

Search for the production of dark matter in association with top-quarks in the single-lepton final state in proton-proton collisions at $\sqrt{s} = 8\text{TeV}$

João A. Gonçalves

Paper may be found here

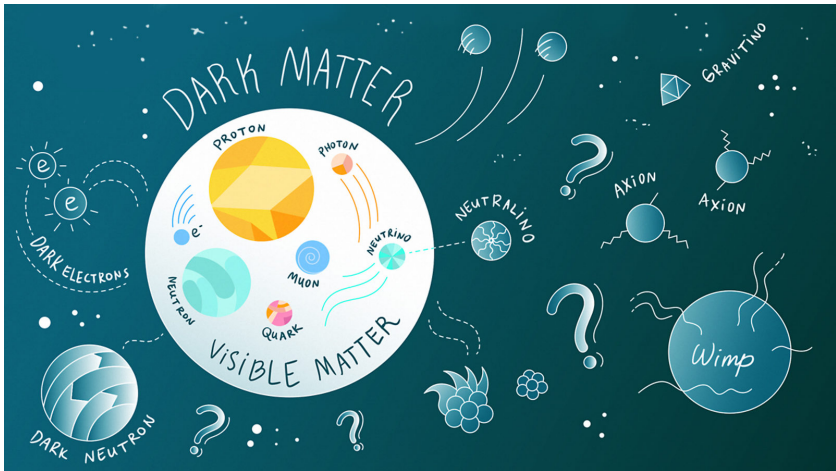
IST
LIP
CMS

05/06/2017

Outline

- 1 Introduction
- 2 The CMS detector
- 3 Data and simulated samples
- 4 Object reconstruction
- 5 Event selection
- 6 Background estimation
- 7 Systematic uncertainties
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Introduction



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Γ represents the type of interaction.

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In this scenario we can have the production of DM particles associated with :

- a monojet.
- a monophoton.
- a W/Z boson.

(1) (3) (4) (5) (6) (7) (8)

(4)

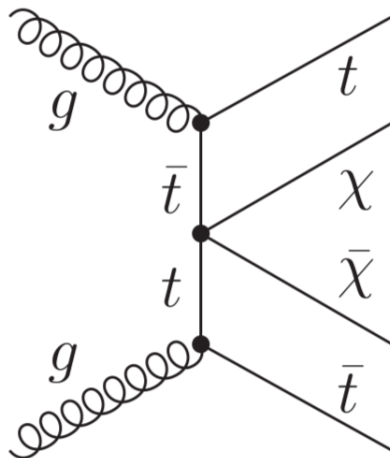
(9)

Introduction

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$$L_{\text{int}} = \frac{m_q}{M_*^3} \bar{q}q\bar{\chi}\chi.$$

Introduction



The CMS detector

CMS DETECTOR

Total weight : 14,000 tonnes
 Overall diameter : 15.0 m
 Overall length : 28.7 m
 Magnetic field : 3.8 T

STEEL RETURN YOKE
 12,500 tonnes

SILICON TRACKERS

Pixel ($100 \times 150 \mu\text{m}$) $\sim 16\text{m}^2 \sim 66\text{M}$ channels
 Microstrips ($80 \times 180 \mu\text{m}$) $\sim 200\text{m}^2 \sim 9.6\text{M}$ channels

SUPERCONDUCTING SOLENOID

Niobium titanium coil carrying $\sim 18,000\text{A}$

MUON CHAMBERS

Barrel: 250 Drift Tube, 480 Resistive Plate Chambers
 Endcaps: 468 Cathode Strip, 432 Resistive Plate Chambers

PRESHOWER

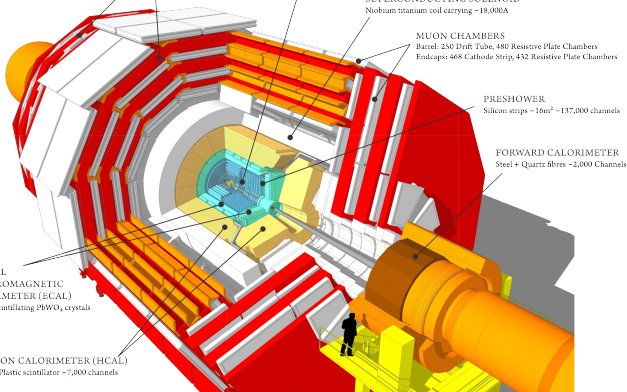
Silicon strips $\sim 16\text{m}^2 \sim 137,000$ channels

