

# Symmergent Gravity

Durmuş Demir

IZMIR INSTITUTE OF TECHNOLOGY

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# Question

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How to reconcile GR with the SM?

Way 1: quantized GR + the SM ☹️

Way 2: classical GR + the SM ☹️

Way 3: ??????? ☹️

What could be **Way 3** ?

► EP is **violated** in quantum theory:

$$-\frac{\hbar^2}{2m_I} \frac{\partial^2}{\partial z^2} \psi(z, t) + m_G g z \psi(z, t) = i\hbar \frac{\partial}{\partial t} \psi(z, t)$$

► EP is **revived** after quantum fluctuations are integrated out:

$$m_I \frac{d^2}{dt^2} \langle z \rangle = -m_G g$$

# An Answer

It is thus possible that:

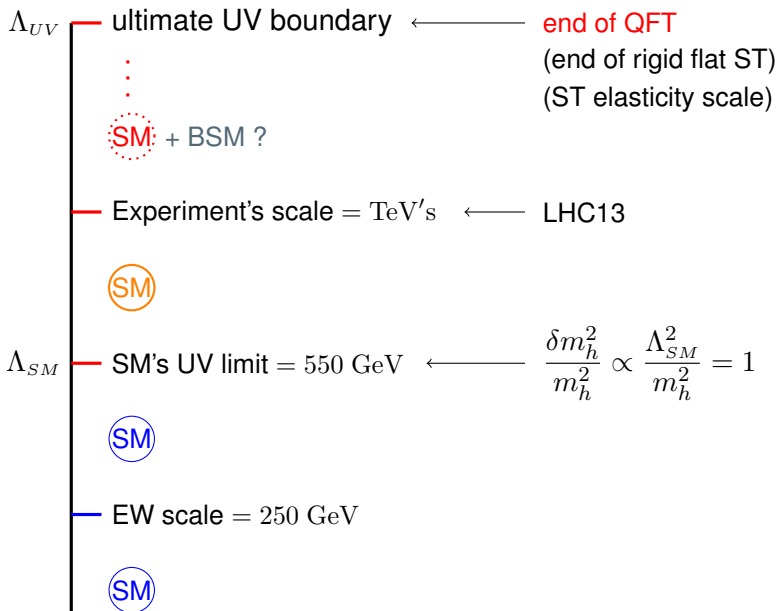
Way 3 = “classical GR + effective QFT”

reveals UV sensitivity  
(log, power-law)

involves long-wavelength fields  
(quasi-classical field theory)

lives in experimental reach  
(testable with data)

# Scales and QFTs

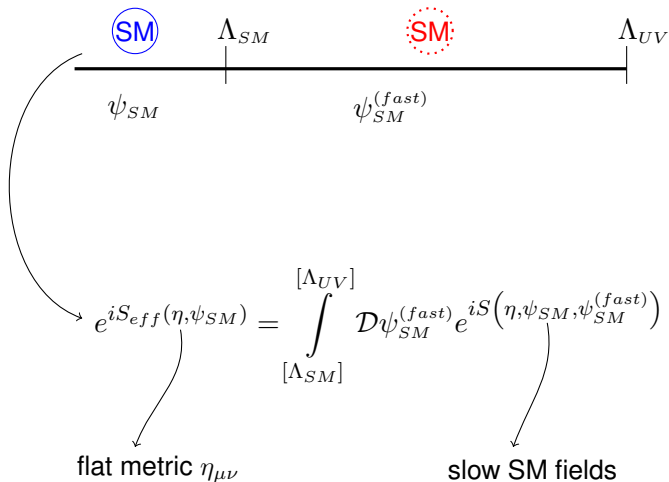


# A Realistic Setup

In line with the status at the LHC experiments:

- ▶ The **SM**, which has already transcended  $\Lambda_{SM}$ , can be assumed to **extend all the way up to  $\Lambda_{UV}$** .
- ▶ Possible **BSM** physics can be revealed by physical consistency.

# The SM Effective Action



# The SM Effective Action

$$S_{eff} = S_{tree}(\psi_{SM}) + \delta S_{log}\left(\psi_{SM}, \Lambda_{SM} \log \frac{\Lambda_{SM}}{\Lambda_{UV}}\right)$$

at the SM scale

$$+ \delta S_{power}\left(H, V_{\mu}, \Lambda_{UV}^2 - \Lambda_{SM}^2, \Lambda_{UV}^2 + \Lambda_{SM}^2\right)$$

at the UV scale

$$\Lambda_{UV}^2 + \Lambda_{SM}^2$$

( $\leadsto$  gravity)

$$\frac{\Lambda_{UV}}{\Lambda_{SM}} = \text{hierarchy}$$

$$\Lambda_{UV}^2 - \Lambda_{SM}^2$$

( $\leadsto$  ???????)



