

# The Case against Ghosts in Fundamental Theory

RP Woodard

University of Florida

PACTS: June 22, 2018

# Three Quick Questions & Answers

## (Witnessing to the heathen)

1. What is a ghost?
  - Particle with negative KE
2. Why should we avoid ghosts?
  - Interacting ghosts blow up the universe!
3. Why do people (here!) nonetheless consider ghosts?
  - They want to quantize gravity
  - Stelle (1977)  $\rightarrow R + R^2 + C^2$  is renormalizable
  - Higher  $\partial$ 's in  $C^2$  give ghosts!



# How Lower Derivatives Work

- Dynamical variable  $q(t)$  & Lagrangian  $L(q, \dot{q})$ 
  - $\frac{d}{dt} \left( \frac{\partial L}{\partial \dot{q}} \right) = \frac{\partial L}{\partial q}$  initial conditions  $\rightarrow$  2 canonical variables
- Canonical formulation
  - $Q = q$  &  $P = \frac{\partial L}{\partial \dot{q}} \rightarrow \dot{q} = v(Q, P)$  (nondegeneracy)
  - $H(Q, P) = P v(Q, P) - L(Q, v(Q, P))$
- Hamilton's equations generate time evolution
  - $\dot{Q} = \frac{\partial H}{\partial P} = v + P \frac{\partial v}{\partial P} - \frac{\partial L}{\partial \dot{q}} \frac{\partial v}{\partial P} = v$  😊
  - $\dot{P} = -\frac{\partial H}{\partial Q} = -P \frac{\partial v}{\partial Q} + \frac{\partial L}{\partial q} + \frac{\partial L}{\partial \dot{q}} \frac{\partial v}{\partial Q} = \frac{\partial L}{\partial q}$  😊
- $H(Q, P)$  can be bounded below

# Higher Derivatives (Ostrogradsky 1850)

- Lagrangian  $L(q, \dot{q}, \ddot{q})$ 
  - $\frac{d}{dt} \left[ \frac{\partial L}{\partial \dot{q}} - \frac{d}{dt} \left( \frac{\partial L}{\partial \ddot{q}} \right) \right] = \frac{\partial L}{\partial q}$       4 initial conditions  $\rightarrow$  4 canonical coordinates
- Canonical Formulation
  - $Q_1 = q, \quad P_1 = \frac{\partial L}{\partial \dot{q}} - \frac{d}{dt} \left( \frac{\partial L}{\partial \ddot{q}} \right), \quad Q_2 = \dot{q}, \quad P_2 = \frac{\partial L}{\partial \ddot{q}} \rightarrow \ddot{q} = a(\vec{Q}, P_2)$  (ND)
  - $H(\vec{Q}, \vec{P}) = P_1 Q_2 + P_2 a(\vec{Q}, P_2) - L(Q_1, Q_2, a(\vec{Q}, P_2))$
- Hamilton's equations generate time evolution
  - $\dot{Q}_1 = \frac{\partial H}{\partial P_1} = Q_2$  😊,  $\dot{Q}_2 = \frac{\partial H}{\partial P_2} = a + P_2 \frac{\partial a}{\partial P_2} - \frac{\partial L}{\partial \ddot{q}} \frac{\partial a}{\partial P_2} = a$  😊
  - $\dot{P}_2 = -\frac{\partial H}{\partial Q_2} = -P_1 - P_2 \frac{\partial a}{\partial Q_2} + \frac{\partial L}{\partial \dot{q}} + \frac{\partial L}{\partial \ddot{q}} \frac{\partial a}{\partial Q_2} = \frac{\partial L}{\partial \dot{q}} - P_1$  😊
  - $\dot{P}_1 = -\frac{\partial H}{\partial Q_1} = -P_2 \frac{\partial a}{\partial Q_1} + \frac{\partial L}{\partial q} + \frac{\partial L}{\partial \ddot{q}} \frac{\partial a}{\partial Q_1} = \frac{\partial L}{\partial q}$  😊
- $H$  is **linear** in  $P_1 \rightarrow$  not bounded below (or above) ⚡

