



Searches for dark matter at CMS

Speaker: Marco Cipriani

on behalf of the CMS collaboration

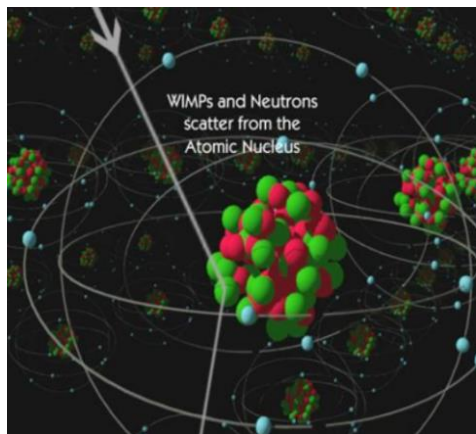
Sapienza Università di Roma & INFN Roma

Phenomenology Symposium 2017

Pittsburgh, 8-10 May 2017

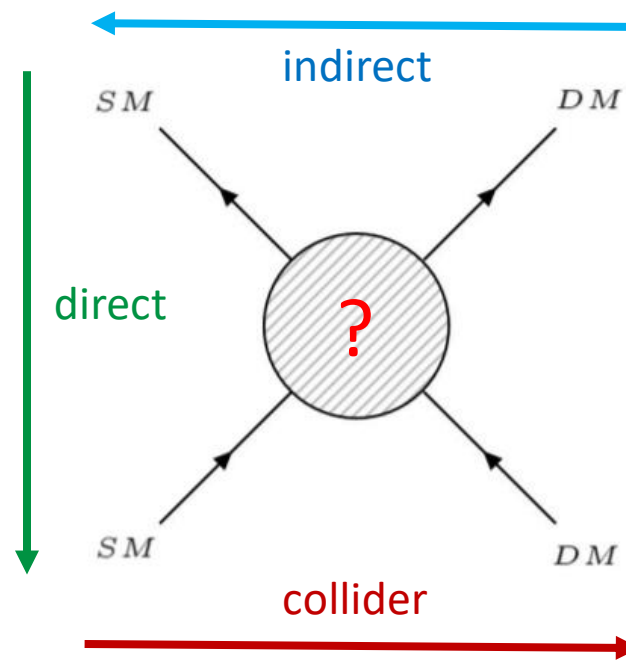
Searching for dark matter (DM)

- ❑ DM probably made of **Weakly Interacting Massive Particles (WIMPs)**
- ❑ 3 types of searches, high degree of complementarity
- ❑ sensitivity depends on interaction details

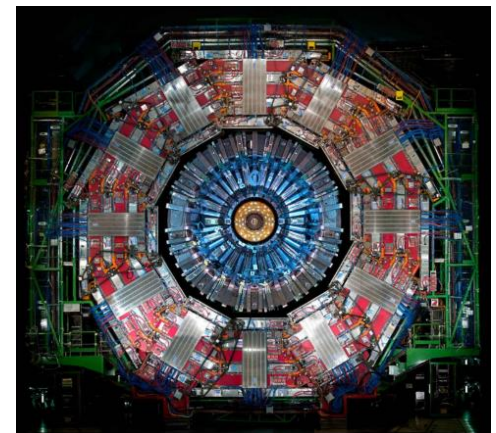


DM scattering on nuclei

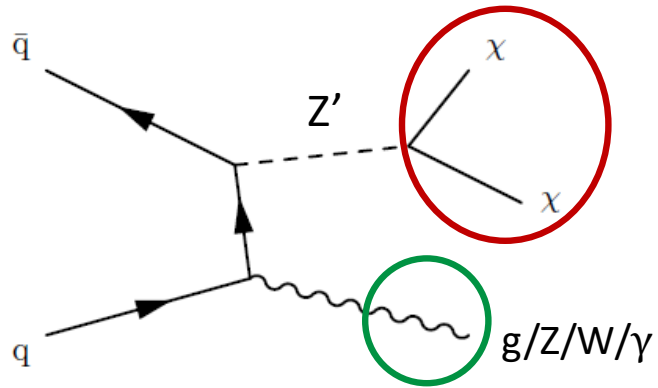
DM annihilation to SM particles



DM production in pp collisions



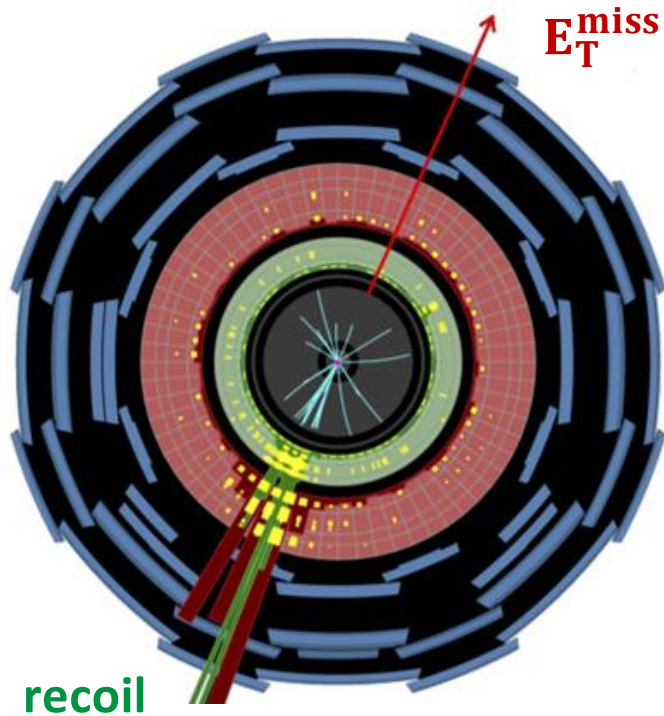
Dark Matter at colliders



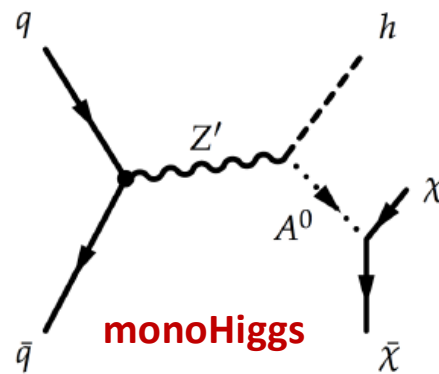
monoJet/Z/W/ γ

monoX searches:

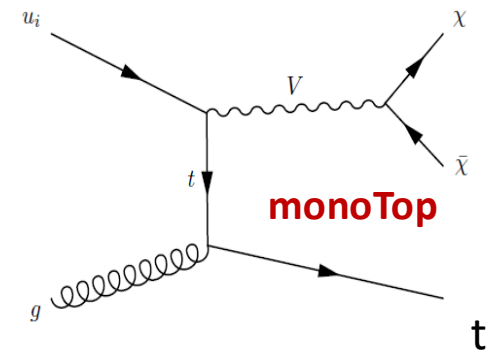
- **DM:** missing transverse energy (E_T^{miss} , **MET**)
- **X:** **Standard Model** (SM) particles from initial state radiation (**ISR**) \rightarrow **trigger** on the event
- **signal:** excess of events in the high MET region



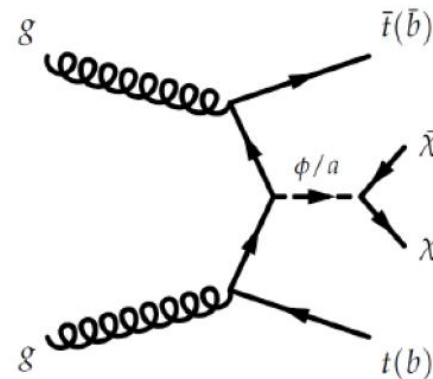
recoil



monoHiggs

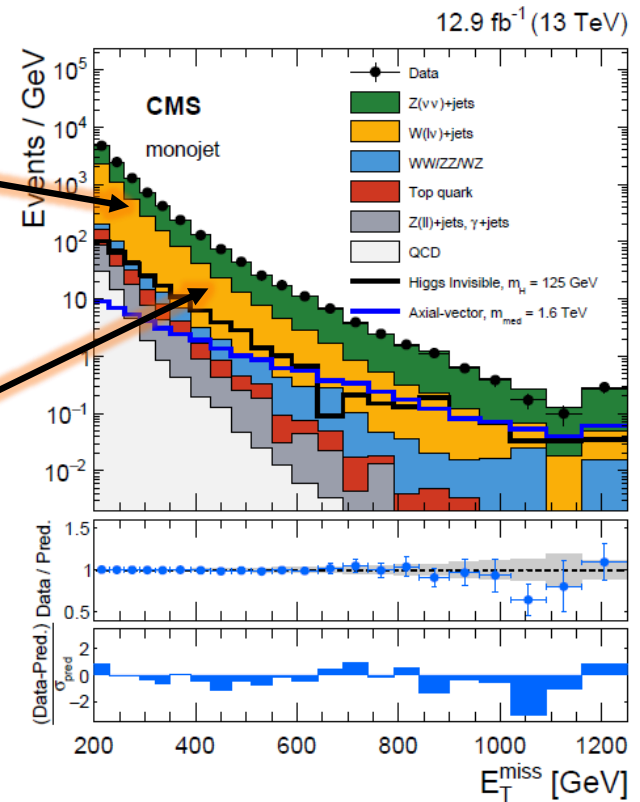
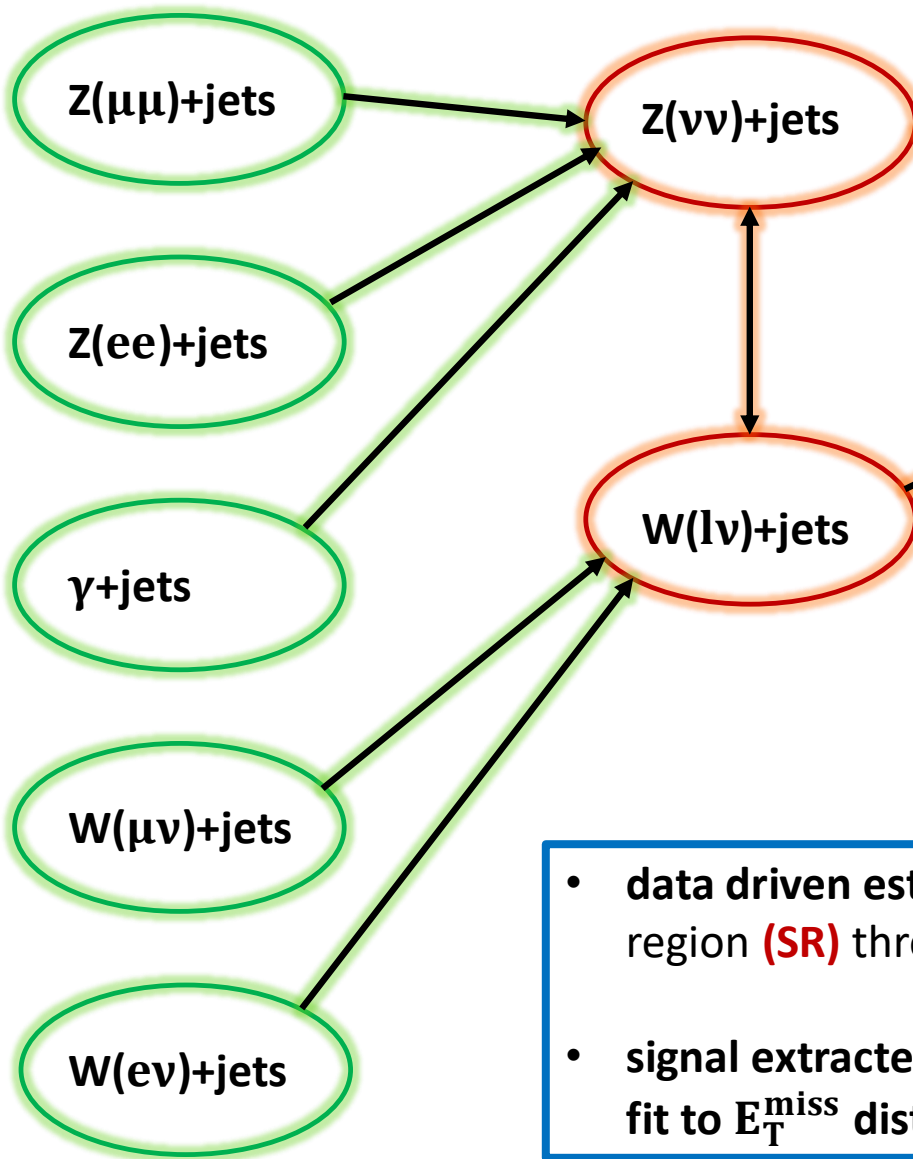


monoTop



DM + heavy quarks

General analysis strategy



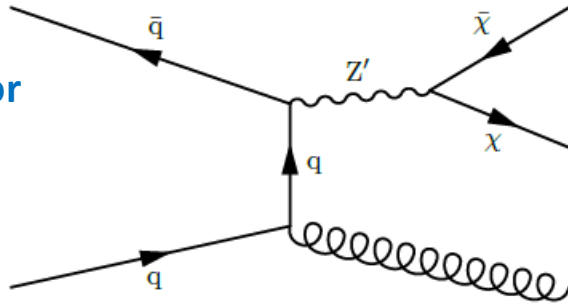
- data driven estimate of **irreducible backgrounds** in signal region (**SR**) through many **control regions (CR)**
- signal extracted through a combined maximum likelihood fit to E_T^{miss} distribution in SR and CR

MonoJet/V overview

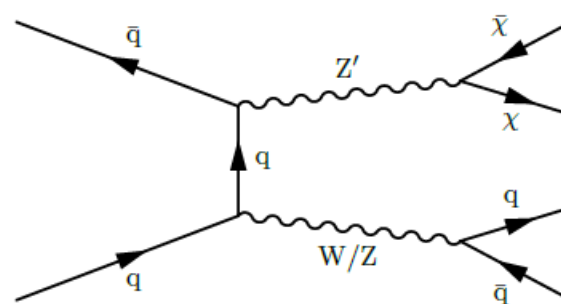
Greatest potential: sensitive to **all** possible mediator coupling structure

- vector, axial-vector, scalar, pseudoscalar
- 2 exclusive category: jet from quark/gluon radiation or W/Z hadronic decay
- can **probe many theoretical scenarios**, unlike other channels

