

QCD from quark, gluon, and meson correlators

Mario Mitter

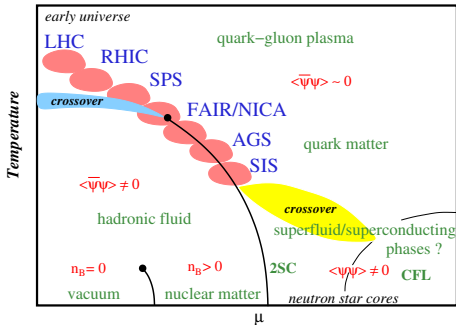
Karl-Franzens-Universität Graz

Sintra, May 2017



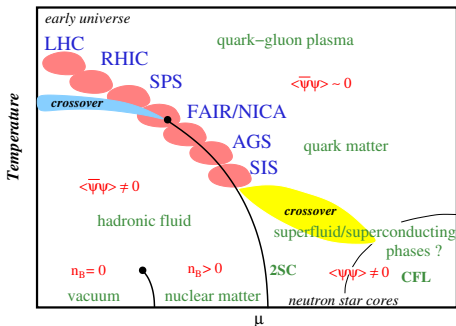
fQCD collaboration - QCD (phase diagram) with FRG:

J. Braun, L. Corell, A. K. Cyrol, W. J. Fu, M. Leonhardt, MM,
J. M. Pawłowski, M. Pospiech, F. Rennecke, N. Strodthoff, N. Wink, ...



fQCD collaboration - QCD (phase diagram) with FRG:

J. Braun, L. Corell, A. K. Cyrol, W. J. Fu, M. Leonhardt, MM,
J. M. Pawłowski, M. Pospiech, F. Rennecke, N. Strodthoff, N. Wink, ...



large part of this effort: vacuum YM-theory and QCD

QCD from the effective action

$$\Gamma[\Phi] = \sum_n \int_{\{p_i\}} \Gamma_{\Phi_1 \dots \Phi_n}^{(n)}(p_1, \dots, p_{n-1}) \Phi^1(p_1) \cdots \Phi^n(-p_1 - \cdots - p_{n-1})$$

QCD from the effective action

$$\Gamma[\Phi] = \sum_n \int_{\{p_i\}} \Gamma_{\Phi_1 \dots \Phi_n}^{(n)}(p_1, \dots, p_{n-1}) \Phi^1(p_1) \cdots \Phi^n(-p_1 - \cdots - p_{n-1})$$

- full information about QFT encoded in $\Gamma[\Phi]$ /correlators:

QCD from the effective action

$$\Gamma[\Phi] = \sum_n \int_{\{p_i\}} \Gamma_{\Phi_1 \dots \Phi_n}^{(n)}(p_1, \dots, p_{n-1}) \Phi^1(p_1) \dots \Phi^n(-p_1 - \dots - p_{n-1})$$

- full information about QFT encoded in $\Gamma[\Phi]$ /correlators:

- ▶ bound state spectrum: pole structure of the $\Gamma^{(n)}$

e.g. [Roberts, Williams, '94], [Alkofer, Smekal, '00], [Eichmann, Sanchis-Alepuz, Williams, Alkofer, Fischer, '16]

- ▶ form factors: photon-particle correlators

e.g. [Cloet, Eichmann, El-Bennich, Klahn, Roberts, '08], [Sanchis-Alepuz, Williams, Alkofer, '13]

⇒ decay constants

e.g. [Maris, Roberts, Tandy, '97], [MM, Pawłowski, Strodthoff, in prep.]

QCD from the effective action

$$\Gamma[\Phi] = \sum_n \int_{\{p_i\}} \Gamma_{\Phi_1 \dots \Phi_n}^{(n)}(p_1, \dots, p_{n-1}) \Phi^1(p_1) \cdots \Phi^n(-p_1 - \cdots - p_{n-1})$$

- full information about QFT encoded in $\Gamma[\Phi]$ /correlators:

- ▶ bound state spectrum: pole structure of the $\Gamma^{(n)}$

e.g. [Roberts, Williams, '94], [Alkofer, Smekal, '00], [Eichmann, Sanchis-Alepuz, Williams, Alkofer, Fischer, '16]

- ▶ form factors: photon-particle correlators

e.g. [Cloet, Eichmann, El-Bennich, Klahn, Roberts, '08], [Sanchis-Alepuz, Williams, Alkofer, '13]

⇒ decay constants

e.g. [Maris, Roberts, Tandy, '97], [MM, Pawłowski, Strodthoff, in prep.]

