

# Observing Geometrical Torsion

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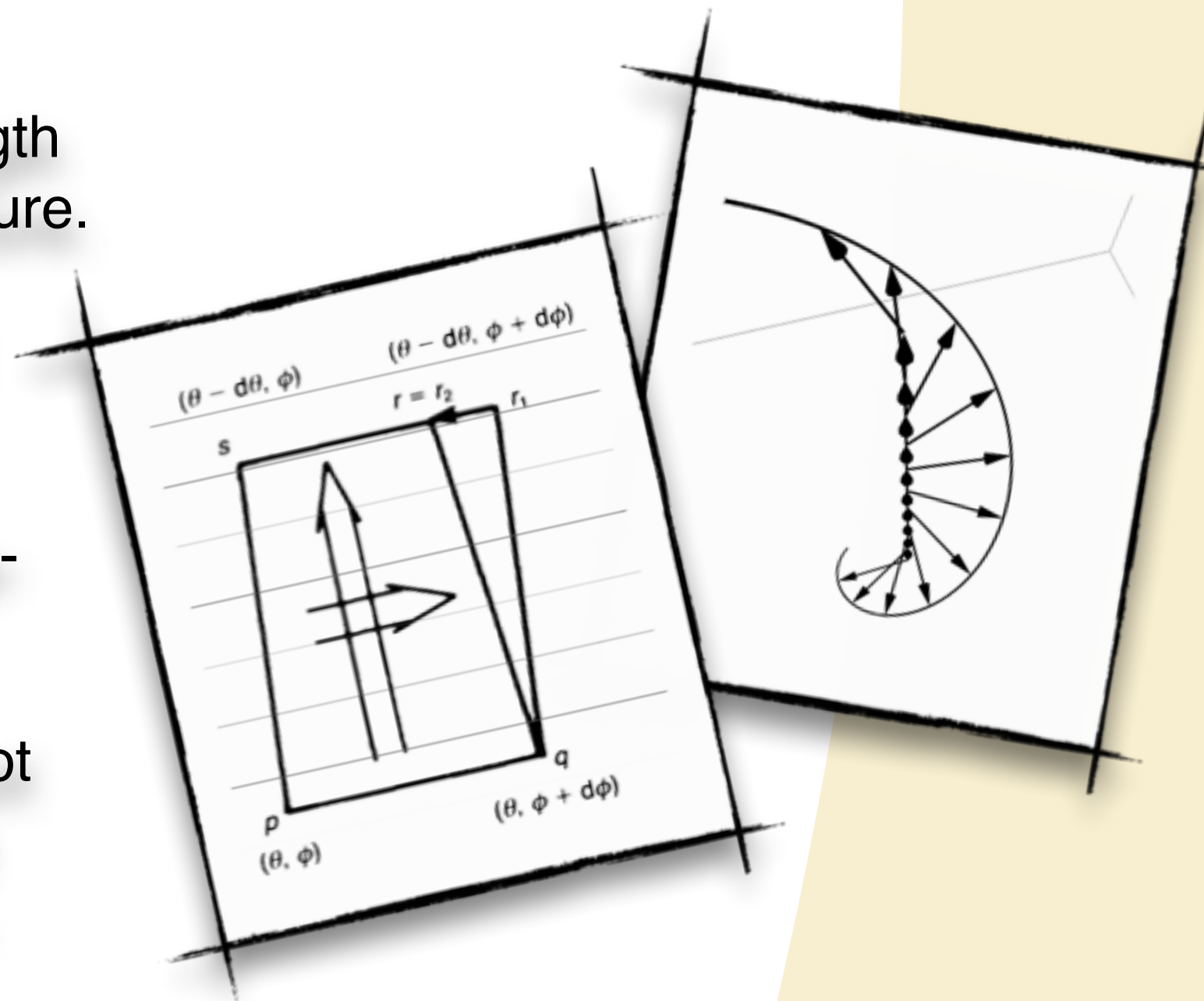
# Content

- Gravity with torsion, “Who, What, When, Where, Why”.
- Weyl invariant gravity + standard model with torsion.
- Possibility of production of torsion waves and detection.
- Perspective on future directions (preliminary results).
- Conclusions



# Some geometrical intuition

- Einstein: Gravity is a geometrical force, its strength given by space-time curvature.
- Cartan: adds an additional geometrical structure, separated from curvature, linked to “twisting” of space-time.
- Misconception: torsion is not just an external field. It is a geometrical universal field.



