

# ***Dynamical Dark Matter from Strongly-Coupled Dark Sectors***

arXiv:1610.04112

*Fei Huang*

***University of Arizona***

**in collaboration with:**

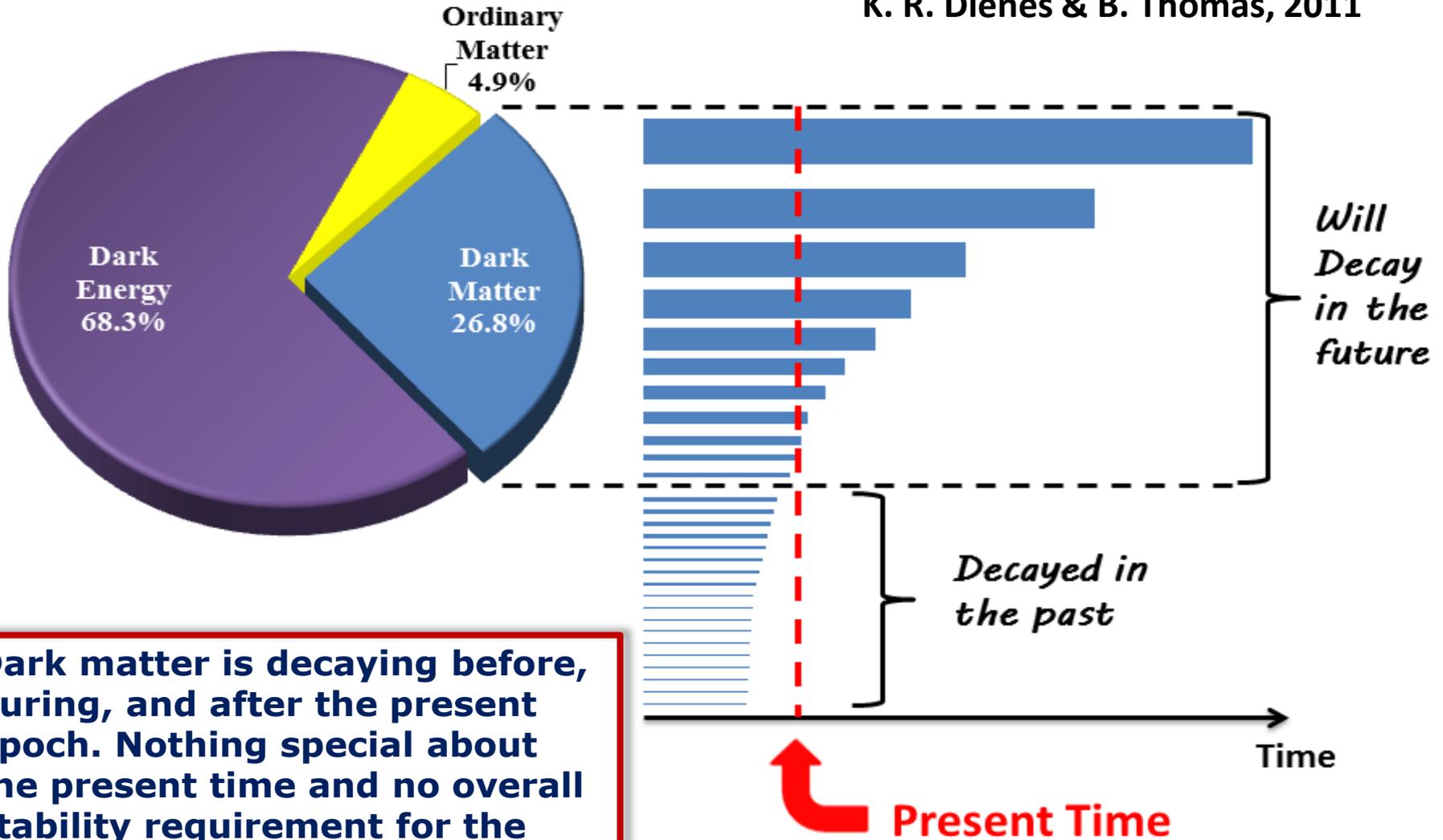
*Keith Dienes, Shufang Su, Brooks Thomas*



**Pheno 2017**

# Basic Picture of Dynamical Dark Matter (DDM)

K. R. Dienes & B. Thomas, 2011



**Dark matter is decaying before, during, and after the present epoch. Nothing special about the present time and no overall stability requirement for the dark sector.**

**DDM requires an *ENSEMBLE* of states whose lifetimes scales with their abundances. Is there any theory that can naturally provide us with such DDM ensembles?**

