The Vev Flip-Flop: Dark Matter Decay between Weak Scale Phase Transitions

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JGU Mainz

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Based on 1608.07578, MJB & Joachim Kopp

- Vev Flip-Flop: The Idea

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- There is a phase in the early universe where this scalar has a vev, so the symmetry is broken and DM can decay

- At a lower temperature the symmetry is restored and DM is then stable

- **Vev Flip-Flop: Model**

Field	Spin	SM	$\mathbb{Z}_3$	mass scale
χ	$\frac{1}{2}$	$(1, 1, 0)$	$\chi \to e^{2\pi i/3} \chi$	TeV
S_3	0	$(1, 3, 0)$	$S_3 \to e^{2\pi i/3} S_3$	100 GeV
Ψ_3	$\frac{1}{2}$	$(1, 3, 0)$	$\Psi_3 \to e^{-2\pi i/3} \Psi_3$	TeV
Ψ'_3	$\frac{1}{2}$	$(1, 3, 0)$	$\Psi'_3 \to e^{-2\pi i/3} \Psi'_3$	TeV

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$$\mathcal{L}_{\text{Yuk}} = y_\chi S_3^\dagger \overline{\chi} \Psi_3 + y_\chi' S_3^\dagger \overline{\chi} \Psi_3'$$

$$+ y_\Psi \epsilon^{ijk} S_3^i \overline{\Psi_3^j} (\Psi_3'^k)^c + h.c.$$

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$$y_\chi \sim y_\chi' \sim 10^{-7} \qquad\qquad y_\Psi \sim 1$$

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$$
\begin{aligned}
V(H,S_3) = &-\mu^2 H^\dagger H + \lambda(H^\dagger H)^2 \\
&- \mu_S^2 S_3^\dagger S_3 + \lambda_S(S_3^\dagger S_3)^2 + \lambda_3(S_3^\dagger T^a S_3)^\dagger S_3^\dagger T^a S_3 \\
&+ \alpha H^\dagger H S_3^\dagger S_3 + \beta H^\dagger \tau^a H S_3^\dagger T^a S_3
\end{aligned}
$$

- ## Vev Flip-Flop: Effective Potential

$$V^{\text{eff}} = V^{\text{tree}} + V^{\text{1-loop}}$$

$$= V^{\text{tree}} + \sum_i \frac{n_i T}{2} \sum_{n=-\infty}^{+\infty} \int \frac{d^3\vec{k}}{(2\pi)^3} \log\left[\vec{k}^2 + \omega_n^2 + m_i^2(h, S)\right]$$

$$V^{\text{eff}} = V^{\text{tree}} + V^{\text{CW}} + V^T + V^{\text{daisy}}$$

- Vev Flip-Flop: Effective Potential

$$V^{\text{eff}} = V^{\text{tree}} + \underline{V^{\text{CW}}} + V^T + V^{\text{daisy}}$$

$$V^{\text{CW}}(h, S) = \sum_i \frac{n_i}{64\pi^2} m_i^4(h, S) \left[\log \frac{m_i^2(h, S)}{\Lambda^2} - \frac{3}{2} \right]$$

$$m_A^2(h, S) = -\mu_S^2 + \frac{1}{2}\alpha h^2 + \lambda_S S^2 \qquad m_Z^2(h, S) = \frac{1}{4}(g^2 + g'^2)h^2$$

$$m_{G^0}^2(h, S) = -\mu^2 + \lambda h^2 + \frac{1}{2}\alpha S^2 \qquad m_\gamma^2(h, S) = 0$$

$$m_{W^\pm}^2(h, S) = \frac{1}{4}g^2(h^2 + 4S^2) \qquad m_t^2(h, S) = \frac{1}{2}y_t^2 h^2$$

$$i \in \{A, G^0, W^\pm, Z, \gamma, t, h - S, S^+ - S^- - G^+\}$$

Coleman & Weinberg, 1973

- **Vev Flip-Flop: Effective Potential**

$$V^{\mathrm{eff}} = V^{\mathrm{tree}} + \underline{V^{\mathrm{CW}}} + V^{T} + V^{\mathrm{daisy}}$$

$$V^{\mathrm{ct}} = -\frac{h^2}{2}\left(\left.\frac{\partial^2 V^{\mathrm{CW}}}{\partial h^2}\right|_{\mathrm{min}} + \frac{3}{2v}\left(\left.\frac{\partial V^{\mathrm{CW}}}{\partial h}\right|_{\mathrm{min}} - v\left.\frac{\partial^2 V^{\mathrm{CW}}}{\partial h^2}\right|_{\mathrm{min}}\right)\right) +$$
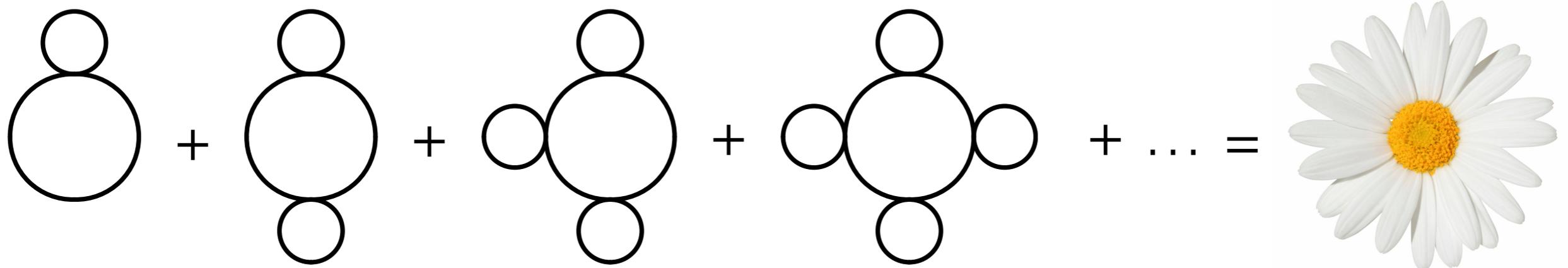$$+ \frac{h^4}{4}\frac{1}{2v^3}\left(\left.\frac{\partial V^{\mathrm{CW}}}{\partial h}\right|_{\mathrm{min}} - v\left.\frac{\partial^2 V^{\mathrm{CW}}}{\partial h^2}\right|_{\mathrm{min}}\right)$$