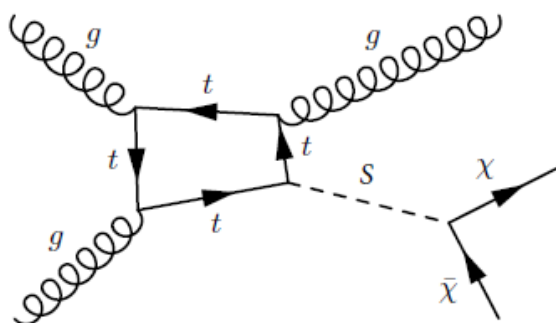
(axial)vector
mediator,
monoJet



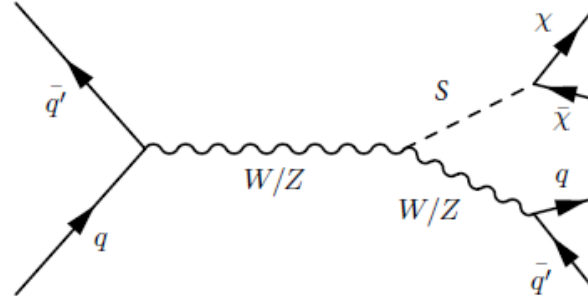
(axial)vector
mediator,
monoV



(pseudo)scalar
mediator,
monoJet




(pseudo)scalar
mediator,
monoV

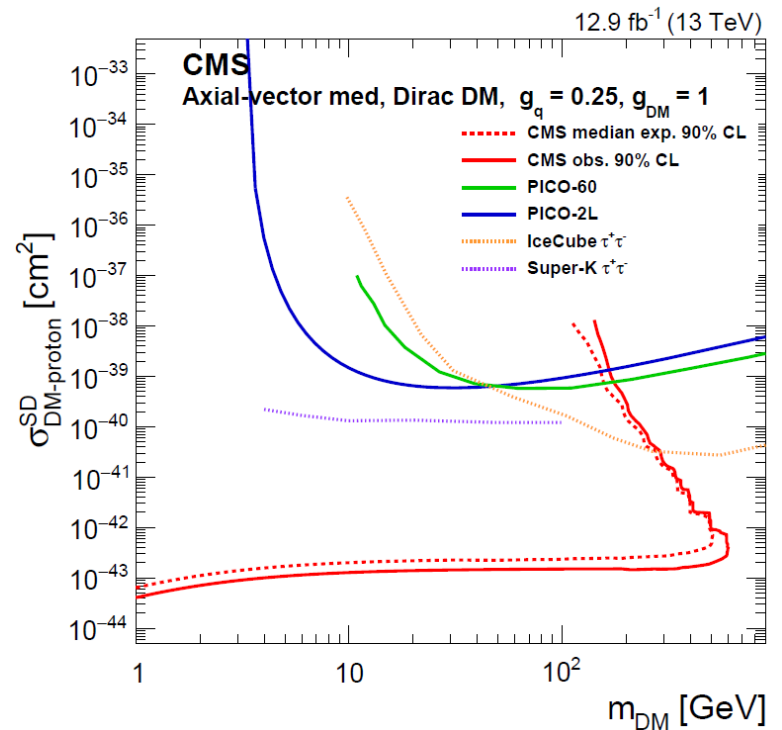
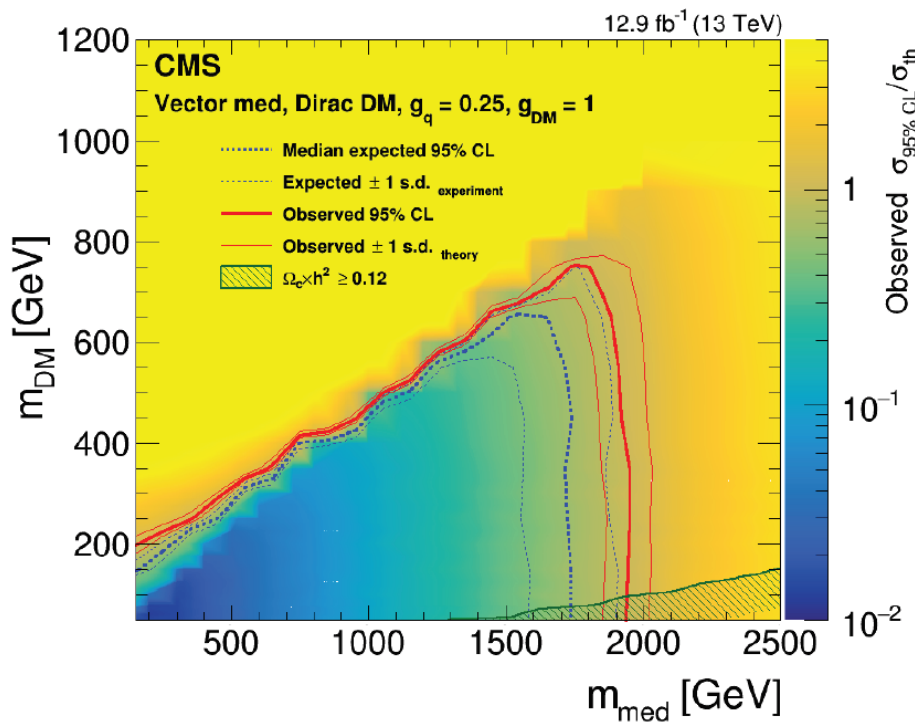


MonoJet/V results

Analysis on full 2016 dataset (36 fb^{-1}) under approval, showing ICHEP results (12.9 fb^{-1})

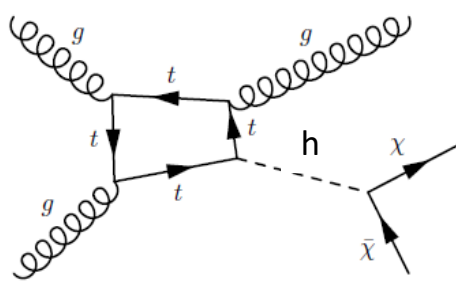
 CMS paper (12.9 fb^{-1}) → arXiv:1703.01651 [hep-ex], submitted to JHEP

- interpretation with **simplified models: (axial)vector, (pseudo)scalar mediators**
- $m_{\text{med}} < 1.95 \text{ TeV}$ for (axial)vector, $< 100(430) \text{ GeV}$ for scalar(pseudoscalar) mediator
 - translated into DM-nuclei cross section plane and compared to direct searches

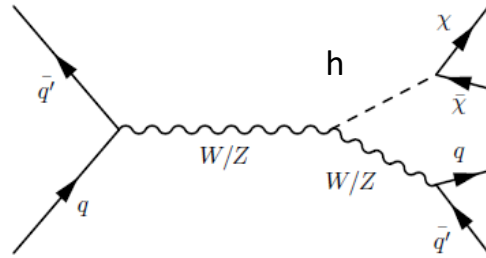


MonoJet Higgs \rightarrow DM reinterpretation

Category	Expected limit	Observed limit	± 1 s.d.	Expected signal composition
Mono-V	0.72	1.17	[0.51–1.02]	39.6% ggH, 6.9% VBF, 32.4% WH, 21.1% ZH
Monojet	0.85	0.48	[0.58–1.27]	71.5% ggH, 20.3% VBF, 4.4% WH, 3.8% ZH
Combined	0.56	0.44	[0.40–0.81]	-



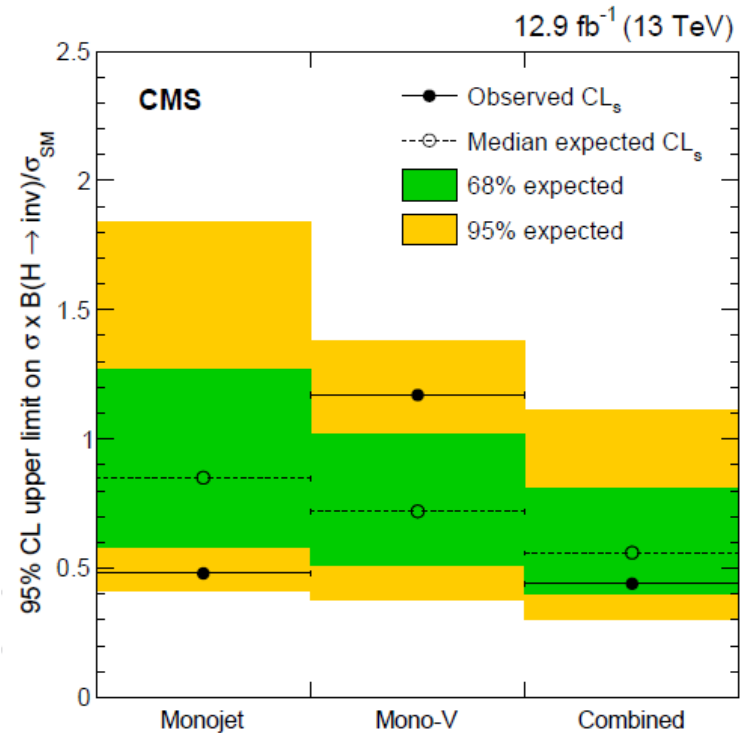
monoJet



monoV

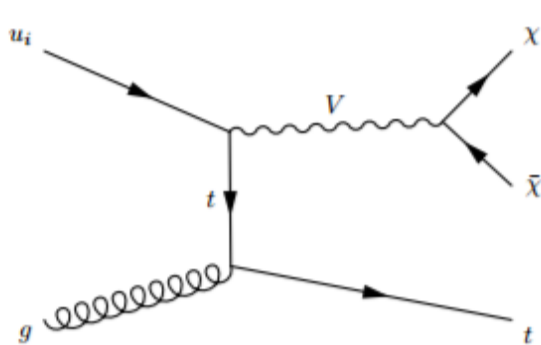
limits on SM Higgs \rightarrow invisible branching ratio

- monoJet/V combined limit: **BR < 0.44**
- current observed limit is **0.24** from **Higgs \rightarrow Inv direct searches**: arXiv:1610.09218v2 [hep-ex]

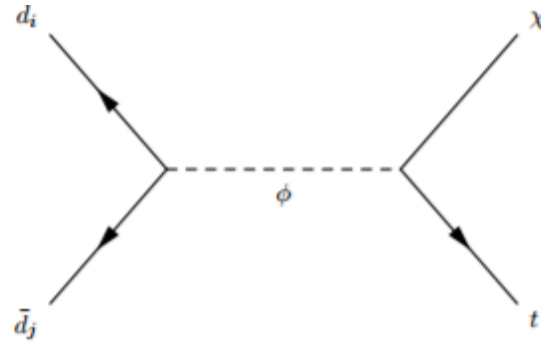


Mono Top

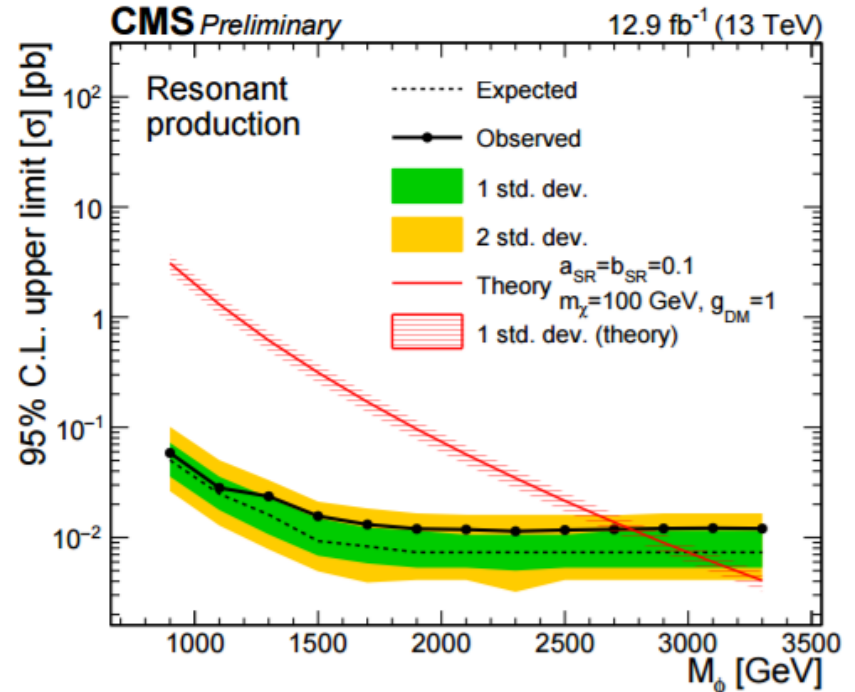
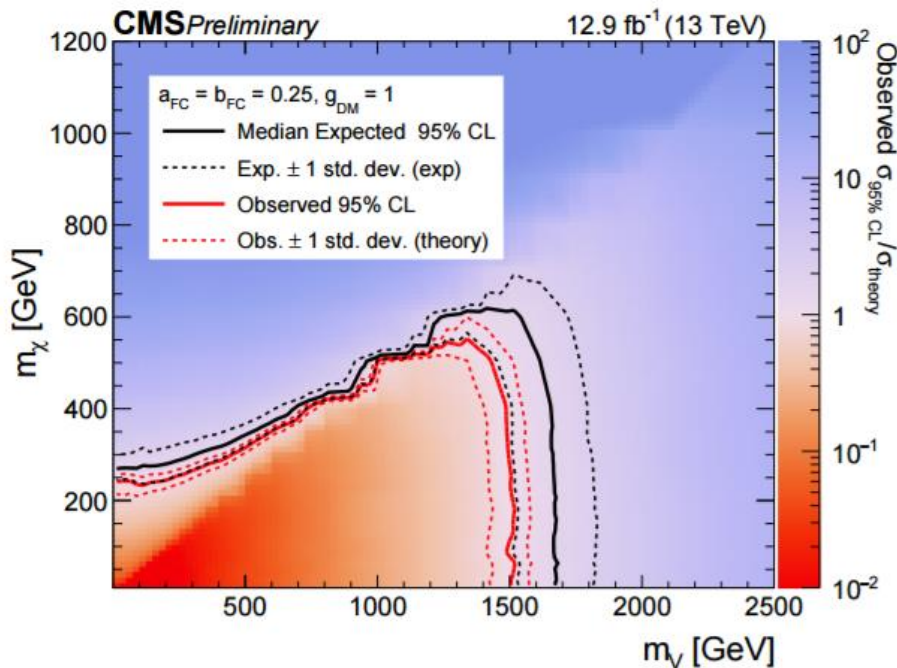
Search for MET and a single top quark decaying hadronically



