

# Surveying the Left-Right Symmetric $SO(10)$ Landscape

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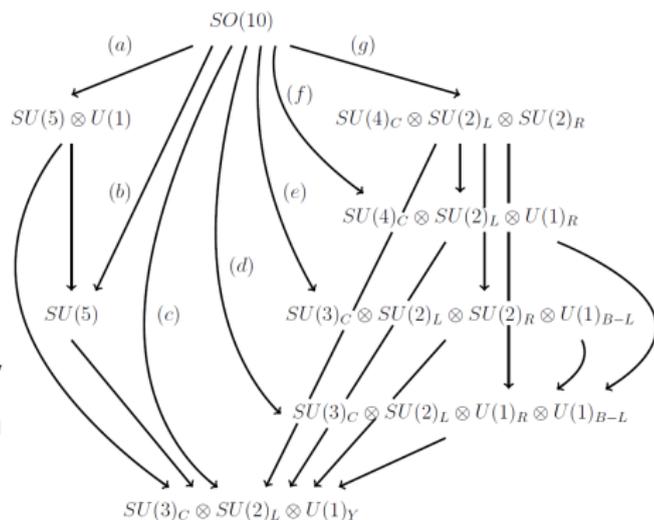
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- current state - lack of new physics signals near the electroweak scale
- GUTs - well motivated extensions of the Standard Model
- sound way how to explore very high scales
- $SO(10)$  - very rich framework, many free parameters  $\rightarrow$  many possible scenarios, but minimal settings still phenomenologically interesting
- $\implies$  a general, automated approach to the model building procedure

- specify theory at high energies  $\leftrightarrow$  symmetries + field content
- series of symmetry breaking steps reducing the symmetry down to the SM gauge group
- fermions - only  $16_F$
- scalars - exotics allowed, each scalar can survive or be integrated out at every breaking scale
- manageable computation  $\leftrightarrow$  up to 5 scalars at every level allowing for the given breaking chain
- duplicates generated are kept, may correspond to different Lagrangians



- model generating algorithm: a set of reps at GUT scale decomposed into reps of the subsequent groups of the chain
- $\implies$  all possible combinations satisfying the mentioned restrictions obtained
- for a chain  $\mathcal{G} \rightarrow \mathcal{F}_1 \rightarrow \dots \rightarrow \mathcal{F}_{m-1} \rightarrow \mathcal{G}_{SM}$  with  $m$  breaking steps the models are characterized as

$$\{\mathcal{M}\} = \left\{ \begin{array}{ll} \text{Chain: } \{\mathcal{G} \rightarrow \dots \rightarrow \mathcal{G}_{SM}\}, & \text{Reps: } \{\mathcal{R}_i^{(0)}\}, \\ \text{Chain: } \{\mathcal{F}_1 \rightarrow \dots \rightarrow \mathcal{G}_{SM}\}, & \text{Reps: } \{\mathcal{R}_i^{(1)}\}, \\ \dots & \dots \\ \text{Chain: } \{\mathcal{G}_{SM}\}, & \text{Reps: } \{\mathcal{R}_i^{(m)}\} \end{array} \right\}$$

$\{\mathcal{R}_i^{(0)}\}$  - the representations at the  $SO(10)$  scale,

$\{\mathcal{R}_i^{(j)}\}$  - a combination of their decompositions at the  $j$ th step

- analysis of all possible scenarios - exhausting  $\rightarrow$  test case
- LR-symmetric models - minimal realisations extensively described in literature, a lot of pheno interest

$\implies$  consider two-step symmetry breaking

$$\begin{array}{c} SO(10) \\ \downarrow \\ SU(3)_C \otimes SU(2)_L \otimes SU(2)_R \otimes U(1)_{B-L} \\ \downarrow \\ SU(3)_C \otimes SU(2)_L \otimes U(1)_Y \end{array}$$

- representations at  $SO(10)$  scale:  $\{3 \times \mathbf{16}_F, \mathbf{10}, \mathbf{45}, \mathbf{126}, \overline{\mathbf{126}}\}$