- **Vev Flip-Flop: Effective Potential**

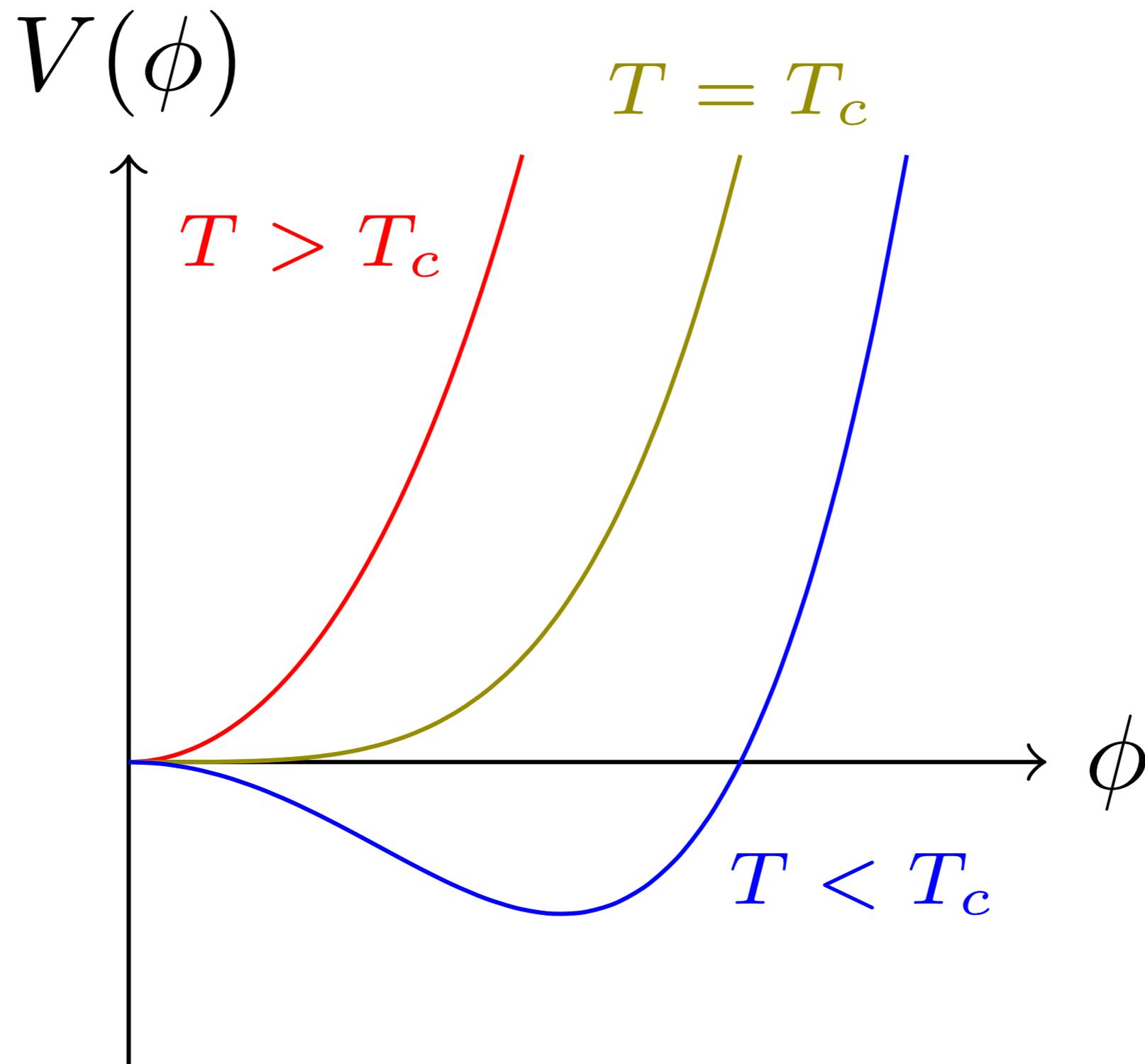$$V^{\text{eff}} = V^{\text{tree}} + V^{\text{CW}} + \underline{V^T} + V^{\text{daisy}}$$

$$V^T(h, S) = \sum_i \frac{n_i T^4}{2\pi^2} \int_0^\infty dx\, x^2 \log\left[1 \pm \exp\left(-\sqrt{x^2 + m_i^2(h, S)/T^2}\right)\right]$$

Dolan & Jackiw, 1974

- Vev Flip-Flop: Effective Potential

$$V^{\mathrm{eff}} = V^{\mathrm{tree}} + V^{\mathrm{CW}} + V^{T} + \underline{V^{\mathrm{daisy}}}$$