# The geometrical field

- As a geometrical field, parallel transport.

$$\nabla_{\dot{\gamma}} \dot{\gamma} = 0$$

$$\frac{d^2 x^\mu}{d\tau^2} = -\Gamma^\mu_{\alpha\beta} \frac{dx^\alpha}{d\tau} \frac{dx^\beta}{d\tau}$$

$$\text{if } \nabla_{\dot{\gamma}} g_{\mu\nu} = 0 \implies$$

$$g_{\mu\nu} = \left\{ \mathcal{P} \left[ \exp \left( \int_{\gamma} \Gamma^\lambda_{\sigma\mu} \dot{\gamma}^\sigma \right) \right] \right\}^2$$

- One gains Weyl symmetry. Curvature and geodesics become invariant under rescaling of proper lengths:

$$d\tau^2 \rightarrow \Omega^2(\tau) d\tau^2, g_{\mu\nu} \rightarrow \Omega^2(x) g_{\mu\nu}$$

$$\dot{\gamma} \rightarrow \Omega^{-1} \dot{\gamma}$$

Space-time dependent rescaling of dimension full quantities



# Coupling to matter

- Minimal couple prescription, in accordance with GR symmetries

$$\nabla_{\mu}(\Lambda(x) \cdot \Phi) = \Lambda(x) \cdot \nabla_{\mu} \Phi ,$$

$$\Lambda(x) \in SO(1, 3)$$

- Matter couple universally to connection:

- Couple microscopically:

$$\partial_{\mu} \Phi \rightarrow \nabla_{\mu} \Phi ,$$

$$\eta^{\mu\nu} \rightarrow g^{\mu\nu}$$

$$\nabla_{\mu}(\Omega(x)\mathbb{1} \cdot \Phi) = \Omega(x)\mathbb{1} \cdot \nabla_{\mu} \Phi ,$$

$$\Omega(x) \in \mathbb{R}^{+}$$

Chiral  
Symmetry

$$\mathcal{L}_{\psi} = \Gamma_{\alpha\beta\lambda} \bar{\psi} \gamma^{[\alpha} \gamma^{\beta} \gamma^{\lambda]} \psi$$

$$\mathcal{L}_{\Phi} = \Phi^{\dagger} T_{\mu} (2\partial^{\mu} + T^{\mu}) \Phi$$

Weyl  
Symmetry



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# The conformal extension of the standard model

- Scalars couple to the torsion trace. At least two scalars, Higgs and an additional dilaton.
- Fermions couple to skew symmetric torsion, but not to the trace. The torsion trace contribution cancels in the action.
- Gauge fields are conformal only in four dimensions. For D generic, conformal invariance breaks gauge symmetry.

$$L = \left( \frac{\Phi^2}{2\alpha^2} R - \frac{1}{2} \bar{\nabla}_\mu \Phi \bar{\nabla}^\mu \Phi - V(\Phi) \right)$$

$$L = \frac{i}{2} (\bar{\psi} \gamma^\mu \bar{\nabla}_\mu \psi - \bar{\nabla}_\mu \bar{\psi} \gamma^\mu \psi)$$

$$F_{\mu\nu} = \bar{\nabla}_{[\mu} A_{\nu]} \stackrel{D \rightarrow 4}{=} (dA)_{\mu\nu}$$

$$L = -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} + L_{int}$$





# The link between torsion and Weyl symmetry

- Why should torsion be linked to Weyl symmetry?