- chirality - complex representations needed  $\rightarrow \mathbf{16}_F$
- anomalies - in our context: ABJ anomaly relevant - satisfied by construction: no exotic fermions,  $\mathbf{16}_F$  of  $SO(10)$  is anomaly free
- symmetry breaking - at every step we require presence of a scalar field, which can break the symmetry in the desired way
- Standard Model - last step of the chain is the SM gauge group with (at least) SM content - in addition to a Higgs doublet also extra scalar fields allowed (such that are not immediately ruled out)
- fermion mixing, CKM matrix - checking specific values beyond our scope, but enough degrees of freedom to fit the parameters must be provided (minimal choice for the LR scale:  $2 \times \{\mathbf{1}, \mathbf{2}, \mathbf{2}, 0\} + \{\mathbf{1}, \mathbf{1}, \mathbf{3}, 0\}$ )

- 1-loop RGE for the gauge coupling  $g$ :  $\mu \frac{d\alpha^{-1}}{d\mu} = -\frac{b}{2\pi}$
- the slope  $b$  is given by:

$$b = \frac{2}{3} \sum_{\text{Fermions}} S(\mathcal{R}_f) d_{\perp}(\mathcal{R}_f) + \frac{1}{3} \sum_{\text{Scalars}} S(\mathcal{R}_s) d_{\perp}(\mathcal{R}_s) - \frac{11}{3} C_2(\mathcal{G})$$

- substituting  $t = \frac{1}{2\pi} \log\left(\frac{\mu}{M_Z}\right)$  and solving RGEs corresponding to 3 SM couplings  $\rightarrow$  3 conditions, in a matrix form:

$$\alpha^{-1} \equiv \begin{pmatrix} \alpha_3^{-1} \\ \alpha_2^{-1} \\ \alpha_1^{-1} \end{pmatrix} = \begin{pmatrix} 1 & b_1^3 & b_2^3 & \dots & b_m^3 \\ 1 & b_1^2 & b_2^2 & \dots & b_m^2 \\ 1 & b_1^1 & b_2^1 & \dots & b_m^1 \end{pmatrix} \cdot \begin{pmatrix} \alpha_{GUT} \\ \Delta t_1 \\ \Delta t_2 \\ \vdots \\ \Delta t_m \end{pmatrix} \equiv B_0 \cdot \Delta t$$

- solution dependent on number of intermediate scales  $m$ 
  - for  $m = 2$  system gives a unique solution for scales and  $\alpha_{GUT}$
  - for  $m > 2 \rightarrow$  underdetermined system,  $m - 2$  free parameters

- 3 generations of right-handed neutrino  $\nu^c$  singlet always present, assumed to live at LR scale  $\rightarrow$  seesaw type I
- if Higgs triplet  $\Delta$  present  $\rightarrow$  seesaw type II
- current experimental bound:  $m_\nu^{\text{exp}} = \sum m_{\nu_i} \lesssim 0.3 \text{ eV}$ ; lower limit from the atmospheric mass splitting  $\sqrt{\Delta m_{\text{atm}}^2} \approx 0.05 \text{ eV}$
- $\rightarrow$  assume a conservative range of the masses

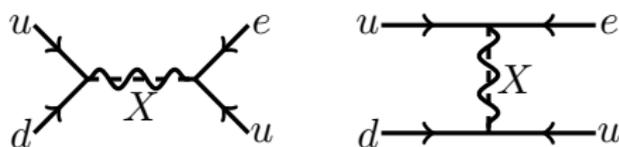
$$0.16 < \frac{m_\nu}{m_\nu^{\text{exp}}} \approx |\lambda_\Delta - y_\nu^2| \left( \frac{2 \times 10^{14} \text{ GeV}}{m_{LR}} \right) < 1$$

- to generate the light neutrino masses, the LR scale  $\approx 10^{14} \text{ GeV}$ , with couplings  $\lambda_\Delta, y_\nu = \mathcal{O}(1)$

- proton decay - typically 6D eff. operators - could be mediated by a scalar or a gauge boson and suppressed by  $M_X^{-2}$

	SM	LR	$SO(10)$
gauge	$\{\mathbf{3}, \mathbf{2}, -\frac{5}{6}\}, \{\mathbf{3}, \mathbf{2}, \frac{1}{6}\}$	$\{\mathbf{3}, \mathbf{2}, \mathbf{2}, -\frac{2}{3}\}$	<b>45, 54</b>
scalar	$\{\mathbf{3}, \mathbf{3}, -\frac{1}{3}\}, \{\mathbf{3}, \mathbf{1}, -\frac{1}{3}\}$	$\{\mathbf{3}, \mathbf{3}, \mathbf{1}, -\frac{2}{3}\}, \{\mathbf{3}, \mathbf{1}, \mathbf{3}, -\frac{2}{3}\}$ $\{\mathbf{3}, \mathbf{1}, \mathbf{1}, -\frac{2}{3}\}$	<b>10, 120</b> <b>126, <math>\overline{126}</math></b>