 + + + + … =

$$V^{\mathrm{daisy}} = -\frac{T}{12\pi}\sum_i n_i\left(\left[m_i^2(h,S)+\Pi_i(T)\right]^{\frac{3}{2}} - \left[m_i^2(h,S)\right]^{\frac{3}{2}}\right)$$

$$\Pi_{S,A,S_R^+,S_I^+,S_R^-,S_I^-} = \frac{T^2}{6}\left(3g^2 + \alpha + 2\lambda_3 + 4\lambda_S\right)$$

Dolan & Jackiw, 1974

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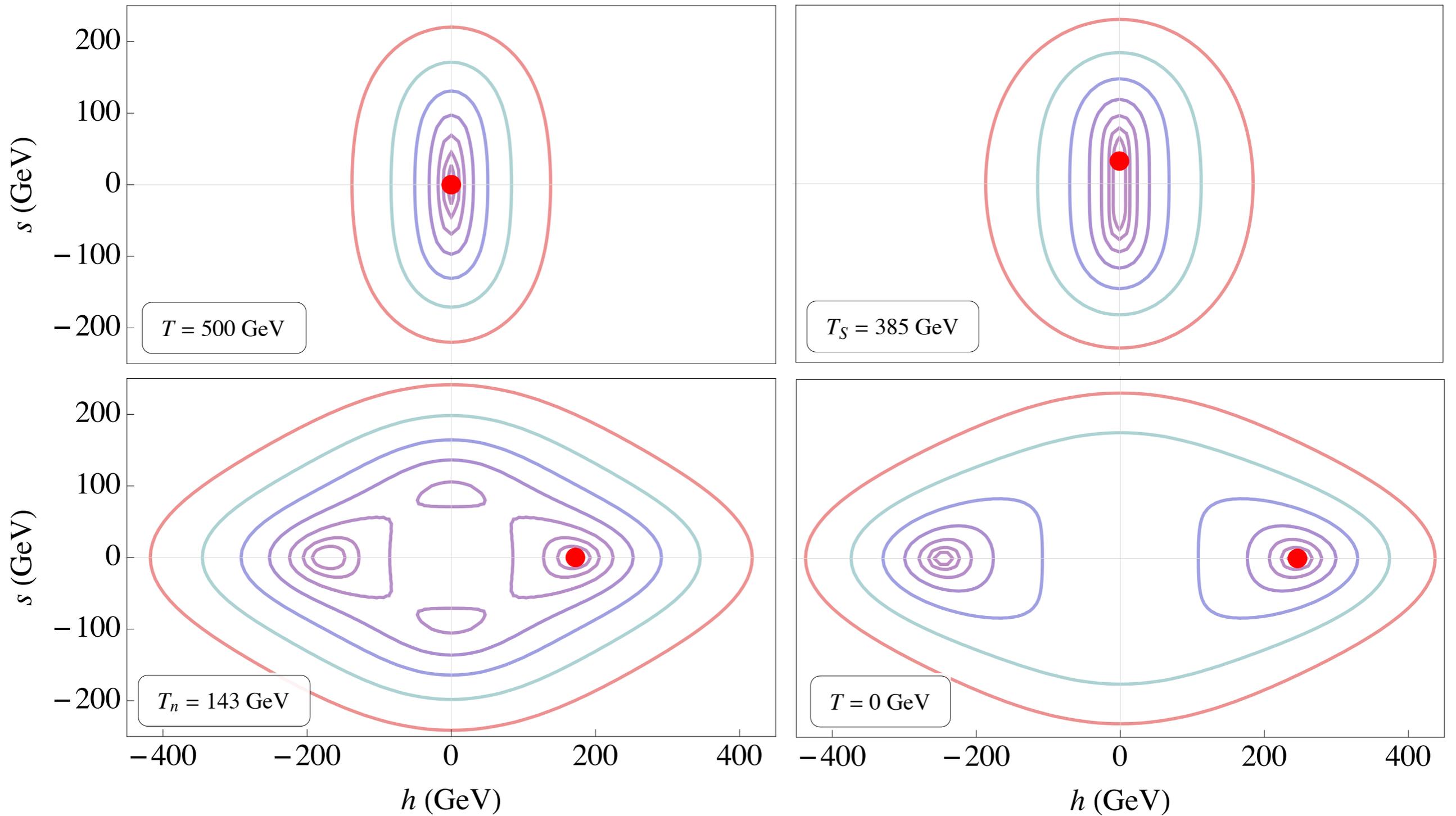
- Vev Flip-Flop: Dark Matter Decay

$$\mathcal{L} \supset y_\chi S_3 \overline{\chi} \Psi_3 + y'_\chi S_3 \overline{\chi} \Psi'_3$$