# Power-Law UV Contributions

$$\delta S_{power} = \delta S_O + \delta S_H + \delta S_V$$

$$\int d^4x \sqrt{-\eta} c_V (\Lambda_{UV}^2 - \Lambda_{SM}^2) \text{Tr}[V_\mu V^\mu]$$

$$-\int d^4x \sqrt{-\eta} c_H (\Lambda_{UV}^2 - \Lambda_{SM}^2) H^\dagger H$$

$$-\int d^4x \sqrt{-\eta} c_O (\Lambda_{UV}^4 - \Lambda_{SM}^4)$$

$$c_{O,H,V} = c_{O,H,V} \left( \frac{\Lambda_{SM}}{\Lambda_{UV}} \right)$$


# Hierarchy Problems


$$\delta S_{power} = \delta S_O + \delta S_H + \delta S_V$$

$$\int d^4x \sqrt{-\eta} c_V (\Lambda_{UV}^2 - \Lambda_{SM}^2) \text{Tr}[V_\mu V^\mu]$$

$$-\int d^4x \sqrt{-\eta} c_H (\Lambda_{UV}^2 - \Lambda_{SM}^2) H^\dagger H$$

hard UV masses to all  
 $U(1)_Y, SU(2)_L, SU(3)_C$   
gauge bosons

  
**big  
hierarchy  
problem**

  
**explicit breaking of  
 $SU(3)_C$  and  $U(1)_{EM}$**

# Curving Away Gauge Boson UV-Masses


$$\delta S_V = \delta S_V - \int d^4x \sqrt{-\eta} \frac{c_V}{2} \text{Tr}[V_{\mu\nu} V^{\mu\nu}] + \int d^4x \sqrt{-\eta} \frac{c_V}{2} \text{Tr}[V_{\mu\nu} V^{\mu\nu}]$$

$$= \int d^4x \sqrt{-\eta} c_V (\Lambda_{UV}^2 - \Lambda_{SM}^2) \text{Tr}[V_\mu V^\mu]$$

$$- \int d^4x \sqrt{-\eta} \frac{c_V}{2} \text{Tr}[V_{\mu\nu} V^{\mu\nu}]$$

$$+ \int d^4x \sqrt{-\eta} c_V \text{Tr}[V^\mu (-D^2 \eta_{\mu\nu} + D_\mu D_\nu + V_{\mu\nu}) V^\nu]$$

$$+ \int d^4x \sqrt{-\eta} c_V \text{Tr}[\partial_\nu (V_\mu V^{\mu\nu})]$$


$$D_\mu = \partial_\mu - V_\mu$$

# Curving Away Gauge Boson UV-Masses

$$\delta S_V \xrightarrow{\eta_{\mu\nu} \rightleftharpoons g_{\mu\nu}} \int d^4x \sqrt{-g} c_V (\Lambda_{UV}^2 - \Lambda_{SM}^2) \text{Tr}[V_\mu V^\mu]$$

$$- \int d^4x \sqrt{-g} \frac{c_V}{2} \text{Tr}[V_{\mu\nu} V^{\mu\nu}]$$

$$+ \int d^4x \sqrt{-g} c_V \text{Tr}[V^\mu (-\mathcal{D}^2 g_{\mu\nu} + \mathcal{D}_\mu \mathcal{D}_\nu + V_{\mu\nu}) V^\nu]$$

$$+ \int d^4x \sqrt{-g} c_V \text{Tr}[\nabla_\nu (V_\mu V^{\mu\nu})]$$


$$\mathcal{D}_\mu = \nabla_\mu - V_\mu$$

# Curving Away Gauge Boson UV-Masses

$$= - \int d^4x \sqrt{-g} \frac{c_V}{2} \text{Tr} [V_{\mu\nu} V^{\mu\nu}]$$

$$+ \int d^4x \sqrt{-g} c_V \text{Tr} [V^\mu (-\mathcal{D}^2 g_{\mu\nu} + \mathcal{D}_\mu \mathcal{D}_\nu + V_{\mu\nu} + (\Lambda_{UV}^2 - \Lambda_{SM}^2) g_{\mu\nu}) V^\nu]$$

$$+ \int d^4x \sqrt{-g} c_V \text{Tr} [\nabla_\nu (V_\mu V^{\mu\nu})]$$

What if this is a fixed value assigned to **curvature**?