- decay channel  $p \rightarrow e^+ \pi^0$   
for a gauge or scalar  
boson mediator  $X$



- estimation of the p decay half-life ( $\tau_p^{\text{exp}} > 1.29 \times 10^{34}$  yr):

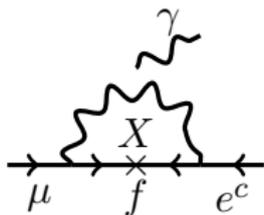
$$\frac{\tau_p^{\text{gauge}}}{\tau_p^{\text{exp}}} \approx \frac{1}{\alpha_X^2} \left( \frac{M_X}{2.6 \times 10^{16} \text{ GeV}} \right)^4, \quad \frac{\tau_p^{\text{scalar}}}{\tau_p^{\text{exp}}} \approx \frac{1}{\lambda^4} \left( \frac{M_X}{7.3 \times 10^{15} \text{ GeV}} \right)^4$$

- also  $n - \bar{n}$  oscillations - 9D eff. operators - scalar mediators, suppression by  $M_X^{-6}$ ; less restrictive than proton decay

- 6D eff. operators - mediated by a scalar or a gauge boson and suppressed by  $M_X^{-2}$

- experimental limits:  $\mathcal{B}(\mu \rightarrow e\gamma)_{\text{exp}} < 5.7 \times 10^{-13}$   
 $\mathcal{B}(\mu \rightarrow eee)_{\text{exp}} < 1.0 \times 10^{-12}$   
 $\mathcal{B}(\mu N \rightarrow eN)_{\text{exp}} < 7 \times 10^{-13}$

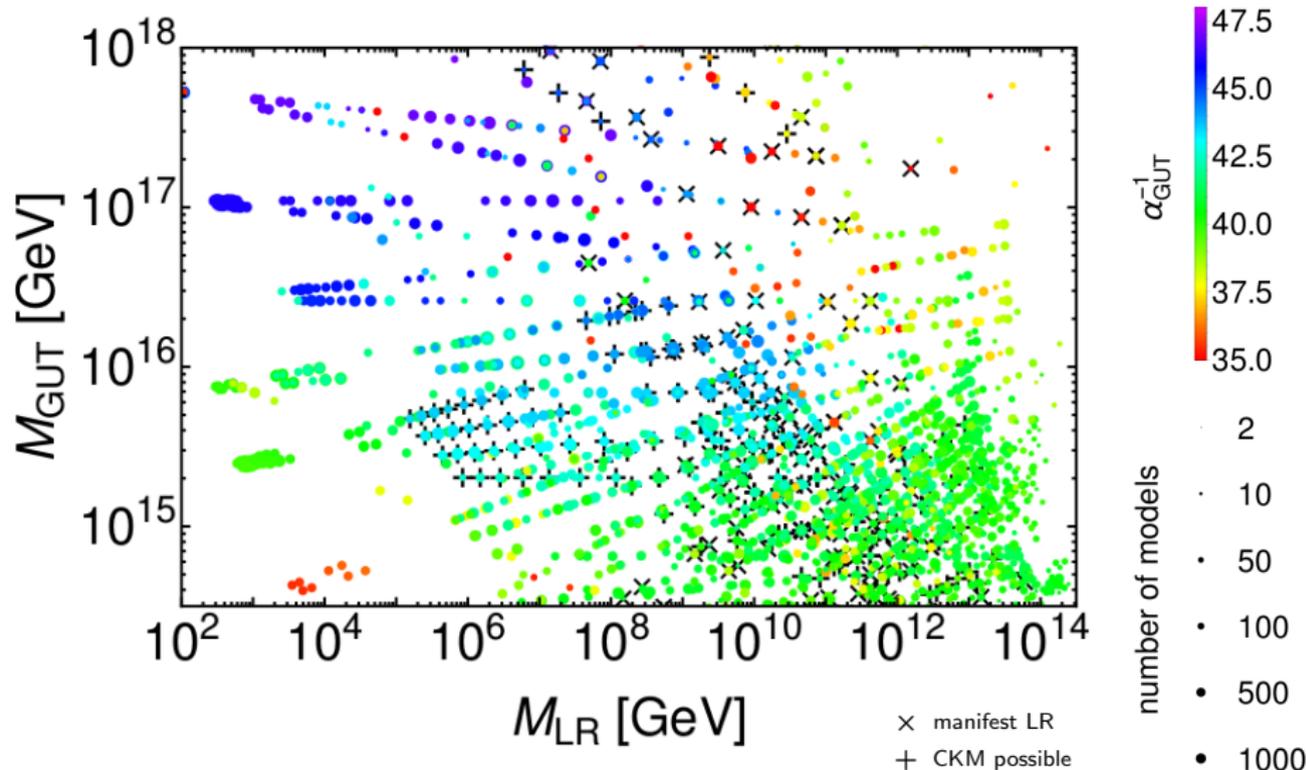
- dominant contributions:  $\frac{\mathcal{B}(\mu \rightarrow e\gamma)}{\mathcal{B}(\mu \rightarrow e\gamma)_{\text{exp}}} \approx |g_{e\mu}|^2 \left(\frac{g_R}{g_L}\right)^4 \left(\frac{2.3 \times 10^4 \text{ GeV}}{M_{LR}}\right)^4$

