

FLAVOUR PHYSICS IN THE LHC ERA

Alexander Lenz Pheno 2017 8.5.2017

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- Message 5: SM/CKM dominance gives important bounds on BSM models
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- Message 8: All relevant observables have to be identified for BSM searches
- Message 9: Alternative paths for BSM searches can be interesting

INTRODUCTION

There are (at least) six kinds (=flavours) of quarks

$$\begin{pmatrix} u \\ d \end{pmatrix} \begin{pmatrix} c \\ s \end{pmatrix} \begin{pmatrix} t \\ b \end{pmatrix} \begin{pmatrix} q = +2/3 \\ q = -1/3 \end{pmatrix}$$

• Proton $p = |uud\rangle$

■ (Heavy) Flavour Physics describes hadrons with a charm- or a bottom-quark

	$D^0 = (\bar{u}c)$	$D^+ = (\bar{d}c)$	$D_s^+ = (\bar{s}c)$	$\Lambda_c = (udc)$
Mass (GeV)	1.86486	1.86962	1.96850	2.28646
Lifetime (ps)	0.4101	1.040	0.500	0.200

	$B_d = (\overline{b}d)$	$B^+ = (ar{b}u)$	$B_s = (ar{b}s)$	$B_c^+ = (\bar{b}c)$	$egin{array}{c} \Lambda_b = (udb) \end{array}$
Mass (GeV)	5.27958	5.27926	5.3667	6.2745	5.6194
Lifetime(ps)	1.519	1.638	1.512	0.500	1.451

NOT COVERED

Charm Physics

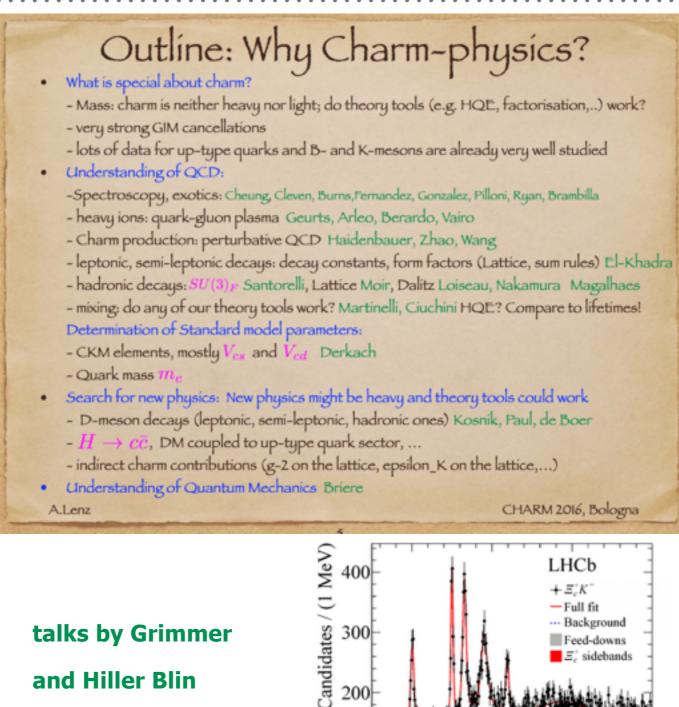
Spectroscopy, Production

Not one, not two, but five new particles discovered with the Large Hadron Collider

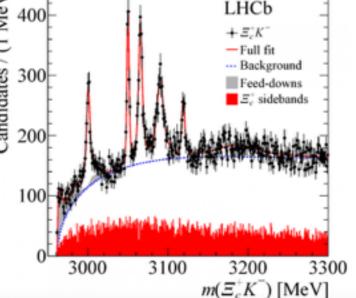
By Jessica Hall on March 24, 2017 at 8:41 am 41 Comments

🕈 😏 G+ 🥳





and Hiller Blin



MOTIVATION FOR FLAVOUR PHYSICS

Baryon Asymmetry in the Universe:

A violation of the CP symmetry - which causes matter and anti-matter to evolve differently with time - seems to be necessary to explain the existence of matter in the Universe.

CP violation has so far only been found in hadron decays, which are experimentally investigated at LHCb and NA62 (CERN), SuperBelle (Japan),...





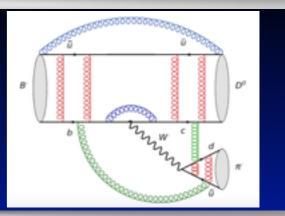
Indirect Search for BSM Physics:

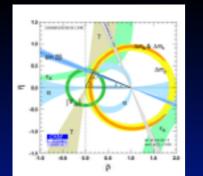
To find hints for Physics beyond the Standard Model we can either use brute force (= higher energies) or more subtle strategies like high precision measurements. New contributions to an observable f are identified via:

$$f^{\rm SM} + f^{\rm NP} = f^{\rm Exp}$$

Understanding QCD:

Hadron decays are strongly affected by QCD (strong interactions) effects, which tend to overshadow the interesting fundamental decay dynamics. Theory tools like **effective theories**, Heavy Quark Expansion, HQET, SCET,...enable a control over QCD-effects and they are used in other fields like Collider Physics, Higgs Physics, DM searches...





Standard Model parameters:

Hadron decays depend strongly on Standard Model parameters like quark masses and CKM couplings (which are the only known source of CP violation in the SM). A precise knowledge of these parameters is needed for all branches of particle physics.

- ► Huge experimental progress: B-factories, Tevatron and LHC
- LHCb: 462 papers
 20689 citations
 till 2016 5fb-1
 see/saw Uli Uwer
 Soeren Prell

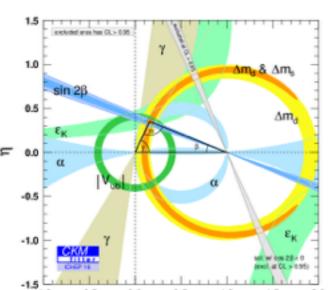


LHC

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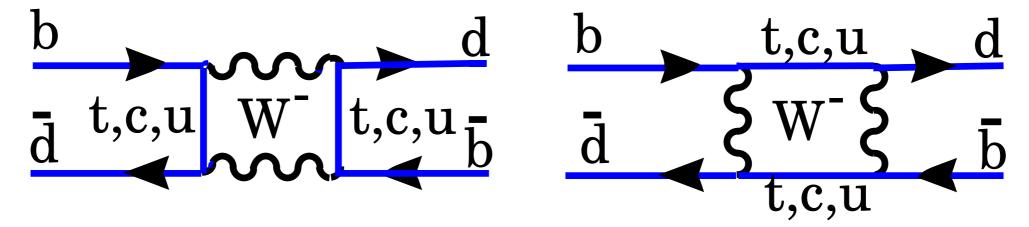
► Message 1: SM and CKM work perfectly



 $\beta^{\text{HFAG}} = 21.9^{\circ} \pm 0.7^{\circ} \text{ vs. } \beta^{\text{CKMfitter}} = 23.74^{\circ} _{-0.98^{\circ}}^{+1.13^{\circ}}$ $\gamma^{\text{HFAG}} = 71.3^{\circ} _{-6.1^{\circ}}^{+5.7^{\circ}} \text{ vs. } \gamma^{\text{CKMfitter}} = 65.33^{\circ} _{-2.54^{\circ}}^{+0.96^{\circ}}$

similar results from UTfit; Eigen et al.; Laiho et al

► Message 2: Many times: $\delta^{\text{Exp.}} < \delta^{\text{Theory}}$ e.g. B-mixing

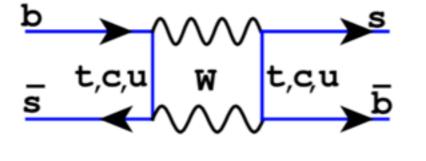


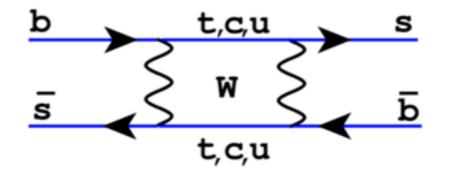
 $|M_{12}|$, $|\Gamma_{12}|$ and $\phi = \arg(-M_{12}/\Gamma_{12})$ can be related to three observables:

- Mass difference: $\Delta M := M_H M_L \approx 2|M_{12}|$ (off-shell) $|M_{12}|$: heavy internal particles: t, SUSY, ...
- Decay rate difference: $\Delta \Gamma := \Gamma_L \Gamma_H \approx 2|\Gamma_{12}| \cos \phi$ (on-shell) $|\Gamma_{12}|$: light internal particles: u, c, ... (almost) no NP!!!
- **Flavor specific/semi-leptonic CP asymmetries:** e.g. $B_q \rightarrow X l \nu$ (semi-leptonic)

$$a_{sl} \equiv a_{fs} = \frac{\Gamma(\overline{B}_q(t) \to f) - \Gamma(B_q(t) \to \overline{f})}{\Gamma(\overline{B}_q(t) \to f) + \Gamma(B_q(t) \to \overline{f})} = \left|\frac{\Gamma_{12}}{M_{12}}\right| \sin\phi$$

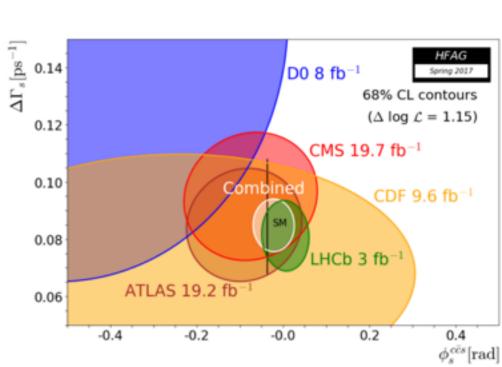
► Message 2: Many times: $\delta^{\text{Exp.}} < \delta^{\text{Theory}}$ e.g. B-mixing





 $a_{sl}^s =$

 $\Delta M_s = 2|M_{12}^s|, \quad \Delta \Gamma_s = 2|\Gamma_{12}^s|\cos\phi_{12}^s,$



 $\frac{\Gamma_{12}^s}{M_{12}^s}$

Preliminary -	 Moriond 	2017
Tim Gershon		

 $\sin\phi_{12}^s$

Observable	SM	Experiment
ΔM_s	$(18.3 \pm 2.7) \mathrm{ps^{-1}}$	$(17.757 \pm 0.021) \mathrm{ps^{-1}}$
$\Delta\Gamma_s$	$(0.088 \pm 0.020) { m ps^{-1}}$	$(0.082\pm0.006){ m ps^{-1}}$
a_{sl}^s	$(2.22\pm0.27)\cdot10^{-5}$	$(-750 \pm 410) \cdot 10^{-5}$
$\Delta \Gamma_s / \Delta M_s$	$48.1(1\pm0.173)\cdot10^{-4}$	$46.2(1\pm0.073)\cdot10^{-4}$

CP violation in the Bs system Marina Artuso, Guennadi Borissov, AL Rev.Mod.Phys. 88 (2016) no.4,045002

Message 3: Higher precision in theory needed

$\Delta\Gamma_s^{\mathrm{SM}}$	This work
Central value	0.088 ps ⁻¹
$\delta(B_{ ilde{R}_2})$	14.8%
$\delta(f_{B_s}\sqrt{B})$	13.9%
$\delta(\mu)$	8.4%
$\delta(V_{cb})$	4.9%
$\delta(\tilde{B}_S)$	2.1%
$\delta(B_{R_0})$	2.1%
$\delta(\bar{z})$	1.1%
$\delta(m_b)$	0.8%
$\delta(B_{ ilde{R}_1})$	0.7%
$\delta(B_{\tilde{R}_3})$	0.6%
$\delta(B_{R_1})$	0.5%
$\delta(B_{R_3})$	0.2%
$\delta(m_s)$	0.1%
$\delta(\gamma)$	0.1%
$\delta(\alpha_s)$	0.1%
$\delta(V_{ub}/V_{cb})$	0.1%
$\delta(\bar{m}_t(\bar{m}_t)$	0.0%
$\sum \delta$	22.8%

• Dim 7 has never been done

-Wingate works on lattice -Rauh, Kirk, AL with QCD sum rules

$$\langle R_2 \rangle = -\frac{2}{3} \left[\frac{M_{B_s}^2}{m_b^{\text{pow}2}} - 1 \right] M_{B_s}^2 f_{B_s}^2 B_{R_2}, \qquad R_2 = \frac{1}{m_b^2} \,\overline{s}_\alpha \overleftarrow{D}_\rho \gamma^\mu (1 - \gamma_5) D^\rho b_\alpha \,\overline{s}_\beta \gamma_\mu (1 - \gamma_5) b_\beta$$

• Dim 6 is done on the lattice - newest result:

B(s)-mixing matrix elements from lattice QCD for the Standard Model and beyond Fermilab Lattice and MILC Collaborations Phys.Rev. D93 (2016) no.11, 113016, arXiv:1602.03560 [hep-lat]

$$\langle Q \rangle \equiv \langle \bar{B}^0_s | Q | B^0_s \rangle = \frac{8}{3} M^2_{B^0_s} f^2_{B_s} B(\mu)$$

$$Q = \bar{s}^{\alpha} \gamma_{\mu} (1 - \gamma_5) b^{\alpha} \times \bar{s}^{\beta} \gamma^{\mu} (1 - \gamma_5) b^{\beta}$$

Also QCD sum rules: **B-mixing at NLO Grozin, Klein, Mannel, Pivovarov** Phys.Rev. D94 (2016) no.3, 034024, arXiv:1606.06054 [hep-ph]

indicates a small tension with experiment

First steps in NNLO-QCD

The phase space analysis for 3 and 4 massive particles in final states Asatrian, Hovhannisyan, Yeghiazaryan Phys.Rev. D86 (2012) 114023, arXiv:1210.7939 [hep-ph]

CP violation in the Bs system Marina Artuso, Guennadi Borissov, AL

Rev.Mod.Phys. 88 (2016) no.4,045002



High experimental precision requires to think again:

- Message 4: Standard assumptions/textbook wisdom might have to be re-considered
 - How large are penguins? How well does SU(3) F work?
 How large can quark hadron duality violation be?

 - How well does QCD-factorisation work? e.g. Bobeth, Gorbahn, Vickers 2014; Bell 2015;...
 - How large can BSM effects be in tree-level decays?

Message 5: SM/CKM dominance gives bounds on BSM models

Meson mixing:

e.g.Charles,Descotes-Genon,Ligeti,Monteil,Papucciand,Trabelsi

Phys.Rev.D89,no. 3, 033016 (2014) [arXiv:1309.2293].

Couplings	NP loop	Scales (in 7	TeV) probed by
Couplings	order	B_d mixing	B_s mixing
$ C_{ij} = V_{ti}V_{tj}^* $	tree level	17	19
(CKM-like)	one loop	1.4	1.5
$ C_{ij} = 1$	tree level	$2 imes 10^3$	5×10^2
(no hierarchy)	one loop	2×10^2	40

TEST OF UNDERLYING THEORY ASSUMPTIONS: DUALITY

1970 Blom, Gilman for e-p scattering
1979 Poggio, Quinn, Weinberg for e+e- to hadrons Basic idea: Sum overall hadrons = quark level
Our definition: duality violation is deviation from HQE

$$\Gamma = \Gamma_0 + \frac{\Lambda^2}{m_b^2}\Gamma_2 + \frac{\Lambda^3}{m_b^3}\Gamma_3 + \frac{\Lambda^4}{m_b^4}\Gamma_4 + \dots$$

Actual expansion parameter is momentum release $\frac{M_b}{M_i^2 - M_f^2}$ Taylor expansion of exp[-1/x] in x does give zero

Channel	Expansion parameter x	Numerical value	$\exp[-1/x]$
	$\frac{\Lambda}{\sqrt{m_b^2 - 4m_c^2}} \approx \frac{\Lambda}{m_b} \left(1 + 2\frac{m_c^2}{m_b^2} \right)$		
$b \to c \bar{u} s$	$\frac{\Lambda}{\sqrt{m_b^2 - m_c^2}} \approx \frac{\Lambda}{m_b} \left(1 + \frac{1}{2} \frac{m_c^2}{m_b^2} \right)$	0.045 - 0.49	$1.9 \cdot 10^{-10} - 0.13$
$b \to u \bar{u} s$	$\frac{\Lambda}{\sqrt{m_b^2 - 4m_u^2}} = \frac{\Lambda}{m_b}$	0.042 - 0.48	$4.2 \cdot 10^{-11} - 0.12$

