

TESTING THE COSMOLOGICAL PRINCIPLE

Subir Sarkar

University of Oxford

Niels Bohr Institute, Copenhagen

None of us can understand why there is a Universe at all, why anything should exist; that is the ultimate question. But while we cannot answer this question, we can at least make progress with the next simpler one, of what the Universe as a whole is like.

Dennis Sciama (1978)

*... the Universe must appear to be the same to all observers wherever they are. This ‘**cosmological principle**’ ...*



Edward Arthur Milne

Rouse Ball Professor of Mathematics & Fellow of Wadham College, Oxford, 1929-50

STANDARD COSMOLOGICAL MODEL

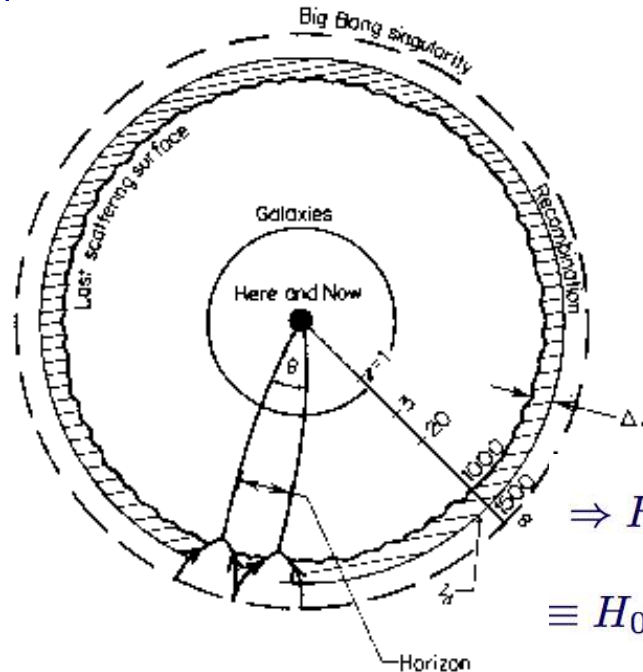
Universe is isotropic + Universe is homogeneous (when averaged on large scales)
 \Rightarrow Maximally-symmetric space-time and ideal fluid energy-momentum tensor

$$ds^2 \equiv g_{\mu\nu} dx^\mu dx^\nu$$

$$= a^2(\eta) [d\eta^2 - d\vec{x}^2]$$

$$a^2(\eta) d\eta^2 \equiv dt^2$$

Robertson-Walker



$$R_{\mu\nu} - \frac{1}{2} R g_{\mu\nu} + \lambda g_{\mu\nu} = 8\pi G_N T_{\mu\nu}$$

Einstein

$$\Rightarrow H^2 = \left(\frac{\dot{a}}{a}\right)^2 = \frac{8\pi G_N \rho_m}{3} - \frac{k}{a^2} + \frac{\Lambda}{3}$$

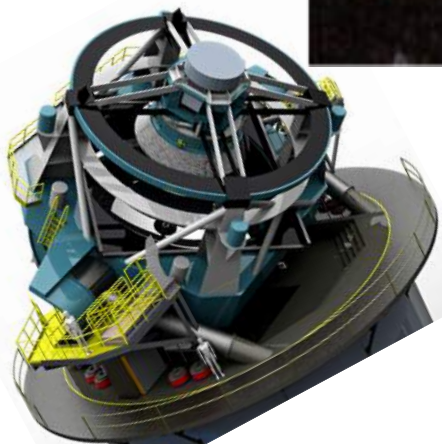
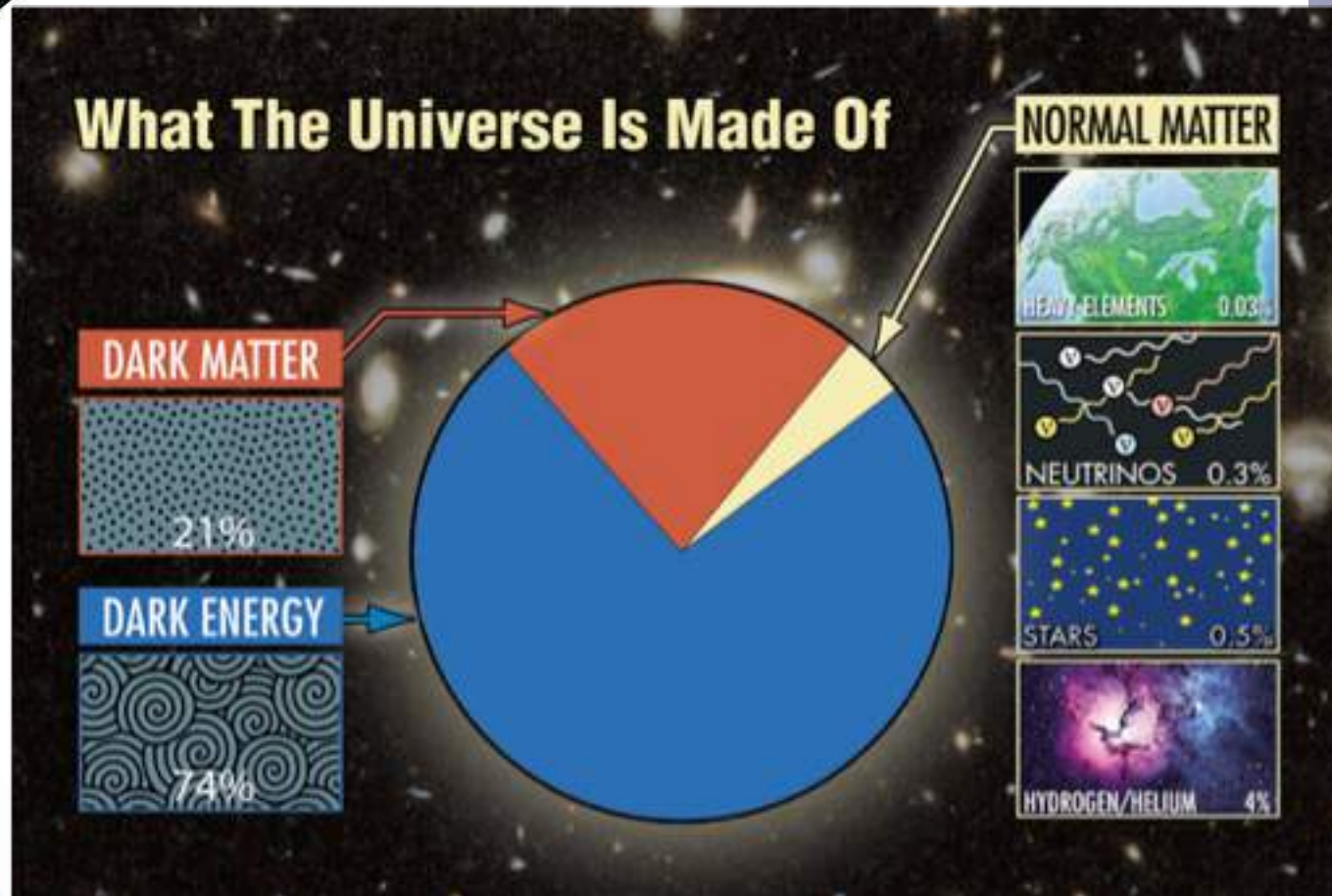
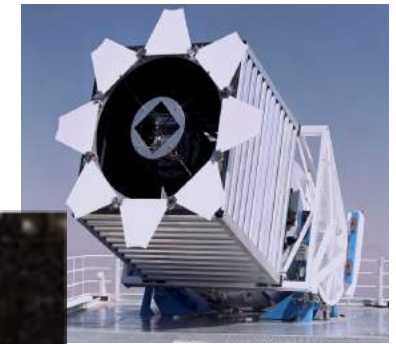
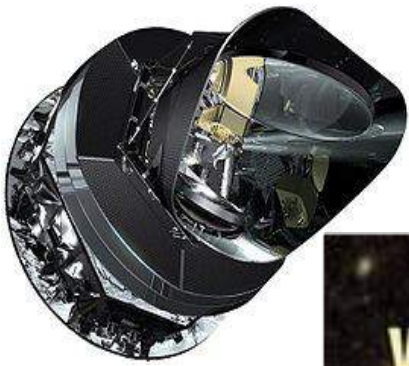
$$\equiv H_0^2 [\Omega_m(1+z)^3 + \Omega_k(1+z)^2 + \Omega_\Lambda]$$

So the **Friedmann-Lemaître equation** \Rightarrow Cosmic sum rule: $\Omega_{\text{matter}} + \Omega_{\text{curvature}} + \Omega_\Lambda = 1$

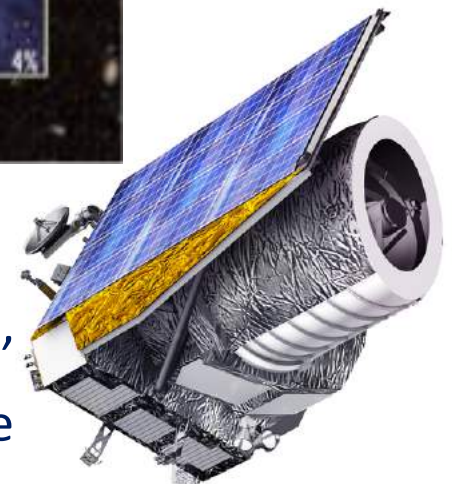
We observe \sim zero curvature (CMB fluctuations) + insufficient matter to make up critical density

\rightarrow **Universe is dominated by dark energy** with: $\Omega_\Lambda = 1 - \Omega_m - \Omega_k \sim 0.7 \Rightarrow \Lambda \sim 2H_0^2$

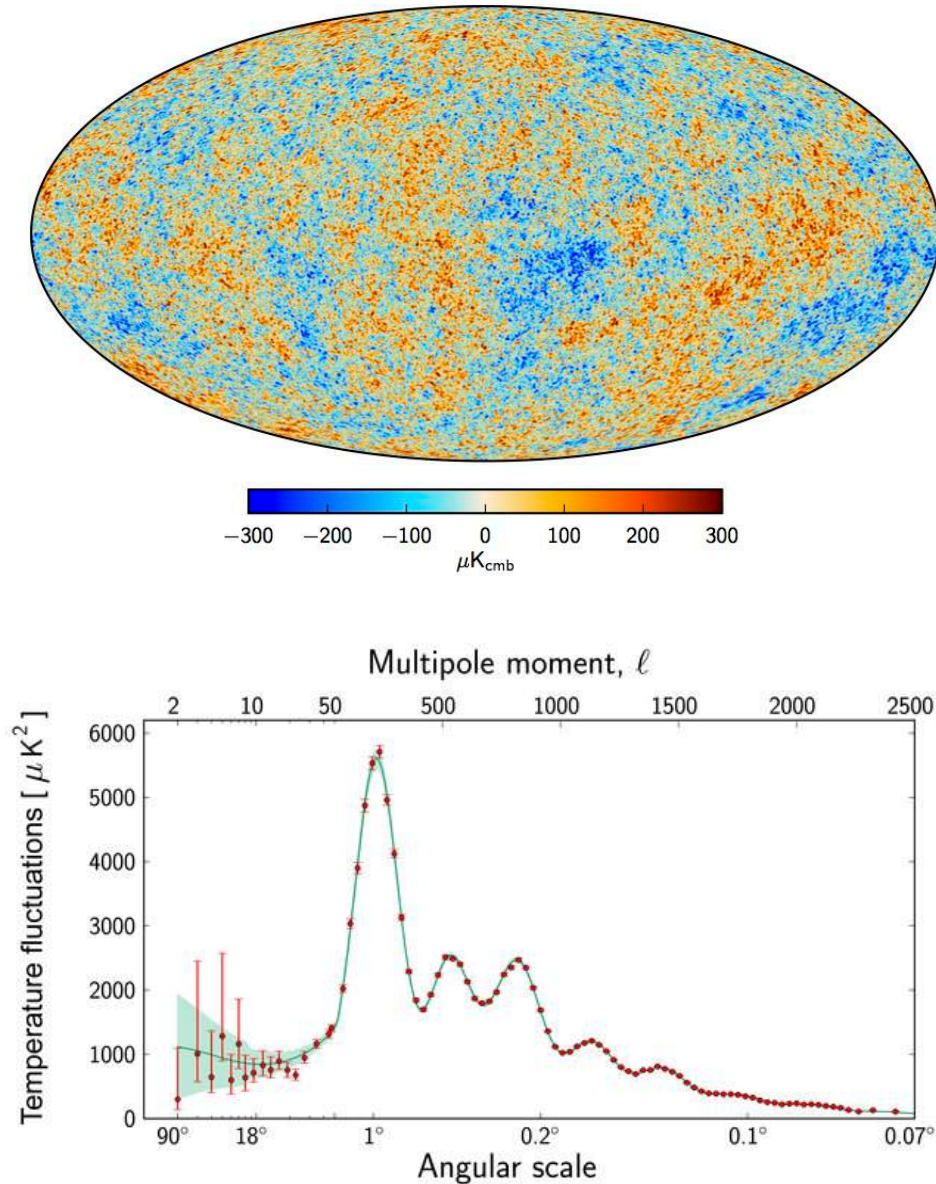
Since 1998 (Riess *et al.*¹, Perlmutter *et al.*²), surveys of cosmologically distant Type Ia supernovae (SNe Ia) have indicated an acceleration of the expansion of the Universe, distant SNe Ia being dimmer than expected in a decelerating Universe. With the assumption that the Universe can be described on average as isotropic and homogeneous, this acceleration implies either the existence of a fluid with negative pressure usually called “Dark Energy”, a constant in the equations of general relativity or modifications of gravity on cosmological scales.