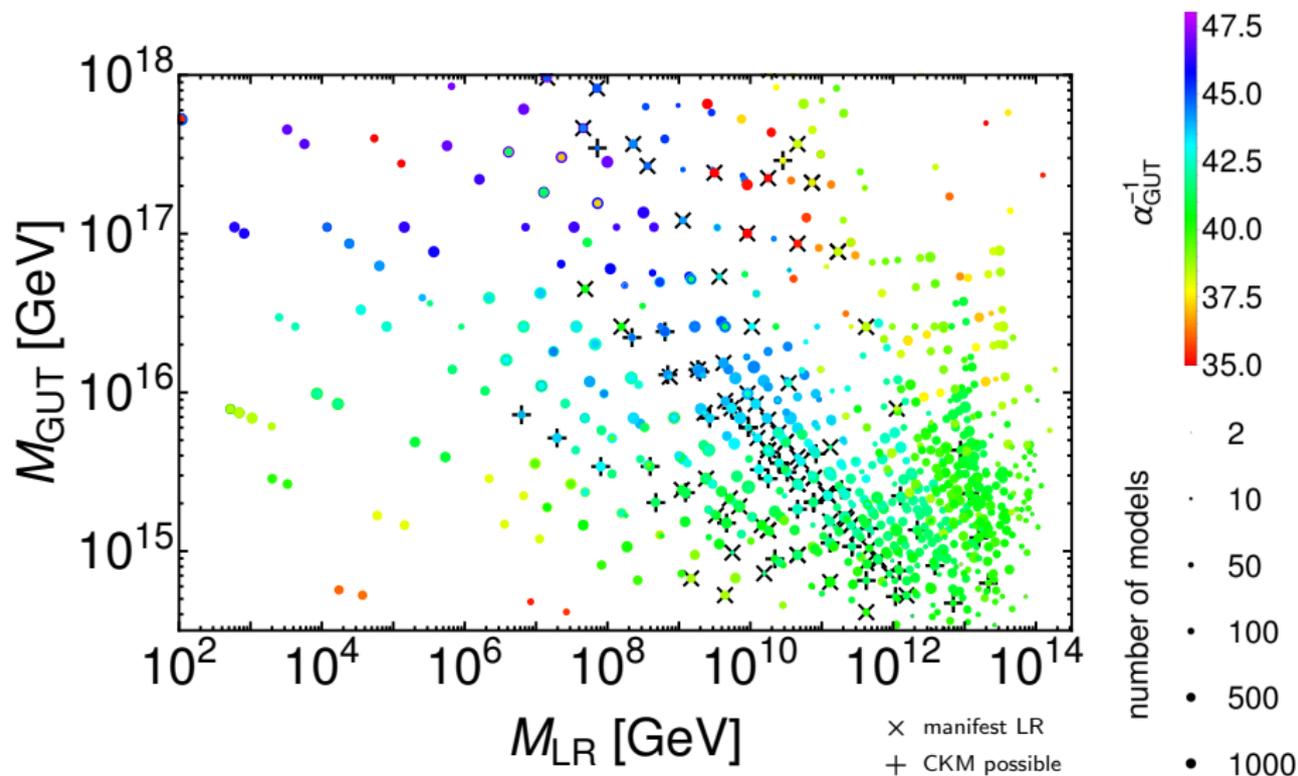


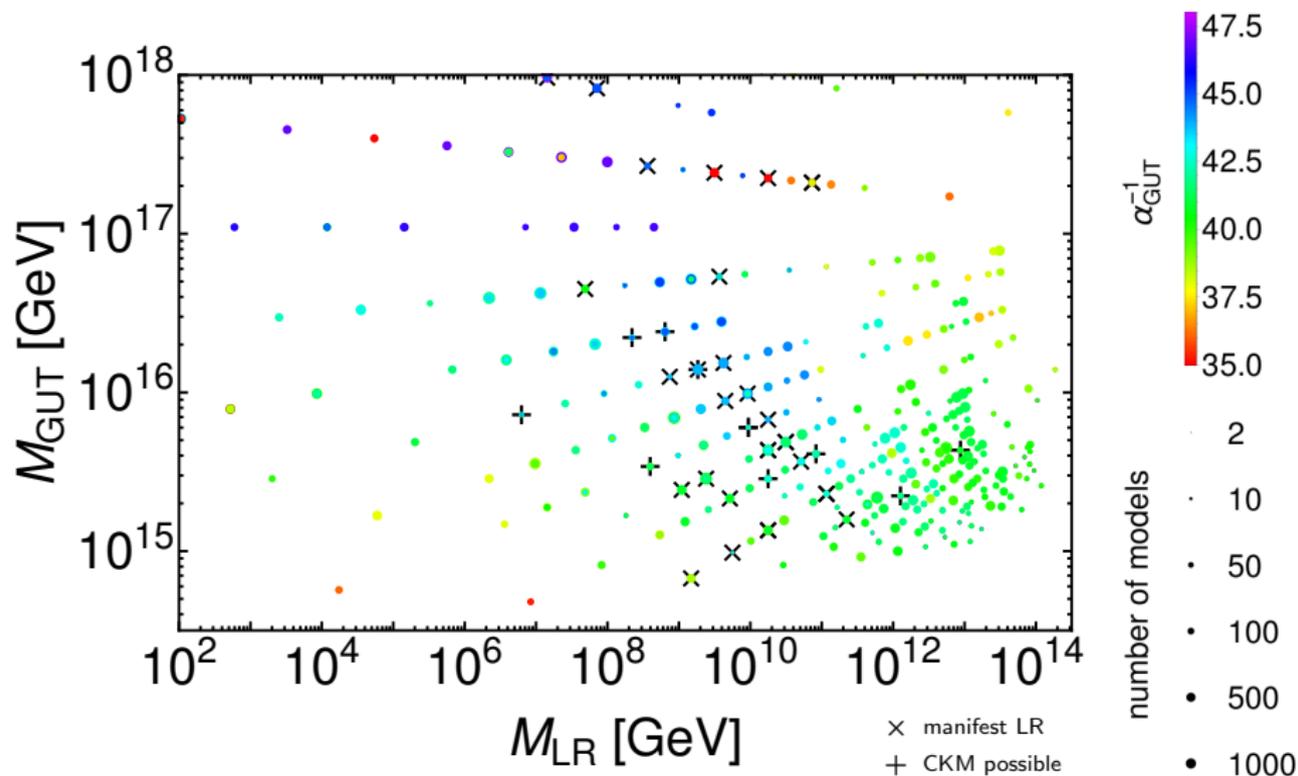
$$\frac{\mathcal{B}(\mu - e)}{\mathcal{B}(\mu - e)_{\text{exp}}} \approx |g_{e\mu}|^2 \left(\frac{g_R}{g_L}\right)^4 \left(\frac{2.1 \times 10^4 \text{ GeV}}{M_{LR}}\right)^4$$

