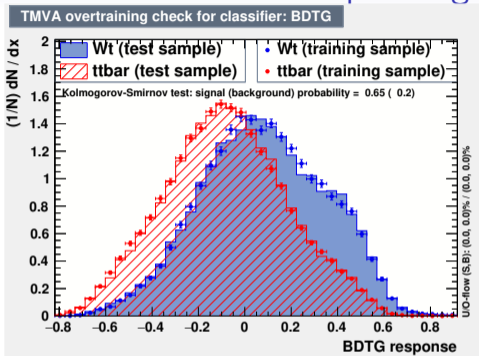
$$\langle S^2 \rangle = \frac{1}{2} \int \frac{(Y_S(y) - Y_B(y))^2}{(Y_S(y) + Y_B(y))} dy$$

- **Cut** on BDT output value of 0.3 (chosen to minimise the total uncertainty)

Variable	$S [10^{-2}]$
$p_T^{\text{sys}}(\ell_1 \ell_2 E_T^{\text{miss}} j_1)$	4.1
$\Delta p_T(\ell_1 \ell_2 j_1, E_T^{\text{miss}})$	2.5
$\sum E_T$	2.3
$\eta(\ell_1 \ell_2 E_T^{\text{miss}} j_1)$	1.3
$\Delta p_T(\ell_1 \ell_2, E_T^{\text{miss}})$	1.1
$p_T^{\text{sys}}(\ell_1 \ell_2 j_1)$	1.0
$C(\ell_1 \ell_2)$	0.9
$m(\ell_2, j_1)$	0.2
$m(\ell_1, j_1)$	0.1

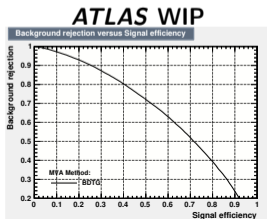
p_T^{sys} is the vectorial sum of p_T ;
 C is centrality (scalar sum of $p_T / \Sigma E$)

Separating the $t\bar{t}$ background



Theoretical considerations

- tW at NLO and $t\bar{t}$ at LO contain interfering diagrams (“doubly-resonant”)
- Different schemes exist to handle this
- **Diagram subtraction** (DS) adds a gauge-invariant term to cancel the $t\bar{t}$ contribution
- **Diagram removal** (DR) discards all diagrams with resonant contributions
- A new scheme **DR2** also keeps the mixed term



$$|\mathcal{M}_{\text{tot}}|^2 = |\mathcal{M}_{\text{nr}}|^2 + 2\text{Re}(\mathcal{M}_{\text{nr}}\mathcal{M}_r) + |\mathcal{M}_r|^2$$

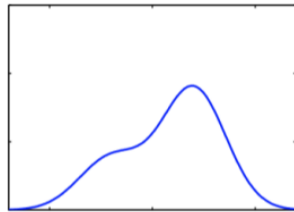
Unfolding

Overview

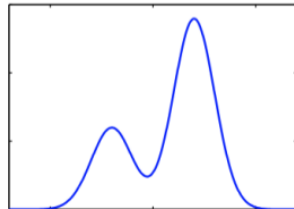
- Our definitions of top quark etc. depend on details of the detector (**reconstruction level**)
- Ideal: present in terms of the **physical** objects
- Compromise: use straightforward **definitions** for objects at **particle level**
- Reco \rightarrow particle level is called **unfolding**

General procedure

- Account for **migration between bins** of reco-level and particle-level quantities
- Invert the response matrix



↑ Folding ↑ ↓ Unfolding ↓



[source \rightarrow]