- ▶ thermodynamic quantities: $\Gamma[\Phi] \propto$ grand potential

- ★ equation of state

e.g. [Herbst, MM, Pawłowski, Schaefer, Stiele, '13]

- ★ fluctuations of conserved charges

e.g. [Fu, Rennecke, Pawłowski, Schaefer, '16]

QCD from the effective action

$$\Gamma[\Phi] = \sum_n \int_{\{p_i\}} \Gamma_{\Phi_1 \dots \Phi_n}^{(n)}(p_1, \dots, p_{n-1}) \Phi^1(p_1) \cdots \Phi^n(-p_1 - \cdots - p_{n-1})$$

- full information about QFT encoded in $\Gamma[\Phi]$ /correlators:

- ▶ bound state spectrum: pole structure of the $\Gamma^{(n)}$

e.g. [Roberts, Williams, '94], [Alkofer, Smekal, '00], [Eichmann, Sanchis-Alepuz, Williams, Alkofer, Fischer, '16]

- ▶ form factors: photon-particle correlators

e.g. [Cloet, Eichmann, El-Bennich, Klahn, Roberts, '08], [Sanchis-Alepuz, Williams, Alkofer, '13]

⇒ decay constants

e.g. [Maris, Roberts, Tandy, '97], [MM, Pawłowski, Strodthoff, in prep.]

- ▶ thermodynamic quantities: $\Gamma[\Phi] \propto$ grand potential

- ★ equation of state

e.g. [Herbst, MM, Pawłowski, Schaefer, Stiele, '13]

- ★ fluctuations of conserved charges

e.g. [Fu, Rennecke, Pawłowski, Schaefer, '16]

- ▶ further quantities: $\Gamma[\Phi] \propto$ eff. potential, propagators, 't Hooft determinant

- ★ chiral condensate(s) $\langle \sigma \rangle$

e.g. [Schaefer, Wambach '04], [Fischer, Luecker, Mueller '11], [MM, Schaefer, '13]

- ★ (dressed) Polyakov loop

e.g. [Braun, Gies, Pawłowski, '07], [Fischer, '09], [Braun, Haas, Marhauser, Pawłowski, '09]

- ★ axial anomaly

e.g. [Grah, Rischke, '13], [MM, Schaefer, '13], [Fejos, '15], [Heller, MM, '15]

- ★ spectral functions

e.g. [Tripolt, Strodthoff, Smekal, Wambach, '14]

Correlation functions with the FRG

- $S[\Phi] = \Gamma_\Lambda[\Phi]$: use only perturbative QCD input
 - ▶ $\alpha_S(\Lambda = \mathcal{O}(10) \text{ GeV})$
 - ▶ $m_q(\Lambda = \mathcal{O}(10) \text{ GeV})$

Correlation functions with the FRG

- $S[\Phi] = \Gamma_\Lambda[\Phi]$: use only perturbative QCD input
 - ▶ $\alpha_S(\Lambda = \mathcal{O}(10) \text{ GeV})$
 - ▶ $m_q(\Lambda = \mathcal{O}(10) \text{ GeV})$
- integration of momentum shells:

$$\partial_k \Gamma_k[A, \bar{c}, c, \bar{q}, q] = \frac{1}{2} \left(\text{Diagram 1} - \text{Diagram 2} - \text{Diagram 3} \right)$$


Correlation functions with the FRG

- $S[\Phi] = \Gamma_\Lambda[\Phi]$: use only perturbative QCD input
 - ▶ $\alpha_S(\Lambda = \mathcal{O}(10) \text{ GeV})$
 - ▶ $m_q(\Lambda = \mathcal{O}(10) \text{ GeV})$
- integration of momentum shells:

$$\partial_k \Gamma_k[A, \bar{c}, c, \bar{q}, q] = \frac{1}{2} \left(\text{Diagram 1} - \text{Diagram 2} - \text{Diagram 3} \right)$$


\Rightarrow full quantum effective action

Correlation functions with the FRG

- $S[\Phi] = \Gamma_\Lambda[\Phi]$: use only perturbative QCD input
 - ▶ $\alpha_S(\Lambda = \mathcal{O}(10) \text{ GeV})$
 - ▶ $m_q(\Lambda = \mathcal{O}(10) \text{ GeV})$
- integration of momentum shells:

$$\partial_k \Gamma_k[A, \bar{c}, c, \bar{q}, q] = \frac{1}{2} \left(\text{Diagram 1} - \text{Diagram 2} - \text{Diagram 3} \right)$$


⇒ full quantum effective action

- gauge-fixed approach (Landau gauge): ghosts appear

Correlation functions with the FRG

- $S[\Phi] = \Gamma_\Lambda[\Phi]$: use only perturbative QCD input
 - ▶ $\alpha_S(\Lambda = \mathcal{O}(10) \text{ GeV})$
 - ▶ $m_q(\Lambda = \mathcal{O}(10) \text{ GeV})$
- integration of momentum shells:

$$\partial_k \Gamma_k[A, \bar{c}, c, \bar{q}, q] = \frac{1}{2} \left(\text{Diagram 1} - \text{Diagram 2} - \text{Diagram 3} \right)$$

\Rightarrow full quantum effective action

- gauge-fixed approach (Landau gauge): ghosts appear
- functional derivatives \Rightarrow equations for correlators
- aim for “apparent convergence” of $\Gamma[\Phi] = \lim_{k \rightarrow 0} \Gamma_k[\Phi]$