Best candidate:

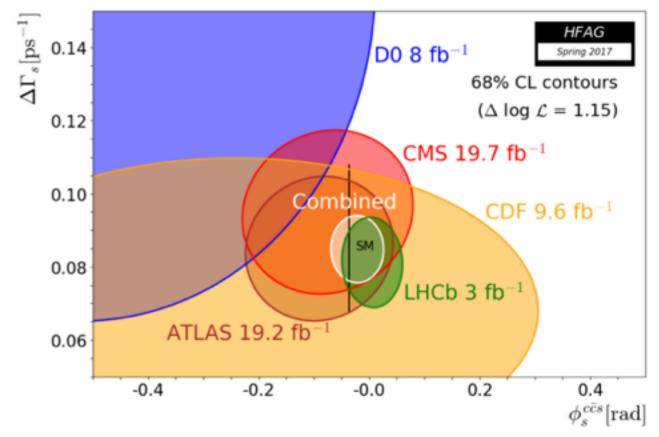
 $h \rightarrow c \bar{c} s$

DUALITY VIOLATION

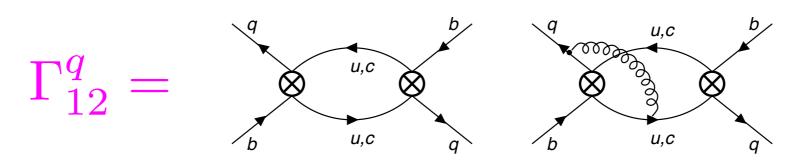
- ➤ Many historic hints for possible duality violation: missing charm puzzle, Λ_b-lifetime, dimuon asymmetry,...
- Duality cannot be proofed QCD solution would be required: test whether duality based predictions agree with experiment
- ► Since Moriond 2012:

size of duality violations is severely constrained by perfect agreement of experiment and theory for

$$\frac{\left(\frac{\Delta\Gamma_s}{\Delta M_s}\right)^{\text{SM}}}{\left(\frac{\Delta\Gamma_s}{\Delta M_s}\right)^{\text{Exp}}} = 0.99 \pm 0.20$$



QUANTIFY THE POSSIBLE SIZE OF DUALITY VIOLATIONS



We expect duality violations to be more pronounced if the final state phase space is becoming smaller

$$\Gamma_{12}^{s,cc} \to \Gamma_{12}^{s,cc}(1+4\delta)$$
,

our ansatz:

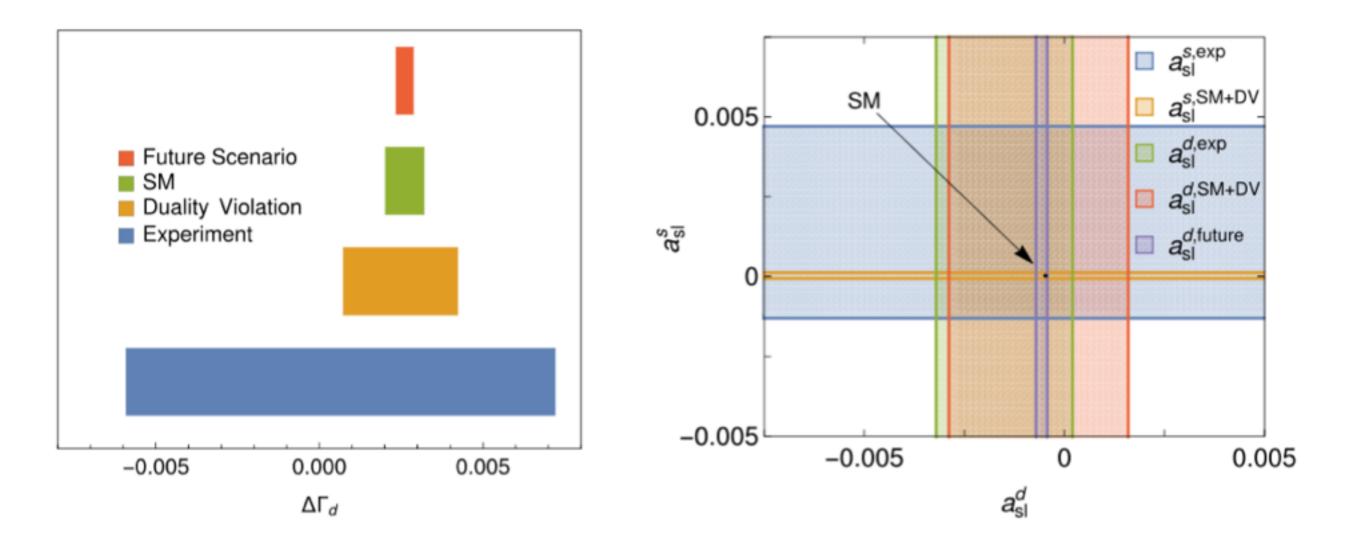
$$\Gamma_{12}^{s,uc} \to \Gamma_{12}^{s,uc}(1+\delta) ,$$

$$\Gamma_{12}^{s,uu} \to \Gamma_{12}^{s,uu}(1+0\delta)$$
.

We get the following dependence of mixing observables

Observable	B_s^0	B^0_d
$\frac{\Delta\Gamma_q}{\Delta M_q}$	$48.1(1+3.95\delta) \cdot 10^{-4}$	$49.5(1+3.76\delta)\cdot 10^{-4}$
$\Delta \Gamma_q$	$0.0880(1+3.95\delta)\mathrm{ps}^{-1}$	$2.61(1+3.759\delta) \cdot 10^{-3} \mathrm{ps}^{-1}$
a^q_{sl}	$2.225(1 - 22.3\delta) \cdot 10^{-5}$	$-4.74(1-24.5\delta)\cdot 10^{-4}$

QUANTIFY THE POSSIBLE SIZE OF DUALITY VIOLATIONS



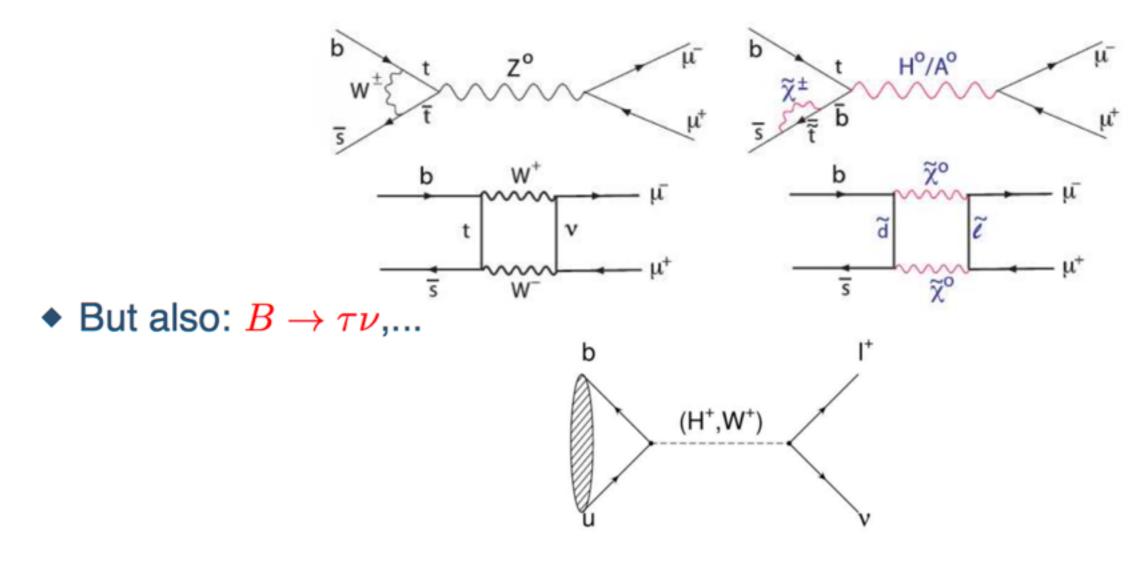
On the ultimate precision of meson mixing observables Thomas Jubb, Matthew Kirk, Alexander Lenz, Gilberto Tetlalmatzi-Xolocotzi Published in Nucl.Phys. B915 (2017) 431-453

