

RG Flows Across Dimensions and Holography

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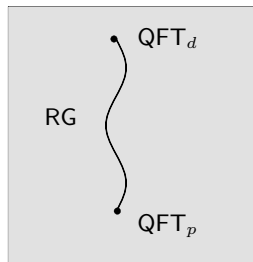
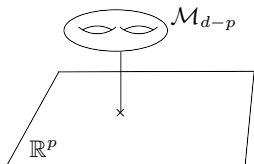
Geometry of String and Gauge Theories, CERN
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+ to appear

Francesco Benini, P. Marcos Cricigno
Francesco Azzurli, Vincent Min, Krzysztof Pilch, Orestis
Vasilakis, Alberto Zaffaroni

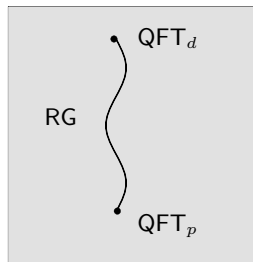
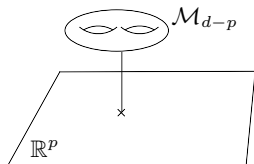
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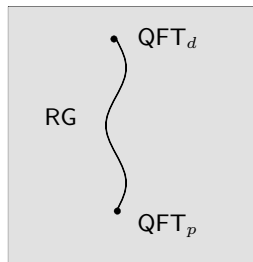
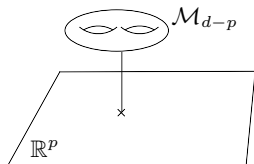


Questions:

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Employ supersymmetry to simplify the problem.

Motivation

- ▶ Construct and explore large classes of interacting p -dimensional superconformal field theories (SCFTs) obtained from a d -dimensional theory “compactified” on the manifold \mathcal{M}_{d-p} . [Vafa-Witten], [Witten], [Bershadsky-Johansen-Sadov-Vafa], [Klemm-Lerche-Mayr-Vafa-Warner], [Maldacena-Núñez], [Kapustin], [Gaiotto], [Gaiotto-Moore-Neitzke], [Dimofte-Gukov-Gaiotto], [Cecotti-Córdova-Vafa], [Gadde-Gukov-Putrov], [Benini-NB], [Bah-Beem-NB-Wecht], ...

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- ▶ This setup leads to interesting “correspondences” between the p -dimensional SCFT and a (“topological”) theory on \mathcal{M}_{d-p} . [Alday-Gaiotto-Tachikawa], [Gadde-Pomoni-Rastelli-Razamat], [Dimofte-Gaiotto-Gukov], [Cecotti-Córdova-Vafa], ...

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- ▶ Obtain insight into the original d -dimensional theory and RG flows across dimensions.
- ▶ The p -dimensional SCFTs typically admit a “large N ” limit and have holographic duals which can be explicitly constructed. New examples of AdS/CFT. [Maldacena-Núñez], ...

Tools

- ▶ Anomaly matching [’t Hooft]
- ▶ Anomaly polynomials [Alvarez-Gaumé-Ginsparg], [Witten], [Harvey-Minasian-Moore], ...
- ▶ Topological twists [Witten], ...
- ▶ a -maximization [Intriligator-Wecht], F -maximization [Jafferis],
[Closset-Dumitrescu-Festuccia-Komargodski-Seiberg], c -extremization [Benini-NB]
- ▶ Supersymmetric localization [Witten], [Moore], [Nekrasov], [Pestun],...
- ▶ Unitarity bounds [Hofman-Maldacena]
- ▶ Wrapped branes [Bershadsky-Sadov-Vafa], [Maldacena-Núñez], [Gauntlett-Waldram et al.], ...
- ▶ Holography [Maldacena], [GKP], [Witten], ...

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Disclaimer: Here I always take \mathcal{M}_{d-p} to be compact and (in holography) with an Einstein metric. Generalizations are possible and very interesting!

