

Identification of detrimental effects to multi-band superconductivity: Application to Sr_2RuO_4

Phys. Rev. B **94**, 104501 (2016)

Journal of Physics: Conf. Series **807**, 052011 (2017)

Aline Ramires¹ and Manfred Sigrist²

¹ Institute for Theoretical Studies - ETH - Zurich

² Institute for Theoretical Physics - ETH - Zurich

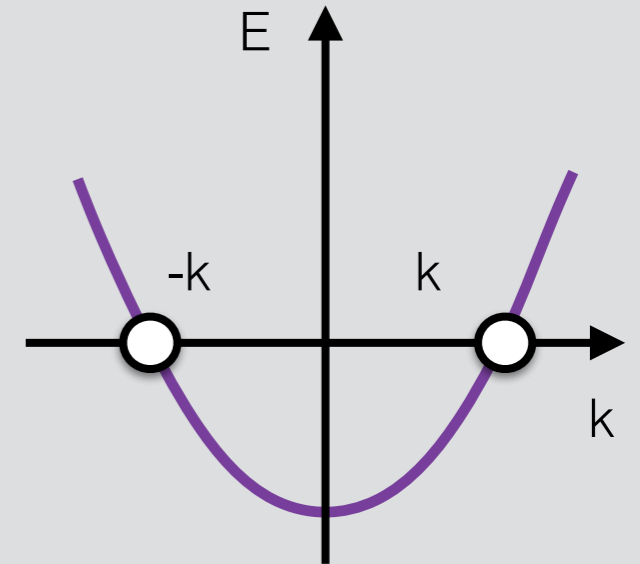
Outline

- 1) **What we already know:**
 - **BCS Theory, Symmetries and Anderson's Theorems**
- 2) **Orbital x Band Basis:**
 - **Introducing the concept of "SC Fitness"**
- 3) **Sanity check:**
 - **1 band + key symmetry breaking fields**
- 4) **Application to Sr_2RuO_4 :**
 - **The most compatible order parameter**
 - **New mechanism for suppression of SC**
- 5) **Final remarks**

What we already know:

BCS Theory: pairing of degenerate electrons around Fermi energy, with zero total momentum

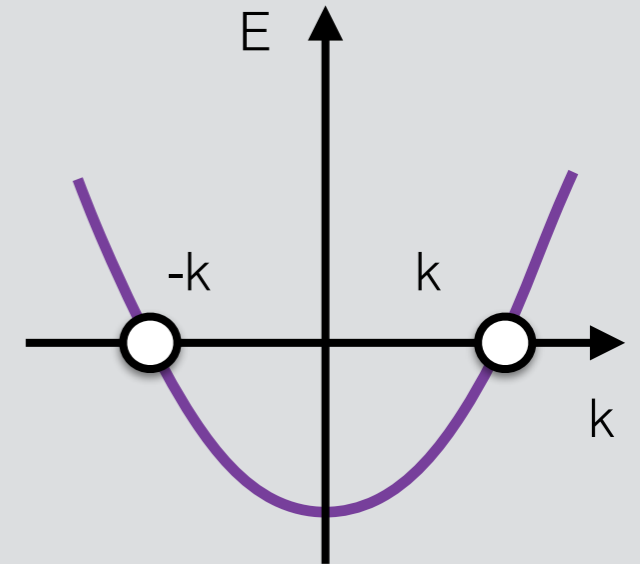
Robust instability: $T_c \sim \omega_c e^{-1/v_{eff} N(0)}$



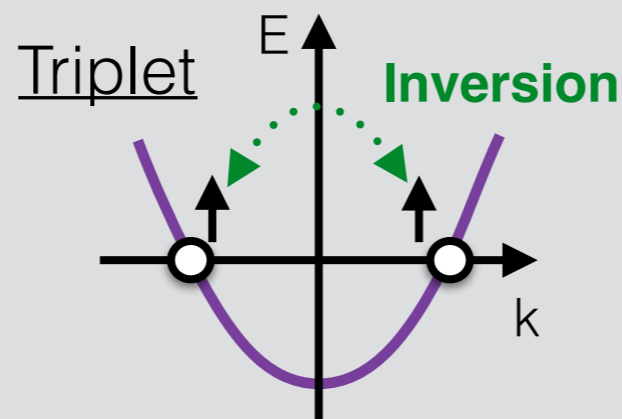
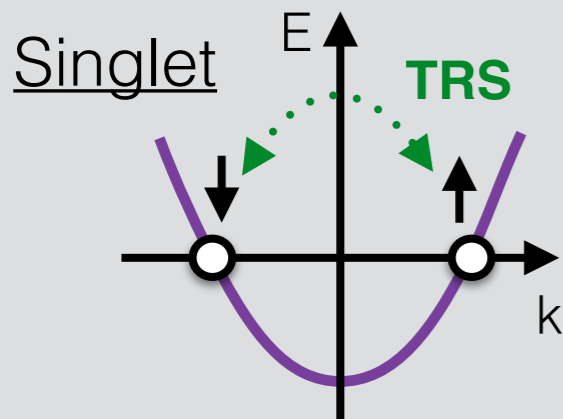
What we already know:

BCS Theory: pairing of degenerate electrons around Fermi energy, with zero total momentum

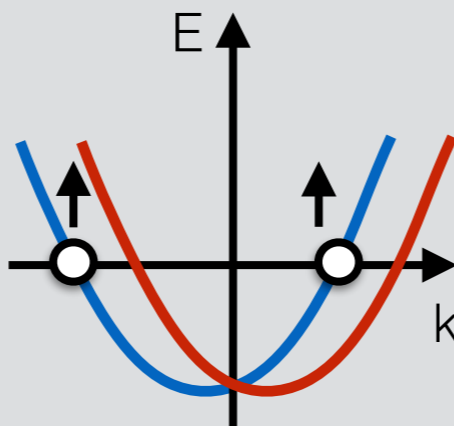
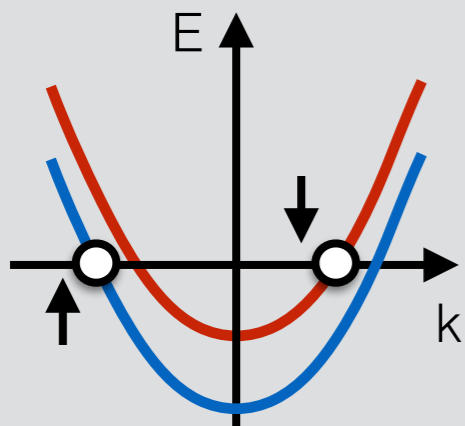
Robust instability: $T_c \sim \omega_c e^{-1/v_{eff} N(0)}$