MOTIVATION FOR BSM SEARCHES WITH FLAVOUR PHYSICS

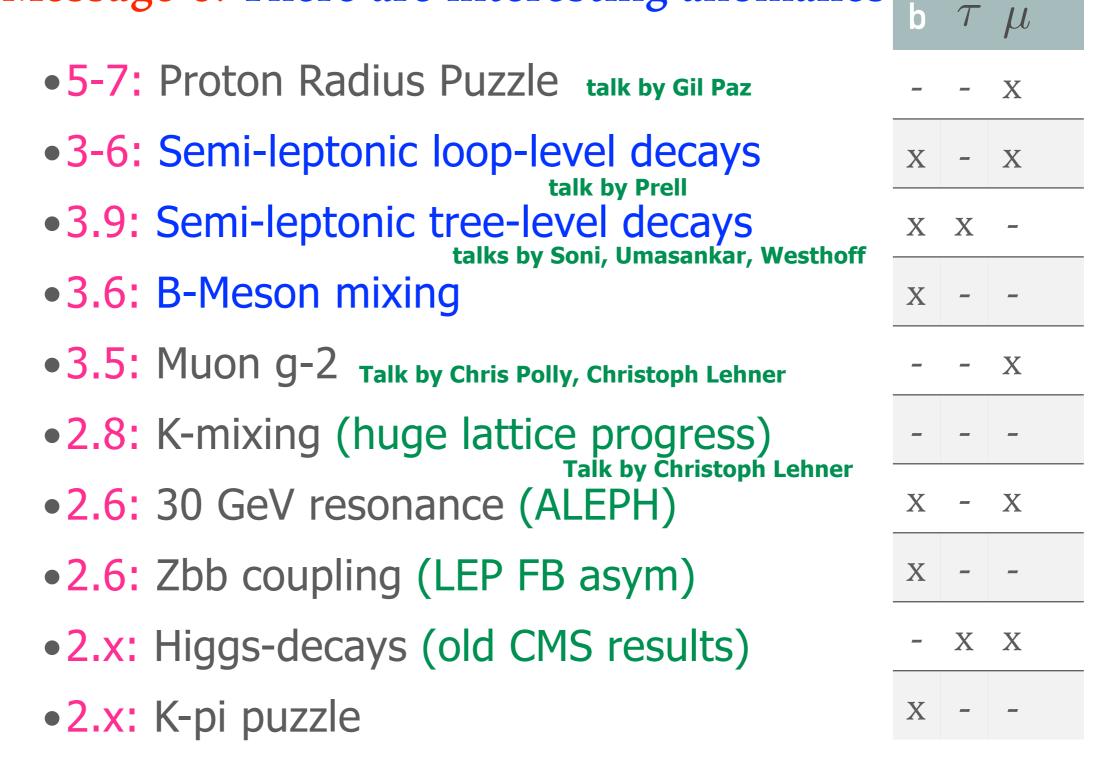
- CP violation till now only found in quark flavour physics
- Theoretically clean: $lpha_s(m_b)pprox 0.2pprox \Lambda/m_b$



- many processes strongly suppressed in the SM due to quantum cor
 - $B_s \rightarrow \mu\mu$ or $b \rightarrow s\gamma$: Flavour Changing Neutral Currents

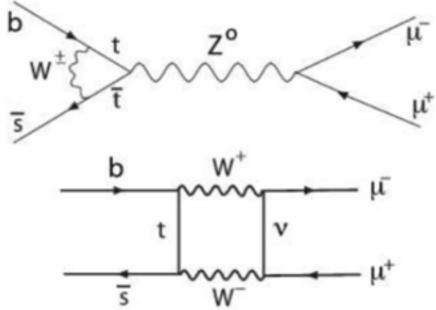


Message 6: There are interesting anomalies



Loop-level (semi) leptonic decays: $b \rightarrow s \mu \mu b$

- very simple hadronic structure
- $B_{d,s} \rightarrow \mu \mu$: decay constant
- $H_b \to H_a \mu \mu$: form factor



Can be determined with lattice, sum rules,... **Talk by Christoph Lehner**

Observables:

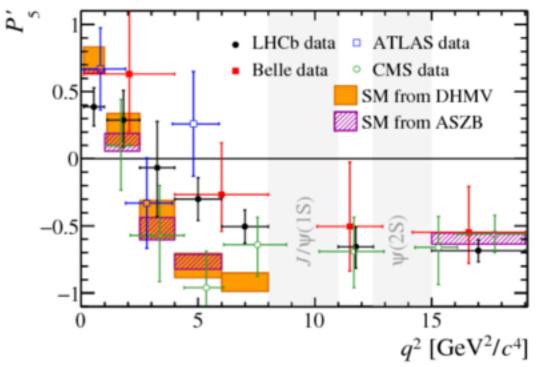
- Branching ratios $Br(B_s \to \phi \mu \mu), Br(B \to K^* \mu \mu),$
- Angular observables, e.g. P'_5 hadronic uncertainties cancel partially

• Ratios $R_K = \frac{Br(B^+ \to K^+ \mu^- \mu^+)}{Br(B^+ \to K^+ e^- e^+)}$ hadronic uncertainties cancel completely

Loop-level (semi) leptonic decays: Pessimistic view

- Hadronic contributions might be larger than expected:
 Rome group; Zwicky,...
 entertaining fights in the community this does not hold for R K
 e.g. Bordone, Isidori, Pattori 1605.07633
- New ATLAS and CMS results are closer to the SM talk by Prell but are consistent with LHCb
- Individual observables do not exceed 3 sigma!

2.9	P_5' [4,6]
2.9	P' ₅ [6,8]
2.6	R_K
2.6	R_{K^*} [1.1,6]
2.3	R_{K^*} [0.045,1.1]
2.2	$Br(B_s \to \phi \mu \mu)$ [2,5]
2.2	$Br(B_s ightarrow \phi \mu \mu)$ [5,8]



Patterns of NP in b->sll transitions in the light of recent data Capdevilla, Crivellin, Descotes-Genon, Matias, Virto 1704.05340

Loop-level (semi) leptonic decays: Optimistic view

all can be fitted in very simple scenario

$$Q_{9V} = \frac{\alpha_e}{4\pi} \left(\bar{s}_L \gamma_\mu b_L \right) \left(\bar{l} \gamma^\mu l \right)$$
$$Q_{10A} = \frac{\alpha_e}{4\pi} \left(\bar{s}_L \gamma_\mu b_L \right) \left(\bar{l} \gamma^\mu \gamma^5 l \right)$$

e.g. just modify the Wilson coefficient C9!

3 *σ* 1704.05447

Ciuchini, Coutinho, Fedele, Franco, Paul, Silvestrini, Valli

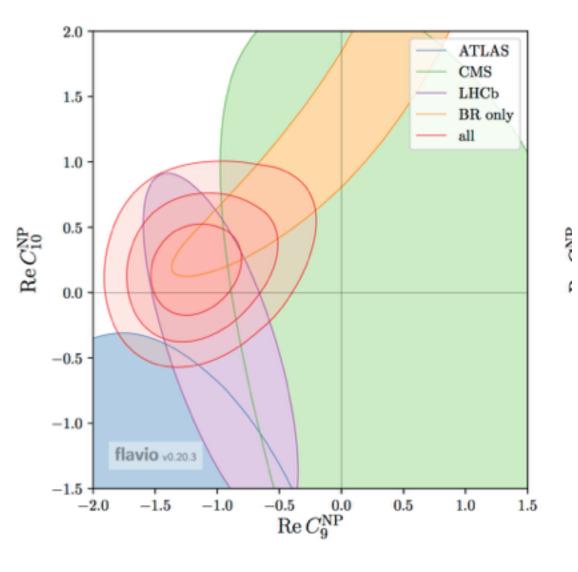
On Flavourful Easter eggs for NP hunger and LFU violation

