



WAYNE STATE
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The Proton Radius Puzzle

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Form Factors

- Matrix element of EM current between nucleon states give rise to two form factors ($q = p_f - p_i$)

$$\langle N(p_f) | \sum_q e_q \bar{q} \gamma^\mu q | N(p_i) \rangle = \bar{u}(p_f) \left[\gamma^\mu F_1(q^2) + \frac{i\sigma_{\mu\nu}}{2m} F_2(q^2) q^\nu \right] u(p_i)$$

- Sachs electric and magnetic form factors

$$G_E(q^2) = F_1(q^2) + \frac{q^2}{4m_p^2} F_2(q^2) \quad G_M(q^2) = F_1(q^2) + F_2(q^2)$$

$$G_E^p(0) = 1$$

$$G_M^p(0) = \mu_p \approx 2.793$$

- The slope of G_E^p

$$\langle r^2 \rangle_E^p = 6 \left. \frac{dG_E^p}{dq^2} \right|_{q^2=0}$$

determines the charge radius $r_E^p \equiv \sqrt{\langle r^2 \rangle_E^p}$

- The proton *magnetic* radius

$$\langle r^2 \rangle_M^p = \frac{6}{G_M^p(0)} \left. \frac{dG_M^p(q^2)}{dq^2} \right|_{q^2=0}$$

Charge radius from atomic physics



- Lamb shift in muonic hydrogen [Pohl et al. Nature **466**, 213 (2010)]
 $r_E^p = 0.84184(67) \text{ fm}$
more recently $r_E^p = 0.84087(39) \text{ fm}$ [Antognini et al. Science **339**, 417 (2013)]
- CODATA value [Mohr et al. RMP **80**, 633 (2008)]
 $r_E^p = 0.87680(690) \text{ fm}$
more recently $r_E^p = 0.87750(510) \text{ fm}$ [Mohr et al. RMP **84**, 1527 (2012)]
extracted mainly from (electronic) hydrogen
- **5σ discrepancy!**
- This is the proton radius puzzle

What could be the reason for the discrepancy?

- What could the reason for the discrepancy?
 - 1) Problem with the electronic extraction?
 - 2) Hadronic Uncertainty?
 - 3) New Physics?

Problem with the electronic extraction?

- Recent development: use of the z expansion based on known analytic properties of form factors
- **The** method for **meson** form factors
[Flavor Lattice Averaging Group, EPJ C **74**, 2890 (2014)]
- Now applied successfully to **baryon** form factors to extract $r_E^p, r_M^p, r_M^n, m_A \dots$

Citation: C. Patrignani *et al.* (Particle Data Group), *Chin. Phys. C*, **40**, 100001 (2016)

p CHARGE RADIUS

This is the rms electric charge radius, $\sqrt{\langle r_E^2 \rangle}$.

<u>VALUE (fm)</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
0.8751 ± 0.0061	MOHR	16	RVUE 2014 CODATA value
0.84087 ± 0.00026 ± 0.00029	ANTOGNINI	13	LASR μp -atom Lamb shift
● ● ● We do not use the following data for averages, fits, limits, etc. ● ● ●			
0.895 ± 0.014 ± 0.014	¹ LEE	15	SPEC Just 2010 Mainz data
0.916 ± 0.024	LEE	15	SPEC World data, no Mainz
0.8775 ± 0.0051	MOHR	12	RVUE 2010 CODATA, $e p$ data
0.875 ± 0.008 ± 0.006	ZHAN	11	SPEC Recoil polarimetry
0.879 ± 0.005 ± 0.006	BERNAUER	10	SPEC $e p \rightarrow e p$ form factor
0.912 ± 0.009 ± 0.007	BORISYUK	10	reanalyzes old $e p$ data
0.871 ± 0.009 ± 0.003	HILL	10	z-expansion reanalysis
0.84184 ± 0.00036 ± 0.00056	POHL	10	LASR See ANTOGNINI 13
0.8768 ± 0.0069	MOHR	08	RVUE 2006 CODATA value
0.844 +0.008 -0.004	BELUSHKIN	07	Dispersion analysis
0.897 ± 0.018	BLUNDEN	05	SICK 03 + 2γ correction
0.8750 ± 0.0068	MOHR	05	RVUE 2002 CODATA value
0.895 ± 0.010 ± 0.013	SICK	03	$e p \rightarrow e p$ reanalysis