Some representative equations

[MM, Strodthoff, Pawlowski, 2014],

[Cyrol, Fister, MM, Strodthoff, Pawlowski, 2016]

$$\partial_t \text{---}^{-1} = \text{---}^{\text{---}} + \text{---}^{\text{---}} + \frac{1}{2} \text{---}^{\text{---}} + \text{---}^{\text{---}} + \text{---}^{\text{---}} - \text{---}^{\text{---}}$$

$$\partial_t \text{---}^{\text{---}} = - \text{---}^{\text{---}} - \text{---}^{\text{---}} - \text{---}^{\text{---}} - \text{---}^{\text{---}} - \frac{1}{2} \text{---}^{\text{---}} - \text{---}^{\text{---}} + 2 \text{---}^{\text{---}} - \text{---}^{\text{---}} + \text{perm.}$$

$$\partial_t \text{---}^{\text{---}} = 2 \text{---}^{\text{---}} - \text{---}^{\text{---}} - \text{---}^{\text{---}} - \text{---}^{\text{---}} - \text{---}^{\text{---}} - \text{---}^{\text{---}} - \text{---}^{\text{---}} - \text{---}^{\text{---}} - \text{---}^{\text{---}} + \text{perm.}$$

Some representative equations

[MM, Strodthoff, Pawlowski, 2014],

[Cyrol, Fister, MM, Strodthoff, Pawlowski, 2016]

$$\partial_t \text{---}^{-1} = \text{---} \text{---} \text{---} + \text{---} \text{---} \text{---} + \frac{1}{2} \text{---} \text{---} \text{---} + \text{---} \text{---} \text{---} + \text{---} \text{---} \text{---} - \text{---} \text{---} \text{---}$$

$$\partial_t \text{---}^{-1} = \text{---} \text{---} \text{---} + \text{---} \text{---} \text{---}$$

$$\partial_t \text{---}^{-1} = \text{---} \text{---} \text{---} - 2 \text{---} \text{---} \text{---} + \frac{1}{2} \text{---} \text{---} \text{---}$$

$$\partial_t \text{---} = - \text{---} \text{---} \text{---} - \text{---} \text{---} \text{---} + \text{perm.}$$

$$\partial_t \text{---} = - \text{---} \text{---} \text{---} - \text{---} \text{---} \text{---} - \text{---} \text{---} \text{---} - \text{---} \text{---} \text{---} + 2 \text{---} \text{---} \text{---} - \text{---} \text{---} \text{---} + \text{perm.}$$

$$\partial_t \text{---} = - \text{---} \text{---} \text{---} + 2 \text{---} \text{---} \text{---} - \text{---} \text{---} \text{---} + \text{perm.}$$

$$\partial_t \text{---} = - \text{---} \text{---} \text{---} - \text{---} \text{---} \text{---} + 2 \text{---} \text{---} \text{---} - \text{---} \text{---} \text{---} + \text{perm.}$$

$$\partial_t \text{---} = 2 \text{---} \text{---} \text{---} - \text{---} \text{---} \text{---} - \text{---} \text{---} \text{---} - \text{---} \text{---} \text{---} - \text{---} \text{---} \text{---} - \text{---} \text{---} \text{---} - \text{---} \text{---} \text{---} + \text{perm.}$$

Some representative equations

[MM, Strodthoff, Pawlowski, 2014],

[Cyrol, Fister, MM, Strodthoff, Pawlowski, 2016]

$$\partial_t \text{---}^{-1} = \text{---} \text{---} \text{---} + \text{---} \text{---} \text{---} + \frac{1}{2} \text{---} \text{---} \text{---} + \text{---} \text{---} \text{---} + \text{---} \text{---} \text{---} - \text{---} \text{---} \text{---}$$

$$\partial_t \text{---}^{-1} = \text{---} \text{---} \text{---} + \text{---} \text{---} \text{---}$$

$$\partial_t \text{---}^{-1} = \text{---} \text{---} \text{---} - 2 \text{---} \text{---} \text{---} + \frac{1}{2} \text{---} \text{---} \text{---}$$

$$\partial_t \text{---} = - \text{---} \text{---} \text{---} - \text{---} \text{---} \text{---} + \text{perm.}$$

$$\partial_t \text{---} = - \text{---} \text{---} \text{---} - \text{---} \text{---} \text{---} - \text{---} \text{---} \text{---} - \text{---} \text{---} \text{---} + 2 \text{---} \text{---} \text{---} - \text{---} \text{---} \text{---} + \text{perm.}$$

$$\partial_t \text{---} = - \text{---} \text{---} \text{---} + 2 \text{---} \text{---} \text{---} - \text{---} \text{---} \text{---} + \text{perm.}$$

$$\partial_t \text{---} = - \text{---} \text{---} \text{---} - \text{---} \text{---} \text{---} + 2 \text{---} \text{---} \text{---} - \text{---} \text{---} \text{---} + \text{perm.}$$

$$\partial_t \text{---} = 2 \text{---} \text{---} \text{---} - \text{---} \text{---} \text{---} - \text{---} \text{---} \text{---} - \text{---} \text{---} \text{---} - \text{---} \text{---} \text{---} - \text{---} \text{---} \text{---} - \text{---} \text{---} \text{---} + \text{perm.}$$

$$\partial_t \text{---}^{-1} = -2 \text{---} \text{---} \text{---} + \text{---} \text{---} \text{---} + \frac{1}{2} \text{---} \text{---} \text{---}$$

$$\partial_t \text{---} = - \text{---} \text{---} \text{---} - \text{---} \text{---} \text{---} - \text{---} \text{---} \text{---} + 2 \text{---} \text{---} \text{---} + \text{perm.}$$