***Yes!***

**Previous works on DDM studied ensembles from KK towers and representations of large-hidden gauge groups.**

**These realizations of DDM all have polynomial scaling relations between abundances, lifetimes and masses.**

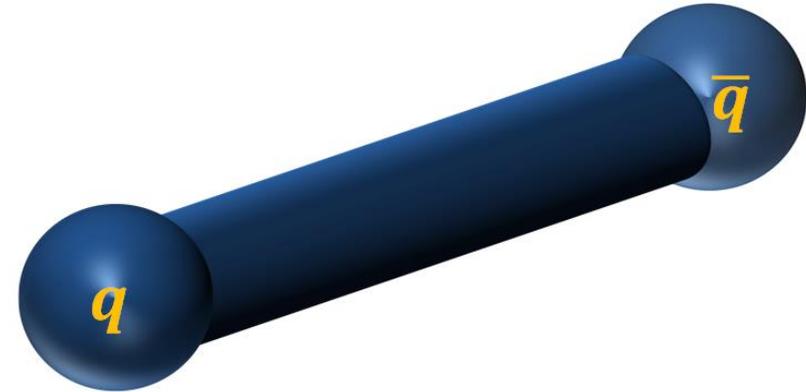
**Are there other possibilities?**

***Yes!***

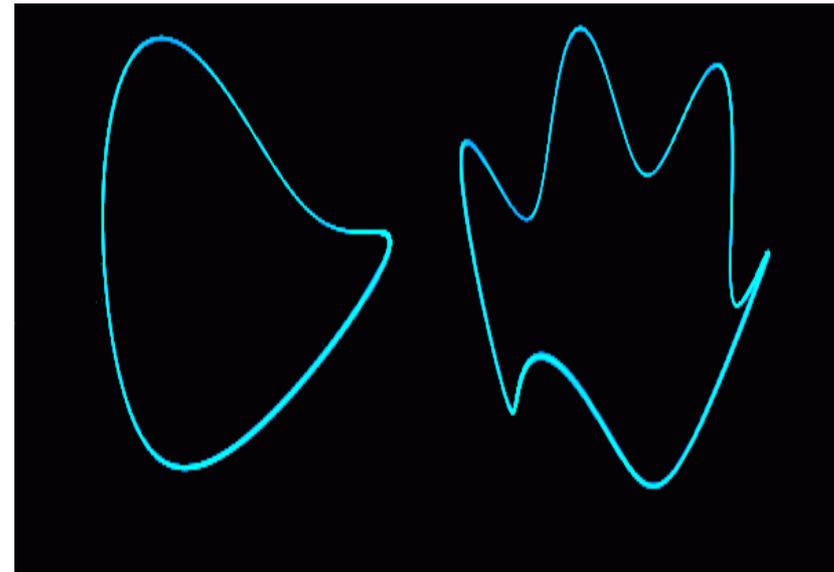
**In fact, we explored something quite different...**

## Strongly-Coupled DDM Ensembles

- Fermions (dark quarks) attached on the ends of a flux tube, charged under a non-Abelian gauge group  $G$ .
- In the confining phase below  $T_c$ , physical d.o.f are composite states (dark “hadrons”).



- Bulk states in Type I string theories.
- Typically neutral with respect to all brane gauge symmetries
- Interact with brane states only gravitationally.
- For brane-localized observers, these states are dark matter.



## Common features:

- Mass spectrum follow linear **Regge** trajectories:

$$M_n^2 = nM_s^2 + M_0^2$$

$M_s \equiv 1/\sqrt{\alpha'}$ : "String scale"

$M_0$ : Lightest mass in the spectrum

- Exponentially growing degeneracy of states (**Hagedorn** behavior)

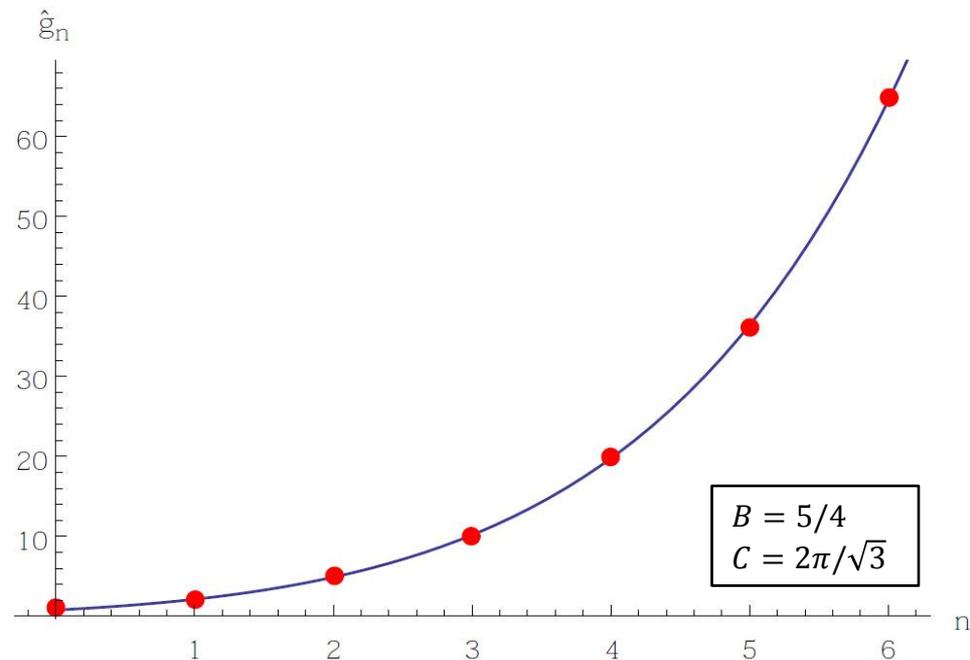
$$g_n \propto n^{-B} e^{C\sqrt{n}} \text{ as } n \rightarrow \infty$$

$$\begin{cases} B = (3 + D_\perp)/4 \\ C = \pi\sqrt{2c/3} \end{cases}$$

$D_\perp$ : Number of transverse dimensions into which string/flux tube can fluctuate

$c \equiv c_{int} + D_\perp$ : Central charge

$M_s$ : string scale  
 $M_0$ : lightest mass  
 $B \sim D_\perp$   
 $C \sim$  central charge  
 $r \equiv M_0/M_s$



$M_s$ : string scale  
 $M_0$ : lightest mass  
 $B \sim D_\perp$   
 $C \sim$  central charge  
 $r \equiv M_0/M_s$

# Common features:

- Mass spectrum follow linear **Regge** trajectories:

$$M_n^2 = nM_s^2 + M_0^2$$

$M_s \equiv 1/\sqrt{\alpha'}$ : "String scale"