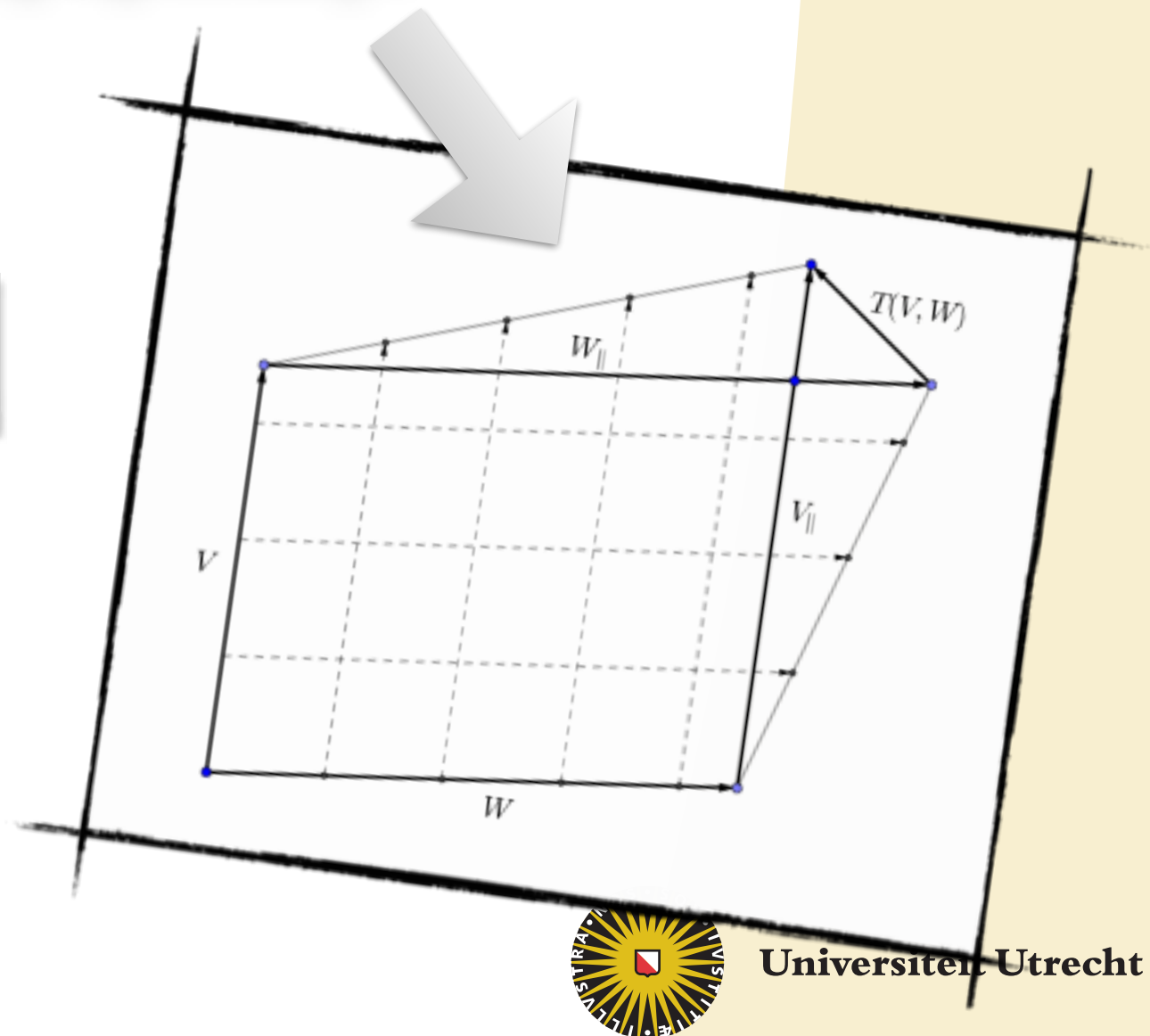
$$e_{\mu}^a e_{\nu}^b \eta_{ab} = g_{\mu\nu}$$

$$e_{\mu}^a \rightarrow e^{\theta(x)} e_{\mu}^a$$

$$\omega_b^a \rightarrow \omega_b^a$$

$$T^a \rightarrow T^a + e^a \wedge d\theta$$

- The torsion trace is naturally linked to scale transformations.
- Transforming torsion and vierbein leaves the Cartan connection invariant.