Anderson's Theorems: the presence of states necessary for Cooper pairing in the singlet and triplet channels is guaranteed by the presence of key symmetries:



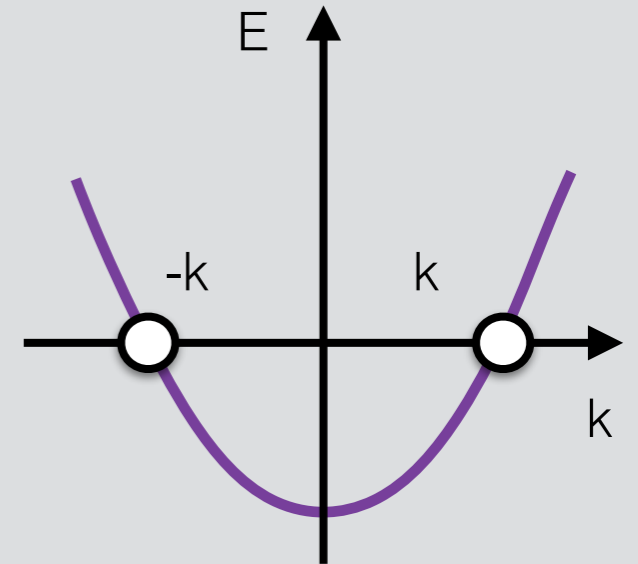
Breaking Key Symmetry



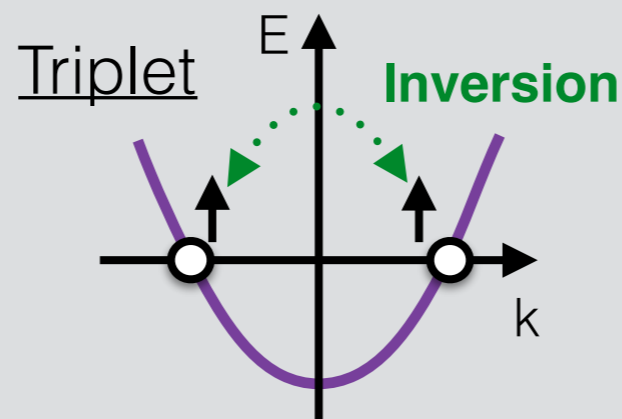
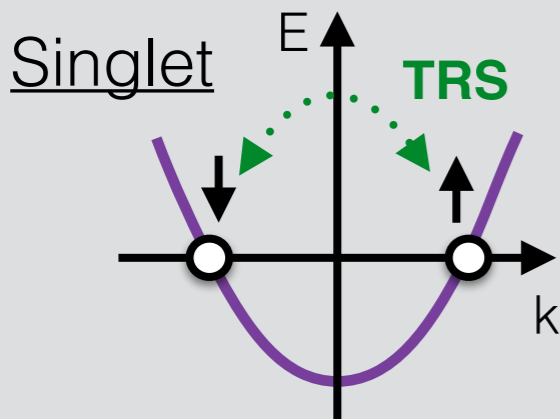
What we already know:

BCS Theory: pairing of degenerate electrons around Fermi energy, with zero total momentum

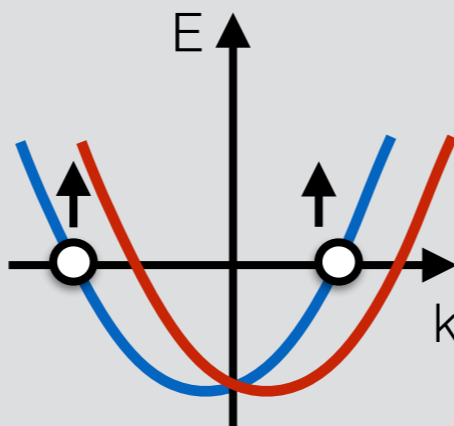
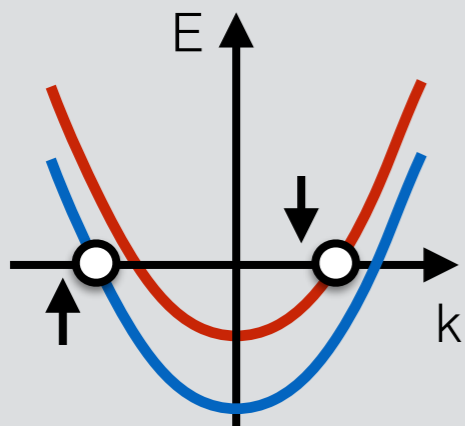
Robust instability: $T_c \sim \omega_c e^{-1/v_{eff} N(0)}$



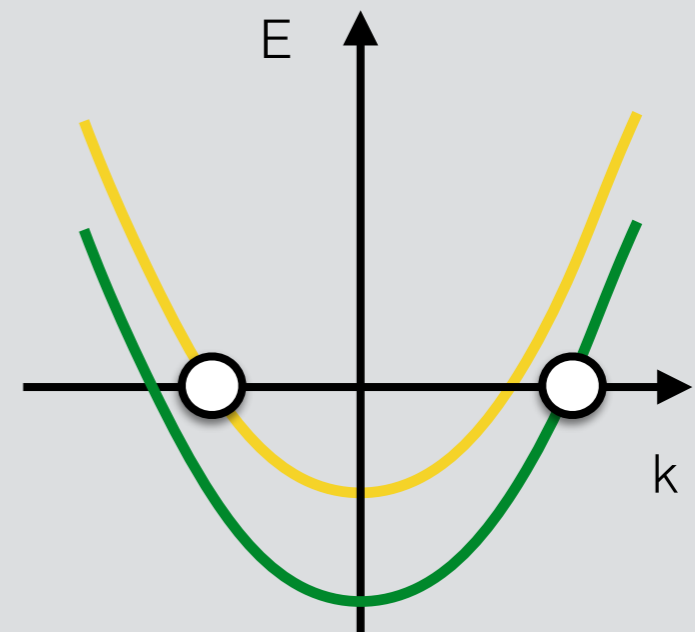
Anderson's Theorems: the presence of states necessary for Cooper pairing in the singlet and triplet channels is guaranteed by the presence of key symmetries:



Breaking Key Symmetry



Multi-Band Systems



"Inter-band pairing is NOT robust"

From the orbital to the band basis...

