

Gravitational Wave Signals of Electroweak Phase Transition Triggered by Dark Matter

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In collaboration with Wei Chao and Jing Shu,

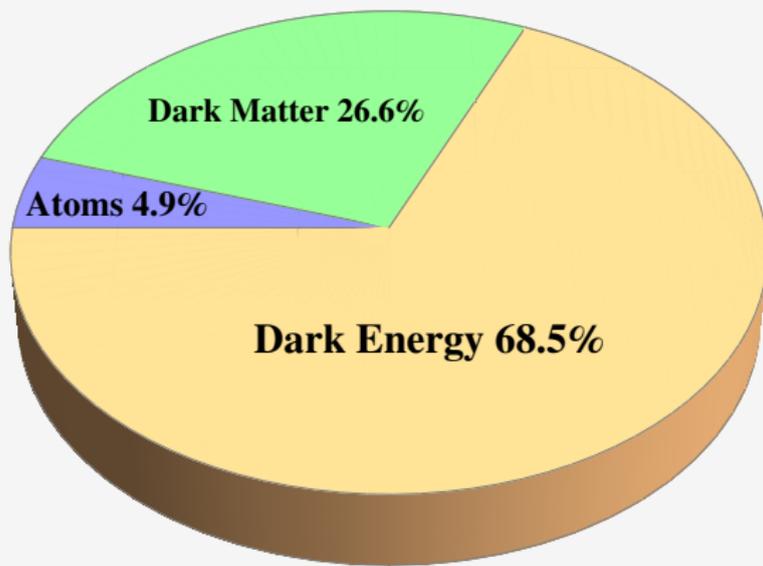
(arXiv:1702.02698)

ITP-CAS

May 8, 2017

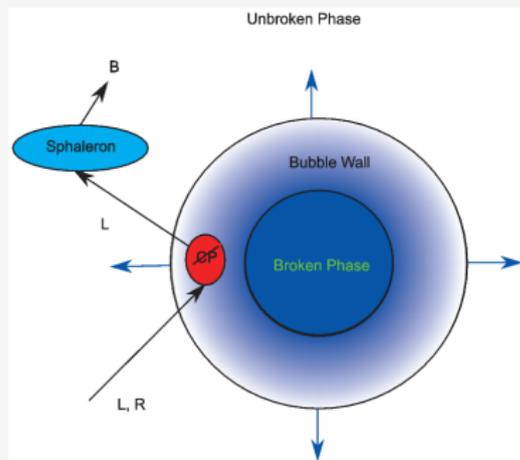
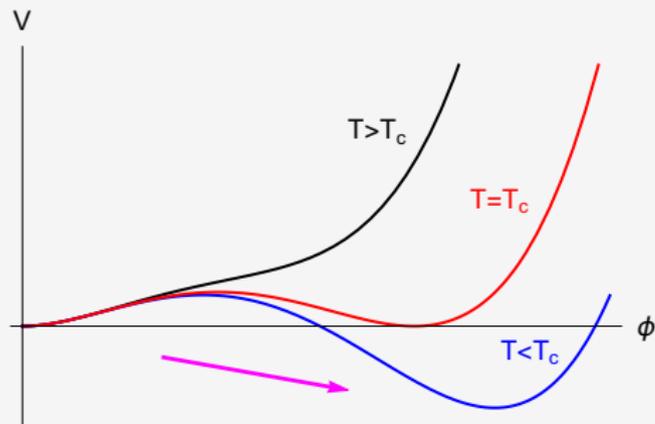


Baryon Asymmetry and Dark Matter



Planck, 2013. *Astron.Astrophys.* (2014)