flavour changing neutral current

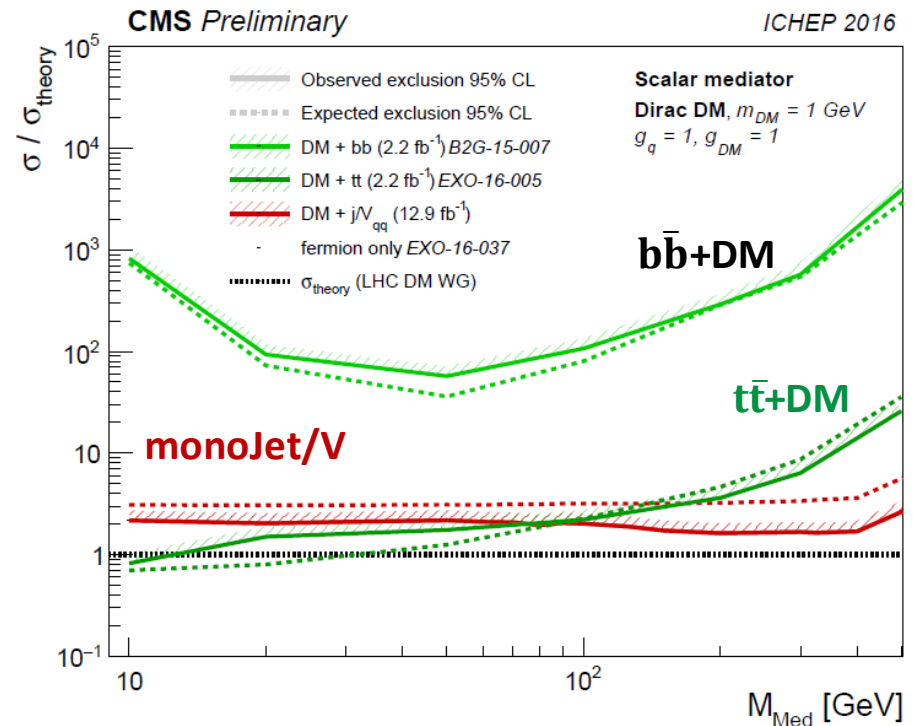
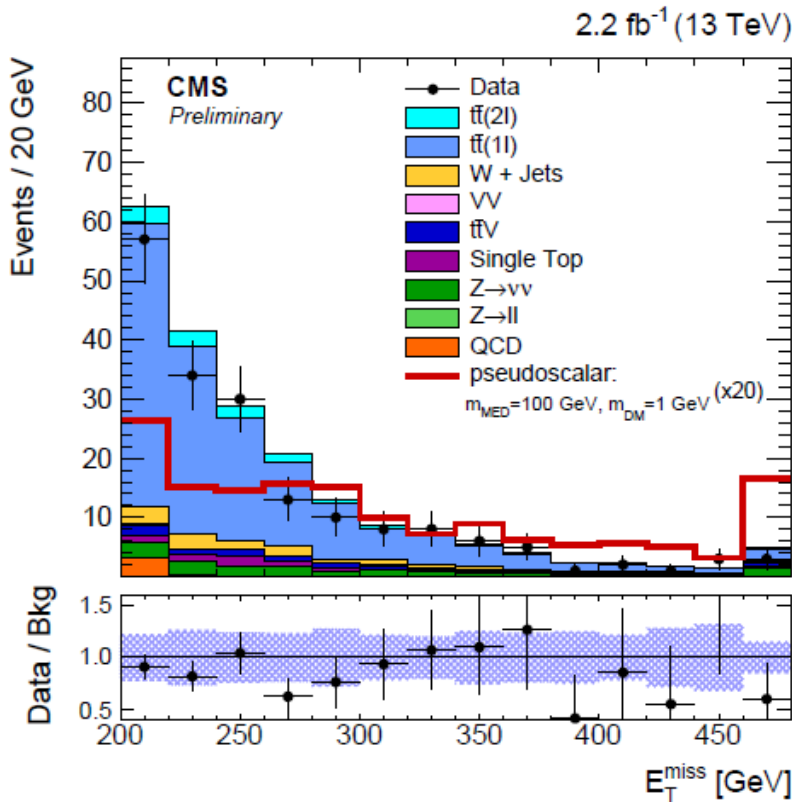
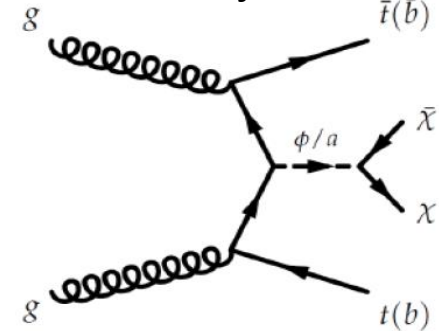


colored charged scalar ϕ



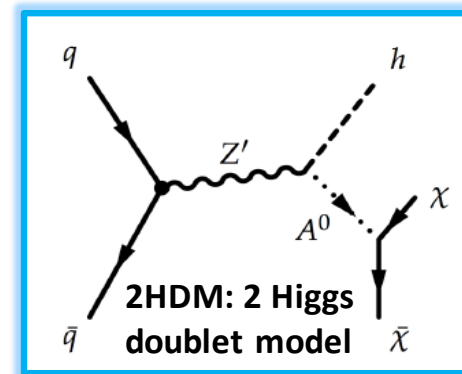
$t\bar{t}$ + Dark Matter

- **top tagging techniques:** identify top quarks decaying into three resolved jets
- **hadronic or semileptonic top decays** considered
- data-driven estimate of **SM $t\bar{t}$ and W/Z+jets backgrounds**
- sensitive to **(pseudo)scalar mediator**



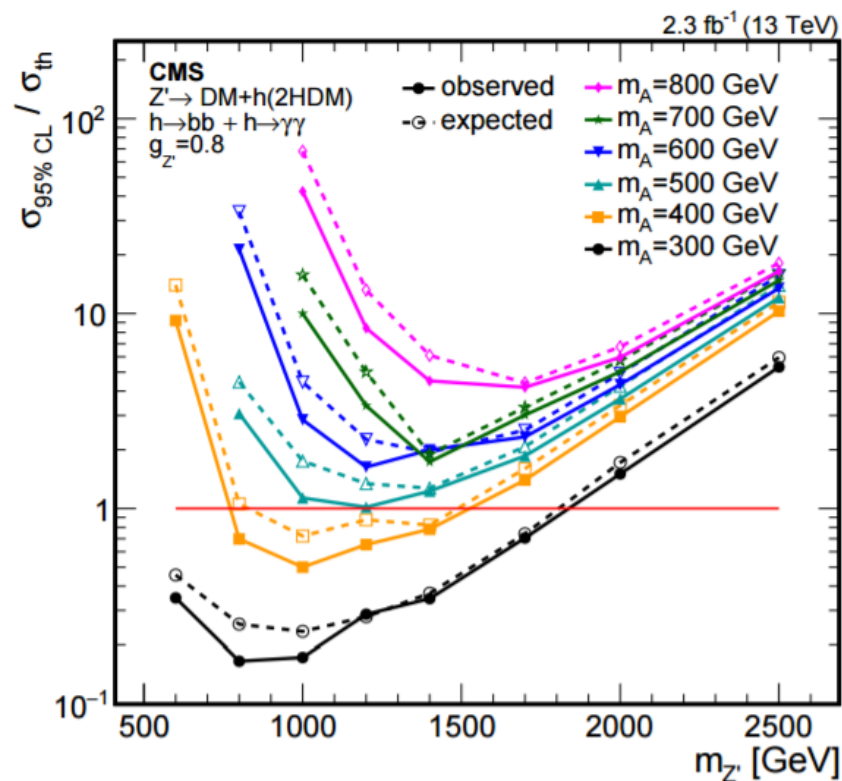
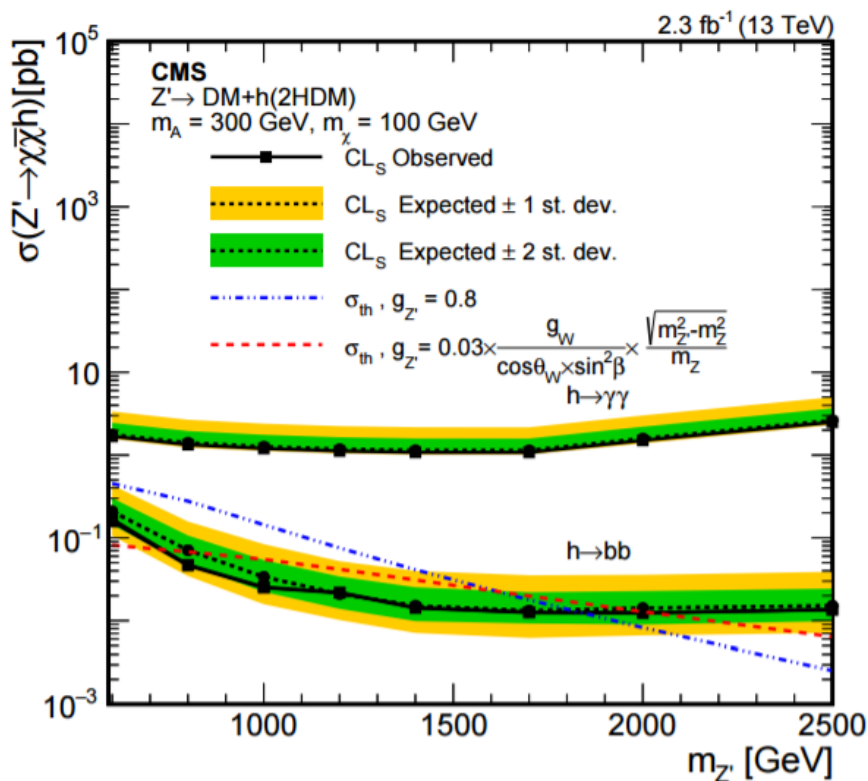
Probing Dark Matter with MonoHiggs

- Higgs ISR suppressed: **probe direct Higgs-DM coupling**
- search for MET+ SM Higgs \rightarrow $b\bar{b}$, $\gamma\gamma$, ZZ , $\tau\bar{\tau}$...



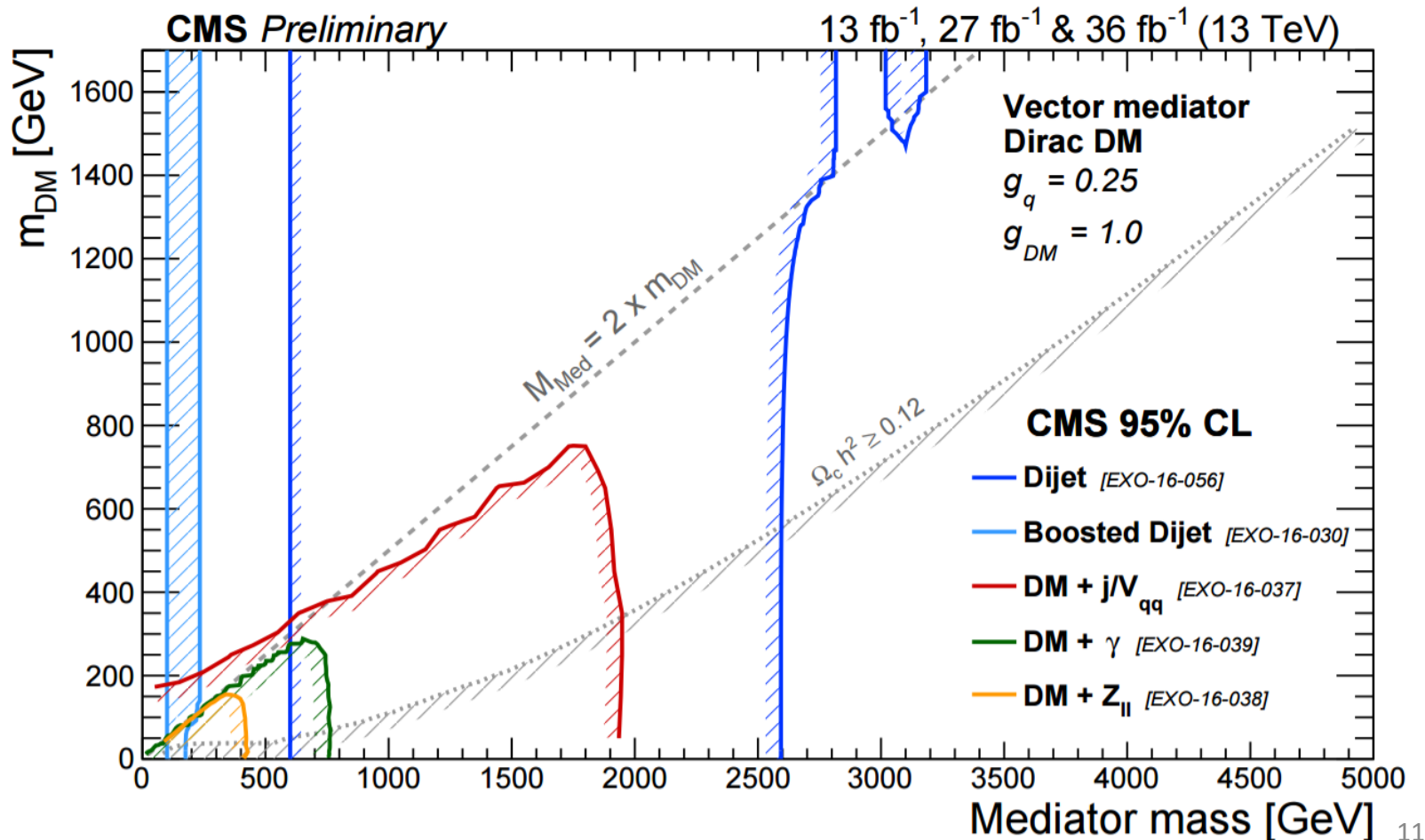
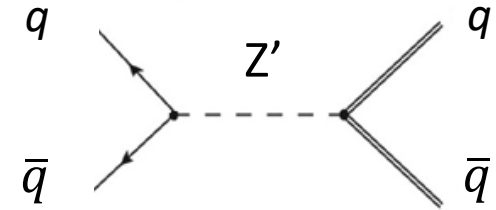
combination of $b\bar{b}$, $\gamma\gamma$ (2.3 fb^{-1} @ 13TeV) submitted to JHEP

- Z' -2HDM interpretation
- $m_{Z'} < 1860 \text{ GeV}$ for $g_{Z'} = 0.8$ and $m_A = 300 \text{ GeV}$



Constraints on DM from dijet search

- most powerful search at a hadron collider
- search for a bump over falling dijet mass spectrum
- limits on mediator mass



What next?