The MF Hamiltonian for a multi-orbital system:

$$H_{MF} = \sum_{\mathbf{k}} \Psi_{\mathbf{k}}^{\dagger} \begin{pmatrix} H_0(\mathbf{k}) & \Delta(\mathbf{k}) \\ \Delta^{\dagger}(\mathbf{k}) & -H_0^*(-\mathbf{k}) \end{pmatrix} \Psi_{\mathbf{k}},$$

in the **orbital basis**:

$$\Psi_{\mathbf{k}}^{\dagger} = (\psi_{\mathbf{k}}^{\dagger}, \psi_{-\mathbf{k}}^T),$$

$$\psi_{\mathbf{k}}^{\dagger} = (a_{1\mathbf{k}\uparrow}^{\dagger} \ a_{1\mathbf{k}\downarrow}^{\dagger}, \dots, a_{n\mathbf{k}\uparrow}^{\dagger} \ a_{n\mathbf{k}\downarrow}^{\dagger})$$

From the orbital to the band basis...

The MF Hamiltonian for a multi-orbital system:

$$H_{MF} = \sum_{\mathbf{k}} \Psi_{\mathbf{k}}^{\dagger} \begin{pmatrix} H_0(\mathbf{k}) & \Delta(\mathbf{k}) \\ \Delta^{\dagger}(\mathbf{k}) & -H_0^*(-\mathbf{k}) \end{pmatrix} \Psi_{\mathbf{k}},$$

in the **orbital basis**:

$$\Psi_{\mathbf{k}}^{\dagger} = (\psi_{\mathbf{k}}^{\dagger}, \psi_{-\mathbf{k}}^T),$$
$$\psi_{\mathbf{k}}^{\dagger} = (a_{1\mathbf{k}\uparrow}^{\dagger} \ a_{1\mathbf{k}\downarrow}^{\dagger}, \dots, a_{n\mathbf{k}\uparrow}^{\dagger} \ a_{n\mathbf{k}\downarrow}^{\dagger})$$

In general the non-interacting Hamiltonian is not diagonal in the orbital basis.
(Inter-orbital hopping, SOC,...)

Unitary transformation to the **band basis**:

$$H_0^B(\mathbf{k}) = U_{\mathbf{k}} H_0(\mathbf{k}) U_{\mathbf{k}}^{\dagger}$$

From the orbital to the band basis...

The MF Hamiltonian for a multi-orbital system:

$$H_{MF} = \sum_{\mathbf{k}} \Psi_{\mathbf{k}}^{\dagger} \begin{pmatrix} H_0(\mathbf{k}) & \Delta(\mathbf{k}) \\ \Delta^{\dagger}(\mathbf{k}) & -H_0^*(-\mathbf{k}) \end{pmatrix} \Psi_{\mathbf{k}},$$

in the **orbital basis**:

$$\Psi_{\mathbf{k}}^{\dagger} = (\psi_{\mathbf{k}}^{\dagger}, \psi_{-\mathbf{k}}^T),$$
$$\psi_{\mathbf{k}}^{\dagger} = (a_{1\mathbf{k}\uparrow}^{\dagger} \ a_{1\mathbf{k}\downarrow}^{\dagger}, \dots, a_{n\mathbf{k}\uparrow}^{\dagger} \ a_{n\mathbf{k}\downarrow}^{\dagger})$$


In general the non-interacting Hamiltonian is not diagonal in the orbital basis.
(Inter-orbital hopping, SOC,...)

Unitary transformation to the **band basis**:

$$H_0^B(\mathbf{k}) = U_{\mathbf{k}} H_0(\mathbf{k}) U_{\mathbf{k}}^{\dagger}$$

$$\Delta^B(\mathbf{k}) = U_{\mathbf{k}} \Delta(\mathbf{k}) U_{-\mathbf{k}}^T$$

Note the different transformation
for the gap matrix!



From the orbital to the band basis...

The MF Hamiltonian for a multi-orbital system:

$$H_{MF} = \sum_{\mathbf{k}} \Psi_{\mathbf{k}}^{\dagger} \begin{pmatrix} H_0(\mathbf{k}) & \Delta(\mathbf{k}) \\ \Delta^{\dagger}(\mathbf{k}) & -H_0^*(-\mathbf{k}) \end{pmatrix} \Psi_{\mathbf{k}},$$

in the **orbital basis**:

$$\Psi_{\mathbf{k}}^{\dagger} = (\psi_{\mathbf{k}}^{\dagger}, \psi_{-\mathbf{k}}^T),$$
$$\psi_{\mathbf{k}}^{\dagger} = (a_{1\mathbf{k}\uparrow}^{\dagger} \ a_{1\mathbf{k}\downarrow}^{\dagger}, \dots, a_{n\mathbf{k}\uparrow}^{\dagger} \ a_{n\mathbf{k}\downarrow}^{\dagger})$$

In general the non-interacting Hamiltonian is not diagonal in the orbital basis.
(Inter-orbital hopping, SOC,...)

Unitary transformation to the **band basis**:

$$H_0^B(\mathbf{k}) = U_{\mathbf{k}} H_0(\mathbf{k}) U_{\mathbf{k}}^{\dagger}$$

In presence of **TRS and Inv**:
2x2 block diagonal

$$\Delta^B(\mathbf{k}) = U_{\mathbf{k}} \Delta(\mathbf{k}) U_{-\mathbf{k}}^T$$

Note the different transformation
for the gap matrix!

From the orbital to the band basis...

The MF Hamiltonian for a multi-orbital system:

$$H_{MF} = \sum_{\mathbf{k}} \Psi_{\mathbf{k}}^{\dagger} \begin{pmatrix} H_0(\mathbf{k}) & \Delta(\mathbf{k}) \\ \Delta^{\dagger}(\mathbf{k}) & -H_0^*(-\mathbf{k}) \end{pmatrix} \Psi_{\mathbf{k}},$$

in the **orbital basis**:

$$\Psi_{\mathbf{k}}^{\dagger} = (\psi_{\mathbf{k}}^{\dagger}, \psi_{-\mathbf{k}}^T),$$
$$\psi_{\mathbf{k}}^{\dagger} = (a_{1\mathbf{k}\uparrow}^{\dagger} \ a_{1\mathbf{k}\downarrow}^{\dagger}, \dots, a_{n\mathbf{k}\uparrow}^{\dagger} \ a_{n\mathbf{k}\downarrow}^{\dagger})$$

In general the non-interacting Hamiltonian is not diagonal in the orbital basis.
(Inter-orbital hopping, SOC,...)

Unitary transformation to the **band basis**:

$$H_0^B(\mathbf{k}) = U_{\mathbf{k}} H_0(\mathbf{k}) U_{\mathbf{k}}^{\dagger}$$

In presence of **TRS and Inv**:
2x2 block diagonal

$$\Delta^B(\mathbf{k}) = U_{\mathbf{k}} \Delta(\mathbf{k}) U_{-\mathbf{k}}^T$$

If only **intra-band pairing**:
2x2 block diagonal

Note the different transformation
for the gap matrix!