$$(\Lambda_{UV}^2 - \Lambda_{SM}^2) g_{\mu\nu} \iff R_{\mu\nu}(g) ?$$

# Curving Away Gauge Boson UV-Masses

$$\underline{(\Lambda_{UV}^2 - \Lambda_{SM}^2)g_{\mu\nu} \iff R_{\mu\nu}(g)}$$

$$-\int d^4x \sqrt{-g} \frac{c_V}{2} \text{Tr}[V_{\mu\nu} V^{\mu\nu}]$$

$$+\int d^4x \sqrt{-g} c_V \text{Tr}[V^\mu (-\mathcal{D}^2 g_{\mu\nu} + \mathcal{D}_\mu \mathcal{D}_\nu + V_{\mu\nu} + R_{\mu\nu}(g)) V^\nu]$$

$$+\int d^4x \sqrt{-g} c_V \text{Tr}[\nabla_\nu (V_\mu V^{\mu\nu})]$$

if  $c_V$  is held unaffected

$$-\int d^4x \sqrt{-g} \frac{c_V}{2} \text{Tr}[V_{\mu\nu} V^{\mu\nu}] + \int d^4x \sqrt{-g} \frac{c_V}{2} \text{Tr}[V_{\mu\nu} V^{\mu\nu}]$$

= 0  $\rightarrow$  massless photon, massless gluon!

$\rightarrow$  log-UV corrections to W, Z masses!

# SM/UV Hierarchy Is Preserved

While  $(\Lambda_{UV}^2 - \Lambda_{SM}^2)g_{\mu\nu} \iff R_{\mu\nu}(g)$

▶  $c_V\left(\frac{\Lambda_{SM}}{\Lambda_{UV}}\right)$  is held unchanged

▶ hence  $\frac{\Lambda_{SM}}{\Lambda_{UV}}$

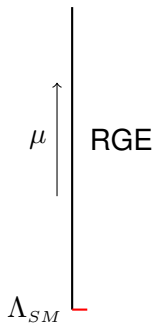
▶ and so are  $c_O\left(\frac{\Lambda_{SM}}{\Lambda_{UV}}\right)$  and  $c_H\left(\frac{\Lambda_{SM}}{\Lambda_{UV}}\right)$

# Log UV Contributions Are As Usual (Dim. Reg.)

While  $(\Lambda_{UV}^2 - \Lambda_{SM}^2)g_{\mu\nu} \iff R_{\mu\nu}(g)$

▶  $\frac{\Lambda_{SM}}{\Lambda_{UV}}$  is held unchanged

▶ so is  $\delta S_{log}(\psi_{SM}, \Lambda_{SM} \log \frac{\Lambda_{SM}}{\Lambda_{UV}})$



▶ **Dim. Reg.:**  $\log \frac{\Lambda_{SM}}{\Lambda_{UV}} \equiv -\frac{1}{\epsilon} - \log \frac{\mu}{\Lambda_{SM}}$



# Power-Law UV Contributions Set Curvature Sector

$$\delta S_{\text{power}} \xrightarrow[\substack{c_O, c_H \text{ held unchanged}}]{(\Lambda_{UV}^2 - \Lambda_{SM}^2) g_{\mu\nu} \mapsto R_{\mu\nu}(g)}$$

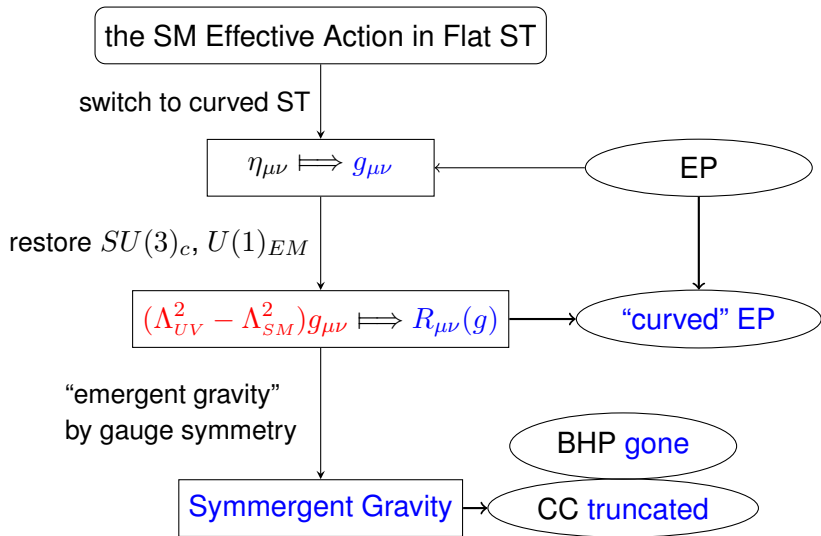
$$- \int d^4x \sqrt{-g} \underbrace{\frac{c_O}{4} (\Lambda_{UV}^2 + \Lambda_{SM}^2)}_{\frac{M_{Pl}^2}{2}} R(g)$$

