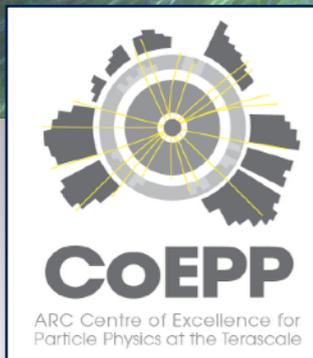
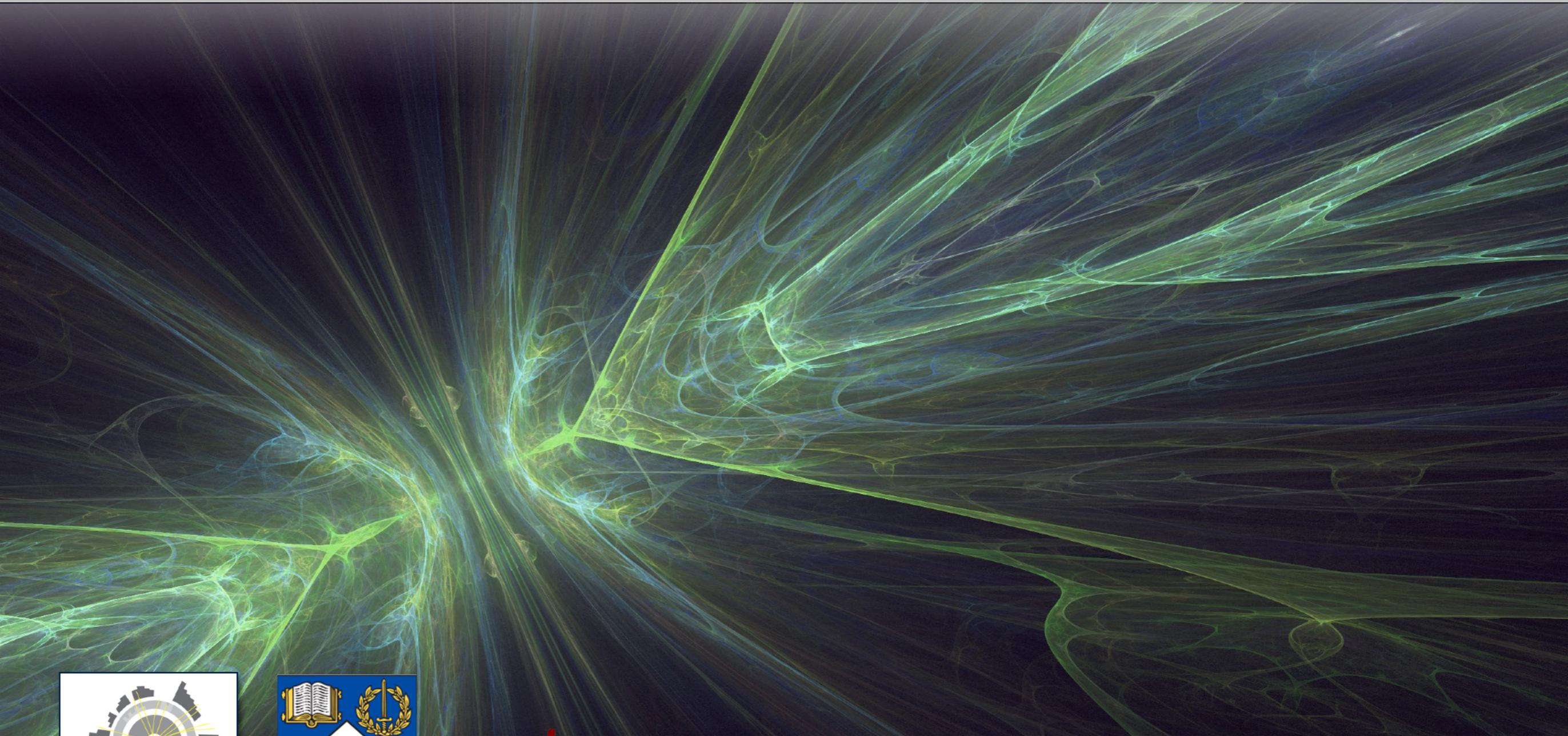


Introduction to Event Generators

Lecture 4: Physics at Hadron Colliders



Peter Skands (Monash University)
11th MCnet School, Lund 2017

PHENO AT THE LHC

What are we really colliding?

Hadrons are composite, with time-dependent structure

Partons within clouds of further partons, constantly being emitted and absorbed

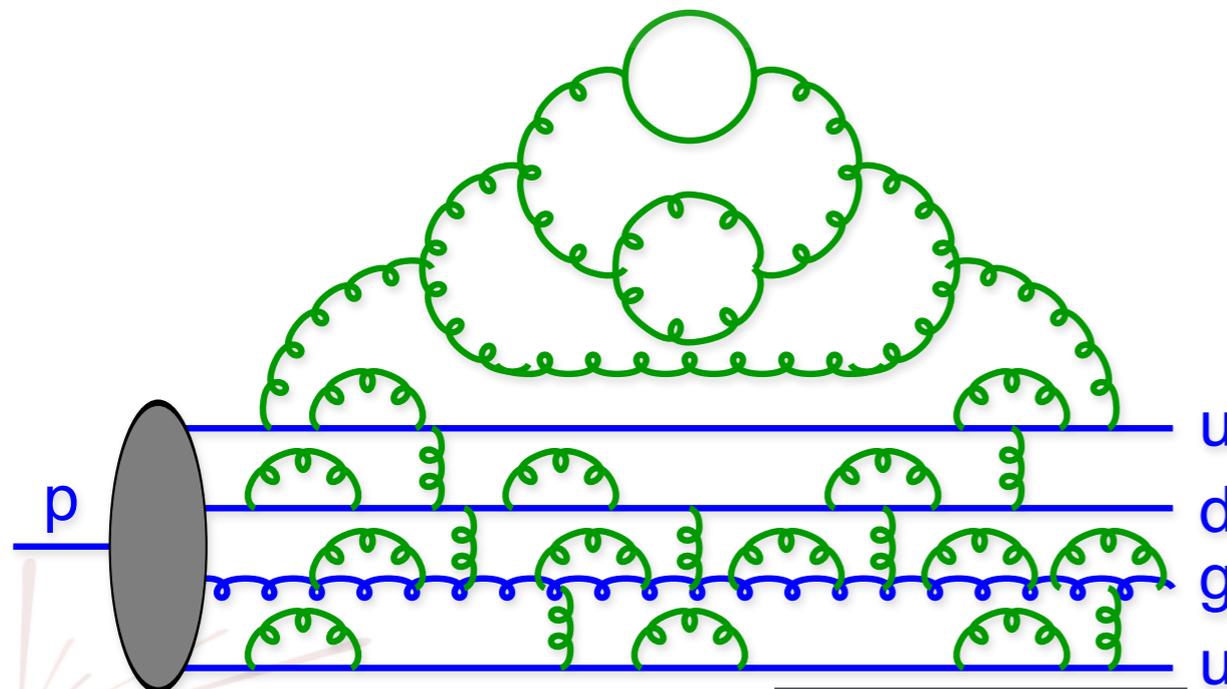
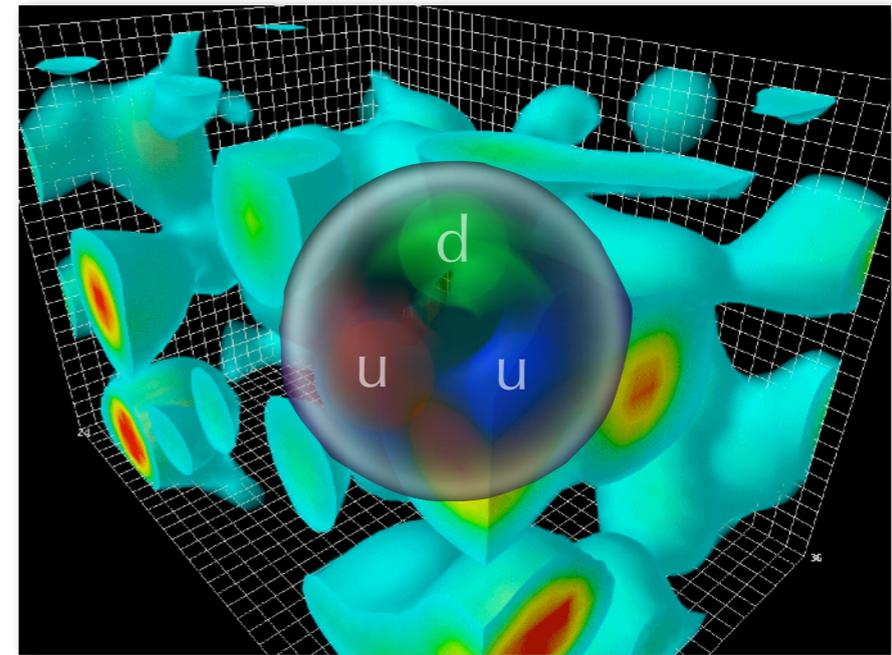


illustration by T. Sjöstrand

(for hadron to remain intact,
virtualities $k^2 < M_h^2$
High-virtuality fluctuations
suppressed by powers of:

$$\frac{\alpha_s M_h^2}{k^2}$$

M_h : mass of hadron
 k^2 : virtuality of fluctuation

SUCH STUFF AS BEAMS ARE MADE OF

Lifetime of typical fluctuation $\sim r_p/c$ (=time it takes light to cross a proton)

$\sim 10^{-23}$ s; Corresponds to a frequency of ~ 500 billion THz

To the LHC, that's slow! (reaches "shutter speeds" thousands of times faster)

$E=h\nu \rightarrow \nu_{\text{LHC}} = 13 \text{ TeV}/h = 3.14$ million billion THz

\rightarrow Protons look "frozen" at moment of collision

But they have a lot more than just two "u" quarks and a "d" inside

Hard to calculate (non-perturbative), so use statistics to parametrise the structure: **parton distribution functions (PDFs)**

@LO: Every so often I will pick a gluon, every so often a quark (antiquark)

Measured at previous colliders (+ now at LHC), as function of energy fraction

Hard scattering knows nothing of the target hadron apart from the fact that it contained the struck parton \rightarrow **factorisation**

[M. Seymour]

HADRON COLLISIONS

Simple question: what does the *average* LHC collision look like?

First question: how many are there?

What is $\sigma_{\text{tot}}(pp)$ at LHC ?

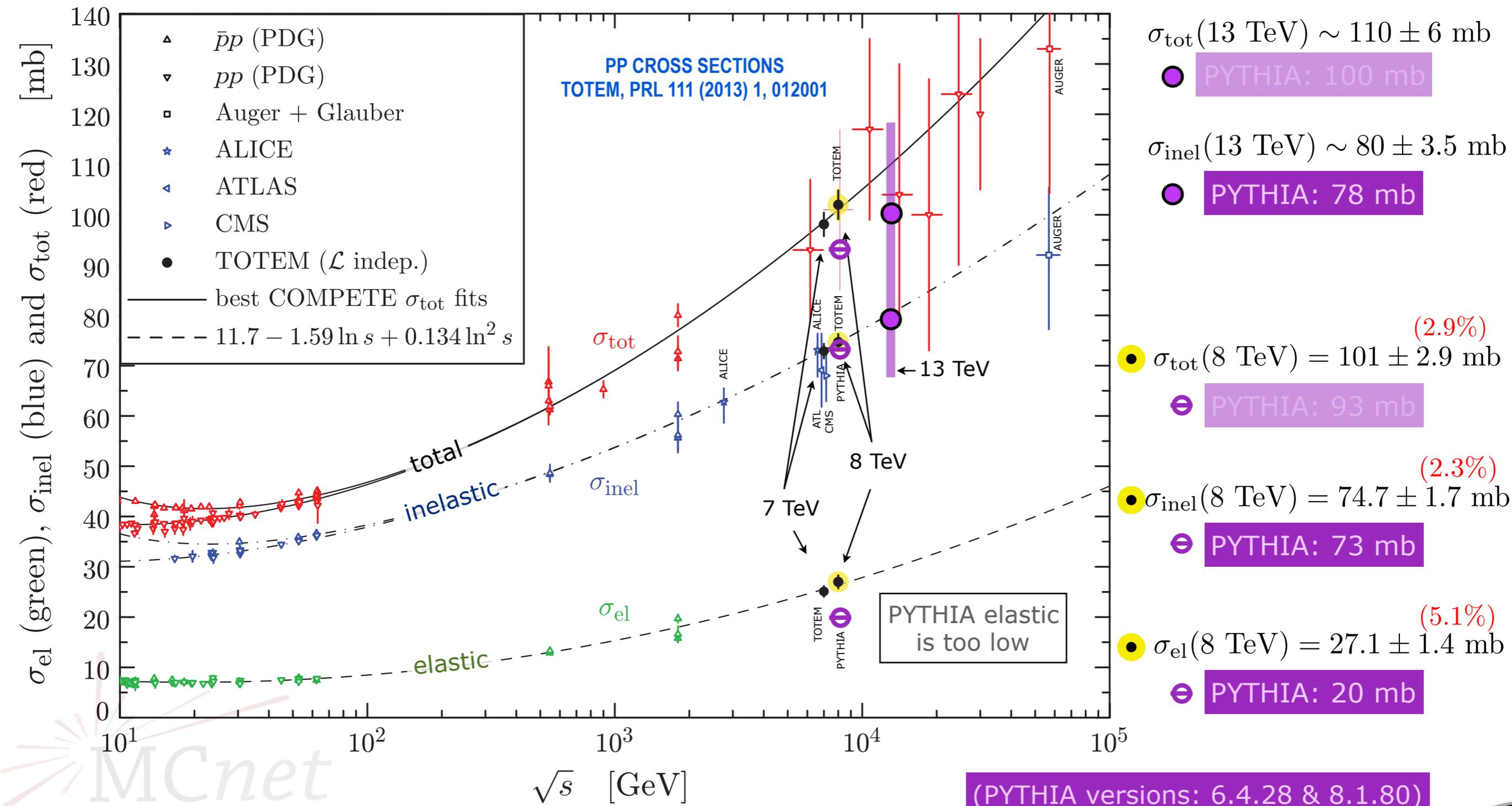
(could we compute it in perturbation theory?)