[Hill, GP PRD **82** 113005 (2010)]

[Lee, Arrington, Hill, PRD **92**, 013013 (2015)]

Citation: C. Patrignani *et al.* (Particle Data Group), *Chin. Phys. C*, **40**, 100001 (2016)

p MAGNETIC RADIUS

This is the rms magnetic radius, $\sqrt{\langle r_M^2 \rangle}$.

<u>VALUE (fm)</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
$0.776 \pm 0.034 \pm 0.017$	¹ LEE	15	SPEC Just 2010 Mainz data
● ● ● We do not use the following data for averages, fits, limits, etc. ● ● ●			
0.914 ± 0.035	LEE	15	SPEC World data, no Mainz
0.87 ± 0.02	EPSTEIN	14	Using ep , $e\pi$, $\pi\pi$ data
$0.867 \pm 0.009 \pm 0.018$	ZHAN	11	SPEC Recoil polarimetry
$0.777 \pm 0.013 \pm 0.010$	BERNAUER	10	SPEC $ep \rightarrow ep$ form factor
$0.876 \pm 0.010 \pm 0.016$	BORISYUK	10	Reanalyzes old $ep \rightarrow ep$ data
0.854 ± 0.005	BELUSHKIN	07	Dispersion analysis

¹ Authors also provide values for a combination of all available data.

[Epstein, GP, Roy PRD **90**, 074027 (2014)]
 [Lee, Arrington, Hill, PRD **92**, 013013 (2015)]

PDG 2016: r_M^n

Citation: C. Patrignani *et al.* (Particle Data Group), Chin. Phys. C, **40**, 100001 (2016)

n MAGNETIC RADIUS

This is the rms magnetic radius, $\sqrt{\langle r_M^2 \rangle}$.

<u>VALUE (fm)</u>	<u>DOCUMENT ID</u>	<u>COMMENT</u>
$0.864^{+0.009}_{-0.008}$ OUR AVERAGE		
0.89 ± 0.03	EPSTEIN	14 Using ep , en , $\pi\pi$ data
$0.862^{+0.009}_{-0.008}$	BELUSHKIN	07 Dispersion analysis

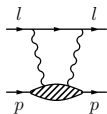
[Epstein, GP, Roy PRD **90**, 074027 (2014)]

The bottom line

- Scattering:
 - World $e - p$ data [Lee, Arrington, Hill '15]
 $r_E^p = 0.918 \pm 0.024$ fm
 - Mainz $e - p$ data [Lee, Arrington, Hill '15]
 $r_E^p = 0.895 \pm 0.020$ fm
 - Proton, neutron and π data [Hill, GP '10]
 $r_E^p = 0.871 \pm 0.009 \pm 0.002 \pm 0.002$ fm
- Muonic hydrogen
 - [Pohl et al. Nature **466**, 213 (2010)]
 $r_E^p = 0.84184(67)$ fm
 - [Antognini et al. Science **339**, 417 (2013)]
 $r_E^p = 0.84087(39)$ fm
- The bottom line:
using z expansion scattering disfavors muonic hydrogen
- Is there a problem with muonic hydrogen *theory*?

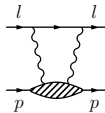
Muonic hydrogen theory

- Is there a problem with muonic hydrogen *theory*?
- Potentially yes!
[Hill, GP PRL **107** 160402 (2011)]
- Muonic hydrogen measures ΔE and translates it to r_E^p
 - [Pohl et al. Nature **466**, 213 (2010) Supplementary information]
 $\Delta E = 206.0573(45) - 5.2262(r_E^p)^2 + 0.0347(r_E^p)^3$ meV
 - [Antognini et al. Science **339**, 417 (2013), Ann. of Phys. **331**, 127]
 $\Delta E = 206.0336(15) - 5.2275(10)(r_E^p)^2 + 0.0332(20)$ meV
- In both cases apart from r_E^p need two-photon exchange