There has been substantial investment in major satellites and telescopes to *measure the parameters* of the 'standard cosmological model' with increasing 'precision'... but surprisingly little interest in ***testing its foundational assumptions***

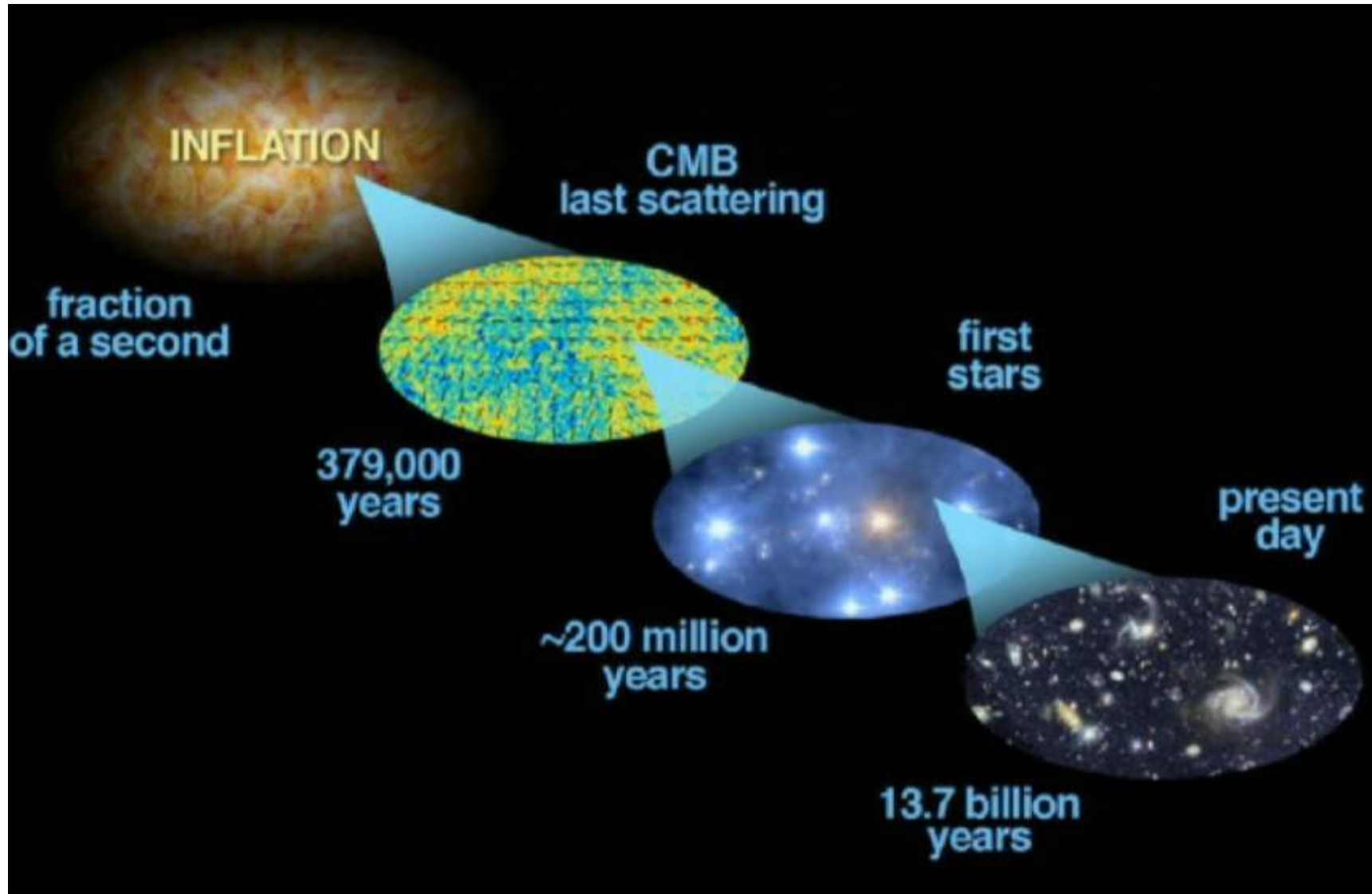


“Data from the Planck satellite show the universe to be highly isotropic” (Wikipedia)



We observe a statistically isotropic *Gaussian* random field of small temperature fluctuations (fully quantified by the 2-point correlations \rightarrow angular power spectrum)

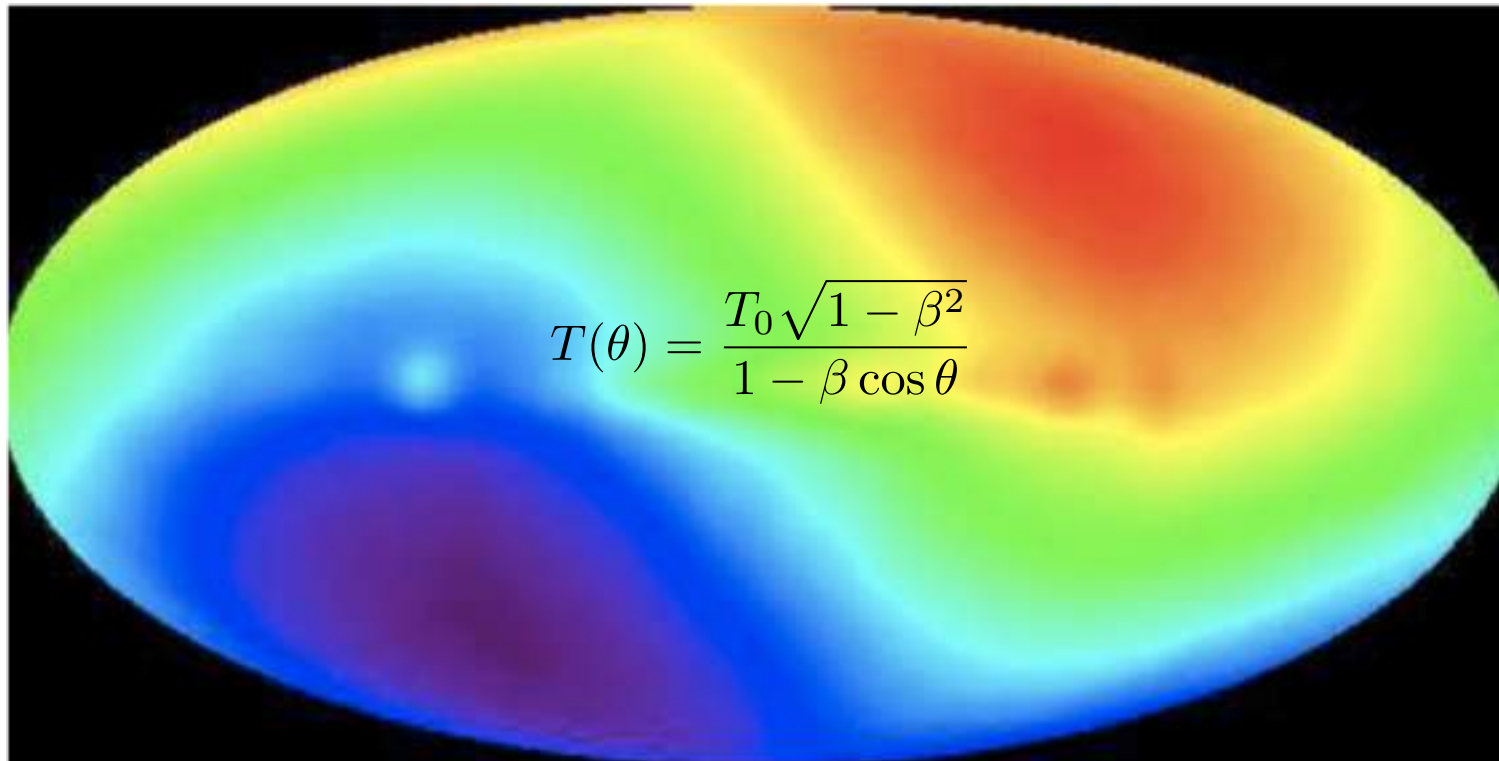
STANDARD MODEL OF STRUCTURE FORMATION



The tiny **CMB temperature fluctuations** are understood as due to **scalar density perturbations** with an \sim scale-invariant spectrum which were generated during an early phase of inflationary expansion ... these perturbations have subsequently grown into the **large-scale structure** of galaxies observed today through **gravitational instability** in a sea of **dark matter**

BUT THE CMB SKY IS IN FACT *VERY* ANISOTROPIC

There is a ~ 100 times *bigger* signal than the fluctuations in the form of a dipole with
 $\Delta T/T \sim 10^{-3}$



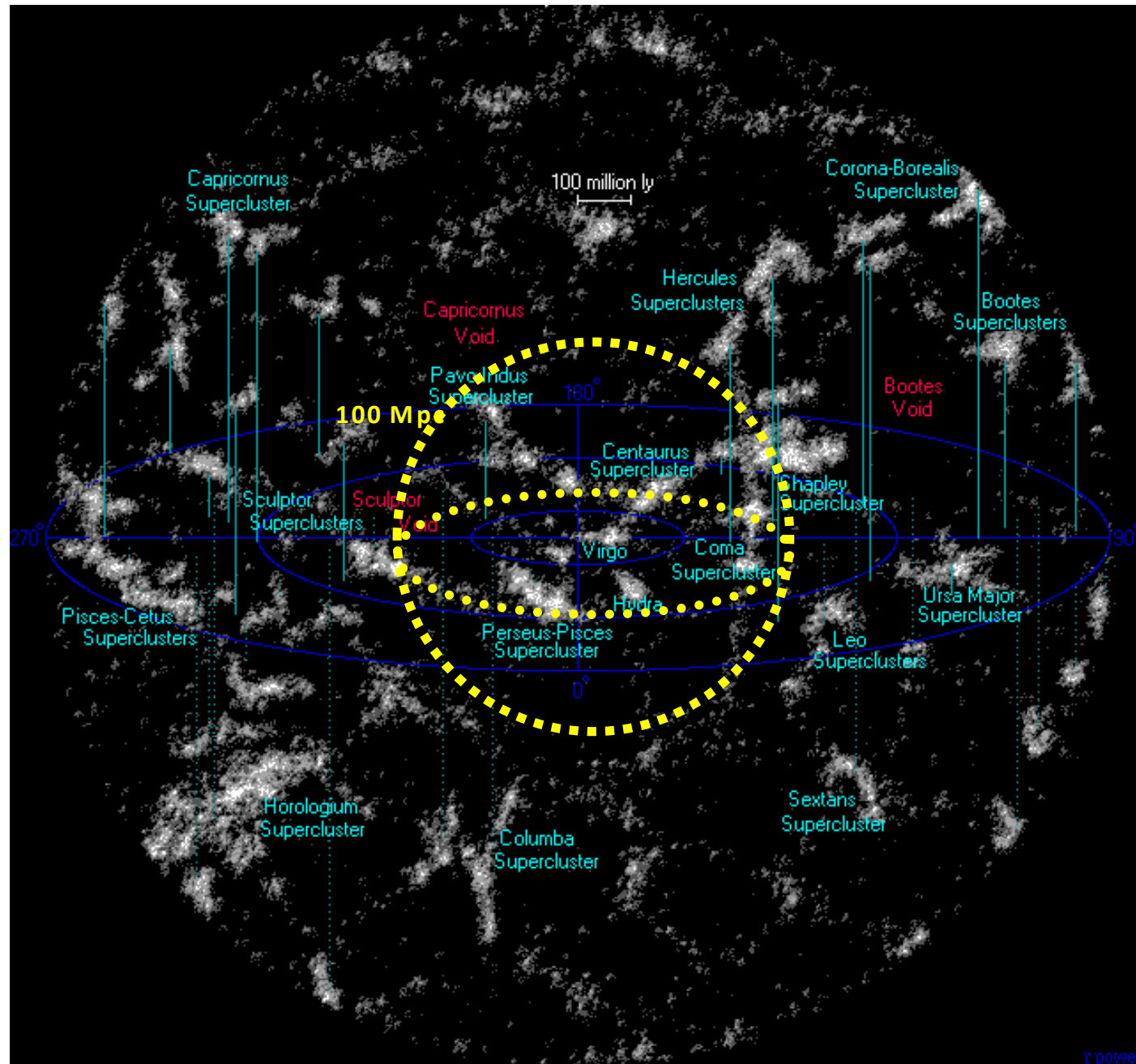
Stewart & Sciamma 1967, Peebles & Wilkinson 1968

This is *interpreted* as due to our motion at 368 km/s wrt the frame in which the CMB is truly isotropic \Rightarrow motion of the Local Group at 627 km/s towards $l=263.85^\circ$, $b=48.25^\circ$

This motion is presumed to be due to ***local inhomogeneity*** in the matter distribution

Its scale – *beyond* which we converge to the CMB frame – is supposedly of $\mathcal{O}(100)$ Mpc (Counts of galaxies in SDSS & WiggleZ surveys are said to scale as $\sim r^3$ on larger scales)

THIS IS WHAT OUR UNIVERSE *ACTUALLY* LOOKS LIKE (OUT TO ~300 Mpc)



Our motion is towards the Shapley supercluster, supposedly due to a 'Great Attractor' beyond ...

THEORY OF PECULIAR VELOCITY FIELDS

In linear perturbation theory, the growth of the density contrast $\delta(x) = [\rho(x) - \bar{\rho}]/\bar{\rho}$ as a function of comoving coordinates and time is governed by:

$$\frac{\partial^2 \delta}{\partial t^2} + 2H(t) \frac{\partial \delta}{\partial t} = 4\pi G_N \bar{\rho} \delta$$

We are interested in the ‘growing mode’ solution – the density contrast grows self-similarly and so does the perturbation potential and its gradient ... so the direction of the acceleration (and its integral – the peculiar velocity) remains *unchanged*.

The peculiar velocity field is related to the density contrast as:

$$v(\mathbf{x}) = \frac{2}{3H_0} \int d^3y \frac{\mathbf{x} - \mathbf{y}}{|\mathbf{x} - \mathbf{y}|^3} \delta(\mathbf{y}),$$

So the peculiar Hubble flow, $\delta H(\mathbf{x}) = H_L(\mathbf{x}) - H_0$ (\Rightarrow trace of the shear tensor), is:

$$\delta H(\mathbf{x}) = \int d^3y \mathbf{v}(\mathbf{y}) \cdot \frac{\mathbf{x} - \mathbf{y}}{|\mathbf{x} - \mathbf{y}|^2} W(\mathbf{x} - \mathbf{y}),$$

where $H_L(\mathbf{x})$ is the **local** value of the Hubble parameter and $W(\mathbf{x} - \mathbf{y})$ is the ‘window function’ (e.g. $\Theta(R - |\mathbf{x} - \mathbf{y}|) (4\pi R^3/3)^{-1}$ for a volume-limited survey, out to distance R)

THEORY OF PECULIAR VELOCITY FIELDS (CONT.)