.... (see next page)

5.7*σ* 1704.05340

Capdevilla, Cvrivellin, Descotes-Genon, Matias, Virto

Patterns of NP in b to all transitions in the light of recent data



arXiv:1703.09189 [pdf, other] Status of the $B \rightarrow K^* \mu^+ \mu^-$ anomaly after Moriond 2017 Wolfgang Altmannshofer, Christoph Niehoff, Peter Stangl, David M. Straub

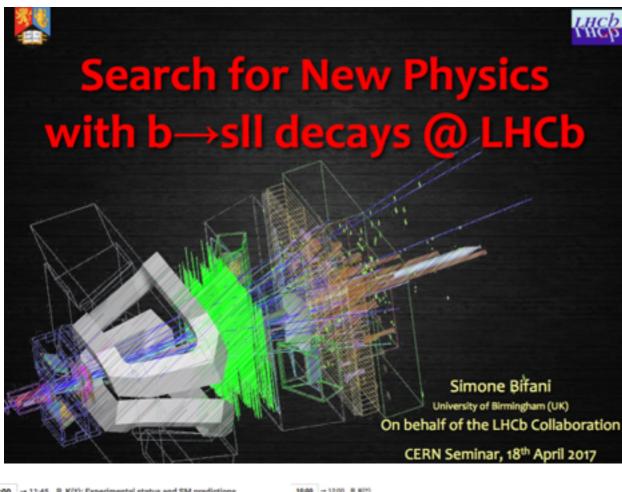
Instant workshop on B meson anomalies

- IT May 2017, 09:00 → 19 May 2017, 15:30 Europe/Zurich
- 4-3-006 TH Conference Room (CERN)

Description In light of recent anomalies in B physics there is an increased interest in the theory community on its implications. As a quick response we are organizing an "Instant workshop on B meson anomalies" at CERN from May 17-May 19 2017.

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411,40	N_N(-). 6	Aperimental status and on predictions		10.00		
	10:00	Experimental status of R_K* and siblings Speaker: Marie Helene Schune (Investe de Paris Sud 11 (78)			10:00	RS and composites Speaker. Guilano Panico (ameniar Autorana in Receipt
	-				10.30	Further thoughts on RS and composites Speaker Mariano Quiros Cartellen (CREA - Instruct only
	10:30	K_K: Quo vadis?				speaker, wanene duros cardelen (color-menuce car
	- 1	R_K: Quo vadis? Speaker: Mitesh Patel (imperial College (Dit))			10:45	Flavor models
						Flavor models Speaker. One Isiduri (inventiat Juenti (21))
	10.00	Overview of SM predictions Speaker: Sebastian Jager (Interest)			11.15	Links MP
						Light NP Speaker Diptimoy Chosh (meanure method)
	11:25	The role of EM corrections				
		The role of EM corrections Speaker: Marzia Bordone (University of Zurich)			11.45	UPV and neutrinos Speaker. John Cargolionis (The University of Network)
4:00 → 15:45	R_K(*): F	its and model building	13:30 -			
	14:00	Angular analysis			13:30	Theory prospects Speaker: Gudrun Hiller: (Subrochs Universitier: Dermund ()
		Angular analysis Speaker: Konstantinos Petridis (University of Inistal (UR))			1	Speaker: Gudhun Hiller (hadmische Universitaer Zertmund (2)
	:				1400	LHCb prospects
	14:30	Global Fits and Global New Physics Patterns				LHCb prospects Speaker: Unk Epede (montal talkor (14))
		Global Fits and Global New Physics Patterns Speaker: Javler Virto (Universitat Geger)			-	
	<u> </u>					Belle 2 prospects Speaker: Bostjan Golob
	14:50	What do we learn from fits				

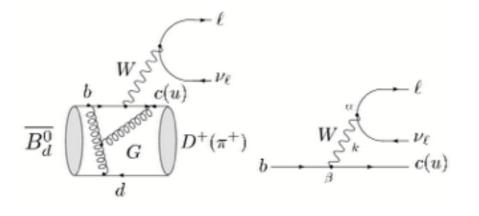
- 1704.05340 Capdevilla, Cvrivellin, Descotes-Genon, Matias, VirtoPatterns of NP in b to all transit
- 1704.05435 Altmannshoher, Stange, Straub Interpreting hints for LEtpon Universality Violation
- 1704.05438 D'Amico, Nardecchia, Panci, Sannino, Stremai, Torre, Urbano Flavour anomalies after
- 1704.05444 Hiller, Nisandzic RK and RK* beyond the SM
- 1704.05446 Geng, Grinstein, Jaeger, Camalich, Ren, Shi Towards the discovery of new physics with
- 1704.05447 Ciuchini, Coutinho, Fedele, Franco, Paul, Silvestrini, Valli On Flavourful Easter eggs f
- 1704.05672 Celis, Fuentes-Martin, Vicente, Virto Gauge-invariant implications of the LHCb measu
- 1704.05835 Becirevic, Sumensari A leptoquark model to accommodate RK and RK*
- T1704.05849 Cai, Gargalionis, Schmidt, Volkas Reconsidering the One Leptoquark solution: flavou
- 1704.06005 Kamenik, Soreq, Zupan Lepton Flavour Universality violation without new sources of
 - 1704.06188 Sala, Straub A new light particle in B decays
 - 1704.06200 Di Chiara, Fowlie, Fraser, Marzo, Marzola, Raidal, Christian Spethmann. Minimal fl
 - 1704.06240 Gosh Explaining RK and RK* anomalies
- 1704.06659 Altmannshofer, Dev, Son I RD(*) anomaly: a possible hint for natural SUSY with R-
- 1704.07397 Alok, Bhattacharya, Datta, Kumar, Kumar, LondON, New physics in b->s mu mu af
- 1704.07347 Alok, Sharma, Kumar, Kumar Lepton-Flavour non-universality in the B-sector: a glo
- 1704.08158 Alonso, Cox, Han, Yamagida Anomaly-free local horizontal symmetry and anomaly-fu
- 1704.08168 Wang, Zhao Implications of the RK and RK* anomalies
- 1704.09015 Admir Greljo, David Marzocca High-pT dilepton tails and flavour physics
- 1705.00915 Cesar Bonilla, Tanmoy Modak, Rahul Srivastava, Jose W. F. Valle U(1)B3–3Lμ gau
 - 1705.00929 Ferruccio Feruglio, Paride Paradisi, Andrea Pattori O