From the orbital to the band basis...

The MF Hamiltonian for a multi-orbital system:

$$H_{MF} = \sum_{\mathbf{k}} \Psi_{\mathbf{k}}^{\dagger} \begin{pmatrix} H_0(\mathbf{k}) & \Delta(\mathbf{k}) \\ \Delta^{\dagger}(\mathbf{k}) & -H_0^*(-\mathbf{k}) \end{pmatrix} \Psi_{\mathbf{k}},$$

in the **orbital basis**:

$$\Psi_{\mathbf{k}}^{\dagger} = (\psi_{\mathbf{k}}^{\dagger}, \psi_{-\mathbf{k}}^T),$$

$$\psi_{\mathbf{k}}^{\dagger} = (a_{1\mathbf{k}\uparrow}^{\dagger} \ a_{1\mathbf{k}\downarrow}^{\dagger}, \dots, a_{n\mathbf{k}\uparrow}^{\dagger} \ a_{n\mathbf{k}\downarrow}^{\dagger})$$

In general the non-interacting Hamiltonian is not diagonal in the orbital basis.
(Inter-orbital hopping, SOC,...)

Unitary transformation to the **band basis**:

$$H_0^B(\mathbf{k}) = U_{\mathbf{k}} H_0(\mathbf{k}) U_{\mathbf{k}}^{\dagger}$$

$$\Delta^B(\mathbf{k}) = U_{\mathbf{k}} \Delta(\mathbf{k}) U_{-\mathbf{k}}^T$$

In presence of **TRS and Inv**:
2x2 block diagonal

If only **intra-band pairing**:
2x2 block diagonal

$$[H_0^B(\mathbf{k}), \Delta^B(\mathbf{k})] = 0$$

Note the different transformation
for the gap matrix!

...and back to the orbital basis:

Key symmetries+
Only intra-band pairing



$$[H_0^B(\mathbf{k}), \Delta^B(\mathbf{k})] = 0$$

Q: What condition it imposes in the **orbital basis**?

$$H_0(\mathbf{k})\Delta(\mathbf{k}) - \Delta(\mathbf{k})H_0^*(-\mathbf{k}) = 0,$$

...and back to the orbital basis:

Key symmetries+
Only intra-band pairing



$$[H_0^B(\mathbf{k}), \Delta^B(\mathbf{k})] = 0$$

Q: What condition it imposes in the **orbital basis**?

$$H_0(\mathbf{k})\Delta(\mathbf{k}) - \Delta(\mathbf{k})H_0^*(-\mathbf{k}) = 0,$$

In general this condition won't be satisfied and this modified commutator is finite, so we introduce the concept of "**SC Fitness**":

$$H_0(\mathbf{k})\Delta(\mathbf{k}) - \Delta(\mathbf{k})H_0^*(-\mathbf{k}) = F(\mathbf{k})(i\sigma_2)$$

Measure of **IN**compatibility of
the gap and underlying
electronic structure.

$$\begin{cases} F(\mathbf{k}) = 0 & \text{Compatible} \\ F(\mathbf{k}) \neq 0 & \text{Not Compatible/} \\ & \text{Detrimental} \end{cases}$$

Sanity check:

Single band problem under external perturbation:

$$H_0(\mathbf{k})\Delta(\mathbf{k}) - \Delta(\mathbf{k})H_0^*(-\mathbf{k}) = F(\mathbf{k})(i\sigma_2)$$

$$H_0(\mathbf{k}) \rightarrow H_0(\mathbf{k}) + \delta H(\mathbf{k})$$

<p>"SC Fitness"</p> <p>$F(\mathbf{k})$</p>	<p>TRS breaking</p> <p>$\delta H_{TR} = -\mathbf{h} \cdot \boldsymbol{\sigma}$</p>	<p>Inv breaking</p> <p>$\delta H_{Inv}(\mathbf{k}) = \mathbf{g}(\mathbf{k}) \cdot \boldsymbol{\sigma}$</p>
<p>Singlet</p> <p>$\Delta_S(\mathbf{k}) = d_0(\mathbf{k})(i\sigma_2)$</p>		
<p>Triplet</p> <p>$\mathbf{z} = S_z=0\rangle = \frac{1}{\sqrt{2}}(\uparrow\downarrow\rangle + \downarrow\uparrow\rangle),$</p> <p>$\mathbf{x} = S_x=0\rangle = \frac{1}{\sqrt{2}}(- \uparrow\uparrow\rangle + \downarrow\downarrow\rangle),$</p> <p>$\mathbf{y} = S_y=0\rangle = \frac{1}{\sqrt{2}}(\uparrow\uparrow\rangle + \downarrow\downarrow\rangle),$</p> <p>$\Delta_T(\mathbf{k}) = (\mathbf{d}_T(\mathbf{k}) \cdot \boldsymbol{\sigma})(i\sigma_2)$</p>		

Sanity check:

Single band problem under external perturbation:

$$H_0(\mathbf{k})\Delta(\mathbf{k}) - \Delta(\mathbf{k})H_0^*(-\mathbf{k}) = F(\mathbf{k})(i\sigma_2)$$

$$H_0(\mathbf{k}) \rightarrow H_0(\mathbf{k}) + \delta H(\mathbf{k})$$

<p>"SC Fitness"</p> <p>$F(\mathbf{k})$</p>	<p>TRS breaking</p> <p>$\delta H_{TR} = -\mathbf{h} \cdot \boldsymbol{\sigma}$</p>	<p>Inv breaking</p> <p>$\delta H_{Inv}(\mathbf{k}) = \mathbf{g}(\mathbf{k}) \cdot \boldsymbol{\sigma}$</p>
<p>Singlet</p> <p>$\Delta_S(\mathbf{k}) = d_0(\mathbf{k})(i\sigma_2)$</p>		
<p>Triplet</p> <p>$\mathbf{z} = S_z=0\rangle = \frac{1}{\sqrt{2}}(\uparrow\downarrow\rangle + \downarrow\uparrow\rangle),$</p> <p>$\mathbf{x} = S_x=0\rangle = \frac{1}{\sqrt{2}}(- \uparrow\uparrow\rangle + \downarrow\downarrow\rangle),$</p> <p>$\mathbf{y} = S_y=0\rangle = \frac{1}{\sqrt{2}}(\uparrow\uparrow\rangle + \downarrow\downarrow\rangle),$</p> <p>$\Delta_T(\mathbf{k}) = (\mathbf{d}_T(\mathbf{k}) \cdot \boldsymbol{\sigma})(i\sigma_2)$</p>		

