A Higgs-Eye View of the Cosmos

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A Higgs-Eye View of the Cosmos

What does cosmology have to do with electroweak symmetry breaking and the Higgs?

Why might the Higgs lead us to understand:

• The nature of dark matter?
• The origin of the matter-antimatter imbalance?
• The ultimate fate of the Universe?

And recent theoretical ideas on how cosmology may explain the origin of the Higgs mass...
Whether in the laboratory or the cosmos, we model nature with fields. Here we will only consider bosonic fields, with no associated direction (scalars).
Out in the Fields

Whether in the laboratory or the cosmos, we model nature with fields

\[ \phi(x, t) \]

5 years ago at the LHC fluctuations of the only known potentially fundamental scalar field were discovered. The infamous Higgs boson!
(By the way, Happy 5\textsuperscript{th} Birthday Higgs Boson!)
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Right where you are sitting now!
Out in the Fields

While the human history of the Higgs field is very interesting, the true, cosmological, history of the Higgs field

\[ H \]

is a story that ties together the physics of superconductors to the mysteries of the origins of matter in the distant past and even determines the ultimate fate of the Universe.

Lets start with the superconductors...
Why are you living inside a Superconductor?
Ginzburg-Landau

- The G-L Theory of superconductivity involves a complex scalar field \( \Phi \) and the photon (magnetic vector potential) \( A \).

\[
F = \left| (\nabla + 2ieA) \Phi \right|^2 + m^2(T) |\Phi|^2 + \lambda |\Phi|^4 + \ldots
\]

- Where the mass depends on the temperature.
At high temperatures the mass-squared is positive:

\[ F = |(\nabla + 2ieA)\Phi|^2 + m^2(T)|\Phi|^2 + \lambda|\Phi|^4 + \ldots \]

- At high temperatures the mass-squared is positive:

- Just a hot metal.
Ginzburg-Landau

\[ F = \left| (\nabla + 2ieA) \Phi \right|^2 \]
\[ + m^2(T) |\Phi|^2 + \lambda |\Phi|^4 + \ldots \]

- At the critical temperature the mass-squared vanishes:

- Strange theory with massless fluctuations.
Below the critical temperature the mass-squared is negative:

\[ F = |(\nabla + 2ieA)\Phi|^2 + m^2(T)|\Phi|^2 + \lambda|\Phi|^4 + \ldots \]

- Below the critical temperature the mass-squared is negative:

  ![Diagram with a potential well]

- Photon has become massive: \( m_A \sim e\langle \Phi \rangle \)
Ginzburg-Landau

• Thus a simplified theoretical picture of a superconductor going through the critical temperature is

• Where in the lab we may control T.
What does this have to do with the Higgs Boson?
Higgs Mechanism

• The Higgs sector of the Standard Model involves the Higgs field and the gauge fields $H$ and $W^a_\mu$.

• The Lagrangian for this theory is

$$\mathcal{L} = \left| (\partial_\mu + ig\sigma^a W^a_\mu) H \right|^2 - m^2(T) |H|^2 - \lambda(T) |H|^4 + \ldots$$

• This is just the relativistic non-Abelian version of Ginzburg-Landau.
Higgs Mechanism

\[ L = \left| \left( \partial_\mu + ig \sigma^a W^a_\mu \right) H \right|^2 \]

\[ - m^2(T) |H|^2 - \lambda(T) |H|^4 + \ldots \]

- At high temperatures the mass-squared is positive:

- Just a hot relativistic gas.
At the critical temperature the mass-squared vanishes:

At the critical temperature the mass-squared vanishes:

Higgs boson is massless at this point.
Below the critical temperature the mass-squared is negative:

\[ \mathcal{L} = \left| \left( \partial_\mu + ig\sigma^a W^a_\mu \right) H \right|^2 - m^2(T) |H|^2 - \lambda(T) |H|^4 + \ldots \]

- Below the critical temperature the mass-squared is negative:
  
- Gauge bosons become massive: \( M_W \sim g\langle H \rangle \)
Thus the Higgs sector of the Standard Model is a relativistic superconductor!
The physical excitations of the Higgs field correspond to the Higgs boson:

Introducing Mr Higgs.