On the Importance of Elect

Tree-level semi leptonic decays

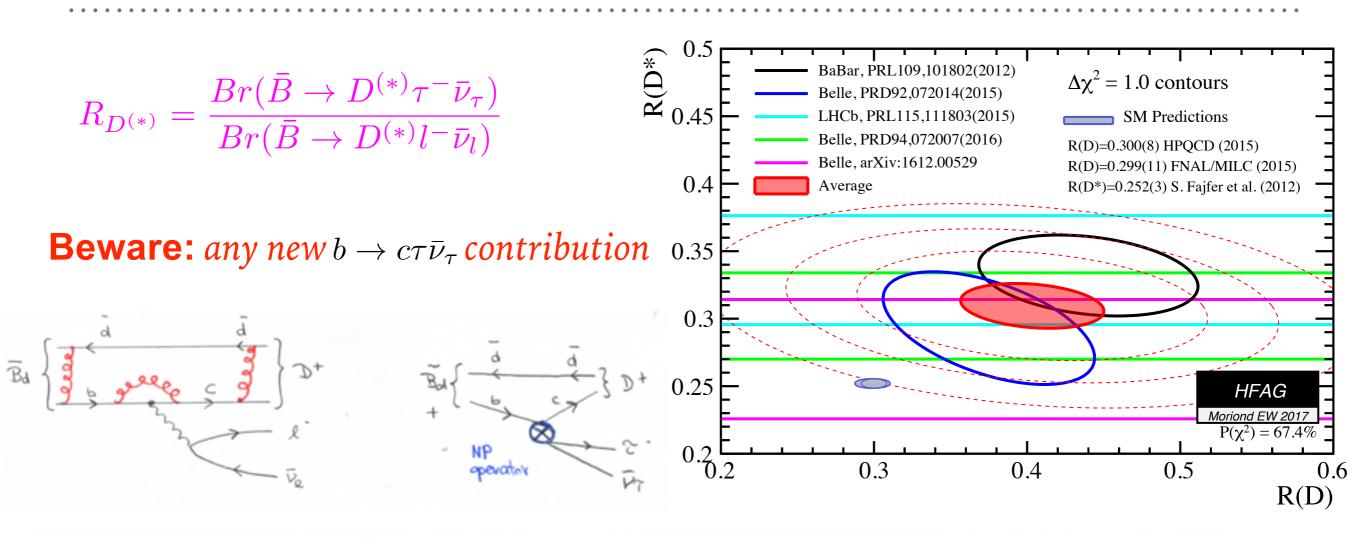
again simple hadronic structure

form factor: lattice, sum rules

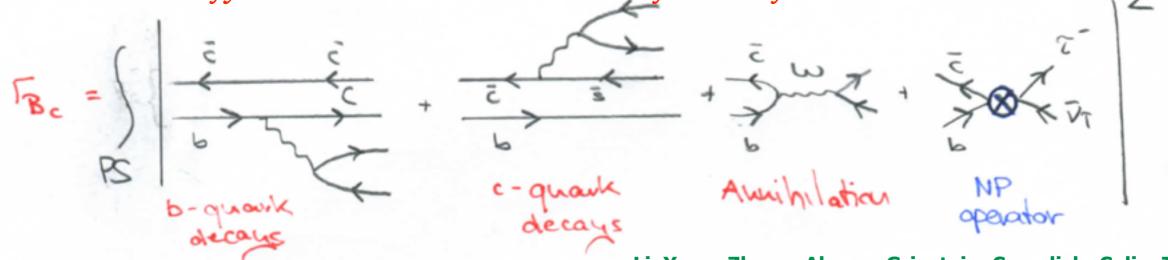


 $3\sigma V_{ub}, V_{cb}$: long standing discrepancy between exclusive and inclusive CKM determination

3.90 $R_{D^{(*)}} = \frac{Br(\bar{B} \to D^{(*)}\tau^-\bar{\nu}_{\tau})}{Br(\bar{B} \to D^{(*)}l^-\bar{\nu}_l)}$ more recent problem *individually:* R_D :2.2 σ R_D^* :3.4 σ **talks by Soni, Umasankar, Westhoff**



will also modify other observables like the lifetime of the Bc meson!



e.g. Li, Yang, Zhang; Alonso, Grinstein, Camalich; Celis, Jung, Li, Pich

B mixing

 $A^{\text{Di-muon}} = C_s a^s_{sl} + C_d a^d_{sl} + \frac{1}{2} C_\Delta \Delta \Gamma_d$ **3.6** σ : D0 result

Evidence for an anomalous like-sign dimuon charge asymmetry V.M.Abazov et al (D0 Collaboration) Phys. Rev. Lett 105 (2010) 081801

Study of CP violating charge asymmetry... V.M.Abazov et al (D0 Collaboration) Phys. Rev. D 89 (2014) 012002 Understanding the anomalous like-sign dca Guennadi Borissov, Boris Hoeneisen Phys. Rev. D 87 (2013) 074020

Effect of Delta Gamma_d on the dimoun asymmetry Uli Nierste Talk at CKM 2014

seems to be the largest individual deviation

2
$$\mathcal{O}$$
: New lattice results $\Delta M_s = 2|M_{12}^s|, \quad \Delta \Gamma_s = 2|\Gamma_{12}^s|\cos\phi_{12}^s, \quad a_{sl}^s = \left|\frac{\Gamma_{12}^s}{M_{12}^s}\right|\sin\phi_{12}^s$

Observable	e SM – conservative	SM – aggressive	Experiment
ΔM_s	$(18.3 \pm 2.7) \text{ ps}^{-1}$	$(20.11 \pm 1.37) \text{ ps}^{-1}$	$(17.757 \pm 0.021) \text{ ps}^{-1}$
$\Delta\Gamma_s$	$(0.088 \pm 0.020) \text{ ps}^{-1}$	$(0.098 \pm 0.014) \text{ ps}^{-1}$	$(0.082 \pm 0.006) \text{ ps}^{-1}$
a_{sl}^s	$(2.22 \pm 0.27) \cdot 10^{-5}$	$(2.27 \pm 0.25) \cdot 10^{-5}$	$(-7.5 \pm 4.1) \cdot 10^{-3}$

B(s)-mixing matrix elements from lattice QCD for the SM and beyond Fermilab Lattice and MILC Collaborations Phys.Rev. D93 (2016) no.11, 113016, arXiv:1602.03560 [hep-lat] On the ultimate precision of meson mixing observables Thomas Jubb, Matthew Kirk, AL, Gilberto Tetlalmatzi-Xolocotzi Nucl.Phys. B915 (2017) 431-453

List of models:

- Z' new U(1) or SU(2)
- Leptoquarks
- W'- new SU(2)
- Composite Models
- WED
- SUSY
- 2HDM

•

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hundreds of papers...

"Qual der Wahl"

agony of choice

or

choice of agony?



A popular BSM model for solving the anomalies related to loop-level (semi) leptonic decays are Z' models:

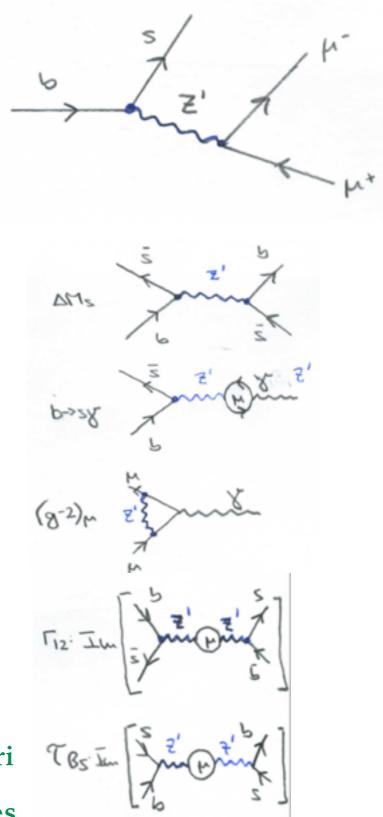
Such a new tree-level transition will also affect many other observables,

most notably **B-mixing at tree-level**,

but also many loop processes.

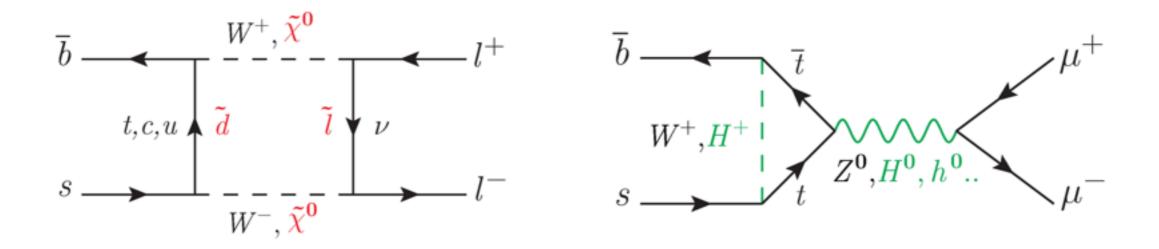
Make sure all relevant bounds are included, e.g. electro-weak precision bounds 1705.00929 Ferruccio Feruglio, Paride Paradisi, Andrea Pattori

On the Importance of Electroweak Corrections for B Anomalies



INDIRECT NP SEARCHES: NON-MAIN STREAM

► Main stream: BSM effects hide in loop processes



► Non-main stream:

Consider BSM effects in non-leptonic tree-level decays?

- typically not considered, but: **Bauer, Dunn; hep-ph/1006.1629** as an explanation for the dimuon asymmetry
- can clearly not be O(100%), but what about 20%?