Landau gauge QCD

- two crucial phenomena: $S\chi$ SB and confinement
- similar scales - hard to disentangle

Landau gauge QCD

- two crucial phenomena: $S\chi$ SB and confinement
- similar scales - hard to disentangle

- quenched matter part

[MM, Strodthoff, Pawłowski, 2014]

- pure YM-theory

[Cyrol, Fister, MM, Pawłowski, Strodthoff, 2016]

- outlook:

[Cyrol, MM, Strodthoff, Pawłowski, in preparation]

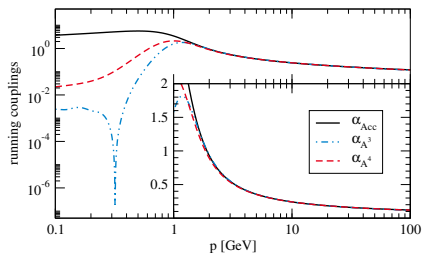
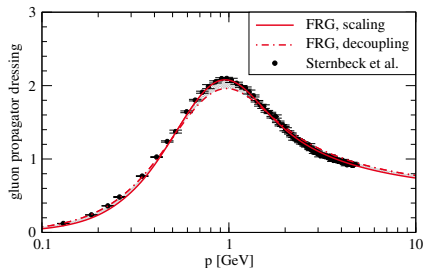
- ▶ unquenching
- ▶ YM-theory at finite temperature $T > 0$

Landau gauge QCD

- two crucial phenomena: $S\chi$ SB and confinement
- similar scales - hard to disentangle
- quenched matter part [MM, Strodthoff, Pawłowski, 2014]
- pure YM-theory [Cyrol, Fister, MM, Pawłowski, Strodthoff, 2016]
- outlook: [Cyrol, MM, Strodthoff, Pawłowski, in preparation]
 - ▶ unquenching
 - ▶ YM-theory at finite temperature $T > 0$
- use results from lattice QCD to gauge truncation

- truncation: momentum dependent dressing functions for all classical tensors
- hardest part of solution: fulfilling the modified STI (\Rightarrow scaling solution)

- truncation: momentum dependent dressing functions for all classical tensors
 - hardest part of solution: fulfilling the modified STI (\Rightarrow scaling solution)
-
- $\Gamma_{AA}^{(2)}(p) \propto Z_A(p) p^2 (\delta^{\mu\nu} - p^\mu p^\nu / p^2)$
 - IR-suppression \Leftrightarrow “confinement”
 - smooth transition to perturbation theory
-
- running couplings
 - degeneracy at large p due to STI
 - test of truncation



lattice data: A. Sternbeck, E. M. Ilgenfritz, M. Muller-Preussker, A. Schiller, and I. L. Bogolubsky, PoS LAT2006, 076.

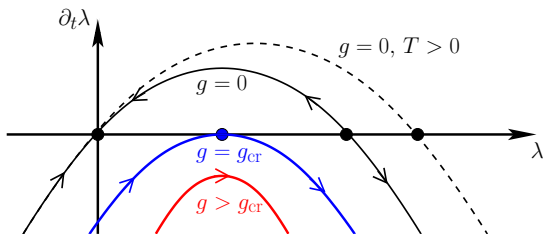
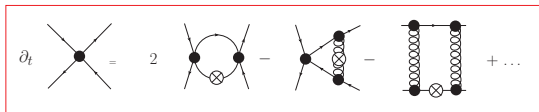
Chiral symmetry breaking

- χ SB \Leftrightarrow resonance in 4-Fermi interaction λ (pion pole):

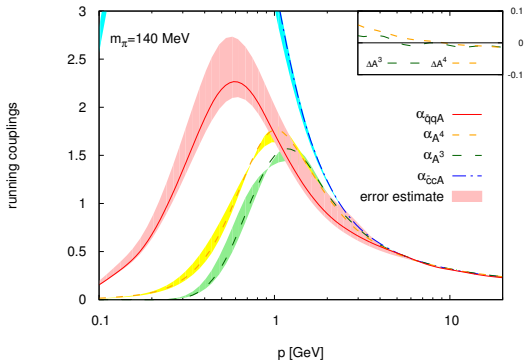
Chiral symmetry breaking

- χ SB \Leftrightarrow resonance in 4-Fermi interaction λ (pion pole):
- β -function of momentum independent 4-Fermi interaction:

$$\partial_t \lambda = 2\lambda + a\lambda^2 + b\lambda\alpha + c\alpha^2, \quad b > 0, \quad a, c \leq 0$$



[Braun, 2011]

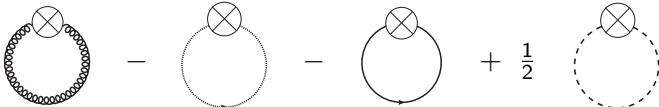


- agreement in perturbative regime required by Slavnov-Taylor identities
- non-degenerate in nonperturbative regime: reflects gluon mass gap
- $\alpha_{\bar{q}Aq} > \alpha_{cr}$: necessary for chiral symmetry breaking
- area above α_{cr} very sensitive to errors
 \Rightarrow use STI in perturbative regime

