

Extracting Dark Matter Information at the LHC

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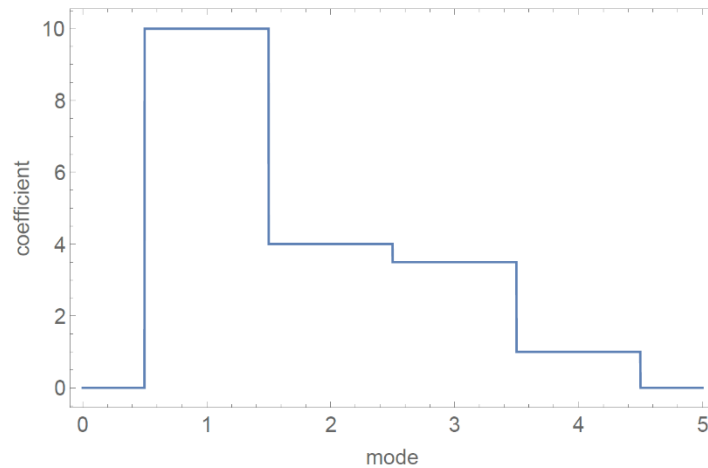
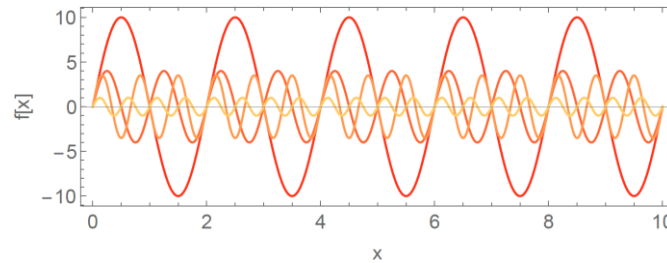
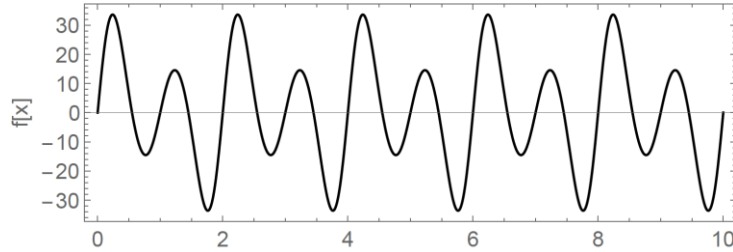
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Signals with Missing Transverse Energy at Colliders

- Motivation : Dark Matter
- We can extract many information of DM once we see a signal excess.
- What kind of information? : DM mass, couplings, interaction type(e.g., s-wave int., p-wave int., whether they can form a bound state, ...)
- How? :
 1. Theory calculation and fitting
 - applicable for any models, but time consuming
 2. Spectral decomposition
 - applicable only for s-channel mediator models, but efficient and less biased