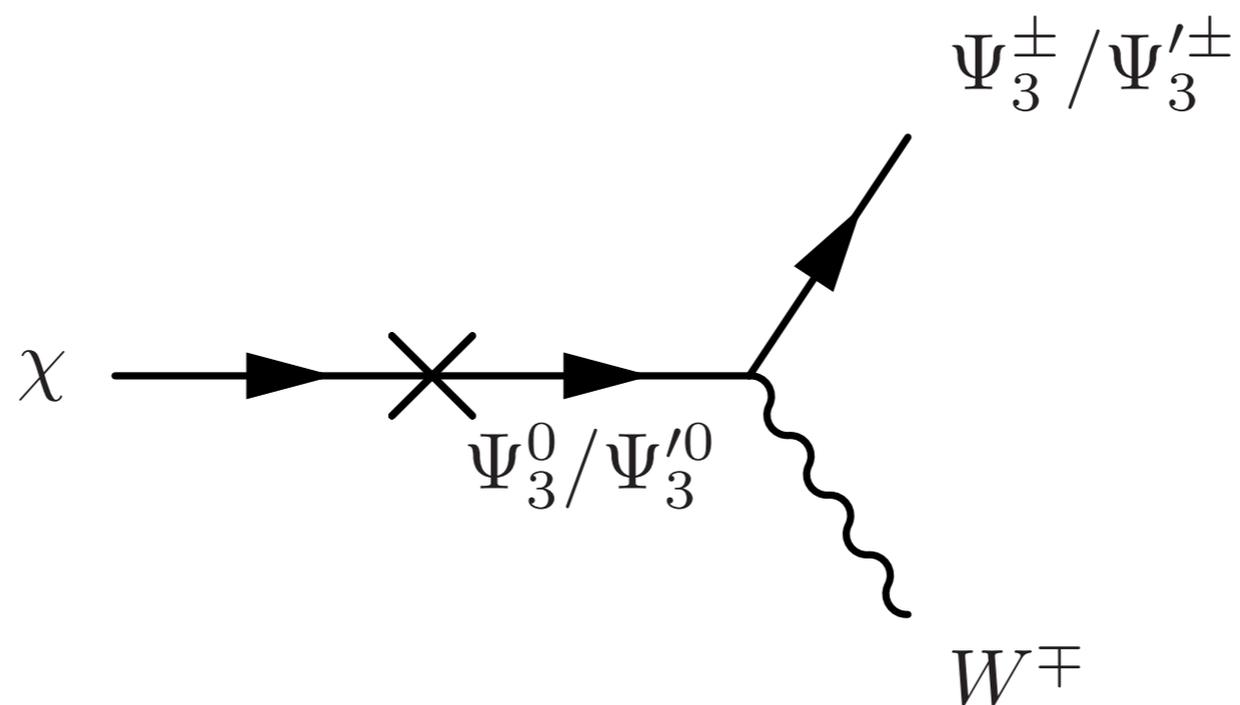- Vev Flip-Flop: Dark Matter Decay

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$$\langle S_3 \rangle \neq 0 \implies \mathcal{L} \supset y_\chi \langle S_3 \rangle \overline{\chi} \Psi_3 + y'_\chi \langle S_3 \rangle \overline{\chi} \Psi'_3$$

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- ## Vev Flip-Flop: Dark Matter Decay

Coupled Boltzmann equations

$$\dot{n}_\chi^j + 3H n_\chi^j = -\frac{\Gamma}{\gamma^j}(n_\chi^j - n_\chi^{j,\mathrm{eq}})$$

$$H^2 = \left(\frac{\dot{a}}{a}\right)^2 = \frac{8\pi G_N}{3}(\rho_{\mathrm{SM}} + \rho_\chi)$$

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In practice we substitute

$$Y^j \equiv \frac{n^j}{s}, \qquad s = g_* \frac{2\pi^2}{45} T^3$$

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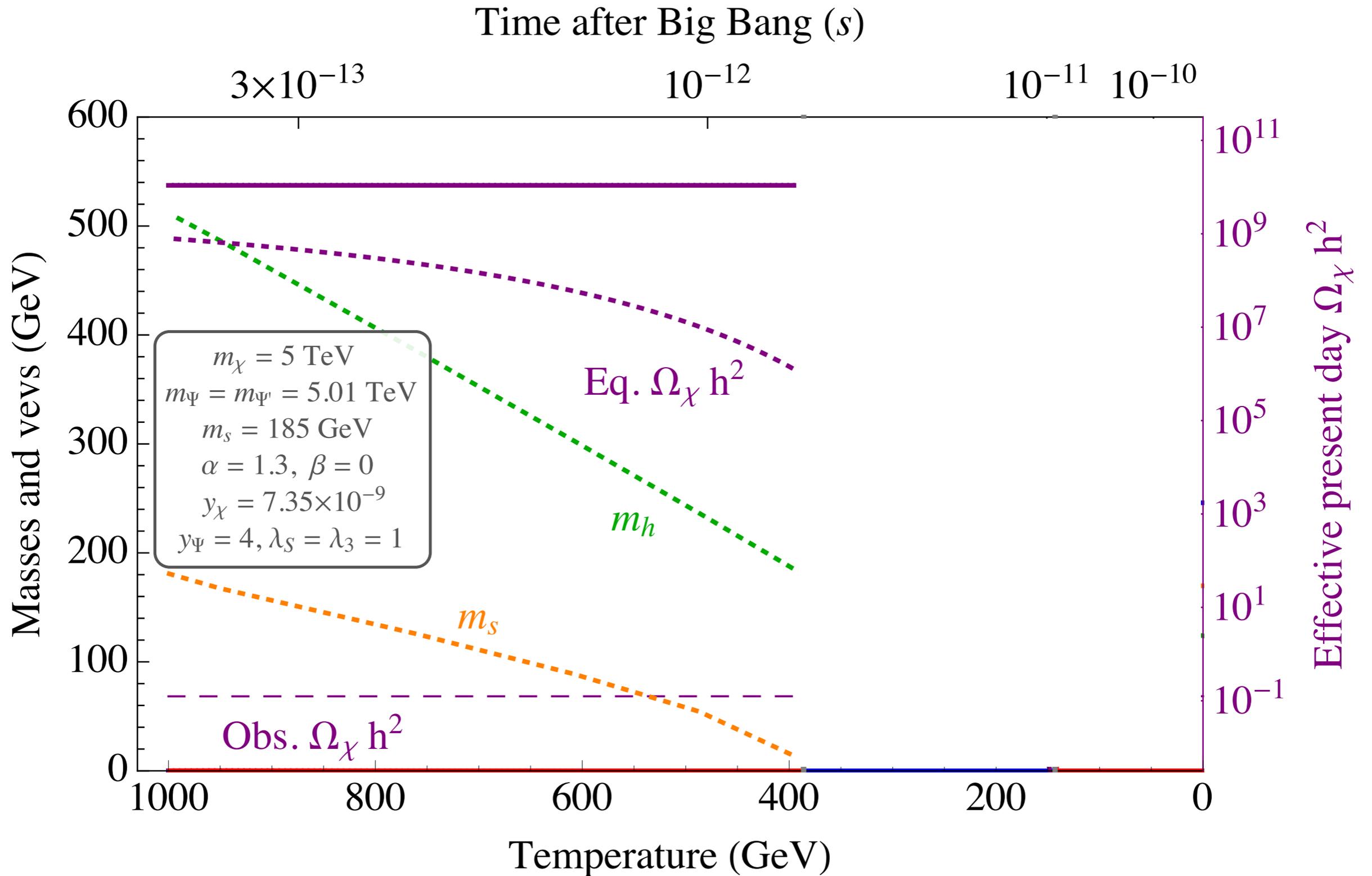$$H^2 = \left(\frac{\dot{a}}{a}\right)^2 = \frac{8\pi G_N}{3}(\rho_{\mathrm{SM}} + \rho_\chi)$$