4-Fermi vertex via dynamical hadronization

[Gies, Wetterich, 2002]

- change of variables: particular 4-Fermi channels \rightarrow meson exchange
- efficient inclusion of pole structure \Rightarrow no spurious singularities
- calculation of low-energy model parameters from QCD

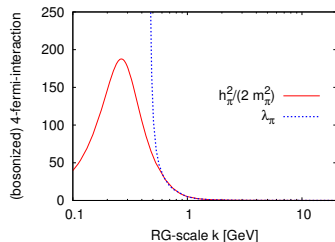
$$\partial_k \Gamma_k = \frac{1}{2} \left(\text{diagram 1} - \text{diagram 2} - \text{diagram 3} + \frac{1}{2} \text{diagram 4} \right)$$


4-Fermi vertex via dynamical hadronization

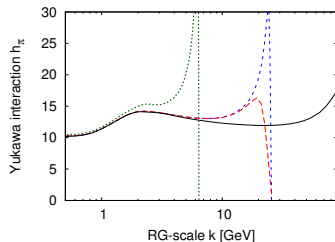
[Gies, Wetterich, 2002]

- change of variables: particular 4-Fermi channels \rightarrow meson exchange
- efficient inclusion of pole structure \Rightarrow no spurious singularities
- calculation of low-energy model parameters from QCD

$$\partial_k \Gamma_k = \frac{1}{2} \left(\text{Diagram 1} - \text{Diagram 2} - \text{Diagram 3} + \frac{1}{2} \text{Diagram 4} \right)$$



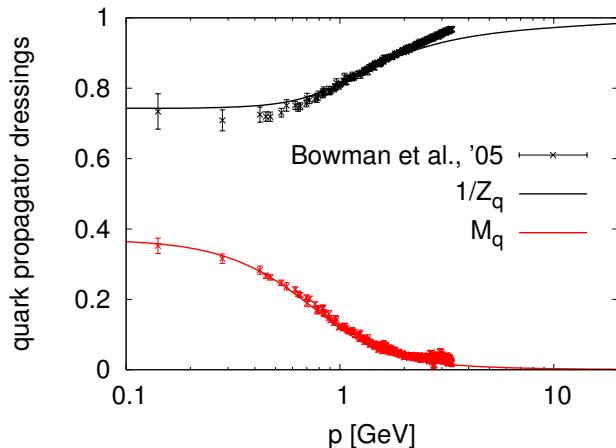
[MM, Strodthoff, Pawłowski, 2014]



[Braun, Fister, Haas, Pawłowski, Rennecke, 2014]

[MM, Strodthoff, Pawłowski, 2014]

- $\Gamma_{\bar{q}q}^{(2)}(p) \propto Z_q(p) (\not{p} + M(p))$

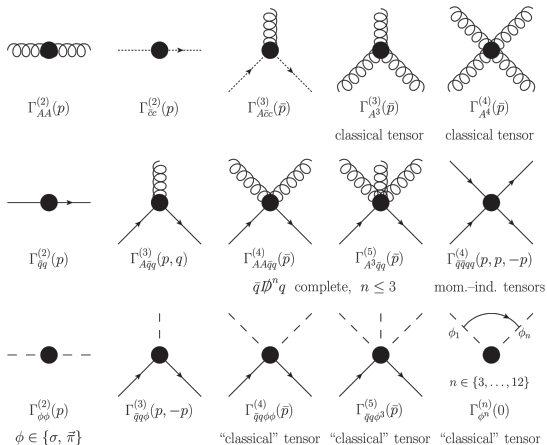


- FRG vs. lattice: bare mass, quenched, scale set via gluon propagator
- not sufficient: need convergence of unquenched QCD at $T, \mu \neq 0$

Outlook: unquenching, $N_f = 2$

[Cyrol, MM, Pawłowski, Strodthoff, in prep.]

extended truncation:

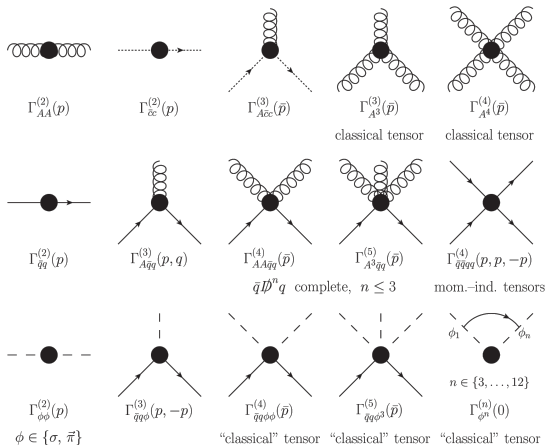


systematics of improving the truncation?