The discovery of this excitation put the entire Standard Model Higgs sector on a very firm footing. Allows us to deduce the quartic interactions (assuming the Standard Model).
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**Introducing Mr Higgs.**

The discovery of this excitation put the entire Standard Model Higgs sector on a very firm footing. Allows us to deduce the quartic interactions (assuming the Standard Model).

**Analogy with superconductors breaks down here, as there is no analogous weakly coupled light excitation, but breaking of gauge symmetry still analogous.**
Introducing Mr Higgs.

The physical excitations of the Higgs field correspond to the Higgs boson:

Now that we finally have this last piece of information, we can take the next steps into the unknown....

Where first?

The discovery of this excitation put the entire Standard Model Higgs sector on a very firm footing. Allows us to deduce the quartic interactions (assuming the Standard Model).
Dark Times

Evidence for dark matter is now overwhelming

- Rotation curves
- CMB
- Large scale structure
- Velocity dispersions
- Gravitational lensing (Bullet Cluster)
- ....

Yet we have no clue what it is at the particle level!
The Higgs portal term is the only renormalizable interaction between a Standard Model particle and a new stable neutral particle

\[ \lambda |H|^2 X^2 \]

This property is known as the “Higgs Portal”. Since we know other stable neutral particles exist (Dark Matter) then the Higgs portal could be the only window into the dark sector!
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Dark Times

This means that we can search for new dark sector states by searching for exotic Higgs decays

Or other quantum modifications of Higgs properties

We don’t know what’s out there, but there is something, and the Higgs is a great place to look.
The laboratory search for dark matter will hit a difficult background, called the “neutrino floor”.

LHC Searches can cover otherwise difficult territory in the search for Higgs boson interactions with dark matter.

A Portal to Hidden Worlds...

**Courtesy:** Phil Harris
The laboratory search for dark matter will hit a difficult background, called the “neutrino floor”.

LHC Searches can cover otherwise difficult territory in the search for Higgs boson interactions with dark matter.

It’s not guaranteed, but there is good reason to believe some of the dark sector may show its face first through the Higgs portal.

A Portal to Hidden Worlds…

The laboratory search for dark matter will hit a difficult background, called the “neutrino floor”.

Courtesy: Phil Harris
What’s the Matter?

To generate more matter than antimatter, would have to satisfy Sakharov conditions:

I demand:
• Baryon-number violation
• CP-violation
• Out-of-equilibrium processes

If the phase transition had been strongly first order we could have satisfied these conditions.

But it was not (we think) ...
What’s the Matter?

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But it was not (we think) ...

Relatively minor modifications of the Higgs sector of Standard Model can actually generate the observed matter-antimatter imbalance. I.e. You and me!
Ginzburg-Landau is just a phenomenological model, with no explanation of parameters. The macroscopic parameters follow from the detailed microscopic BCS theory (Gor’kov) and there are no surprises.

The order parameter at zero temperature is of the typical scale associated with underlying microscopic parameters.
Performing the same exercise with the Higgs field.

We can look to see if the symmetry breaking is like Ginzburg-Landau

• Direct analogy, would have no light Higgs boson: Experimentally excluded.

• Perhaps not directly analogous, but similar composite story: Study the Higgs...
The size of the Higgs boson...

- Probing deep into the Higgs boson. Light purple lines show current constraints.

Measurements are different ways of probing the “compositeness of the Higgs”.

No substructure has yet shown up...

\[ \lambda_h \approx 10^{-17} \text{ m} \]

\[ \lambda_{10 \text{ TeV}} \approx 10^{-19} \text{ m} \]
The Elephant in the Room

We expect the Higgs model is phenomenological, just like G-L. But something totally different seems to be going on.

There is a hierarchy between the model parameters and the microscopic parameters. Furthermore, this hierarchy is not protected by any symmetry: Quantum corrections do not respect such a hierarchy.
The Elephant in the Room

- Microscopic theory at high energy.
- If states couple to the Higgs Boson:

\[ H \xrightarrow{\text{Quantum Correction}} H^\dagger \rightarrow m_H^2 \sim \text{Loop} \times \Lambda_{NP}^2 \]

- If \( \Lambda_{NP} \gg \text{TeV} \) then howcome the Higgs mass, and the weak scale are below TeV?

- One answer: \( V_H = \tilde{M}_H^2 |H|^2 + \Lambda_{NP}^2 |H|^2 \)

\[ \tilde{M}_H^2 + \Lambda_{NP}^2 \sim \text{small} \]
The Elephant in the Room

• Microscopic theory at high energy.

