

Interference effects in vacuum hadron polarization...

Vladimír Šauli
NIP, Rez near Prague
sauli@ujf.cas.cz

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Outline

- Introduction
- Theory of $e^+e^- \rightarrow \mu^+\mu^-$ for KLOE
- Results and extraction of Π_h
- Running QED coupling at energy of ρ (KLOE 2016)

Introduction

based on : "Hadronic vacuum polarization in $e^+e^- \rightarrow \mu^+\mu^-$ process below 3 GeV, arXiv 1704.01887"

Test of SM is based on comparison between theory and experiment.

$a_\mu, \alpha_{M_Z}, \dots$

Π_h plays crucial role

$$\Pi^{\mu\nu}(x) = \sum_q \text{Tr} \langle j_q^\mu(x) j_q^\nu(0) \rangle$$

$$\Pi^{\mu\nu}(q) = P_T^{\mu\nu} \Pi_h(q^2) = \text{Tr} \int_k \Gamma^\mu(k, p) S_q(k) \gamma^\mu S_q(p)$$

Theory of $e^+e^- \rightarrow \mu^+\mu^-$ for KLOE

Interference effects between hadronic and leptonic vacuum polarization are observed in close vicinity of narrow resonances: $J/\Psi, \Psi, \Upsilon$ as well as ϕ and ω .

They have unique interference pattern and appears in all predominantly QED process in the timelike momentum channel. 1963 estimate : In vicinity of resonances the charges changes like :

$$e^2 \rightarrow e^2 \left(1 - \frac{3m_V \Gamma_{ll} \alpha^{-1}}{m_V^2 - s + im\Gamma} \right), \quad (1)$$

Recently ϕ meson sector measured by KLOE in 2004 and SND 2001.

Experiment. data: F. A. Ambrosino, et. al. Phys. Lett. B 608 (2005). see also M.N. Achasov NPA 675, 2000, PRL 86 (2001)

Using $+\gamma$ process the running coupling has been measured by KLOE 2015.

To see the effect, the high precision measurement is required

$$\epsilon(= \sigma_{stat})(m_\phi) = 0.1nb \text{ with KLOE experimental setup, } \sigma(m_\phi) \equiv 40nb.$$

Using SM, the cross section reads (in next to (next) leading order)

$$\sigma(s) = \frac{4\pi C_t}{|1 - \Pi(s)|^2} \left[\sigma_A(s) \left(2 - \beta_\mu^2 \left(1 - \frac{C_t^2}{3} \right) \right) + \sigma_B(s) \right], \quad (2)$$

where $C_t = \cos(\theta_{min})$ with $\theta_{min} = 50^\circ$ ($\theta_{max} = 140^\circ$), which is KLOE experimental cut on polar scattering angle between μ^- and e^- particles and $\beta_\mu = \sqrt{1 - 4m_\mu^2/s}$.

$$\sigma_B(s) = -\frac{\alpha^3}{4\pi s} (1 - \beta_\mu^2) \ln \frac{1 + \beta_\mu}{1 - \beta_\mu}. \quad (3)$$

$\sigma_A(s)$ - the dominant term, listed completely in Arbuzov1997 (9702262), collects all leading logs of Dirac and Pauli form factors and the known soft photon contributions for which we take $\ln \frac{\Delta\epsilon}{\epsilon} = 0.05$ (15 MeV cut on c.m.s. soft photon energy at ϕ peak).

Theory of $e^+e^- \rightarrow \mu^+\mu^-$ for KLOE

Nontrivial information is in $\Pi(s) = \Pi_l(s) + \Pi_h(s)$

which completes $\alpha(s) = \frac{\alpha}{1-\Pi(s)}$

$\alpha = \alpha(0) = 1/137.0359991390$

$\Pi_l(s)$ - includes complete one loop + leading log in $(\alpha)^2$

hadronic vacuum polarization Π_h -driven by strong coupling QCD

$$\Pi_h(s) = \frac{s}{4\pi^2\alpha} \int_{4m_\pi^2}^{\infty} d\omega \frac{\sigma_h(\omega) \left[\frac{\alpha}{\alpha(\omega)} \right]^2}{\omega - s + i\epsilon}. \quad (4)$$

$\sigma_h = \sigma(e^+e^- \rightarrow \text{hadrons})$

Theory of $e^+e^- \rightarrow \mu^+\mu^-$ for KLOE

For σ_h dominant below 3GeV we use

this millennium published data for exclusive channels:

$e^+e^- \rightarrow \pi\pi, K^+K^-, K_LK_S, \pi\pi\pi$ and $\eta\gamma$

from **SND, CMD, CMD2, KLOE, BABAR, BESSIII** collab.