$M_0$ : Lightest mass in the spectrum

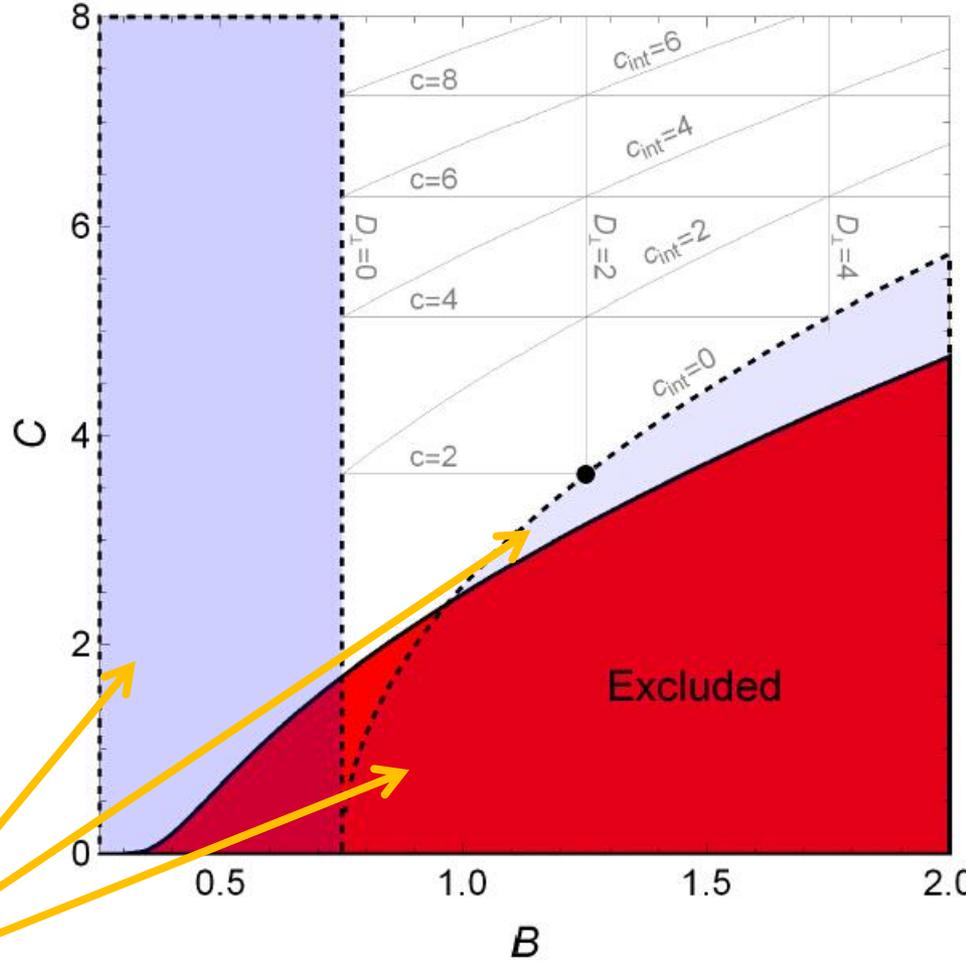
- Exponentially growing degeneracy of states (**Hagedorn** behavior)

$$g_n \propto n^{-B} e^{C\sqrt{n}} \text{ as } n \rightarrow \infty$$

$$\begin{cases} B = (3 + D_\perp)/4 \\ C = \pi\sqrt{2c/3} \end{cases}$$

$D_\perp$ : Number of transverse dimensions into which string/flux tube can fluctuate

$c \equiv c_{int} + D_\perp$ : Central charge



Excluded by  $\begin{cases} D_\perp > 0 \\ c_{int} \geq 0 \\ g_1 > 1 \end{cases}$

*Just to be self-consistent!*

# Cosmological Abundances

In DDM, each component has an abundance:  $\Omega_n(t)$

Total dark matter abundance:  $\Omega_{tot}(t) = \sum_n g_n \Omega_n(t)$

Finite!

Exponentially growing!

$M_s$ : string scale
$M_0$ : lightest mass
$T_c$ : critical temp
$B \sim D_\perp$
$C \sim$ central charge
$r \equiv M_0/M_s$
$s \equiv T_c/M_s$

$\Omega_n$  has to fall sufficiently fast!

Assuming initial abundance distribution follows Boltzmann distribution:

$$\Omega_n(t_c) \sim e^{-M_n/T_c}$$

Boltzmann suppression!

To have a finite total abundance:

$$\begin{cases} \Omega_n(t_c) \sim e^{-M_n/T_c} \\ g_n \sim n^{-B} e^{C\sqrt{n}} \end{cases} \Rightarrow \frac{M_s}{T_c} \equiv \frac{1}{s} \geq C$$

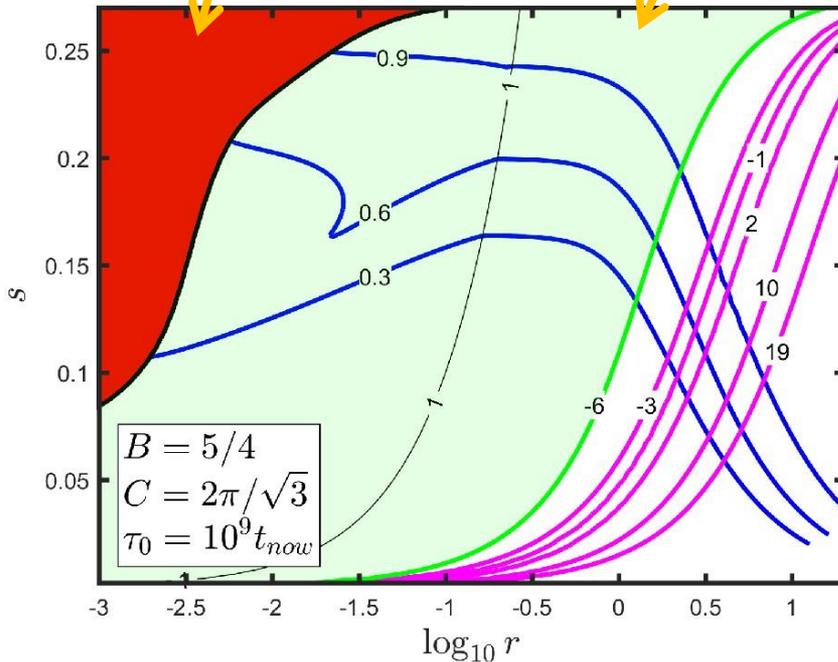
Confinement temperature must be below or at Hagedorn temperature  
 $T_H \equiv M_s/C$

# Parameter Space

$M_s$ : string scale  
 $M_0$ : lightest mass  
 $T_c$ : critical temp  
 $B \sim D_\perp$   
 $C \sim$  central charge  
 $r \equiv M_0/M_s$   
 $s \equiv T_c/M_s$   
 $\tau_0$ : longest lifetime

Excluded by  $t_{LB}$  and EoS constraints

Excluded by BBN, structure formation constraints:  $M_0 > \mathcal{O}(\text{keV})$