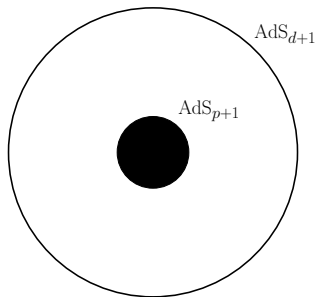
[*Anderson-Beem-NB-Rastelli*]; [*Gaiotto-Maldacena*], [*Bah*]

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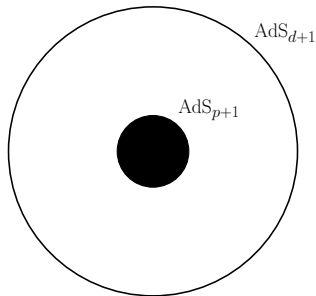
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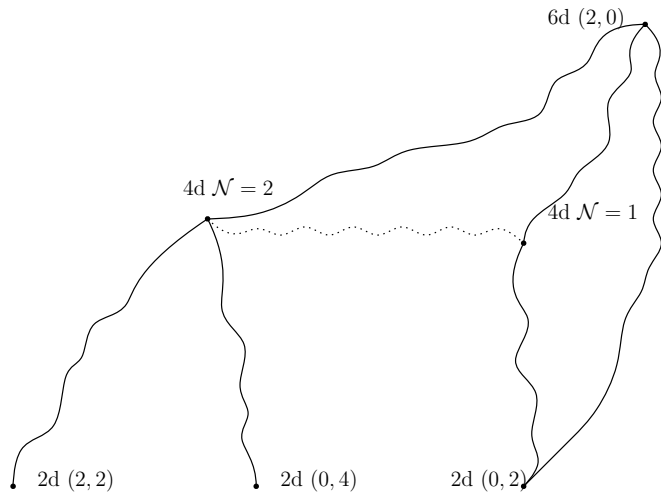
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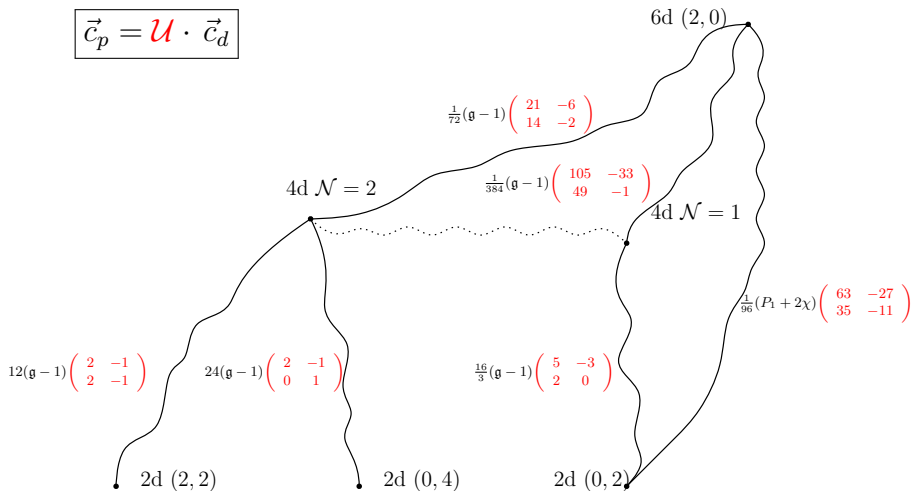
Numerous holographic RG flows of this type can be explicitly constructed and embedded in string theory by using branes wrapped on \mathcal{M}_{d-p} .

Universal flows across even dimensions



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$$\vec{C}_p = \mathcal{U} \cdot \vec{C}_d$$

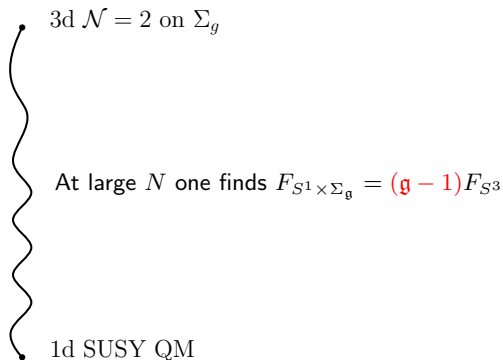


Universal flows across odd dimensions

- ▶ No 't Hooft anomalies \Rightarrow technically harder.