THE TOTAL CROSS SECTION

$$\sigma_{\text{tot}}(s) = \sigma_{\text{el}}(s) + \sigma_{\text{inel}}(s) \propto s^{0.08} \text{ or } \ln^2(s) ?$$

Donnachie-Landshoff

Froissart-Martin Bound

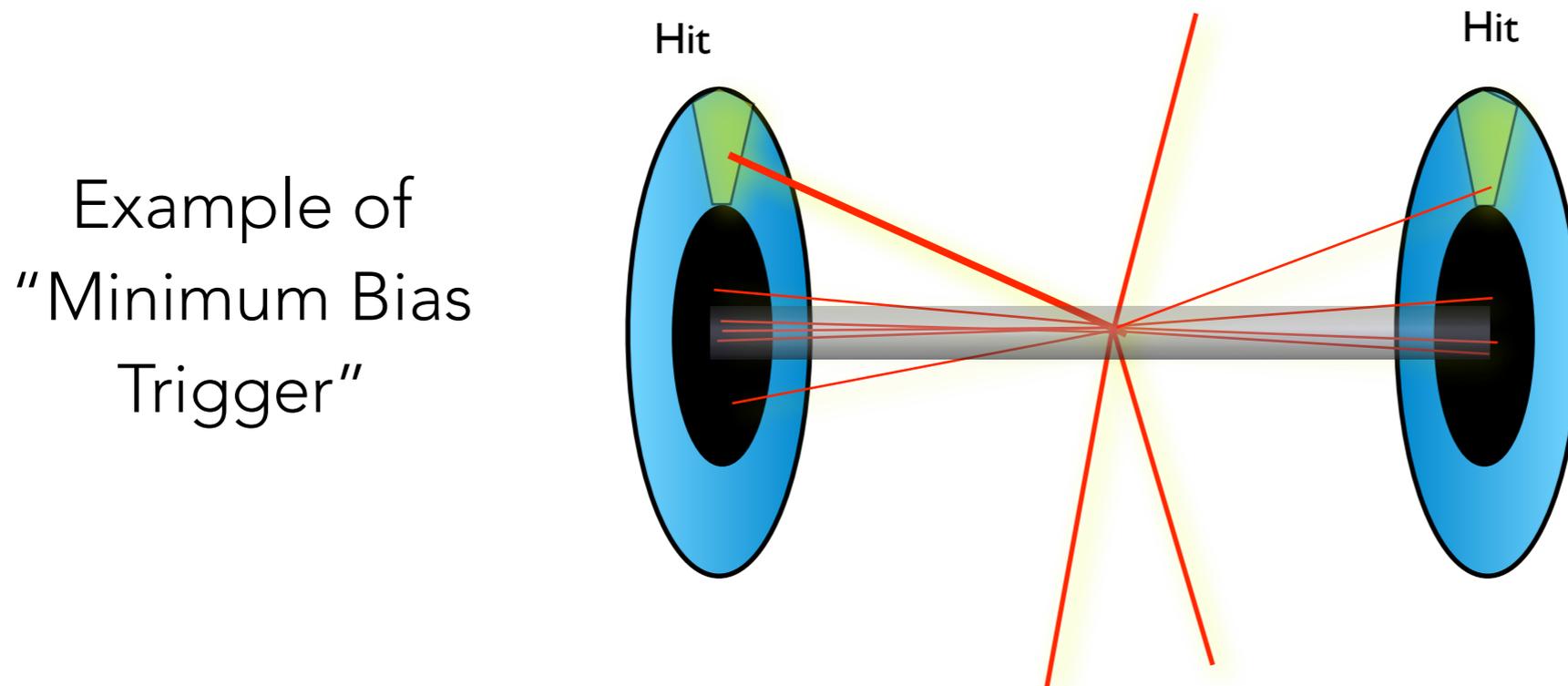


HADRON COLLISIONS

Simple question: what does the *average* LHC collision look like?

First question: how many are there? What is $\sigma_{\text{tot}}(pp)$ at LHC ?

Around 100mb (of which about half is "inelastic, non-diffractive")



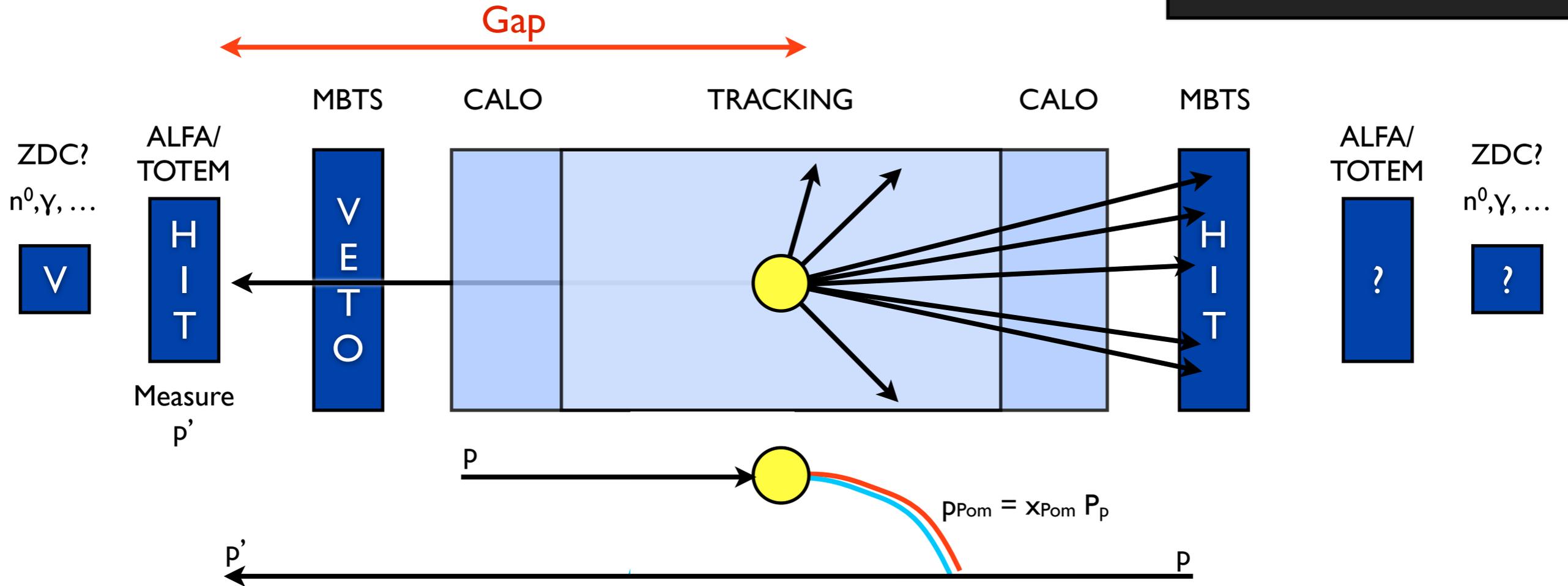
Minimal trigger requirement

At least one hit in some simple and efficient hit counters (typically at large η)
(Double-sided trigger requirement suppresses "single diffraction")

(ASIDE: WHAT IS DIFFRACTION?)

Single Diffraction

Glueball-Proton Collider
with variable E_{CM}

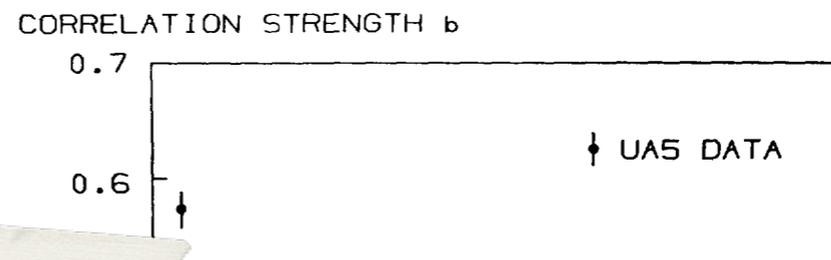
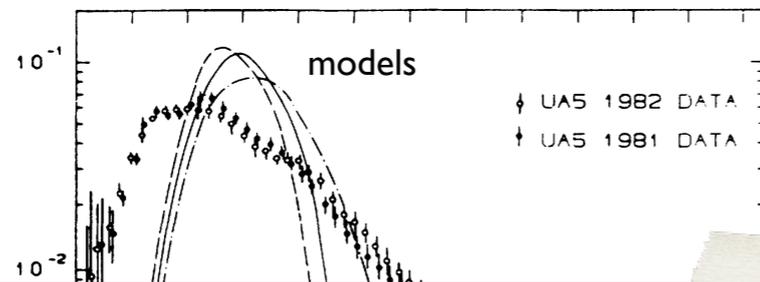


Also:

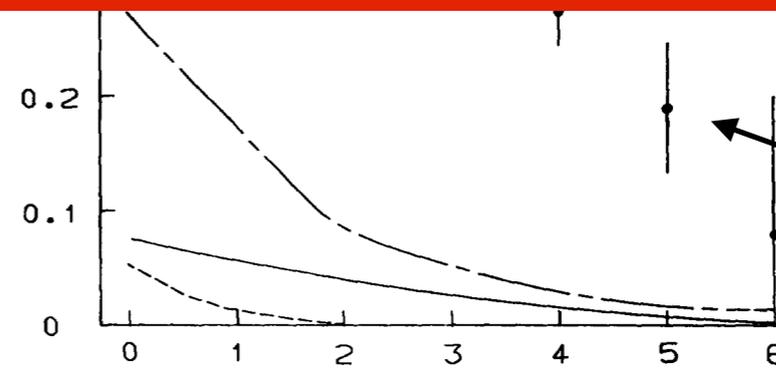
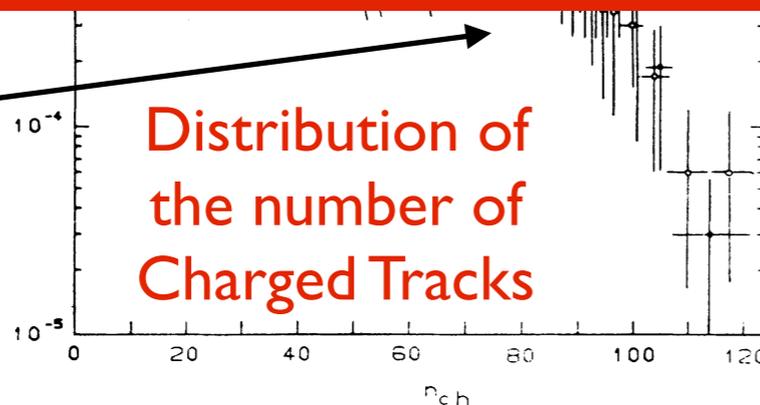
“Double Diffraction”: both protons explode; defined by gap inbetween

“Central Diffraction”: two protons + a central (exclusive) system

MC vs Hadron Collisions



Do not be scared of the failure of physical models
(typically points to more interesting physics)



some mechanism for generating much bigger fluctuations in multiplicity (here: of charged tracks)

Distribution of the number of Charged Tracks

Correlation Strength (forward-backward)

some global (quantum) number tells the entire event to fluctuate up or down?