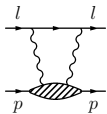
Two photon exchange

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Two photon exchange

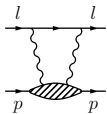
- In both cases apart from r_E^p we have two-photon exchange



$$\begin{aligned}
 W^{\mu\nu} &= \frac{1}{2} \sum_s i \int d^4x e^{iq \cdot x} \langle \mathbf{k}, s | T \{ J_{\text{e.m.}}^\mu(x) J_{\text{e.m.}}^\nu(0) \} | \mathbf{k}, s \rangle \\
 &= \left(-g^{\mu\nu} + \frac{q^\mu q^\nu}{q^2} \right) W_1 + \left(k^\mu - \frac{k \cdot q q^\mu}{q^2} \right) \left(k^\nu - \frac{k \cdot q q^\nu}{q^2} \right) W_2
 \end{aligned}$$

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 \end{aligned}$$

- Dispersion relations ($\nu = 2k \cdot q$, $Q^2 = -q^2$)

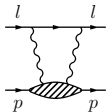
$$W_1(\nu, Q^2) = W_1(0, Q^2) + \frac{\nu^2}{\pi} \int_{\nu_{\text{cut}}(Q^2)^2}^{\infty} d\nu'^2 \frac{\text{Im} W_1(\nu', Q^2)}{\nu'^2(\nu'^2 - \nu^2)}$$

$$W_2(\nu, Q^2) = \frac{1}{\pi} \int_{\nu_{\text{cut}}(Q^2)^2}^{\infty} d\nu'^2 \frac{\text{Im} W_2(\nu', Q^2)}{\nu'^2 - \nu^2}$$

- W_1 requires subtraction...

Two photon exchange

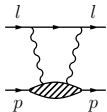
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- Imaginary part of TPE related to data:
form factors, structure functions

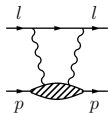
Two photon exchange

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- Imaginary part of TPE related to data:
form factors, structure functions
- Cannot reproduce it from its imaginary part:
Dispersion relation requires subtraction
- Need poorly constrained non-perturbative function $W_1(0, Q^2)$
- Calculable in small Q^2 limit using NRQED
[Hill, GP, PRL **107** 160402 (2011)]

Two Photon Exchange: large Q^2 limit



- Calculable in *large* Q^2 limit using Operator Product Expansion (OPE) [J. C. Collins, NPB **149**, 90 (1979)]
The photon “sees” the quarks and gluons inside the proton

$$W_1(0, Q^2) = c/Q^2 + \mathcal{O}(1/Q^4)$$

- Result was used to estimate two photon exchange effects
- c calculated in [J. C. Collins, NPB **149**, 90 (1979)]

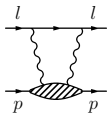
RENORMALIZATION OF THE COTTINGHAM FORMULA

John C. COLLINS *

Joseph Henry Laboratories, Princeton University, Princeton, New Jersey 08540, USA

Received 23 October 1978

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- Was it?

Large Q^2 behavior

- In 1978 Collins calculated EM corrections to the nucleon mass
The mass only depends on spin-0 operators (q quark, $G^{\mu\nu}$ gluon)

$$\langle P | m_q \bar{q}q | P \rangle, \quad \langle P | G^{\mu\nu} G_{\mu\nu} | P \rangle$$

	Quark	Gluon
Spin-0	Collins '78	Collins '78

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- For $W_1(0, Q^2)$ you need also spin-2 operators

$$\langle P | \bar{q} (iD^\mu \gamma^\nu + iD^\nu \gamma^\mu - \frac{1}{4} i \not{D} g^{\mu\nu}) q | P \rangle, \quad \langle P | G^{\mu\alpha} G_\alpha^\nu - \frac{1}{4} G^{\alpha\beta} G_{\alpha\beta} g^{\mu\nu} | P \rangle$$

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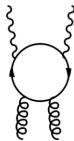
- Need to calculate the spin-2 contribution
[Hill, GP arXiv:1611.09917, to appear in PRD]

	Quark	Gluon
Spin-0	Collins '78	Collins '78
Spin-2	Hill, GP '16	Hill, GP '16

- Collins's result is not enough for muonic hydrogen!