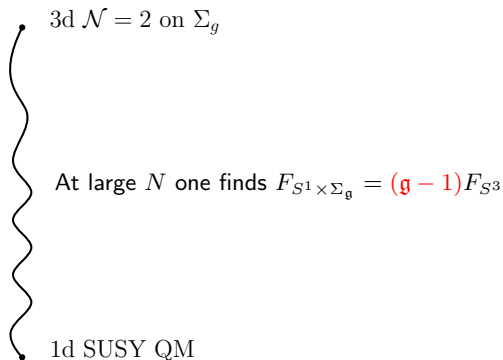
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- ▶ In holography this maps to a (magnetically charged) black hole in AdS_4 .
Microscopic counting of the BH entropy!

4d $\mathcal{N} = 1$ SCFTs on Σ_g

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Consider a general 4d $\mathcal{N} = 1$ SCFTs on $\mathbb{R}^2 \times \Sigma_g$ and perform a (partial) “topological twist”, i.e. use the R-symmetry to cancel the space-time curvature

[Witten]

$$\mathcal{A}_\mu^R = -\frac{1}{4}\omega_\mu, \quad \rightarrow \quad \tilde{\nabla}_\mu \epsilon = \left(\partial_\mu + \frac{1}{4}\omega_\mu + \mathcal{A}_\mu^R \right) \epsilon = \partial_\mu \epsilon = 0.$$

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General background fluxes in the Cartan of G_F do not break additional supersymmetry. The construction depends on the details of the 4d $\mathcal{N} = 1$ theory, i.e. **non-universal** \rightarrow **Rich families of 2d $\mathcal{N} = (0, 2)$ SCFTs.**

[Almuhairi-Polchinski], [Kutasov-Lin], [Franco-Lee-Vafa et al.], [Schäfer-Nameki-Weigand],

[Amariti-Cassia-Penati]...

4d $\mathcal{N} = 1$ SCFTs on Σ_g

The anomaly polynomials in 4d and 2d are

$$\mathcal{I}_6 = \frac{k_{RRR}}{6} c_1(\mathcal{F}_R^{(4)})^3 - \frac{k_R}{24} c_1(\mathcal{F}_R^{(4)}) p_1(\mathcal{T}_4) + I_6^F ,$$

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Superconformal Ward identities relate conformal and 't Hooft anomalies

[Anselmi-Freedman-Grisaru-Johansen]

$$a = \frac{9}{32} k_{RRR} - \frac{3}{32} k_R , \quad c = \frac{9}{32} k_{RRR} - \frac{5}{32} k_R , \quad c_r = 3k_{RR} , \quad c_r - c_l = k .$$

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Use ($\kappa = 1$ for S^2 , $\kappa = 0$ for T^2 , and $\kappa = -1$ for H^2)

$$\mathcal{F}_R^{(4)} \rightarrow \mathcal{F}_R^{(2)} - \frac{\kappa}{2} t_g, \quad \text{and} \quad \mathcal{I}_4 = \int_{\Sigma_g} \mathcal{I}_6,$$

to extract the 2d conformal anomalies

$$\boxed{\begin{pmatrix} c_r \\ c_l \end{pmatrix} = \frac{16}{3} (g-1) \begin{pmatrix} 5 & -3 \\ 2 & 0 \end{pmatrix} \begin{pmatrix} a \\ c \end{pmatrix}}$$

4d $\mathcal{N} = 1$ SCFTs on Σ_g

Comments:

- ▶ The R-charges have to be properly quantized (i.e. rational) in order to be able to perform the universal twist.
- ▶ Notice that 't Hooft anomaly matching amounts to

$$k_{RR} = (g-1)k_{RRR}, \quad k = (g-1)k_R.$$

- ▶ Unitarity implies that $g > 1$. In addition we should have

$$\frac{3}{5} < \frac{a}{c}, \quad \text{compatible with} \quad \frac{1}{2} \leq \frac{a}{c} \leq \frac{3}{2} \quad [\text{Hofman-Maldacena}].$$

Are there any interacting 4d $\mathcal{N} = 1$ SCFTs with $\frac{1}{2} < \frac{a}{c} < \frac{3}{5}$?

- ▶ For $a = c$ we have $c_r = c_l = \frac{32}{3}(g-1)a$. To be tested holographically!