Electroweak Baryogenesis: The Standard Picture

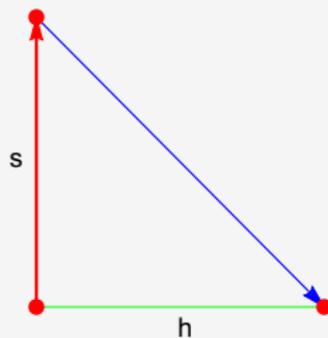


$T \approx 100\text{GeV} \approx 10^{15}\text{K}$

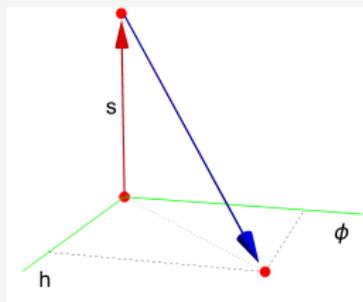
Gravitational Waves (mHz level)

Minimal Extensions of the SM: Dark Matter and EWPT

SM + S (stringent experimental tension)



SM + S + S



SM + DM + Scalar: Model at $T = 0$

A simplified $Z_2 \times Z_2$ symmetric potential:

$$V_0 = -\frac{1}{2}\mu_\Phi^2\Phi^2 + \frac{1}{4}\lambda_\Phi\Phi^4 - \frac{1}{2}\mu_S^2S^2 + \frac{1}{4}\lambda_S S^4 \\ -\mu^2 H^\dagger H + \lambda(H^\dagger H)^2 + \lambda_1 S^2 H^\dagger H + \lambda_2 \Phi^2 H^\dagger H \\ + \lambda_3 S^2 \Phi^2,$$

- $v_S = 0 \Rightarrow$ s is stable
- h, ϕ mix: $\hat{h} = c_\theta h + s_\theta \phi, \quad \hat{\phi} = -s_\theta h + c_\theta \phi. \quad (c_\theta > 0.86 \text{ at } 95\% \text{ CL})$

Physical Parameters:

$\mu_\Phi, \mu_S, \mu, \lambda_\Phi, \lambda_S, \lambda, \lambda_1, \lambda_2, \lambda_3$



$v_H, m_H, v_\Phi, m_{\hat{\phi}}, m_S, \lambda_S, \theta, \lambda_1, \lambda_3$



7 free parameters

Dark Matter: Turn Off SI Scattering at Tree Level

DM-Nucleon scattering cross section:

$$\sigma_n = \frac{\mu^2 m_n^2}{\pi v_{\text{EW}}^2 m_S^2} \left| \frac{c_\theta a_{\hat{h}}}{m_{\hat{h}}^2} - \frac{s_\theta a_{\hat{\phi}}}{m_{\hat{\phi}}^2} \right|^2 \left(\frac{2}{9} + \frac{7}{9} \sum_{q=u,d,s} f_{T_q}^n \right)^2$$

Two sets of diagrams:

$$\begin{aligned} a_{\hat{h}} : s^2 \hat{h} & : 2\lambda_3 v_\phi s_\theta + \lambda_1 v_H c_\theta \\ a_{\hat{\phi}} : s^2 \hat{\phi} & : 2\lambda_3 v_\phi c_\theta - \lambda_1 v_H s_\theta \end{aligned}$$

Zoomin to the parameter space with a **negligible** direct detection signal:

$$\lambda_3 = \frac{v_H \lambda_1 (m_{\hat{h}}^2 \tan \theta + m_{\hat{\phi}}^2 \cot \theta)}{2v_\Phi (m_{\hat{h}}^2 - m_{\hat{\phi}}^2)} \Rightarrow \boxed{6 \text{ free parameters}}$$

Effective Potential: Gauge Dependence

$$V_{\text{eff}}^T(\phi) = V^{T=0}(\phi) + V_{\text{CW}}^{T=0}(\phi) + \frac{T^4}{2\pi^2} \left[\sum_{\text{scalars}} J_B\left(\frac{M^2(\xi)}{T^2}\right) + 3 \sum_{\text{gauge}} J_B\left(\frac{\mu^2}{T^2}\right) - \sum_{\text{gauge}} J_B\left(\frac{\xi\mu^2}{T^2}\right) - 4 \sum_{\text{fermions}} n_C^f J_F\left(\frac{m_f^2}{T^2}\right) \right],$$

ξ : Gauge fixing parameter in R_ξ -gauge.