• If states couple to the Higgs Boson:

\[ m_H^2 \sim \text{Loop} \times \Lambda_{NP}^2 \]

• If \( \Lambda_{NP} \gg 1 \) then loops are large, and the weak scale are set.

• One answer:

\[ V_H = \tilde{M}_H^2 |H|^2 + \Lambda_{NP}^2 |H|^2 \]

\[ \tilde{M}_H^2 + \Lambda_{NP}^2 \sim \text{small} \]

But this requires fine-tuning the fundamental parameters. Why would nature do that?
Problem solved! An entirely new spacetime symmetry, Supersymmetry, protects the Higgs!

Unifies fields into superfields, transforming into one another! Most important component: The “Stop Squark”. The heavier, the more fine-tuned...
Supersymmetry

The last time a spacetime symmetry was discovered...
The last time a spacetime symmetry was discovered...

Supersymmetry... However...

Someone didn’t show up to the party... yet...

\[ 1s=13 \text{ TeV} \]

\[ 0L \ 36.1 \ \text{fb}^{-1} \ [\text{CONF-2017-020}] \]

\[ 1L \ 36.1 \ \text{fb}^{-1} \ [\text{CONF-2017-037}] \]

\[ 2L \ 36.1 \ \text{fb}^{-1} \ [\text{CONF-2017-034}] \]

\[ \text{Monojet} \ 3.2 \ \text{fb}^{-1} \ [1604.07773] \]

\[ \text{Run 1} \ [1506.08616] \]
Circa 1980’s

Problem solved! The Higgs boson is like a pion, naturally light because of spontaneously broken global symmetry!

\[ H = \overline{f}f \sim f \]

Then no hierarchy problem, since no fundamental scalar! Realise the Higgs boson analogously to the pion in QCD.

\[ \rho = \overline{f}f \sim \Lambda \]

Should also get other heavy resonances then!
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Realise the Higgs boson analogously to the pion in QCD.

Should also get other heavy resonances then!

\[
\begin{align*}
    f &\sim f \\
    \therefore &\sim \phi \\
    \therefore &\sim g_\rho = 4\pi
\end{align*}
\]

Circa 1980's

However… Someone didn’t show up to the party… yet…

Should also get other heavy resonances then!
Ok, so back to fine-tuning...

Analogous to taking a superconductor and fine-tuning the temperature to just at the phase transition, such that

\[ m_\Phi(T) \ll m_\Phi(0) \]

It can exist, physically, and we can tune parameters in a lab, but why would nature?
Essentially, it seems like the Universe is just like a Transition Edge Sensor:

A tiny change of fundamental parameters would completely change the weak scale!

*Taken from 1309.5383*
Essentially, it seems like the Universe is just like a Transition Edge Sensor:

\[ \delta T \]

How could nature have engineered such a scenario?

A tiny change of fundamental parameters would completely change the weak scale!
The Relaxion

Rather than fighting with this conundrum, maybe embrace it? Perhaps this is precisely what happened in the early Universe?

How could this have happened in the primordial Universe?

Brand new set of theoretical ideas emerging...

Graham, Kaplan, Rajendran, 2015
The Relaxion

Rather than fighting with this conundrum, maybe embrace it? Perhaps this is precisely what happened in the early Universe?

\[ m_H(t) \ll m_H(t_0) \]

How could this have happened in the primordial Universe?

• Graham, Kaplan, Rajendran, 2015
The Relaxion

• Radically different take on the hierarchy problem.
• Basic ingredients: Add a new scalar field.

\[ \mathcal{L} \sim (M^2 - g\phi)|H|^2 \]

- Higgs mass takes large value \( M \gg 125 \text{ GeV} \)
- Relaxion “scans” Higgs mass
- Relaxion wants to roll due to small terms in potential

\[ -gM^2\phi + \frac{\phi}{32\pi^2f} \tilde{G}G \]

- Relaxion couples to gluons.
The Relaxion

• Radically different take on the hierarchy problem.
• Basic ingredients

\[ \mathcal{L} \sim (M^2 - g\phi)|H|^2 \]

\[ -gM^2 \phi \]

\[ + f_\pi^3 \lambda_q \langle h \rangle \cos \left( \frac{\phi}{f} \right) \]

Thus in terms of the Higgs vacuum expectation value the potential is
The Relaxion

- Radically different take on the hierarchy problem.
- Basic ingredients

\[ \mathcal{L} \sim (M^2 - g\phi) |H|^2 \]

Once it has rolled far enough, Higgs will develop a small VEV.