Many analysis on full 2016 dataset (36 fb⁻¹) being approved: signal region still blinded

Not just a bare recast of lower luminosity analyses

- treatment of theoretical uncertainties
- better constraints from control regions with more available data
- additional interpretations

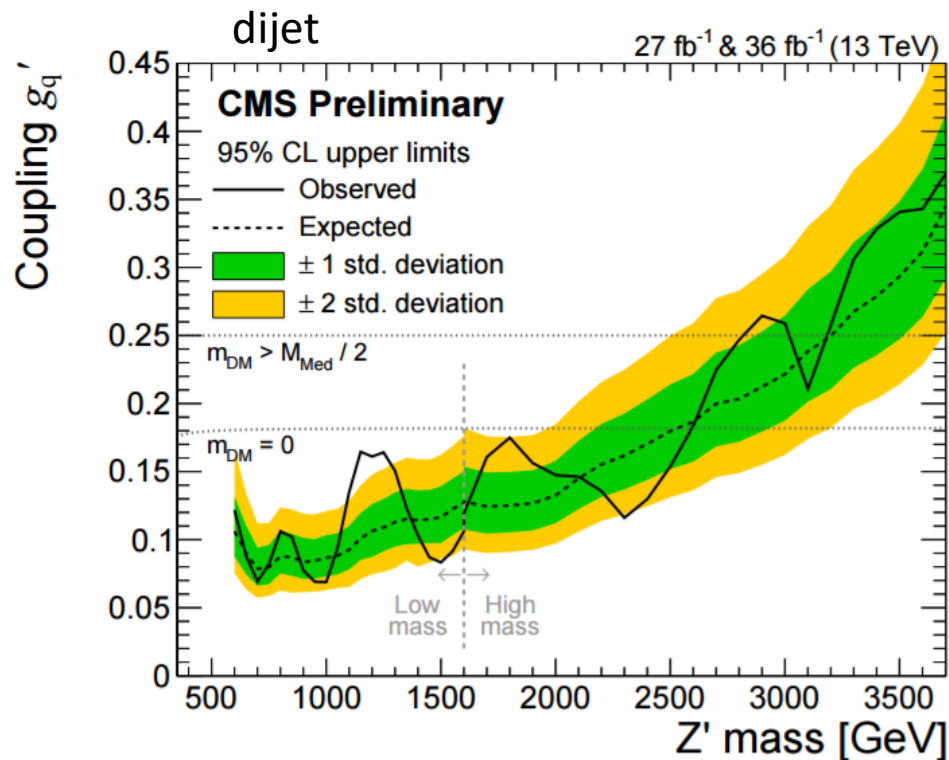
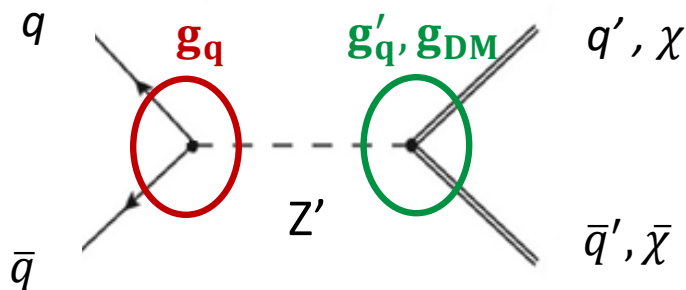
Still entering the high luminosity era

- other 100 fb⁻¹ expected in 2017-18

Limits scaling with couplings not trivial

- coupling → mediator width → kinematics

Perform scan in both mass and couplings phase space



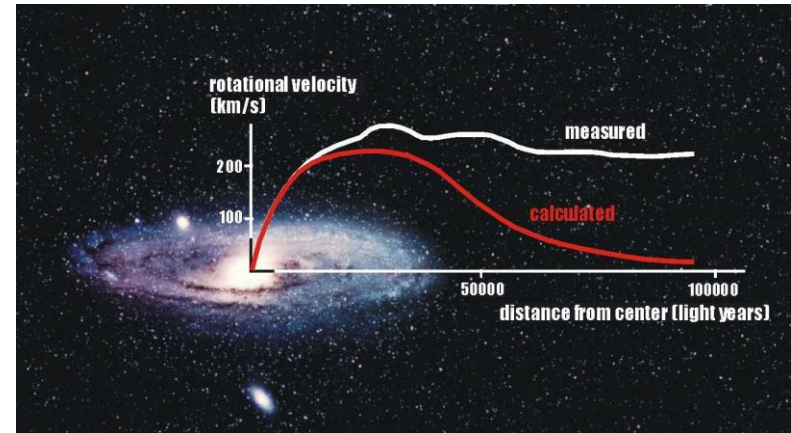
Summary and outlook

- **CMS has an extensive physics program focused on dark matter**
 - **searches performed in all possible final states**
 - monoJet/V, monoHiggs, SingleTop/ $t\bar{t}$ +DM presented
 - probe many theoretical scenarios
 - **no evidence for DM production yet**
 - **current public results mainly based on 2015 or ICHEP 2016 datasets**
 - **many analyses (top,Higgs) still limited by low statistics**
 - a lot to gain from an always larger datasets
- **generally speaking, LHC DM physics program still has a lot of potential**
 - energy bound (almost) reached, but **high luminosity era approaching**
 - possibility to perform an unprecedented **scan in mass-coupling phase space**
 - opportunity to **probe unexplored physics scenarios**

BACKUP

What do we know about dark matter?

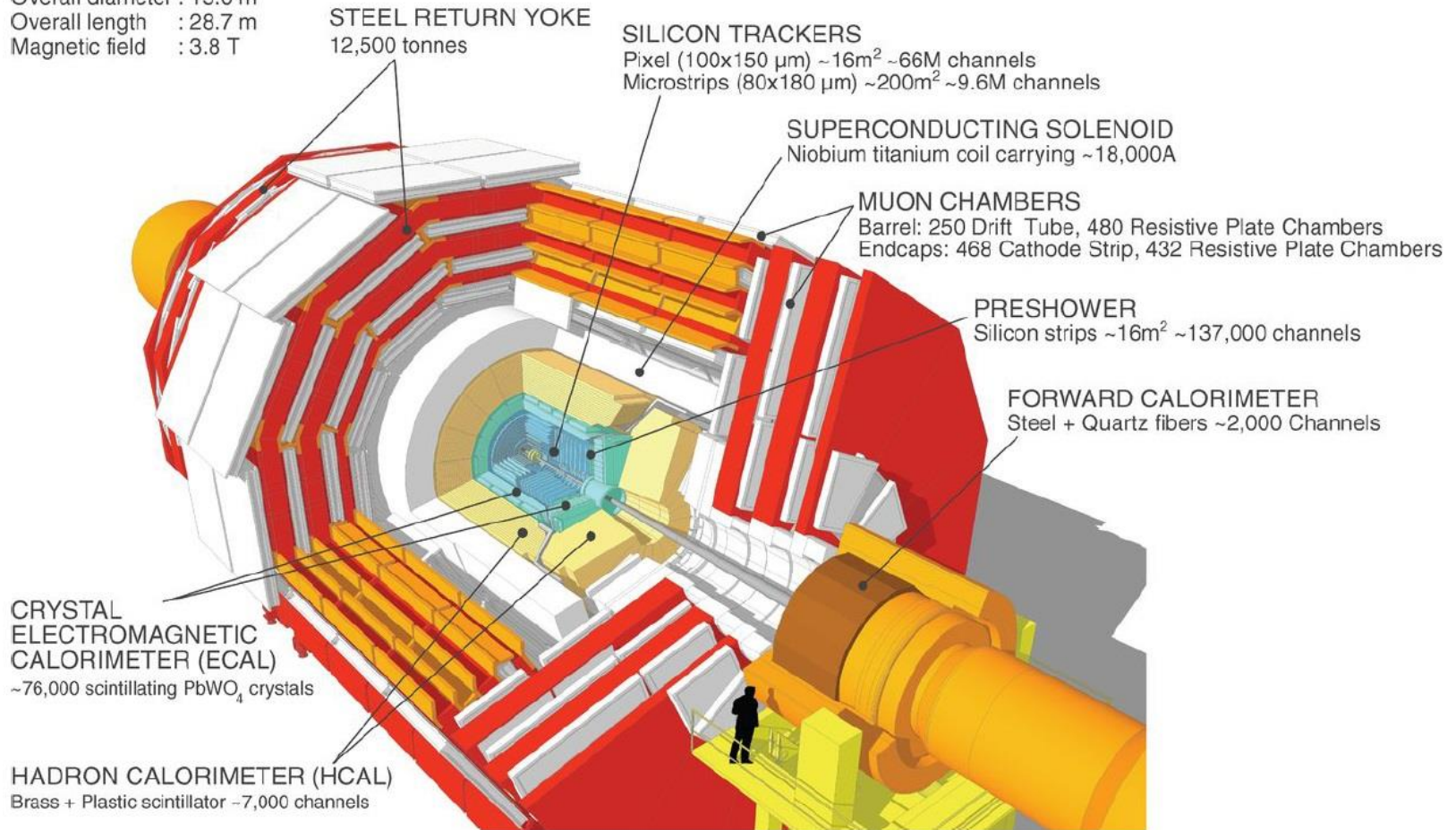
- ❑ many evidences of dark matter (DM) from **astrophysical observations**
- ❑ **particle nature of DM** a **big assumption**, though very natural
- ❑ properties of DM:
 - **gravitational influence** on ordinary matter
 - **neutral** under electromagnetic or strong interaction
 - **stable** on universe lifetime scale
- ❑ assume **DM interacts weakly with Standard Model (SM) particles**
 - Weakly Interacting Massive Particle (WIMP)



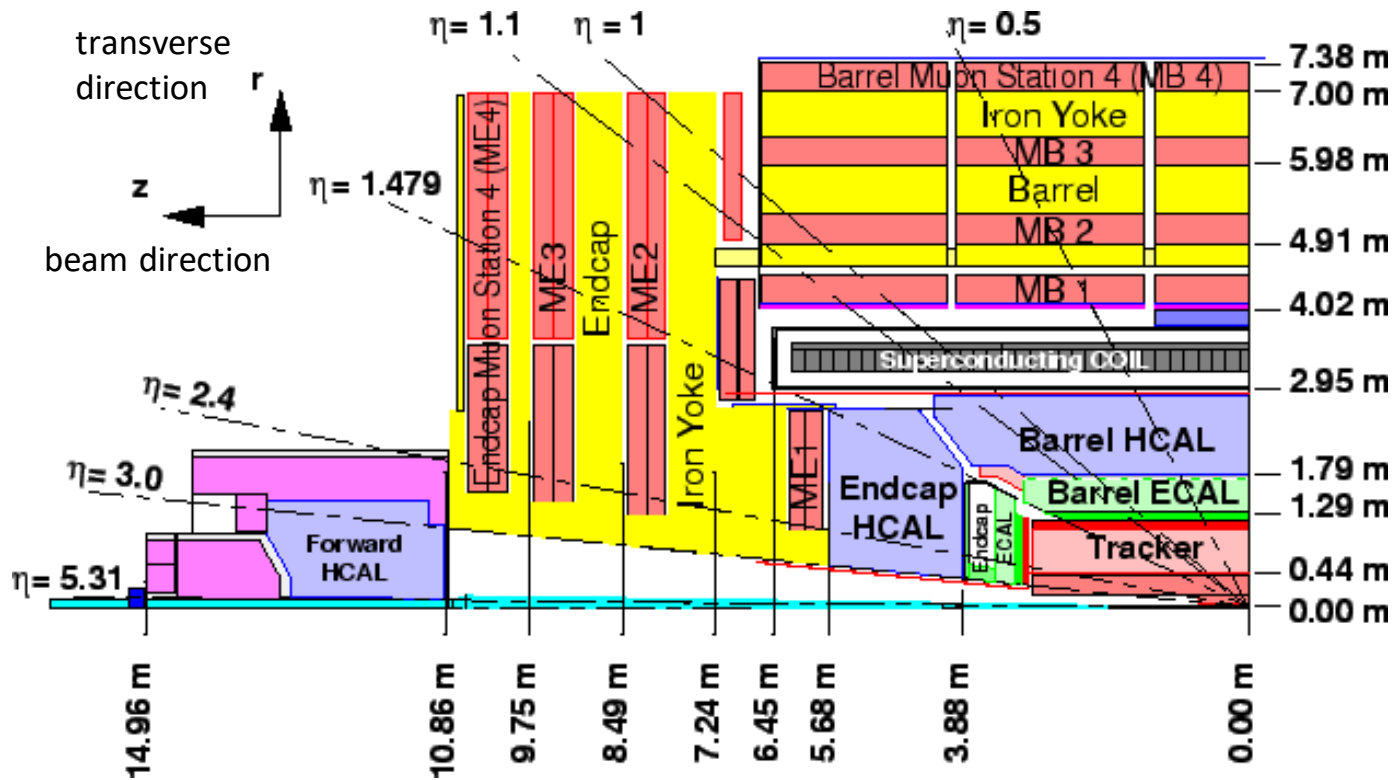
COMPACT MUON SOLENOID

CMS DETECTOR

Total weight : 14,000 tonnes
Overall diameter : 15.0 m
Overall length : 28.7 m
Magnetic field : 3.8 T



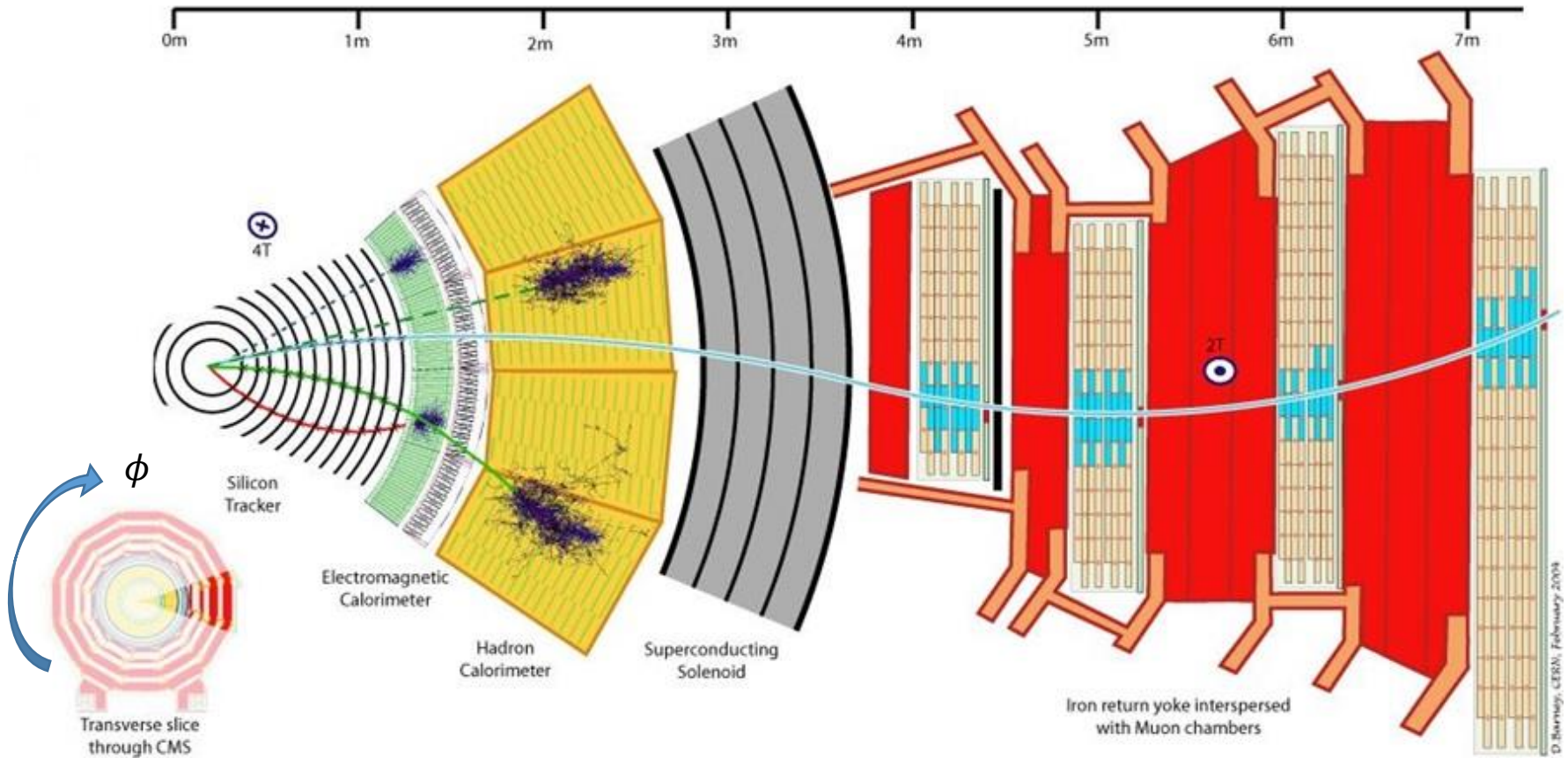
CMS longitudinal view



$$\eta = -\ln \operatorname{tg} \frac{\theta}{2}$$

η differences are Lorentz invariant for high energy particles

Particles reconstruction in CMS



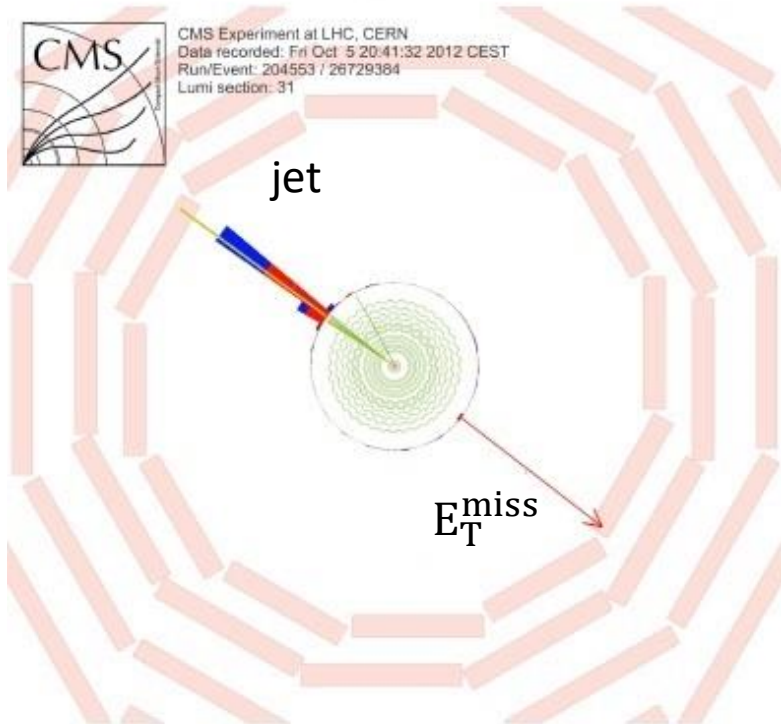
- muon
- electron
- charged hadron (e.g. π^+)
- - - neutral hadron (e.g. neutron)
- - - photon