Outlook: unquenching, $N_f = 2$

[Cyrol, MM, Pawłowski, Strodthoff, in prep.]

extended truncation:



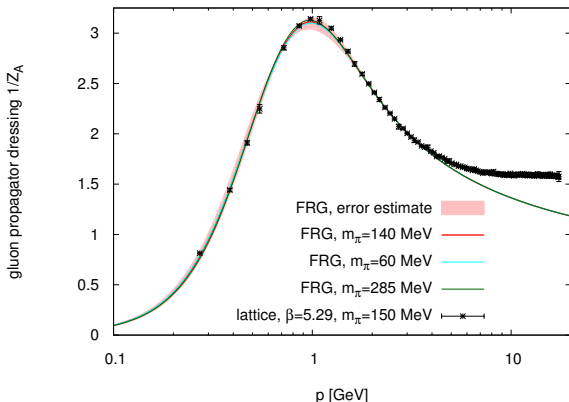
systematics of improving the truncation?

\Rightarrow BRST-invariant operators, e.g. $\bar{\psi}\not{D}^n\psi$

Outlook: gluon propagator

[Cyrol, MM, Pawłowski, Strodthoff, in prep.]

- $\Gamma_{AA}(p) = Z_A(p) p^2 (\delta^{\mu\nu} - p^\mu p^\nu / p^2)$



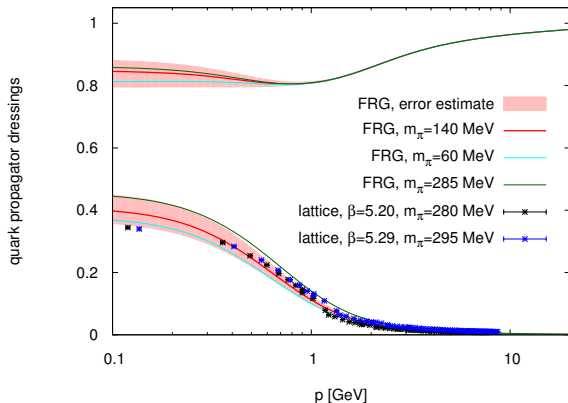
- infrared suppression \Leftrightarrow “confinement”
- insensitive to pion mass
- transition to perturbation theory
- scaling solution

lattice data: A. Sternbeck, K. Maltman, M. Müller-Preussker, L. von Smekal, PoS LATTICE2012 (2012) 243.

Outlook: quark propagator

[Cyrol, MM, Pawłowski, Strodthoff, in prep.]

- $\Gamma_{\bar{q}q}(p) = Z_q(p) (\not{p} + M(p))$



- wave function renormalisation Z_q
- mass function M_q
- sensitive to $\bar{q}qA$ -interaction, relative scales
- lattice Z_q ?

lattice data: O. Oliveira, A. Kzlersu, P. J. Silva, J.-I. Skullerud, A. Sternbeck, A. G. Williams, arXiv:1605.09632 [hep-lat].

Outlook: Quark-gluon interactions

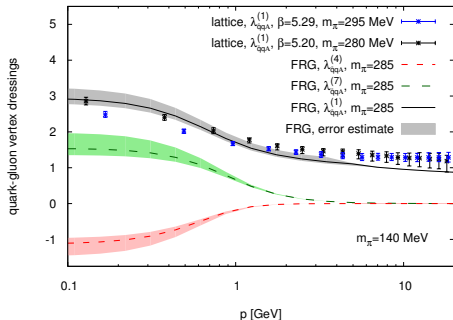
[Cyrol, MM, Pawłowski, Strodthoff, in prep.]

- quark-gluon interaction most crucial for chiral symmetry breaking
- transverse tensor basis (8 tensors), e.g. γ^μ , $i(\not{p} + \not{q}) \gamma^\mu$, $\frac{1}{2} [\not{p}, \not{q}] \gamma^\mu$
- $\lambda^{(i)}(p, q) \rightarrow \lambda^{(i)}(p^2, q^2, p \cdot q)$

Outlook: Quark-gluon interactions

[Cyrol, MM, Pawłowski, Strodthoff, in prep.]

- quark-gluon interaction most crucial for chiral symmetry breaking
- transverse tensor basis (8 tensors), e.g. γ^μ , $i(\not{p} + \not{q}) \gamma^\mu$, $\frac{1}{2} [\not{p}, \not{q}] \gamma^\mu$
- $\lambda^{(i)}(p, q) \rightarrow \lambda^{(i)}(p^2, q^2, p \cdot q)$

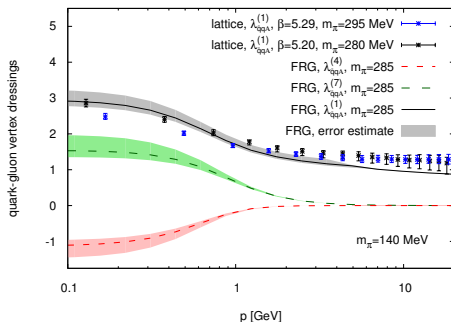


- 3 leading tensors:
 - ▶ classical tensor:
constrained by STI
at large momenta
 - ▶ chirally symmetric
 - ▶ break chiral symmetry
- systematic lattice error?

Outlook: Quark-gluon interactions

[Cyrol, MM, Pawłowski, Strodthoff, in prep.]