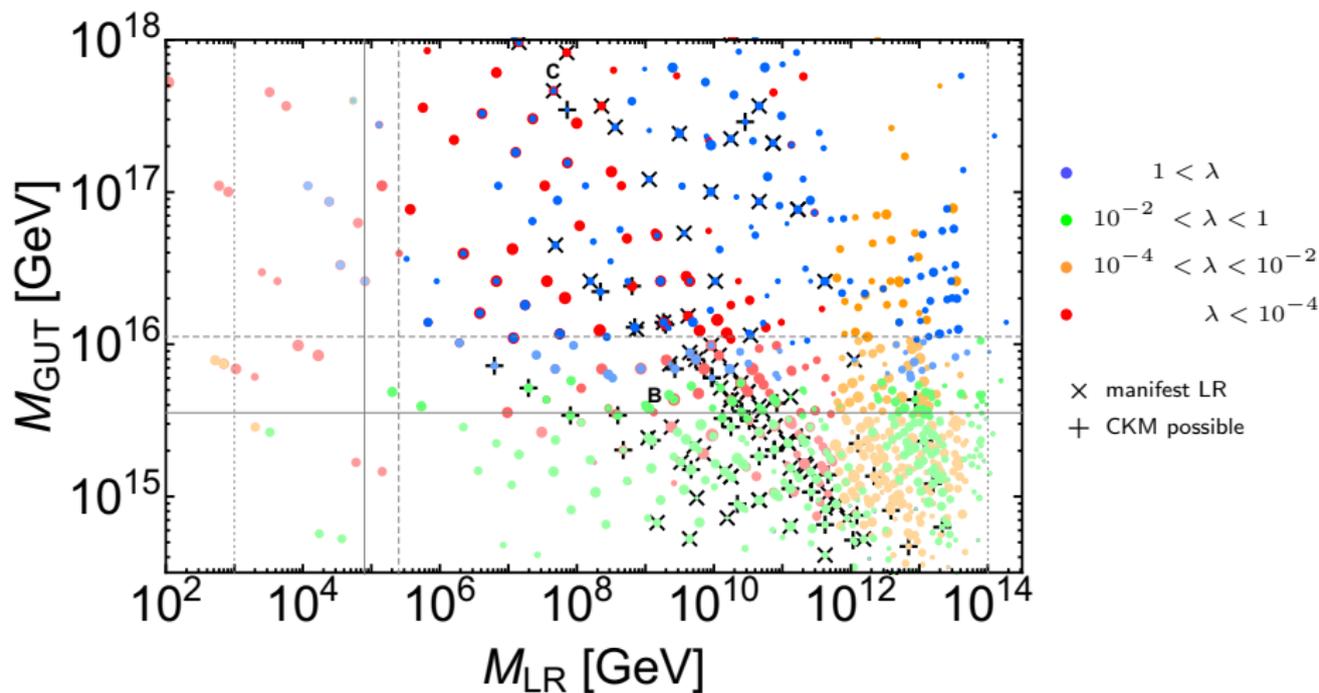
$$\frac{\mathcal{B}(\mu \rightarrow eee)}{\mathcal{B}(\mu \rightarrow eee)_{\text{exp}}} \approx |g_{e\mu}|^2 \left(\frac{9.1 \times 10^4 \text{ GeV}}{M_{LR}}\right)^4$$

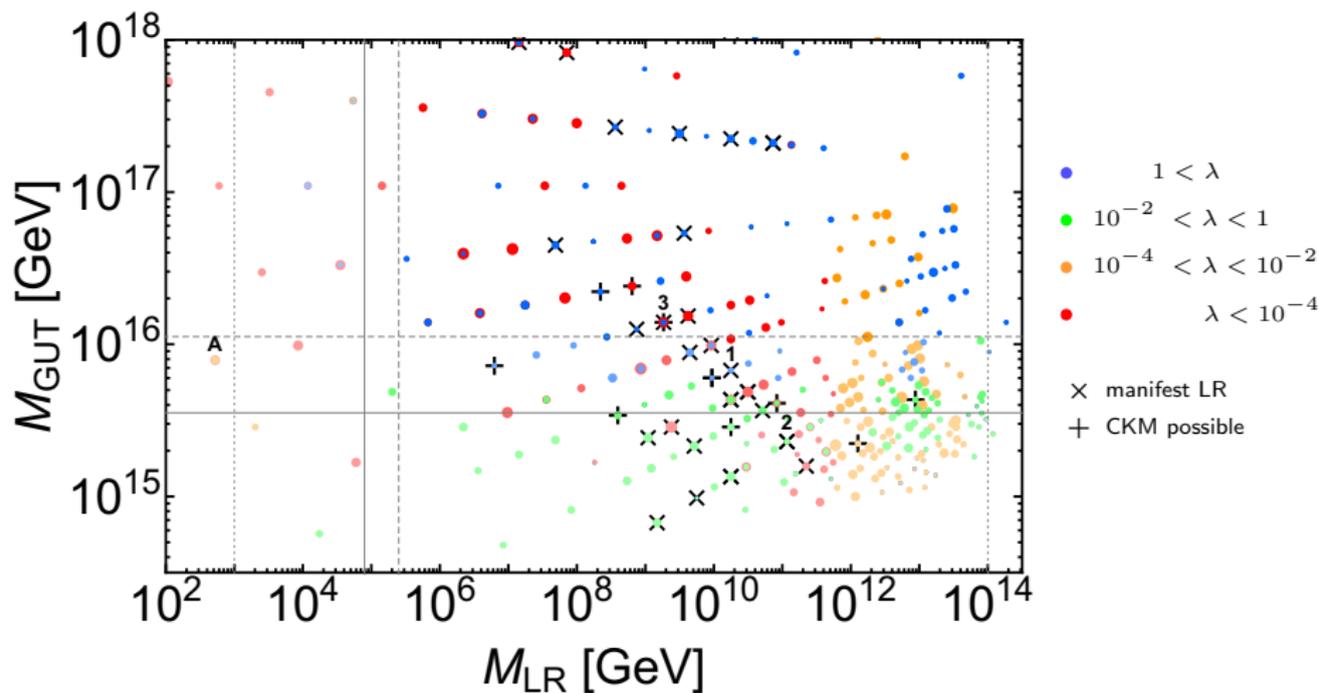
- $\mu \rightarrow e\gamma$  and  $\mu - e$  mediated dominantly by  $W_R$
- dominant mediator for  $\mu \rightarrow eee$  is the triplet Higgs