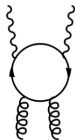
Large Q^2 behavior

- Requires 1-loop calculation



Large Q^2 behavior

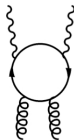
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- Doing that, we found a mistake in Collins 1978 spin-0 calculation...

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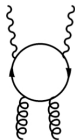
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- Collins didn't calculate the spin-0 gluon contribution directly
He extracted it from another calculation

Large Q^2 behavior

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- For three light quark u, d, s
Correct result: $\sum_q e_q^2 = \left(\frac{2}{3}\right)^2 + \left(\frac{1}{3}\right)^2 + \left(\frac{1}{3}\right)^2 = \frac{2}{3}$
Collins: $\sum_q = 3$
Too large by a factor of 4.5...

Large Q^2 behavior

	Quark	Gluon	
Spin-0	Collins '78	Collins '78	Hill, GP '16
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- Even worse, quark spin-0 and gluon spin-0 come with opposite signs
After correcting the mistake they largely cancel
 $W_1(0, Q^2)$ is **dominated** by spin-2 contribution
- Lesson: It is important to do a full calculation

Large Q^2 behavior

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since gluon contribution is the same at lowest order in isospin breaking

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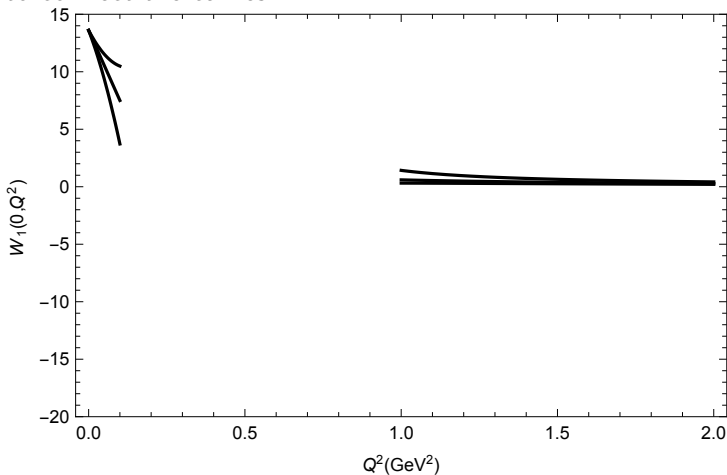
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- Some good news: The mistake has no effect on $m_n - m_p$
since gluon contribution is the same at lowest order in isospin breaking
- Flip side: You cannot use $m_n - m_p$ to constrain muonic hydrogen

Two Photon Exchange: Modeling

- “Aggressive” modeling: use OPE for $Q^2 \geq 1 \text{ GeV}^2$
Use NRQED for $Q^2 \leq 0.1 \text{ GeV}^2$
Model unknown Q^4 and $1/Q^4$

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Use NRQED for $Q^2 \leq 0.1 \text{ GeV}^2$
Model unknown Q^4 and $1/Q^4$
- How to connect the curves?

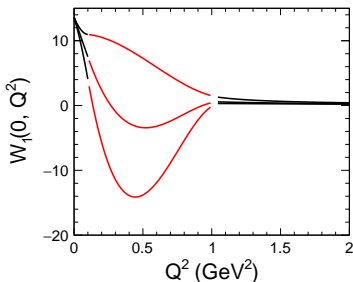


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- Interpolating:

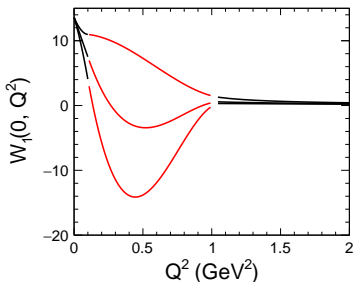
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- Energy contribution: $\delta E(2S) W_1(0, Q^2) \in [-0.046 \text{ meV}, -0.021 \text{ meV}]$
To explain the puzzle need this to be $\sim -0.3 \text{ meV}$
- Caveats: OPE might be only valid for larger Q^2
 $W_1(0, Q^2)$ might be different than the interpolated lines