Particle flow algorithm:

- combine information from all subdetectors to build particles

Seeing the invisible

event candidate for monoJet



- **DM does not interact** with the detector
- observed through **momentum imbalance** in transverse plane
- **single DM particle kinematics not accessible**, only the pair total transverse momentum (i.e. the mediator \vec{p}_T)
- **irreducible backgrounds** due to SM processes involving **neutrinos**

$$\vec{E}_T^{\text{miss}} = - \sum_{\text{visible particles}} \vec{p}_T$$

\vec{p}_T : momentum in transverse plane

Theory

Lagrangian for spin 1 mediator

arXiv:1603.04156v1

$$\mathcal{L}_{\text{vector}} = -g_{\text{DM}} Z'_{\mu} \bar{\chi} \gamma^{\mu} \chi - g_q \sum_{q=u,d,s,c,b,t} Z'_{\mu} \bar{q} \gamma^{\mu} q ,$$

$$\mathcal{L}_{\text{axial-vector}} = -g_{\text{DM}} Z'_{\mu} \bar{\chi} \gamma^{\mu} \gamma_5 \chi - g_q \sum_{q=u,d,s,c,b,t} Z'_{\mu} \bar{q} \gamma^{\mu} \gamma_5 q$$

Partial widths

$$\Gamma_{\text{vector}}^{\chi\bar{\chi}} = \frac{g_{\text{DM}}^2 M_{\text{med}}}{12\pi} (1 - 4z_{\text{DM}})^{1/2} (1 + 2z_{\text{DM}}) ,$$

$$\Gamma_{\text{vector}}^{q\bar{q}} = \frac{g_q^2 M_{\text{med}}}{4\pi} (1 - 4z_q)^{1/2} (1 + 2z_q) ,$$

$$\Gamma_{\text{axial-vector}}^{\chi\bar{\chi}} = \frac{g_{\text{DM}}^2 M_{\text{med}}}{12\pi} (1 - 4z_{\text{DM}})^{3/2} ,$$

$$\Gamma_{\text{axial-vector}}^{q\bar{q}} = \frac{g_q^2 M_{\text{med}}}{4\pi} (1 - 4z_q)^{3/2} .$$

$$z_{\text{DM},q} = m_{\text{DM},q}^2 / M_{\text{med}}^2$$

Theory

Lagrangian for spin 0 mediator

arXiv:1603.04156v1

$$\mathcal{L}_{\text{scalar}} = -g_{\text{DM}} \phi \bar{\chi} \chi - g_q \frac{\phi}{\sqrt{2}} \sum_{q=u,d,s,c,b,t} y_q \bar{q} q,$$

$$\mathcal{L}_{\text{pseudo-scalar}} = -ig_{\text{DM}} \phi \bar{\chi} \gamma_5 \chi - ig_q \frac{\phi}{\sqrt{2}} \sum_{q=u,d,s,c,b,t} y_q \bar{q} \gamma_5 q,$$

Partial widths

$$\Gamma_{\text{scalar}}^{\chi\bar{\chi}} = \frac{g_{\text{DM}}^2 M_{\text{med}}}{8\pi} (1 - 4z_{\text{DM}}^2)^{3/2},$$

$$\Gamma_{\text{scalar}}^{q\bar{q}} = \frac{3g_q^2 y_q^2 M_{\text{med}}}{16\pi} (1 - 4z_q^2)^{3/2},$$

$$\Gamma_{\text{scalar}}^{gg} = \frac{\alpha_s^2 g_q^2 M_{\text{med}}^3}{32\pi^3 v^2} |f_{\text{scalar}}(4z_t)|^2,$$

$$\Gamma_{\text{pseudo-scalar}}^{\chi\bar{\chi}} = \frac{g_{\text{DM}}^2 M_{\text{med}}}{8\pi} (1 - 4z_{\text{DM}}^2)^{1/2},$$

$$\Gamma_{\text{pseudo-scalar}}^{q\bar{q}} = \frac{3g_q^2 y_q^2 M_{\text{med}}}{16\pi} (1 - 4z_q^2)^{1/2},$$

$$\Gamma_{\text{pseudo-scalar}}^{gg} = \frac{\alpha_s^2 g_q^2 M_{\text{med}}^3}{32\pi^3 v^2} |f_{\text{pseudo-scalar}}(4z_t)|^2.$$

$$z_{\text{DM},q} = m_{\text{DM},q}^2 / M_{\text{med}}^2$$

$$y_q = \sqrt{2} m_q / v$$

$$v \simeq 246 \text{ GeV}$$

$$f_{\text{scalar}}(\tau) = \tau \left[1 + (1 - \tau) \arctan^2 \left(\frac{1}{\sqrt{\tau - 1}} \right) \right],$$

$$f_{\text{pseudo-scalar}}(\tau) = \tau \arctan^2 \left(\frac{1}{\sqrt{\tau - 1}} \right).$$

Comparison to indirect searches

- **vector, scalar** mediator \rightarrow **spin independent** (SI) cross section
- **axial-vector, pseudoscalar** mediator \rightarrow **spin dependent** (SD) cross section

DM-nucleon scattering cross section

$$\sigma_{\text{SI}} = \frac{f^2(g_q)g_{\text{DM}}^2\mu_{n\chi}^2}{\pi M_{\text{med}}^4}, \quad \mu_{n\chi} = m_n m_{\text{DM}} / (m_n + m_{\text{DM}}) \quad m_n \simeq 0.939 \text{ GeV}$$

vector: $\sigma_{\text{SI}} \simeq 6.9 \times 10^{-41} \text{ cm}^2 \cdot \left(\frac{g_q g_{\text{DM}}}{0.25}\right)^2 \left(\frac{1 \text{ TeV}}{M_{\text{med}}}\right)^4 \left(\frac{\mu_{n\chi}}{1 \text{ GeV}}\right)^2$

scalar: $\sigma_{\text{SI}} \simeq 6.9 \times 10^{-43} \text{ cm}^2 \cdot \left(\frac{g_q g_{\text{DM}}}{1}\right)^2 \left(\frac{125 \text{ GeV}}{M_{\text{med}}}\right)^4 \left(\frac{\mu_{n\chi}}{1 \text{ GeV}}\right)^2$

$$\sigma_{\text{SD}} = \frac{3f^2(g_q)g_{\text{DM}}^2\mu_{n\chi}^2}{\pi M_{\text{med}}^4}.$$

axial-vector: $\sigma^{\text{SD}} \simeq 2.4 \times 10^{-42} \text{ cm}^2 \cdot \left(\frac{g_q g_{\text{DM}}}{0.25}\right)^2 \left(\frac{1 \text{ TeV}}{M_{\text{med}}}\right)^4 \left(\frac{\mu_{n\chi}}{1 \text{ GeV}}\right)^2.$

pseudoscalar: highly suppressed

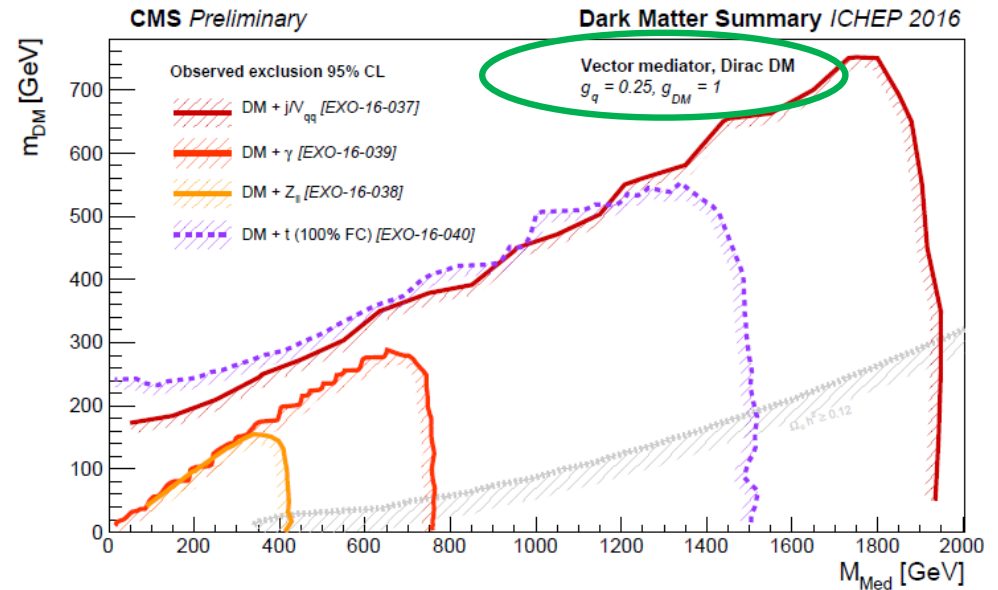
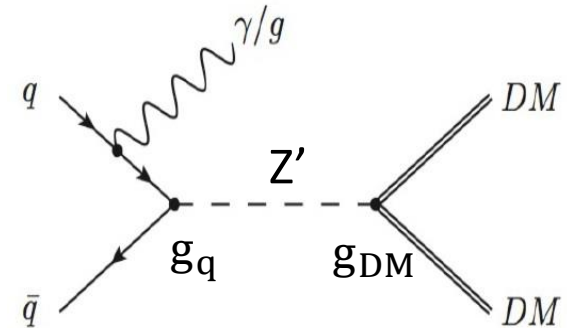
Theoretical overview

Interpretation with simplified models

- ❑ Dark Matter Forum prescriptions → arXiv:1507.00966
- ❑ benchmark of Run2 interpretation
- ❑ new mediator connecting SM and DM
- ❑ free parameters: m_{DM} , M_{med} , g_{DM} , g_q

Assumptions:

- DM is a Dirac fermion
- DM produced on-shell in pairs
- minimal decay width for mediator
- minimal flavour violation
- $g_{DM} = 1$ and $g_q = 0.25$



limits strongly depends on the couplings choice and model

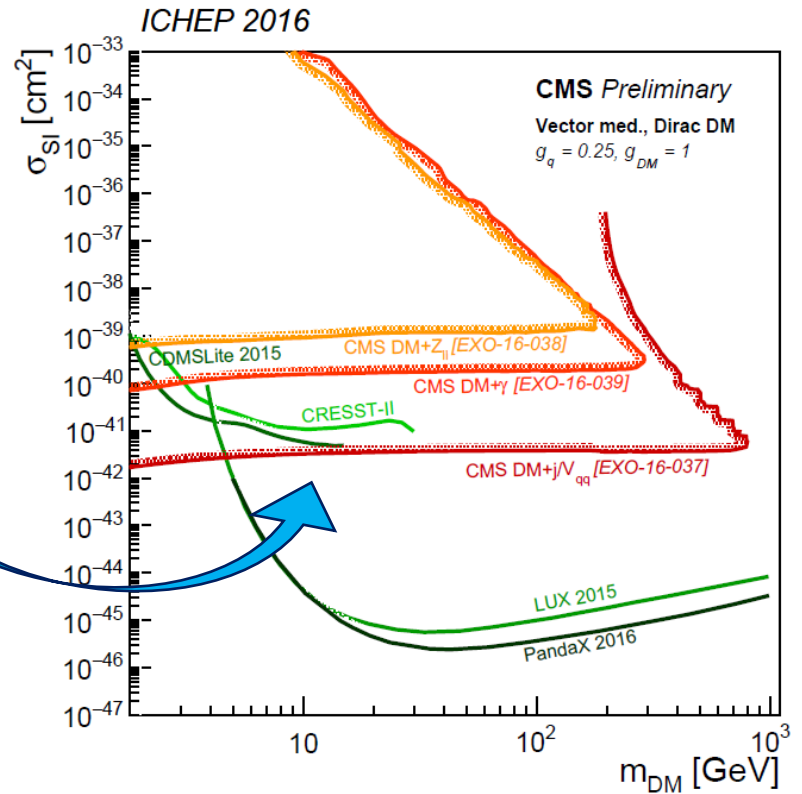
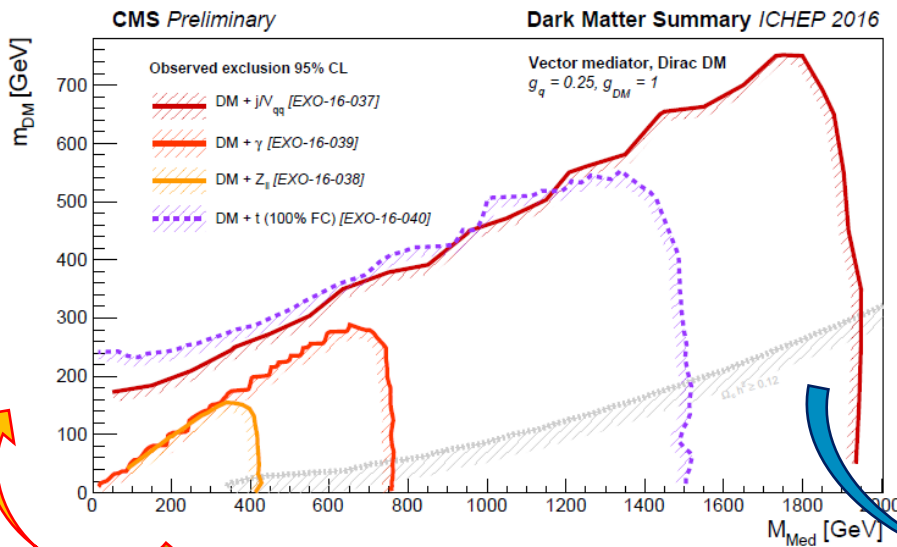
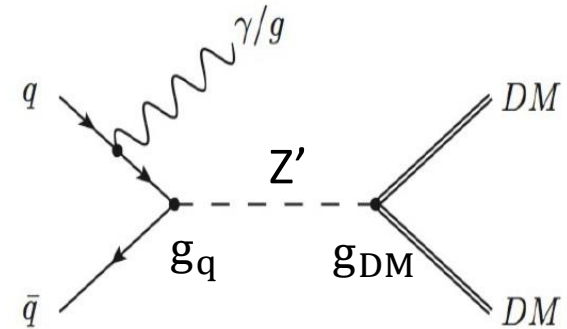
→ change in couplings affects mediator's width

