ATLAS measurements of the Higgs boson decay to a pair of τ leptons

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Vyjezdny seminar UCJF

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Standard Model (I)

- ▶ Particle interactions: electromagnetic, weak and strong
- Mass constituents fermions: quarks and leptons
- ► Force carriers bosons: photon, gluon, W and Z bosons



The structure of the Standard Model (I)

Fundamental principle: Local gauge invariance Quantum Electrodynamics (QED) – U(1)

- Free Dirac equation: $i\gamma^{\mu}\partial_{\mu}\psi m\psi = 0$
- ► Lagrangian formalism: $L = i \bar{\psi} \gamma^{\mu} \partial_{\mu} \psi m \bar{\psi} \psi$
- Local gauge transformation: $\psi(x) \rightarrow e^{i\omega(x)}\psi(x)$
- ► term $\partial_{\mu}\omega(x)$ breaks invariance of *L* $\partial_{\mu}\psi(x) \rightarrow e^{i\omega(x)}\partial_{\mu}\psi(x) + ie^{i\omega(x)}\psi(x)\partial_{\mu}\omega(x)$
- ► Invariance of L under local gauge transformations can be accomplished by itroducing a gauge field A_{μ} , which transforms as $A_{\mu}(x) \rightarrow A_{\mu}(x) + \frac{1}{e}\partial_{\mu}\omega(x)$
- Can be formally achieved by the construction of a "modified" derivative

$$\partial_{\mu}
ightarrow D_{\mu} = \partial_{\mu} - \textit{ieA}_{\mu}$$
 (covariant derivative)

The structure of the Standard Model (II)

► Lagrangian of QED:

$$L = -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} + i\bar{\psi}\gamma^{\mu}\partial_{\mu}\psi - m\bar{\psi}\psi + e\bar{\psi}\gamma^{\mu}\psi A_{\mu} \rightarrow -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} + i\bar{\psi}\gamma^{\mu}\partial_{\mu}\psi - m\bar{\psi}\psi - \bar{\psi}\gamma^{\mu}\psi\partial_{\mu}\omega + e\bar{\psi}\gamma^{\mu}\psi A_{\mu} + \bar{\psi}\gamma^{\mu}\psi\partial_{\mu}\omega$$
where $F_{\mu\nu}$ is the usual field strength tensor: $F_{\mu\nu} = \partial_{\mu}A_{\nu} - \partial_{\nu}A_{\mu}$

Note:

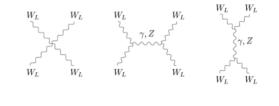
- Interacting term necessary for gauge invariance
- A mass term $(\frac{1}{2}m^2A_{\mu}A^{\mu})$ for the gauge field A_{μ} would violate gauge invariance

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Problems at that stage

Similarly, weak and electromagnetic interactions are described by theory with the symmetry group $SU(2) \times U(1)$

- Such theory results into massless gauge bosons, but we know that W and Z posses some mass
- Divergences in the theory (scattering of W bosons)

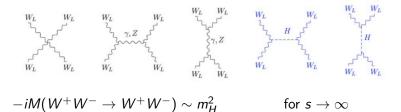


$$-iM(W^+W^- o W^+W^-) \sim rac{s}{M_W^2} ext{ for } s o \infty$$

Higgs mechanism (I)

Solution to **both** problems:

- Create mass via spontaneous breaking of electroweak symmetry
- Introduce a scalar particle that regulates the WW scattering amplitude
 - \blacktriangleright Higgs boson guarantees unitarity (if its mass is $< \sim 1 \mbox{ TeV})$



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Higgs mechanism (II)

- Introduce a new complex doublet field Φ = (^{φ⁺}_{φ⁰}) and add its kinetic term and potential into Lagrangian
 L = L_{EW} + (D_μΦ[†])(D^μΦ) - V(Φ)
 D_μ = ∂_μ - ig₁YB_μ - ig₂Aⁱ_μ ^{σⁱ}/₂
 V(Φ) = -μ²Φ[†]Φ + λ(Φ[†]Φ)²
- B_{μ} and A_{μ}^{i} are linear combination of photon, W^{\pm} and Z bosons fields
- Min. energy (vacuum) for const. field $\Phi^{\dagger}\Phi = \frac{\mu^2}{2\lambda} = \frac{v^2}{2}$
- After gauge transform one scalar field survives (Higgs field), which effectively generates mass terms via the interaction with other particles