- a framework generating GUT models in an automated way developed
- aim: scanning through a large model landscapes, considering various symmetry breaking chains and field contents
- procedure applied to  $SO(10)$  models with a single intermediate scale respecting the LR symmetry
- experimental bounds applied to the model landscape restricting phenomenologically incompatible scenarios
- note: a rough scan of large model landscape  $\implies$  2-loop contributions to RGEs, threshold corrections as well as model details neglected
- however, described mechanism is universal  $\rightarrow$  future plans: application to different types of GUT models, more precise calculations

Thank You for attention!

based on [hep-ph/1705.xxxxx](#), to appear within a few days

# Backup Slides

- Abelian breaking of type  $U(1)_A \otimes U(1)_B \rightarrow U(1)_C$  appears in  $SO(10)$
- at scale  $t_{mix}$  the boundary condition for the gauge coupling is:

$$\alpha_C^{-1}(t_{mix}) = r_A^2 \alpha_A^{-1}(t_{mix}) + r_B^2 \alpha_B^{-1}(t_{mix}),$$

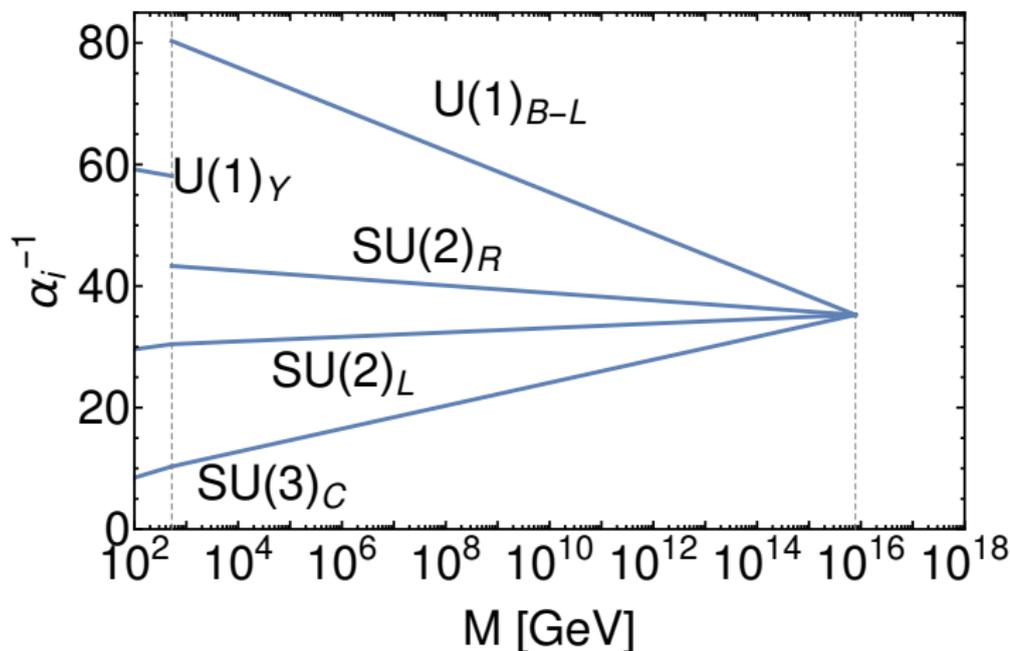
where  $r_{A,B} = Q_{A,B}^v / \sqrt{(Q_A^v)^2 + (Q_B^v)^2}$  such that  $r_A^2 + r_B^2 = 1$