**Blue lines:** tower fraction,  $\eta \equiv 1 - \frac{\max\{g_n \Omega_n(t)\}}{\Omega_{tot}(t)}$

**Magenta lines:** string scale,  $\log_{10} M_s/\text{GeV}$

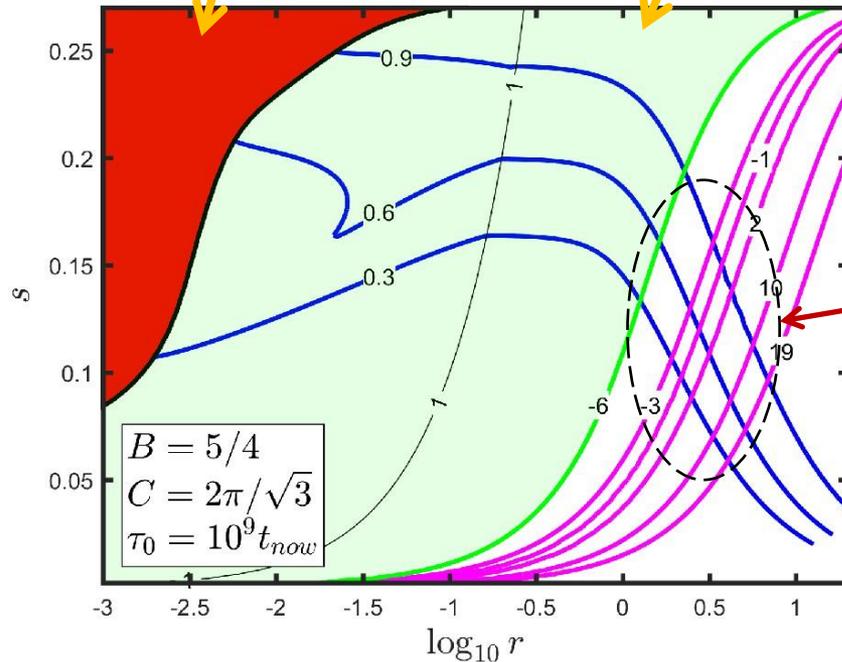
**Green line:** lightest mass,  $\log_{10} M_0/\text{GeV}$

# Parameter Space

$M_s$ : string scale  
 $M_0$ : lightest mass  
 $T_c$ : critical temp  
 $B \sim D_\perp$   
 $C \sim$  central charge  
 $r \equiv M_0/M_s$   
 $s \equiv T_c/M_s$   
 $\tau_0$ : longest lifetime

Excluded by  $t_{LB}$  and EoS constraints

Excluded by BBN, structure formation constraints:  $M_0 > \mathcal{O}(\text{keV})$



Region of interest for DDM:

- $\eta(t_{now})$  varies from 0 to 1
- $M_s$  ranges from keV to Planck scale

**Blue lines:** tower fraction,  $\eta \equiv 1 - \frac{\max\{g_n \Omega_n(t)\}}{\Omega_{tot}(t)}$

**Magenta lines:** string scale,  $\log_{10} M_s/\text{GeV}$

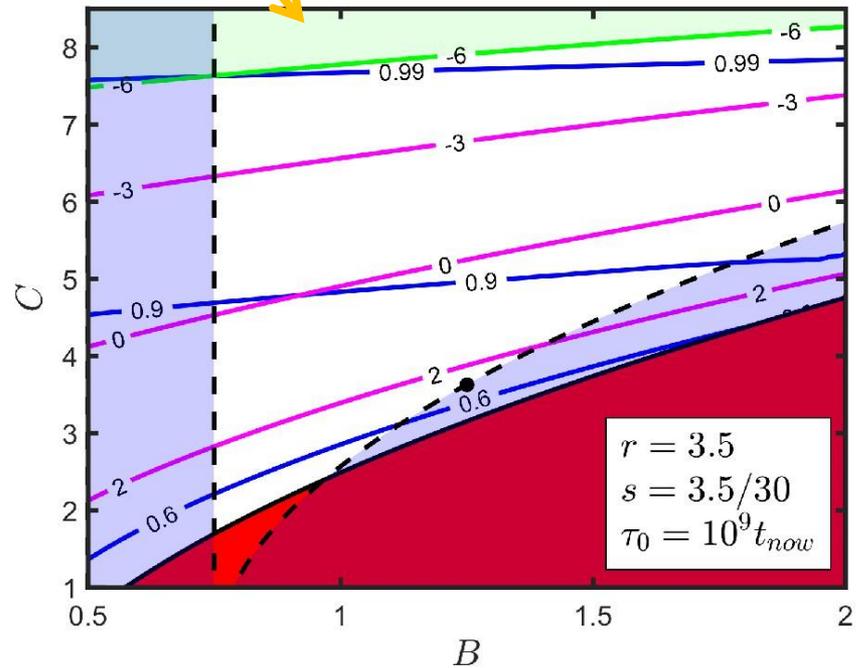
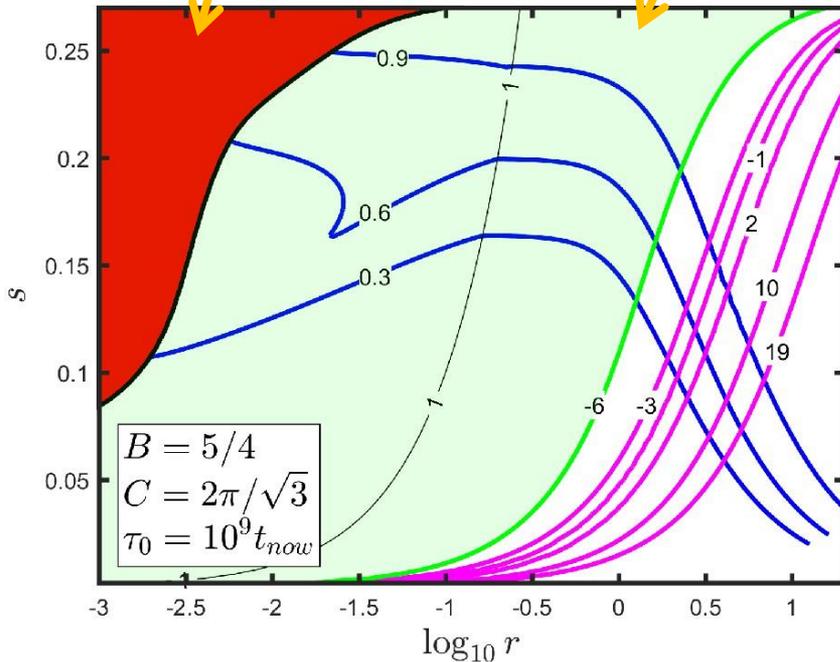
**Green line:** lightest mass,  $\log_{10} M_0/\text{GeV}$