Rewrite in terms of the Fourier transform $\delta(\mathbf{k}) \equiv (2\pi)^{3/2} \int d^3x \delta(\mathbf{x}) e^{i\mathbf{k}\cdot\mathbf{x}}$:

$$\frac{\delta H}{H_0} = \int \frac{d^3k}{(2\pi)^{3/2}} \delta(k) \mathcal{W}_H(kR) e^{ik \cdot x}, \quad \mathcal{W}_H(x) = \frac{3}{x^3} \left(\sin x - \int_0^x dy \frac{\sin y}{y} \right)$$

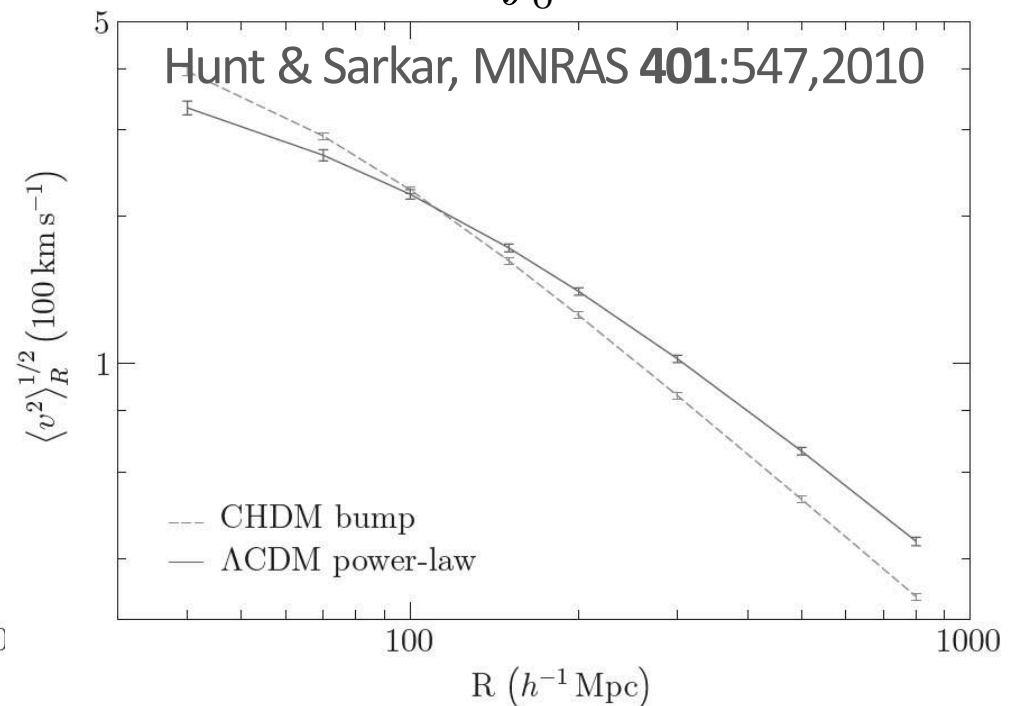
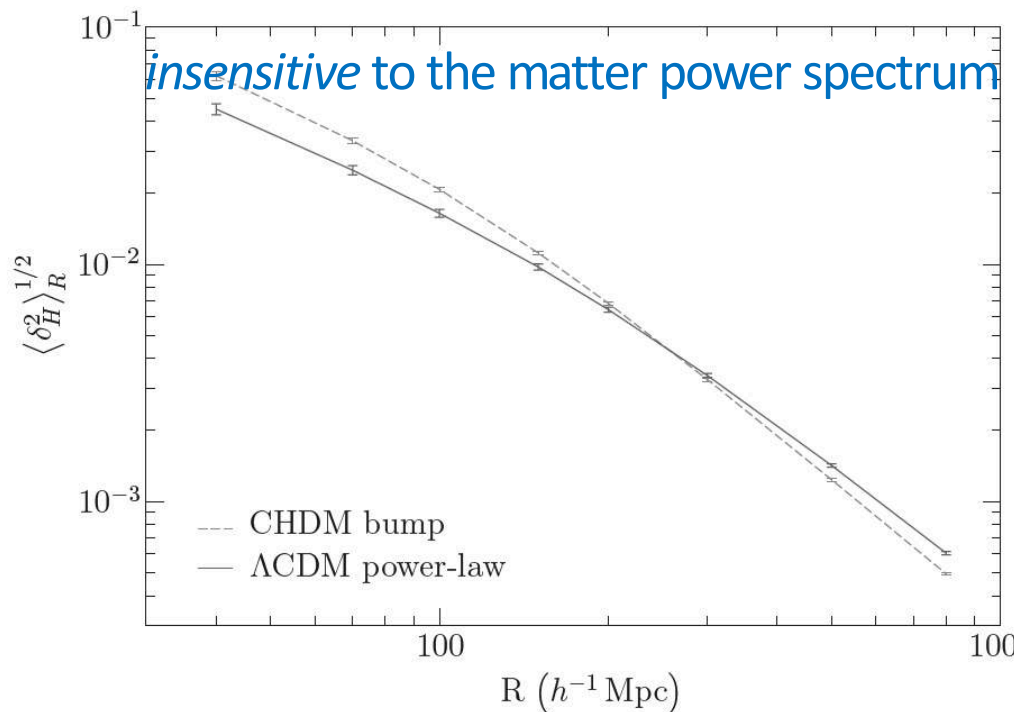
Window function

Then the RMS fluctuation in the local Hubble constant $\delta_H \equiv \langle (\delta H/H_0)^2 \rangle^{1/2}$ is:

$$\delta_H^2 = \frac{f^2}{2\pi^2} \int_0^\infty k^2 dk P(k) \mathcal{W}^2(kR), \quad P(k) \equiv |\delta(k)|^2, \quad f \simeq \Omega_m^{4/7} + \frac{\Omega_\Lambda}{70} \left(1 + \frac{\Omega_m}{2} \right)$$

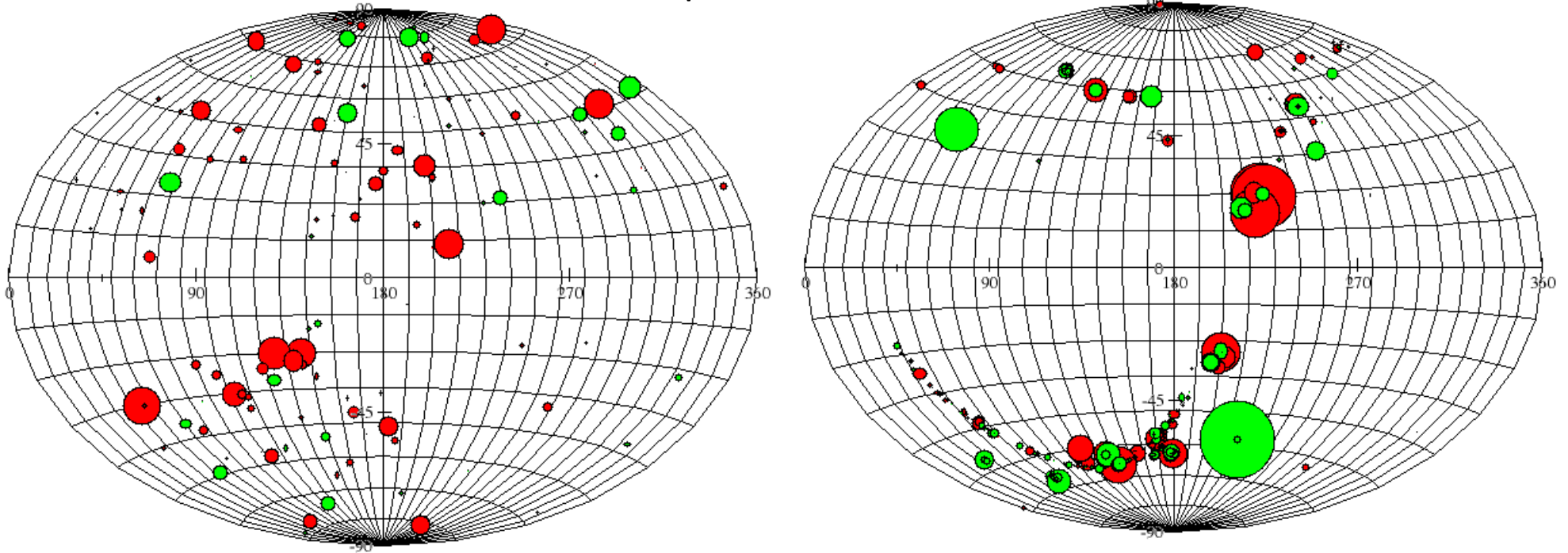
Growth rate

Similarly the variance of the peculiar velocity is: $\langle v^2 \rangle_R = \frac{f^2 H_0^2}{2\pi^2} \int_0^\infty dk P(k) \mathcal{W}^2(kR)$



UNION 2 COMPILATION OF 557 SNE IA

Aitoff-Hammer plot, Galactic coordinates



Left panel: The red spots represent the data points for $z < 0.06$ with distance moduli μ_{data} bigger than the values μ_{CDM} predicted by ΛCDM , and the green spots are those with μ_{data} less than μ_{CDM} ; the spot size is a relative measure of the discrepancy. A dipole anisotropy is visible around the direction $b = -30^\circ$, $l = 96^\circ$ (red points) and its opposite direction $b = 30^\circ$, $l = 276^\circ$ (small green points), which is the direction of the CMB dipole. **Right panel:** Same plot for $z > 0.06$

Colin, Mohayaee, Sarkar & Shafieloo, MNRAS **414**:264,2011

Use this to do *tomography* of the local Hubble flow by asking if the supernovae are at the expected distances: any residuals \Rightarrow 'peculiar velocity' flow in local universe

METHOD OF RESIDUALS AND SMOOTHING

$$q_i(z_i, \theta_i, \phi_i) = \frac{\mu_i(z_i, \theta_i, \phi_i) - \tilde{\mu}_i(z_i, \theta_i, \phi_i)}{\sigma_i(z_i, \theta_i, \phi_i)} \quad \text{Calculation of Residuals}$$

$$Q(\theta, \phi) = \sum_{i=1}^N q_i(z_i, \theta_i, \phi_i) W(\theta, \phi, \theta_i, \phi_i) \quad \text{2D smoothing on unit sphere}$$

$$W(\theta, \phi, \theta_i, \phi_i) = \frac{1}{\sqrt{2\pi}\delta} \exp \left[-\frac{L(\theta, \phi, \theta_i, \phi_i)^2}{2\delta^2} \right] \quad \text{Window function}$$

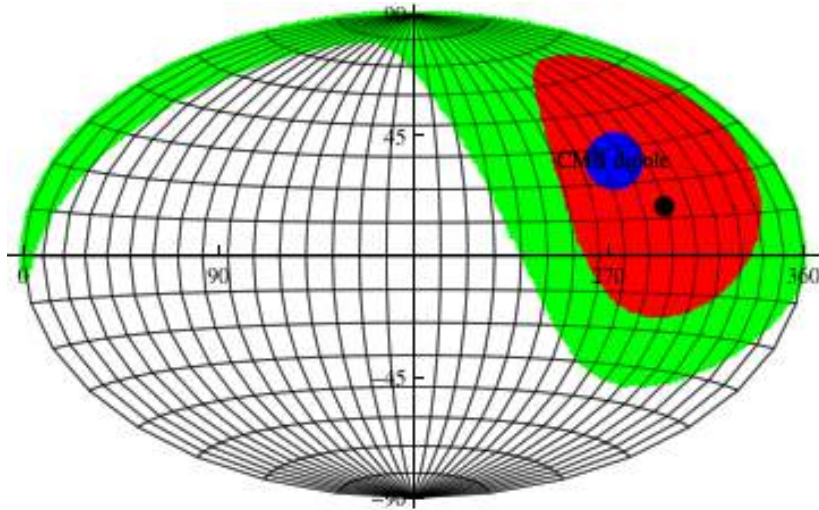
$$L(\theta, \phi, \theta_i, \phi_i) = 2 \arcsin \frac{R}{2}, \quad R = \left([\sin(\theta_i) \cos(\phi_i) - \sin(\theta) \cos(\phi)]^2 + [\sin(\theta_i) \sin(\phi_i) - \sin(\theta) \sin(\phi)]^2 + [\cos(\theta_i) - \cos(\theta)]^2 \right)^{1/2}$$

$$\Delta Q_{\text{data}} = Q(\theta_{\text{max}}, \phi_{\text{max}}) - Q(\theta_{\text{min}}, \phi_{\text{min}}) \quad \text{Statistical measure}$$

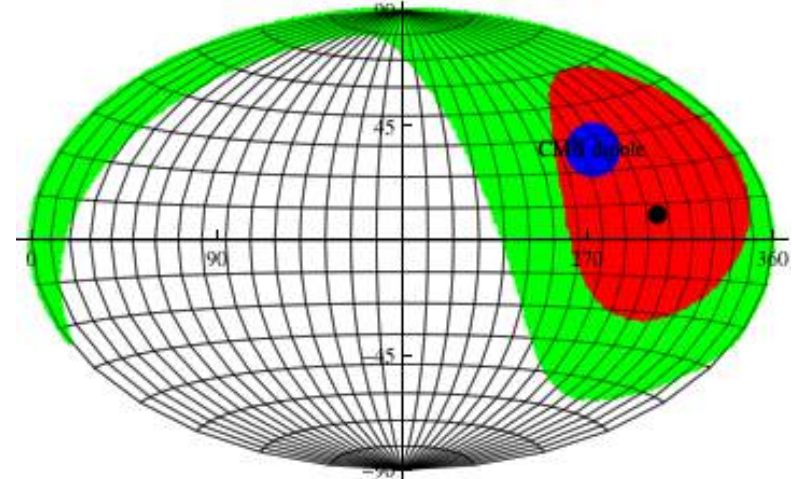
Calculate for the data (as well as for Monte Carlo simulations of isotropic distribution, in order to obtain p-value), using a ratio method to *minimise* systematic uncertainties