- therefore, the matrix equation takes the form

$$\begin{pmatrix} \alpha_3^{-1} \\ \alpha_2^{-1} \\ \alpha_1^{-1} \end{pmatrix} = (r_A^2 B_A + r_B^2 B_B + B_C) \cdot \Delta t.$$

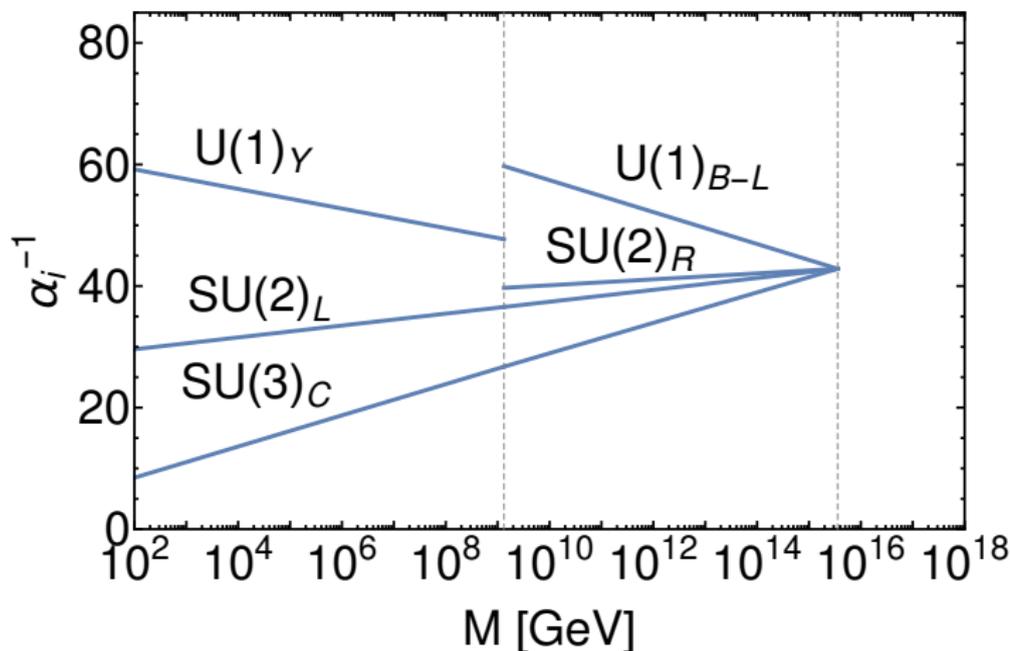
$$\mathcal{R}_{LR}^A \supset \{1, 2, 2, 0\}, \{3, 2, 2, \frac{1}{2}\}, \{\bar{3}, 2, 2, 1\}, \{\bar{3}, 1, 3, -\frac{1}{2}\}, \{1, 1, 3, -\frac{3}{2}\}$$

$$\mathcal{R}_{SM}^A \supset \{1, 2, -\frac{1}{2}\}$$



$$\mathcal{R}_{LR}^B \supset \{1, 2, 2, 0\}, \{3, 2, 2, \frac{1}{2}\}, \{1, 2, 2, 0\}, \{1, 1, 3, -\frac{3}{2}\}, \{1, 1, 3, \frac{3}{2}\}$$

$$\mathcal{R}_{SM}^B \supset \{1, 2, \frac{1}{2}\}, \{1, 2, -\frac{1}{2}\}, \{1, 2, \frac{1}{2}\}, \{1, 2, -\frac{1}{2}\}$$



$$\mathcal{R}_{LR}^C \supset \{1, 2, 2, 0\}, \{8, 1, 1, 0\}, \{\bar{3}, 1, 1, 1\}, \{1, 1, 3, -\frac{3}{2}\} \{1, 3, 1, \frac{3}{2}\}$$

$$\mathcal{R}_{SM}^C \supset \{1, 2, -\frac{1}{2}\}, \{1, 2, \frac{1}{2}\}$$