# Parameter Space

$M_s$ : string scale  
 $M_0$ : lightest mass  
 $T_c$ : critical temp  
 $B \sim D_\perp$   
 $C \sim$  central charge  
 $r \equiv M_0/M_s$   
 $s \equiv T_c/M_s$   
 $\tau_0$ : longest lifetime

Excluded by  $t_{LB}$  and EoS constraints

Excluded by BBN, structure formation constraints:  $M_0 > \mathcal{O}(\text{keV})$



**Blue lines:** tower fraction,  $\eta \equiv 1 - \frac{\max\{g_n \Omega_n(t)\}}{\Omega_{tot}(t)}$

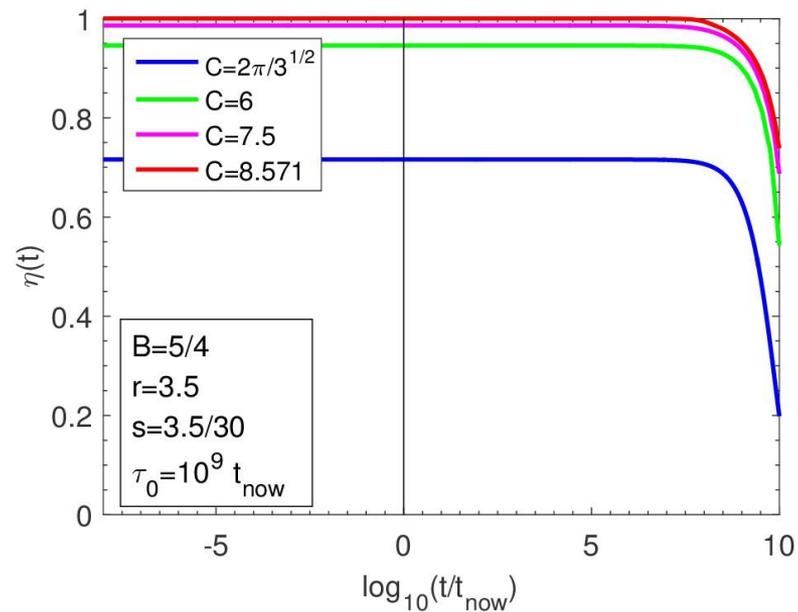
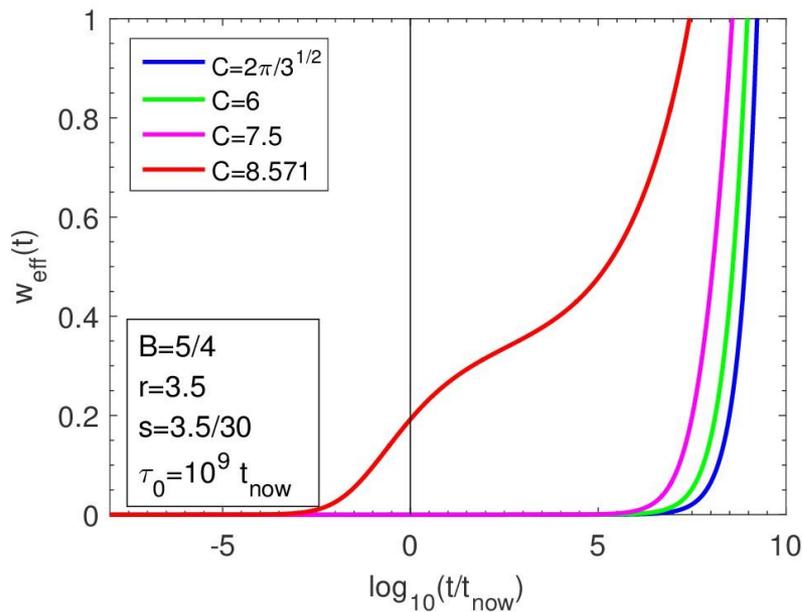
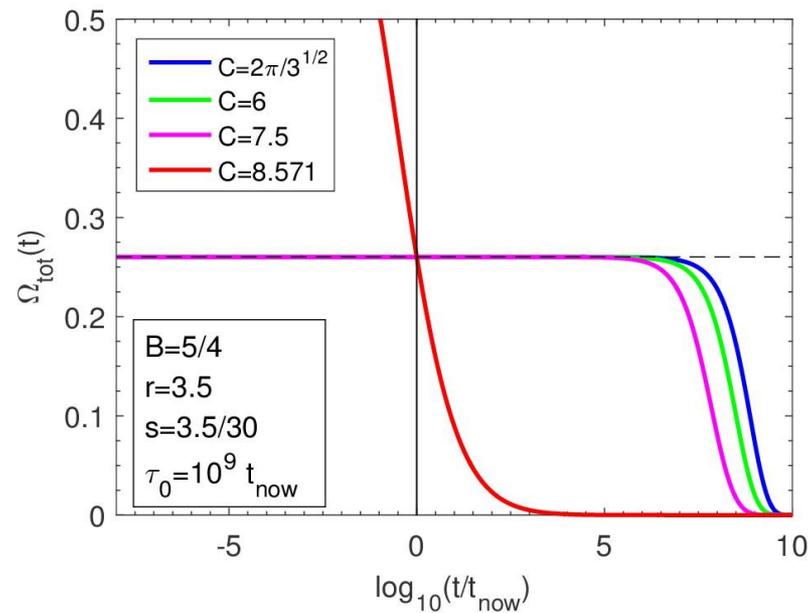
**Magenta lines:** string scale,  $\log_{10} M_s/\text{GeV}$

**Green line:** lightest mass,  $\log_{10} M_0/\text{GeV}$

Large portion of parameter space that satisfies internal consistency requirements is also viable under the phenomenological constraints!