NP IN TREE-LEVEL DECAYS

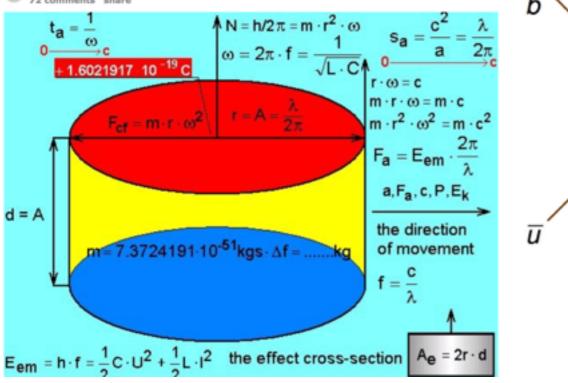


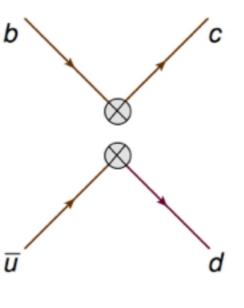
Main Chapters
Introduction
Why Time Dilation must be impossible



$$\hat{\mathcal{H}}_{eff} = \frac{V_{cb}V_{ud}^*}{\sqrt{2}} \left(C_1 \hat{Q}_1 + C_2 \hat{Q}_2 \right)$$

INAGE Anyone else get these weird emails from Gabor Fekete? (Ling submitted 2 years ago by Astrokiwi Astrophysics 72 comments share





Do a systematic study of tree-level observables that are both well known in experiment and theory

$C_{1,2}^{SM} \to C_{1,2}^{SM} + \Delta C_{1,2}$

CAA

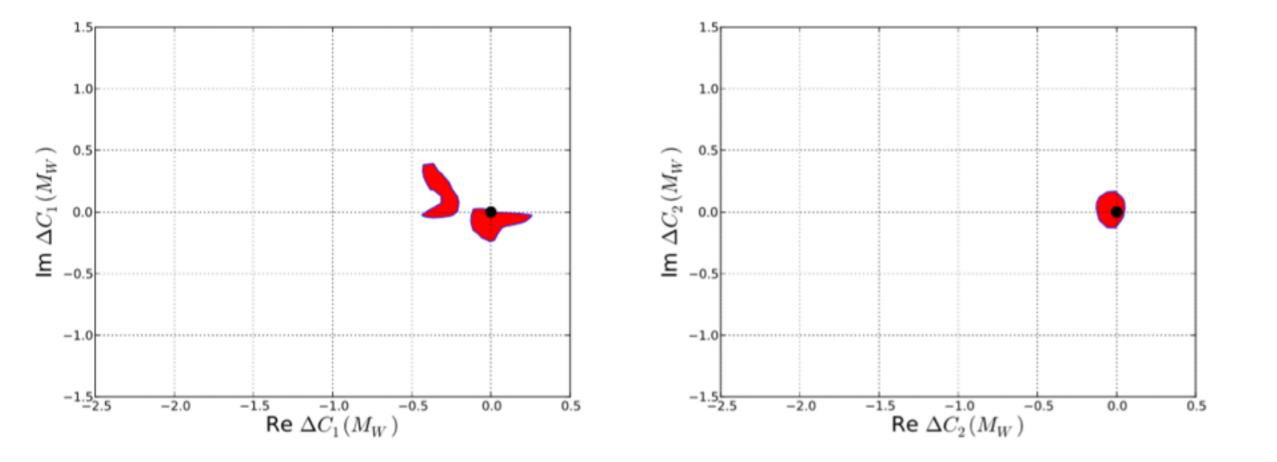
C A A

4.3	Constraints from $b \to u \bar{u} d$ transitions		
	4.3.1	$R_{\pi\pi}$	
	4.3.2	$S_{\pi\pi}$ and $S_{\rho\pi}$	
	4.3.3	$R_{ ho ho}$	
4.4	Constraints from $b \to c \bar{u} d$ transitions $\ . \ . \ .$		
	4.4.1	$\bar{B}^0 \rightarrow D^{*+} \pi^-$	
	4.4.2	S_{D^*h}	
4.5	Observables constraining $b \to c \bar c d$ transitions $% c \bar c d \bar c$		
	4.5.1	M^d_{12}	
	4.5.2	$B \to X_d \gamma$	
4.6	Constr	caints from $b \to c\bar{c}s$ transitions $\ldots \ldots$	
	4.6.1	$\bar{B} \to X_s \gamma$	

4.7 Constraints using multiple channels observables: a_{sl}^s , a_{sl}^d and $\Delta\Gamma$

NP IN TREE-LEVEL DECAYS

Result:



What does this mean? Is this an important effect?

NP IN TREE-LEVEL DECAYS

► Decay rate difference of neutral Bd mesons, $\Delta\Gamma_d$, can be enhanced by several 100% work triggered by D0 di-muon asymmetry - Borissov work triggered ATLAS measurement of $\Delta\Gamma_d$ - Borissov

Extraction of CKM angleγ can be modified by several degrees SM precision: 1 ppm
NP effects in tree-level decay and the precision of

Experimental precision: now 6deg, future 1 deg

NP effects in tree-level decay and the precision of γ Brod, Lenz, Tetlamatzi-Xolocotzi Alexander Lenz Rev.Mod.Phys. 88 (2016) no.4,045002

► More profound analysis in progress

AL, Tetlalmatzi-Xolocotzi

till now only SM Dirac structures

NP IN RARE B DECAYS AND MIXING

Is there a connection between mixing and rare decays? Charming new physics in rare B-decays and mixing Jaeger, Kirk, Lenz, Leslie arXiv: 1701.09183

Consider NP in tree-level b -> ccs traditions with general Dirac structures

This affects rare decays and mixing/lifetimes:

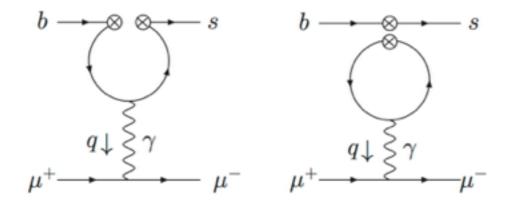


FIG. 1. Leading Feynman diagrams for CBSM contributions to rare and semileptonic decays. With our choice of Fierzordering, only the diagram on the left is relevant.

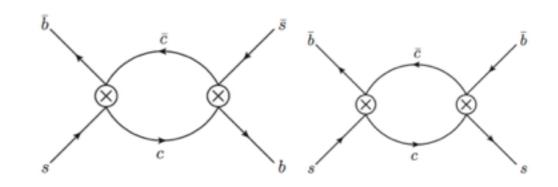


FIG. 2. Leading Feynman diagrams for CBSM contributions to the width difference $\Delta\Gamma_s$ (left) and the lifetime ratio $\tau(B_s)/\tau(B_d)$ (right).

NP IN RARE B DECAYS

C_7 and C_9 depend on different new physics contributions!

$$\Delta C_9^{\text{eff}}(q^2) = \left(C_{1,2}^c - \frac{C_{3,4}^c}{2}\right)h - \frac{2}{9}C_{3,4}^c,$$

$$\Delta C_7^{\text{eff}}(q^2) = \frac{m_c}{m_b} \left[\left(4C_{9,10}^c - C_{7,8}^c\right)y + \frac{4C_{5,6}^c - C_{7,8}^c}{6} \right]$$

with $C^c_{x,y} = 3 \Delta C^c_x + \Delta C^c_y$ and the loop functions

New physics contributions to rare b decays are now q2 dependent!

$$egin{aligned} h(q^2,m_c,\mu) &= -rac{4}{9} \left[\ln rac{m_c^2}{\mu^2} - rac{2}{3} + (2+z)a(z) - z
ight], \ y(q^2,m_c,\mu) &= -rac{1}{3} \left[\ln rac{m_c^2}{\mu^2} - rac{3}{2} + 2a(z)
ight], \end{aligned}$$

where $a(z) = \sqrt{|z-1|} \arctan \frac{1}{\sqrt{z-1}}$ and $z = 4m_c^2/q^2$.