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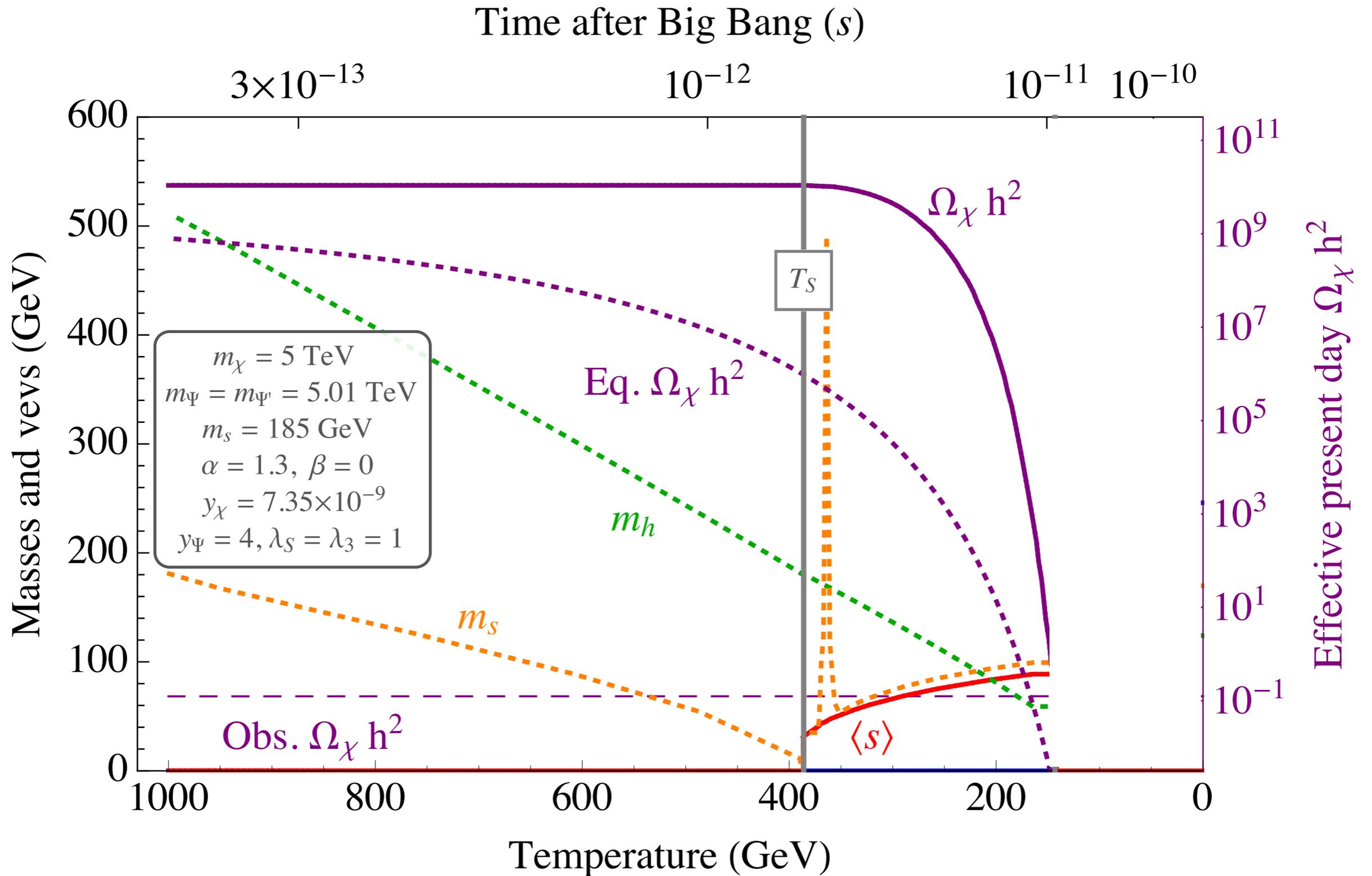$$Y^j \equiv \frac{n^j}{s}, \qquad s = g_* \frac{2\pi^2}{45} T^3$$