THERE IS A DIPOLE IN THE SNE IA VELOCITY FIELD TOWARDS THE CMB DIPOLE

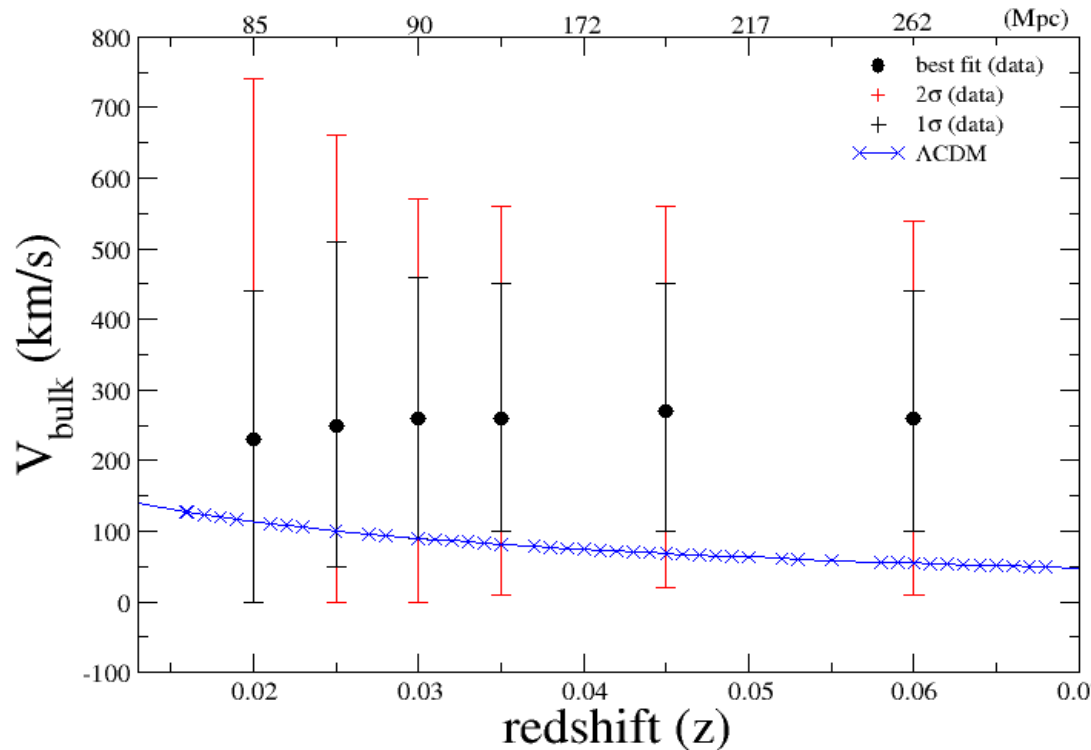
$0.015 < z < 0.045, v = 270 \text{ km/s}, l = 291, b = 15$



$0.015 < z < 0.06, v = 260 \text{ km/s}, l = 298, b = 8$



Colin et al, MNRAS 414:264,2011



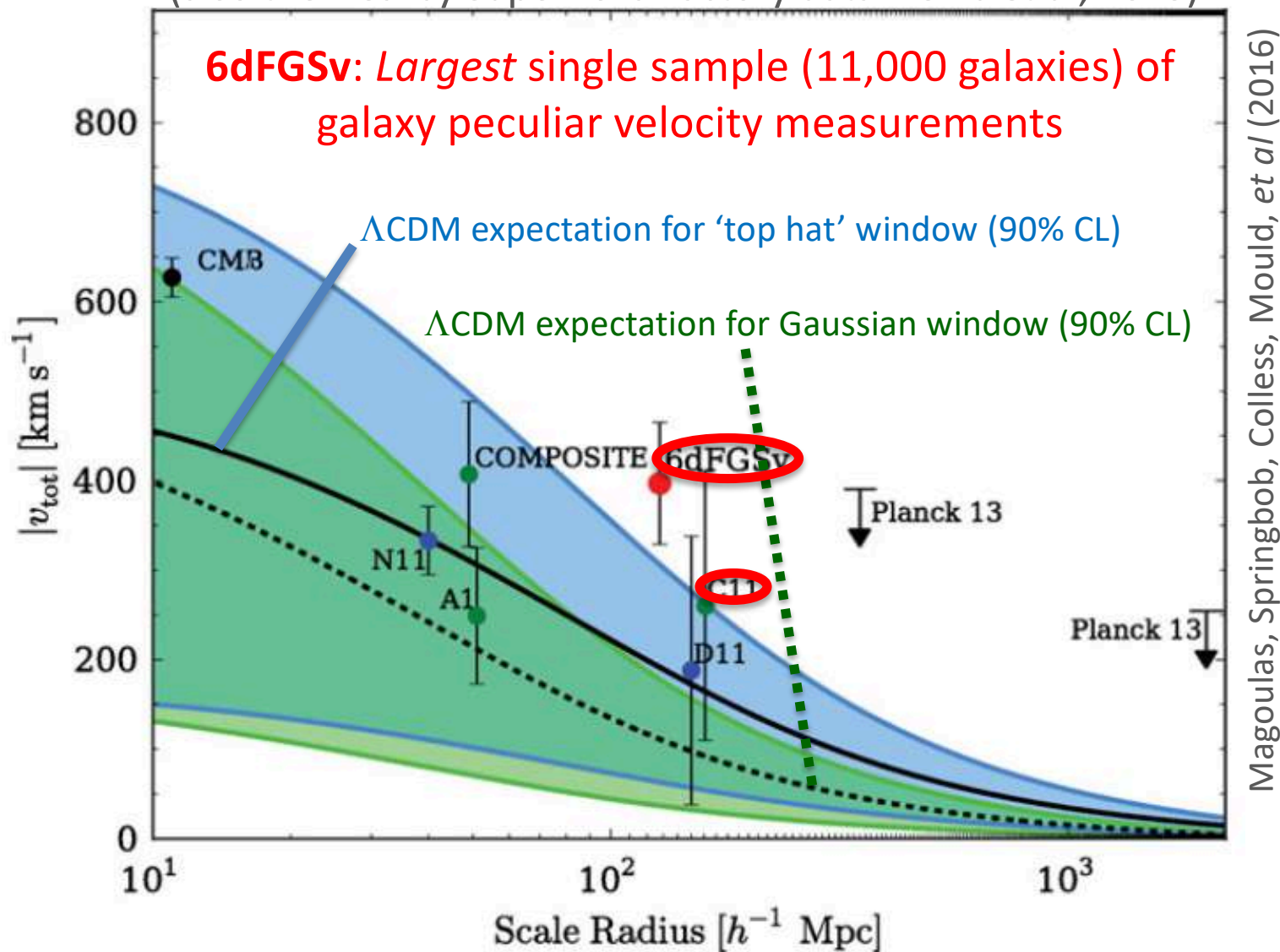
This is $\sim 1.5\sigma$ higher than expected for the standard Λ CDM model and extends *beyond* Shapley supercluster (260 Mpc)

Consistent with Watkins *et al* (2009) who found a bulk flow of $416 \pm 78 \text{ km/s}$ towards $b = 60 \pm 6^\circ, l = 282 \pm 11^\circ$ extending up to $\sim 100 h^{-1} \text{ Mpc}$

There is *no* convergence to CMB frame, well beyond 'scale of homogeneity'!

OUR RESULT HAS BEEN *CONFIRMED* BY THE 6-DEGREE FIELD GALAXY SURVEY

(also the Nearby Supernova Factory data: Feind *et al*, 2013)



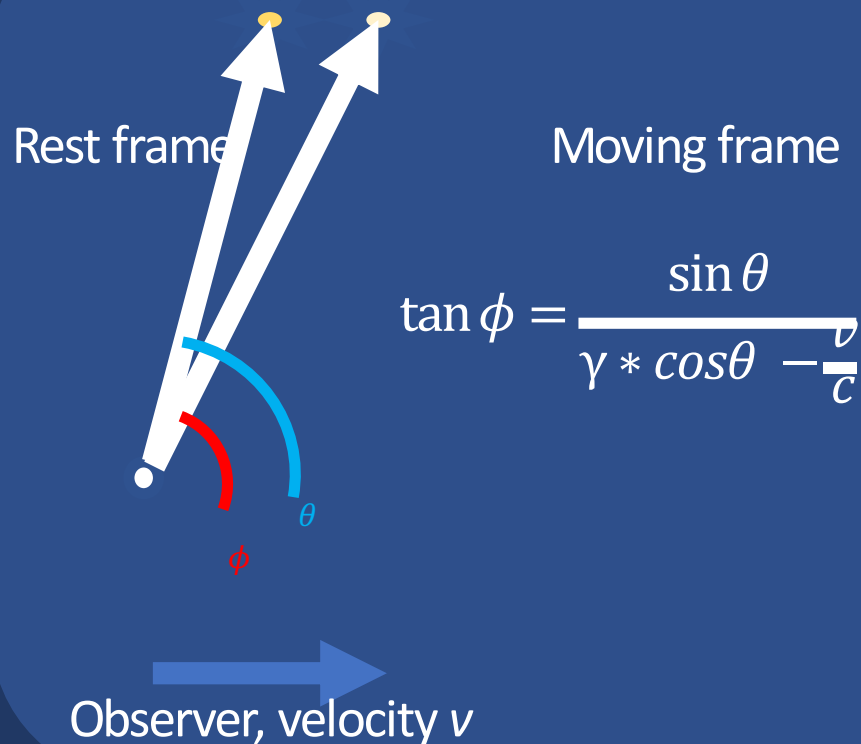
We find that in the 'Dark Sky' ΛCDM simulations, *less than 1%* of Milky Way-like observers experience a bulk flow as large as is observed, extending out as far as is seen

Rameez, Mohayaee, Sarkar & Colin, MNRAS **477**:1722,2018

A MOVING OBSERVER \rightarrow *KINEMATIC DIPOLE*

$$\sigma(\theta)_{obs} = \sigma_{rest} \left[1 + \left[2 + x(1 + \alpha) \right] \frac{v}{c} \cos(\theta) \right]$$

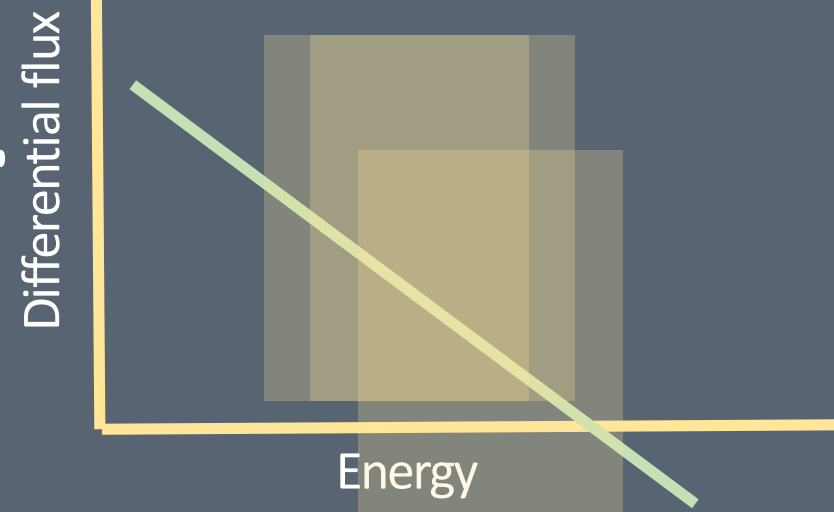
Aberration



+

Doppler boosting

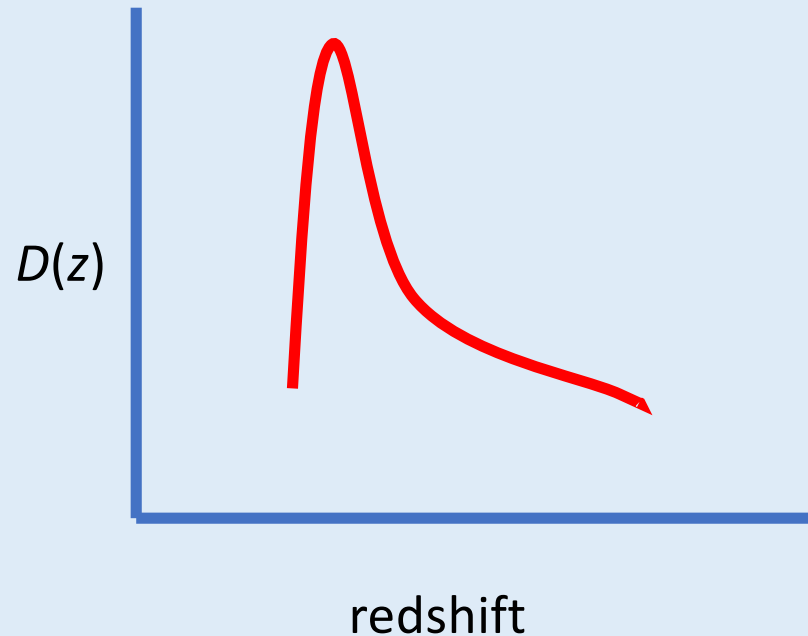
$$\phi \propto E^{-x} \text{ -ve power law}$$



Flux limited catalog \rightarrow more sources in direction of motion

DIPOLES IN A CATALOGUE OF GALAXIES

All-sky catalogue with N sources
with redshift distribution $D(z)$ from
a directionally unbiased survey



$$\vec{\delta} = \vec{\mathcal{K}}(\vec{v}_{obs}, x, \alpha) + \vec{\mathcal{R}}(N) + \vec{\mathcal{S}}(D(z))$$

$\vec{\mathcal{K}} \rightarrow$ The kinematic dipole: *independent*
of source distance, but depends on
source spectrum, source flux
function, observer velocity