FIG. 3. Charged-multiplicity distribution at 540 GeV, UA5 results (Ref. 32) vs simple models: dashed low p_T only, full including hard scatterings, dash-dotted also including initial- and final-state radiation.

HARD INTERACTIONS IN HADRON COLLISIONS

1983: the "Pedestal Effect"

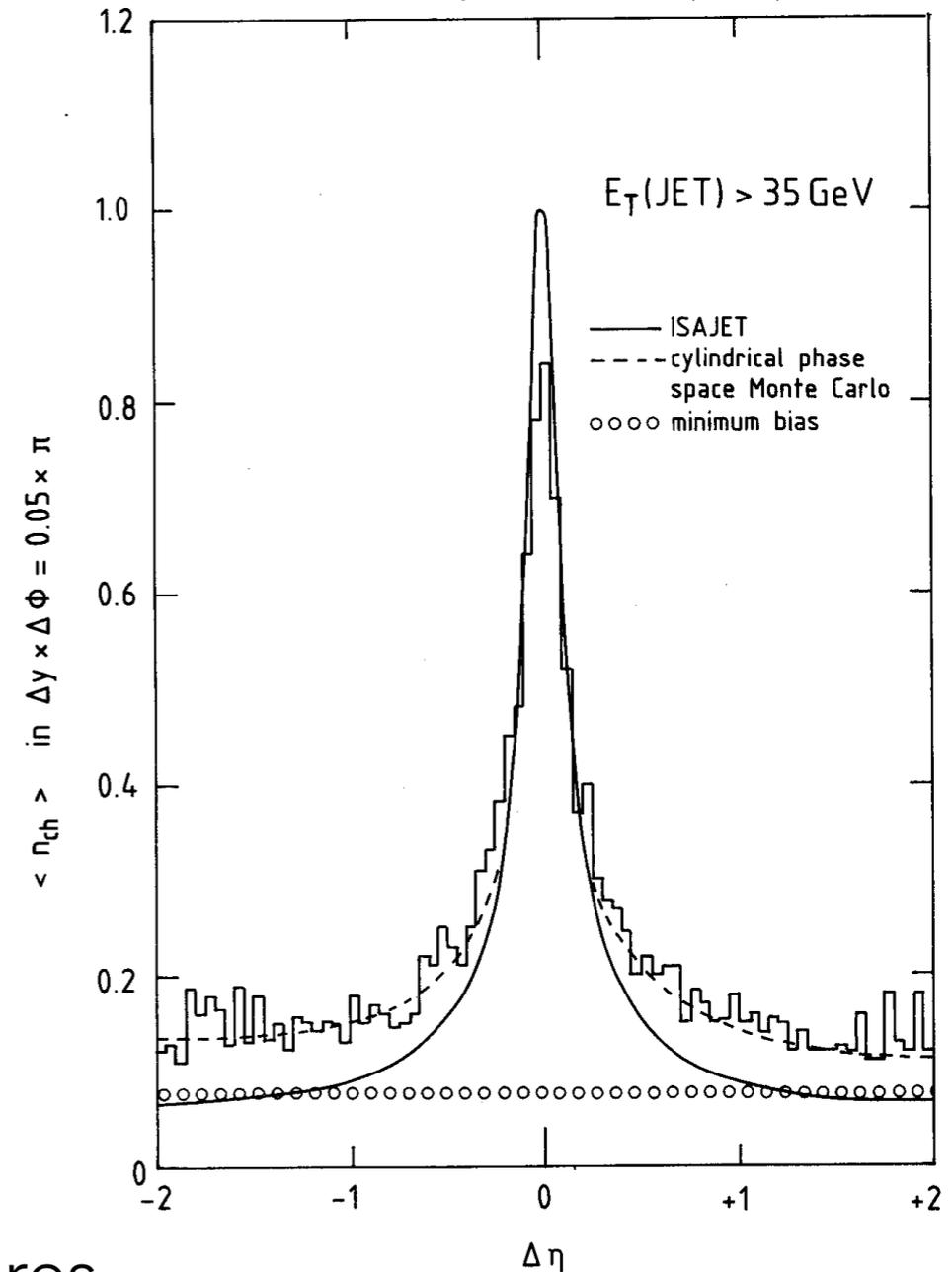
UA1: $p\bar{p}$ at $\sqrt{s} = 540$ GeV

Studies of jets with E_T up to 100 GeV

"Outside the [jet], a constant E_T plateau is observed, whose height is independent of the jet E_T . Its value is substantially higher than the one observed for minimum bias events."

In hadron collisions, hard jets sit on "pedestals" of increased particle production extending far from the jet cores.

Phys. Lett. B 132 (1983) 214-222

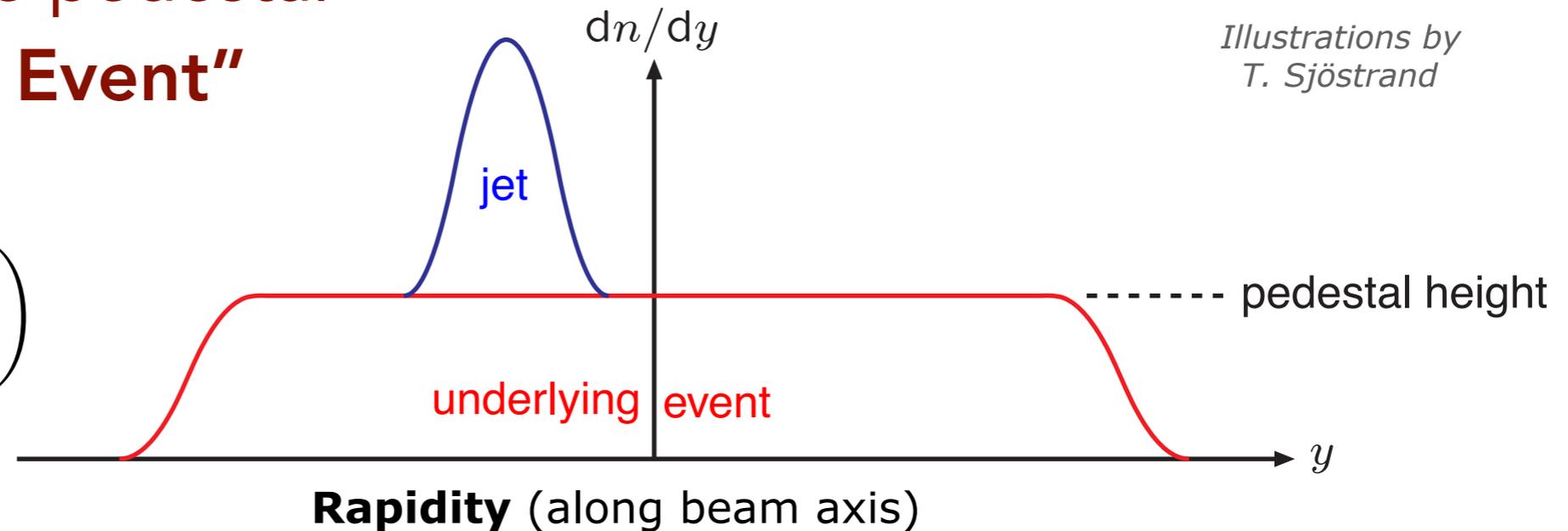


DISSECTING THE PEDESTAL

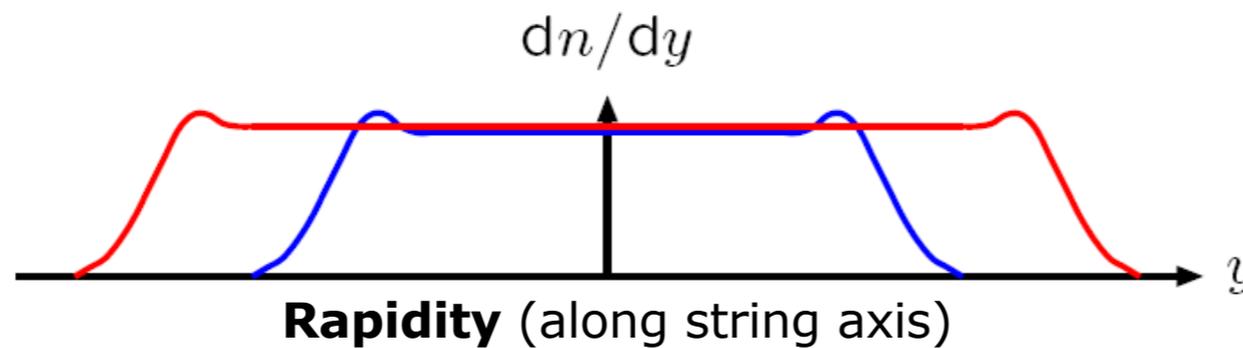
Today, we call the pedestal
"the Underlying Event"

*Illustrations by
T. Sjöstrand*

$$y = \frac{1}{2} \ln \left(\frac{E + p_z}{E - p_z} \right)$$



Looks like something we've seen before ... ?



(but pedestal too high to be just one string ...)

FROM HARD TO SOFT

Factorization and IR safety

Main tools for jet calculations

Corrections suppressed by powers of $\Lambda_{\text{QCD}}/Q_{\text{Hard}}$

Soft QCD / Minimum-Bias

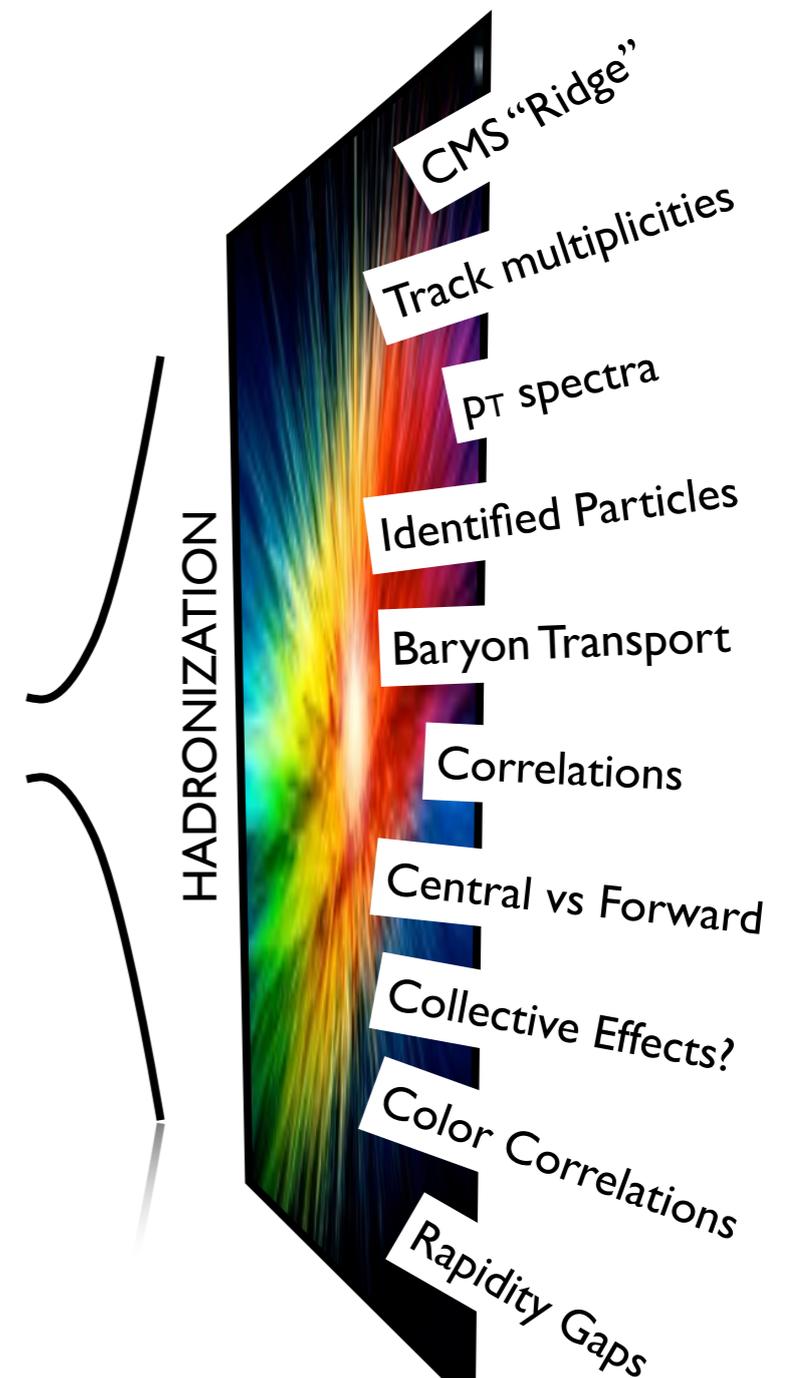
NO HARD SCALE

Typical Q scales $\sim \Lambda_{\text{QCD}}$
Extremely sensitive to IR effects
→ Excellent LAB for studying IR effects