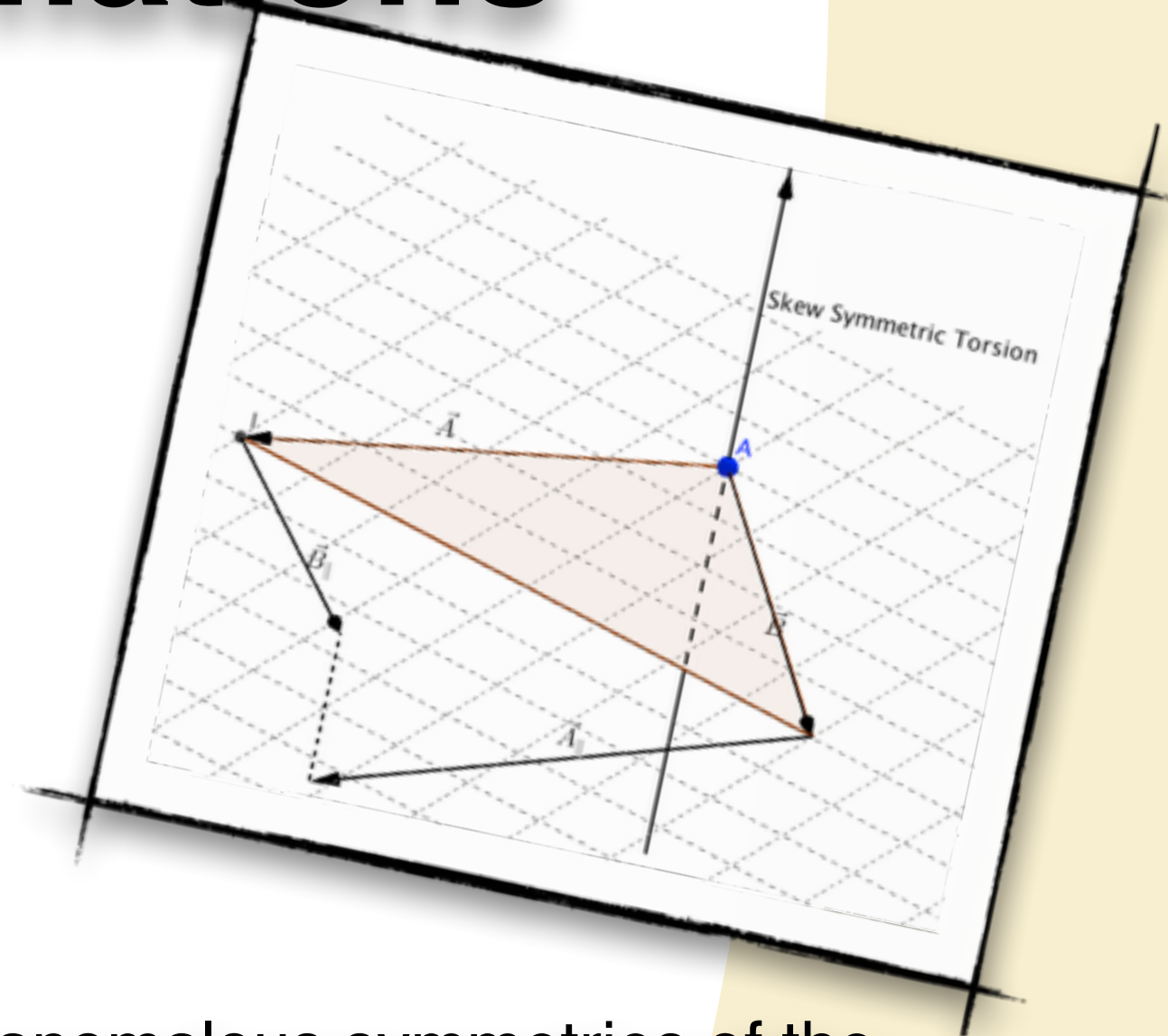
# Link to chiral transformations

- Geometrical breaking of the “right hand rule”.
- Skew symmetric torsion, couples to chiral fermionic current.



Einstein-Cartan-Kibble-Schiama  
Theory

- Geometrical fields, linked to two anomalous symmetries of the standard model. Maybe we can learn more about it by studying it carefully.