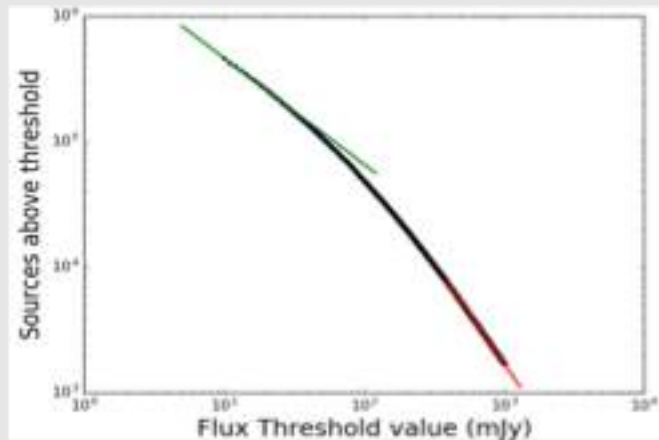
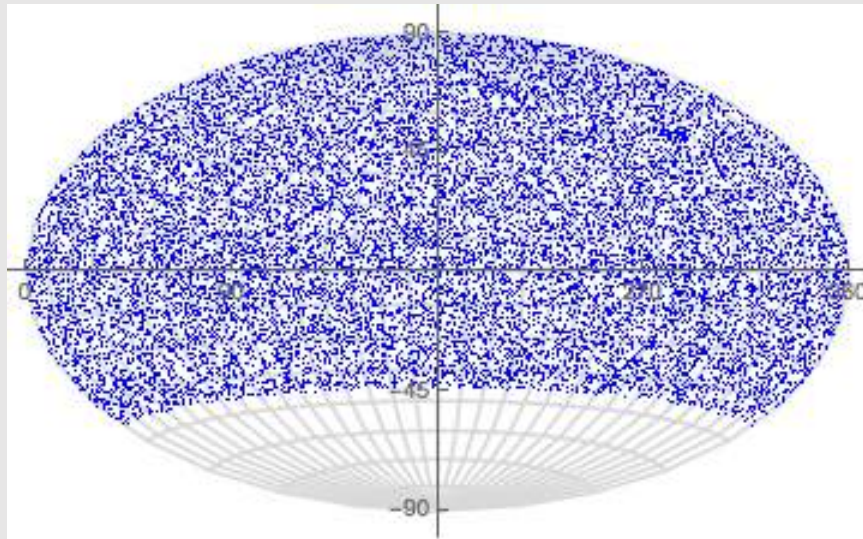
$\vec{\mathcal{R}} \rightarrow$ The random dipole: $\propto 1/\sqrt{N}$
isotropically distributed

$\vec{\mathcal{S}} \rightarrow$ The dipole component of an actual
anisotropy in the distribution of
sources in the cosmic rest frame
(significant for shallow surveys)

Radio sources: NVSS + SUMSS, 600,000 galaxies $z \sim 1$, $\vec{\mathcal{S}}(D(z)) \rightarrow 0$
Colin, Mohayaee, Rameez & Sarkar, MNRAS **471**:1045,2017

Wide Field Infrared Survey Explorer, 2,400,000 galaxies, $z \sim 0.14$, $\vec{\mathcal{S}}(D(z))$ significant
Rameez, Mohayaee, Sarkar & Colin MNRAS **477**:1722,2018

THE NRAO VLA SKY SURVEY (NVSS)

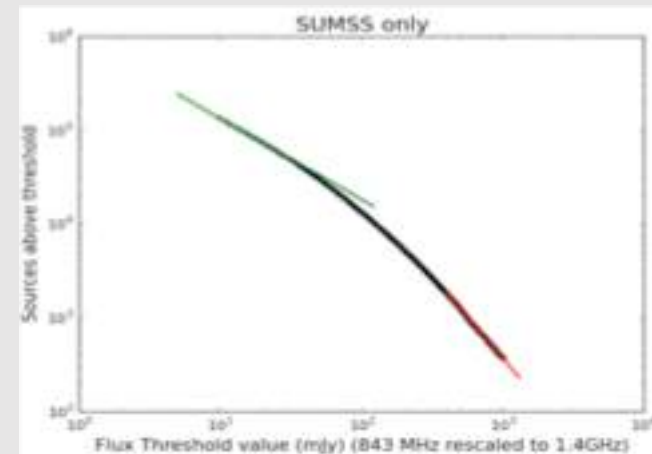
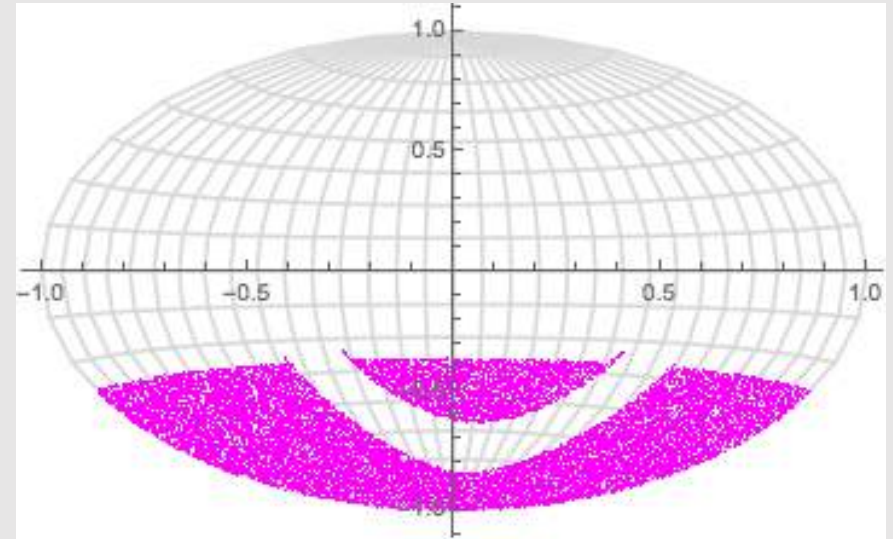


1.4 GHz survey (down to Dec = -40.4°)
National Radio Astronomy Observatory

1,773,488 sources >2.5 mJy
(complete above 10 mJy)

Most are believed to be at $z \gtrsim 1$

SYDNEY UNIVERSITY MOLONGLO SKY SURVEY (SUMSS)



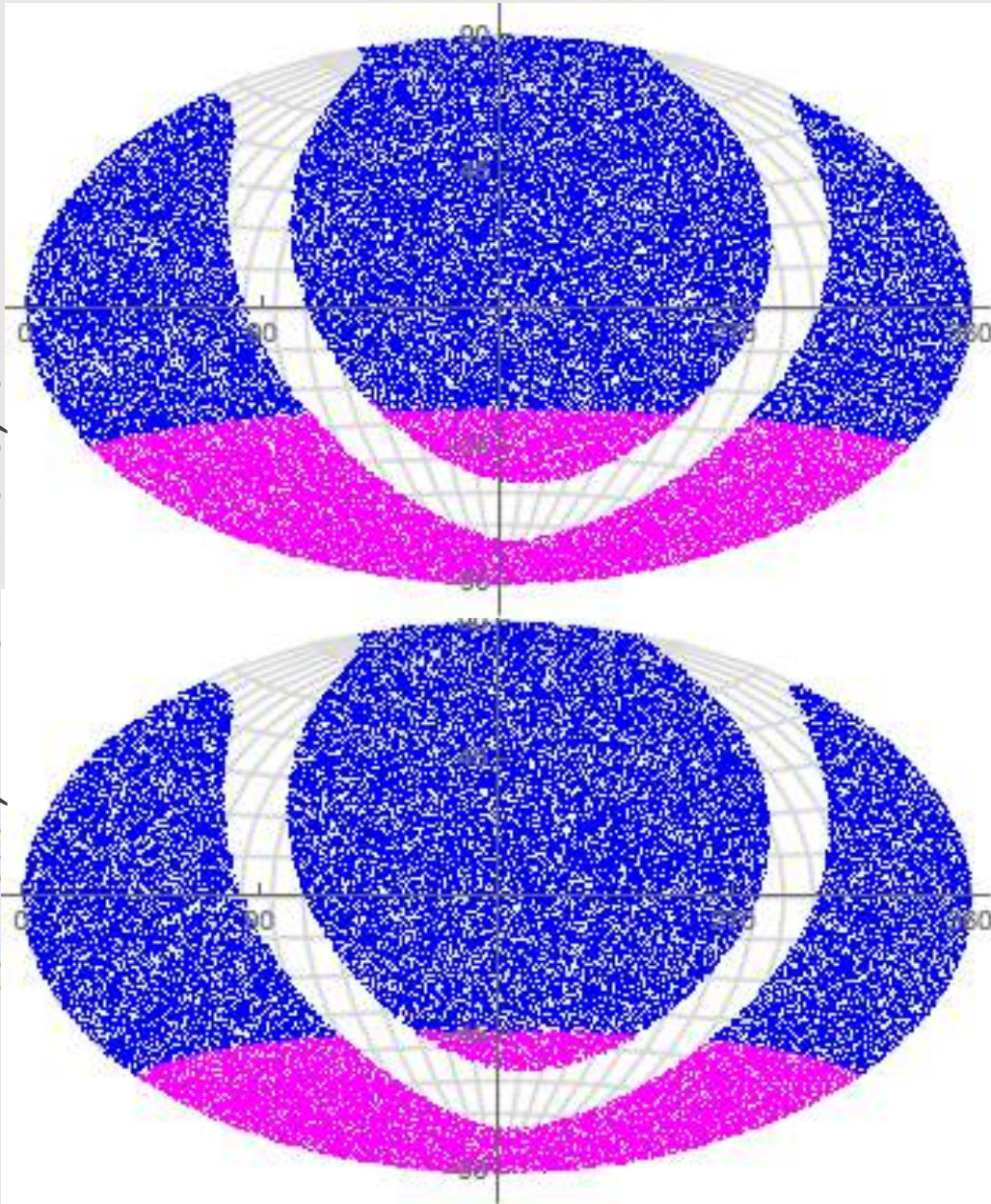
843 MHz survey (Dec $< -30.0^\circ$)
Molonglo Observatory Synthesis telescope

211,050 sources (with similar sensitivity and
resolution to NVSS catalogue)

... Similar expected redshift distribution

THE NVSUMSS-COMBINED ALL SKY CATALOG

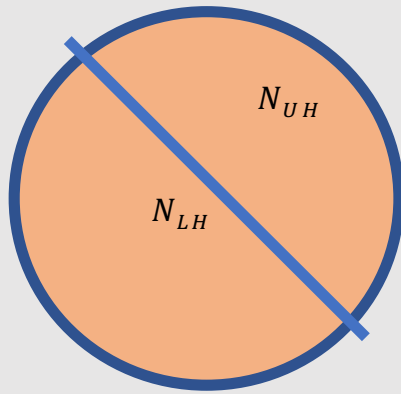
Colin et al, MNRAS 471:1045, 2017



- Rescale SUMSS fluxes by $(843/1400)^{-0.75} \sim 1.46$ to match with NVSS (within $\sim 1\%$)
- Remove Galactic Plane at $\pm 10^\circ$ (also super-galactic plane)
- Remove NVSS sources below, and SUMSS sources above, dec -30 (or -40)
- Apply common threshold flux cut to both samples
- Remove *any* nearby sources (common with 2MRS & LRS)

ESTIMATORS FOR THE DIPOLE

$$\vec{D}_H = \hat{z} * \frac{N_{UH} - N_{LH}}{N_{UH} + N_{LH}}$$



- Vary the direction of the hemispheres until maximum asymmetry is observed

- Easy visualisation

- But high bias (find by Monte Carlo)

$$\vec{D}_C = \frac{1}{N} \sum_{i=1}^N \hat{r}_i$$

- Add up the unit vectors corresponding to directions in the sky for every source

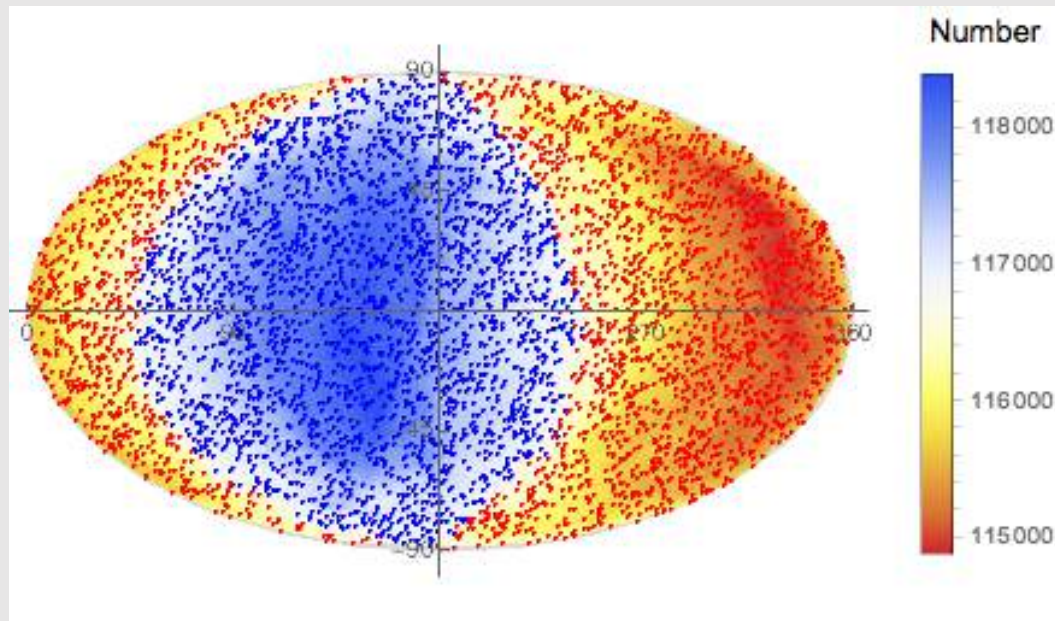
- Relatively low bias and statistical error $1/\sqrt{N}$

$$\vec{D}_C = \frac{\hat{z}}{N} \int_{\phi=0}^{\phi=2\pi} \int_{\theta=0}^{\theta=\pi} \sigma(\theta) \cos\theta \sin\theta d\theta d\phi$$

$$\vec{D}_H = \frac{\hat{z}}{N} \int_{\phi=0}^{\phi=2\pi} \int_{\theta=0}^{\theta=\pi} \sigma(\theta) \frac{|\cos\theta|}{\cos\theta} \sin\theta d\theta d\phi$$

OUR PECULIAR VELOCITY WRT RADIO SOURCES

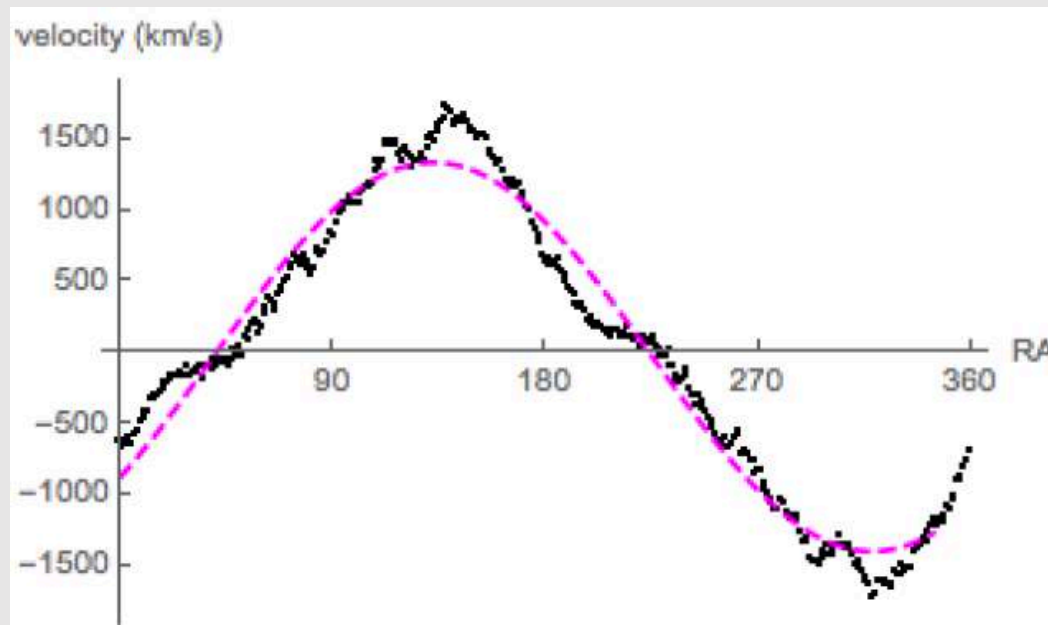
Colin et al, MNRAS 471:1045, 2017



Velocity $\sim 1355 \pm 174$ km/s
(with the 3D linear estimator)

Direction within 10° of CMB
dipole (but **4 times faster**)!