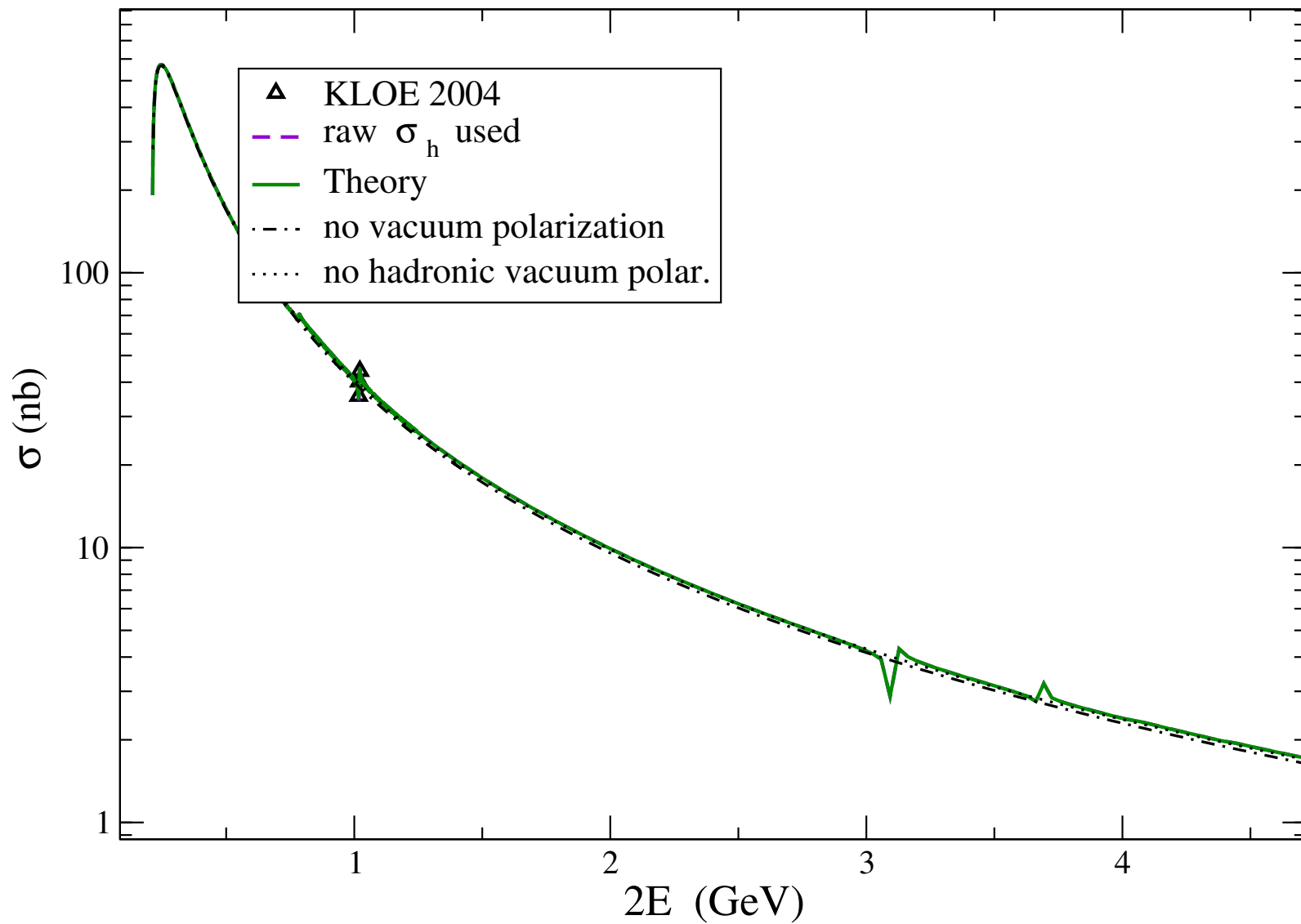
For each channel we find fit valid for all s and generate quasidata, from which we calculate Π_h .

Existing fits for σ_i and ϵ_i are explored and refitted.