Experimental test

- How to test?
- New experiment: $\mu - p$ scattering
MUSE (MUon proton Scattering Experiment) at PSI
[R. Gilman et al. (MUSE Collaboration), arXiv:1303.2160]



- Need to connect muon-proton scattering and muonic hydrogen
can use a new effective field theory: QED-NRQED
[Hill, Lee, GP, Mikhail P. Solon, PRD **87** 053017 (2013)]
[Steven P. Dye, Matthew Gonderinger, GP, PRD **94** 013006 (2016)]

Conclusions

Conclusions

- Proton radius puzzle: $> 5\sigma$ discrepancy between
 - r_E^p from muonic hydrogen
 - r_E^p from hydrogen and $e - p$ scattering
- Recent muonic deuterium results find similar discrepancies
[Pohl et al. Science **353**, 669 (2016)]
- After 6 years the origin is still not clear
 - 1) Is it a problem with the electronic extraction?
 - 2) Is it a hadronic uncertainty?
 - 3) is it new physics?
- Motivates a reevaluation of our understanding of the proton

Conclusions

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 - 1) Extraction of proton radii from scattering:
Using the z expansion disfavors the muonic hydrogen value

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Can improve the modeling of two photon exchange effects

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 - 1) Extraction of proton radii from scattering:
Using the z expansion disfavors the muonic hydrogen value
 - 2) The first *full* and *correct* evaluation of large Q^2 behavior of forward virtual Compton tensor
Can improve the modeling of two photon exchange effects
- Motivates a direct connection between muon-proton scattering (MUSE experiment) and muonic hydrogen using a new effective field theory: QED-NRQED
[Steven P. Dye, Matthew Gonderinger, GP, *in progress*]

Conclusions

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 - 1) Extraction of proton radii from scattering:
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 - 2) The first *full* and *correct* evaluation of large Q^2 behavior of forward virtual Compton tensor
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- Motivates a direct connection between muon-proton scattering (MUSE experiment) and muonic hydrogen using a new effective field theory: QED-NRQED
[Steven P. Dye, Matthew Gonderinger, GP, *in progress*]
- Much more work to do!
- Thank you

Backup

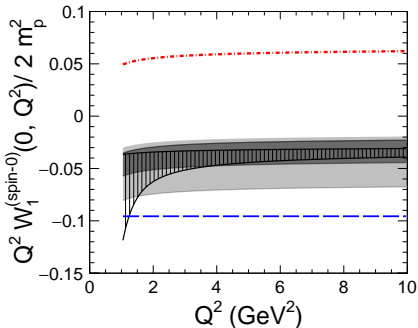
Large Q^2 behavior: Spin 0 contribution

- The *correct* spin 0 result

$$\frac{Q^2}{2m_p^2} W_1^{(\text{spin}=0)}(0, Q^2) = -2 \sum_q e_q^2 f_q^{(0)} + \left(\sum_q e_q^2 \right) \frac{\alpha_s}{12\pi} \tilde{f}_g^{(0)}$$

- quark and gluon matrix elements related by

$$1 = (1 - \gamma_m) \sum_q f_q^{(0)} - \frac{\beta}{2g} \tilde{f}_g^{(0)}$$



- Dashed blue: quark
- Dash-dotted red: gluon
- Vertical stripes: perturbative uncertainty $Q/2 < \mu < 2Q$
- Solid bands: hadronic uncertainties on matrix elements