$\sim \infty$ statistics for min-bias

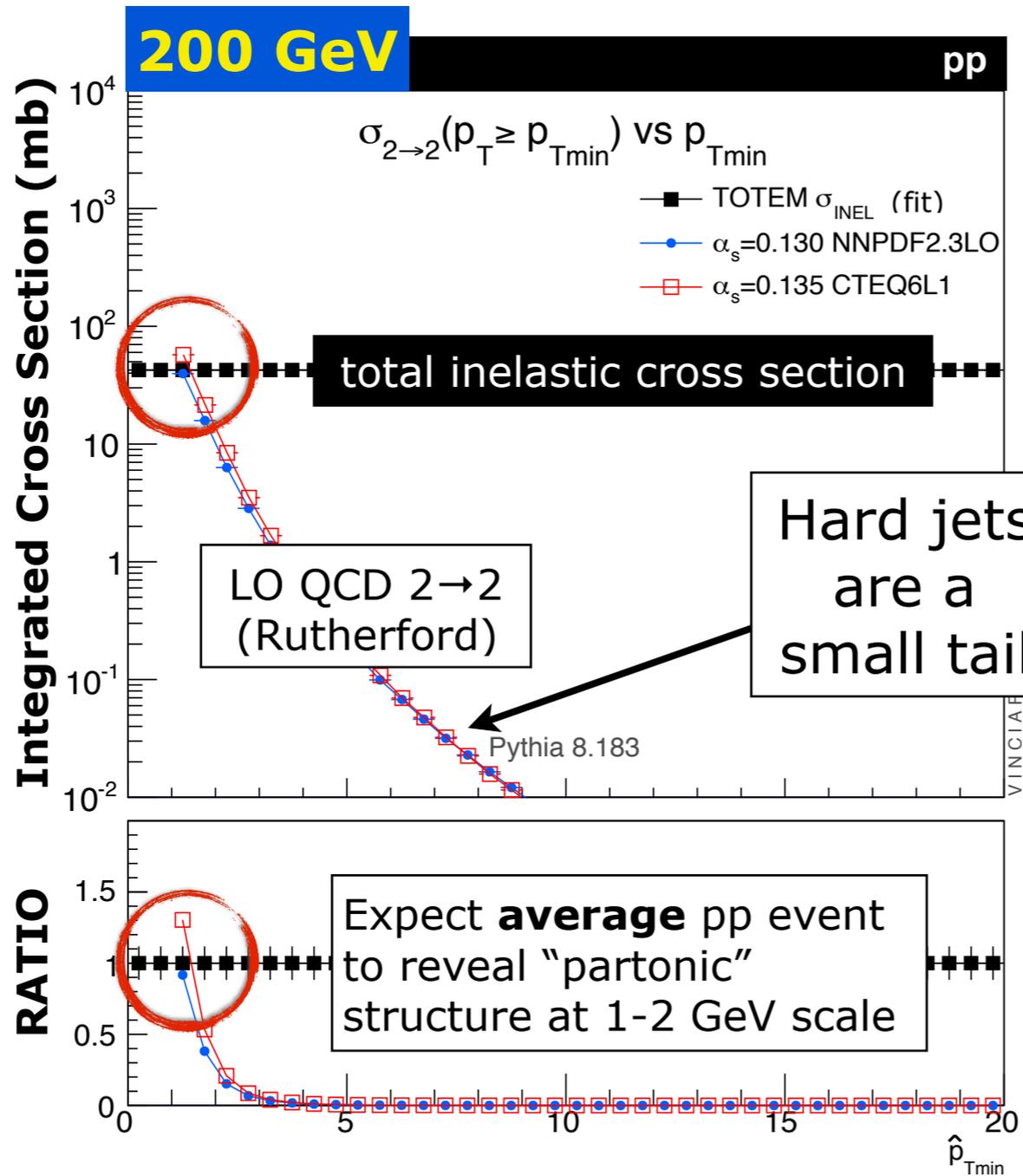
→ Access tails, limits

Universality: Recycling PU ↔ MB ↔ UE



IS THERE NO HARD SCALE?

Compare total (inelastic) hadron-hadron cross section to calculated parton-parton (LO QCD 2→2) cross section



Hard jets are a small tail

Leading-Order pQCD

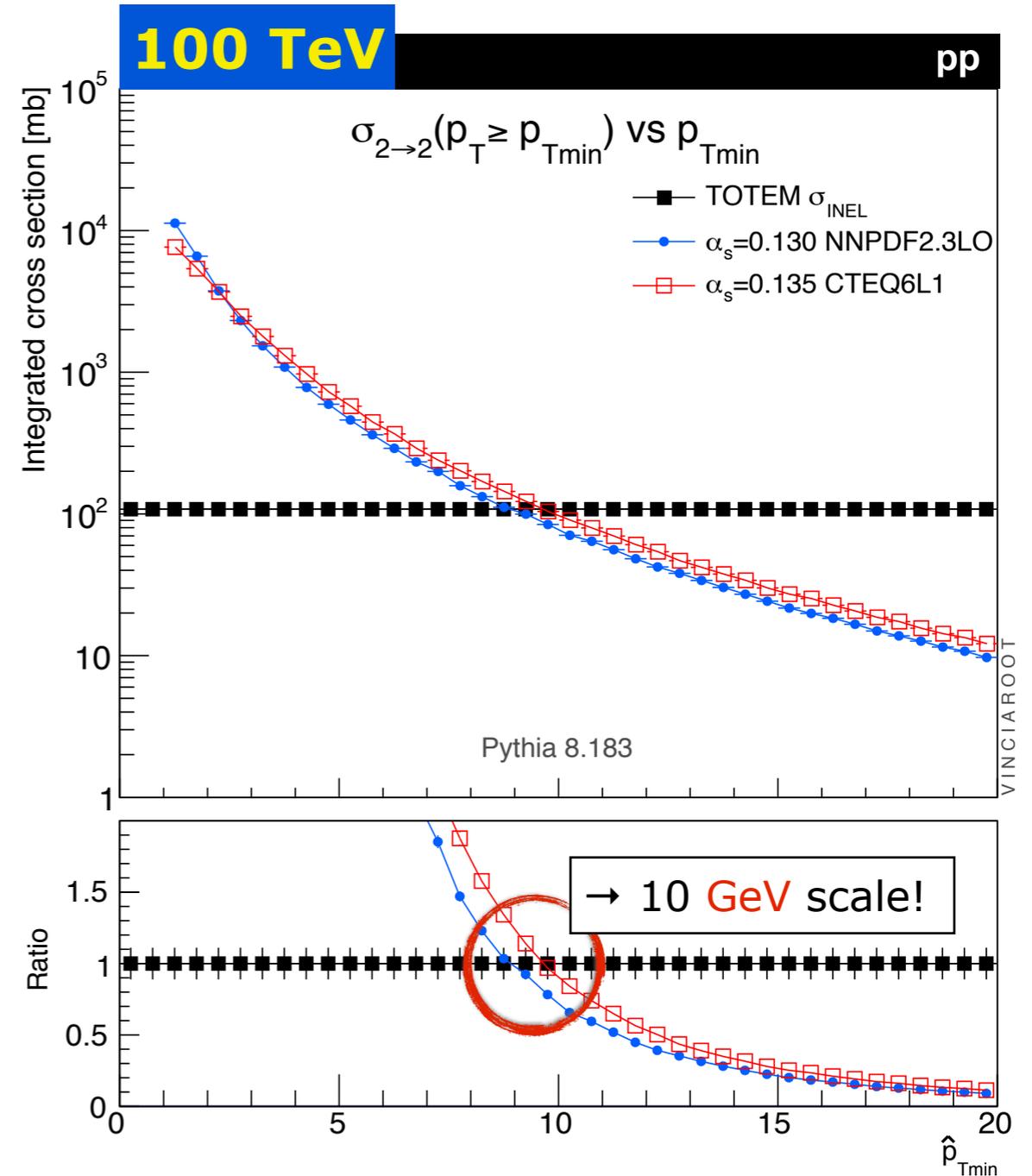
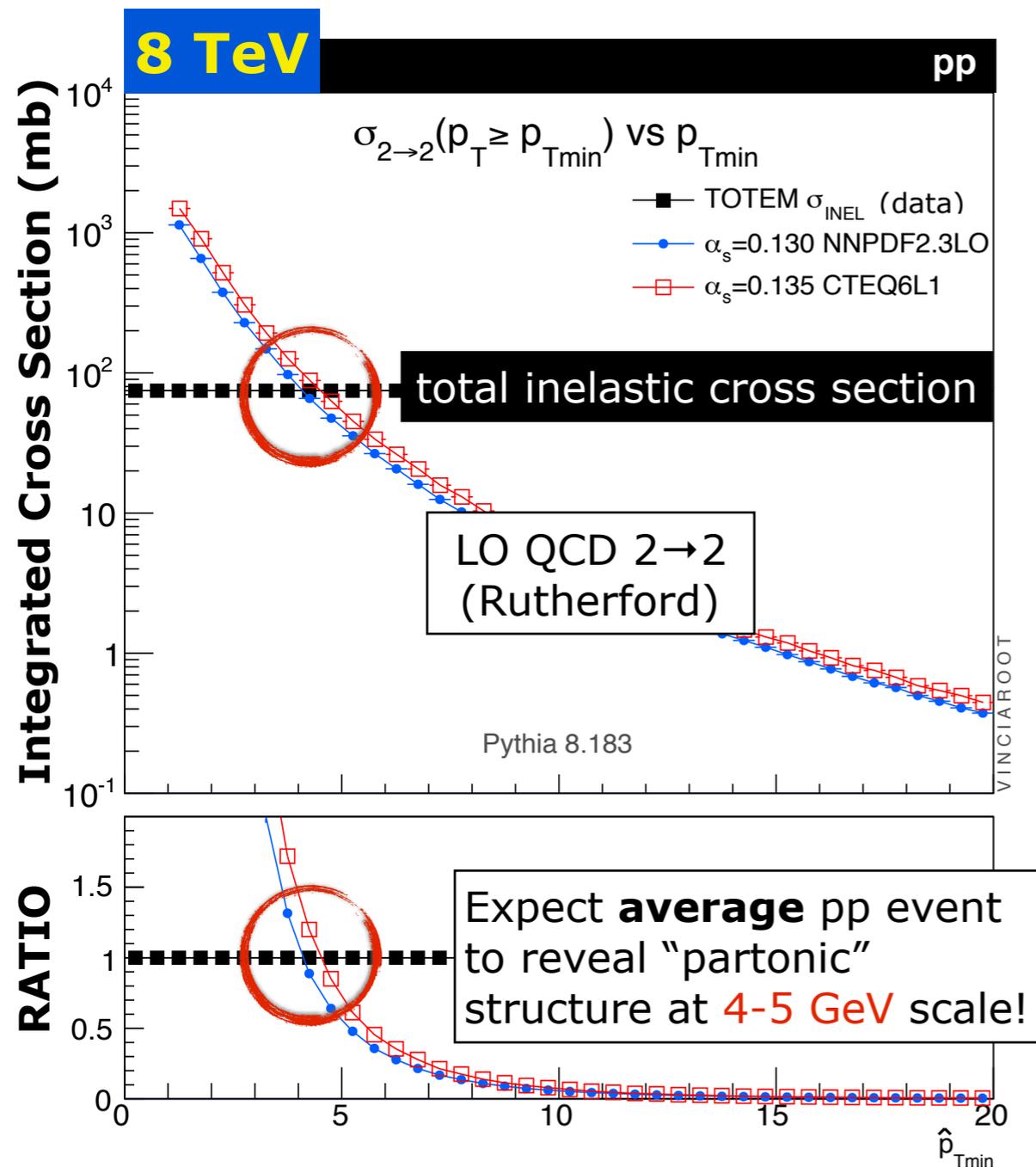
$$\int_{p_{\perp, \min}^2} dp_{\perp}^2 \frac{d\sigma_{Dijet}}{dp_{\perp}^2}$$

$$d\sigma_{2 \rightarrow 2} \propto \frac{dp_{\perp}^2}{p_{\perp}^4} \otimes \text{PDFs}$$