# Dynamical Torsion

- If a field couples through covariant derivatives

$$\mathcal{D}_\mu \Phi \in \mathcal{L}$$

$$\mathcal{F}_{\mu\nu} = [\mathcal{D}_\mu, \mathcal{D}_\nu] \in \Gamma_{1-loop}$$

- If torsion couples to matter in a universal manner, through Weyl and chiral charge, one loop effect will turn it dynamical.

$$\Gamma_{1-loop} \subset \left\{ \begin{array}{l} (\partial_\mu T_\nu - \partial_\nu T_\mu)^2 ; \\ (\partial_\mu \Sigma_\nu^* - \partial_\nu \Sigma_\mu^*)^2 \end{array} \right.$$

$$\Sigma^{*\mu} = \epsilon^{\mu\alpha\beta\gamma} \Gamma_{[\alpha\beta\gamma]}$$

+ $\mathcal{O}(R^2)$  operators

+Kinetic term  
for third  
Irreducible component

- Observations: in Colliders, via EFT interactions. Geometrically, via geodesics displacement.



# Sources (Torsion trace)

- If our intuition is correct, torsion trace couples to fields dilatation current.
- If scale symmetry is realised, dilatation current is conserved. Energy momentum is traceless. No torsion production.
- But we live in the broken phase, effectively there is a scalar mode, satisfying:

$$T_{\mu}^{\mu} = \partial_{\mu} D^{\mu}$$

$$\implies D^{\mu} = x^{\nu} T_{\nu}^{\mu}$$

$$T_{\mu} = \partial_{\mu} \theta$$

$$\square \theta = \frac{1}{M_P^2} T_{\mu}^{\mu} + \mathcal{O}(\theta^2)$$

Gravitational waves:

$$\square h_{ij} = \frac{1}{M_P^2} T_{ij} + \mathcal{O}(h^2)$$



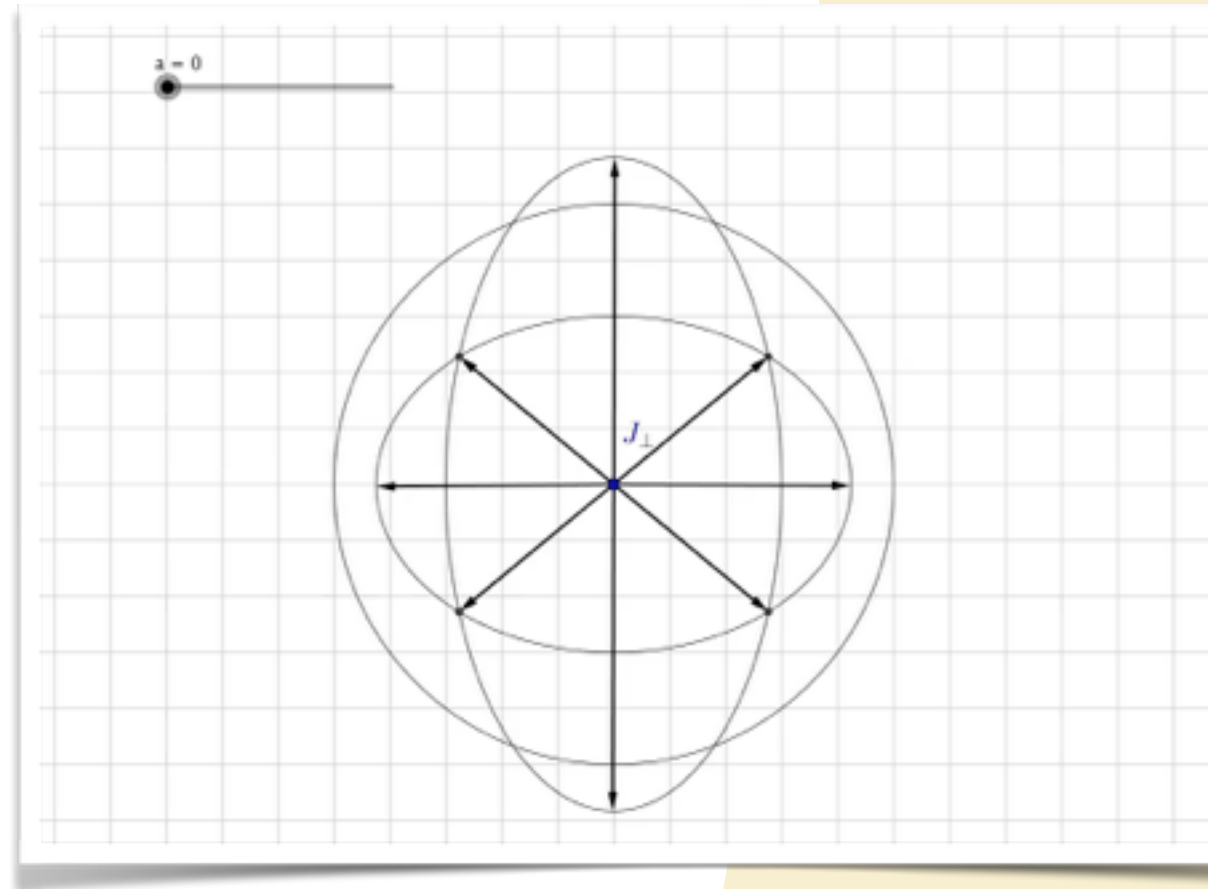
# Geometrical Detection

- Jacobi fields: give the displacement of nearby parallel geodesics, no skew symmetric torsion

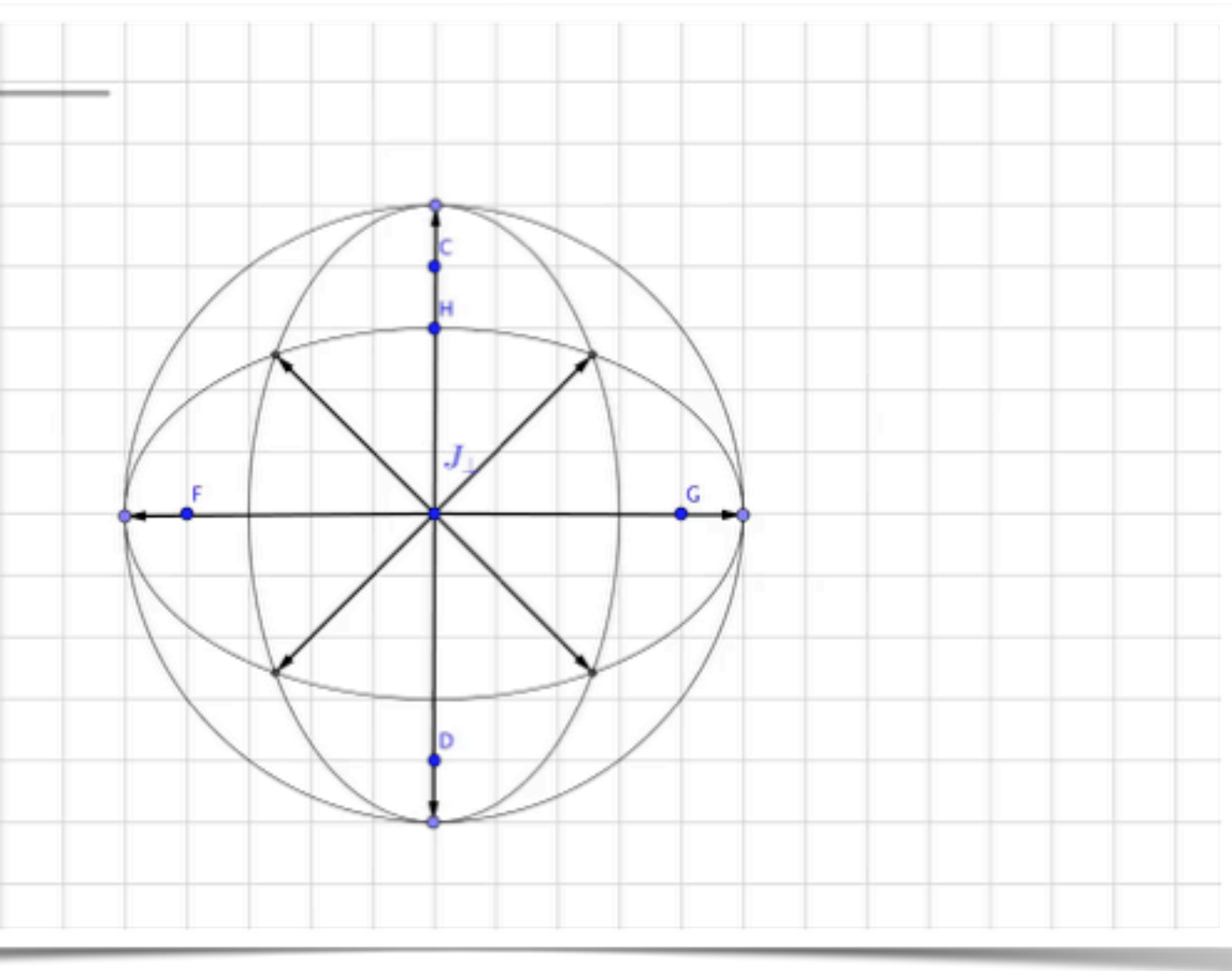
$$\nabla_{\dot{\gamma}} \nabla_{\dot{\gamma}} J_{\perp} + 2\nabla_{\dot{\gamma}} T(\dot{\gamma}, J_{\perp}) = R(\dot{\gamma}, J_{\perp})\dot{\gamma}$$

