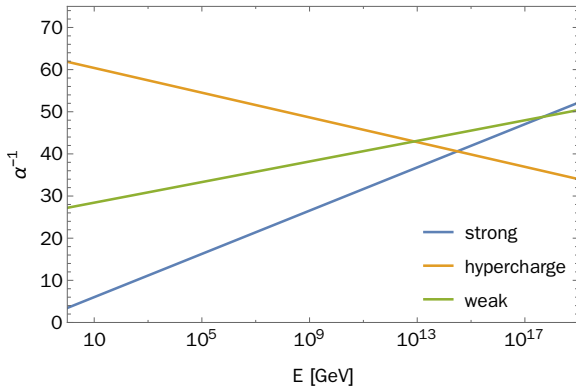


# The group $SO(10)$ in the Grand Unified Theory

Kateřina Jarkovská

Výjezdni seminář ÚČJF

- Standard model gauge group  $G_{SM} = SU(3) \times SU(2) \times U(1)$
- SM running couplings (one loop level, without scalar fields):



$$\alpha^{-1}(M) = \alpha^{-1}(M_0) + \frac{1}{2\pi} \left( -\frac{11}{3} C_2^{adj} + \frac{2}{3} T_2(R) \right) \ln \left( \frac{M}{M_0} \right)$$

# Special orthogonal group $SO(n)$

## Definition

If  $B : V \times V \rightarrow \mathbb{F}$  is a symmetrical positive definite bilinear form. Then we define Orthogonal group

$$O(n, \mathbb{F}) := \{M \in GL(V) : B(Mv, Mw) = B(v, w) \text{ for } v, w \in V\},$$

where  $\mathbb{F} = \mathbb{R}$  or  $\mathbb{C}$ .

$$O(n) \cong \{M \in \mathbb{R}^{n \times n} : M^T M = I_n\}$$

$$SO(n) = O(n) \cap \{M \in \mathbb{R}^{n \times n} : \det M = 1\}$$

- $SO(n)$  is connected and compact  $\Rightarrow$  expressed by an exponential of generators from special orthogonal algebra  $\mathfrak{so}(n)$ .

$$\mathfrak{so}(n) := \{X \in \mathbb{R}^{n \times n} : X^T + X = 0, \text{Tr}(X) = 0\}$$

## Definition

Representation  $(\phi, V)$  of a Lie algebra  $\mathfrak{g}$  is a linear map

$$\phi : \mathfrak{g} \rightarrow \mathfrak{gl}(V)$$

which preserves Lie bracket:

$$[\phi(X), \phi(Y)] = \phi([X, Y])$$

for all  $X, Y \in \mathfrak{g}$ . If  $\langle \cdot, \cdot \rangle$  is a Hermitian scalar product on  $V$  and

$$\forall X \in \mathfrak{g}, v_1, v_2 \in V : \langle v_1, \phi(X)v_2 \rangle = -\langle \phi(X)v_1, v_2 \rangle,$$

then  $(\phi, V)$  is an unitary representation.

- Representation of  $\mathfrak{so}(n) \longleftrightarrow$  representation of a double cover  $Spin(n) \longleftrightarrow$  (single, double valued) representation of  $SO(n)$

# Spinor representation of $\mathfrak{so}(2n)$

## Definition

Clifford algebra  $Cliff(2n, \mathbb{C})$  is an associative algebra generated by an identity 1 and vectors from an orthonormal basis of  $\mathbb{C}^{2n}$  satisfying

$$e_i e_j + e_j e_i = 2\delta_{ij}$$

for  $i, j \in \{1, \dots, 2n\}$

- Witt's basis is for  $j \in \{1, \dots, n\}$

$$w_j := \frac{1}{2}(e_{2j-1} + ie_{2j})$$

$$\bar{w}_j := \frac{1}{2}(e_{2j-1} - ie_{2j})$$

- $\{w_j, w_k\} = \{\bar{w}_j, \bar{w}_k\} = 0, \quad \{w_j, \bar{w}_k\} = \delta_{jk}$

- Exterior algebra

$$\Lambda\mathbb{C}^n = \mathbb{C}^n \oplus \Lambda^1\mathbb{C}^n \oplus \Lambda^2\mathbb{C}^n \dots \oplus \Lambda^n\mathbb{C}^n$$

where  $\Lambda\mathbb{C}^j$  is generated by

$$w_{k_1} \wedge w_{k_2} \cdots \wedge w_{k_j} = \text{sign}(\sigma) w_{\sigma(k_1)} \wedge w_{\sigma(k_2)} \cdots \wedge w_{\sigma(k_j)}$$