Spectral Decomposition

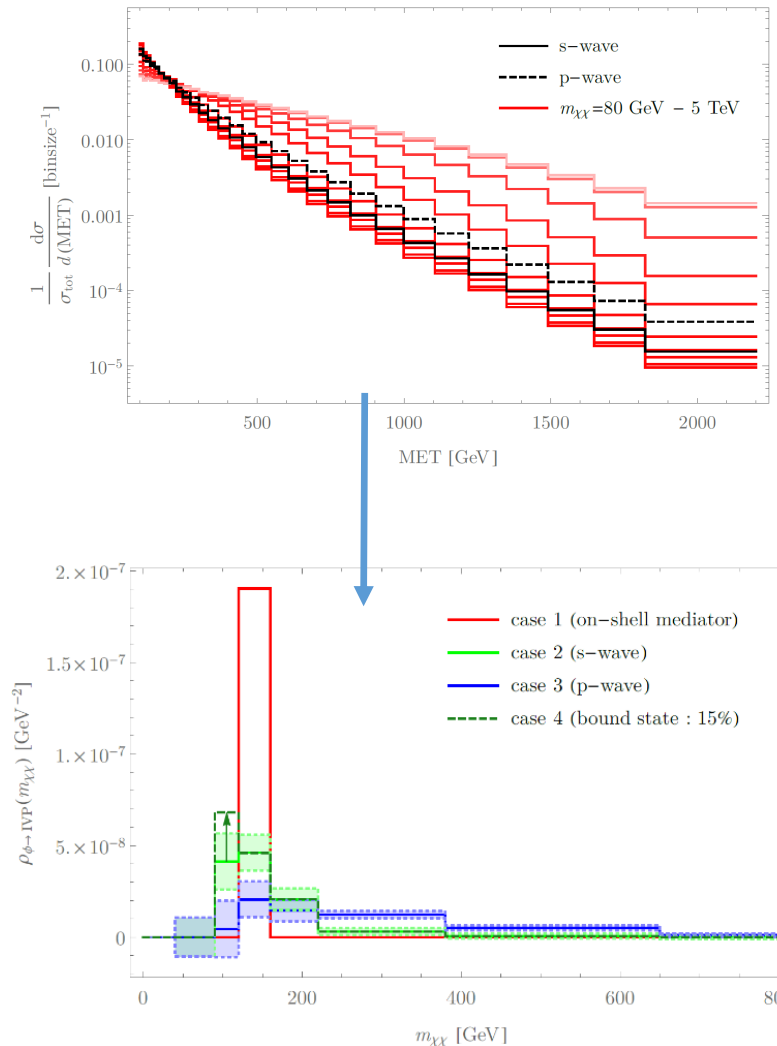
- As a classical example, whenever we see a periodic function, ...



Basis functions : $\sin(n\pi x/L)$ w/ $n = \text{integer}$
Coefficients : amplitude of each mode

Spectral Decomposition

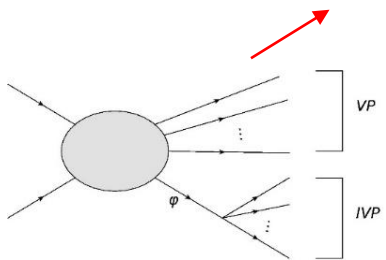
- I am going to do the similar thing to MET distribution.



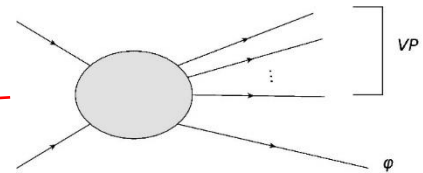
Basis functions in this case?
Physical meaning of coefficients?

Spectral Decomposition

- Decomposition is guaranteed by following equations if the mediator is s-channel.



$$\begin{aligned}
 \hat{\sigma}_{\text{full}} &= \frac{1}{2s} \int d\Phi_{\text{VP}} d\Phi_{\text{IVP}} |\mathcal{M}_\phi|^2 |G_\phi(p_\phi^2, M_\phi^2)|^2 |\mathcal{M}_{\phi \rightarrow \text{IVP}}|^2 (2\pi)^4 \delta^{(4)}\left(\sum_{i \in \text{ext}} P_i\right) \\
 &\times \int dp_\phi^0 \frac{d^3 \vec{p}_\phi}{(2\pi)^3} (2\pi)^3 \delta^{(4)}\left(p_\phi - \sum_{i \in \text{IVP}} p_i\right) \\
 &= \int dm_\phi^2 (\hat{\sigma}_\phi)_{p_\phi^2 = m_\phi^2} |G_\phi(m_\phi^2, M_\phi^2)|^2 \\
 &\times \frac{1}{2\pi} \int d\Phi_{\text{IVP}} |\mathcal{M}_{\phi \rightarrow \text{IVP}}|^2 (2\pi)^4 \delta^{(4)}\left(p_\phi - \sum_{i \in \text{IVP}} p_i\right) \\
 &= \int dm_\phi^2 (\hat{\sigma}_\phi)_{p_\phi^2 = m_\phi^2} |G_\phi(m_\phi^2, M_\phi^2)|^2 \frac{1}{\pi} (-\text{Im}\Sigma_{\phi \rightarrow \text{IVP}}) \\
 &= \int dm_\phi^2 (\hat{\sigma}_\phi)_{p_\phi^2 = m_\phi^2} \times \rho_{\phi \rightarrow \text{DM}}(S_0, M_\phi^2),
 \end{aligned}$$



Narrow Width Approximation
 $(\sigma_{\text{full}} = \sigma_\phi \text{ Br.})$

(Partial spectral density)

Spectral Decomposition

- Decomposition is guaranteed by following equations if the mediator is s-channel.