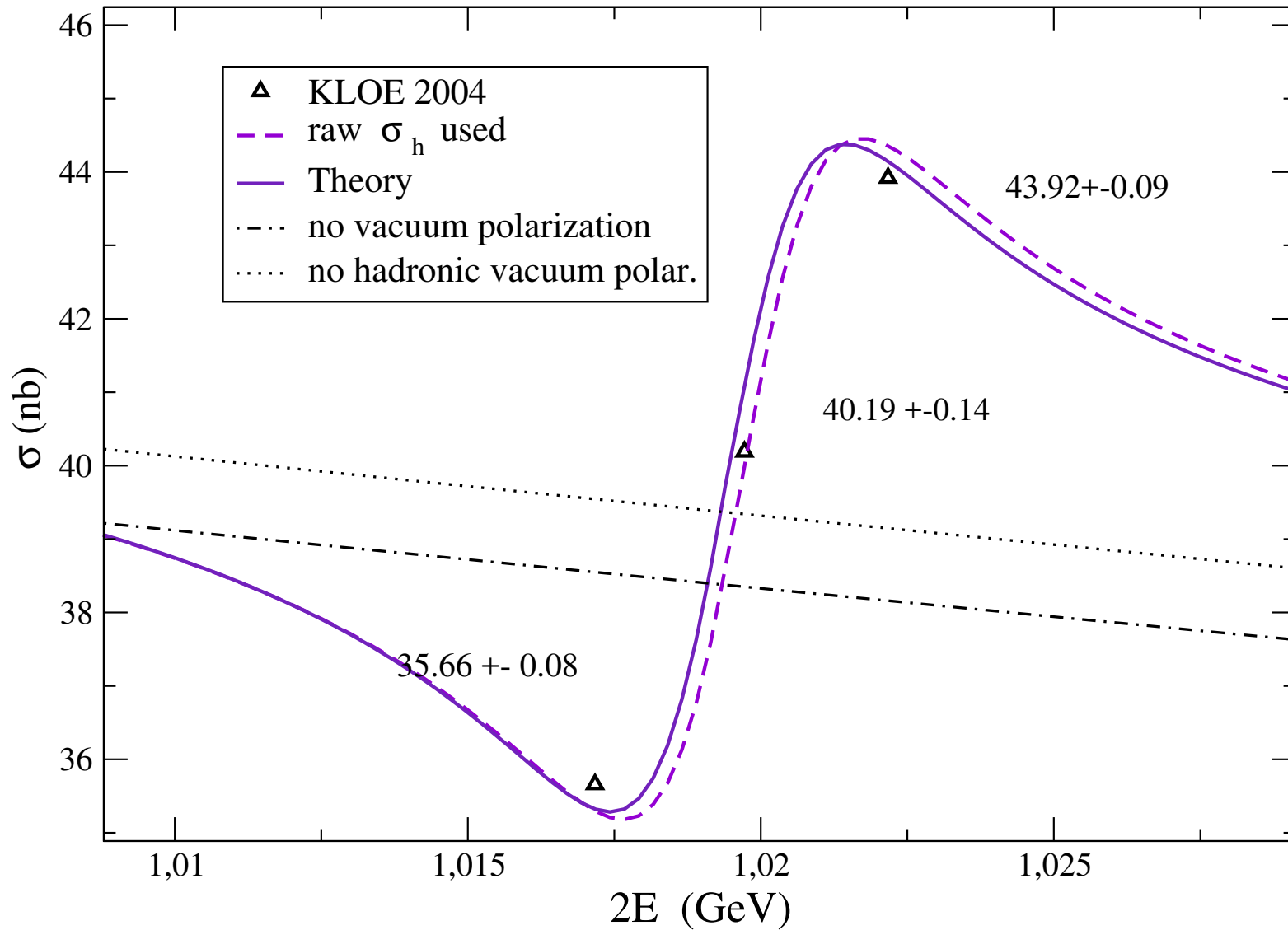
$\sigma_h^0 = \sigma_h(s)(|1 - \Pi(s)|)^2$ calculated

Above 3 GeV all important narrow quarkonia are included also 1 loops PT QCD for b,c are added for all s (almost no effect)