Interesting RGE effects

Charming new physics in rare B-decays and mixing Jaeger, Kirk, Lenz, Leslie arXiv: 1701.09183

$$\begin{pmatrix} \Delta C_1(\mu) \\ \Delta C_2(\mu) \\ \Delta C_3(\mu) \\ \Delta C_4(\mu) \\ \Delta C_7^{\text{eff}}(\mu) \\ \Delta C_9(\mu) \end{pmatrix} = \begin{pmatrix} 1.12 & -0.27 & 0 & 0 \\ -0.27 & 1.12 & 0 & 0 \\ 0 & 0 & 0.92 & 0 \\ 0 & 0 & 0.33 & 1.92 \\ 0.02 & -0.19 & -0.01 & -0.13 \\ 8.65 & 2.00 & -4.33 & -1.95 \end{pmatrix} \begin{pmatrix} \Delta C_1(\mu_0) \\ \Delta C_2(\mu_0) \\ \Delta C_3(\mu_0) \\ \Delta C_4(\mu_0) \end{pmatrix}.$$

NP IN MIXING & LIFETIMES

The decay rate difference of neutral B mesons is again a very strong constraint

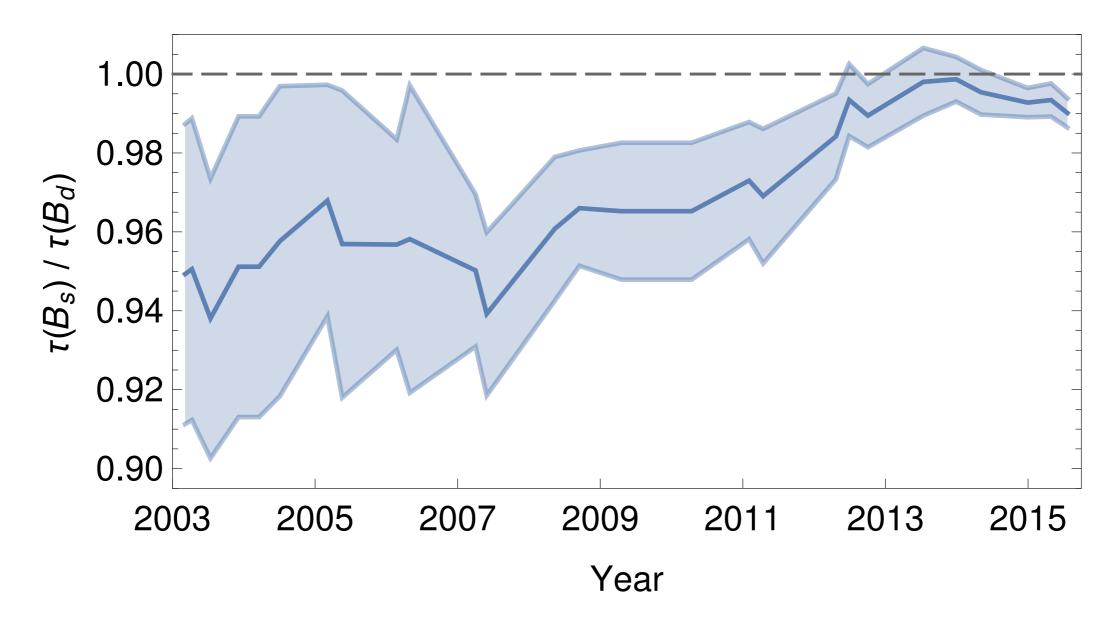
The lifetime ratio of neutral B mesons is more or less exactly one in the SM -> still more precise experimental values needed

$$\begin{split} \Gamma_{12}^{cc} &= -G_F^2 (V_{cs}^* V_{cb})^2 m_b^2 M_{B_s} f_{B_s}^2 \frac{\sqrt{1-4x_c^2}}{576\pi} \times \\ & \left\{ \begin{bmatrix} 16(1-x_c^2)(4C_2^{c,2}+C_4^{c,2})+8(1-4x_c^2) \times \\ (12C_1^{c,2}+8C_1^cC_2^c+2C_3^cC_4^c+3C_3^{c,2})-192x_c^2 \times \\ (3C_1^cC_3^c+C_1^cC_4^c+C_2^cC_3^c+C_2^cC_4^c) \end{bmatrix} B+2(1+2x_c^2) \times \\ & \left(4C_2^{c,2}-8C_1^cC_2^c-12C_1^{c,2}-3C_3^{c,2}-2C_3^cC_4^c+C_4^{c,2})\tilde{B}_S' \right\} \end{split}$$

$$\begin{pmatrix} \frac{\tau_{B_s}}{\tau_{B_d}} \end{pmatrix}_{\rm NP} = G_F^2 |V_{cb} V_{cs}|^2 m_b^2 M_{B_s} f_{B_s}^2 \tau_{B_s} \frac{\sqrt{1 - 4x_c^2}}{144\pi} \times \\ \begin{cases} \left(1 - x_c^2\right) \left[\left(4C_{1,2}^{c,2} + C_{3,4}^{c,2}\right)B_1 + 6\left(4C_2^{c,2} + C_4^{c,2}\right)\epsilon_1 \right] \\ -12x_c^2 \left(C_{1,2}^c C_{3,4}^c B_1 + 6C_2^c C_4^c \epsilon_1\right) - \left(1 + 2x_c^2\right) \times \\ \left[\left(4C_{1,2}^{c,2} + C_{3,4}^{c,2}\right)B_2 + 6\left(4C_2^{c,2} + C_4^{c,2}\right)\epsilon_2 \right] \end{cases}, \quad (13)$$

Charming new physics in rare B-decays and mixing Jaeger, Kirk, Lenz, Leslie arXiv: 1701.09183

NP IN LIFETIMES

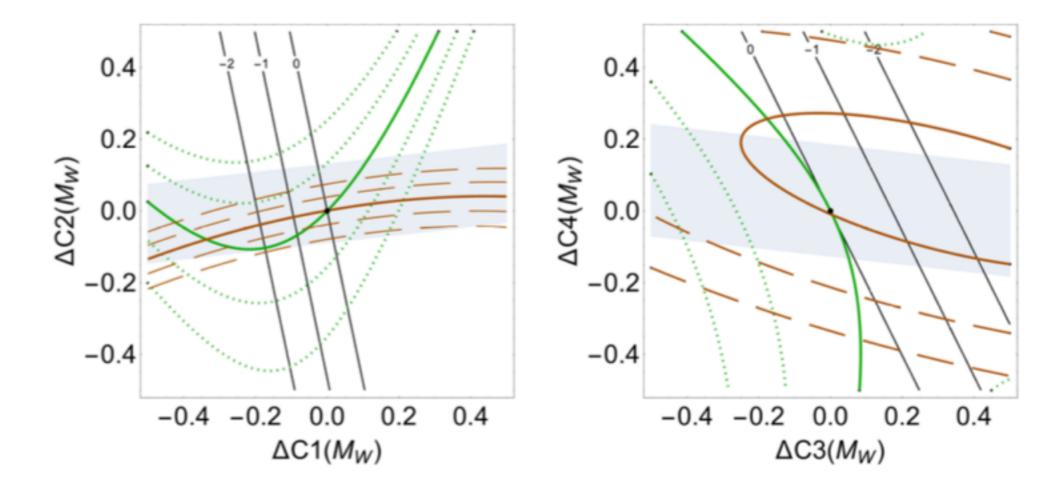


Current experimental precision is not sufficient

B_s^0	1.505 ± 0.005 ps	$B_s^0/B^0 = 0.990 \pm 0.004$
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NP IN RARE B DECAYS & MIXING

> Deviation in some rare B-decays can be explained without violating other bounds



- it is possible that NP in rare decays is q²dependent! (more profound study in progress: Jaeger, Kirk, Lenz, Leslie)
- A UV complete model will typically produce loop-contributions to our anomalous observables (main-stream) but in general also new tree-level contributions (non main stream) - both have to be considered in the end!

CONTENT

Introduction - Motivation for Flavour Physics

- Message 1: SM and CKM work perfectly
- Message 2: Many times experimental error is smaller than theoretical one
- Message 3: !!!Higher precision in theory needed!!!
- Message 4: !!!Standard assumptions might have to be reconsidered!!!
- Message 5: SM/CKM dominance gives important bounds on BSM models
- Message 6: There are very interesting anomalies
- ► Message 7: We are still waiting for a single 5 sigma deviation
- Message 8: Many observables have to be considered for BSM searches
- Message 9: Alternative paths for BSM searches can be interesting

END

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