quartic contribution  
to vacuum energy  
turns into **EH term!**

$$- \int d^4x \sqrt{-g} \frac{c_H}{4} H^\dagger H R(g)$$

big hierarchy  
problem is **gone!**

# “Symmergent Gravity”



# Symmergent Gravity vs. Induced Gravity

Table 1: Contrasting symmergent gravity with Sakharov's induced gravity for  $M_{Pl} \cong \Lambda_{UV}$ .

	Gravity Sector	Cosmological Constant	Higgs Mass
Induced Gravity	EH + HC	$\Lambda_{UV}^4/M_{Pl}^2 \cong M_{Pl}^2$	$\cong M_{Pl}$
Symmergent Gravity	EH	$\Lambda_{SM}^4/M_{Pl}^2 \cong m_\nu^2$	$\cong \Lambda_{SM}$

# Alas! The SM Alone Can't Lead To Proper Gravity

In the SM (at one loop):

$$c_O \cong \frac{(n_b - n_f)}{64\pi^2} = \frac{-62}{64\pi^2}$$

not eligible for  
inducing gravity

introduce **new fields**  
to make  $c_O \gtrsim 1$   
so that  $\Lambda_{UV} \lesssim M_{Pl}$

# A BSM Sector Is Required

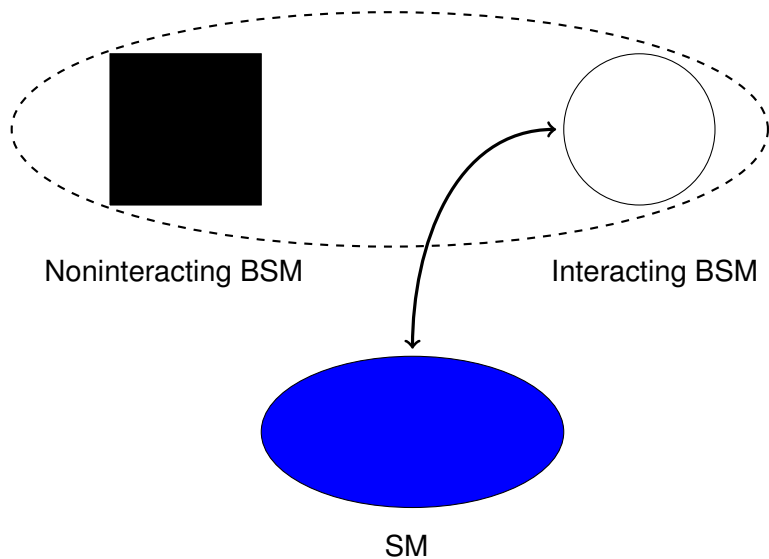
- ▶ In general, BSM **does not have to interact** with the SM but **it can**.

- ▶ All it has to do is to provide some

$$n_b^{BSM} - n_f^{BSM} \gtrsim 128\pi^2 + 62 \approx 1325$$

more bosons than fermions to insure  $\Lambda_{UV} \lesssim M_{Pl}$ .

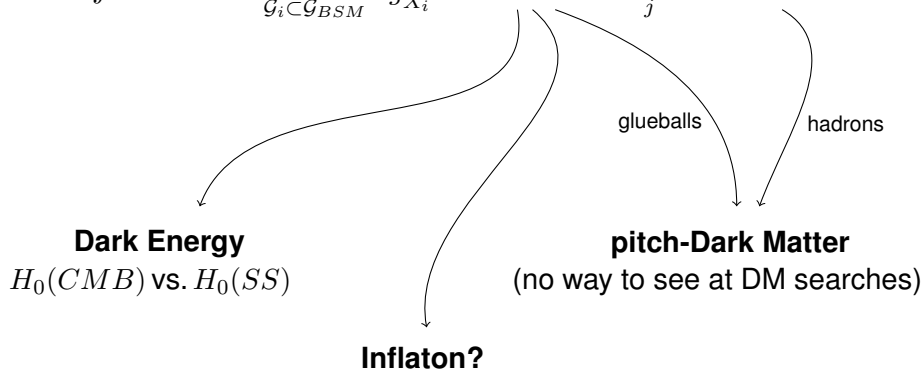
# BSM Can Come In Two Types



# Noninteracting BSM Is Sweet Home For “Dark Stuff”

Non-Abelian gauge theories like  $\mathcal{G}_{BSM} = SO(51), SU(26), E(8)^3, \dots$