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Higgs mechanism (III)

- ► Finally, we get Lagrangian with mass terms for W[±] and Z and a new particle with spin 0 Higgs boson
- ► The same Higgs doublet which generates W[±] and Z masses is sufficient to give masses to the fermions (leptons and quarks)
- The strength of the Higgs interaction is proportional to the particle's mass

$$g_{WWH} = gm_W$$

$$g_{ZZH} = \frac{g}{2\cos\theta_W}m_Z$$

$$g_{ffH} = -\frac{g}{2}\frac{m_f}{m_W}$$

The theory predicts all Higgs boson properties except of its mass

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- ▶ 1967: The theory of electroweak interaction was formulated
- ▶ 1973: Neutral currents neutrino scattering, Gargamelle
- ▶ 1983: W and Z bosons were discovered at SPS (CERN)
- ▶ 2012: Higgs boson was discovered at LHC (CERN)
- But, is it really SM Higgs boson?
 - Measurements of coupling constants HWW, HZZ, $H\tau\tau$, ...
 - Spin and parity measurements
 - ▶ BSM? $H \rightarrow \tau e$ or $H \rightarrow \tau \mu$ Prague involved

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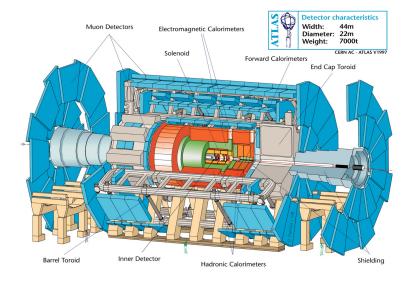
LHC – Large Hadron Collider

- Laboratory CERN, Geneva
- Proton-proton collider
- Circumference 27 km
- Run-1
 - ▶ 2010-2012
 - Int. luminosity $25 fb^{-1}$
 - ► Energy 7, 8 TeV
- Run-2
 - ► 2015 -
 - ▶ Int. luminosity 38*fb*⁻¹
 - ► Energy 13 TeV



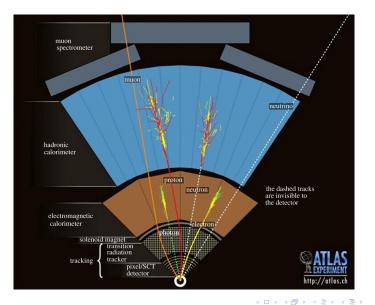
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Experiment ATLAS (I)



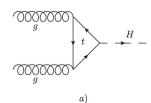
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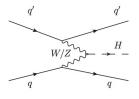
Experiment ATLAS (II)



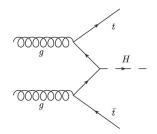
Higgs production at LHC

Processes: ggH – dominant, VBF (10x less frequent), VH, ttH





b)

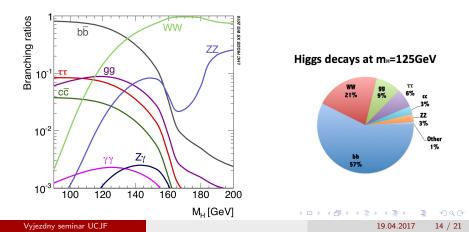


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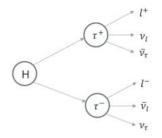
Higgs boson decays

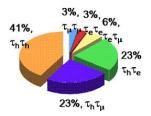
- $H
 ightarrow b ar{b}$ dominant, but very difficult to measure
- Discovery: $H \rightarrow ZZ^*$, $H \rightarrow WW^*$, $H \rightarrow \gamma\gamma$
- Important to measure Higgs coupling to fermions: $H \rightarrow \tau \tau$
 - \blacktriangleright 5 σ significance when combining ATLAS and CMS Run-1 results



$H \rightarrow \tau^+ \tau^-$ decay modes

- Di-lepton: $H \to \tau^+ \tau^- \to \ell^+ \ell^- 4\nu$
- ▶ Semi-lepton: $H \rightarrow \tau^+ \tau^- \rightarrow n\pi \ell 3\nu \ (1 3\pi)$
- Hadronic: $H \rightarrow \tau^+ \tau^- \rightarrow n\pi 2\nu \ (2-6\pi)$