Unfolding: choice of binning

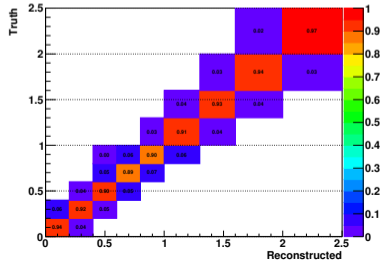
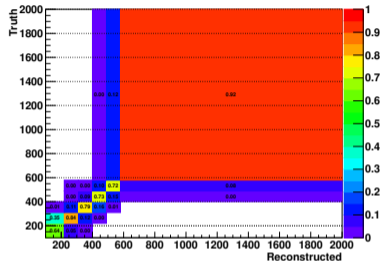
Binning goals

- **Stable:** $\sim 60\%$ of entries on the diagonal
- **Balanced:** $\lesssim 20\%$ statistical uncertainty per bin

Binning procedure

- Identical for particle and reco levels (negligible bias)
- First and last bins chosen based on statistical uncertainty
- Other bins are a **multiple of the resolution** of the variable in question

Note: top plot illustrates $m(\ell\nu b)$, bottom $|\eta(b)|$



Unfolding: specific procedure

Unfolding procedure

- Iterative Bayes-inspired (D'Agostini) in RooUnfold
- Correct for efficiency to reco fiducial event [E^{-1}]
- Then iterative matrix unfolding
- Then correct for unmatched/non-fiducial events [A]

$$\frac{d\sigma}{dX} = \frac{N_i^{\text{ufd}}}{L\Delta_i} \quad (\Delta_i = \text{bin width})$$

Number of iterations

- Chose $N_{it.}$ to minimise χ^2 difference between testing and unfolded samples (per variable)
- Larger $N_{it.}$ emphasise finer-detailed features that are more likely to be statistical noise

Variable	$N_{\text{iterations}}$	
$E(\ell_1\ell_2b)$	5	
$m_T(\ell_1\ell_2\nu b)$	5	*
$m(\ell_1\ell_2b)$	3	
$m(\ell_1b)$	5	†
$m(\ell_2b)$	3	
$E(b)$	6	‡
$p_T^{\text{sys}}(\ell_1\ell_2\nu b)$	10	

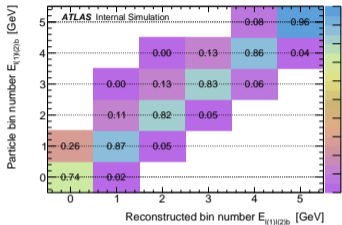
*: sensitive to DR vs. DS scheme

†: sensitive to spin correlation

‡: sensitive to theory modelling

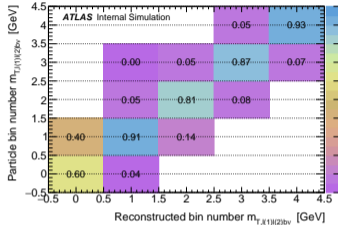
Normalised migration matrices

$$E(\ell_1 \ell_2 b)$$



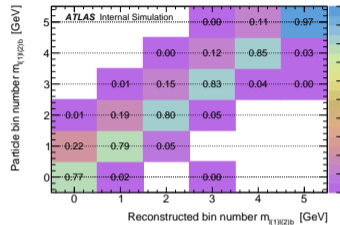
ATLAS WIP

$$m_T(\ell_1 \ell_2 \nu b)$$

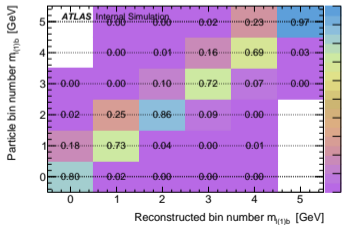


ATLAS WIP

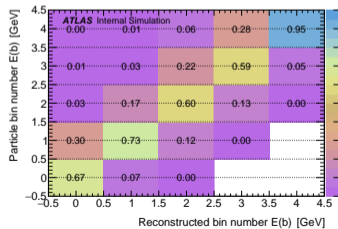
$$m(\ell_1 \ell_2 b)$$



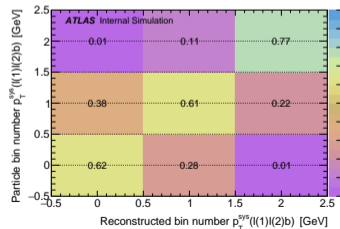
ATLAS WIP



$$m(\ell_1 b)$$



$$E(b)$$



$$p_T^{\text{sys}}(\ell_1 \ell_2 b)$$

Propagation of systematics

X^0 := nominal, X^\pm := variation; unfolding always uses a response matrix made with tW^0