Statistical significance: 99.75%
 $\Rightarrow 2.81\sigma$ (by Monte Carlo)

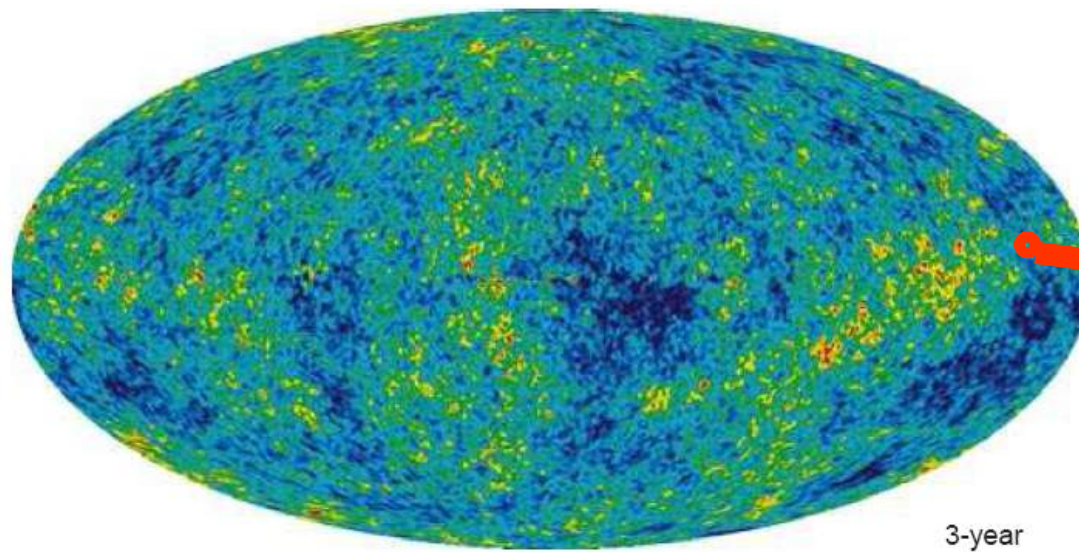


Confirms claim by Singal (2011)
which was criticized subsequently
(Gibelyou & Huterer 2012, Rubart &
Schwarz 2013, Nusser & Tiwari 2015)

We have addressed *all* the concerns
but this strange anomaly remains!

Look forward to data from SKA

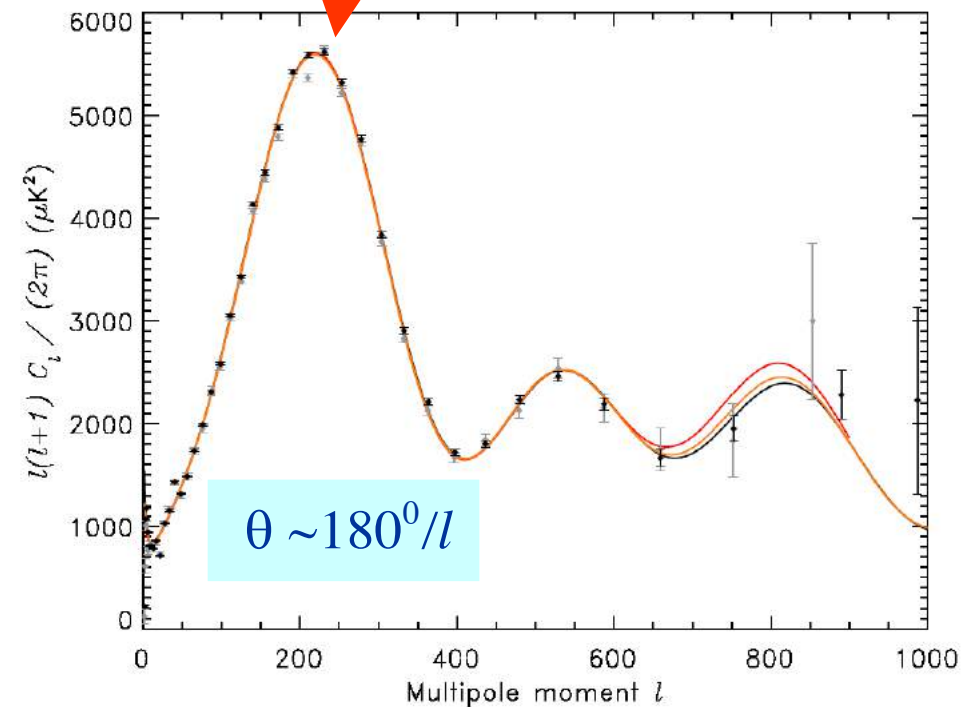
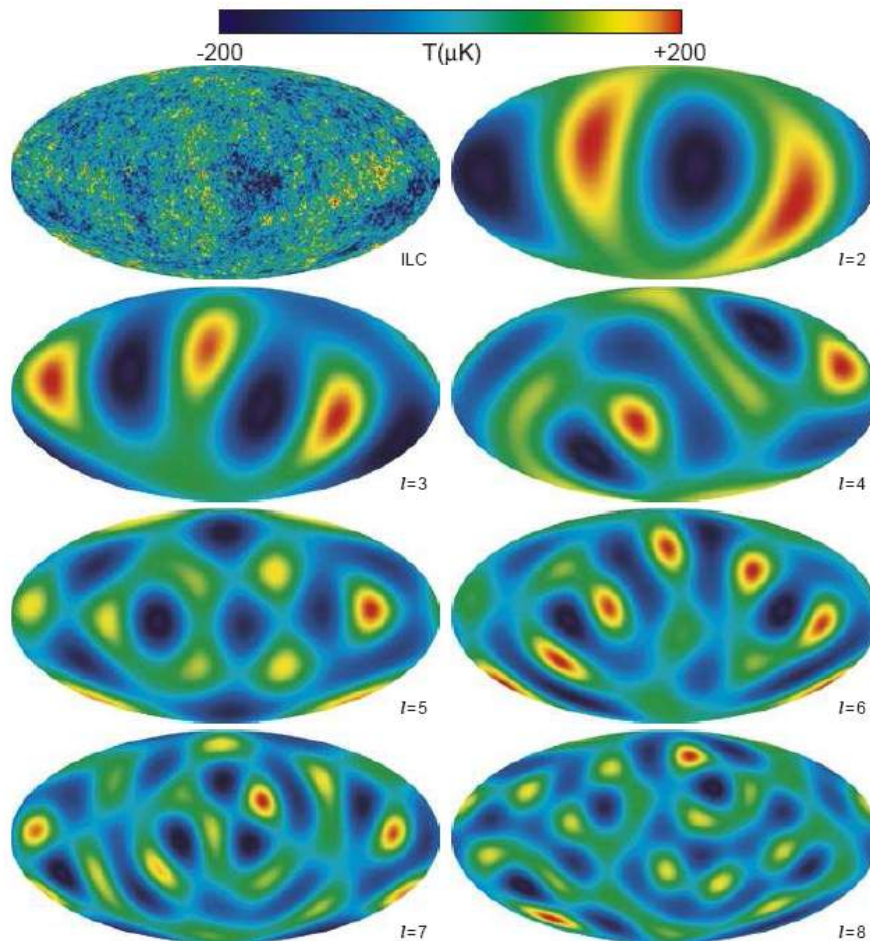
Coherent oscillations in photon-baryon plasma, excited by density perturbations on *super-horizon* scales ...



(Hubble radius at t_{rec})

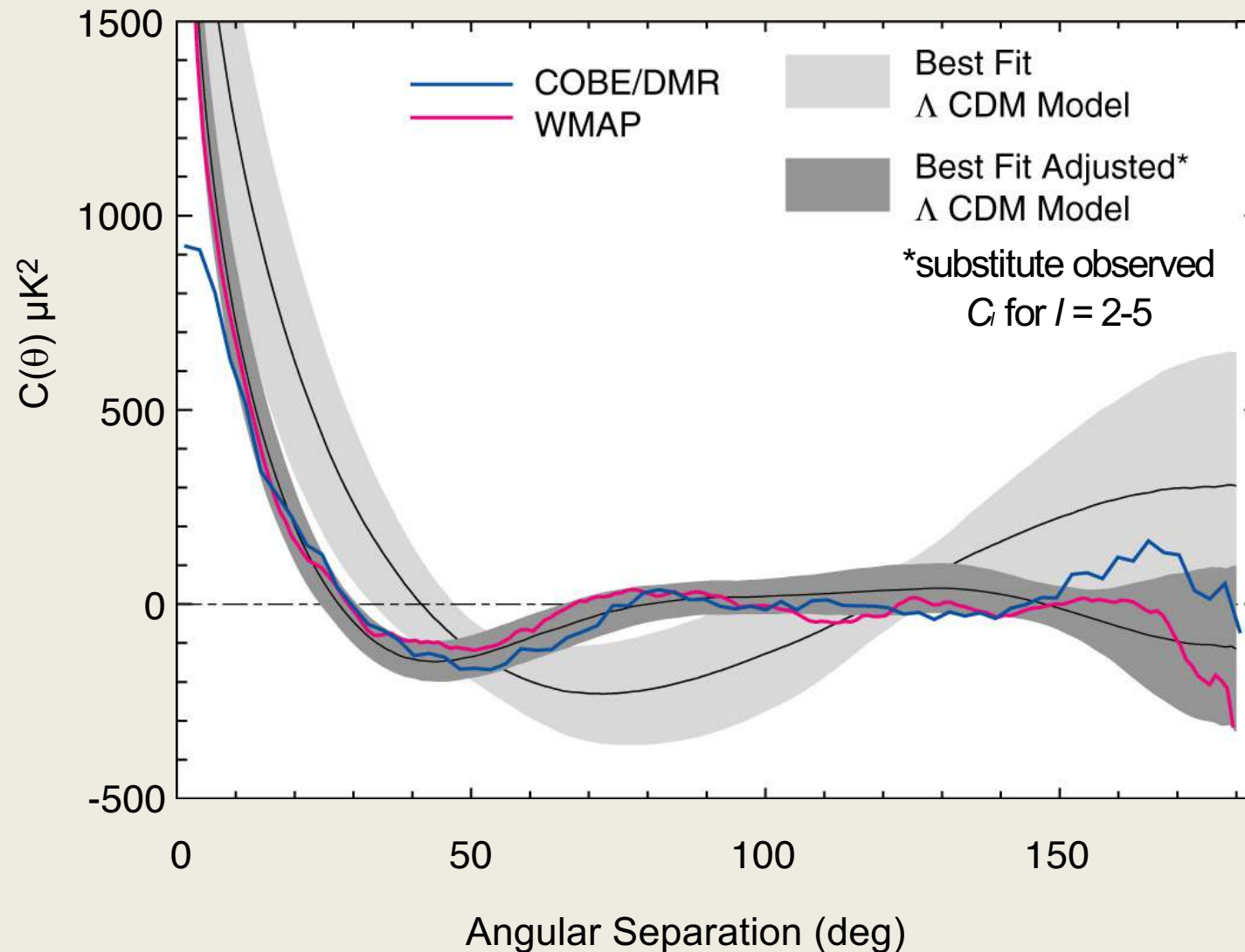
$$\Delta T(\mathbf{n}) = \sum a_{lm} Y_{lm}(\mathbf{n})$$

$$C_l \equiv \frac{1}{2l+1} \sum |a_{lm}|^2$$



The lack of power on large angular scales is most striking, although it is claimed to be *not* unlikely taking cosmic variance and foreground subtraction uncertainties into account