- Linearised equation:

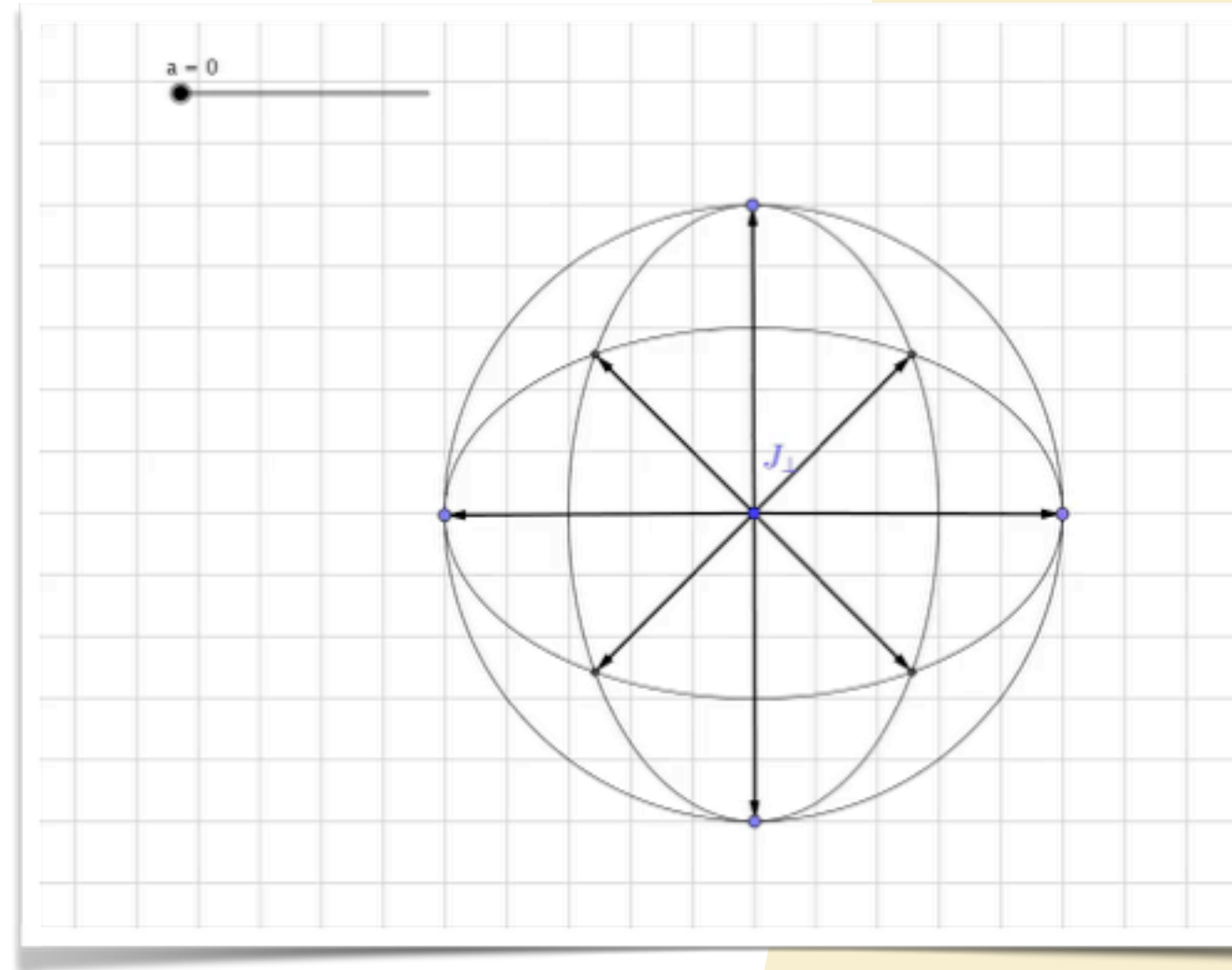
$$\frac{d^2 J^i}{dt^2} = \begin{cases} \frac{\ddot{h}_j^i}{2} J^j \\ J^j \partial_j T^i \\ -2\dot{Q}^i_{0j} J^j \end{cases} \rightarrow (\vec{J} \cdot \vec{\nabla})\vec{T} = (\vec{J} \cdot \vec{\nabla})\vec{\nabla}\theta + (\vec{J} \cdot \vec{\nabla})\vec{\nabla} \times \vec{\tau}$$