Spin = 1 mediator?

$$\hat{\sigma}_{\text{full}} = \frac{1}{2s} \int d\Phi_{\text{VP}} d\Phi_{\text{IVP}} |\mathcal{M}_\phi|^2 |G_\phi(p_\phi^2, M_\phi^2)|^2 |\mathcal{M}_{\phi \rightarrow \text{IVP}}|^2 (2\pi)^4 \delta^{(4)}\left(\sum_{i \in \text{ext}} P_i\right) \\ \times \int dp_\phi^0 \frac{d^3 \vec{p}_\phi}{(2\pi)^3} (2\pi)^3 \delta^{(4)}\left(p_\phi - \sum_{i \in \text{IVP}} p_i\right)$$

$$G_T(p_\phi^2, M_\phi^2) \left(-g_{\mu\nu} + \frac{p_{\phi\mu} p_{\phi\nu}}{p_\phi^2}\right) + G_L(p_\phi^2, M_\phi^2) \frac{p_{\phi\mu} p_{\phi\nu}}{p_\phi^2}$$

$$G_T(p_\phi^2, M_\phi^2) = -i/(k^2 - M_\phi^2 - \Pi_T) \text{ and } G_L(p_\phi^2, M_\phi^2) = i/(M_\phi^2 + \Pi_L)$$

$$\Pi_{\mu\nu} = \left(-g_{\mu\nu} + \frac{p_{\phi\mu} p_{\phi\nu}}{p_\phi^2}\right) \Pi_T + \frac{p_{\phi\mu} p_{\phi\nu}}{p_\phi^2} \Pi_L$$

$$\left| \mathcal{M}_\phi^\mu G_\phi(p_\phi^2, M_\phi^2) \left(-g_{\mu\nu} + \frac{p_{\phi\mu} p_{\phi\nu}}{p_\phi^2}\right) \mathcal{M}_{\phi \rightarrow \text{IVP}}^\nu \right|^2 \\ = \left| \sum_\lambda \mathcal{M}_\phi^\mu G_\phi(p_\phi^2, M_\phi^2) (\epsilon_{\lambda,\mu}^* \epsilon_{\lambda,\nu}) \mathcal{M}_{\phi \rightarrow \text{IVP}}^\nu \right|^2 \\ \rightarrow \frac{1}{3} \sum_{\lambda, \lambda'} \left| \left(\mathcal{M}_\phi^\mu \epsilon_{\lambda,\mu}^* \right) G_\phi(p_\phi^2, M_\phi^2) \left(\mathcal{M}_{\phi \rightarrow \text{IVP}}^\nu \epsilon_{\lambda',\nu} \right) \right|^2$$

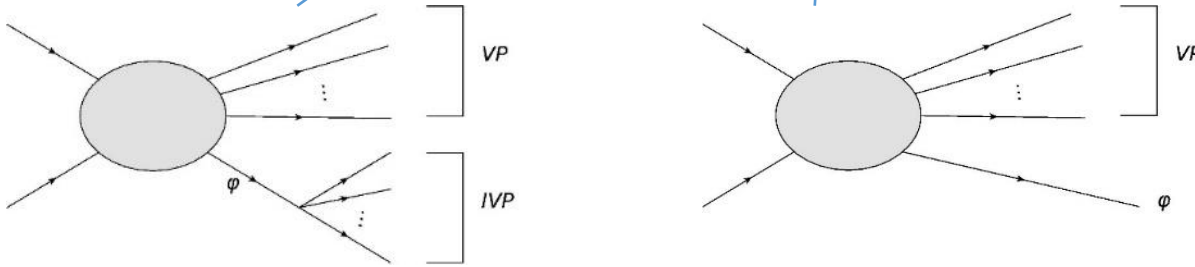
Nonzero only when the interaction is axial (involving gamma 5)