Then axion potential turns on and Relaxion stops rolling

\[ -gM^2 \phi \]

+ \[ f_\pi^3 \lambda_q \langle h \rangle \cos \left( \frac{\phi}{f} \right) \]

In early Universe Relaxion rolls
Relaxion starts at the top of potential. Starts rolling down.

Scans Higgs mass-squared while it rolls, slowly cancelling against large mass-squared.
The Relaxion

- Cosmological evolution

At some point relaxion crosses critical value at which Higgs mass-squared becomes zero.

After this mass-squared becomes negative:
- Higgs gets a vev
- Quarks get mass
- Axion potential turns on
Soon after axion potential turns on (while Higgs vev is still very small), relaxion becomes trapped and stops rolling. Thus Higgs vev becomes stuck at this finely-tuned point!
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Thus Higgs vev becomes stuck at this finely-tuned point!

Can choose “g” parameter such that field stops when \(<h>\) is still very small. This is a parameter choice, not a tuning, since radiatively stable.

\[
\frac{\partial V}{\partial \phi} \sim g M^2 - \frac{f^3 \lambda_q <H>}{f} \sin \left( \frac{\phi}{f} \right) = 0
\]
Soon after axion potential turns on (while Higgs vev is still very small), relaxion becomes trapped and stops rolling. Thus Higgs vev becomes stuck at this finely-tuned point!

Can choose "$g$" parameter such that field stops when $\langle h \rangle$ is still very small. This is a parameter choice, not a tuning, since radiatively stable.

All of this could happen during a very long period of inflation...

$$\frac{\partial V}{\partial \phi} \sim g M^2 - \frac{f^3 \lambda_q \langle H \rangle}{f} \sin \left( \frac{\phi}{f} \right) = 0$$
This theory is not fully fleshed out yet, and there are many details to work out, but it is the first example of a scenario wherein cosmological dynamics explains a finely-tuned electroweak scale.

Watch this space...
The End of the Universe

To finish our cosmological exploration of the weak scale, let’s look to the distant future...

The Higgs potential may have yet another minimum that I wasn’t telling you about. May be at lower energies...
The fate of the Universe...

The Higgs potential has another minimum...

From current LHC measurements it looks like our current location on the Higgs potential is unstable!
The fate of the Universe...

The Higgs potential has another minimum...

As it currently stands, in the distant future our very vacuum will decay with a 95% lower bound of $10^{58}$ Years. (1707.08124).

From current LHC measurements it looks like our current location on the Higgs potential is unstable!
The Truth About the Higgs

Mark Twain: “Truth is stranger than fiction, but it is because Fiction is obliged to stick to possibilities; Truth isn't.”

LHC measurements, taken at face value, are suggesting that:

- We live in a (relativistic) finely-tuned superconductor, sitting just at the threshold of the phase transition, in an unstable state, set to decay to a new Universe in the future.
The Truth About the Higgs

The only way to know if this is our world is to measure the shape of the scalar potential.

Measuring the Higgs self-coupling is the only way to measure the structure of the Higgs potential.

Discovering the Higgs was difficult enough, now we want to know how it behaves in private...

All of these questions (microscopic theory of EW symmetry breaking, phase transition, matter-antimatter imbalance, fine-tuning, Higgs portal) boil down to the shape and overall nature of the Higgs scalar potential.

Why don’t we just go and measure it?

Discovering the Higgs was difficult enough, now we want to know how it behaves in private...
The Truth About the Higgs

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The Higgs Potential...

At colliders gain sensitivity to self-interactions by searching for pair production:

To robustly measure this we will have to look to future colliders, and only then will we really know the shape of the Higgs potential and the superconducting phase transition!
Back to work.

We now have in our hands a little particle that holds the key to the biggest questions we could ask:

- Where did matter come from?
- What is the microscopic description of EW symmetry breaking?
- What is the ultimate fate of our Universe?
- Maybe even what is the dark matter?*

H
We are only at the beginning of this story.

To piece together the pieces of the Higgs puzzle we will increasingly have to interface collider measurements with cosmology. The right now, with the distant past, and the ultimate future.
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