$O(T^2)$ terms are ξ -independent & Equivalent to including thermal masses at tree level:

$$\begin{aligned}\Pi_h &= \left\{ \frac{3g^2 + g'^2}{16} + \frac{\lambda}{2} + \frac{h_t}{4} + \frac{\lambda_1 + \lambda_2}{12} \right\} T^2, \\ \Pi_s &= \left\{ \frac{\lambda_s}{4} + \frac{\lambda_1}{3} + \frac{\lambda_3}{6} \right\} T^2, \\ \Pi_\phi &= \left\{ \frac{\lambda_\phi}{4} + \frac{\lambda_2}{3} + \frac{\lambda_3}{6} \right\} T^2.\end{aligned}$$

More on this topic: H.Patel, M. Ramsey-Musolf, JHEP 1107 (2011) 029.

Analysis of $\mathcal{O}(T^2)$ Terms

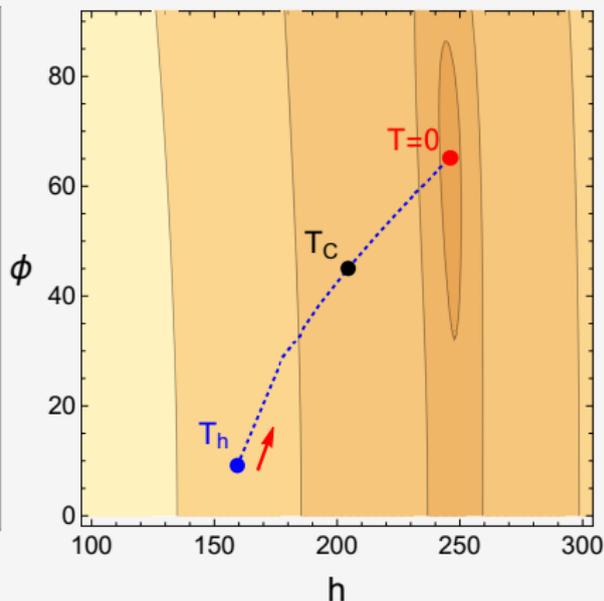
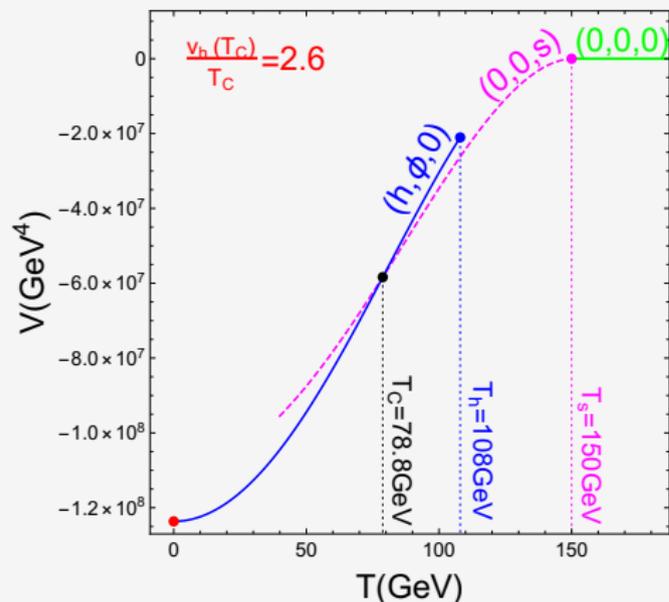
Generic Form:

$$\begin{aligned} V_{\text{eff}}^T(\phi) = & D_h(T^2 - T_h^2)h^2 + \Lambda_h h^4 \\ & + D_\phi(T^2 - T_\phi^2)\phi^2 + \Lambda_\phi \phi^4 \\ & + D_s(T^2 - T_s^2)s^2 + \Lambda_s s^4 \\ & + \delta_{h\phi} h^2 \phi^2 + \delta_{hs} h^2 s^2 + \delta_{\phi s} \phi^2 s^2 \end{aligned}$$