Sanity check:

Single band problem under external perturbation:

$$H_0(\mathbf{k})\Delta(\mathbf{k}) - \Delta(\mathbf{k})H_0^*(-\mathbf{k}) = F(\mathbf{k})(i\sigma_2)$$

$$H_0(\mathbf{k}) \rightarrow H_0(\mathbf{k}) + \delta H(\mathbf{k})$$

<p>"SC Fitness"</p> <p>$F(\mathbf{k})$</p>	<p>TRS breaking</p> <p>$\delta H_{TR} = -\mathbf{h} \cdot \boldsymbol{\sigma}$</p>	<p>Inv breaking</p> <p>$\delta H_{Inv}(\mathbf{k}) = \mathbf{g}(\mathbf{k}) \cdot \boldsymbol{\sigma}$</p>
<p>Singlet</p> <p>$\Delta_S(\mathbf{k}) = d_0(\mathbf{k})(i\sigma_2)$</p>	<p>$-2d_0(\mathbf{k})\mathbf{h} \cdot \boldsymbol{\sigma}$</p> <p>Field in any direction.</p>	<p>0.</p> <p>Not detrimental!</p>
<p>Triplet</p> <p>$\mathbf{z} = S_z=0\rangle = \frac{1}{\sqrt{2}}(\uparrow\downarrow\rangle + \downarrow\uparrow\rangle),$</p> <p>$\mathbf{x} = S_x=0\rangle = \frac{1}{\sqrt{2}}(- \uparrow\uparrow\rangle + \downarrow\downarrow\rangle),$</p> <p>$\mathbf{y} = S_y=0\rangle = \frac{1}{\sqrt{2}}(\uparrow\uparrow\rangle + \downarrow\downarrow\rangle),$</p> <p>$\Delta_T(\mathbf{k}) = (\mathbf{d}_T(\mathbf{k}) \cdot \boldsymbol{\sigma})(i\sigma_2)$</p>	<p>$-2\mathbf{h} \cdot \mathbf{d}_T(\mathbf{k})I_2$</p> <p>Only for field with component in the d-vector direction.</p>	<p>$2i(\mathbf{g}(\mathbf{k}) \times \mathbf{d}_T(\mathbf{k})) \cdot \boldsymbol{\sigma}$</p> <p>Only for field with component perpendicular to d-vector.</p>

Sanity check:

Single band problem under external perturbation:

$$H_0(\mathbf{k})\Delta(\mathbf{k}) - \Delta(\mathbf{k})H_0^*(-\mathbf{k}) = F(\mathbf{k})(i\sigma_2)$$

$$H_0(\mathbf{k}) \rightarrow H_0(\mathbf{k}) + \delta H(\mathbf{k})$$

<p>"SC Fitness"</p> <p>$F(\mathbf{k})$</p>	<p>TRS breaking</p> <p>$\delta H_{TR} = -\mathbf{h} \cdot \boldsymbol{\sigma}$</p>	<p>Inv breaking</p> <p>$\delta H_{Inv}(\mathbf{k}) = \mathbf{g}(\mathbf{k}) \cdot \boldsymbol{\sigma}$</p>
<p>Singlet</p> <p>$\Delta_S(\mathbf{k}) = d_0(\mathbf{k})(i\sigma_2)$</p>	<p>$-2d_0(\mathbf{k})\mathbf{h} \cdot \boldsymbol{\sigma}$</p> <p>Field in any direction.</p>	<p>0.</p> <p>Not detrimental!</p>
<p>Triplet</p> <p>$\mathbf{z} = S_z=0\rangle = \frac{1}{\sqrt{2}}(\uparrow\downarrow\rangle + \downarrow\uparrow\rangle),$</p> <p>$\mathbf{x} = S_x=0\rangle = \frac{1}{\sqrt{2}}(- \uparrow\uparrow\rangle + \downarrow\downarrow\rangle),$</p> <p>$\mathbf{y} = S_y=0\rangle = \frac{1}{\sqrt{2}}(\uparrow\uparrow\rangle + \downarrow\downarrow\rangle),$</p> <p>$\Delta_T(\mathbf{k}) = (\mathbf{d}_T(\mathbf{k}) \cdot \boldsymbol{\sigma})(i\sigma_2)$</p>	<p>$-2\mathbf{h} \cdot \mathbf{d}_T(\mathbf{k})I_2$</p> <p>Only for field with component in the d-vector direction.</p>	<p>$2i(\mathbf{g}(\mathbf{k}) \times \mathbf{d}_T(\mathbf{k})) \cdot \boldsymbol{\sigma}$</p> <p>Only for field with component perpendicular to d-vector.</p>

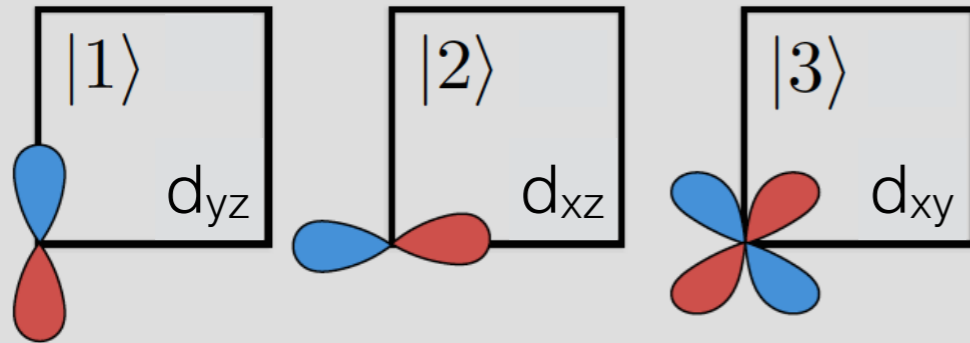
Unified closed form for the effect of perturbations on the critical temperature:

$$T_c \sim T_c^0 \left(1 - \frac{7\xi(3)}{64(\pi T_c^0)^2} \left\langle \text{Tr} |\hat{F}(\mathbf{k})|^2 \right\rangle_{FS} \right)$$

* perturbative in the external symmetry breaking fields

Application to Sr_2RuO_4 :

Basic ingredients:



Hamiltonian including **Inter-orbital hopping (IOH)** and **Spin-orbit coupling (SOC)**

$$H_{SRO}(\mathbf{k}) = \begin{pmatrix} \epsilon_{1\mathbf{k}}I_2 & t_{\mathbf{k}}I_2 + i\eta\sigma_3 & -i\eta\sigma_2 \\ t_{\mathbf{k}}I_2 - i\eta\sigma_3 & \epsilon_{2\mathbf{k}}I_2 & i\eta\sigma_1 \\ i\eta\sigma_2 & -i\eta\sigma_1 & \epsilon_{3\mathbf{k}}I_2 \end{pmatrix}$$

in the orbital basis:

$$\Psi_{\mathbf{k}}^\dagger = (c_{1\mathbf{k}\uparrow}^\dagger, c_{1\mathbf{k}\downarrow}^\dagger, c_{2\mathbf{k}\uparrow}^\dagger, c_{2\mathbf{k}\downarrow}^\dagger, c_{3\mathbf{k}\uparrow}^\dagger, c_{3\mathbf{k}\downarrow}^\dagger),$$

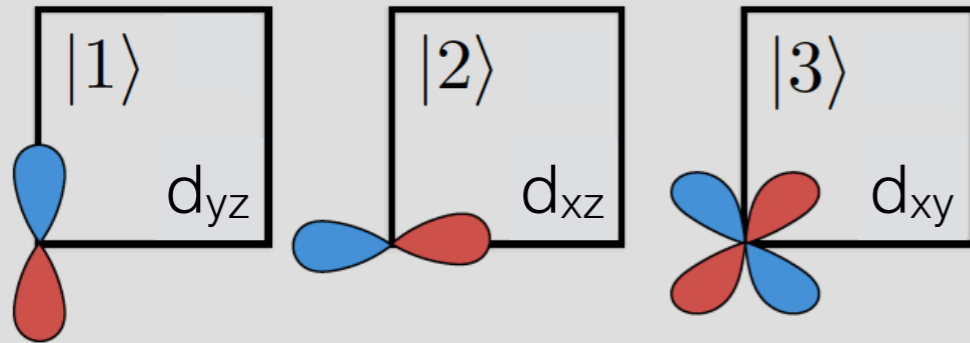
The gap matrix (only intra-orbital pairing):

$$\Delta_a(\mathbf{k}) = \begin{cases} d_{a0}(\mathbf{k})(i\sigma_2) \\ \mathbf{d}_a(\mathbf{k}) \cdot \boldsymbol{\sigma}(i\sigma_2) \end{cases}$$

$a = 1, 2, 3$

Application to Sr₂RuO₄:

Basic ingredients:



Hamiltonian including **Inter-orbital hopping (IOH)** and **Spin-orbit coupling (SOC)**

$$H_{SRO}(\mathbf{k}) = \begin{pmatrix} \epsilon_{1\mathbf{k}}I_2 & t_{\mathbf{k}}I_2 + i\eta\sigma_3 & -i\eta\sigma_2 \\ t_{\mathbf{k}}I_2 - i\eta\sigma_3 & \epsilon_{2\mathbf{k}}I_2 & i\eta\sigma_1 \\ i\eta\sigma_2 & -i\eta\sigma_1 & \epsilon_{3\mathbf{k}}I_2 \end{pmatrix}$$

in the orbital basis:

$$\Psi_{\mathbf{k}}^\dagger = (c_{1\mathbf{k}\uparrow}^\dagger, c_{1\mathbf{k}\downarrow}^\dagger, c_{2\mathbf{k}\uparrow}^\dagger, c_{2\mathbf{k}\downarrow}^\dagger, c_{3\mathbf{k}\uparrow}^\dagger, c_{3\mathbf{k}\downarrow}^\dagger),$$

The gap matrix (only intra-orbital pairing):

$$\Delta_a(\mathbf{k}) = \begin{cases} d_{a0}(\mathbf{k})(i\sigma_2) \\ \mathbf{d}_a(\mathbf{k}) \cdot \boldsymbol{\sigma}(i\sigma_2) \end{cases} \\ a = 1, 2, 3$$

Q: How **IOE** shape the SC state?

$$[H_0(\mathbf{k}) + \delta H(\mathbf{k}), \hat{\Delta}(\mathbf{k})]^* = \hat{F}(\mathbf{k})(i\sigma_2)$$

Example: **IOH**

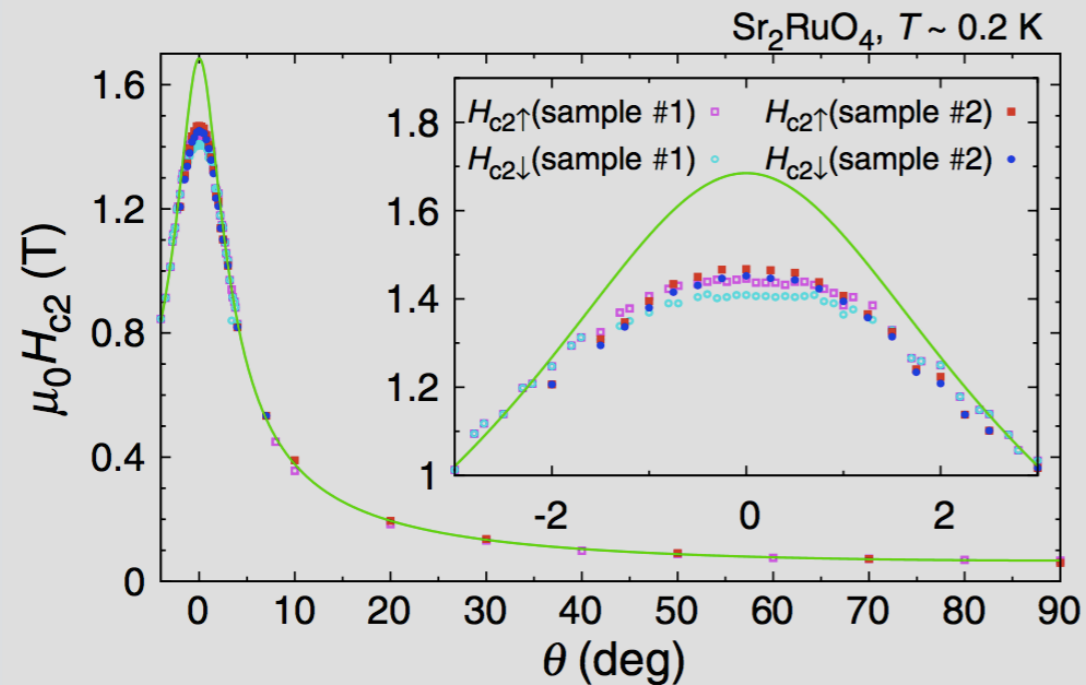
$$\hat{F}_{IOH} = it(\hat{d}_{1j} - \hat{d}_{2j})\lambda_2 \otimes \sigma_j \\ j = 0, x, y, z$$