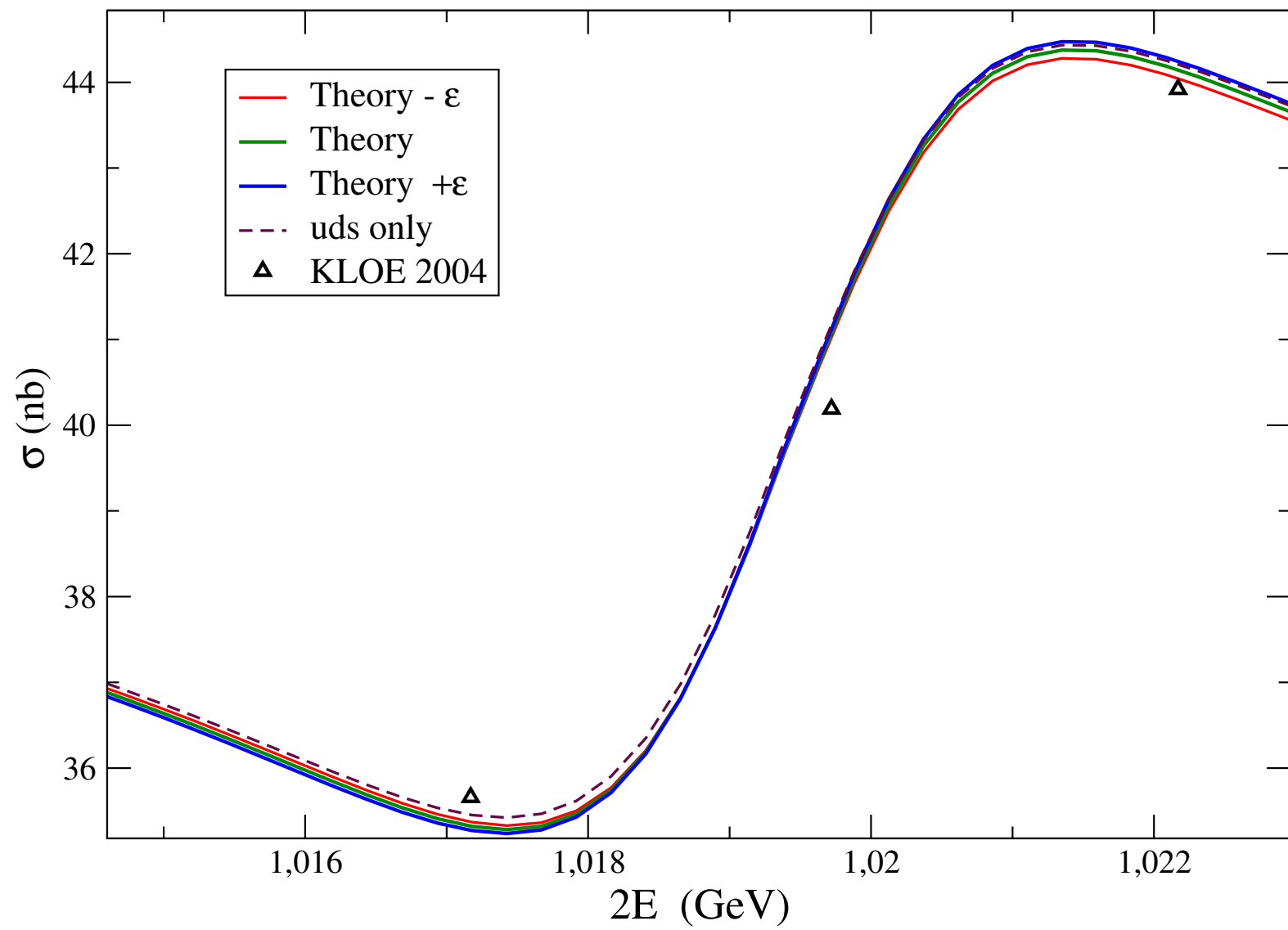
$e^+e^- \rightarrow \mu^+\mu^-(\gamma)$ for KLOE



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$e^+e^- \rightarrow \mu^+\mu^-(\gamma)$ for KLOE



Fits of σ_h

Sum of (un)dressed BW with complex phases

Gounaris-Sakurai VMD, Kuhn-Santamaria model

BABAR 2012 fit $\pi\pi$

modified BABAR fits 2013 for K^+K^-

modified SND fit for $K^L K^S$

own fits for 3π and $\eta\gamma$

fits of errors $\epsilon(s)$ of $\sigma_h(s)$ overestimating fit formula

$$\epsilon[nb] = c\sqrt{\sigma_h[nb]} ; c = 0.8nb^{1/2} \quad (5)$$

such that $\epsilon > \sigma_{TOT} = \sqrt{\sigma_{syst}^2 + \sigma_{stat}^2}$, where σ_{TOT} are the best data in given kinematical region.