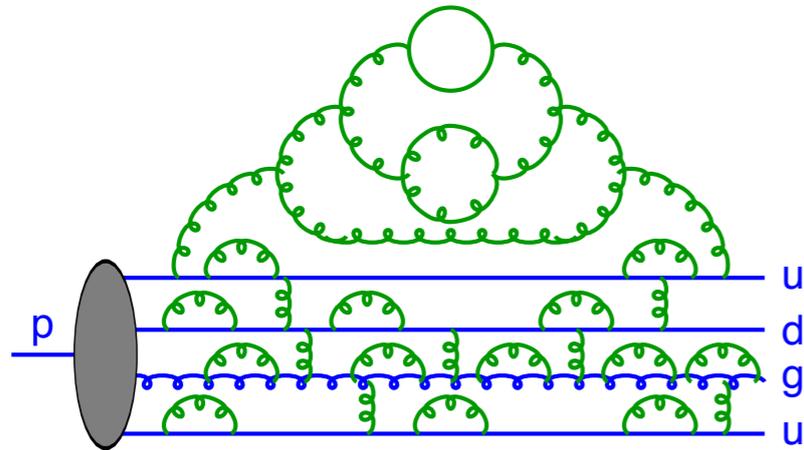


→ 8 TEV → 100 TEV

→ Trivial calculation indicates hard scales in min-bias



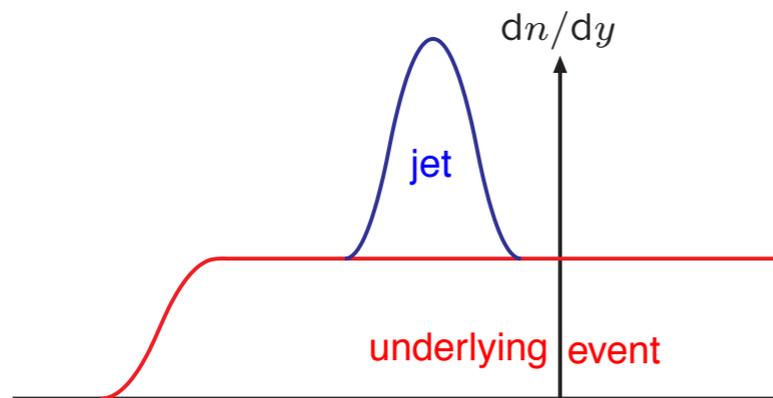
SUMMARY FOR NOW: WE KNOW 3 THINGS



1) Hadrons are composite

Factorisation: hard interaction picks out a single parton; what about the rest?

At some level, multiple-parton-interactions must occur (only a question of how often)



2) Events with a hard trigger are accompanied by an "underlying event"

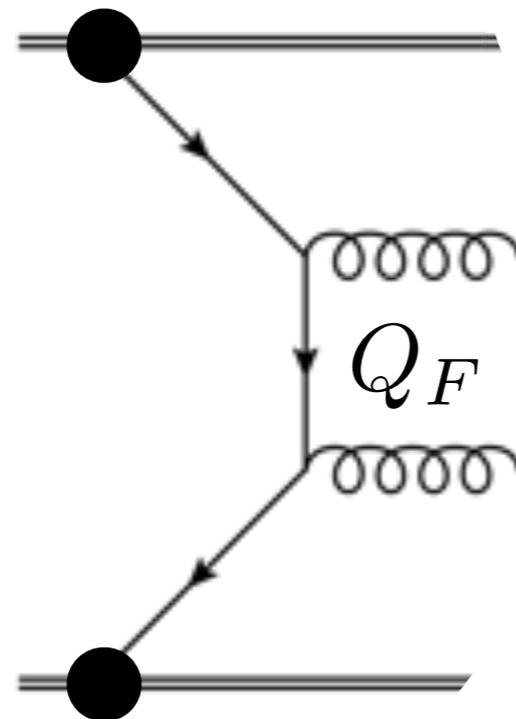
Looks too high to be just one string

Multiple colour exchanges ?

3) Simple calculations indicate the presence of (semi)hard scales even when no hard trigger is imposed ("minimum bias")

PHYSICS OF THE PEDESTAL

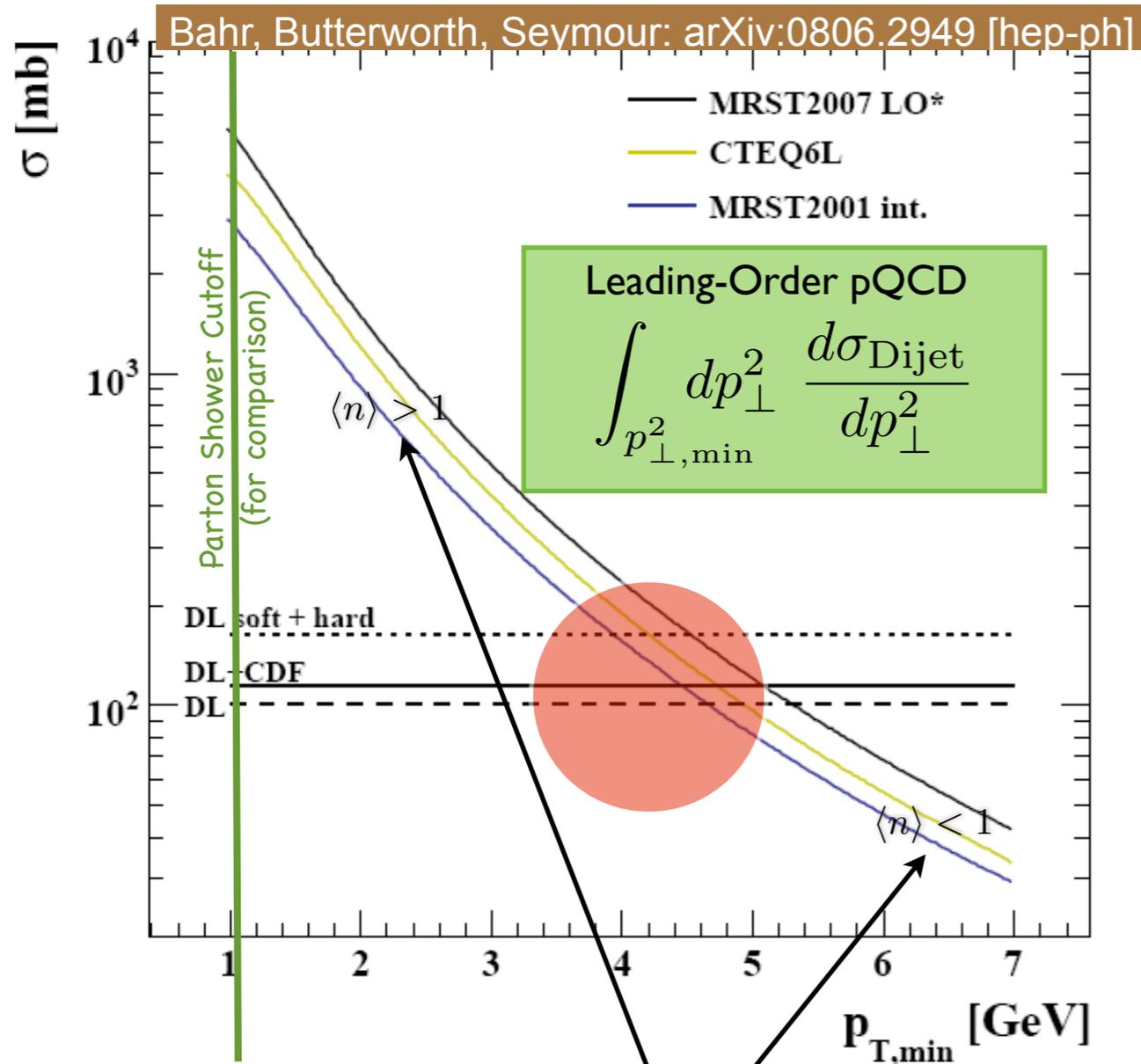
Factorization: Subdivide Calculation



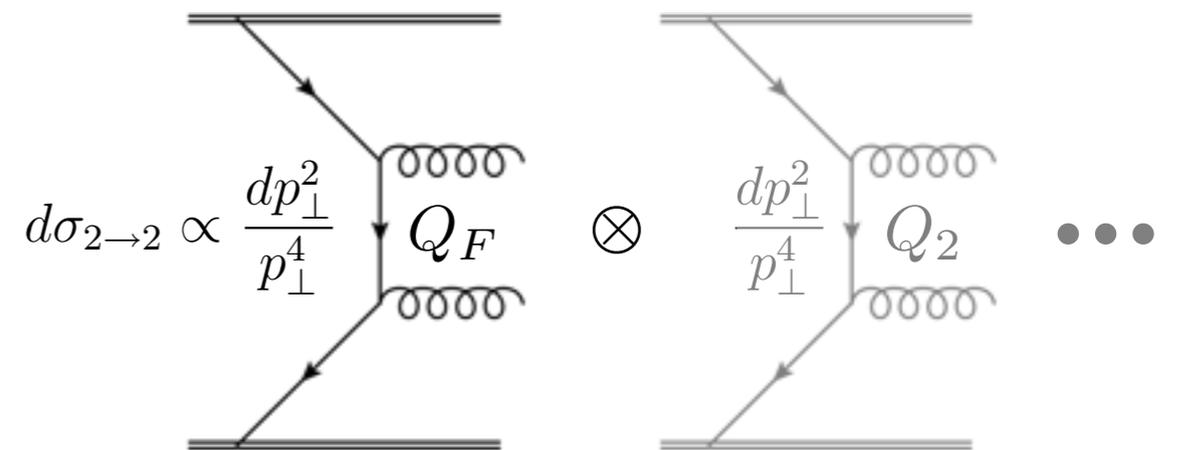
- Multiple Parton Interactions* go beyond existing theorems
- perturbative short-distance physics in Underlying Event
 - Need to generalize factorization to MPI

Multiple Parton Interactions

= Allow several parton-parton interactions per hadron-hadron collision. Requires extended factorization ansatz.



Earliest MC model ("old" PYTHIA 6 model)
Sjöstrand, van Zijl PRD36 (1987) 2019



Lesson from bremsstrahlung in pQCD:
divergences \rightarrow fixed-order breaks down
Perturbation theory still ok, with
resummation (unitarity)

\rightarrow Resum dijets?
Yes \rightarrow MPI!

$$\sigma_{2 \rightarrow 2}(p_{\perp \min}) = \langle n \rangle(p_{\perp \min}) \sigma_{\text{tot}}$$

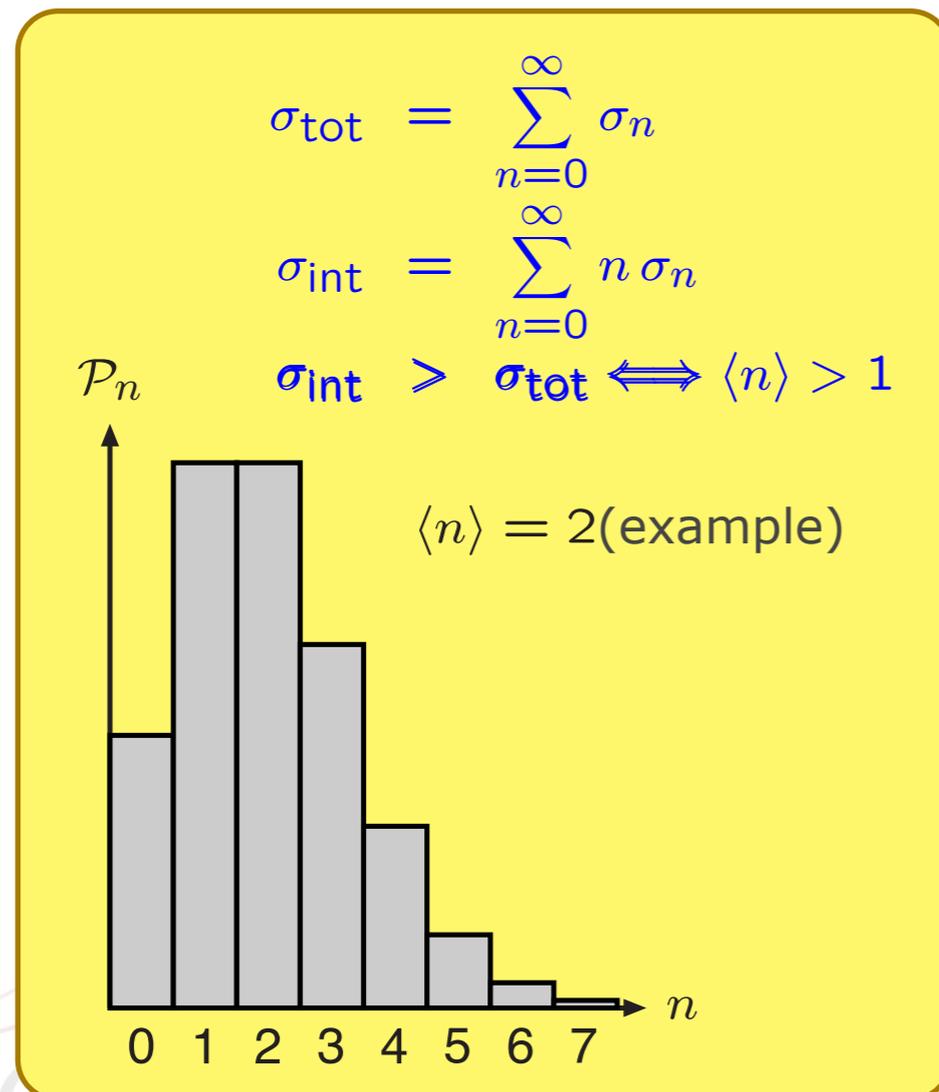
Parton-Parton Cross Section

Hadron-Hadron Cross Section

HOW MANY?