→ chance probability of ~1%



$$C(\theta) = \langle T(n_i)T(n_j) \rangle$$

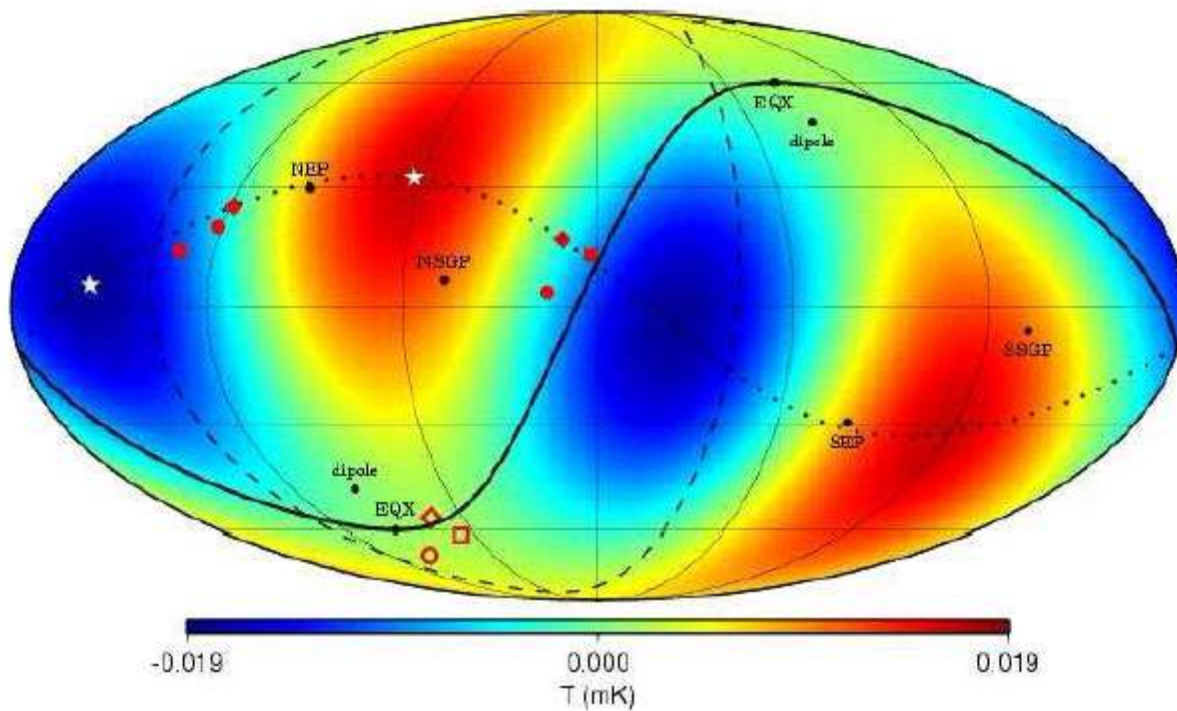
$$n_i \cdot n_j = \cos \theta$$

$$S = \int_{60^\circ}^{180^\circ} C(\theta)^2 d\theta$$

A posteriori
likelihood of
observed S is only
(0.15 - 0.3) %

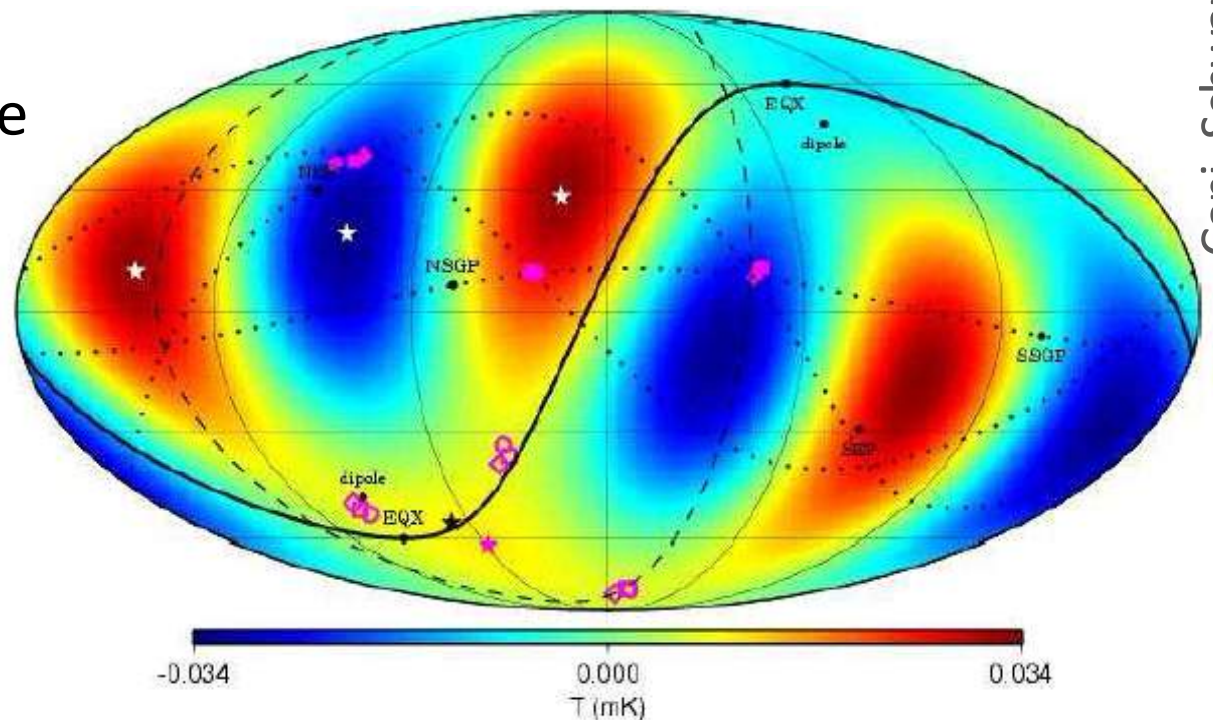
Moreover there is an unexpected alignment of low multipoles, a 'cold spot', and an asymmetry between the North and South ecliptic hemispheres

Curious alignment of quadrupole
and octupole (along the ecliptic)
Power concentrated in plane tilted
by $\sim 30^\circ$ from the Galactic plane
($m = \pm l$ in suitable coord. system)



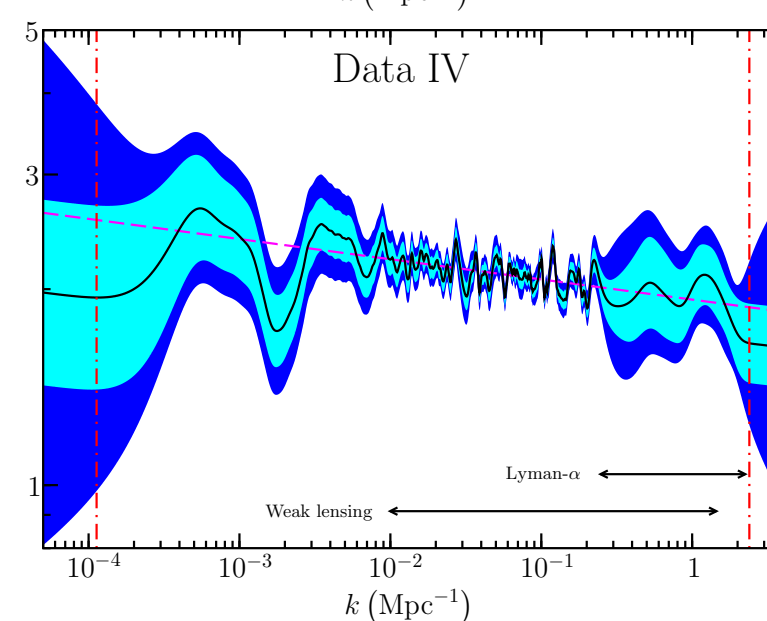
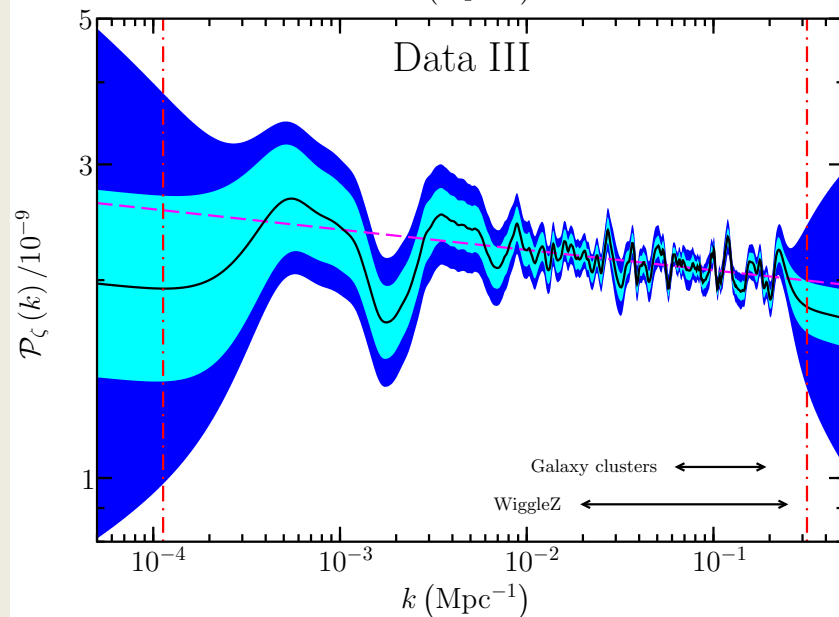
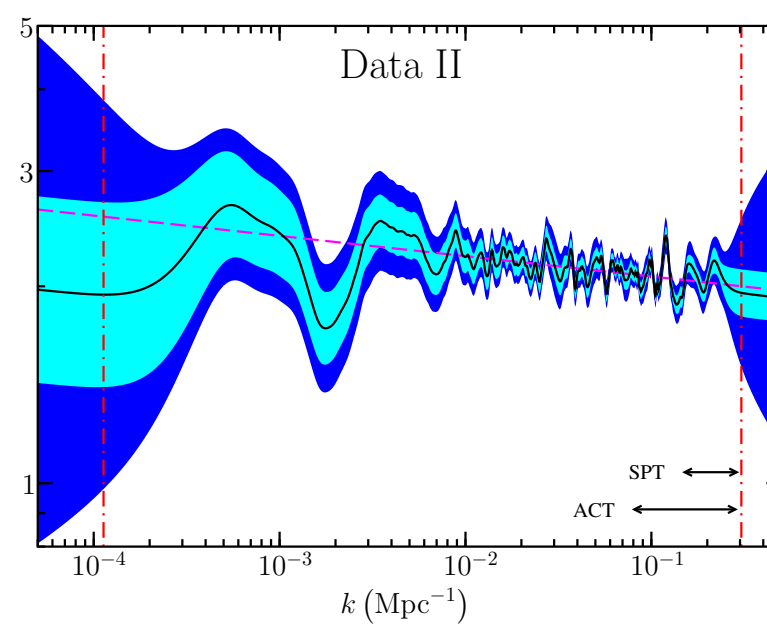
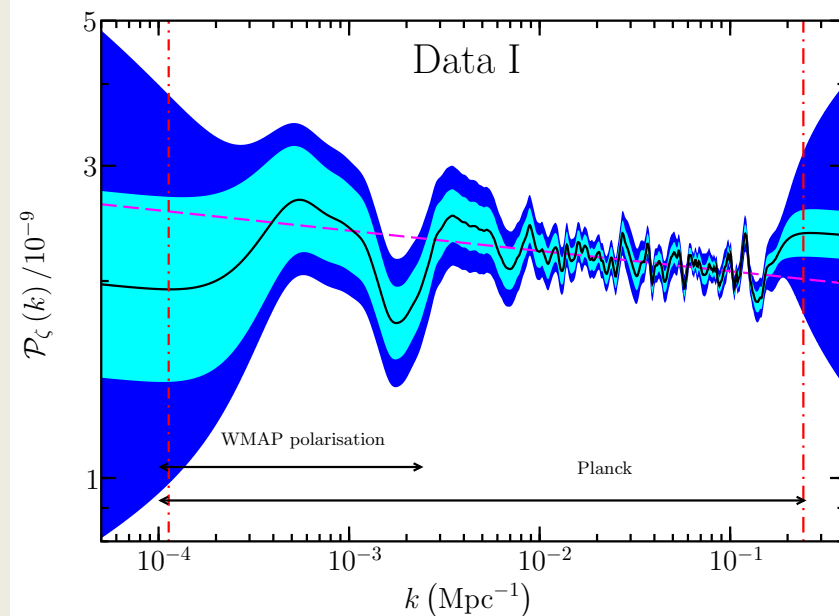
Probability of low quadrupole
+ alignment + “planarity”:
 $\sim 4 \times 10^{-5}$

Tegmark *et al* (2003, 2004)



Copi, Schwarz, Starkman (2004)

The primordial spectrum of perturbations can be deconvoluted from CMB & LSS data *non-parametrically*, using ‘Tikhonov regularisation’ (Hunt & Sarkar, JCAP **12**:052,2015)



This is the
assumed
power-law
spectrum
which is
supposedly
the best-fit
to data:
 $n_s = 0.969$

Comparison with Monte Carlo simulations shows $\sim 2\sigma$ deviations from a power-law spectrum

We can also consider a **direction-dependent** component of the power spectrum of the CMB fluctuations, which is also allowed to vary with the scale (wave number):

$$\mathcal{P}(\mathbf{k}) = \mathcal{P}(k) + \sqrt{4\pi} \sum_{LM} g_{LM}(k) Y_{LM}(\hat{\mathbf{k}})$$