Signal-only uncertainties

- tW matrix element generator, parton shower/had. generator, DS vs. DR, initial-/final-state radiation tune
- $\text{unfold}(tW^\pm) - \text{truth}(tW^\pm)$
- Unfold variations with nominal, cf. truth
- Difference* at particle level is the uncertainty

Other uncertainties

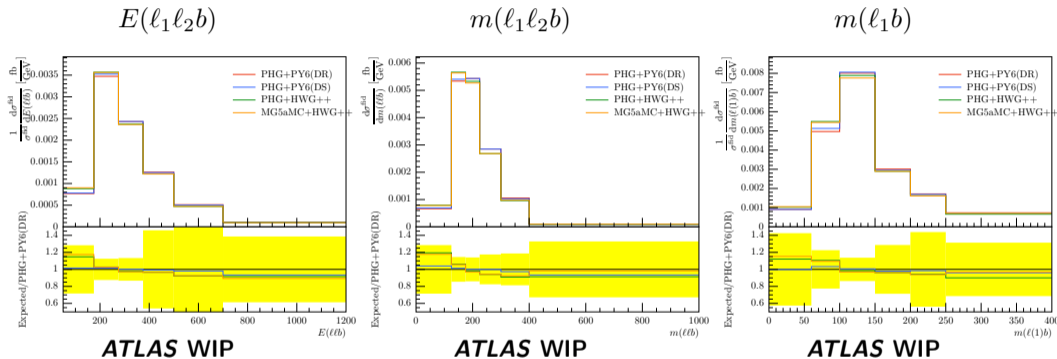
- The background and signal are both varied
- $\text{unfold}(tW^\pm + t\bar{t}^\pm - t\bar{t}^0) - \text{truth}(tW^\pm)$

Background-only uncertainties

- $t\bar{t}$ matrix element gen., parton shower/had. generator, initial-/final-state radiation tune
- $\text{unfold}(tW^0 + t\bar{t}^\pm - t\bar{t}^0) - \text{truth}(tW^0)$
- Difference between background nominal and variation is added to nominal before unfolding
- This is unfolded and compared to truth

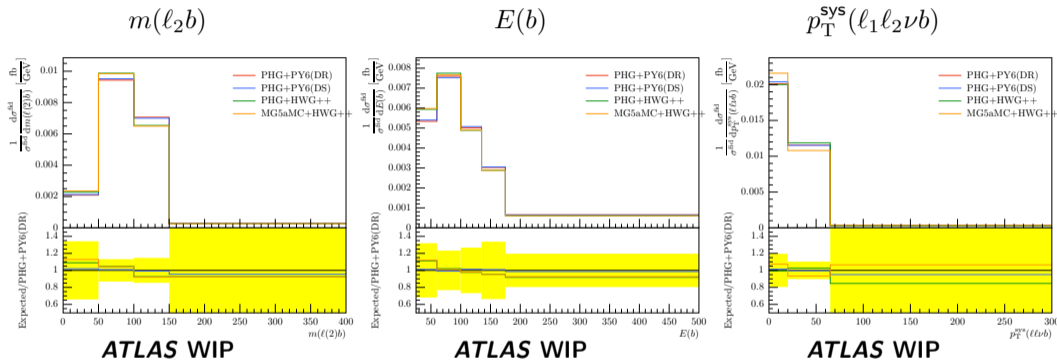
* for 2-point uncertainties we take half of the difference between the up and down variations