Naively $\langle n_{2 \rightarrow 2}(p_{\perp \min}) \rangle = \frac{\sigma_{2 \rightarrow 2}(p_{\perp \min})}{\sigma_{\text{tot}}}$

If the interactions are assumed \sim independent (naive factorization) \rightarrow Poisson



$$\mathcal{P}_n = \frac{\langle n \rangle^n}{n!} e^{-\langle n \rangle}$$

Real Life

Color screening: $\sigma_{2 \rightarrow 2} \rightarrow 0$ for $p_{\perp} \rightarrow 0$

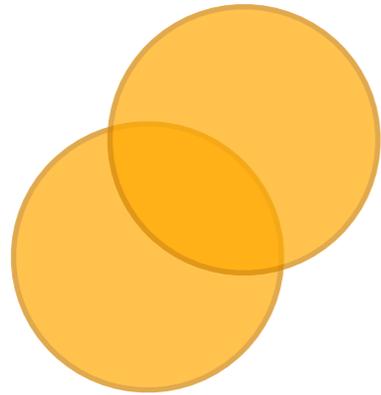
Momentum conservation suppresses high-n tail

Impact-parameter dependence

+ physical correlations

\rightarrow not simple product

IMPACT PARAMETER

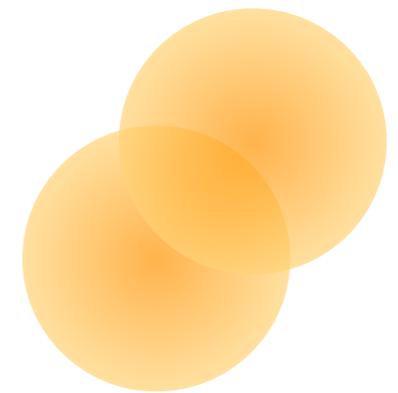


1. **Simple Geometry** (in impact-parameter plane)

Simplest idea: smear PDFs across a uniform disk of size πr_p^2
→ simple geometric overlap factor ≤ 1 in dijet cross section
Some collisions have the full overlap, others only partial
→ Poisson distribution with different mean $\langle n \rangle$ at each b

2. More realistic **Proton b-shape**

Smear PDFs across a non-uniform disk
MC models use Gaussians or **more**/less peaked
Overlap factor = convolution of two such distributions



→ Poisson distribution with different mean $\langle n \rangle$ at each b
“Lumpy Peaks” → large matter overlap enhancements, higher $\langle n \rangle$

Note: this is an *effective* description. Not the actual proton mass density.
E.g., peak in overlap function ($\gg 1$) can represent unlikely configurations with huge overlap enhancement. Typically use total σ_{inel} as normalization.

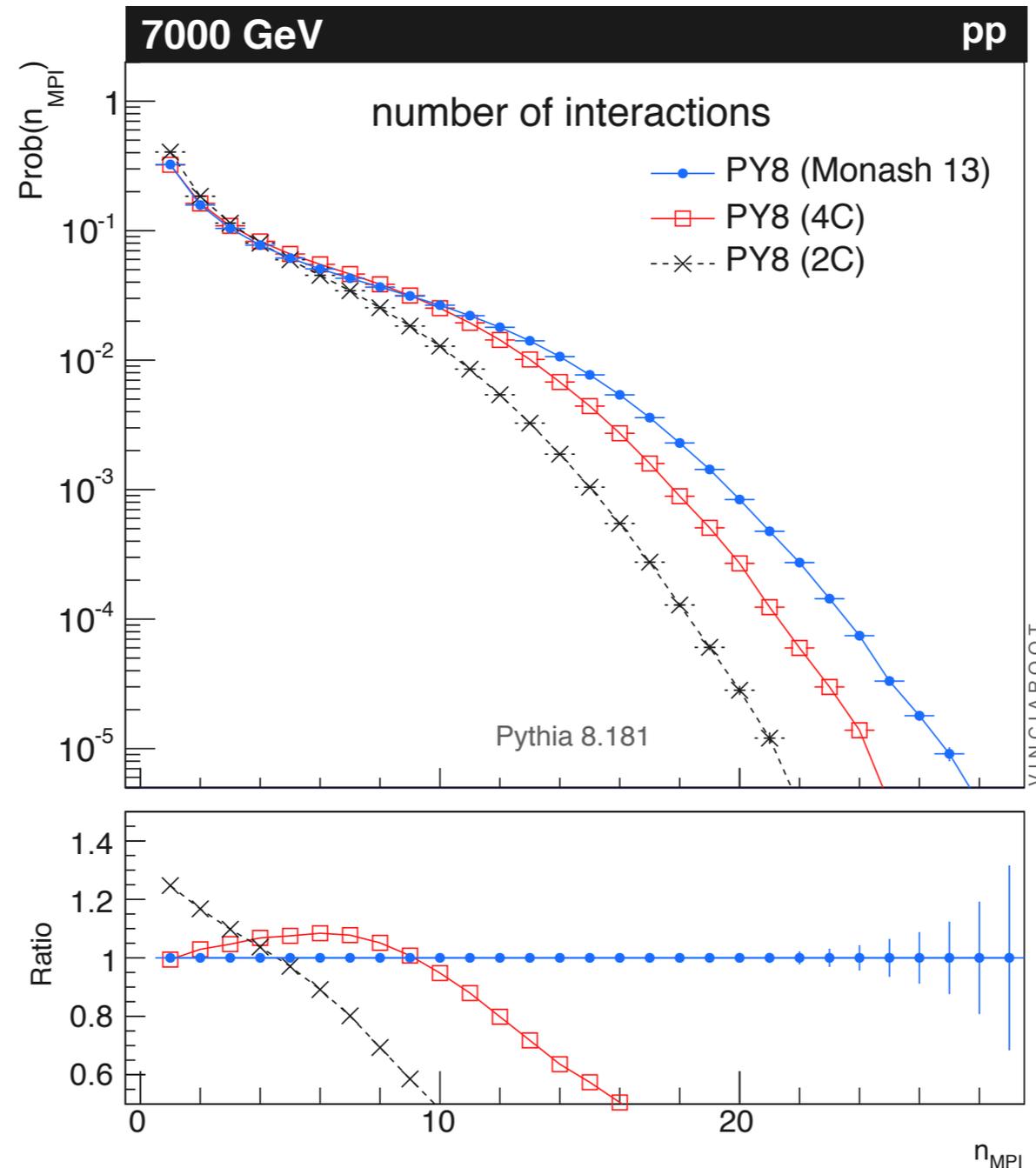
NUMBER OF MPI

*

Minimum-Bias pp collisions at 7 TeV

Averaged over all
pp impact
parameters

(Really:
averaged over all
pp overlap
enhancement
factors)



*note: can be
arbitrarily soft

1: A SIMPLE MODEL

A minimal model incorporating single-parton factorization, perturbative unitarity, and energy-and-momentum conservation

Take literally $\sigma_{2 \rightarrow 2}(p_{\perp \min}) = \langle n \rangle(p_{\perp \min}) \sigma_{\text{tot}}$

Parton-Parton Cross Section Hadron-Hadron Cross Section

1. Choose $p_{T \min}$ cutoff

= main tuning parameter

2. Interpret $\langle n \rangle(p_{T \min})$ as mean of Poisson distribution

Equivalent to assuming all parton-parton interactions equivalent and independent ~ each take an instantaneous “snapshot” of the proton

3. Generate n parton-parton interactions (pQCD $2 \rightarrow 2$)

Veto if total beam momentum exceeded \rightarrow overall (E,p) cons

4. Add impact-parameter dependence $\rightarrow \langle n \rangle = \langle n \rangle(b)$ Ordinary CTEQ, MSTW, NNPDF, ...

Assume factorization of transverse and longitudinal d.o.f., \rightarrow PDFs : $f(x,b) = f(x)g(b)$

b distribution \propto EM form factor \rightarrow **JIMMY model** Butterworth, Forshaw, Seymour Z.Phys. C72 (1996) 637

Constant of proportionality = second main tuning parameter

5. Add separate class of “soft” (zero- p_T) interactions representing

interactions with $p_T < p_{T \min}$ and require $\sigma_{\text{soft}} + \sigma_{\text{hard}} = \sigma_{\text{tot}}$

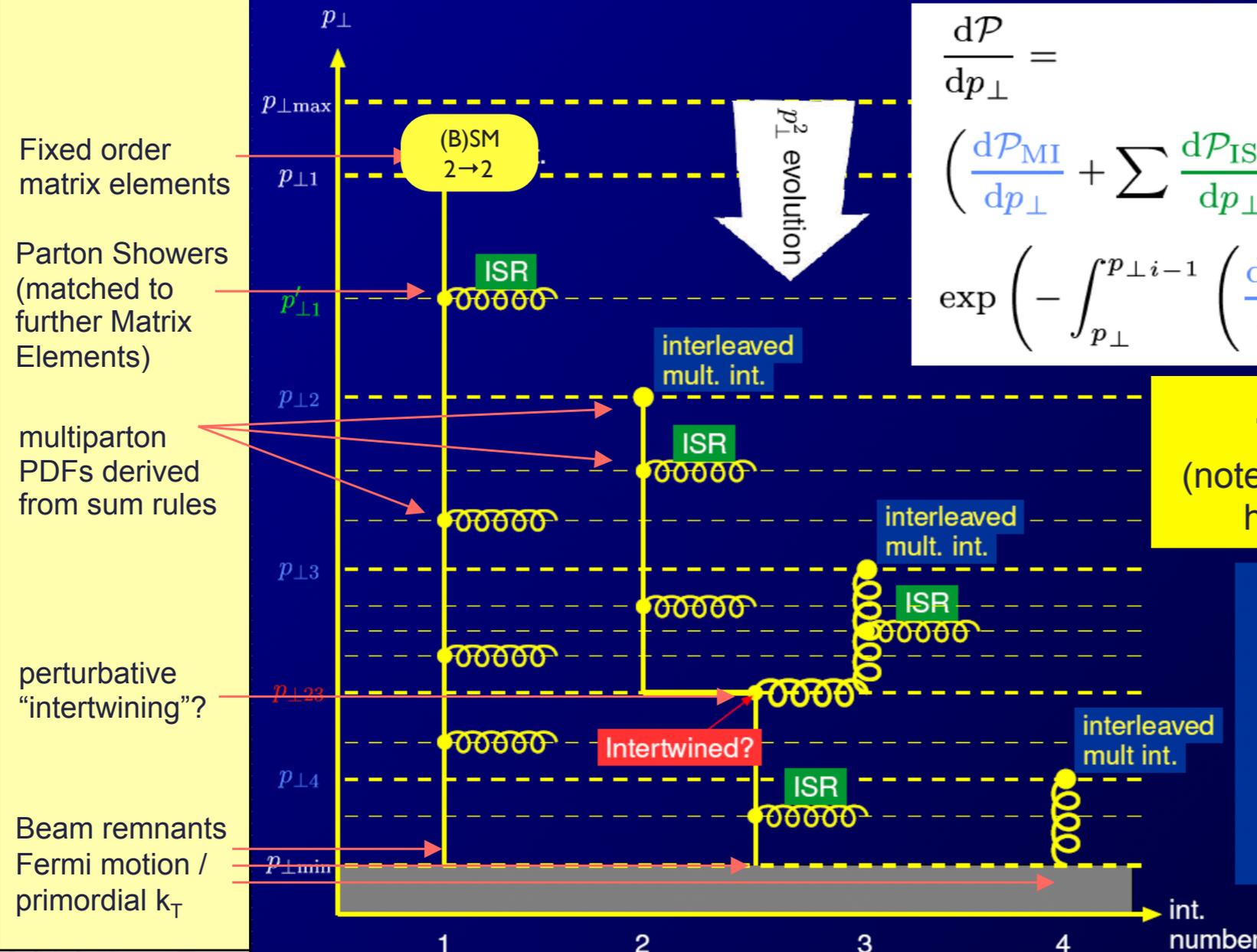
\rightarrow **Herwig++ model** Bähr et al, arXiv:0905.4671