Spectral Decomposition

Kallan-Lehmann
Spectral density

$$\hat{\sigma}_{\text{full}} = \int dm_{\chi\chi}^2 \hat{\sigma}_{\phi}(m_{\chi\chi}) \times \rho_{\phi \rightarrow \text{IVP}}(m_{\chi\chi}^2, M_{\phi}^2)$$

virtual mass of the mediator



$$\frac{d\sigma_{\text{full}}}{d\mathcal{E}_T} = \int_0^{m_{\text{max}}^2} dm_{\chi\chi}^2 \frac{d\sigma_{\phi}(m_{\chi\chi})}{d\mathcal{E}_T} \times \rho_{\phi \rightarrow \text{IVP}}(m_{\chi\chi}, M_{\phi}^2)$$

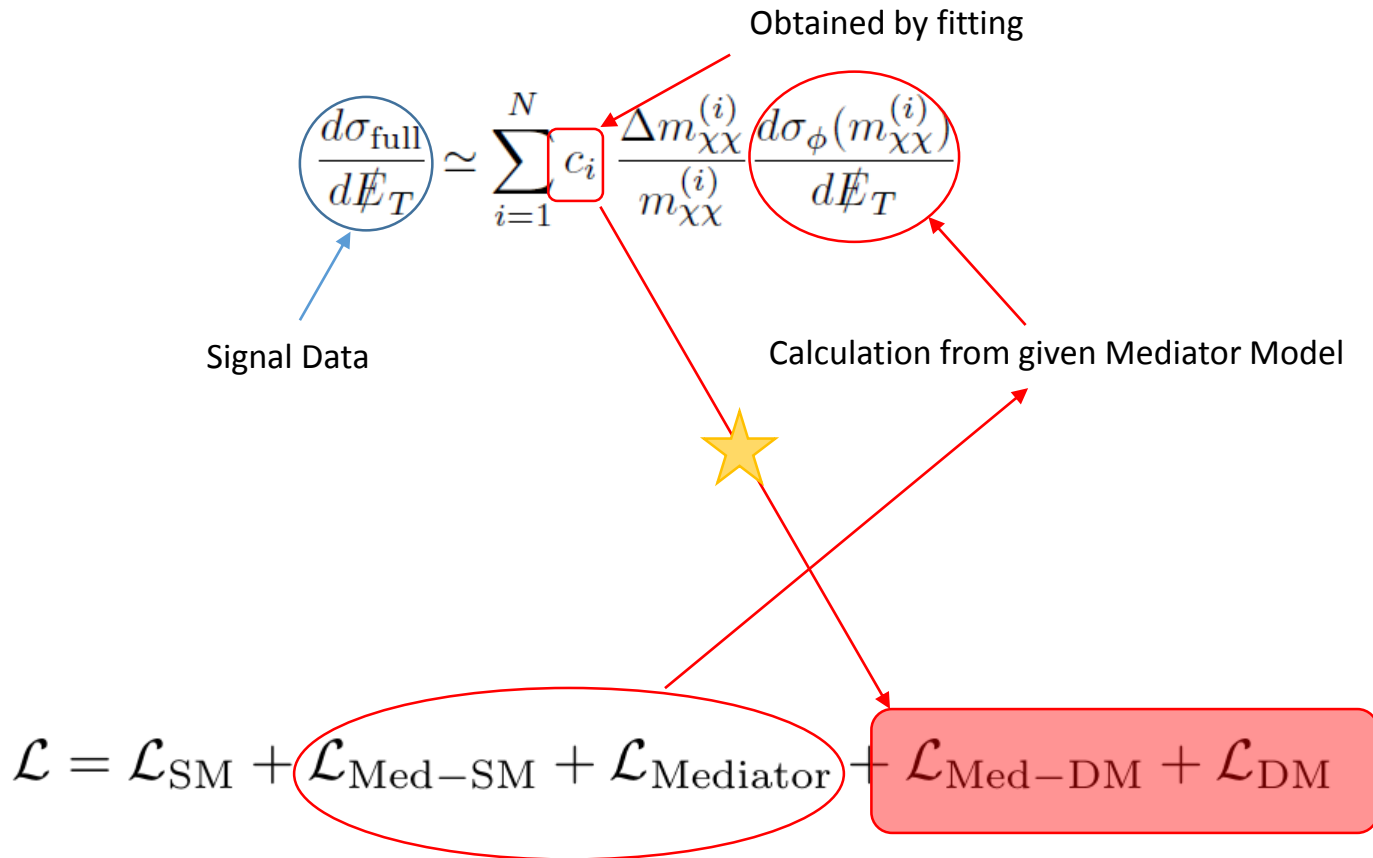
$$\frac{d\sigma_{\text{full}}}{d\mathcal{E}_T} \simeq \sum_{i=1}^N c_i \frac{\Delta m_{\chi\chi}^{(i)}}{m_{\chi\chi}^{(i)}} \frac{d\sigma_{\phi}(m_{\chi\chi}^{(i)})}{d\mathcal{E}_T}$$