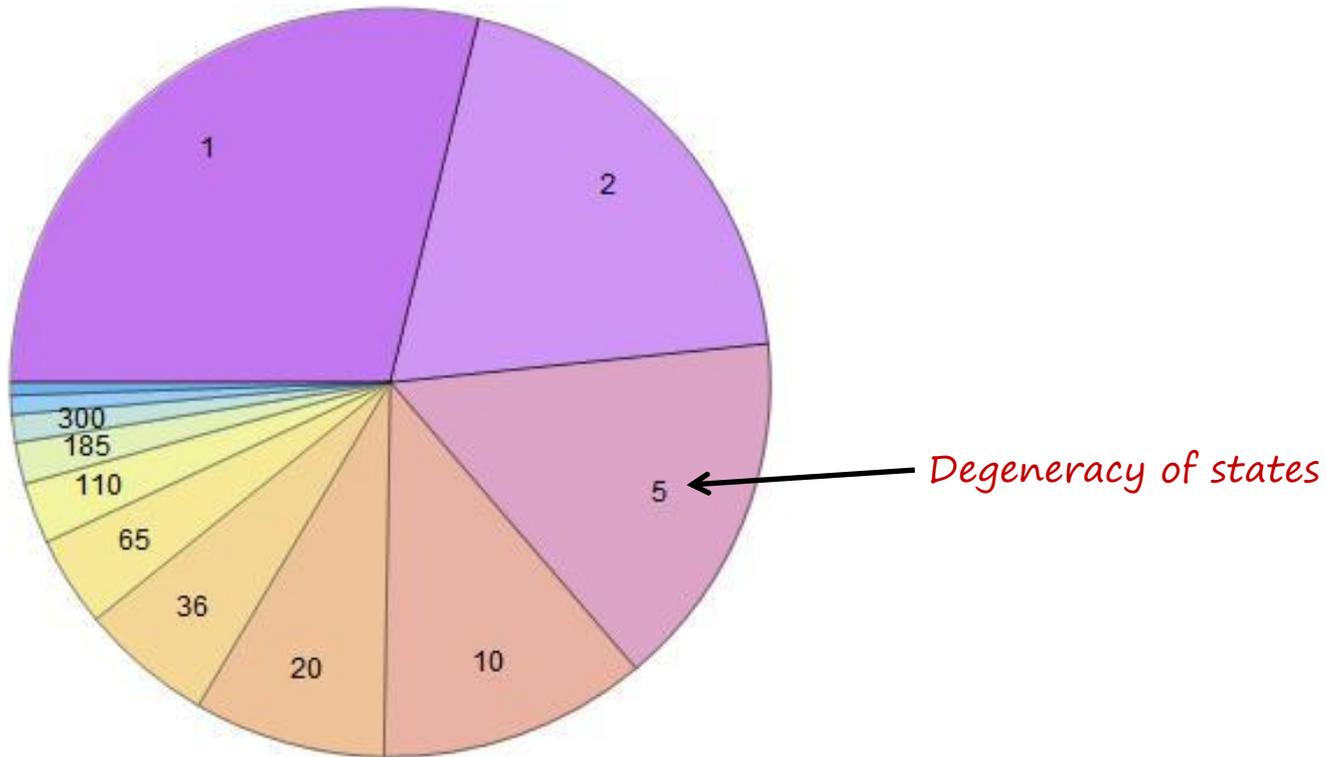
# Time-Evolution of $\Omega_{\text{tot}}(t)$ , $w_{\text{eff}}(t)$ and $\eta(t)$

- $\Omega_{\text{tot}}(t_{\text{now}}) = 0.26$ , all states eventually decay
- Larger  $C \rightarrow$  higher degeneracy  $\rightarrow$  more shorter-lived excited states  $\rightarrow$  more decays
- Evolution well behaved ( $\Omega_{\text{tot}}$  changes little when going back in time,  $w_{\text{eff}}(t_{\text{now}})$  close to zero) if Boltzmann suppression is stronger than Hagedorn behavior. **Red** curve shows the opposite situation.
- $\eta$  significantly above zero, really **DDM-like**, with many states in the ensemble carrying significant abundance.



# Abundance Distribution at present time

$r/s=30, r=3.5$



$M_0=531.9$  GeV,  $T_c=17.7$  GeV,  $M_s=152.0$  GeV

*Fundamental scales*

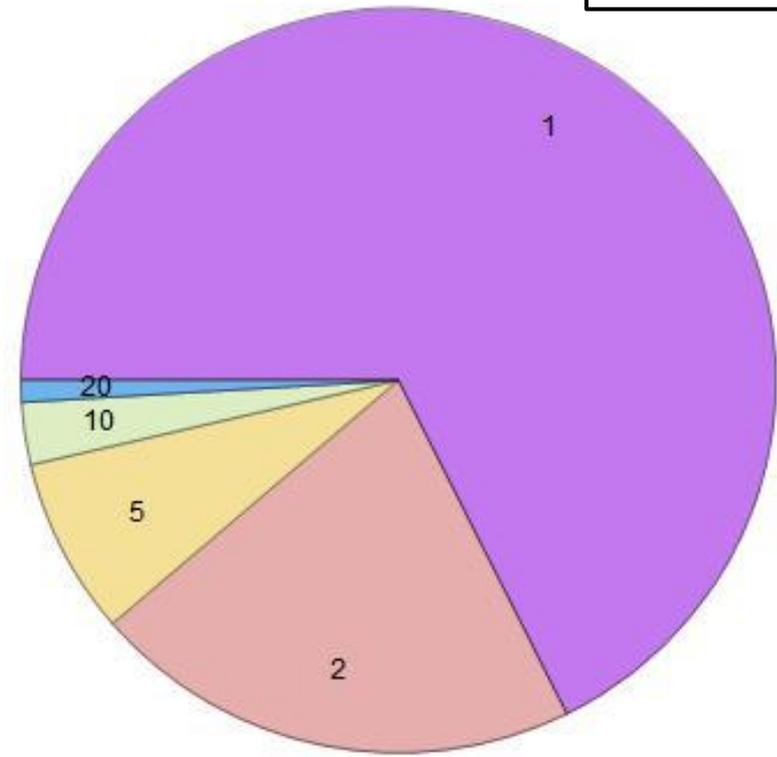
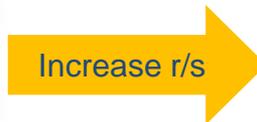
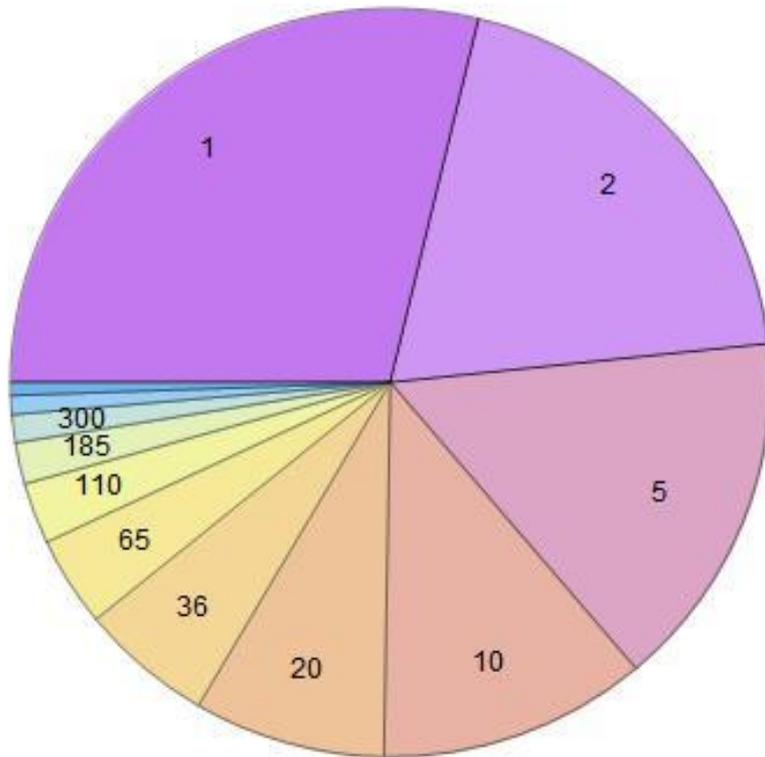
# Abundance Distribution at present time