- quark-gluon interaction most crucial for chiral symmetry breaking
- transverse tensor basis (8 tensors), e.g. γ^μ , $i(\not{p} + \not{q}) \gamma^\mu$, $\frac{1}{2} [\not{p}, \not{q}] \gamma^\mu$
- $\lambda^{(i)}(p, q) \rightarrow \lambda^{(i)}(p^2, q^2, p \cdot q)$



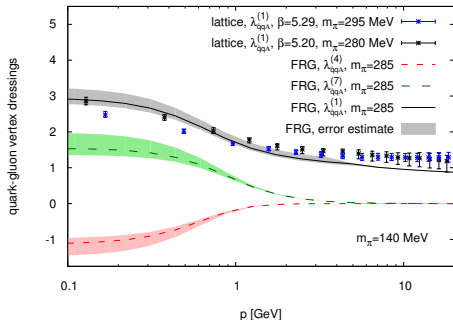
- 3 leading tensors:
 - ▶ classical tensor: constrained by STI at large momenta
 - ▶ **chirally symmetric**
 - ▶ **break chiral symmetry**
- systematic lattice error?

- **chirally symmetric tensors** from operator $\bar{q} \not{D}^3 q$ worsen result
- counteracted by tensor structures in $\Gamma_{AA\bar{q}q}^{(4)}$ and $\Gamma_{A^3\bar{q}q}^{(5)}$ from $\bar{q} \not{D}^3 q$

Outlook: Quark-gluon interactions

[Cyrol, MM, Pawłowski, Strodthoff, in prep.]

- quark-gluon interaction most crucial for chiral symmetry breaking
- transverse tensor basis (8 tensors), e.g. γ^μ , $i(\not{p} + \not{q}) \gamma^\mu$, $\frac{1}{2} [\not{p}, \not{q}] \gamma^\mu$
- $\lambda^{(i)}(p, q) \rightarrow \lambda^{(i)}(p^2, q^2, p \cdot q)$



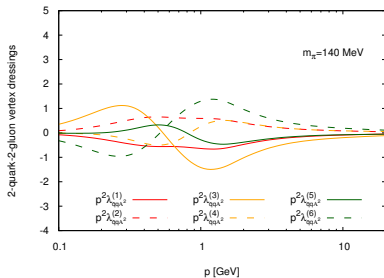
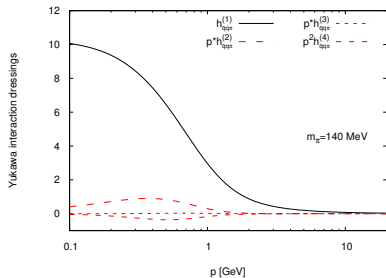
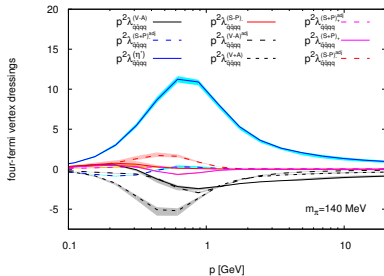
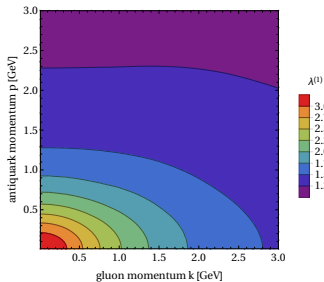
- 3 leading tensors:
 - ▶ classical tensor: constrained by STI at large momenta
 - ▶ chirally symmetric
 - ▶ break chiral symmetry
- systematic lattice error?

- chirally symmetric tensors from operator $\bar{q} \not{D}^3 q$ worsen result
 - counteracted by tensor structures in $\Gamma_{AA\bar{q}q}^{(4)}$ and $\Gamma_{A^3\bar{q}q}^{(5)}$ from $\bar{q} \not{D}^3 q$
- \Rightarrow expansion in BRST-invariant operators improves convergence?

lattice data: O. Oliveira, A. Kzlersu, P. J. Silva, J.-I. Skullerud, A. Sternbeck, A. G. Williams, arXiv:1605.09632 [hep-lat].

Outlook: more correlators

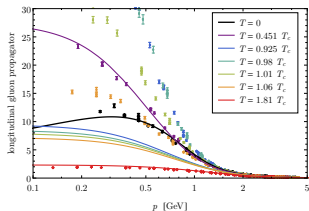
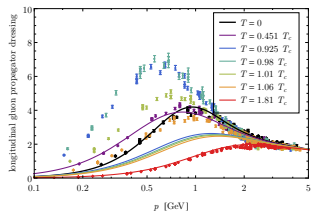
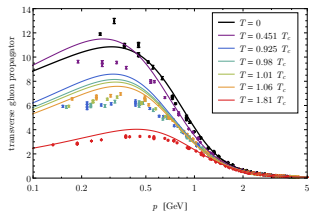
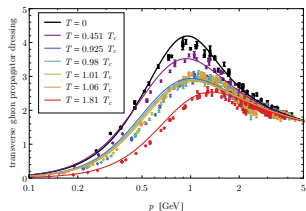
[Cyrol, MM, Pawłowski, Strodthoff, in prep.]