coefficient

basis function

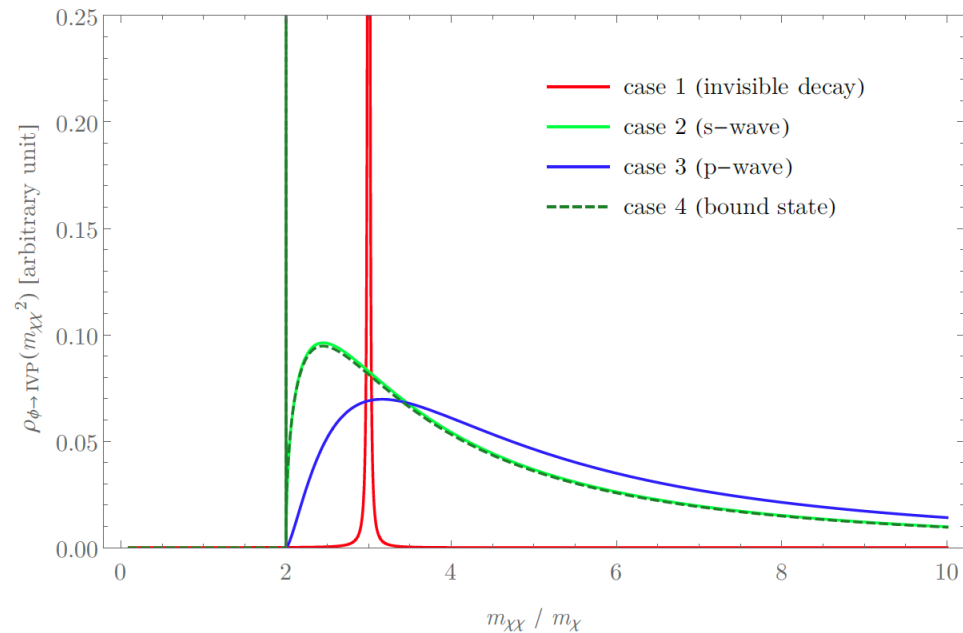
(Mediator model dependent)

Spectral Decomposition



Physical meaning of coefficients = spectral density

	Mediator	Interaction
Case 1	On-shell($M_\phi > 2m_\chi$)	Resonance
Case 2	Off-shell($M_\phi < 2m_\chi$)	S-wave
Case 3	Off-shell($M_\phi < 2m_\chi$)	P-wave
Case 4	Off-shell($M_\phi < 2m_\chi$)	S-wave & bound state (long range force)



Example : Mono-jet channel w/ Simplified Model

1. Simplified model


$$\mathcal{L} = \mathcal{L}_{\text{SM}} + \frac{1}{2} \partial^\mu \phi \partial_\mu \phi - \frac{1}{2} m_\phi^2 \phi^2 + \frac{k}{\Lambda} \phi G^{a\mu\nu} G_{\mu\nu}^a \\ + \mathcal{L}_{1/2\text{-DM}} + \mathcal{L}_{0\text{-DM}}$$

$$\mathcal{L}_{1/2\text{-DM}} = \bar{\chi}(i\not{\partial} - m_\chi)\chi + y_\chi \phi \bar{\chi}\chi$$