$r/s=30, r=3.5$

$r/s=50, r=3.5$

$$r \equiv M_0/M_S$$

$$s \equiv T_c/M_S$$



$M_0=531.9 \text{ GeV}, T_c=17.7 \text{ GeV}, M_S=152.0 \text{ GeV}$

$M_0=42.9 \times 10^{10} \text{ GeV}, T_c=0.9 \times 10^{10} \text{ GeV}, M_S=12.2 \times 10^{10} \text{ GeV}$

Increasing  $r/s$  while holding  $r$  fixed amounts to decreasing  $s$ , leading to a stronger Boltzmann suppression for heavier states, such that

- Tower fraction is decreased, less DDM-like
  - Corresponding mass scales are increased
- } Tower fraction and mass scales are **INVERSELY CORRELATED!**

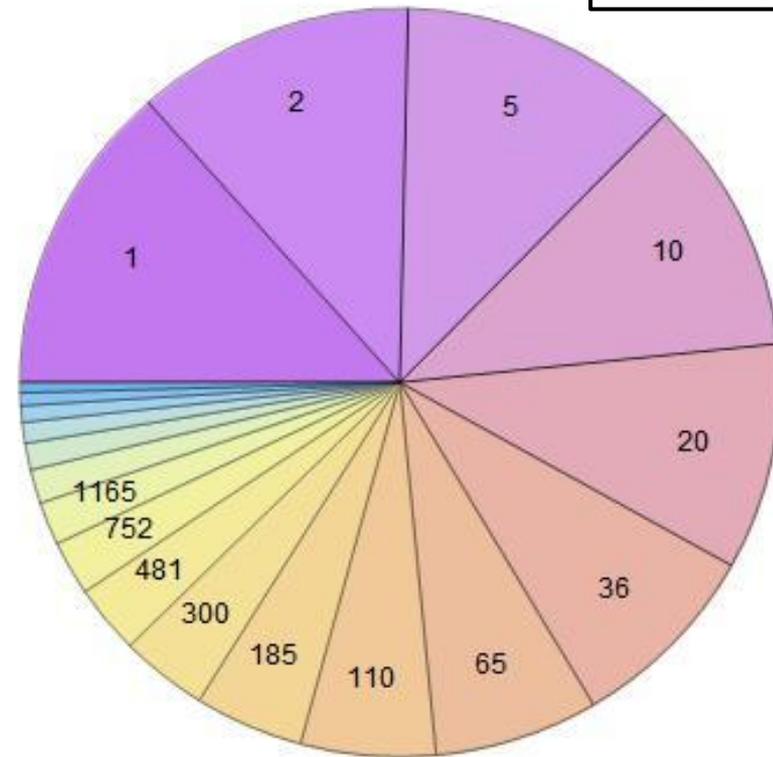
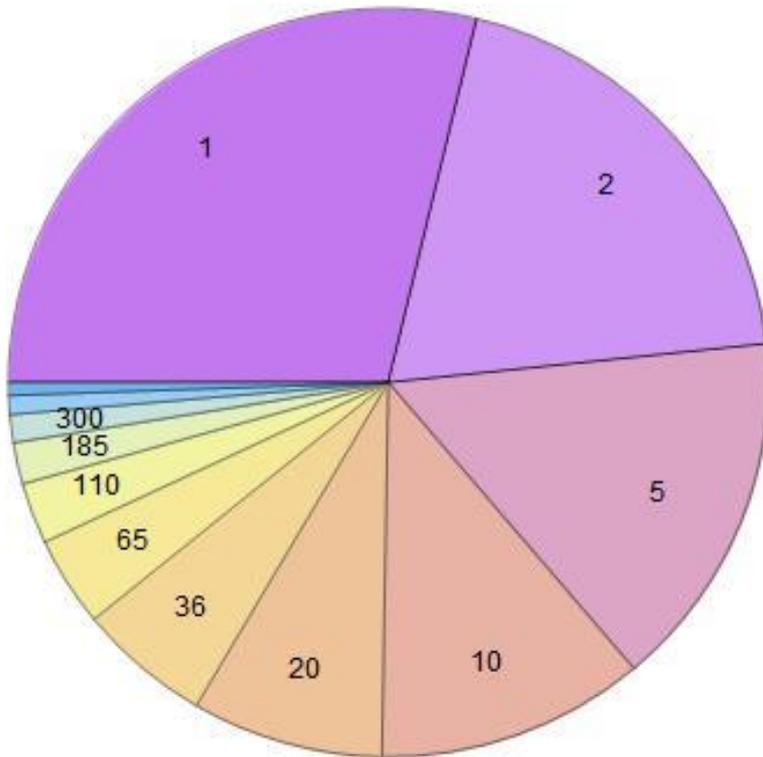
# Abundance Distribution at present time