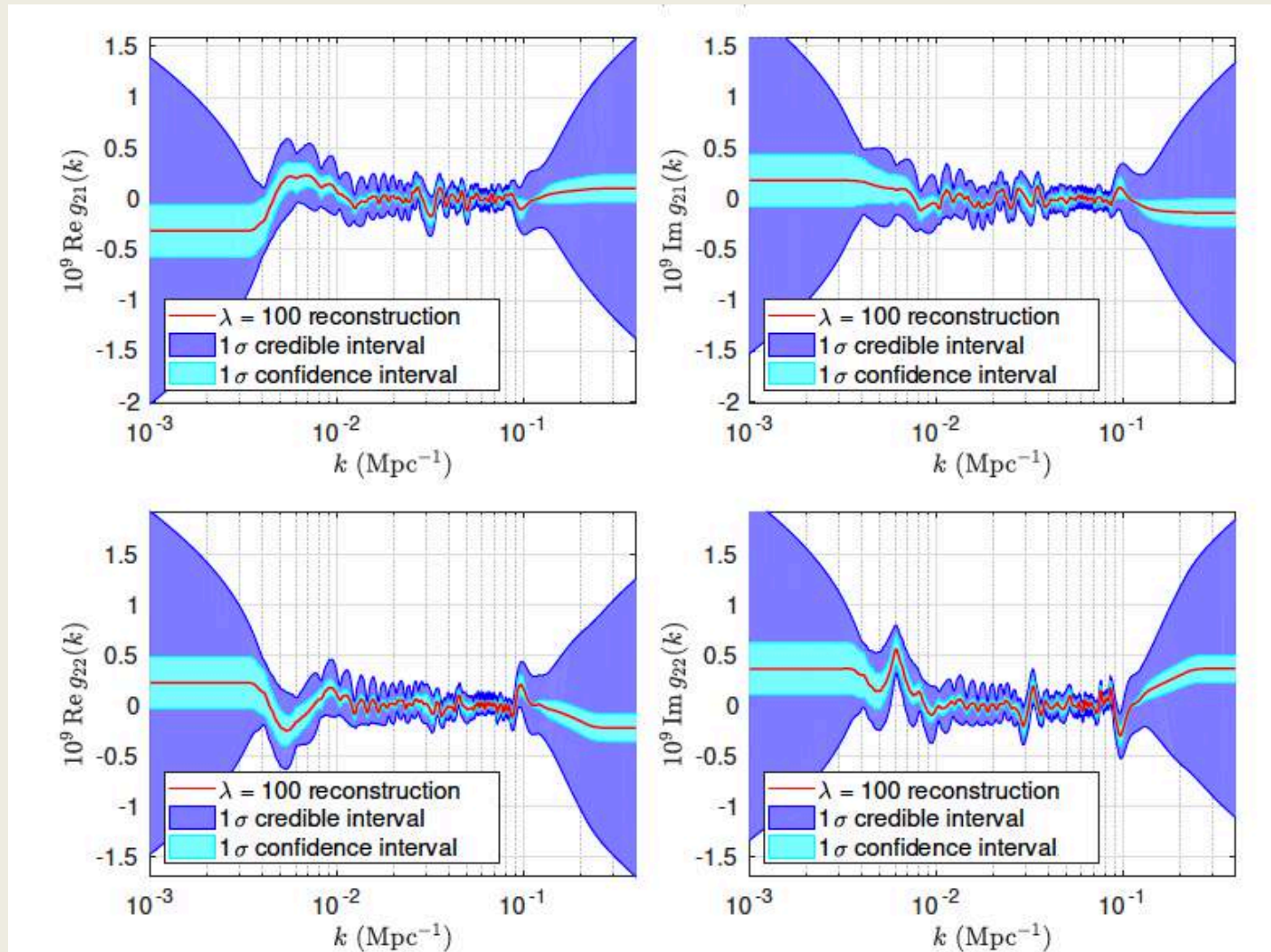
... and focus on the **quadrupole** modulation (NB: Density field is *real*, hence symmetry requires L to be *even* – see Hajian & Souradeep 2005, Pullen & Kamionkowski 2007)

$$\mathcal{P}(\mathbf{k}) = \mathcal{P}(k) + \sqrt{4\pi} \sum_{M=-2}^2 g_{2M}(k) Y_{2M}(\hat{\mathbf{k}})$$

We compute these ‘**bipolar spherical harmonics**’ for the Planck DR2-2015 SMICA map, and estimate the noise covariance from *Planck Full Focal Plane 9* simulations

Previous work by: Groeneboom & Eriksen (2009), Kim & Komatsu (2013); Theoretical models by: Ford (1989), Chibisov (1989), Ackerman *et al* (2007), Pitrou *et al* (2008), Himmetoglu *et al* (2009), Watanabe *et al* (2009), Bartolo *et al* (2013, 2018), ...

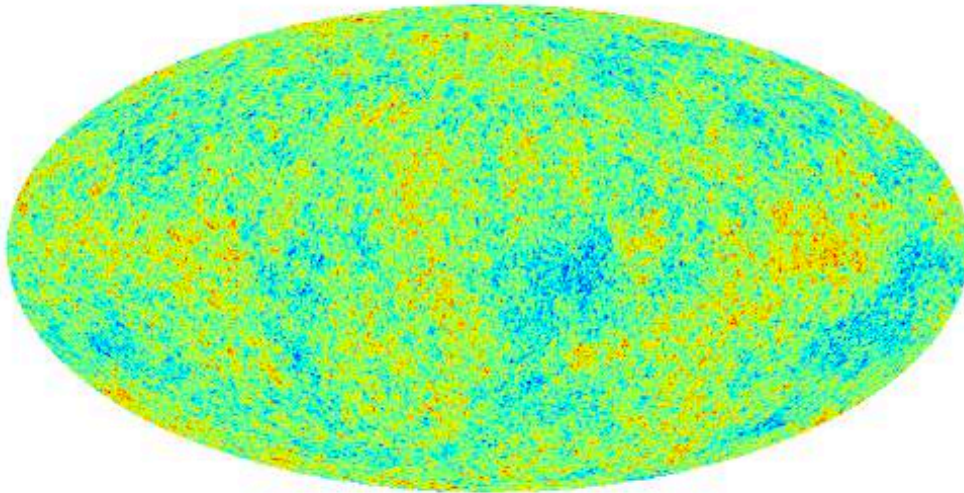
When a *constant quadrupolar modulation* is fitted to Planck data in the range $0.005 \leq k/\text{Mpc}^{-1} \leq 0.008$, its **preferred directions** are found to be *related* to the **cosmic hemispherical asymmetry**, and the **CMB dipole**



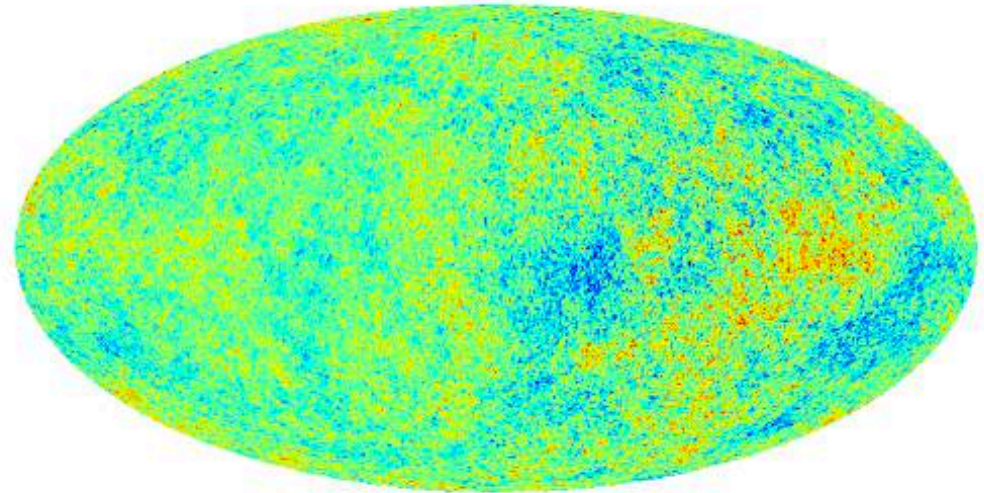
Durakovic et al, JCAP 02:012,2018

The significance is 2.1σ with a test statistic sensitive only to the amplitude of the modulation ... but with a statistic sensitive also to the direction, it rises to 6.9σ !

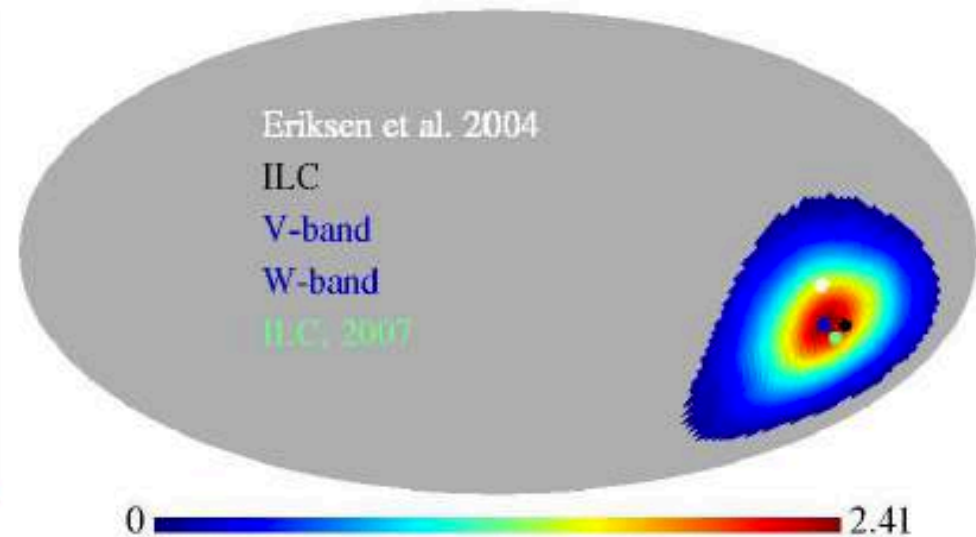
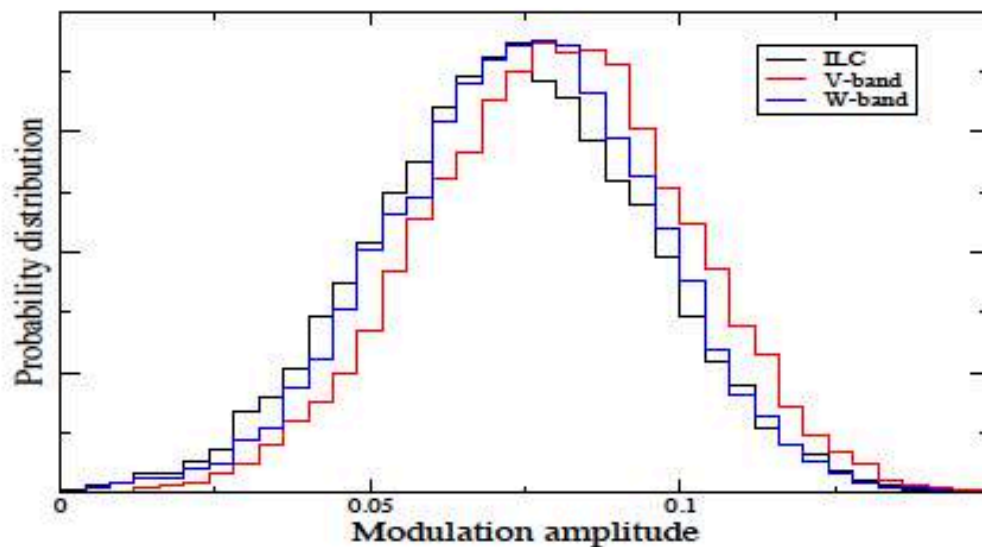
Eriksen *et al* (2004) found that the CMB fluctuations are stronger in one hemisphere of the sky than in the other ($@3\sigma$) ... as if the perturbations are modulated by a *dipole*



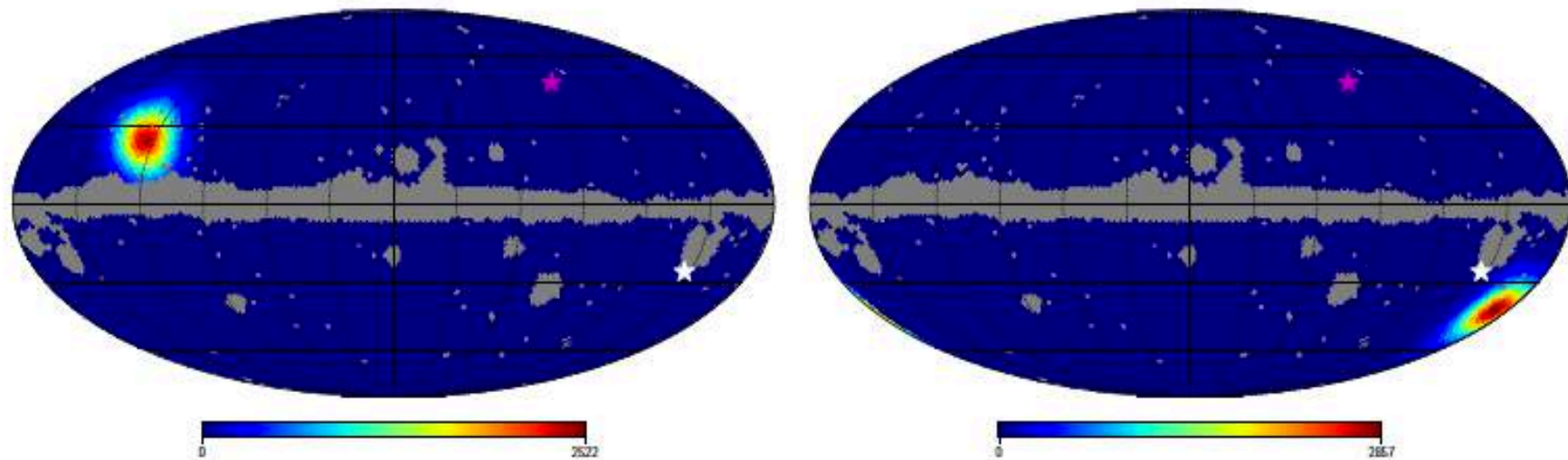
$$T_{\text{iso}}(\hat{\mathbf{n}})$$



$$T_{\text{iso}}(\hat{\mathbf{n}})(1 + A\hat{\mathbf{p}} \cdot \hat{\mathbf{n}})$$



Alignments on the sky : What does this imply for inflation?



Hot ($10^9 g_A = 0.76 \pm 0.22$) quadrupole modulation (left panel),
and cold ($10^9 g_A = -0.82 \pm 0.21$) modulation (right panel).
The magenta and white stars indicate the direction of the CMB
dipole and of the hemispherical asymmetry respectively.

For $k = 0.005\text{-}0.008 \text{ Mpc}^{-1}$:		Angular distances to:	
Amp. $10^9 g_A$	Direction (l, b)	CMB dipole ($264^\circ, 48^\circ$)	Hemisph. asym. ($213^\circ, -26^\circ$)
0.76 ± 0.22	$(128^{+14}_{-14}, 25^{+11}_{-9})$	97°	97°
-0.82 ± 0.21	$(191^{+15}_{-14}, -41^{+10}_{-11})$	110°	24°

CONCLUSIONS

- The distribution of directions of radio galaxies in NVSS and SUMSS is at 2.8σ tension with the kinematic interpretation of the CMB dipole as being due to our motion at 369 km/s w.r.t the CMB rest frame
- There is a *scale-dependent* quadrupolar modulation of CMB anisotropy ... the direction is \sim orthogonal to the CMB dipole but \sim aligned with that of the 'hemispherical asymmetry'

Could all this be an indication of new horizon-scale physics,
e.g. pre-inflationary relics lurking just *outside* our present horizon?
(Gunn 1988, Paczynski & Piran 1990, Turner 1991: 'A tilted universe')

The 'standard' assumptions of *exact* isotropy and homogeneity are *questionable* – data from forthcoming surveys (Euclid, LSST, SKA *etc*) will hopefully provide sufficiently large datasets to enable definitive tests

NB: In the framework of the Λ CDM model, we are a <1% likely observer