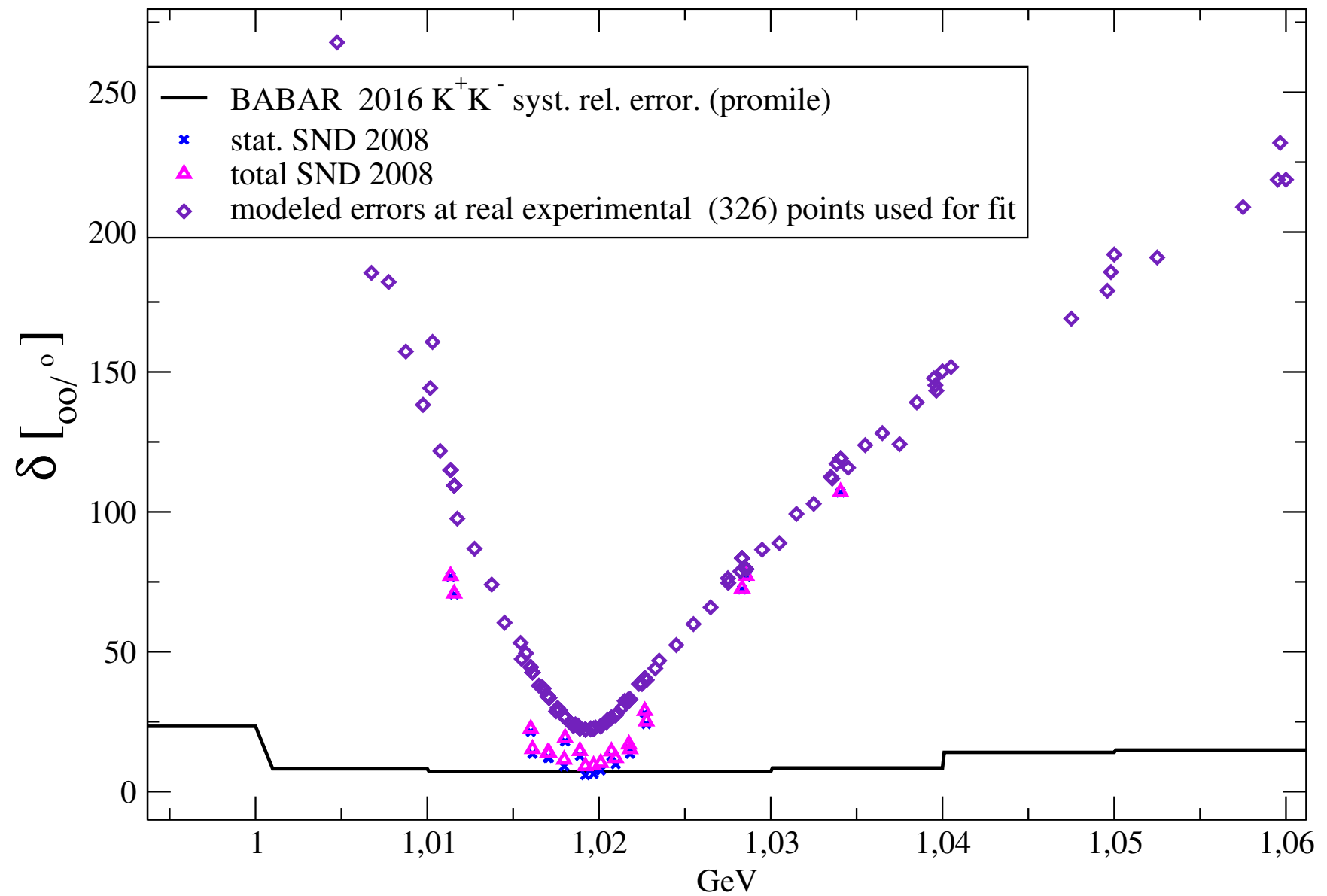
Requirements for minimized χ^2

$\chi^2 \simeq 1$ for single data, $\chi^2(\sigma_{stat} \rightarrow \epsilon) \simeq 1$ for combined data, note usually $\chi^2(\sigma_{stat} \rightarrow \epsilon) \ll 1$ for single data