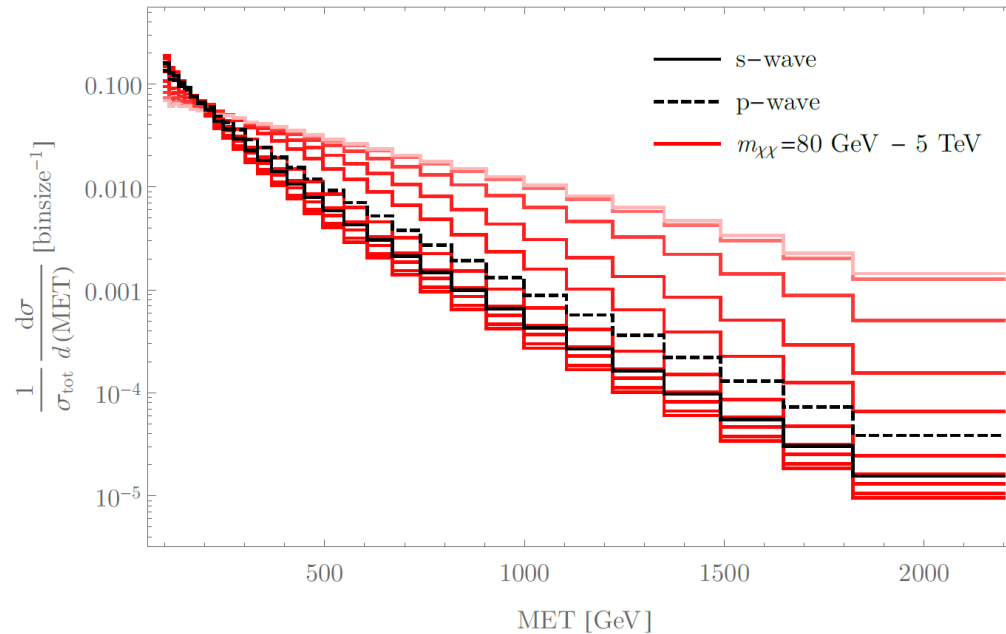
$$\mathcal{L}_{0\text{-DM}} = \frac{1}{2} \partial^\mu s \partial_\mu s - \frac{1}{2} m_s^2 s^2 + \frac{1}{2} m_s g_s \phi s^2$$

2. Parton level

3. Tree level


$$\frac{d\sigma_{\text{full}}}{d\cancel{E}_T} \simeq \sum_{i=1}^N c_i \frac{\Delta m_{\chi\chi}^{(i)}}{m_{\chi\chi}^{(i)}} \frac{d\sigma_\phi(m_{\chi\chi}^{(i)})}{d\cancel{E}_T}$$