2: INTERLEAVED EVOLUTION

Sjöstrand, P.S., JHEP 0403 (2004) 053; EPJ C39 (2005) 129

Add exclusivity progressively by evolving *everything* downwards.



$$\frac{d\mathcal{P}}{dp_{\perp}} = \left(\frac{d\mathcal{P}_{\text{MI}}}{dp_{\perp}} + \sum \frac{d\mathcal{P}_{\text{ISR}}}{dp_{\perp}} + \sum \frac{d\mathcal{P}_{\text{JI}}}{dp_{\perp}} \right) \times \exp \left(- \int_{p_{\perp}}^{p_{\perp}^i-1} \left(\frac{d\mathcal{P}_{\text{MI}}}{dp'_{\perp}} + \sum \frac{d\mathcal{P}_{\text{ISR}}}{dp'_{\perp}} + \sum \frac{d\mathcal{P}_{\text{JI}}}{dp'_{\perp}} \right) dp'_{\perp} \right)$$

→ Underlying Event
(note: interactions correlated in colour: hadronization not independent)

~ "Finegraining"
→ correlations between all perturbative activity at successively smaller scales

MC vs Hadron Collisions

Fluctuations in n_{mpi} \rightarrow Bigger (global) fluctuations

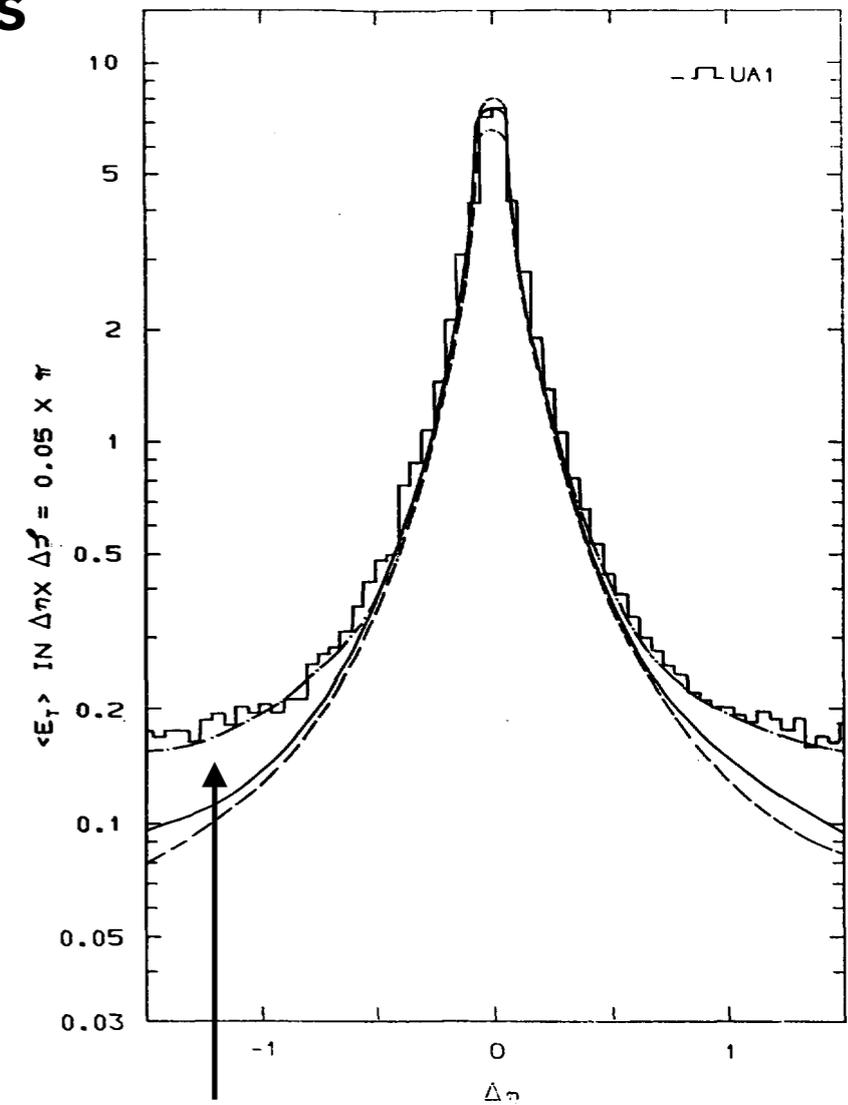
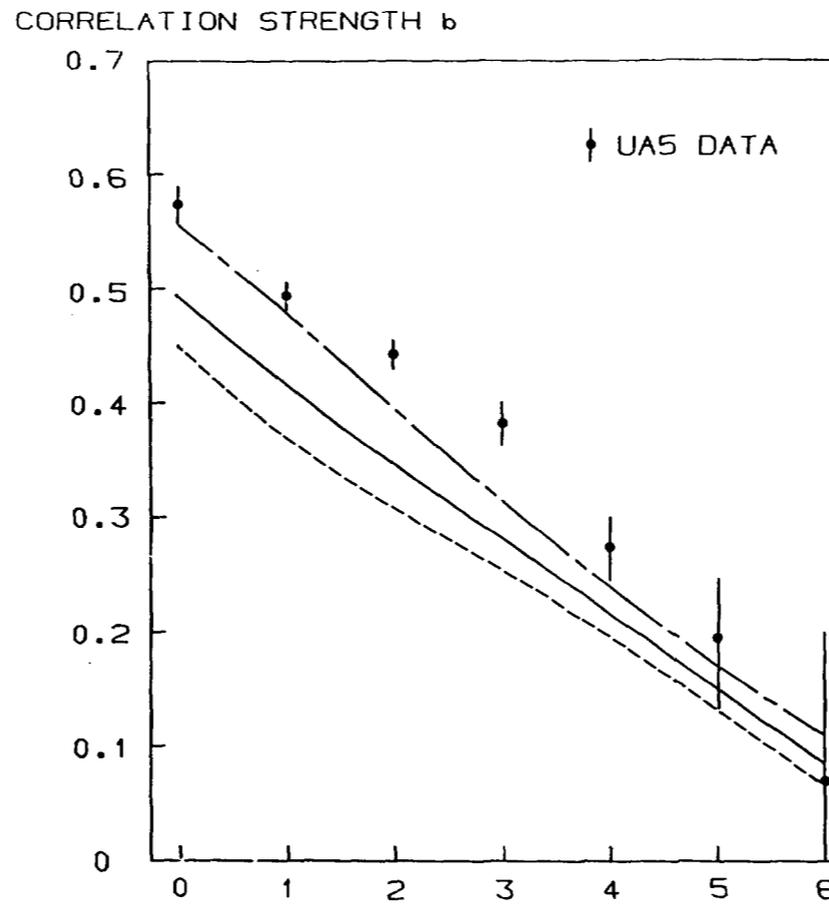
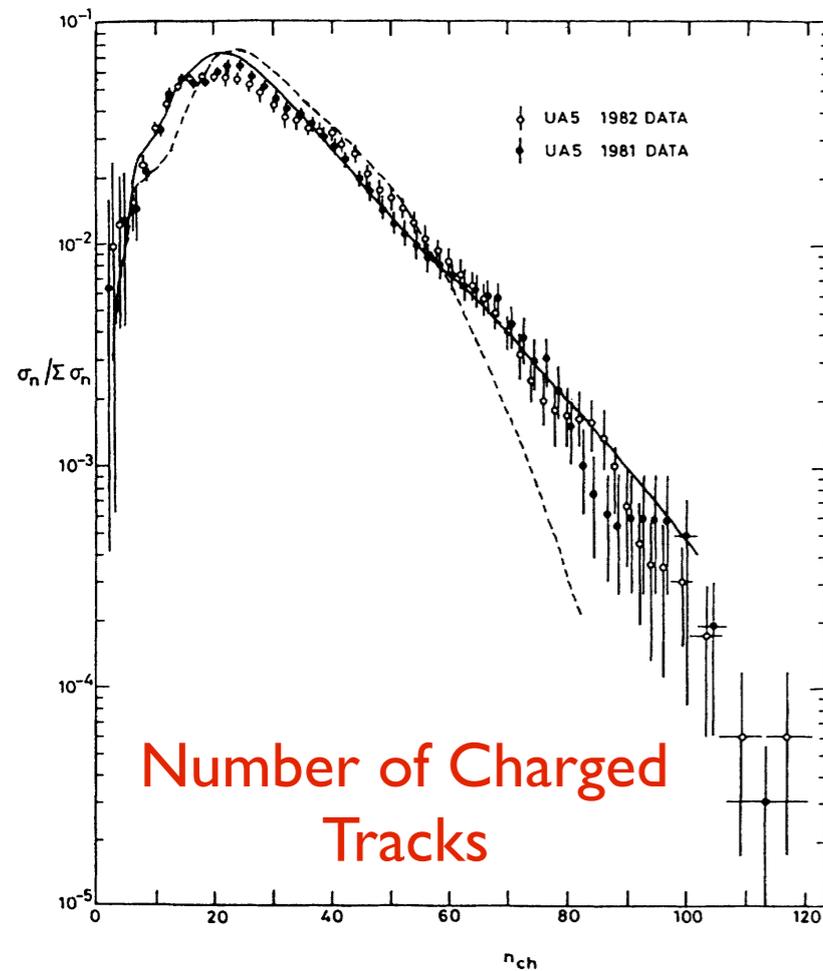


FIG. 12. Charged-multiplicity distribution at 540 GeV, UA5 results (Ref. 32) vs multiple-interaction model with variable impact parameter: solid line, double-Gaussian matter distribution; dashed line, with fix impact parameter [i.e., $\bar{O}_0(b)$].

Impact-parameter dependence \rightarrow UE

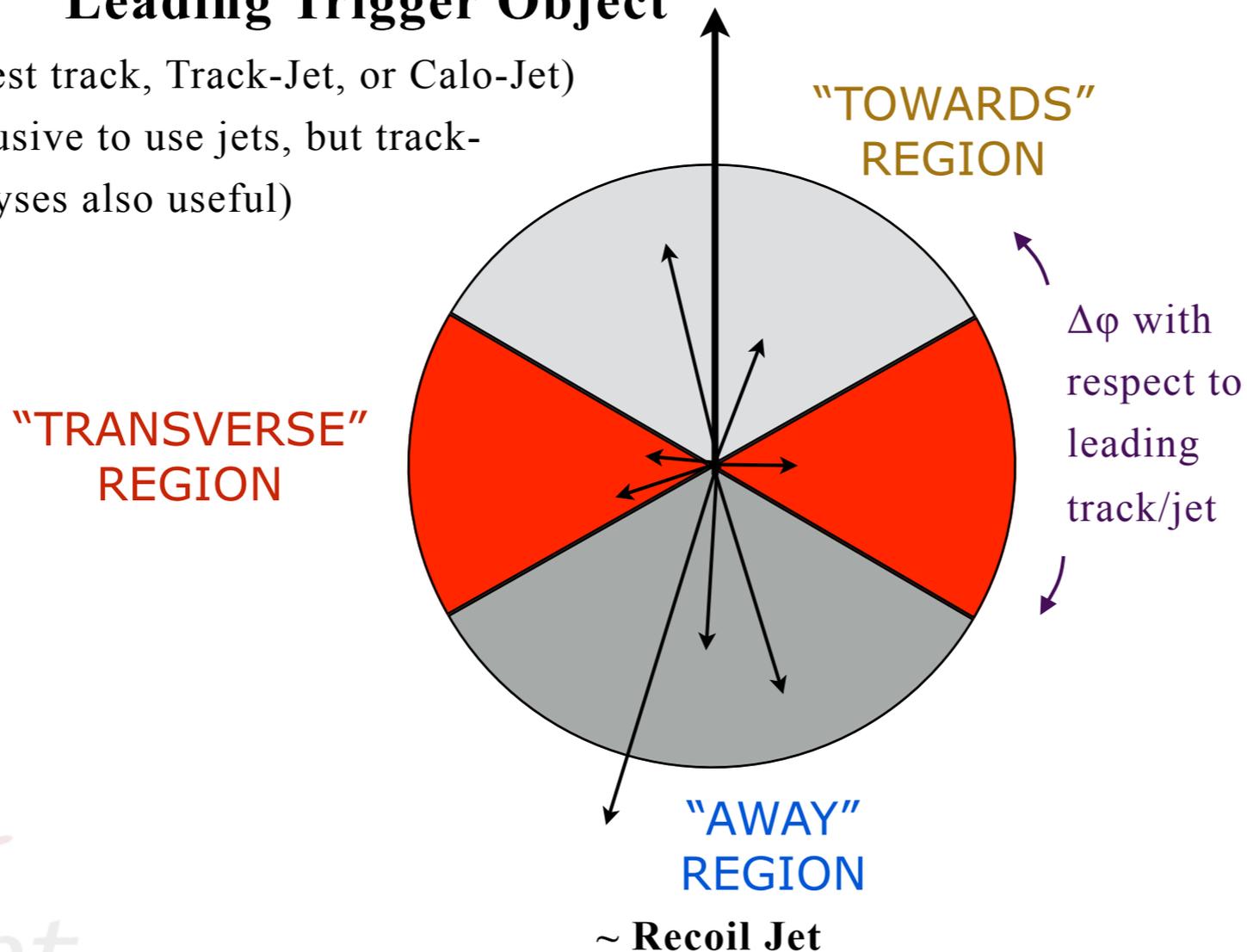
CHARACTERISING THE UNDERLYING EVENT

(The "Rick Field" UE Plots - the same Field as in Field-Feynman)