Normalised differential cross-sections



- DR vs. DS: relative slope would suggest possibility to differentiate the schemes

Normalised differential cross-sections



- Theory modelling: differentiation between Herwig and Pythia samples

Future of the work

General direction

- This will be the first differential measurement in the tW channel
- Future measurements will include additional **features** and more data

Full reconstruction

- Due to the ν , the system is kinematically underconstrained
- **Reconstruction** of tW system possible
- **Top EFT**: lepton association
- At particle level: pseudo-top

Diagram interference

- The $t\bar{t}$ and tW interference issue is an active area of study
- $WWbb$ predictions are being studied
- $tW + t\bar{t}$ **fiducial measurement**
- Would need to look at extra jet bins

Additional channels

- So far we've only looked at the dilepton final state
- **Single lepton** channel is also interesting

(backup)

Triggers

	2015	2016
Electrons	HLT_e24_lhmedium_L1EM20VH or HLT_e60_lhmedium or HLT_e120_lhloose	HLT_e26_lhtight_nod0_ivarloose or HLT_e60_lhmedium_nod0 or HLT_e140_lhloose_nod0
Muons	HLT_mu20_iloose_L1MU15 or HLT_mu50	HLT_mu26_ivarmedium or HLT_mu50

At least one selected lepton must be matched to the trigger object

Note that the dataset is from the top working group's dilepton derivation

BDT cut optimisation

- Studies to optimise value of BDT cut: look at total syst. and stat. uncertainty per bin
- Optimise for normalised cross-section- choose value of 0.1

	Bin 1	Bin 2	Bin 3	Bin 4	Bin 5	Bin 6	simple sum
<hr/>							
$\frac{1}{\sigma} \frac{d\sigma}{dm_{\ell\ell\nu b}}$, $r_{\text{BDT}} > 0.2$							
Stat.	0.151	0.057	0.078	0.128	0.203	0.185	
Total syst.	0.154	0.052	0.063	0.141	0.118	0.12	
Total uncert.	0.216	0.077	0.100	0.190	0.235	0.221	1.039
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$\frac{1}{\sigma} \frac{d\sigma}{dm_{\ell\ell\nu b}}$, $r_{\text{BDT}} > 0.1$							
Stat.	0.14	0.053	0.07	0.11	0.175	0.171	
Total syst.	0.159	0.036	0.058	0.089	0.096	0.168	
Total uncert.	0.212	0.064	0.091	0.141	0.200	0.240	0.948
<hr/>							
$\frac{1}{\sigma} \frac{d\sigma}{dm_{\ell\ell\nu b}}$, $r_{\text{BDT}} > 0$							
Stat.	0.133	0.05	0.064	0.101	0.163	0.161	
Total syst.	0.162	0.038	0.051	0.105	0.154	0.177	
Total uncert.	0.210	0.063	0.082	0.146	0.224	0.239	0.963
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Particle level definitions

Leptons are electrons and muons from W or Z boson decays, dressed by any non-hadronic photons in a ΔR cone of 0.1

Neutrinos are stable particles not from the decay of hadrons

Jets are anti- k_t jets with parameter 0.4, made from all stable particles excluding the leptons and neutrinos (and the photons used in dressing). **b-jets** are jets associated with a heavy flavour hadron with $p_T > 5$ GeV using ghost matching¹.

The **pre-selection** requires two leptons with $p_T > 20$ GeV and $|\eta| < 2.5$, vetoing a third such lepton. One lepton must have $p_T > 27$ GeV. Exactly one b -jet with $p_T > 25$ GeV and $|\eta| < 2.5$ must be present. There is **no event-level cut** such as on E_T^{miss} or $m_{\ell\ell}$.

¹the hadron momentum is rescaled to be vanishingly small before adding them as input to the clustering algorithm; any jet containing such a "ghost" hadron is b-tagged