# How would it look like?



- Trace wave.



- Q wave.



# Perspective: Conformal symmetry breaking

- At the conformal point we can construct a self consistent Weyl invariant theory.
- What is such UV theory? Maybe  $SO(2,4)$  local. Symmetry breaking:

$$SO(2,4) \rightarrow SO(1,4) \rightarrow SO(1,3)$$

$$\mathcal{M} = \Omega^{-1} d\Omega, \quad \Omega \in SO(2,4)/SO(1,3)$$

$$\begin{aligned} [K^a, P^b] &= 2i\eta^{ab} D - 2i\Sigma^{ab}, \\ P^b, \Sigma^{ab}, K^a, D &\in SO(2,4) \end{aligned}$$

$$e^a_\mu = e^\theta \delta^a_\mu, \quad T^a = e^a \wedge d\theta$$

Low energy degrees of freedom





# Perspective: Weyl anomaly

- Conformally invariant theories present quantum anomalies in the Weyl symmetry Ward identities

$$\langle T_{\mu}^{\mu} \rangle \neq 0$$

$$\langle T_{\mu}^{\mu} \rangle = \langle \nabla_{\mu} \Pi^{\mu} \rangle$$

- Local anomaly proportional to topological invariant, Gauss-Bonnet integral

$$\langle T_{\mu}^{\mu} \rangle \propto \epsilon^{\alpha\beta\gamma\delta} \epsilon_{\mu\nu\lambda\sigma} R^{\mu\nu}{}_{\alpha\beta} R^{\lambda\sigma}{}_{\gamma\delta}$$

$$\langle \nabla_{\mu} \Pi^{\mu} \rangle \propto \epsilon^{\alpha\beta\gamma\delta} \epsilon_{\mu\nu\lambda\sigma} R^{\mu\nu}{}_{\alpha\beta} R^{\lambda\sigma}{}_{\gamma\delta}$$

- Solution valid at the conformal fixed point

$$\beta(\{\lambda_i\}) \rightarrow 0, \mu \rightarrow \infty$$

If trace is added in the geometry



# Conclusions

- Study of torsion field yields rich theoretical extensions of gravity and the standard model.
- Extend the symmetry of the theory to include Weyl symmetry (a gravitational gauge symmetry) and chiral transformations (which does not extend to space-time).
- In principle, direct detection is possible, but likely difficult.
- “Well defined” UV completions for these models (know UV theory). Can we use extended symmetries to learn more about quantum gravity?



# Thanks for attention

Questions?