Outlook: YM correlators I at $T \neq 0$ [Cyrol, MM, Pawłowski, Strodthoff, in prep.]

full splitting of classical tensors into electric and magnetic parts

Zeroth mode correlation functions

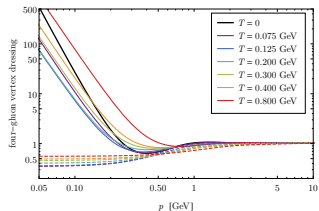
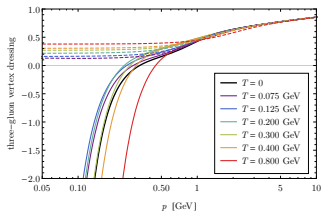
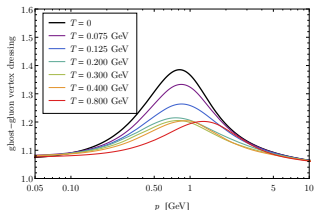
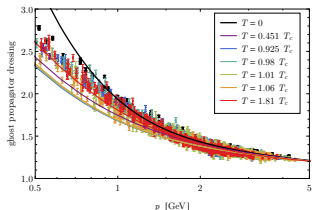


lattice data: A. Maas, J. M. Pawłowski, L. von Smekal, D. Spielmann, Phys.Rev. D85 (2012) 034037.

Outlook: YM correlators II at $T \neq 0$ [Cyrol, MM, Pawłowski, Strodthoff, in prep.]

full splitting of classical tensors into electric and magnetic parts

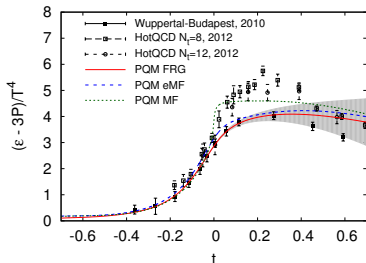
Zeroth mode correlation functions (solid lines: magnetic dressings; dashed lines: dressings with two electric legs)



lattice data: A. Maas, J. M. Pawłowski, L. von Smekal, D. Spielmann, Phys.Rev. D85 (2012) 034037.

“QCD-enhanced” models

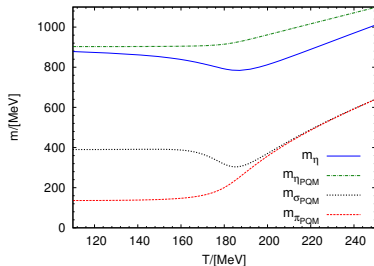
- equation of state
- $N_f = 2 + 1$ PQM model with FRG
- unquenched Polyakov-loop potential from [Braun, Haas, Marhauser, Pawłowski, '11]



[Herbst, MM, Pawłowski, Schaefer, Stiele, '13]

[Haas, Stiele, Braun, Pawłowski, Schaffner-Bielich, '13]

- η' -meson screening mass
- $N_f = 2$ PQM model, extended MF
- running 't Hooft determinant from [MM, Pawłowski, Strothoff, '14]



[Heller, MM, '15]

Status and Outlook

- vacuum QCD well under control
 - ▶ QCD from $\alpha_S(\Lambda = \mathcal{O}(10) \text{ GeV})$ and $m_q(\Lambda = \mathcal{O}(10) \text{ GeV})$
 - ▶ good agreement with lattice correlators
 - ▶ more checks on convergence of vertex expansion

Status and Outlook

- vacuum QCD well under control
 - ▶ QCD from $\alpha_S(\Lambda = \mathcal{O}(10) \text{ GeV})$ and $m_q(\Lambda = \mathcal{O}(10) \text{ GeV})$
 - ▶ good agreement with lattice correlators
 - ▶ more checks on convergence of vertex expansion
- first tests at finite temperature in YM-theory very encouraging

Status and Outlook

- vacuum QCD well under control
 - ▶ QCD from $\alpha_S(\Lambda = \mathcal{O}(10) \text{ GeV})$ and $m_q(\Lambda = \mathcal{O}(10) \text{ GeV})$
 - ▶ good agreement with lattice correlators
 - ▶ more checks on convergence of vertex expansion
- first tests at finite temperature in YM-theory very encouraging
- observables (EOS, fluctuations, order parameters, spectrum):
 - ▶ accessible through $\Gamma[\Phi]$ /correlators
 - ▶ calculation tools well-tested in simpler models
 - ▶ ready for QCD applications

Status and Outlook

- vacuum QCD well under control
 - ▶ QCD from $\alpha_S(\Lambda = \mathcal{O}(10) \text{ GeV})$ and $m_q(\Lambda = \mathcal{O}(10) \text{ GeV})$
 - ▶ good agreement with lattice correlators
 - ▶ more checks on convergence of vertex expansion
- first tests at finite temperature in YM-theory very encouraging
- observables (EOS, fluctuations, order parameters, spectrum):
 - ▶ accessible through $\Gamma[\Phi]$ /correlators
 - ▶ calculation tools well-tested in simpler models
 - ▶ ready for QCD applications
- further applications:
 - ▶ input for “QCD-enhanced” models
 - ▶ other strongly-interacting theories