And find $\dot{T}$

$$\dot{\rho}_{\mathrm{SM}} = 4\rho_{\mathrm{SM}}\frac{\dot{T}}{T} = m_\chi \sum_j \frac{\Gamma}{\gamma^j}(n^j_\chi - n^{j,\mathrm{eq}}_\chi) - 4H\rho_{\mathrm{SM}}$$

$$\frac{dT}{dt} = \sum_j (n^j_\chi - n^{j,\mathrm{eq}}_\chi)\frac{m_\chi \Gamma T}{4\gamma^j \rho_{\mathrm{SM}}} - HT$$
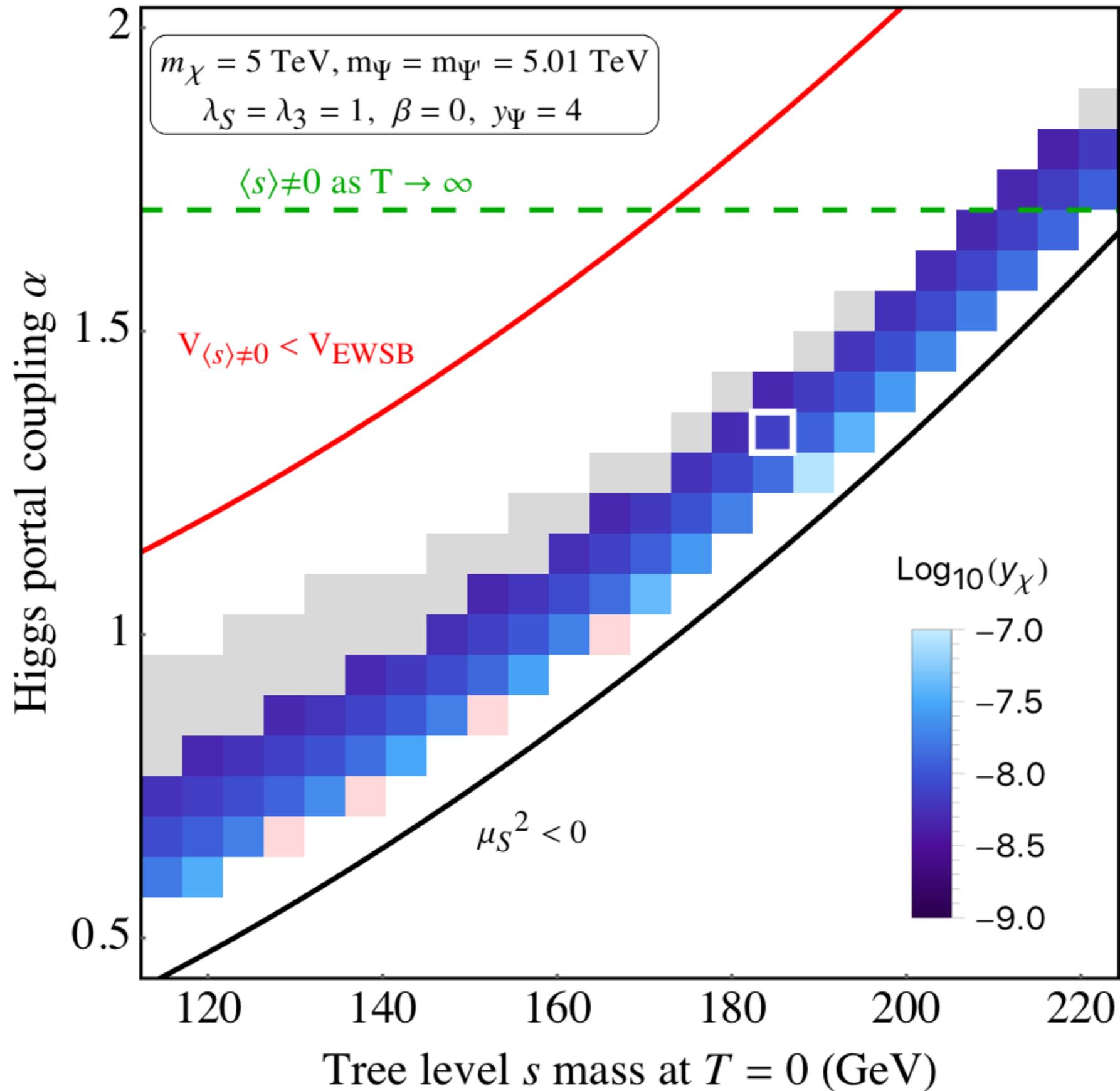
- ## Vev Flip-Flop: Dark Matter Decay

Time after Big Bang (s)