- Basics: $\Lambda_h > 0$, $\Lambda_\phi > 0$, $\Lambda_s > 0$, $\delta_{h\phi} > 0$, $\Lambda_{hs} > 0$, $\Lambda_{\phi s} > 0$.
- Symmetry Restoration: $D_h > 0$, $D_\phi > 0$, $D_s > 0$, $T_h^2 > 0$, $T_\phi^2 > 0$, $T_s^2 > 0$,
- Two minima: $(0, 0, s)$ and $(h, \phi, 0)$ (Hessian Positive Definite)
- DM develops a vev first: $T_s > T_h$, $T_s > T_\phi$
- EW minimum $(h, \phi, 0)$ stable at $T = 0$
- ▶ These two minima will cross at some $T(\equiv T_c)$ in the phase diagram.

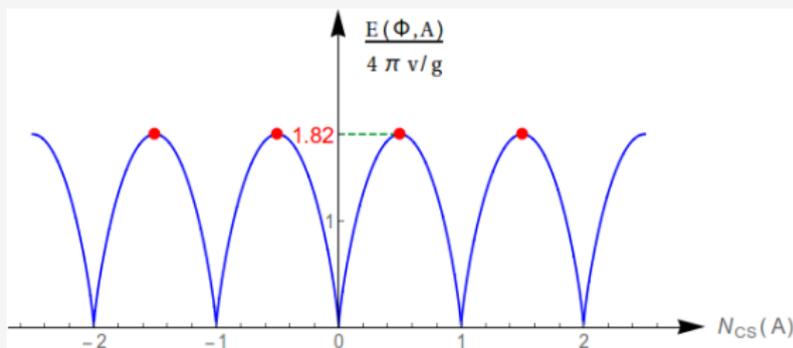
Phase History: A Benchmark Point

$v_\Phi = 65\text{GeV}$, $m_{\hat{\phi}} = 82\text{GeV}$, $m_S = 71\text{GeV}$, $\lambda_S = 0.015$, $\theta = 0.12$, $\lambda_1 = 0.046$ and $\lambda_3 = 0.57$

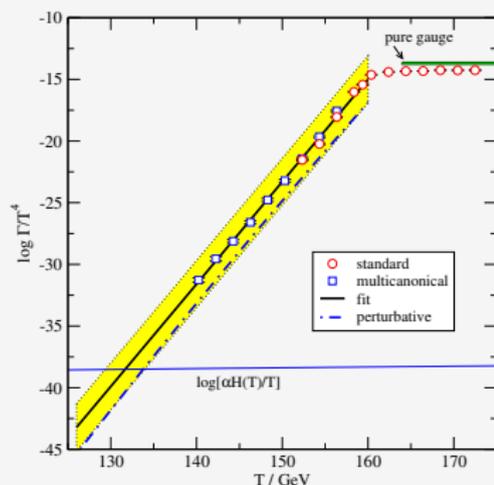


◆ The first step could also be a 1st order phase transition with a more complete V_{eff} .

Avoid Baryon Washout



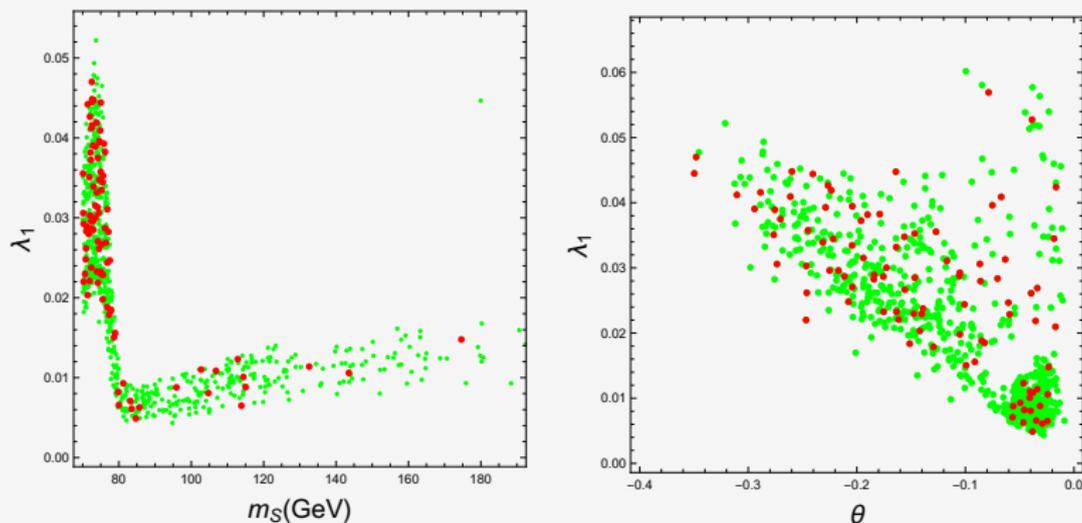
$$\Gamma^{\text{sym}} \approx 6 \times (18 \pm 3) \alpha_W^5 T^4, \quad \Gamma^{\text{brok}} \sim T^4 \exp\left(-\frac{E_{\text{sph}}}{T}\right)$$