FORWARD CALORIMETER

Steel + Quartz fibres $\sim 2,000$ Channels

CRYSTAL
 ELECTROMAGNETIC
 CALORIMETER (ECAL)
 $\sim 76,000$ scintillating PbWO₄ crystals

HADRON CALORIMETER (HCAL)
 Brass + Plastic scintillator $\sim 7,000$ channels



More on the detector and kinematic coordinate system at [13]

Data and simulated samples

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- Single-muon trigger $p_T > 24\text{GeV}$.
- Efficiencies of these triggers in data and simulation are compared and measured using a tag-probe method : [14].



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- Leading order matrix element generator using the CTEQ6L1 PDF's : [16].
- The dominant backgrounds are $t\bar{t}$ +jets, $t\bar{t} + \gamma/W/Z$, W +jets, single top, diboson and Drell-Yan events.

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- The WW background is generated with PYTHIA v6.424 [\[18\]](#).

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- Cross sections of $t\bar{t} + jets$ [20] and $W/Z + jets$ [21] are calculated at NNLO.
- Other backgrounds are calculated at NLO (single top [22], $t\bar{t} + Z$ [23], $t\bar{t} + W$ [24], $t\bar{t} + \gamma$ [25] and diboson [26]).

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- A particle-flow (PF) based event reconstruction, taking into account information from all subdetectros, is used. [27] [28]

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- Primary vertices are reconstructed using a deterministic annealing filter algorithm, and are defined as the vertex with the largest sum of squares of the p_T of the tracks associated to it. [29]

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- Electrons and muon efficiencies are measured via the tag-and-probe technique using inclusive samples of leptonically decaying Z bosons, from data and simulation.




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Jets and CSV

- Jets are reconstructed from PF candidates, that are clustered with the anti- k_T algorithm [\[33\]](#) with $R=0.5$ using FASTJET. [\[34\]](#).

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- A jet is b-tagged if the discriminant output is greater than 0.679 and $|\eta| < 2.4$.



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- $E_T^{miss} > 160$ GeV.

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- Other backgrounds include single top, Drell-Yan and diboson. QCD multijet background is negligible within these requirements.

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- $M_T > 160 \text{ GeV}$ cut is applied.
- The dominant background becomes dileptonic $t\bar{t}$ events, with one lepton unobserved.

Event selection

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- The variable M_{T2}^W is used to further reduce this background.

$$M_{T2}^W = \min \left(m_y \text{ consistent with: } \left\{ \begin{array}{l} \vec{p}_1^T + \vec{p}_2^T = \vec{p}_T^{\text{miss}}, p_1^2 = 0, (p_1 + p_\ell)^2 = p_2^2 = M_W^2, \\ (p_1 + p_\ell + p_{b1})^2 = (p_2 + p_{b2})^2 = m_y^2 \end{array} \right\} \right)$$

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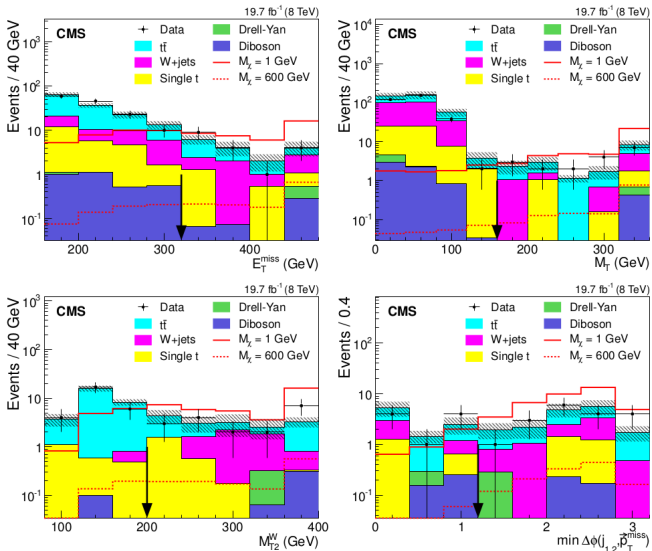
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- These selection cuts are optimized based on the expected significance for DM masses between 1 and 1000 GeV.