3×10^{-13} 10^{-12} 10^{-11} 10^{-10}

Eq. Ω_χ h^2

$m_\chi = 5$ TeV
$m_\Psi = m_{\Psi'} = 5.01$ TeV
$m_s = 185$ GeV
$\alpha = 1.3$, $\beta = 0$
$y_\chi = 7.35 \times 10^{-9}$
$y_\Psi = 4$, $\lambda_S = \lambda_3 = 1$

m_h

m_s

Obs. Ω_χ h^2

- # Vev Flip-Flop: Dark Matter Decay
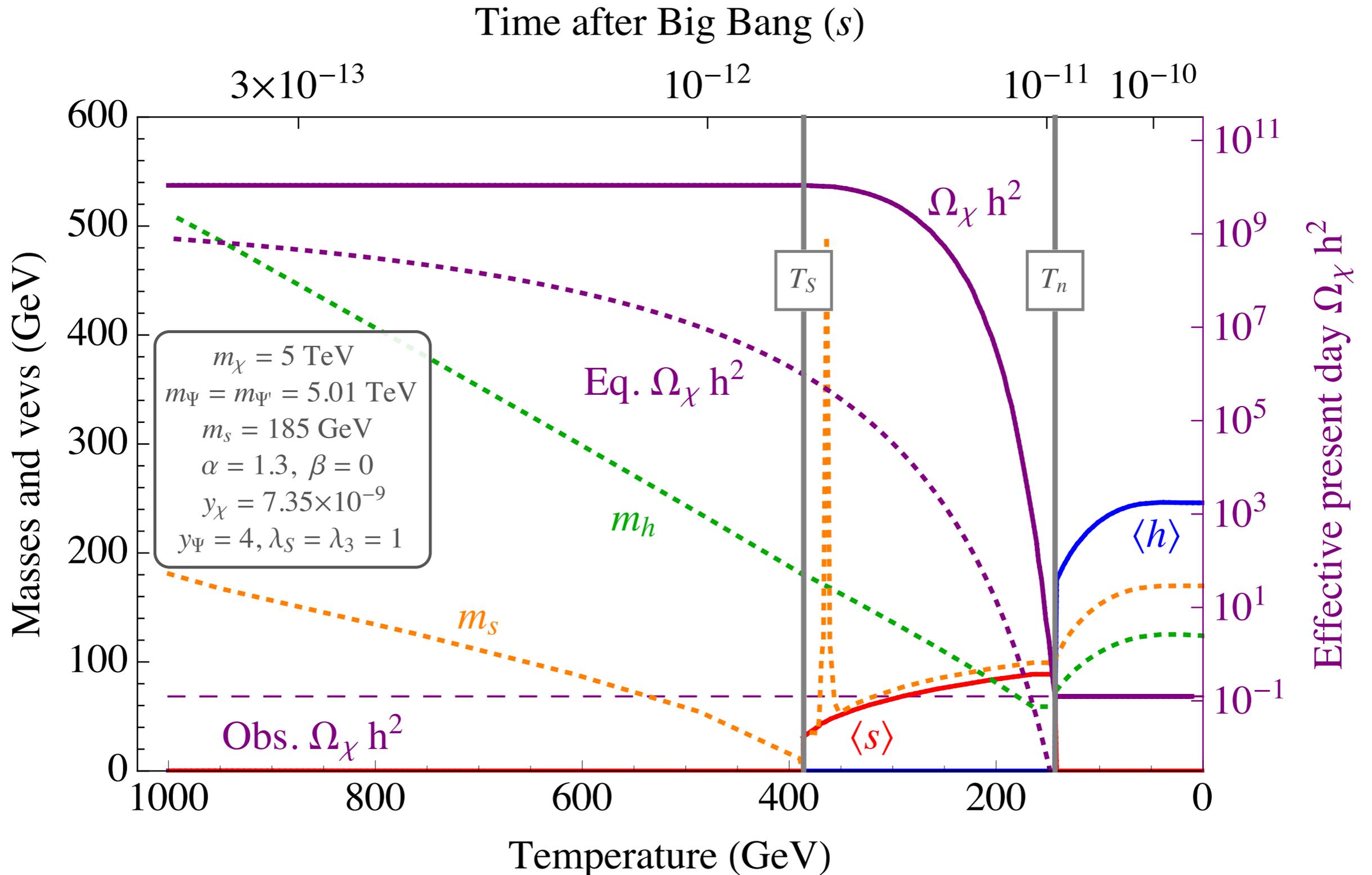
- # Vev Flip-Flop: Parameter Space

- Vev Flip-Flop: Possible tests

Direct and Indirect Detection

Detection of χ hindered by small couplings

Best prospect from subdominant population of $s, \Psi_3^0, \Psi_3'^0$

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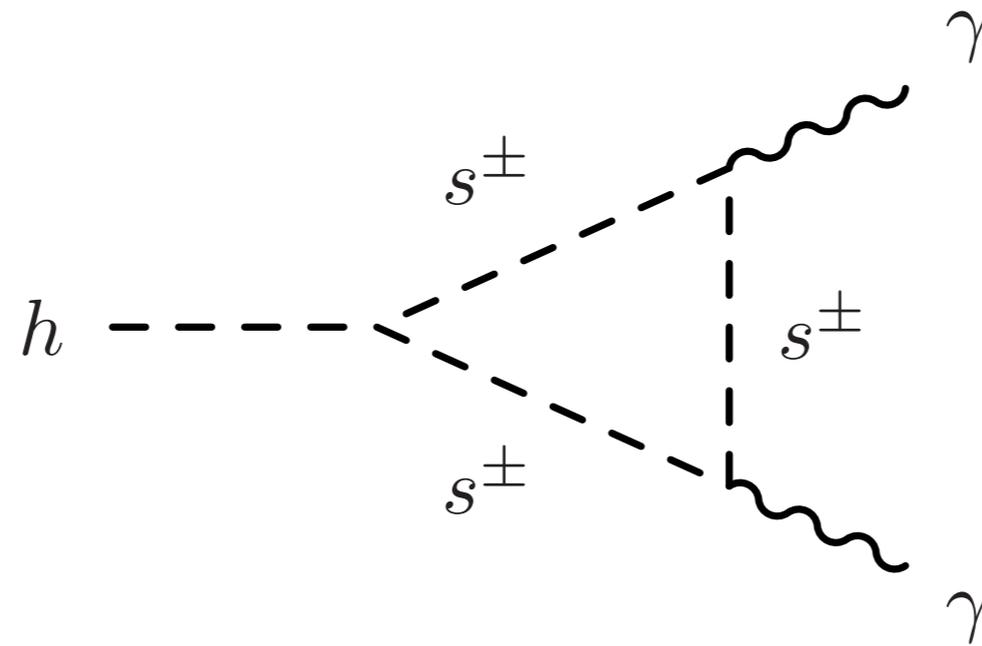Best prospect from subdominant population of $s, \Psi_3^0, \Psi_3'^0$