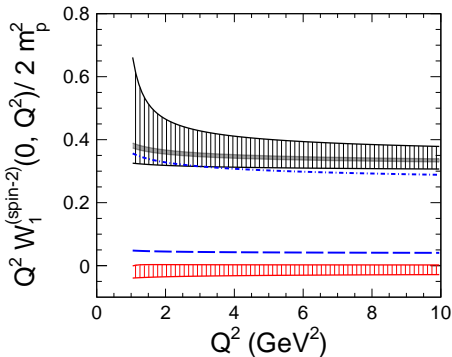
Large Q^2 behavior: Spin 2 contribution

- The *new* spin 2 result

$$\frac{Q^2}{2m_p^2} W_1^{(\text{spin-2})}(0, Q^2) = 2 \sum_q e_q^2 f_q^{(2)}(\mu) + \left(\sum_q e_q^2 \right) \frac{\alpha_s}{4\pi} \left(-\frac{5}{3} + \frac{4}{3} \log \frac{Q^2}{\mu^2} \right) f_g^{(2)}(\mu)$$

- quark and gluon matrix elements related by

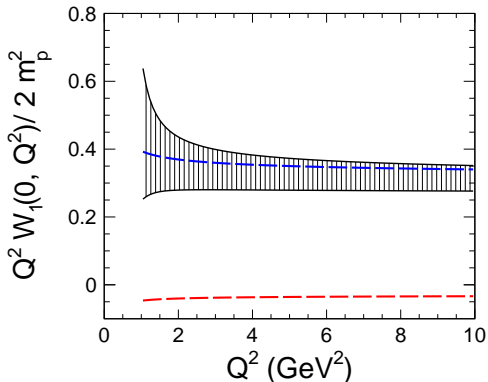
$$\sum f_q^{(2)}(\mu) + f_g^{(2)}(\mu) = 1$$



- Dashed blue: down quark
- Dash-dotted blue: up quark
- Red: Gluon contribution
- Vertical stripes: perturbative uncertainty $Q/2 < \mu < 2Q$
- Solid bands: hadronic uncertainty

Large Q^2 behavior: Total contribution

- The total contribution



- Dashed red: spin 0
- Dashed blue: spin 2
- Vertical stripes: total contribution with perturbative and hadronic errors added in quadrature

Two Photon exchange: small Q^2 and large Q^2

- Using NRQED we have control over low Q^2

$$W_1(0, Q^2) = 2a_p(2+a_p) + \frac{Q^2}{m_p^2} \left\{ \frac{2m_p^3 \bar{\beta}}{\alpha} - a_p - \frac{2}{3} \left[(1+a_p)^2 m_p^2 (r_M^p)^2 - m_p^2 (r_E^p)^2 \right] \right\} + \mathcal{O}(Q^4)$$

- Using OPE we have control over the high Q^2

$$\frac{Q^2}{2m_p^2} W_1^{(\text{spin}-0)}(0, Q^2) = -2 \sum_q e_q^2 f_q^{(0)} + \left(\sum_q e_q^2 \right) \frac{\alpha_s}{12\pi} \tilde{f}_g^{(0)}$$

$$\frac{Q^2}{2m_p^2} W_1^{(\text{spin}-2)}(0, Q^2) = 2 \sum_q e_q^2 f_q^{(2)}(\mu) + \left(\sum_q e_q^2 \right) \frac{\alpha_s}{4\pi} \left(-\frac{5}{3} + \frac{4}{3} \log \frac{Q^2}{\mu^2} \right) f_g^{(2)}(\mu)$$

- The problem, like the joke, is how to make a whole fish from a head and a tail...
- Before this work we had only the low Q^2 knowing the large Q^2 allows to connect the dots

Gluon operators



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 - 9 components of traceless symmetric tensor: $G^{\mu\alpha} G_{\alpha}^{\nu} - \frac{1}{4} G^{\alpha\beta} G_{\alpha\beta} g^{\mu\nu}$
chromomagnetic stress-energy tensor
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$$\text{For example } O^{0123} = G^{01} G^{23} + G^{03} G^{21} = E^1 B^1 - E^3 B^3$$

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- For protons: $\langle \text{Proton} | O^{\mu\alpha\nu\beta} | \text{Proton} \rangle = 0$
 What about $\langle \text{Medium} | O^{\mu\alpha\nu\beta} | \text{Medium} \rangle$?
 Solution looking for a problem...