The holographic dual

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There is a BPS black string solution of the 5d theory [Klemm-Sabra], [Benini-NB]

$$ds^2 = e^{2f(r)}(-dt^2 + dz^2 + dr^2) + e^{2h(r)} \frac{dx^2 + dy^2}{y^2}, \quad A = \frac{dx}{y}$$

- ▶ Analytic solution for $f(r)$ and $h(r)$. Here $\Sigma_{\mathfrak{g}} = \mathbb{H}^2/\Gamma$.
- ▶ Asymptotically locally AdS_5 background at $r \rightarrow \infty$.
- ▶ $AdS_3 \times \Sigma_{\mathfrak{g}}$ background at $r \rightarrow 0$.

The holographic dual

Using standard holographic technology one finds

$$c_r = c_l = \frac{3L_{AdS_3}}{2G_N^{(3)}} = \frac{32}{3}(\mathfrak{g} - 1) \frac{\pi L_{AdS_5}^3}{8G_N^{(5)}} = \frac{32}{3}(\mathfrak{g} - 1)a.$$

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Uplift of this simple solution to many **distinct** string and M-theory backgrounds by using supergravity uplift formulas. [Gauntlett-Varela]

Particular examples:

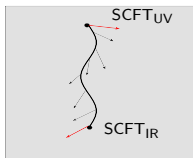
- ▶ IIB compactifications on SE_5 manifolds, i.e. 4d $\mathcal{N} = 1$ quiver gauge theories. [Klebanov-Witten], [Morrison-Plesser], [Gauntlett-Martelli-Sparks-Waldram]...
- ▶ M-theory compactifications of the Maldacena-Núñez type, i.e. 4d $\mathcal{N} = 1$ "class \mathcal{S} " SCFTs. [Bah-Beem-NB-Wecht]

4d $\mathcal{N} = 1$ on Σ_g

More general twists. Turn on fluxes for $G_F \rightarrow$ **non-universal flows**

$$\mathcal{A}_{\text{back}} = \mathcal{A}^R + b_i \mathcal{A}^i, \quad \rightarrow \quad R_{IR} = R_{UV} + \epsilon_i(b) F^i$$

Mixing between the flavor and R-symmetries along the RG flow!

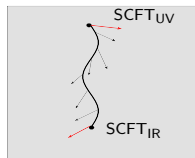


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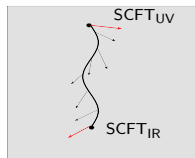
$$\Omega_\mu^{\text{tr}}(t) = \Omega_\mu + \sum_{i \neq R} t_i J_\mu^i .$$

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Construct a trial $c_r^{\text{tr}}(t)$ from the anomaly of the trial R-symmetry. For unitary SCFTs one finds **c-extremization** [Benini-NB]

$$\frac{\partial c_r^{\text{tr}}(t^*)}{\partial t^i} = 0, \quad \forall i \neq R, \quad \rightarrow \quad c_r^{\text{tr}}(t^*) = c_r .$$

Similar to a -maximization in 4d and F -maximization in 3d. [Intriligator-Wecht],

[Jafferis], [Closset-Dumitrescu-Festuccia-Komargodski-Seiberg]

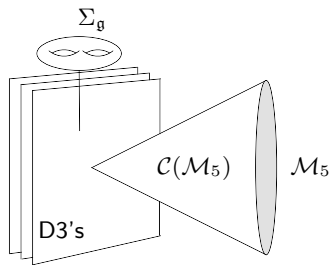
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Explicit examples from string theory: 4d $\mathcal{N} = 1$ SCFTs from D3-branes at conical singularities [Klebanov-Witten], [Morrison-Plesser],...

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- ▶ Transverse space $\mathbb{R}_+ \times \mathcal{M}_5 = \mathcal{C}(\mathcal{M}_5)$.
- ▶ \mathcal{M}_5 is Sasaki-Einstein $\Rightarrow \mathcal{C}(\mathcal{M}_5)$ is Calabi-Yau.
- ▶ $\mathcal{M}_5 = S^5 \Rightarrow \mathcal{N} = 4$ SYM.
- ▶ $\mathcal{M}_5 = Y^{1,0} \Rightarrow \mathcal{N} = 1$ KW.
- ▶ More examples based on $Y^{p,q}$
[Gauntlett-Martelli-Sparks-Waldram] and dP_n surfaces.