Drell-Yan production of $s^\pm$ pairs, followed by $s^\pm \to s + W^{*\pm}$

Small mass splitting implies soft decay products

Mono-X is several orders of magnitude too weak
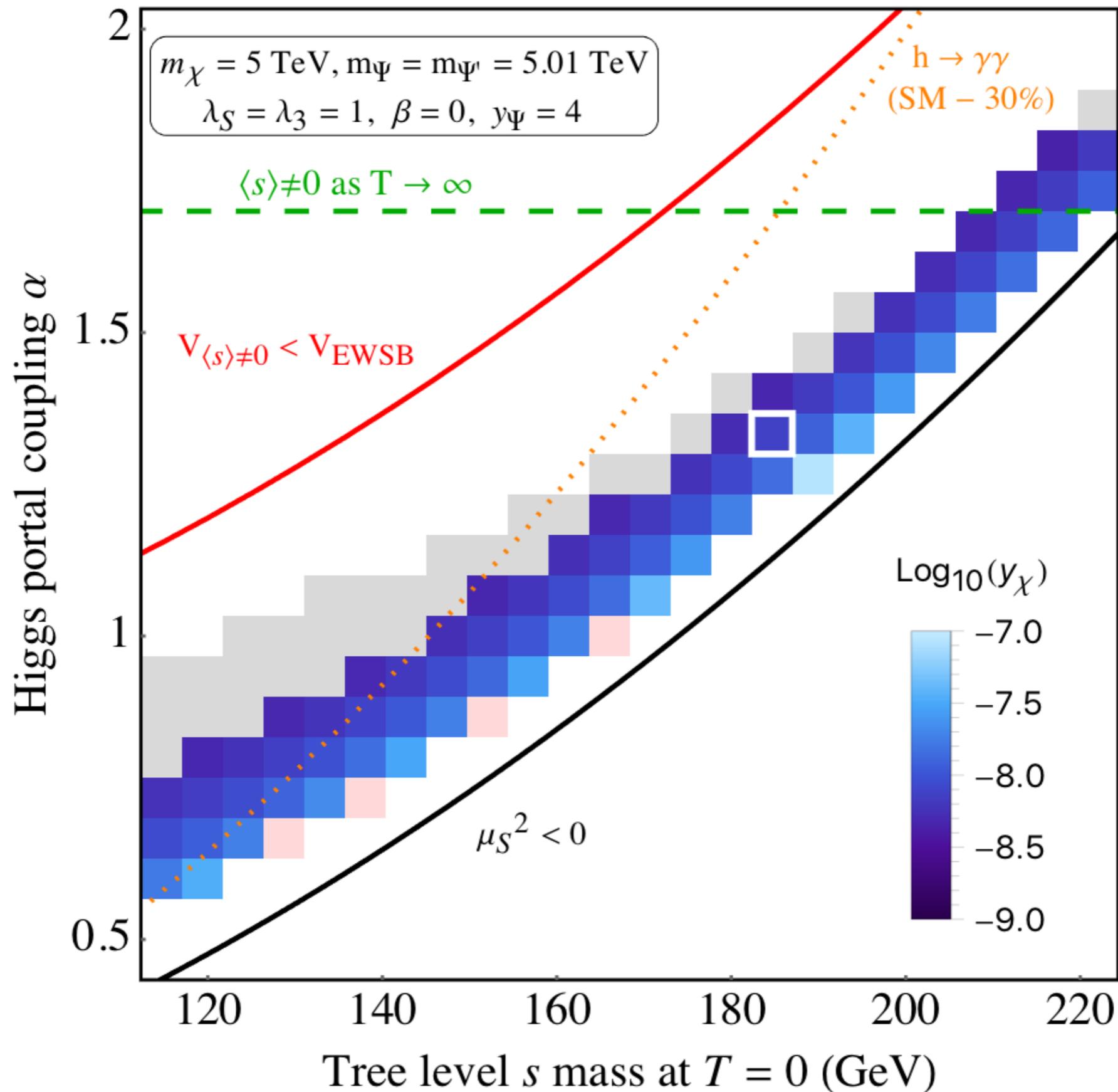
Possibilities in disappearing charged tracks

- Vev Flip-Flop: Possible tests

$$\Gamma(h \to \gamma\gamma) = \frac{v^2}{16\pi m_h} |F^{\gamma\gamma}|^2$$

$$|F^{\gamma\gamma}|^2 = (|F_W^{\gamma\gamma}| - |F_f^{\gamma\gamma}| - |F_{S_3}^{\gamma\gamma}|)^2$$

Vev Flip-Flop: Parameter Space

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- Presented a successful toy model

- Which works in a large parameter space

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- Outlook:
  - More realisations,
  - LHC phenomenology,
  - gravitational wave signals,
  - baryogenesis,…