Sphaleron rate from lattice in SM,
Phys.Rev.Lett,113, 141602 (2014).

Strongly first order EWPT criteria: $\frac{v_h(T_C)}{T_C} > 1$ (ξ -dependence)

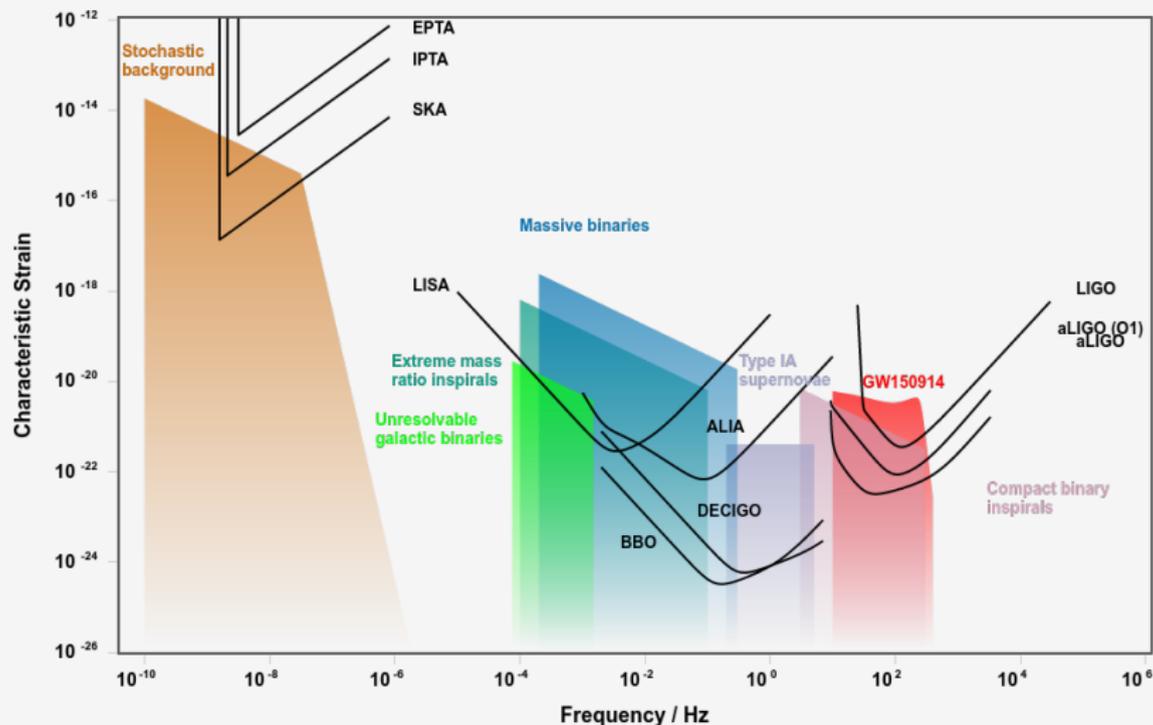
A Survey of Parameter Space



All gives previous EWPT pattern and no SI DM detection signal at tree level.