Example : Mono-jet channel w/ Simplified Model



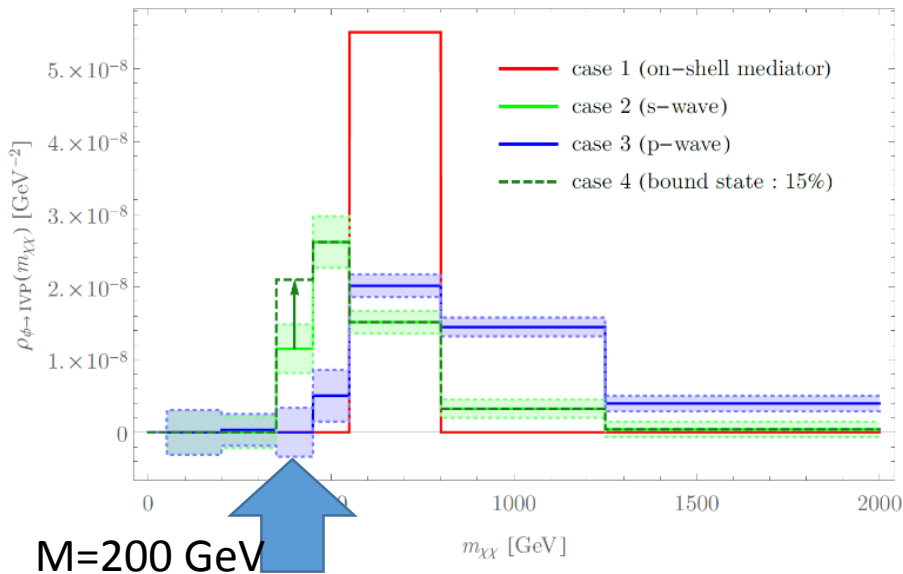
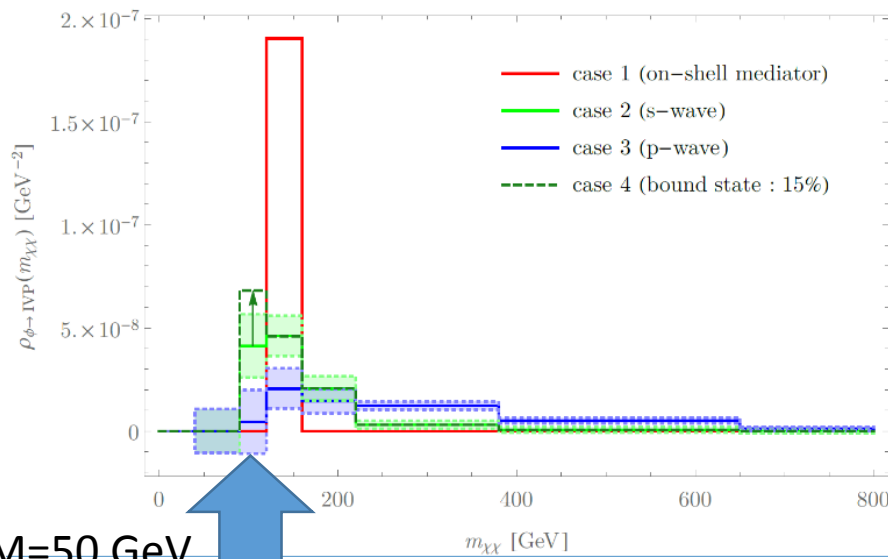
Lessons

1. As m_ϕ increases, events are more distributed at hard MET (PT of jet) region.
2. At hard PT region, m_ϕ dependence of the slope becomes weaker.
(, but S/B ratio becomes bigger. => some balance on PT cut)

Example : Mono-jet channel w/ Simplified Model

$$L=3 \text{ ab}^{-1}, \quad S/\sqrt{S+B} = 10$$

(If the current significance ~ 1 ,
then probably it would be 10 for HL-LHC)



Summary

1. Spectral decomposition approach is suggested as a method to extract information of dark matter.
2. We applied it to the monojet channel with simplified model. It can be applied to any other channels, but only with s-channel mediator models.

And you can play with other kinematics, as well.

Thank you!

Back Up

luminosity function

$$\sigma_\phi = \sum_{i,j} \int dx_1 dx_2 P_i(x_1) P_j(x_2) \hat{\sigma}_\phi^{(ij)}$$

$$= \sum_{i,j} \int dm_{ij} dy \frac{dL_{ij}}{dm_{ij} dy} \hat{\sigma}_\phi^{(ij)}$$

$$\frac{dL_{ij}}{dm_{ij} dy} = \frac{2m_{ij}}{S} P_i\left(\frac{m_{ij}}{\sqrt{S}} e^y\right) P_j\left(\frac{m_{ij}}{\sqrt{S}} e^{-y}\right)$$

$$\frac{d\sigma_\phi}{d\cancel{p}_T} = \sum_{i,j} \int_{\cancel{p}_T + \sqrt{\cancel{p}_T^2 + m_\phi^2}}^{\sqrt{S}} dm_{ij} f_{ij}(m_{ij}) \frac{d\hat{\sigma}_\phi^{(ij)}}{d\cancel{p}_T}$$

$$\frac{\frac{d\sigma_\phi}{d\cancel{p}_T}(\cancel{p}_T, m_{ij}, m_\phi)}{\frac{d\sigma_\phi}{d\cancel{p}_T}(0, m_{ij}, m_\phi)} \sim \frac{f_{ij}(\cancel{p}_T + \sqrt{\cancel{p}_T^2 + m_\phi^2})}{f_{ij}(m_\phi)}$$