# Why this is bad

- No guaranteed problem w/o interactions
  - Problem is energy flow from  $KE < 0$  to  $KE > 0$
- No guaranteed problem w/o continuum DoF's
  - Instability is driven by vast  $d^3k$  UV phase space
    - Overwhelms even the weakest nonzero coupling
  - Decay is instantaneous
    - $\tau \neq 0$  results only come from imposing a UV cutoff
- Power & simplicity of the result
  - Requires only non-degenerate HD's
    - Non-perturbative & independent of interactions
  - This is the strongest constraint on Fundamental Theory!
    - “Newton got it right about  $F = ma$ ”

# Common Misconceptions

- “No problem at any constant  $q(t)$ ”
  - Problem is pathological **time dependence**
- “Quantization might help”
  - This is a large phase space problem
- “Problem is unitarity, not instability”
  - Regards negative KE  $a^\dagger(\vec{k})$  as positive KE  $a(\vec{k})$
- “High mass ghosts decouple at low energies”
  - They actually couple more strongly!
- “No problem if HD’s confined to interactions”
  - Problem is non-perturbative
- “No problem from entire functions of  $\partial^2$ ”
  - Only works perturbatively in Euclidean momentum space

# Alternate Quantizations Sacrifice Classical Correspondence Limit

- $L = \frac{1}{2}m\dot{q}^2 - \frac{1}{2}m\omega^2 q^2 \rightarrow H = \frac{1}{2m}P^2 + \frac{1}{2}m\omega^2 Q^2$ 
  - $a = \sqrt{\frac{m\omega}{2\hbar}} \left[ Q + \frac{i}{m\omega} P \right] \rightarrow \sqrt{\frac{m\omega}{2\hbar}} \left[ q + \frac{\hbar}{m\omega} \frac{\partial}{\partial q} \right]$
  - $a^\dagger = \sqrt{\frac{m\omega}{2\hbar}} \left[ Q - \frac{i}{m\omega} P \right] \rightarrow \sqrt{\frac{m\omega}{2\hbar}} \left[ q - \frac{\hbar}{m\omega} \frac{\partial}{\partial q} \right]$
- Normal Quantization:  $\Omega(q) \propto \exp \left[ -\frac{m\omega}{2\hbar} q^2 \right]$ 
  - $H \cdot (a^\dagger)^N \Omega = (N + \frac{1}{2})\hbar\omega \times (a^\dagger)^N \Omega$
- Alternate Quantization:  $\bar{\Omega}(q) \propto \exp \left[ +\frac{m\omega}{2\hbar} q^2 \right]$ 
  - $H \cdot a^N \bar{\Omega} = -(N + \frac{1}{2})\hbar\omega \times a^N \bar{\Omega}$
- ONLY data from low E gravity is classical GR
  - **Dangerous to give this up**
    - You won't get a local, metric theory  $\rightarrow$  causality? Strong fields? Cosmology?
    - IF everything worked  $\rightarrow$  just START with this and forget about HDG!

# Only Hope is Constraints

- Constraints compromise non-degeneracy
- $R \rightarrow f(R)$  gravity ok
  - Ostrogradsky  $\rightarrow$  new DoF of opposite KE
  - But Newtonian potential has negative KE in GR
  - Hence new  $f(R)$  DoF has positive KE
    - NB This is not a counter-example to Ostrogradsky!
- But there are only so many gauge symmetries
- Could always try for ad hoc constraints
  - But at odds with interacting QFT
  - Same field carries both  $\pm$  DoF's



# Lessons from Pop Culture

- “You can’t always get what you want”
  - Face it:  $C^2$  just isn’t viable as a fundamental theory
- “But if you try, sometimes you just might find, that you get what you need”
  - $C \ln(\square)C$  occurs in  $\Gamma_{1loop}$
  - Coefficient finite & fixed
  - Stronger in the IR than  $C^2$



# Conclusions

- Ostrogradsky Thm is the strongest constraint on fundamental theory
- Need to distinguish effective field theory from fundamental theory
  - Fundamental ghosts present at *all* scales
  - Nonlocal EFT effects stronger than local
- Alternate quantization schemes discard the Correspondence Principle
  - This is not acceptable for gravity!