Event selection

Variable distributions



Background estimation

- Scale factors are applied to $t\bar{t}$ + jets and W + jets events dominant backgrounds to correct discrepancies between data and simulation.

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- SF's are obtained by matching data to the M_T distribution in the CR1, and the E_T^{miss} in the CR2.
- The resulting SF's were $1.11 \pm 0.02(stat)$ and $1.26 \pm 0.06(stat)$ for $t\bar{t}$ and W respectively.



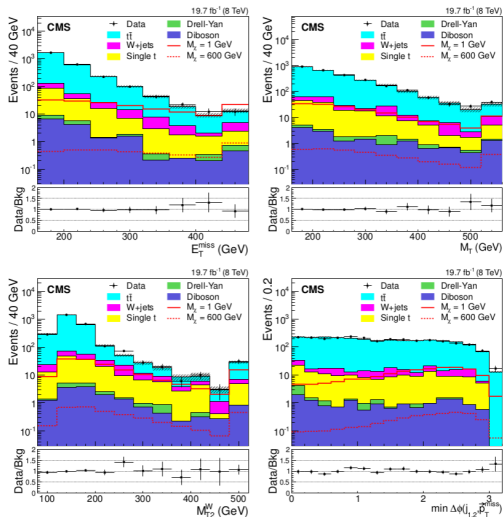
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- These scale factors are later propagated to the signal region to estimate the background.



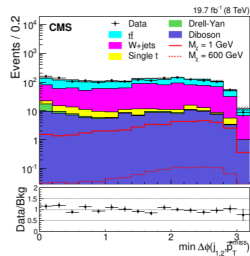
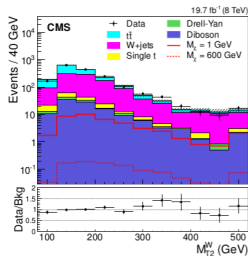
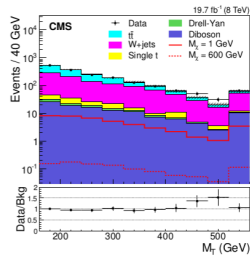
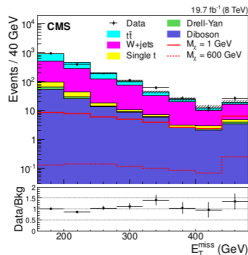
Background Estimation

Distributions in control region 1



Background Estimation

Distributions in control region 2



Systematic uncertainties

Source of systematic uncertainties	Relative uncertainty on total background (%)
50% normalization uncert. of other bkg in deriving SFs	10
SF_{W+jets} (CR tests)	13
$t\bar{t}$ +jets top-quark p_T reweighting	3.9
Jet energy scale	4.0
Jet energy resolution	3.0
b-tagging correction factor (heavy flavour)	1.0
b-tagging correction factor (light flavour)	1.8
Pileup model	2.0
PDF	2.6

Results

Source	Yield ($\pm\text{stat} \pm\text{syst}$)
$t\bar{t}$	$8.2 \pm 0.6 \pm 1.9$
W	$5.2 \pm 1.8 \pm 2.1$
Single top	$2.3 \pm 1.1 \pm 1.1$
Diboson	$0.5 \pm 0.2 \pm 0.2$
Drell–Yan	$0.3 \pm 0.3 \pm 0.1$
Total Bkg	$16.4 \pm 2.2 \pm 2.9$
Data	18

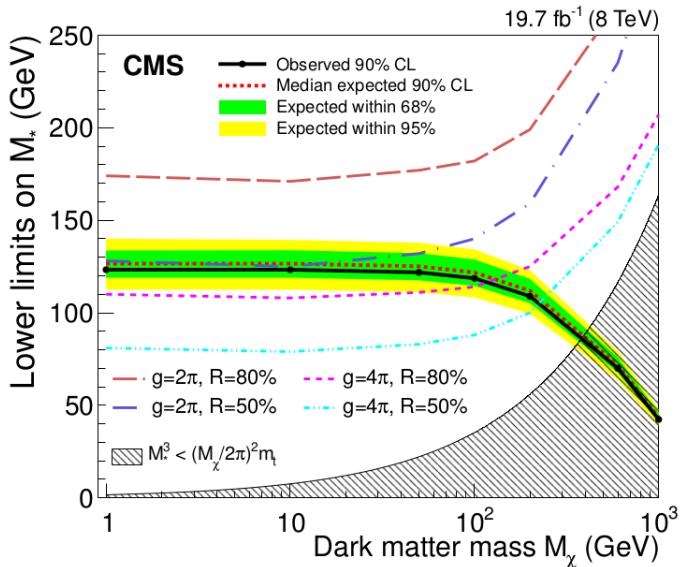
For $M_\chi=1$ and interaction scale $M_* = 100$ GeV.
No excess of events is observed.