$r/s=30, r=3.5$

$r/s=30, r=4$

$$r \equiv M_0/M_s$$

$$s \equiv T_c/M_s$$



$M_0=531.9 \text{ GeV}, T_c=17.7 \text{ GeV}, M_s=152.0 \text{ GeV}$

$M_0=248.5 \text{ GeV}, T_c=8.3 \text{ GeV}, M_s=62.1 \text{ GeV}$

Increasing  $r$  while holding  $r/s$  fixed increases  $s$  while decreases mass ratio between heavier and lighter states, leading to weaker Boltzmann suppression

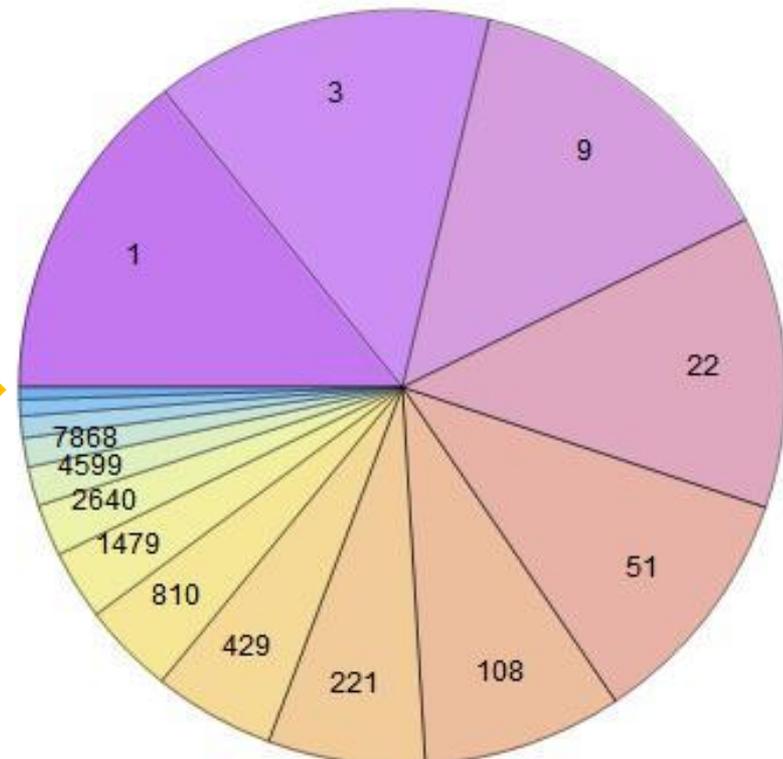
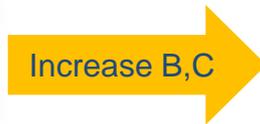
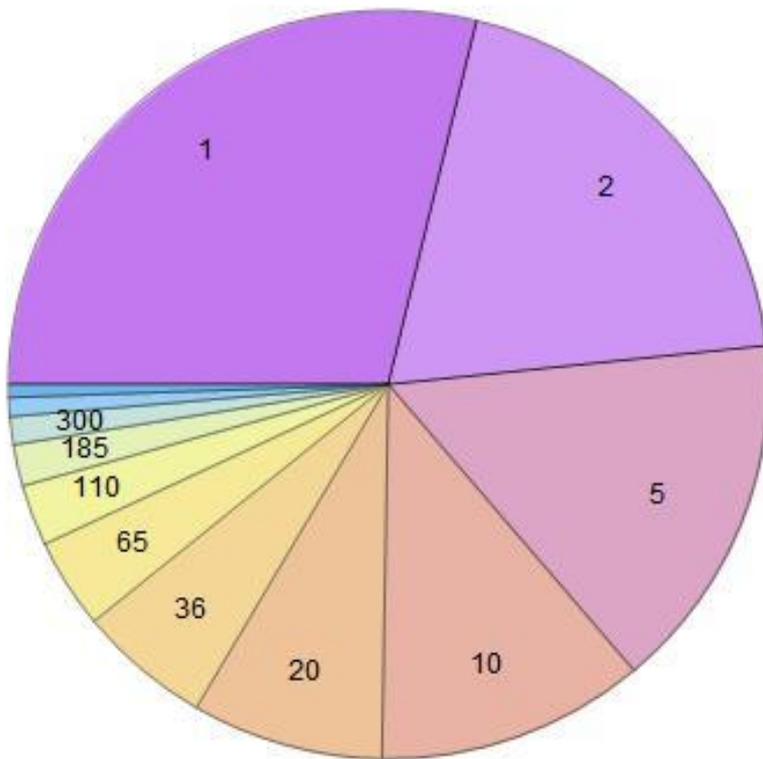
- Heavier states become more abundant, more DDM-like
  - Corresponding mass scales are decreased
- Correlation appears again!*

# Abundance Distribution at present time

$r/s=30, r=3.5$

$r/s=30, r=3.5$   
 $(C=\sqrt{2}\pi, B=3/2)$

$$g_n \sim e^{C\sqrt{n}}$$



$M_0=531.9 \text{ GeV}, T_c=17.7 \text{ GeV}, M_s=152.0 \text{ GeV}$

$M_0=264.5 \text{ GeV}, T_c=8.8 \text{ GeV}, M_s=75.6 \text{ GeV}$

Increasing  $C$  increases degeneracy thus allowing more states to contribute.

- Higher levels become more abundant, more DDM-like
- Corresponding mass scales are decreased

Still, *inversely correlated!*

Ground state need NOT to be the most abundant one at  $t_{now}$

## Conclusions

- DDM from strongly coupled dark sector which exhibits Hagedorn behavior is potentially viable since a natural Boltzmann distribution provides an exponential suppression.
- Evolution of  $\Omega_{tot}(t)$ ,  $w_{eff}(t)$  and  $\eta(t)$  show nontrivial time dependence.
- A large parameter space is explored, fundamental scales, i.e.,  $M_0$ ,  $M_S$  and  $T_c$ , can be dialed from keV/MeV all the way to  $M_{Pl}$ .
- Lower-mass scenarios are more DDM-like, with many states contributing.
- Higher-mass scenarios are less DDM-like, thus more traditional
- DDM framework smoothly incorporates entire transition between different behaviors

## Future Prospects

- A full phenomenological study of this model.
- Including intra-ensemble decay in which heavier states can cascade down the tower and decay into lighter states.
- A detailed study of DDM freeze-out, how it affects DDM abundance distribution.
- How DDM affects small structure.