- $\Omega_c h^2 \in (0.03, 0.12)$
- $\Omega_c h^2 \in (0.03, 0.12)$ and $\frac{v_h(T_C)}{T_C} > 1$

Gravitational Waves: Experimental



C J Moore et al. *Class. Quantum Grav.* 32 (2015) 015014.

Gravitational Waves: A Benchmark Point



$$\Omega_{\text{GW}} h^2 \simeq \Omega_{\phi} h^2 + \Omega_{\text{sw}} h^2 + \Omega_{\text{turb}} h^2$$

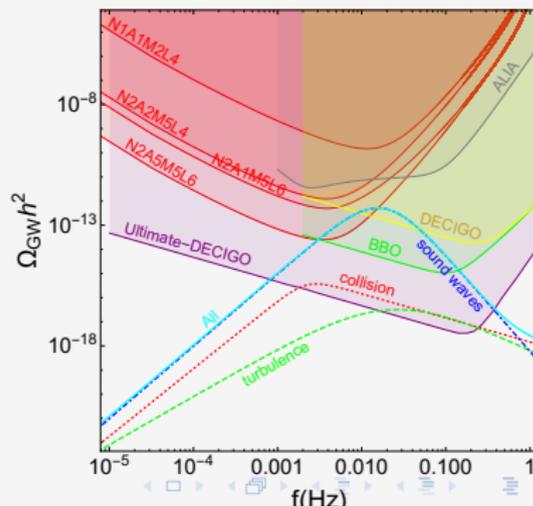
Three sources:

- Bubble Collision
- Sound Waves
- Turbulence

T_n : bubble nucleation temperature

α : strength of phase transition

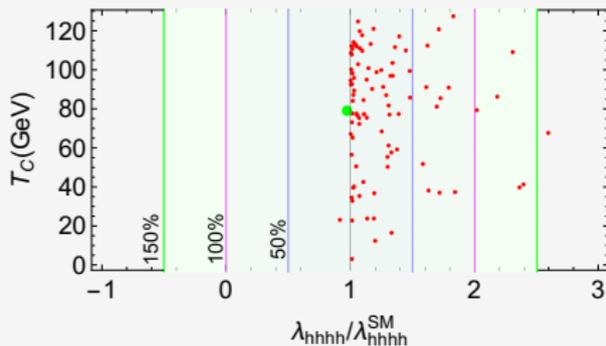
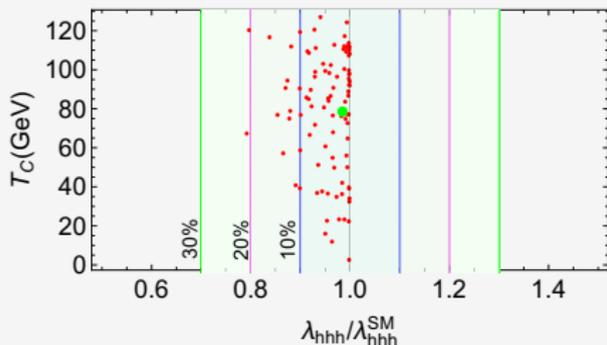
β : inverse duration of phase transition



Higgs Couplings

$$\lambda_{\hat{h}\hat{h}\hat{h}} = -\frac{3m_{\hat{h}}^2}{v_H} \left[1 - \frac{3\theta^2}{2} + \mathcal{O}(\theta^3) \right]$$

$$\lambda_{\hat{h}\hat{h}\hat{h}\hat{h}} = -\frac{3m_{\hat{h}}^2}{v_H^2} \left[1 + \left(\frac{m_{\hat{\phi}}^2}{m_{\hat{h}}^2} - 3 \right) \theta^2 + \mathcal{O}(\theta^3) \right]$$



- $\Omega_{ch^2} \in (0.03, 0.12)$ and $\frac{v_h(T_C)}{T_C} > 1$
- GW benchmark point.

Conclusion

- Dark matter itself could have experienced a phase transition.
- Dark matter might have played a role in EWPT.
- Gravitational waves thus produced can be probed by LISA, etc.,.
- Higgs self-couplings measurement serves as a complementary test.

Thanks