Results

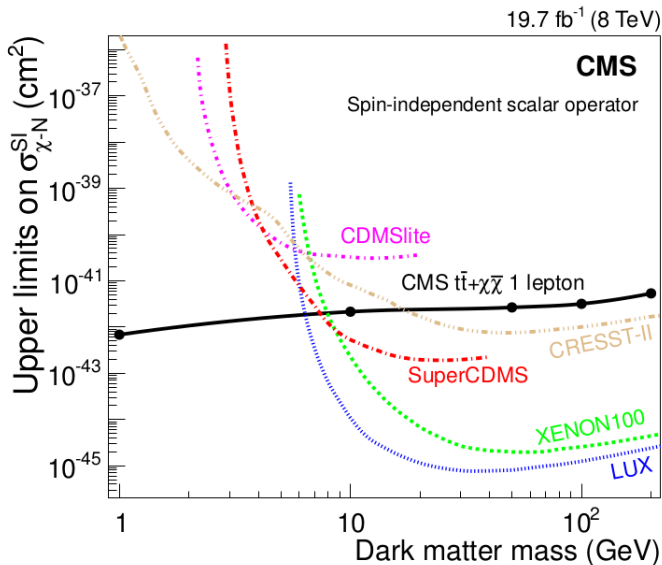
M_χ (GeV)	Yield (\pm stat \pm syst)	Signal efficiency (%) (\pm stat \pm syst)	$\sigma_{\text{exp}}^{\text{lim}}$ (fb)	$\sigma_{\text{obs}}^{\text{lim}}$ (fb)
1	$38.3 \pm 0.7 \pm 2.1$	$1.01 \pm 0.02 \pm 0.05$	47^{+21}_{-13}	55
10	$37.8 \pm 0.7 \pm 2.1$	$1.01 \pm 0.02 \pm 0.05$	46^{+21}_{-13}	54
50	$35.1 \pm 0.6 \pm 1.9$	$1.20 \pm 0.02 \pm 0.06$	39^{+18}_{-11}	45
100	$30.1 \pm 0.4 \pm 1.7$	$1.46 \pm 0.02 \pm 0.07$	32^{+14}_{-9}	37
200	$18.0 \pm 0.2 \pm 1.0$	$1.73 \pm 0.02 \pm 0.08$	27^{+12}_{-8}	32
600	$1.26 \pm 0.02 \pm 0.07$	$2.40 \pm 0.03 \pm 0.11$	19^{+9}_{-6}	23
1000	$0.062 \pm 0.001 \pm 0.003$	$2.76 \pm 0.04 \pm 0.13$	17^{+8}_{-5}	20

Confidence levels are obtained through a modified-frequentist CL_s method [\[47\]](#) [\[48\]](#), with both statistical and systematic uncertainties taken into account.

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- Lower limits on the interaction scale are set, between dark matter particles and top quarks in an EFT scalar theory.
- In the case of an s-channel mediator the limits for which the EFT approximation is valid, are interpreted as limits on the dark matter-nucleon scattering cross sections.

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- A search for the production of dark matter particles in association with top quarks in single-lepton events with the CMS detector at the LHC, using proton-proton collision data.
- Cross sections larger than 20 to 55 fb are excluded at 90% CL, for a dark matter mass from 1 to 1000 GeV.
- Lower limits on the interaction scale are set, between dark matter particles and top quarks in an EFT scalar theory.
- In the case of an s-channel mediator the limits for which the EFT approximation is valid, are interpreted as limits on the dark matter-nucleon scattering cross sections.
- For masses lower than 6 GeV, more stringent limits are obtained than from direct dark matter searches.

Thank you for your attention!
Any questions?