$F(\mathbf{k})$	IOH	SOC	IOH +SOC
Δ_S	$\hat{d}_{10} = \hat{d}_{20}$	$\hat{d}_{10} = \hat{d}_{20} = \hat{d}_{30}$	✓
Δ_{Tx}	$\hat{d}_{1x} = \hat{d}_{2x}$	$-\hat{d}_{1x} = \hat{d}_{2x} = \hat{d}_{3z}$	✗
Δ_{Ty}	$\hat{d}_{1y} = \hat{d}_{2y}$	$\hat{d}_{1y} = -\hat{d}_{2y} = \hat{d}_{3y}$	✗
Δ_{Tz}	$\hat{d}_{1z} = \hat{d}_{2z}$	$\hat{d}_{1z} = \hat{d}_{2z} = -\hat{d}_{3z}$	✓

**Only compatible triplet state:
d-vector in the z-direction!**

Sr₂RuO₄: upper critical field

I) Anisotropy:



Effective mass model (EMM):

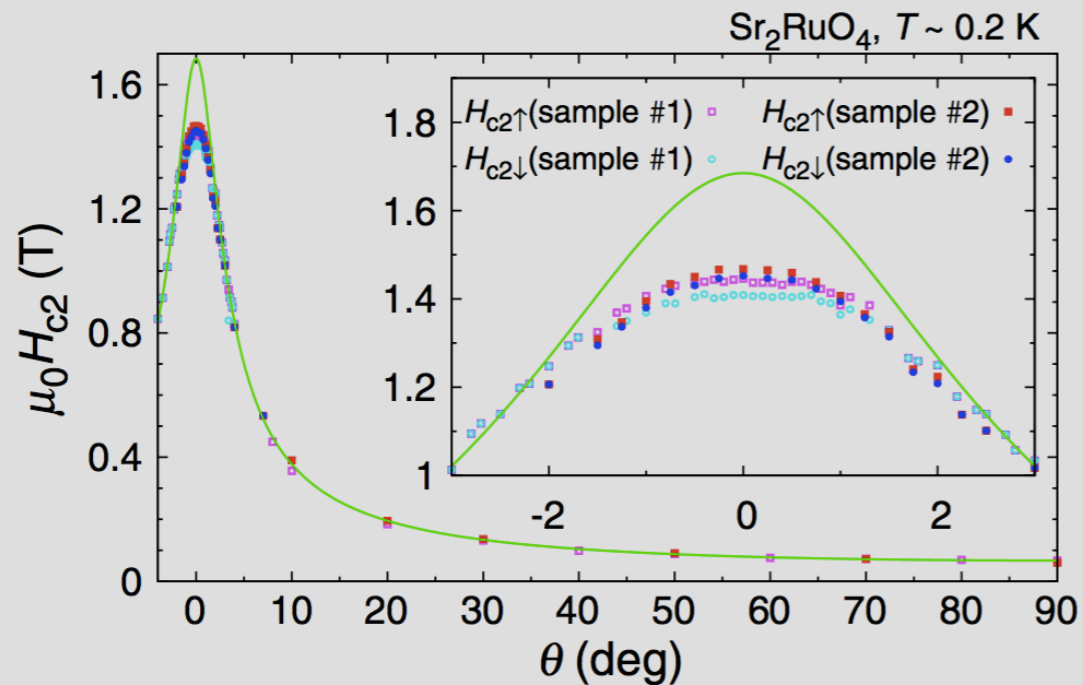
$$h_{c2,EMM}(\theta) = \frac{h_{c2}(90^\circ)}{\sqrt{\sin^2 \theta + \cos^2 \theta / \Gamma^2}},$$

$$\Gamma = \sqrt{M/m} \approx 25$$

Deviations from EMM for angles within 2° from the plane.

Sr₂RuO₄: upper critical field

I) Anisotropy:



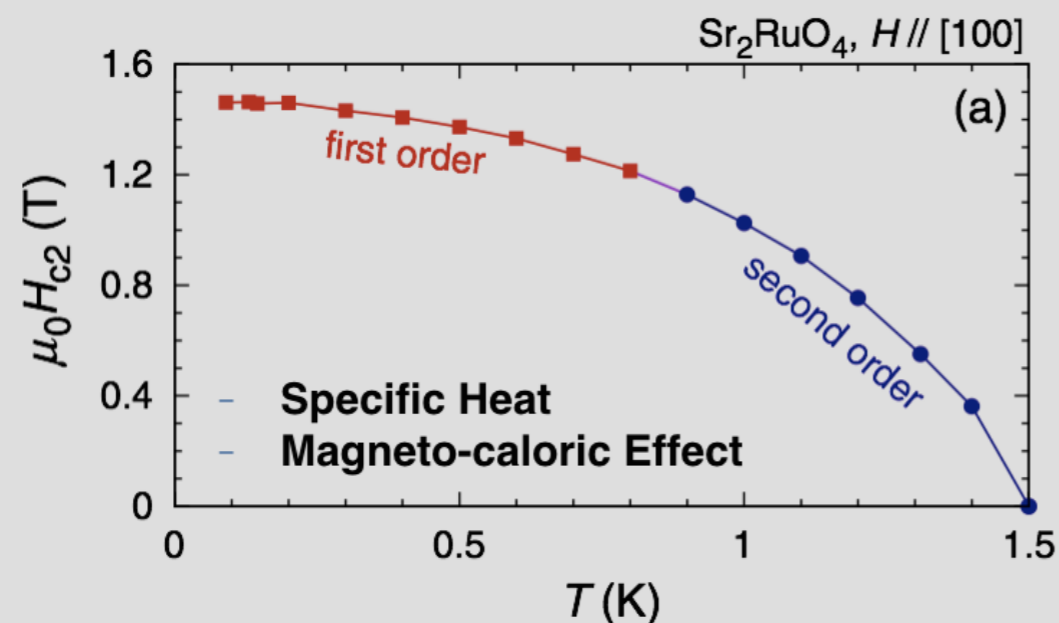
Effective mass model (EMM):

$$h_{c2,EMM}(\theta) = \frac{h_{c2}(90^\circ)}{\sqrt{\sin^2 \theta + \cos^2 \theta / \Gamma^2}},$$

$$\Gamma = \sqrt{M/m} \approx 25$$

Deviations from EMM for angles within 2° from the plane.

II) Transition 2nd → 1st order:



For in plane fields transition becomes **1st order** below 0.8K.

1st order → Pauli limiting effect? NO:
d-vector is along z-direction.
No change in the Knight shift across Tc.

New mechanism?!