One can analyze in the same way the Leigh-Strassler $\mathcal{N} = 1$ mass deformation of $\mathcal{N} = 4$ SYM. [Gubser-Freedman-Pilch-Warner], [NB-Pilch-Vasilakis]

A new feature: “Baryonic” and R-symmetries can mix along such RG flows.

Holographic description

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For $b=0 \rightarrow$ universal twist.

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IIB supergravity: Solve the IIB supergravity supersymmetry variations

$$ds_{10}^2 = ds_{\text{AdS}_3}^2 + \frac{v^2 + v + 1}{4v} ds_{\Sigma_g}^2 + \frac{v^2 + v + 1}{4(v+1)} \left[d\theta^2 + \sin^2 \theta d\phi^2 + \frac{1}{v} (dw^2 + \sin^2 w d\nu^2) \right] + \frac{1}{4} \left(d\psi - \cos \theta d\phi - \cos w d\nu - \frac{dx_1}{x_2} \right)^2,$$
$$G_{(5)} = \text{vol}_{\text{AdS}_3} \wedge \left(\frac{(v+1)^2}{2v} \text{vol}_{\Sigma_g} + \frac{1}{2(v+1)} \left(v^2 \text{vol}_{S_{\theta\phi}^2} + \frac{1}{v} \text{vol}_{S_{w\nu}^2} \right) \right),$$

where $v = \frac{1+4pb}{1-4pb}$.

Holographic description

Example: $Y^{p,0}$ on $\Sigma_{g>1}$ with baryonic flux b (\mathbb{Z}_p orbifold of the KW theory).

Field theory: Apply anomaly matching and c -extremization to find

$$c_r = \frac{32}{3}(g-1)a(Y^{p,0}) + 24(g-1)p^3 b^2 N^2 - 2p(g-1)$$

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- ▶ α -twist with $\mathcal{N} = (2, 2)$. Background flux for $U(1)_R$. The central charges are

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Notice that $\frac{1}{3}\alpha + \frac{4}{3}\beta$ is the "universal" twist with 2d $\mathcal{N} = (0, 2)$ supersymmetry.

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- ▶ Unitarity implies that $g > 1$ and

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Reproduce the CFT central charges holographically.

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- ▶ No AdS_3 solution for the β -twist. Maybe no normalizable vacuum?
Similarities with the $\mathcal{N} = (4,4)$ twist of 4d $\mathcal{N} = 4$ SYM where one finds a 2d σ -model onto the Hitchin moduli space [Bershadsky-Johansen-Sadov-Vafa]?

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$$c_r = \frac{21}{128} (P_1 + 2\chi) a_{6d} = \frac{21}{256\pi^2} \text{Vol}(\mathcal{M}_4) a_{6d}$$

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There are many other possible twists for the $(2, 0)$ theory on various \mathcal{M}_4 . Generically one has to rely on c -extremization to compute the correct central charges. [Benini-NB], [Gadde-Gukov-Putrov]

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Turn a background field **only** for \mathcal{A}_μ^R to preserve 2 supercharges in 1d.

To study this theory use the "topologically twisted index" [Benini-Zaffaroni], [Closset-Kim], [Benini-Hristov-Zaffaroni]. Computed by localization in the same spirit as the S^3 partition function. [Kapustin-Willet-Yaakov], [Drukker-Marino-Putrov], [Jafferis-Klebanov-Pufu-Safdi]

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At large N a simple relation between the two partition functions [Hosseini-Zaffaroni]

$$F_{S^1 \times \Sigma_g}(\Delta_I, \mathbf{n}_I) = (\mathfrak{g} - 1)F_{S^3}(\Delta_I/\pi) + \sum_I \left(\frac{\mathbf{n}_I}{1 - \mathfrak{g}} - \frac{\Delta_I}{\pi} \right) \frac{\pi}{2} \partial_{\Delta_I} F_{S^3}(\Delta_I/\pi)$$

Where: $\Delta_I \rightarrow$ chemical potentials for global symmetries; $\mathbf{n}_I \rightarrow$ magnetic charges.