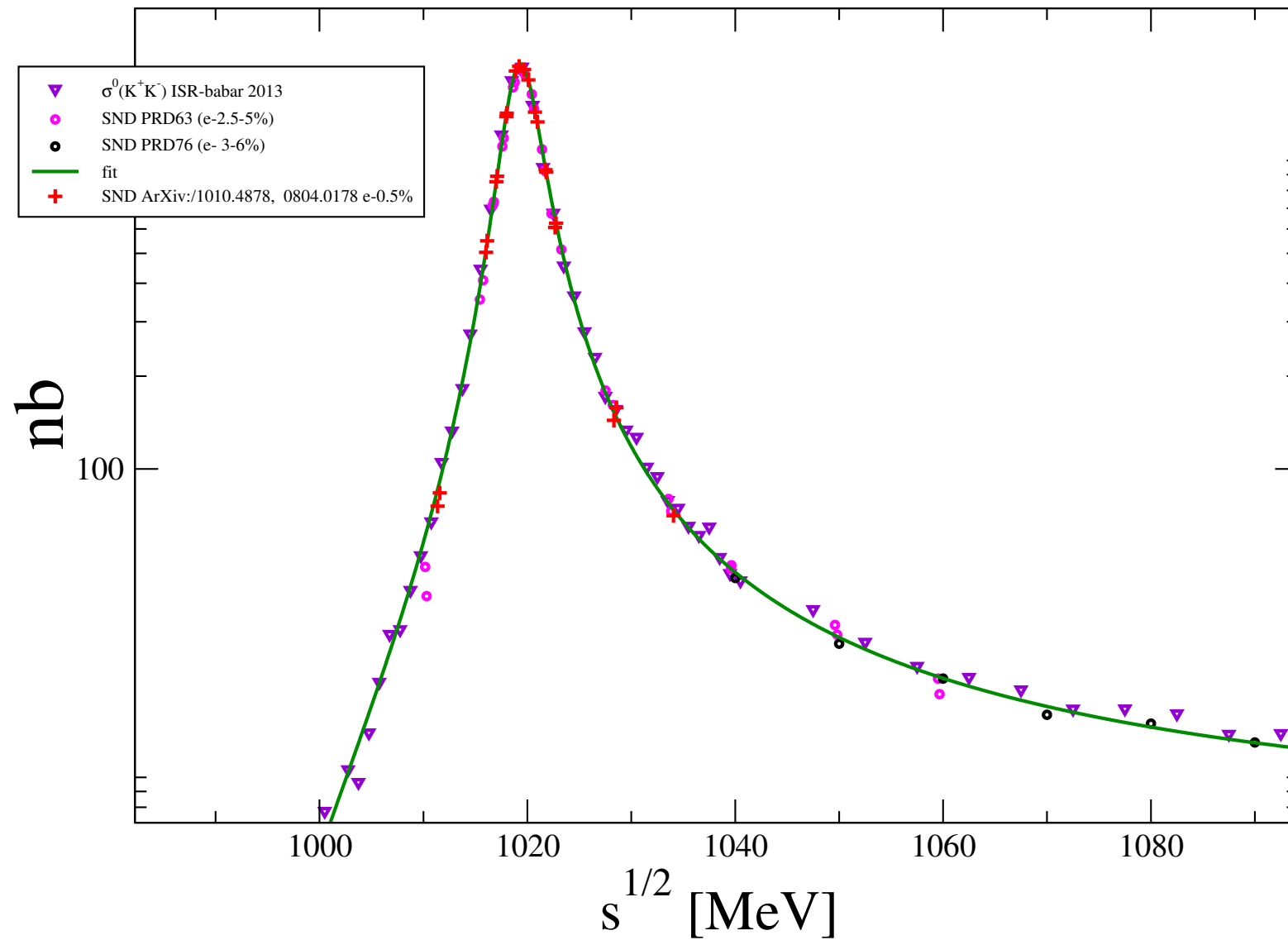
Example $\chi^2(\sigma_{stat} \rightarrow \epsilon) = 0.123$ for BABAR 2013 K^+K^- and $\chi^2(\sigma_{stat} \rightarrow \epsilon) = 0.3$ for combined data