Under magnetic field:

Given a multi-orbital system:

$$\delta H_{Orb} = -\mathbf{L} \cdot \mathbf{h}_o$$

Under magnetic field:

Given a multi-orbital system:

$$\delta H_{Orb} = -\mathbf{L} \cdot \mathbf{h}_o$$

Evaluating the **"SC Fitness"** for:

$$\left\{ \begin{array}{l} \mathbf{h}_o // \hat{x} \\ \hat{d}_{z1} = \hat{d}_{z2} = -\hat{d}_{z3} = d_z \end{array} \right.$$

$$\hat{F}_{Orb} = -2ih_{ox}d_z\lambda_7 \otimes \sigma_3$$

Now in plane magnetic is also detrimental to d-vector along z-direction due to IOE.

Under magnetic field:

Given a multi-orbital system:

$$\delta H_{Orb} = -\mathbf{L} \cdot \mathbf{h}_o$$

Evaluating the **"SC Fitness"** for:

$$\begin{cases} \mathbf{h}_o // \hat{x} \\ \hat{d}_{z1} = \hat{d}_{z2} = -\hat{d}_{z3} = d_z \end{cases}$$

$$\hat{F}_{Orb} = -2ih_{ox}d_z\lambda_7 \otimes \sigma_3$$

Now in plane magnetic is also detrimental to d-vector along z-direction due to IOE.

Qualitative distinct from the single-band scenario:

$$\hat{F}_{Spin} = -2\mathbf{h} \cdot \mathbf{d}_T I_2$$

Under magnetic field:

Given a multi-orbital system:

$$\delta H_{Orb} = -\mathbf{L} \cdot \mathbf{h}_o$$

Evaluating the **"SC Fitness"** for:

$$\begin{cases} \mathbf{h}_o // \hat{x} \\ \hat{d}_{z1} = \hat{d}_{z2} = -\hat{d}_{z3} = d_z \end{cases}$$

$$\hat{F}_{Orb} = -2ih_{ox}d_z\lambda_7 \otimes \sigma_3$$

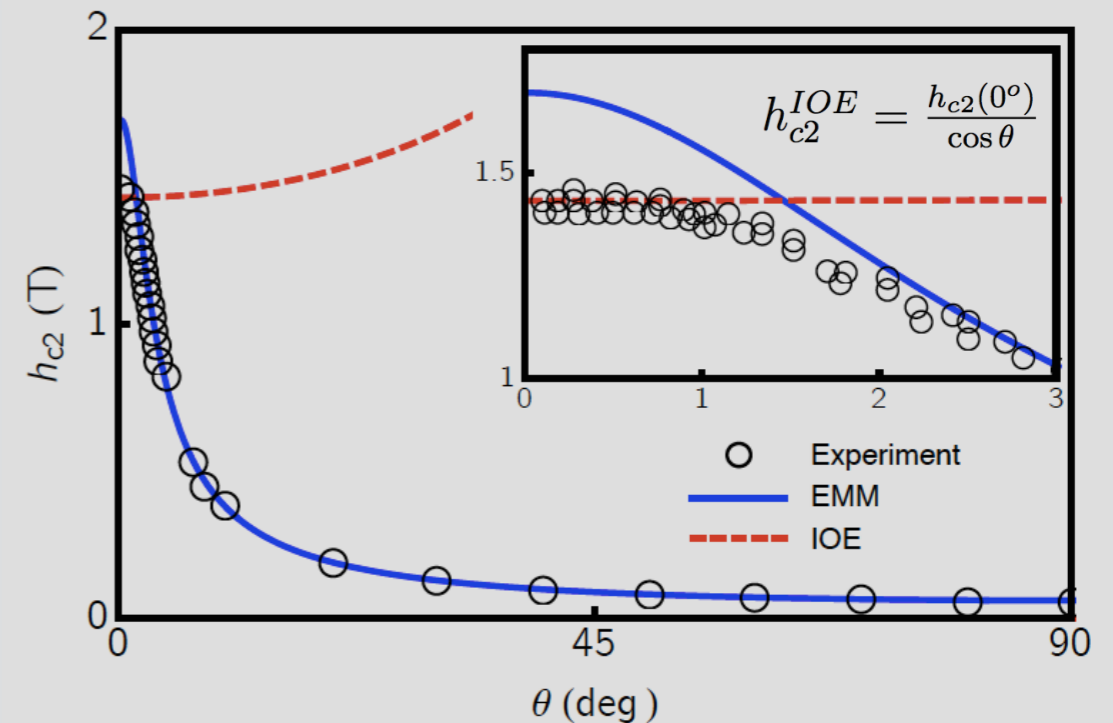
Now in plane magnetic is also detrimental to d-vector along z-direction due to IOE.

Qualitative distinct from the single-band scenario:

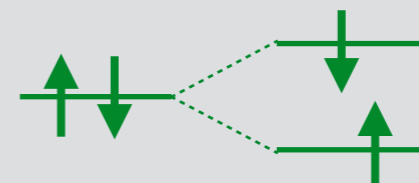
$$\hat{F}_{Spin} = -2\mathbf{h} \cdot \mathbf{d}_T I_2$$

INTER ORBITAL EFFECT

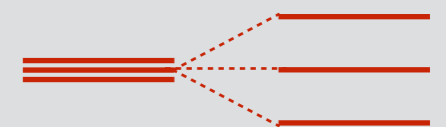
- **Change in anisotropy**
- **1st order transition**



$$\left. \begin{aligned} E_C &= -\frac{N(0)|\Delta|^2}{2} \\ E_M &= -\frac{\Delta\chi_{Orb}h_x^2}{2} \end{aligned} \right\} h_{c2}^{IOE} \approx \sqrt{\frac{N(0)}{\Delta\chi_{Orb}} \frac{T_c^0}{\cos \theta}}$$



Pauli PM Limiting



IOE Limiting

Final Remarks

- We developed the concept of **SC Fitness**:

$$H_0(\mathbf{k})\Delta(\mathbf{k}) - \Delta(\mathbf{k})H_0^*(-\mathbf{k}) = F(\mathbf{k})(i\sigma_2)$$

- Measure of the **instability of the SC state** in presence of symmetry breaking fields;
- **Application to Sr₂RuO₄**:
 - Find the **interplay of IOH and SOC** favours the **d-vector along the z-direction**;
 - Find **new mechanism** to suppress SC in multi-orbital systems:
Inter-Orbital Effect;
- **Future work**:
 - Detailed microscopic calculation for the IOE;
 - Application to other materials and model systems;

SC Fitness 2.0: Useful tool for engineering the normal state in order to obtain the desired (exotic/topological) SC state.