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The universal twist amounts to: $\mathbf{n}_I = (1 - \mathfrak{g})\Delta_I/\pi$

$$F_{S^1 \times \Sigma_g} = (\mathfrak{g} - 1)F_{S^3}$$

Black holes in AdS_4

A simple supersymmetric BH solution of 4d minimal gauged supergravity

[Romans], [Caldarelli-Klemm]

$$ds_4^2 = - \left(\rho - \frac{1}{2\rho} \right)^2 dt^2 + \left(\rho - \frac{1}{2\rho} \right)^{-2} d\rho^2 + \rho^2 ds_{\Sigma_g}^2, \quad F = \frac{1}{2\sqrt{2}} \text{vol}_{\Sigma_g}$$

Asymptotic to AdS_4 for $\rho \rightarrow \infty$ and to $AdS_2 \times \Sigma_g$ for $\rho \rightarrow \frac{1}{\sqrt{2}}$.

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Uplift to M-theory: [Gauntlett-Kim-Waldram], [Gauntlett-Varela]

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In massive IIA: **a new solution!** Deformation of a recently constructed AdS_4 vacuum in massive IIA [Guarino-Jafferis-Varela], [Fluder-Sparks]

$$ds_{10}^2 = e^{2\lambda} L^2 (ds_4^2 + ds_6^2)$$

with

$$ds_6^2 = \omega_0^2 [e^{\varphi-2\phi} X^{-1} d\alpha^2 + \sin^2 \alpha (\Delta_1^{-1} ds_{KE_4}^2 + X^{-1} \Delta_2^{-1} \eta^2)]$$
$$e^{2\lambda} \equiv (\cos(2\alpha) + 3)^{1/2} (\cos(2\alpha) + 5)^{1/8},$$

and $L, \omega_0, \varphi, \phi$ - constants. Nontrivial F_2, H_3 , and F_4 fluxes in massive IIA.

Large class of new massive IIA solutions with CFT duals.

3d $\mathcal{N} = 2$ on $\Sigma_{\mathfrak{g}}$

For the black hole entropy one finds

$$S_{\text{BH}} = \frac{\pi}{2G_N^{(4)}}(\mathfrak{g} - 1) = (\mathfrak{g} - 1)F_{S^3}$$

Thus $F_{S^1 \times \Sigma_{\mathfrak{g}}}$ computes the entropy of these BHs!

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Comments:

- ▶ At large N one has $F_{S^1 \times \Sigma_g} \sim N^{3/2}$ in M-theory and $F_{S^1 \times \Sigma_g} \sim N^{5/3}$ in massive IIA.
- ▶ To generalize this setup turn on background flux for the G_F global symmetry of the CFT. Find the correct R-symmetry in the IR by extremizing the "twisted index". [Benini-Hristov-Zaffaroni]
- ▶ The holographic dual construction is realized by M2-branes at the tip of a conical singularity wrapping Σ_g . A large landscape of AdS_2 vacua in M-theory. [Azzurli-NB-Crichigno-Min-Zaffaroni]
- ▶ One can also add electric charges to these 4d black holes and modify correspondingly the "twisted index" [Benini-Hristov-Zaffaroni], [in progress]

Other flows

- ▶ This general setup is applicable to SCFTs in other dimensions and on other manifolds. A prominent example is the 6d $(2, 0)$ theory on 4-manifolds. [Benini-NB], [Gadde-Gukov-Putrov], [Ganor]...
- ▶ The field theory tools are less developed for 5d $\mathcal{N} = 1$ SCFTs on $\mathbb{R}^3 \times \Sigma_g$ or $\mathbb{R}^2 \times \mathcal{M}_3$ (absence of anomalies or localization results).
- ▶ A holographic analysis in "minimal" gauged supergravity always yields an AdS vacuum. Suggestive results and "predictions" from holography.
- ▶ Interpretation in terms of wrapped branes in string/M-theory.

Summary

- ▶ Evidence for many new SCFTs arising from RG flows across dimensions.
- ▶ Understanding of some of their properties through nonperturbative QFT tools.
- ▶ Explicit "top-down" holographic constructions that are dual to the SCFTs at hand.
- ▶ Useful spin-off: microscopic entropy counting for AdS black holes.

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- ▶ Holographic/geometric dual of c -extremization. Is there some "generalized volume" minimization? [Martelli-Sparks-Yau]

THANK YOU!