- Spinor representation on exterior algebra  $\Lambda\mathbb{C}^n$  is given by an action

$$w_j : \Lambda\mathbb{C}^n \rightarrow \Lambda\mathbb{C}^n$$

$$w_{k_1} \wedge \cdots \wedge w_{k_p} \mapsto w_j \wedge w_{k_1} \wedge \cdots \wedge w_{k_p}$$

and

$$\bar{w}_j : \Lambda\mathbb{C}^n \rightarrow \Lambda\mathbb{C}^n$$

$$w_{k_1} \wedge \cdots \wedge w_{k_p} \mapsto \sum_{s=1}^p \delta_{j s} (-1)^{s+1} w_{k_1} \wedge \cdots \wedge w_{k_{s-1}} \wedge w_{k_{s+1}} \wedge \cdots \wedge w_{k_p}$$

- $\dim \Lambda\mathbb{C}^n = 2^n$

- $\mathfrak{so}(2n) \subset \text{Cliff}(2n, \mathbb{C})$  is generated by

$$\{w_j w_k, \bar{w}_j \bar{w}_k, \bar{w}_j w_k\}_{j,k=1}^n = \{[e_j, e_k]\}_{j,k=1}^n$$

- Action of  $\mathfrak{so}(2n)$  preserves parity of the grades  $\Rightarrow$  spinor representation of  $\mathfrak{so}(2n)$  is reducible  $\Rightarrow$  decomposition into irreducible subspaces:

$$\Lambda \mathbb{C}^n = \Lambda^{\text{even}} \mathbb{C}^n \oplus \Lambda^{\text{odd}} \mathbb{C}^n$$

where

$$\Lambda^{\text{even}} \mathbb{C}^n = \mathbb{C}^n \oplus \Lambda^2 \mathbb{C}^n \oplus \dots$$

and

$$\Lambda^{\text{odd}} \mathbb{C}^n = \Lambda^1 \mathbb{C}^n \oplus \Lambda^3 \mathbb{C}^n \dots$$

- fermionic annihilation and creation operators  $\{f_j^\dagger, f_j\}_{j=1}^n$ :

$$\{f_j^\dagger, f_k^\dagger\} = \{f_j, f_k\} = 0, \quad \{f_j^\dagger, f_k\} = \delta_{jk}$$

- vacuum state  $|0\rangle$ :  $f_j|0\rangle = 0$  for all  $j \in \{1, \dots, n\}$
- fermionic Fock space

$$\mathcal{F}_n^+ = \left\langle |0\rangle, f_{j_1}^\dagger|0\rangle, f_{j_1}^\dagger f_{j_2}^\dagger|0\rangle, \dots, f_{j_1}^\dagger f_{j_2}^\dagger \dots f_{j_n}^\dagger|0\rangle \right\rangle$$

## Fermionic realization:

$$\begin{aligned} w_j &\longleftrightarrow f_j^\dagger \\ \bar{w}_j &\longleftrightarrow f_j \\ \mathfrak{so}(2n) &\cong \langle w_j w_k, \bar{w}_j \bar{w}_k, \bar{w}_j w_k \rangle \longleftrightarrow \langle f_j f_k, f_j^\dagger f_k^\dagger, f_j^\dagger f_k \rangle \\ \Lambda \mathbb{C}^n &\longleftrightarrow \mathcal{F}_n^+ \\ \Lambda^{\text{even}} \mathbb{C}^n &\longleftrightarrow \mathcal{F}_{\text{even}}^+ := \langle |0\rangle, f_{j_1}^\dagger f_{j_2}^\dagger|0\rangle, \dots \rangle \\ \Lambda^{\text{odd}} \mathbb{C}^n &\longleftrightarrow \mathcal{F}_{\text{odd}}^+ := \langle f_{j_1}^\dagger|0\rangle, f_{j_1}^\dagger f_{j_2}^\dagger f_{j_3}^\dagger|0\rangle, \dots \rangle \end{aligned}$$



# Standard model representation

- For one generation ( $Q = T_3 + Y$ ):

$$\begin{pmatrix} \nu_L \\ e_L \end{pmatrix} \longleftrightarrow \mathbb{C} \otimes \mathbb{C}^2 \otimes \mathbb{C}_{-\frac{1}{2}} \longleftrightarrow (1, 2, -\frac{1}{2})$$

$$\begin{pmatrix} u_L^R & u_L^G & u_L^B \\ d_L^R & d_L^G & d_L^B \end{pmatrix} \longleftrightarrow \mathbb{C}^3 \otimes \mathbb{C}^2 \otimes \mathbb{C}_{\frac{1}{6}} \longleftrightarrow (3, 2, \frac{1}{6})$$

$$(e_R) \longleftrightarrow \mathbb{C} \otimes \mathbb{C} \otimes \mathbb{C}_{-1} \longleftrightarrow (1, 1, -1)$$