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$H \to \tau^+ \tau^- \to \ell^+ \ell^- 4\nu$

- Signal selection and cutflow Tomas' presentation
- ▶ Dominant background processes $H \rightarrow \tau^+ \tau^- \rightarrow \ell^+ \ell^- 4\nu$:
 - $\blacktriangleright \ Z \to \tau^+ \tau^- \to \ell^+ \ell^- 4\nu$
 - $Z \to \ell^+ \ell^-$
 - $t \rightarrow bW (\rightarrow \ell \nu)$
- Mass reconstruction
 - 4 neutrinos in the final state
 - not possible to fully reconstruct invariant mass
 - Collinear approximation
 - neutrinos parallel to tau leptons
 - mass calculated using MET
 - MMC (missing mass calculator)
 - consider probability angular distribution of the decay products
 - search for the most probable decay configuration for given lepton momenta and MET

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Statistical analysis of the Higgs boson decay (I)

Poisson distribution applied to a histogram with i bins

$$P(n_i|
u_i(oldsymbol{ heta})) = rac{
u_i(oldsymbol{ heta})^{n_i}}{n_i!}e^{-
u_i(oldsymbol{ heta})}$$

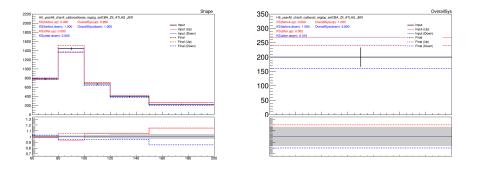
- Probability to observe n_i events when $\nu_i = \mu s_i + b_i$ are expected
- s_i and b_i are the SM expected values for signal and background bin content (may be function of θ)
- Nuisance parameters θ are related to statistical and systematic uncertainties and to the normalisation of background contributions measured in control data samples
- Pol is signal strength $\mu = \frac{\sigma}{\sigma_{SM}}$
- ► Over 100 nuisance parameters, their values θ and stat. deviations σ measured by *combined performance group* → Gauss term:

$$L = \sum_{i} P(n_i | \mu s_i(\theta) + b_i(\theta)) G(\theta, \sigma))$$

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Statistical analysis of the Higgs boson decay (II)

- Impact of different sources of systematic uncertainties is expressed in terms of relative changes of the expected event yields
- Uncertainty is obtained by varying a given experimental or theoretical quantity by ±1 standard deviation around the nominal value

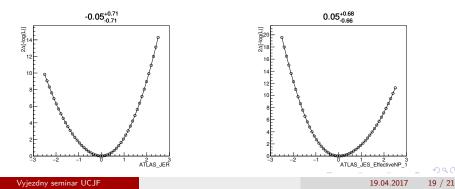


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Statistical analysis of the Higgs boson decay (III)

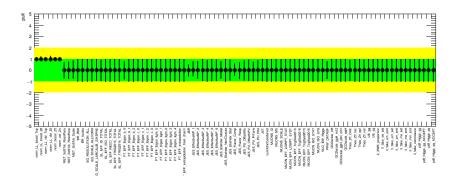
• Profile likelihood ratio: $\lambda(\mu) = \frac{L(\mu,\hat{\theta}(\mu))}{L(\hat{\mu},\hat{\theta}(\mu))}$

- Conditional max likelihood estimator L(μ, θ̂(μ)) provides max likelihood estimate of the NP θ for a given value of μ
- Unconditional max likelihood estimator L(μ̂, θ̂(μ)) corresponds to the best fit result
- Negative log-likelihood: $-2\ln\lambda(\mu)$



Statistical analysis of the Higgs boson decay (IV)

Variety of checks are performing to identify unwanted nuisance parameter's behaviour, such as strong correlations, NP pulled far from their nominal values in the fit, NP which exhibit double-minima or non-parabolic behaviour in negative log-likelihood distribution



- Higgs boson decay to a pair of τ leptons is proven above 5σ significance when combining ATLAS and CMS Run-1 results
- The goal in Run-2 is to improve precision, with so far acquired data in 2015+2016 ATLAS should be able to reach 5σ. Hope to be ready for summer conferences
- Apart from Higgs boson decays predicted by SM, searches for other decay modes as well as for eventual other Higgs bosons (e.g. MSSM Higgs) continue

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