→ more details: arXiv:1603.04156v1

Theoretical overview

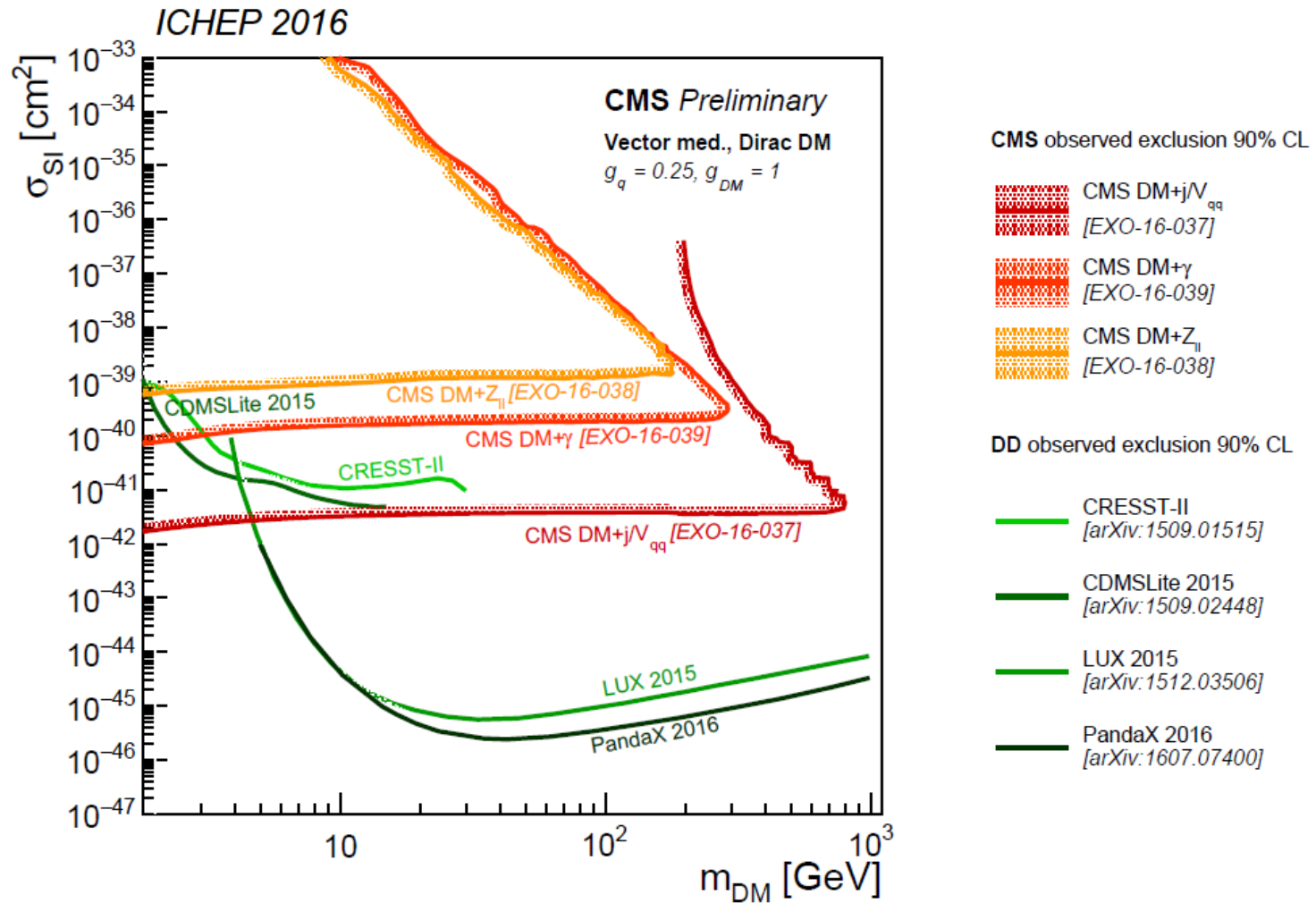
Limits typically @ 95% CL in the m_{DM}, M_{med} plane

- switch to limits on cross section as a function of m_{DM}
- 90% CL used to compare with direct searches



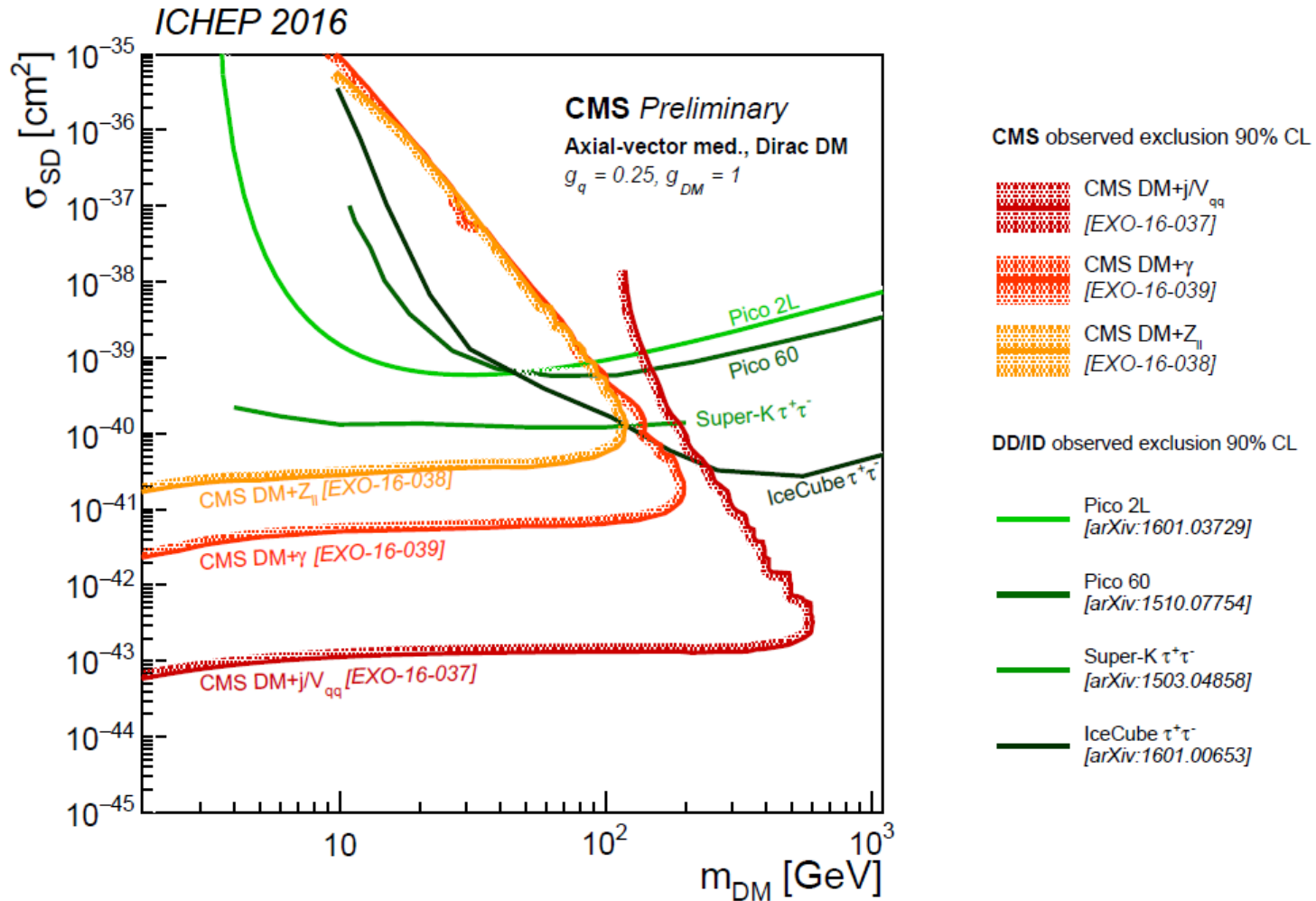
collider searches limits flat with respect to m_{DM} and sensitive also to low m_{DM}

ICHEP summary plots



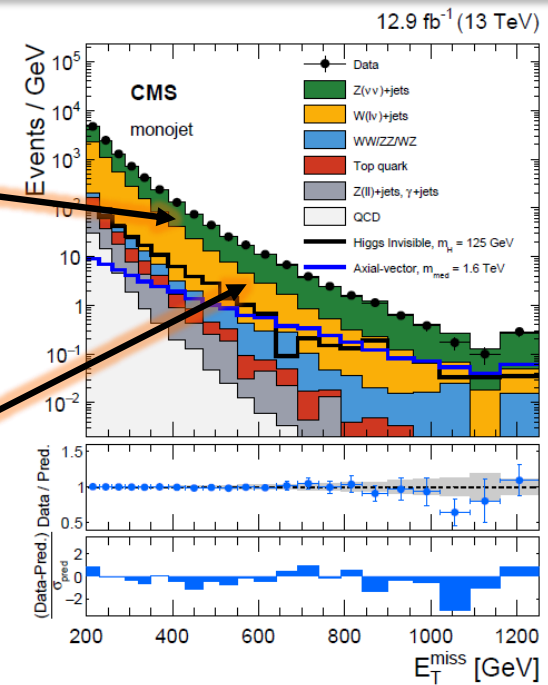
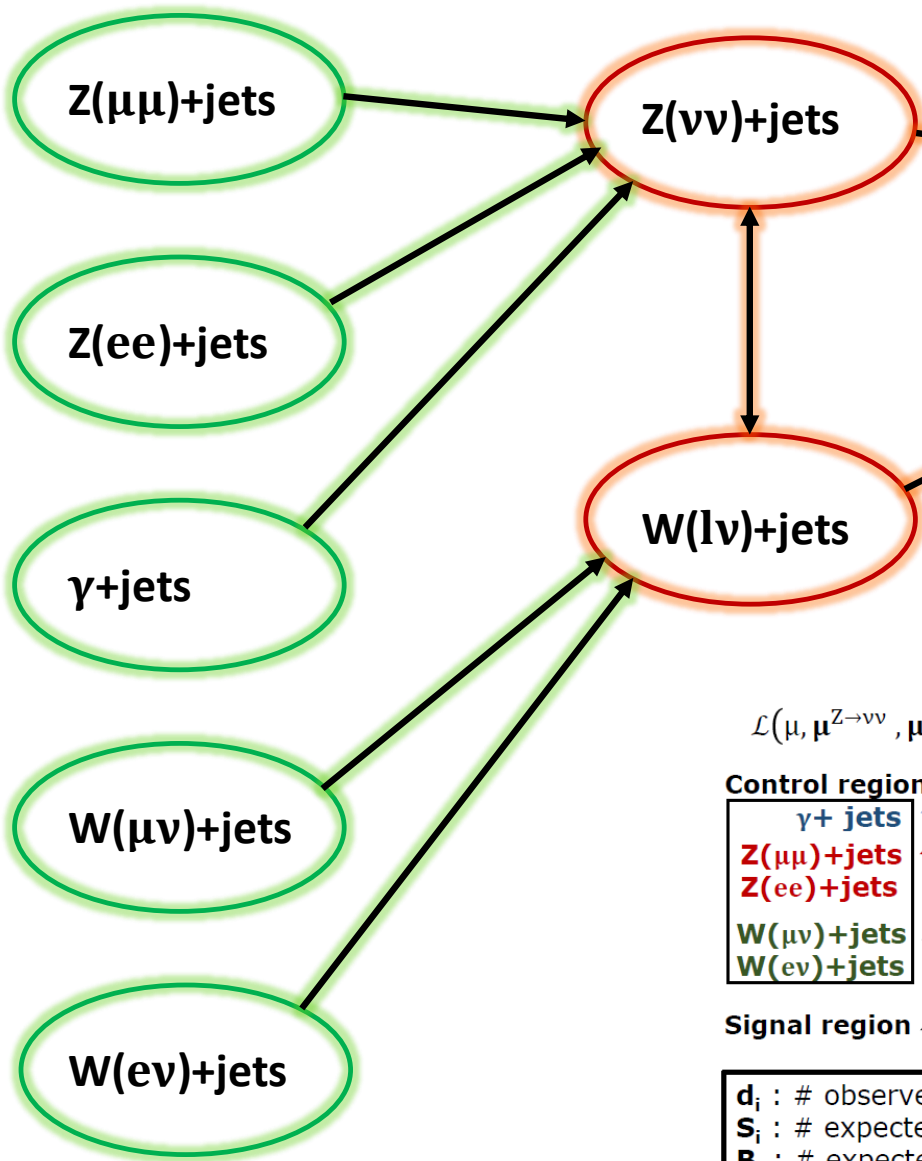
monojet most sensitive channel for vector mediator
direct searches more sensitive than collider searches for $m_{DM} > \text{few GeV}$

ICHEP summary plots



monojet most sensitive channel for axial-vector mediator
collider searches more sensitive than direct searches everywhere

MonoJet background estimate



$$\mathcal{L}(\mu, \mu^{Z \rightarrow \nu\nu}, \mu^{W \rightarrow l\nu}, \theta) = \prod_i \text{Poisson} \left(d_i^Y \mid B_i^Y(\theta) + \frac{\mu_i^{Z \rightarrow \nu\nu}}{R_i^Y(\theta)} \right) \times \prod_i \text{Poisson} \left(d_i^Z \mid B_i^Z(\theta) + \frac{\mu_i^{Z \rightarrow \nu\nu}}{R_i^Z(\theta)} \right) \times \prod_i \text{Poisson} \left(d_i^W \mid B_i^W(\theta) + \frac{\mu_i^{W \rightarrow l\nu}}{R_i^W(\theta)} \right) \times \prod_i \text{Poisson} \left(d_i \mid B_i(\theta) + \mu_i^{W \rightarrow l\nu} + \mu_i^{Z \rightarrow \nu\nu} + \mu S_i(\theta) \right)$$

exp. events for Z(vν):
 # exp. events for W(lν)

- Control regions**
- γ + jets
 - Z(μμ)+jets
 - Z(ee)+jets
 - W(μν)+jets
 - W(eν)+jets
- Signal region**

d_i : # observed events in i-th bin
 S_i : # expected signal events
 B_i : # expected background events (MC)
 R_i : scale factors from CR to SR

μ : signal strength
 θ : nuisance parameters (systematic uncertainties)