Fits of $\sigma_{K^+K^-}, \dots$

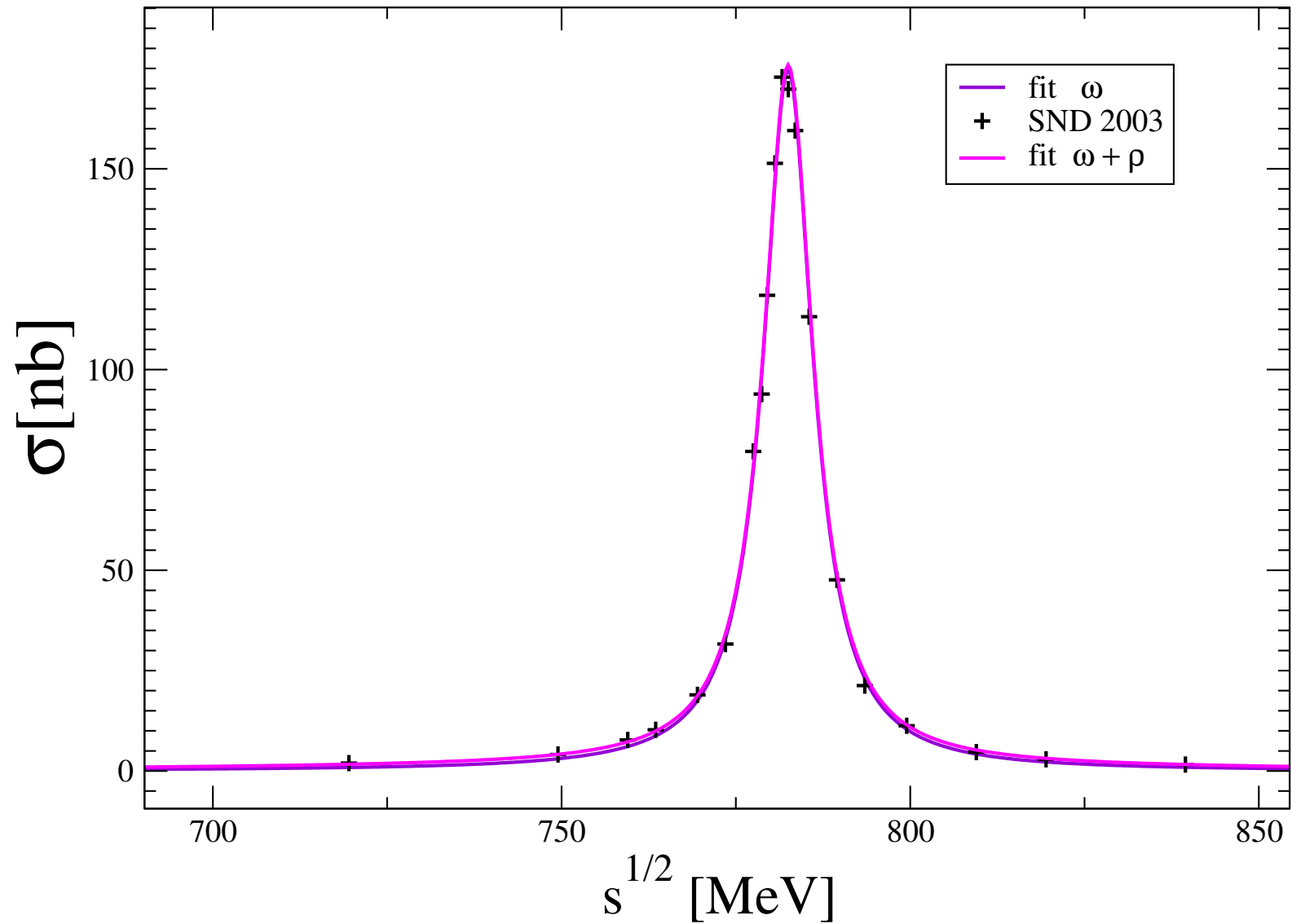
Relative errors of $K^+ K^-$ mode



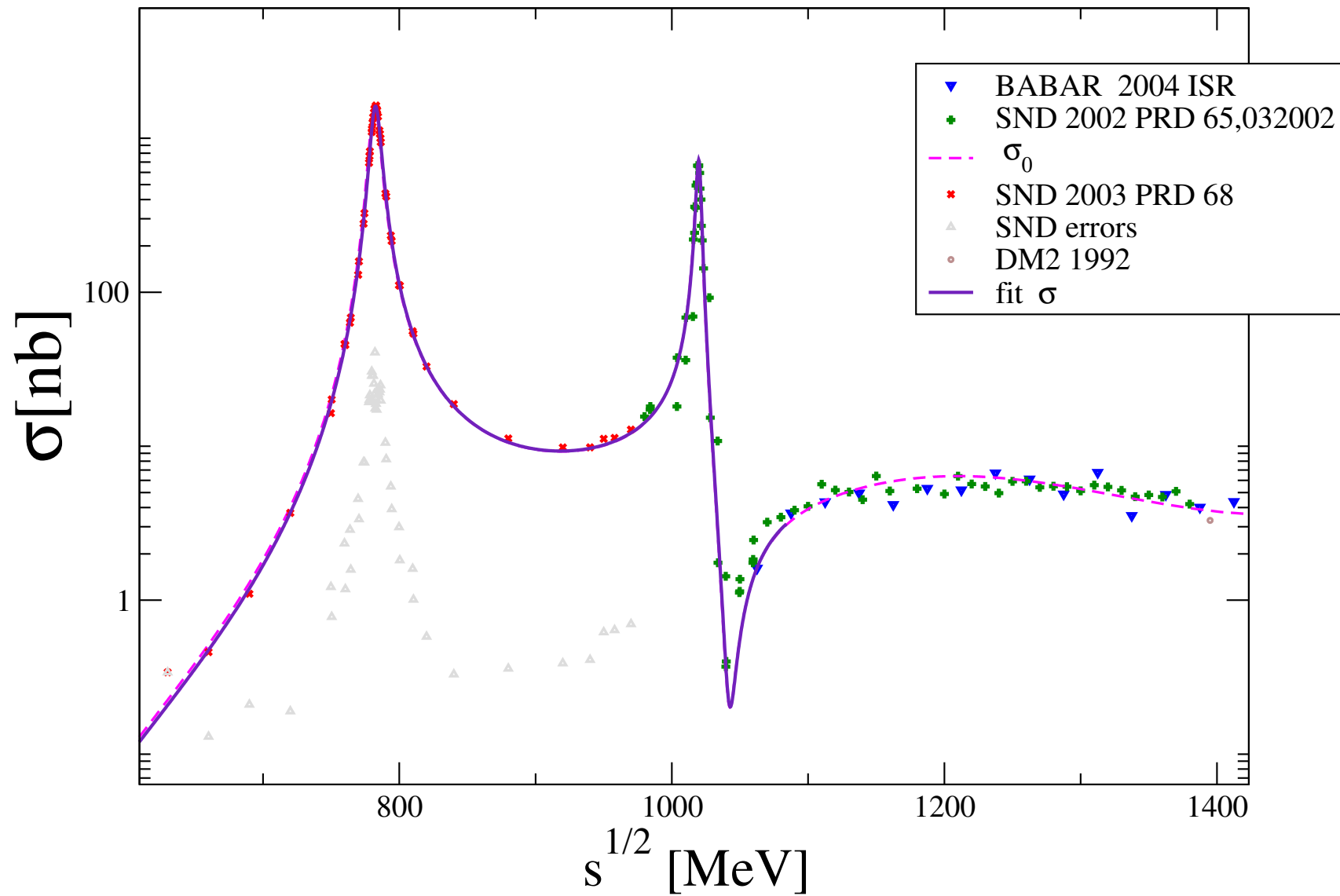
$$\sigma(e^+e^- \rightarrow K^+K^-)$$



$\sigma(e^+e^- \rightarrow \eta\gamma)$



$$\sigma (e^+ e^- \rightarrow \pi\pi\pi)$$

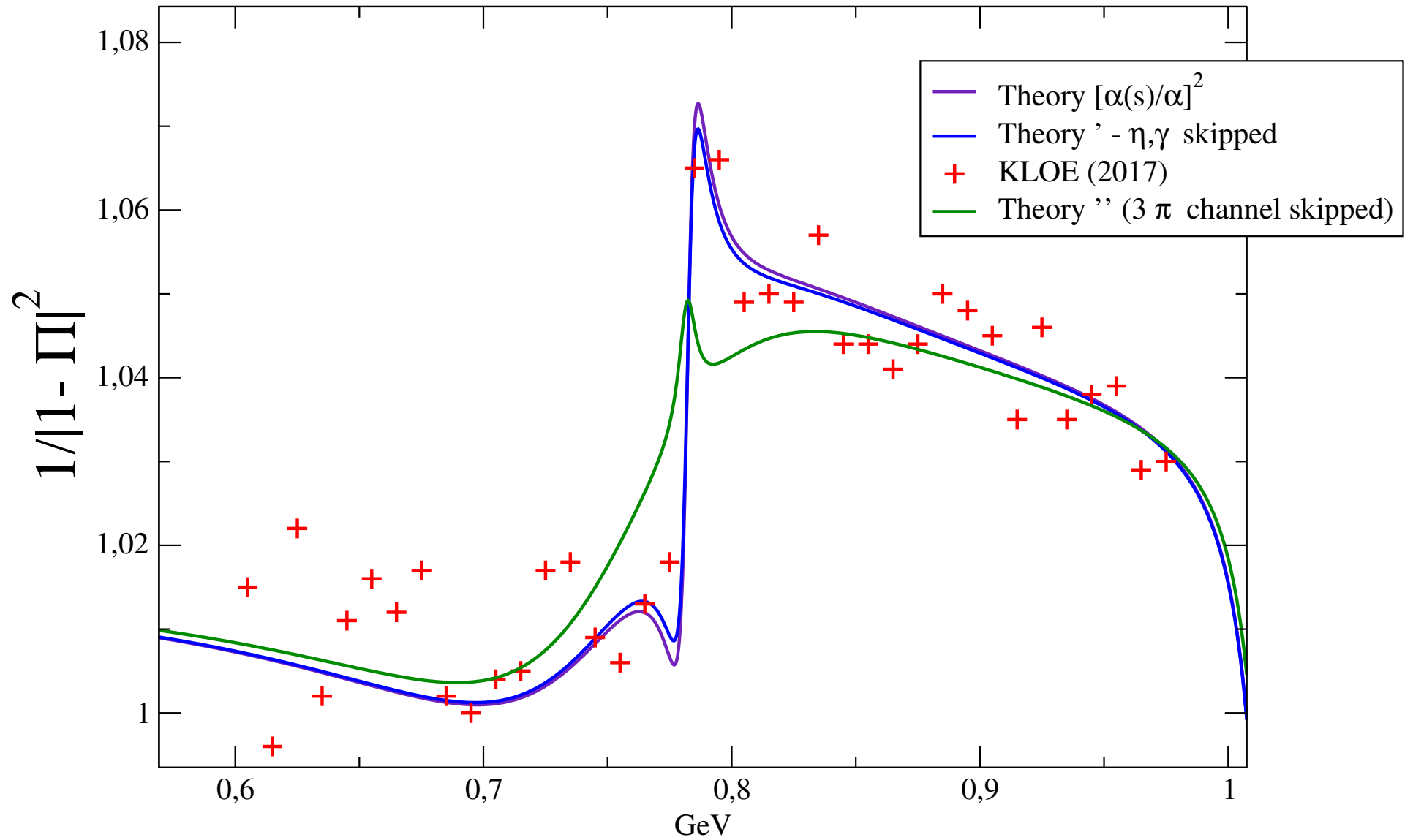


Running α at ω

Never measured precisely in muon pair production

2016 KLOE collaboration studied process $ee \rightarrow \mu\mu\gamma(\gamma)$

KLOE 2016 α from $ee \rightarrow \mu\mu\gamma(\gamma)$



Conclusion

α_{QED} is complex at the timelike scale and runs in accordance with SM with

small, but statistically significant tension between Standard Theory and $e^+e^- \rightarrow \mu^+\mu^-$ 2004 KLOE experiment.

Explanation:

1. σ_h is systematically overestimated (few percentage, s -dependent subtraction would be needed)
2. Analyticity is not an exact property of Π_h (there can be a contribution with different singular structure than a real axis poles and real branch points, their contribution should not be larger than 5%)
3. Speaker, you are listening now, made a bad mistake
4. Something else
5. Combination of 1,2,3,4