$$(\nu_R) \longleftrightarrow \mathbb{C} \otimes \mathbb{C} \otimes \mathbb{C}_0 \longleftrightarrow (1, 1, 0)$$

$$(u_R^R \quad u_R^G \quad u_R^B) \longleftrightarrow \mathbb{C}^3 \otimes \mathbb{C} \otimes \mathbb{C}_{\frac{2}{3}} \longleftrightarrow (3, 1, \frac{2}{3})$$

$$(d_R^R \quad d_R^G \quad d_R^B) \longleftrightarrow \mathbb{C}^3 \otimes \mathbb{C} \otimes \mathbb{C}_{-\frac{1}{3}} \longleftrightarrow (3, 1, -\frac{1}{3})$$

$$G_{SM} = SU(3) \times SU(2) \times U(1)$$

- Fermion representation of one generation:

$$F = (1, 2, -\frac{1}{2}) \oplus (3, 2, \frac{1}{6}) \oplus (1, 1, 0) \oplus (1, 1, -1) \oplus (3, 1, \frac{2}{3}) \oplus (3, 1, -\frac{1}{3})$$

- Antifermion representation of one generation:

$$F^* = (1, 2, \frac{1}{2}) \oplus (\bar{3}, 2, -\frac{1}{6}) \oplus (1, 1, 0) \oplus (1, 1, 1) \oplus (\bar{3}, 1, -\frac{2}{3}) \oplus (\bar{3}, 1, \frac{1}{3})$$

- Standard model representation of one fermion and antifermion generation:

$$F \oplus F^*$$

- $\dim F \oplus F^* = 32 = 2^5$

# SO(10) in the Grand Unified Theory

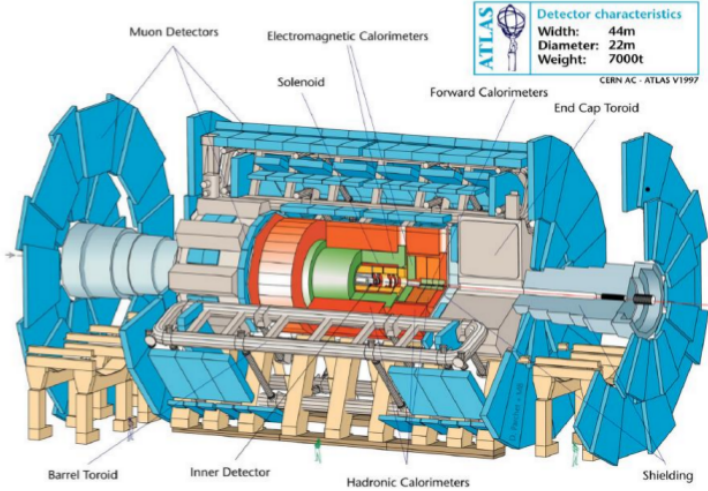
- $\dim F \oplus F^* = 2^5 = \dim \Lambda \mathbb{C}^5$
- spinor representation on  $\Lambda^{odd} \mathbb{C}^5, \Lambda^{even} \mathbb{C}^5$
- Left-handed Weyl spinors:

$$\Lambda^{even} \mathbb{C}^5 \cong (\bar{\nu}_L) \oplus (e_L^+) \oplus \begin{pmatrix} u_L^R & u_L^G & u_L^B \\ d_L^R & d_L^G & d_L^B \end{pmatrix} \oplus (\bar{u}_L^R \quad \bar{u}_L^G \quad \bar{u}_L^B) \oplus \begin{pmatrix} \nu_L \\ e_L \end{pmatrix} \oplus (\bar{d}_L^R \quad \bar{d}_L^G \quad \bar{d}_L^B)$$

- Right-handed Weyl spinors:

$$\Lambda^{odd} \mathbb{C}^5 \cong (e_R^+ \quad \bar{\nu}_R) \oplus (d_R^R \quad d_R^G \quad d_R^B) \oplus (e_R) \oplus \begin{pmatrix} \bar{u}_L^R & \bar{u}_L^G & \bar{u}_L^B \\ \bar{d}_L^R & \bar{d}_L^G & \bar{d}_L^B \end{pmatrix} \oplus \begin{pmatrix} u_R^R & u_R^G & u_R^B \end{pmatrix} \oplus (\nu_R)$$





## References:

- Baez, J., Huerta, J.: The Algebra of Grand Unified Theories. *Bulletin of American Mathematical Society* year 2010, 47: p. 483-552.
- Georgi, H.: *Lie algebras in Particle Physics*. ABP, 1999, ISBN: 0-7382-0233-9
- Langacker, P.: Grand Unified Theories and proton decay. *Physics Reports* year 1981, 72: 185-385.
- Souček, V.: *Reprezentace Lieových grup a algeber*, 2002.