MonoJet event selection

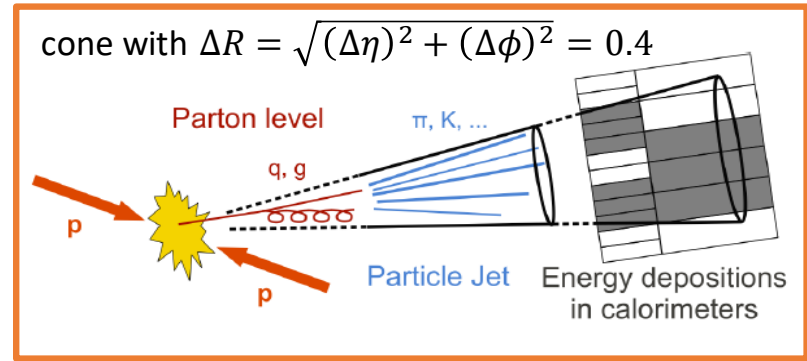
Common selection (both monoJet/V)

- ✓ MET > 200 GeV (recoil)
- ✓ leading **AK4 jet**: $p_T > 100$ GeV, $|\eta| < 2.5$

Background suppression

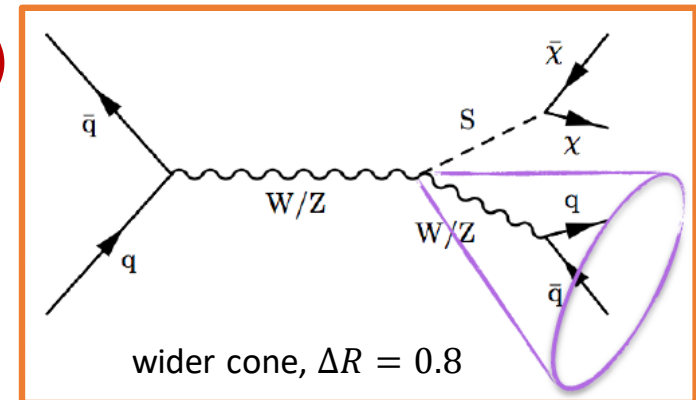
- ✓ leading jet noise cleaning
 - ✓ $\Delta\phi(\text{jets}, \text{MET}) > 0.5$
 - ✓ veto on photons and charged leptons
 - ✓ veto on b-jets
- beam backgrounds
→ QCD
→ EWK backgrounds
→ top

< 10 % of total bkg
after selection



Specific selection for monoV (boosted topology)

- ✓ leading **AK8 jet**: $p_T > 250$ GeV, $|\eta| < 2.4$
- ✓ MET > 250 GeV
- ✓ **V-tagging** techniques
 - boson pruned mass in [65,105] GeV
 - N-subjettiness $\tau_2 / \tau_1 < 0.6$



N-subjettiness

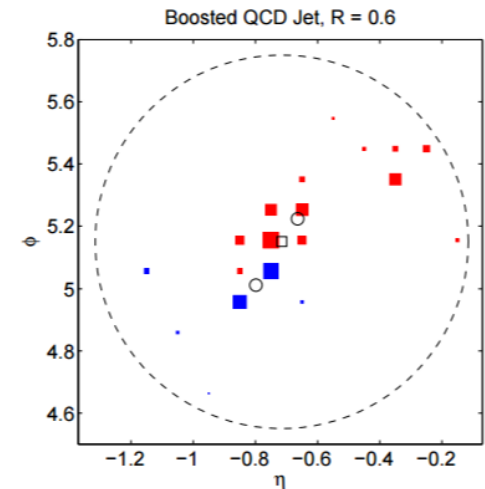
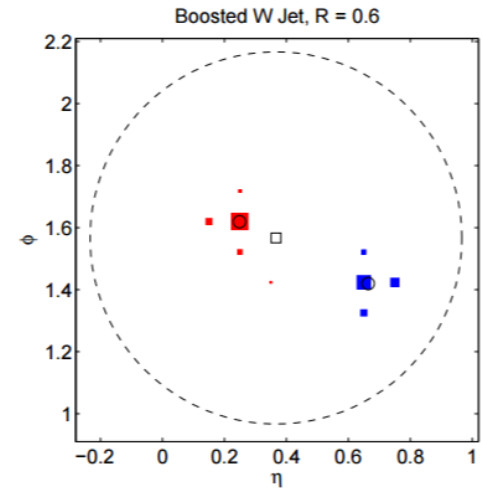
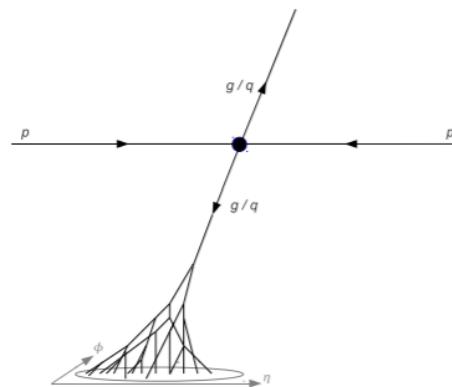
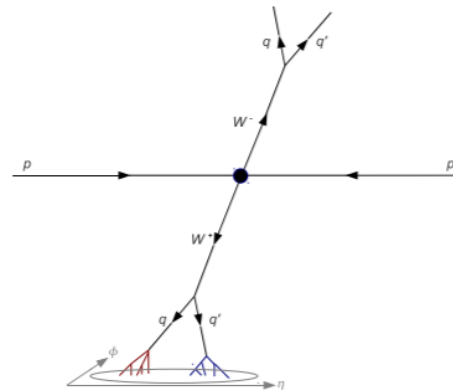
k runs over the constituent particles in a given jet, $p_{T,k}$ are their transverse momenta and $\Delta R_{J,k}$ is the distance in rapidity-azimuth plane between a candidate subjet J and a constituent particle k

$$\tau_N = \frac{1}{d_0} \sum_k p_{T,k} \min \{ \Delta R_{1,k}, \Delta R_{2,k}, \dots, \Delta R_{N,k} \}$$

$$\Delta R_{J,k} = \sqrt{(\Delta\eta)^2 + (\Delta\phi)^2}$$

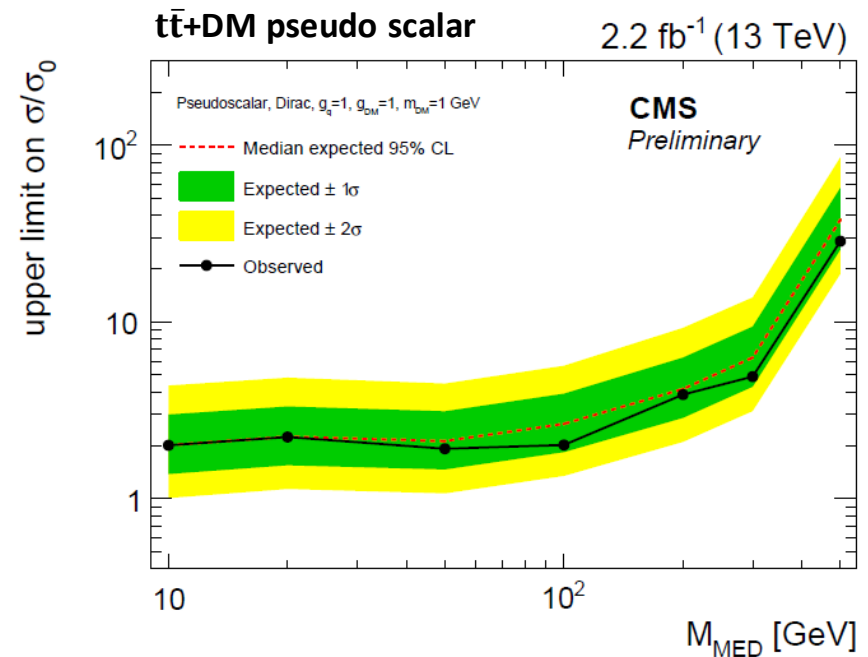
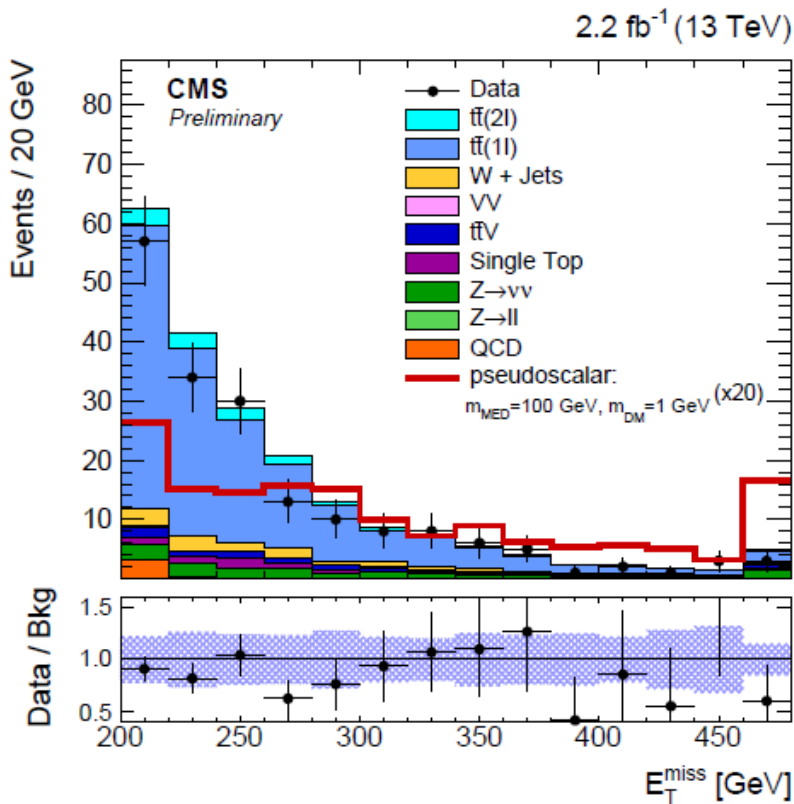
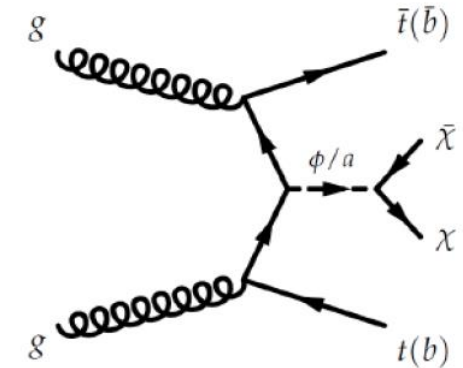
$$d_0 = \sum_k p_{T,k} R_0$$

R_0 : the jet radius used in the original jet clustering algorithm

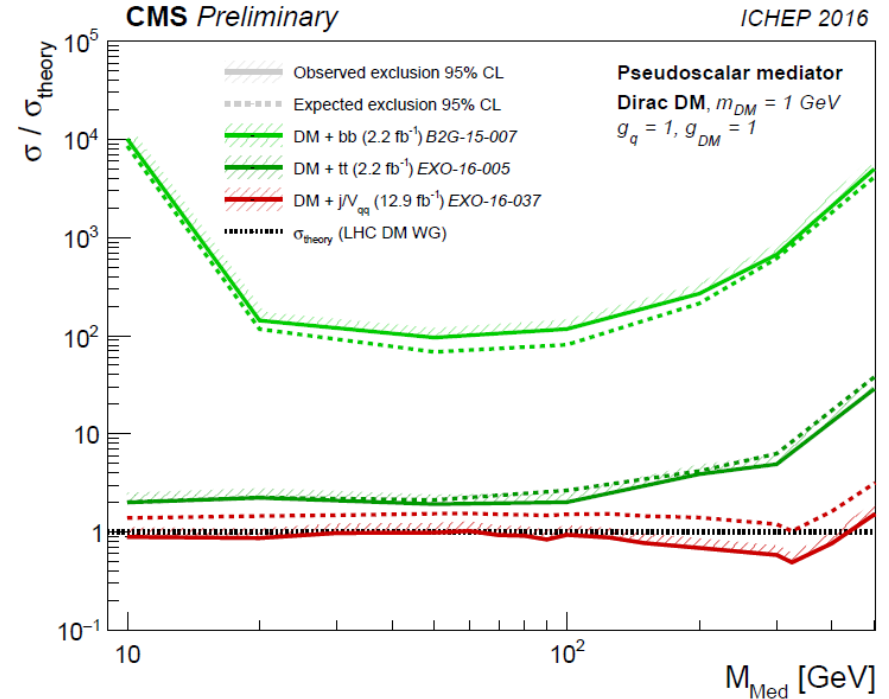
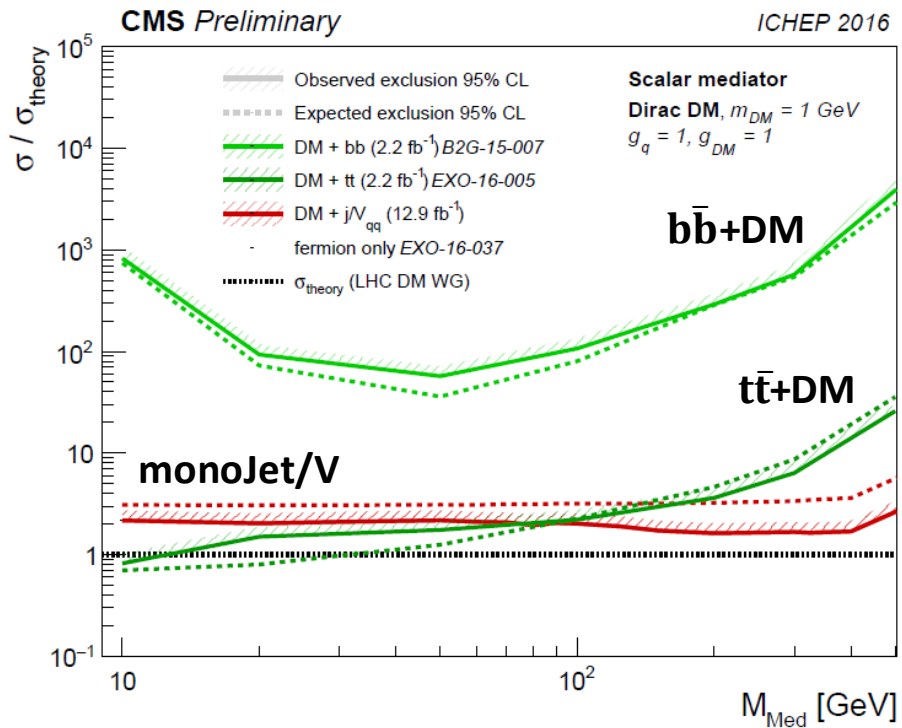


$t\bar{t}$ + Dark Matter

- **top tagging techniques:** identify top quarks decaying into three resolved jets
- **hadronic or semileptonic top decays** considered
- data-driven estimate of **SM $t\bar{t}$ and W/Z+jets backgrounds**



Comparison of Monojet and $t\bar{t}$ +DM



monoJet and **$t\bar{t}$ +DM** most sensitive channels for (pseudo)scalar mediator