$$S^{(ni)} = \int d^4x \sqrt{-\eta} \left\{ - \sum_{\mathcal{G}_i \subset \mathcal{G}_{BSM}} \frac{1}{2g_{X_i}^2} \text{Tr} \{ X_{\mu\nu}^i X_i^{\mu\nu} \} + \sum_j \bar{\chi}_j (i\not{D} - m_{\chi_j}) \chi_j \right\}$$



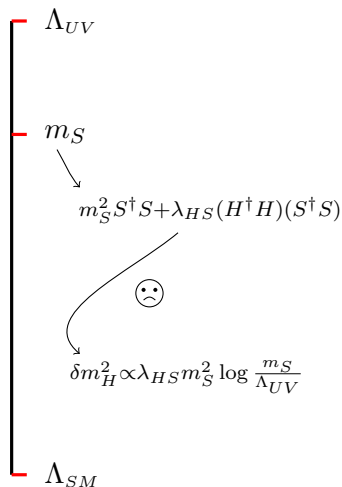
# Interacting BSM Is Needed

- ▶ No evidence for interacting BSM at LHC and DM searches.
- ▶ But it is necessary for Strong CP, Flavor, Baryogenesis, Inflation, ... whose realizations can involve some heavy scalars  $S$ .
- ▶  $\delta m_H^2$  can be suppressed only if

$$\lambda_{HS} \lesssim \Lambda_{SM}^2 / m_S^2$$

as obeyed by phenomena above.  
(SUSY, Xtra Dim, Comp. don't!)

- ▶ Majorana neutrinos:  $\Lambda_{UV} = m_N$   
(with a crowded noninteracting BSM)





# IMPLICATION 1

SM+Gravity+Noninteracting-BSM:

$$S_{eff}^{(SM+BSM)} = \int d^4x \sqrt{-g} \left\{ -\frac{1}{2} M_{Pl}^2 R(g) - \frac{c_H}{4} H^\dagger H R(g) \right\}$$
$$+ S_{tree}(\psi_{SM}) + \delta S_{log}(\psi_{SM}, \Lambda_{SM} \log \frac{\Lambda_{SM}}{\Lambda_{UV}})$$
$$+ S_{tree}^{(BSM)}(\psi_{BSM}) + \delta S_{log}^{(BSM)}(\psi_{BSM}, \Lambda_{BSM} \log \frac{\Lambda_{BSM}}{\Lambda_{UV}})$$

this setup can account for all of the available  
data and constraints except for one thing: **CCP!**

# IMPLICATION 2

- ▶ The real **challenge** is to understand how CC can be reduced from  $\Lambda_{BSM}^4/M_{Pl}^2$  down to  $H_0^2$ .
- ▶ Solution of the CCP can **reveal** structure of the BSM.  $CC_{SM} + CC_{BSM} \cong H_0^2$  can put highly non-trivial constraints on BSM (coupling to SM is not a must).
- ▶ The CCP may well be a **door** to sought new physics. (BSM, if noninteracting, can lie at ...TeV's, GeV's, MeV's, ...)

# IMPLICATION 3

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- ▶ LHC or other colliders **do not have to discover** any new particle.
- ▶ If they **discover any** then BSM needs be structured accordingly.

# In Conclusion

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- ▶ We should analyze any given experimental signal without any **theoretical prejudice**.
- ▶ We should keep in mind that an experimental signal **does not** have to arise.
- ▶ We should focus on **the CCP** to reveal structure of the BSM.
- ▶ We should explore **trans-UV** physics in regard to SM+BSM setup.

# Thank You For Your Attention

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## References:

- ▶ *Naturalizing Gravity of the Quantum Fields, and the Hierarchy Problem*, [arXiv:1703.05733](#)
- ▶ *Curvature-Restored Gauge Invariance and Ultraviolet Naturalness*, [arXiv:1605.00377](#)
- ▶ *A Mechanism of Ultraviolet Naturalness*, [arXiv:1510.05570](#)

# Acknowledgement

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