$t\bar{t}$ +DM has a **better S/B** ratio (note the different luminosity in the plot)

direct searches have much less or no sensitivity to this case

Probing Dark Matter with MonoHiggs

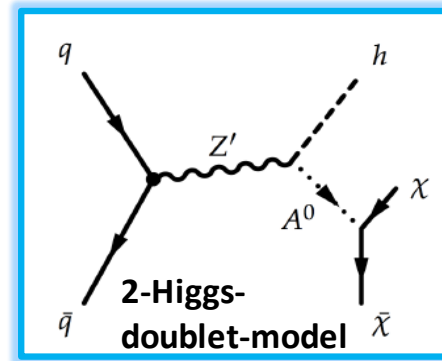
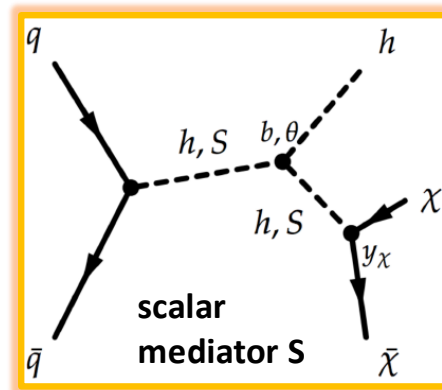
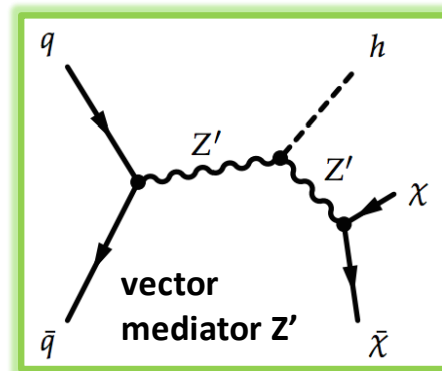
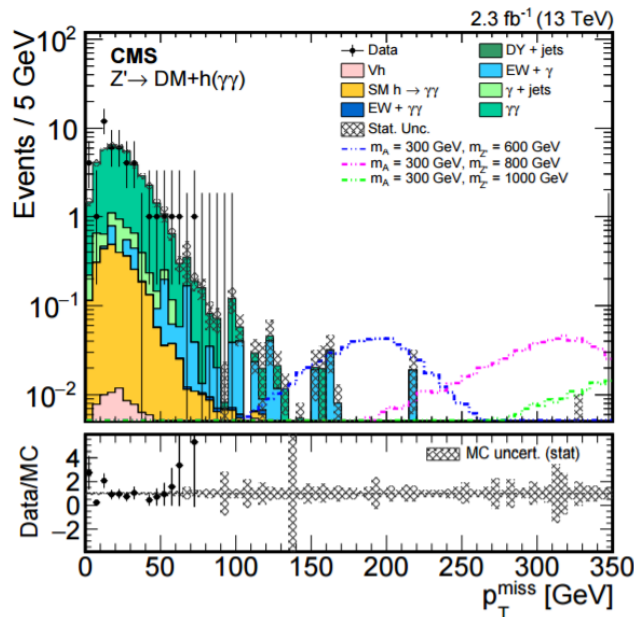
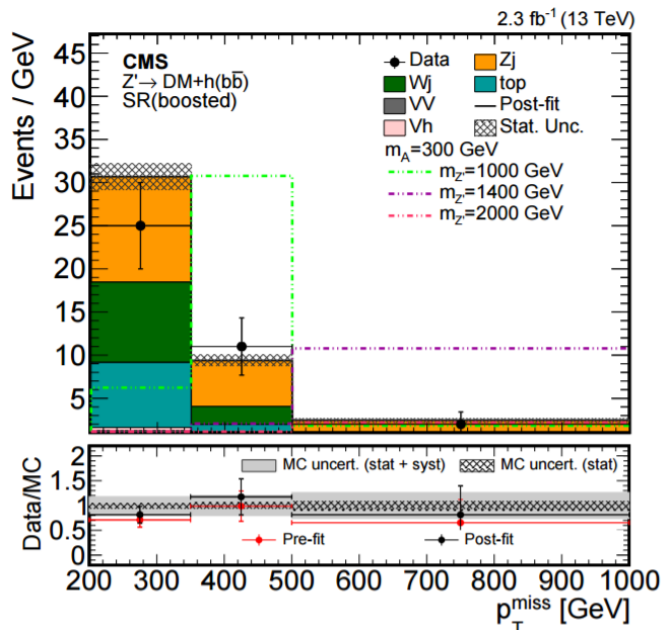
- **Higgs ISR suppressed** (coupling \propto mass)
- search for MET+ SM Higgs \rightarrow $bb, \gamma\gamma, ZZ, \tau\bar{\tau} \dots$

➤ $H \rightarrow b\bar{b}$

- **b-tagging** (resolved and boosted category)

➤ $H \rightarrow \gamma\gamma$

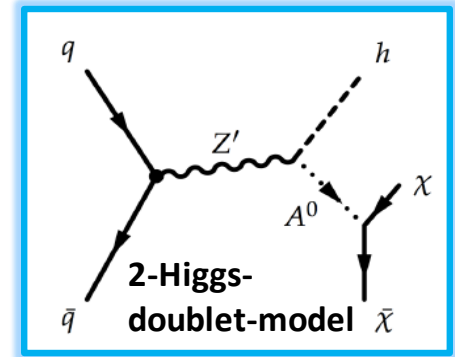
- **SM Higgs** included as **resonant background**
- non resonant bkg from fit to $m_{\gamma\gamma}$ in low- p_T^{miss} CR



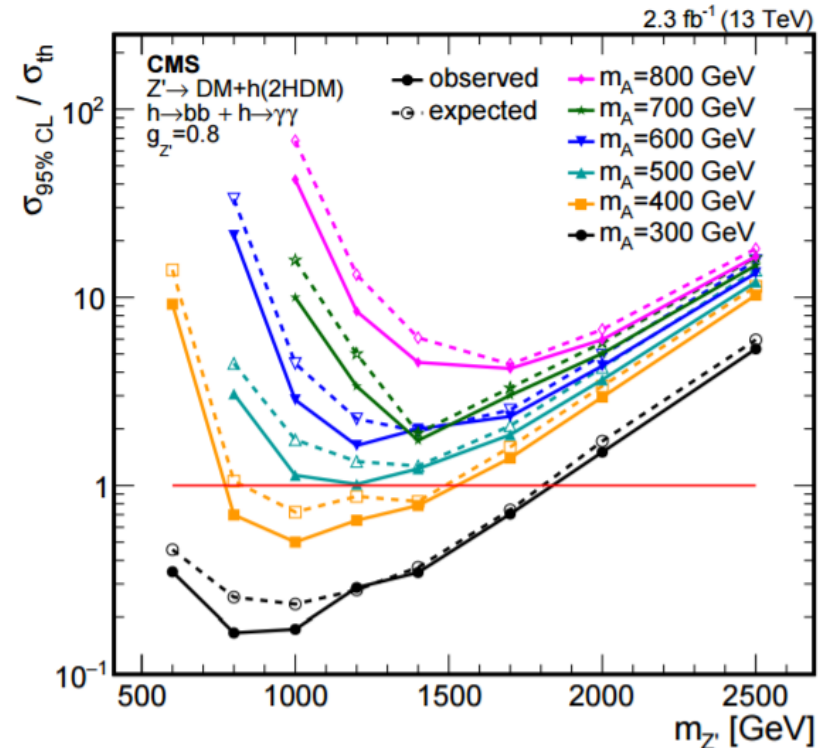
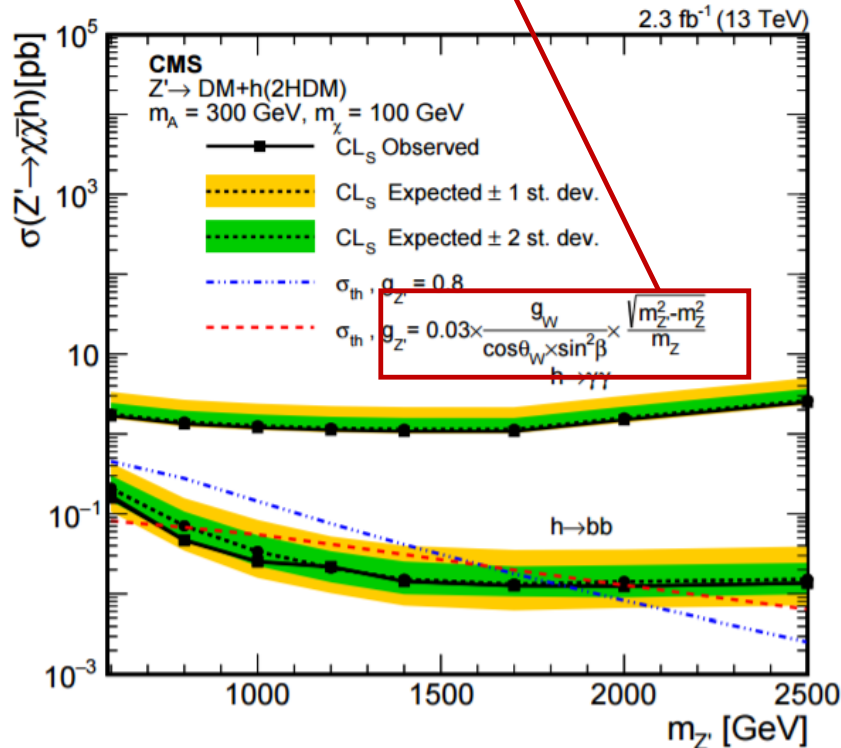
Mono-Higgs search results

combination of $b\bar{b}, \gamma\gamma$ (2.3 fb^{-1} @ 13TeV) submitted to JHEP

- Z' -2HDM interpretation
- $m_{Z'} < 1860 \text{ GeV}$ for $g_{Z'} = 0.8$ and $m_A = 300 \text{ GeV}$
- $m_{Z'} < 2040 \text{ GeV}$ for constrained $g_{Z'}$ and $m_A = 300 \text{ GeV}$

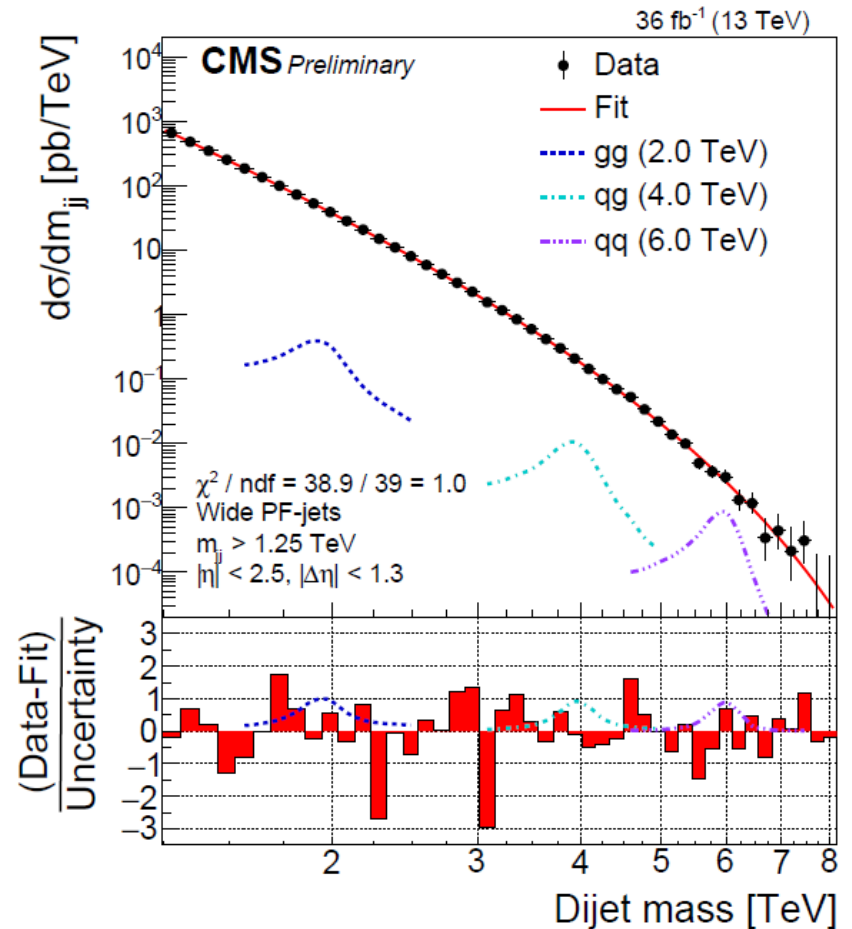
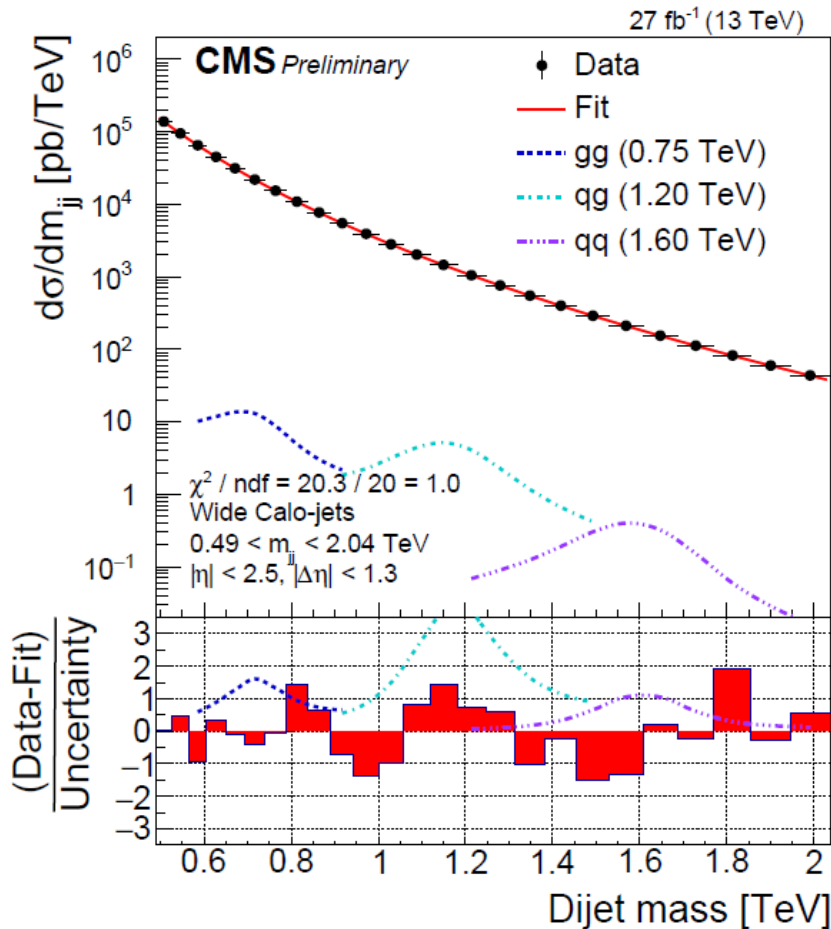
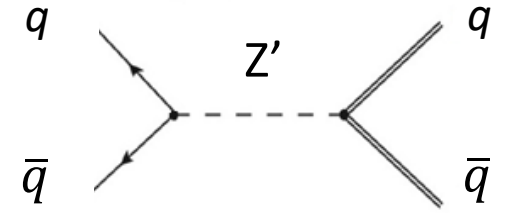


$g'_{Z'}$ constrained from global EWK fit and dijet searches



Dijet search

- **most powerful search at a hadron collider**
- search for a bump over falling dijet mass spectrum
- low mass and high mass search



Dijet search

- constrain on many models predicting resonances decaying into two jets
- constrain also simplified models with **(axial) vector mediator coupling only to quarks and DM**