There are many UE variables.
The most important is $\langle \Sigma p_T \rangle$ in the "Transverse Region"

Leading Trigger Object

(e.g., hardest track, Track-Jet, or Calo-Jet)
(more inclusive to use jets, but track-based analyses also useful)



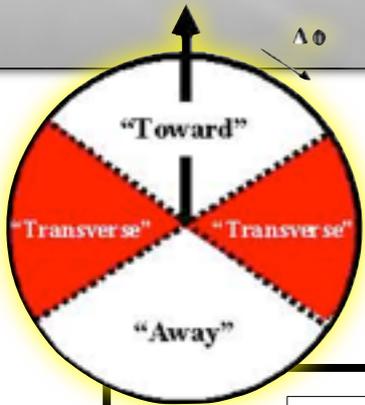
"Transverse Region" (TRNS)

Sensitive to activity at right angles to the hardest jets

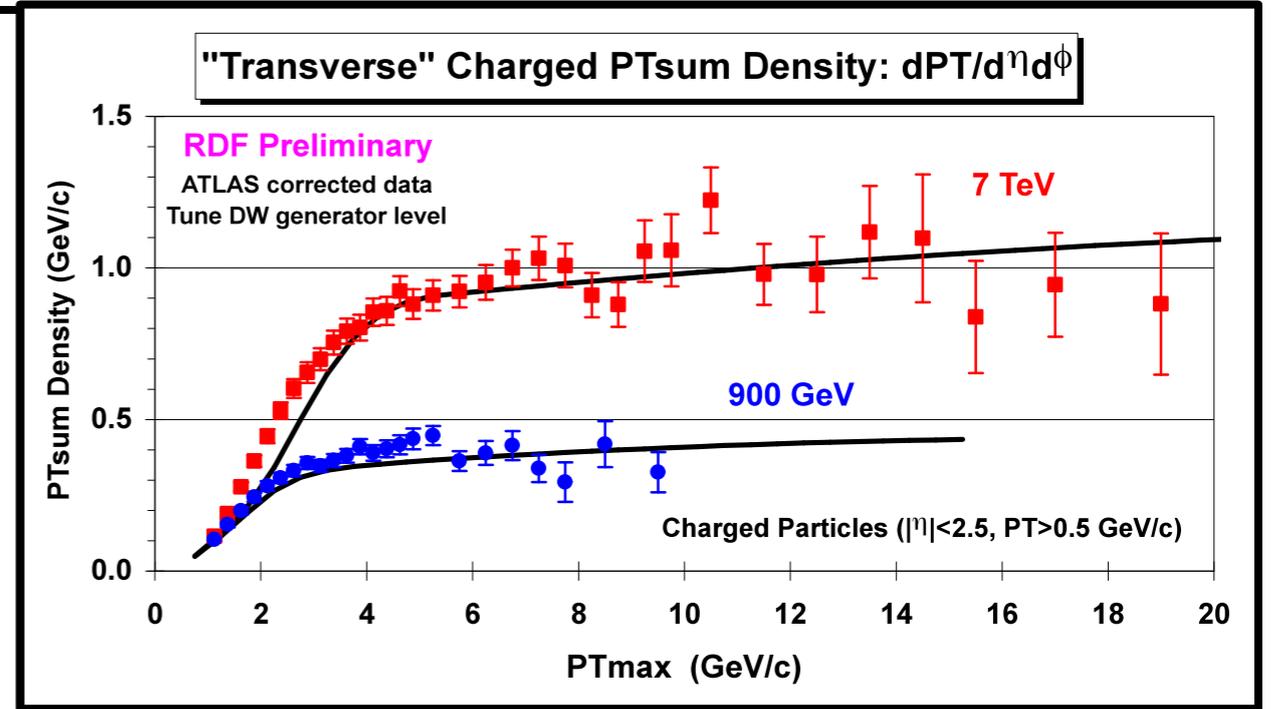
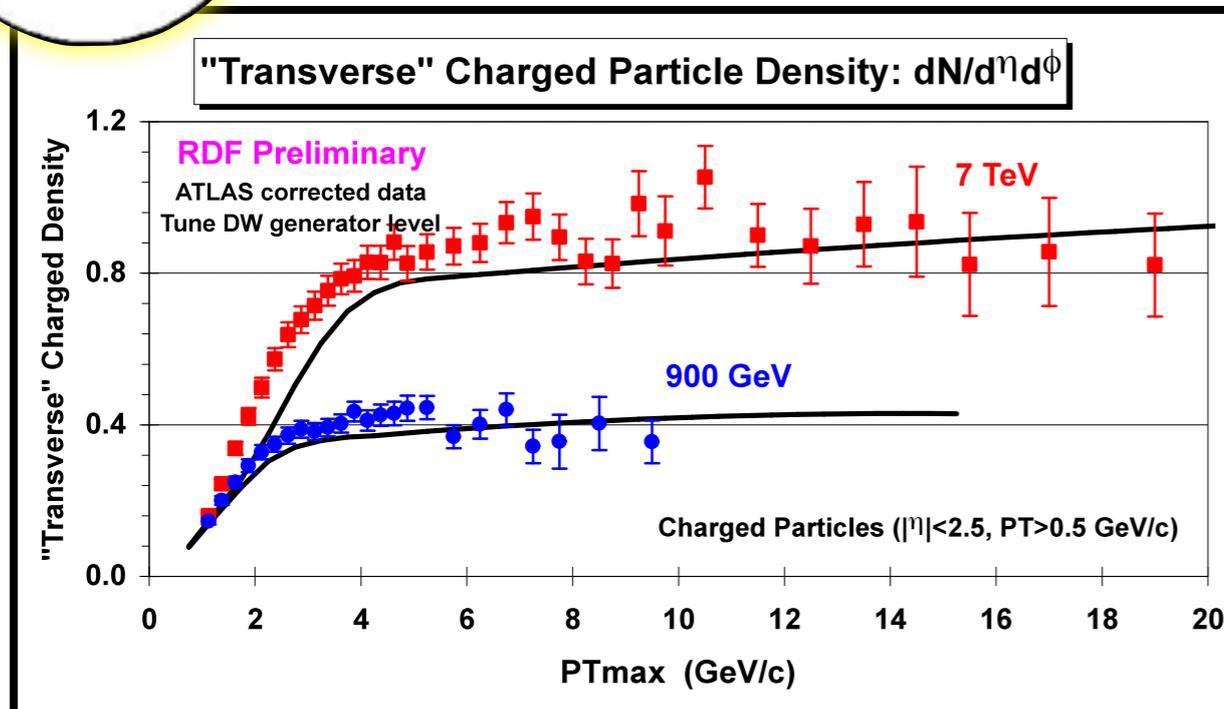
→ Useful definition of Underlying Event

THE PEDESTAL

(NOW CALLED THE UNDERLYING EVENT)



LHC from 900 to 7000 GeV - ATLAS



Track Density (TRANS)

(Not Infrared Safe)

Large Non-factorizable Corrections

Prediction off by $\approx 10\%$

Sum(pT) Density (TRANS)

(more) Infrared Safe

Large Non-factorizable Corrections

Prediction off by $< 10\%$

Truth is in the eye of the beholder:

R. Field: "See, I told you!"

Y. Gehrstein: "they have to fudge it again"

MIN-BIAS VS UNDERLYING EVENT

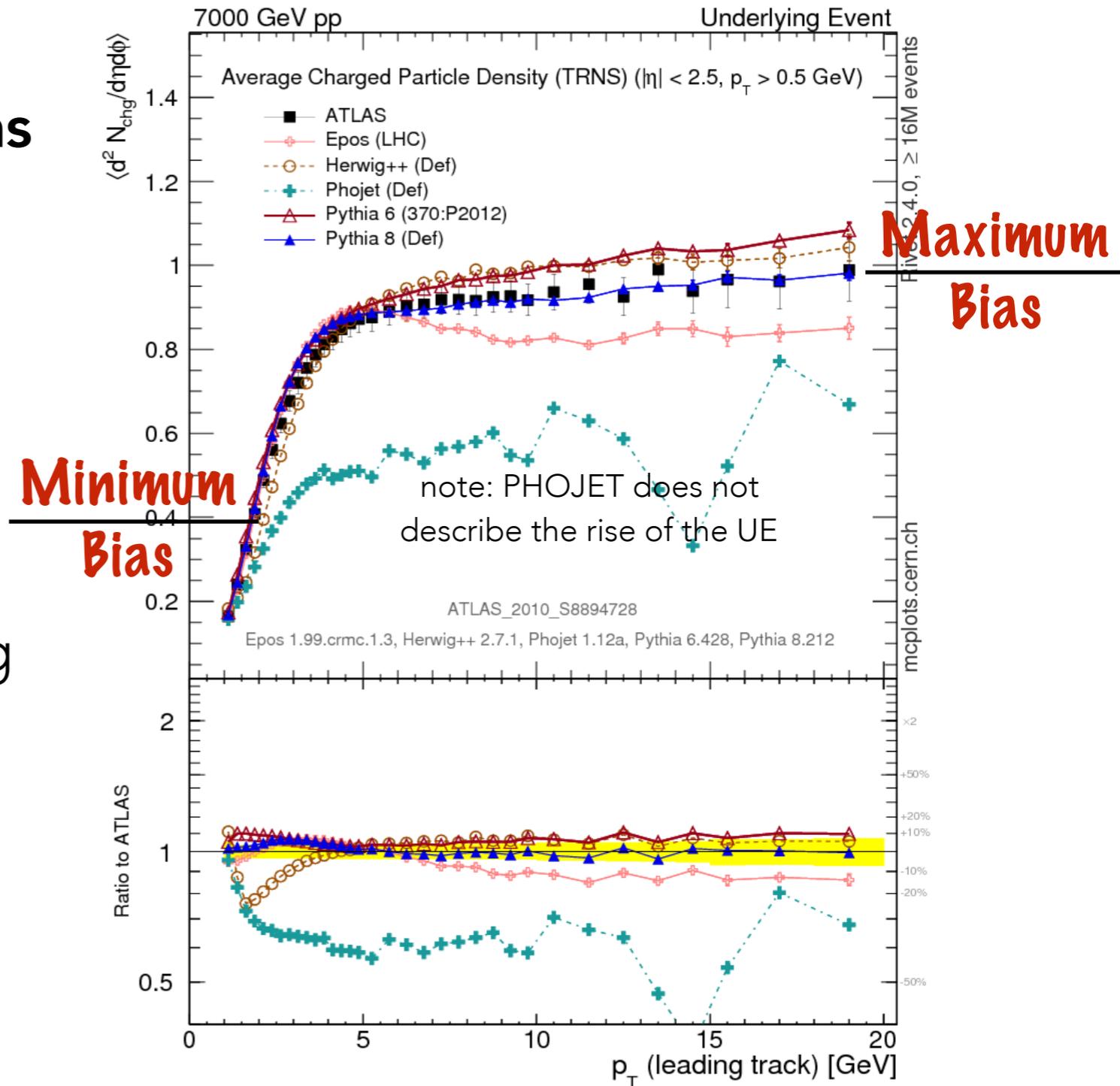
Tautology:

A jet trigger provides a bias
(→ subsample of minimum-bias)

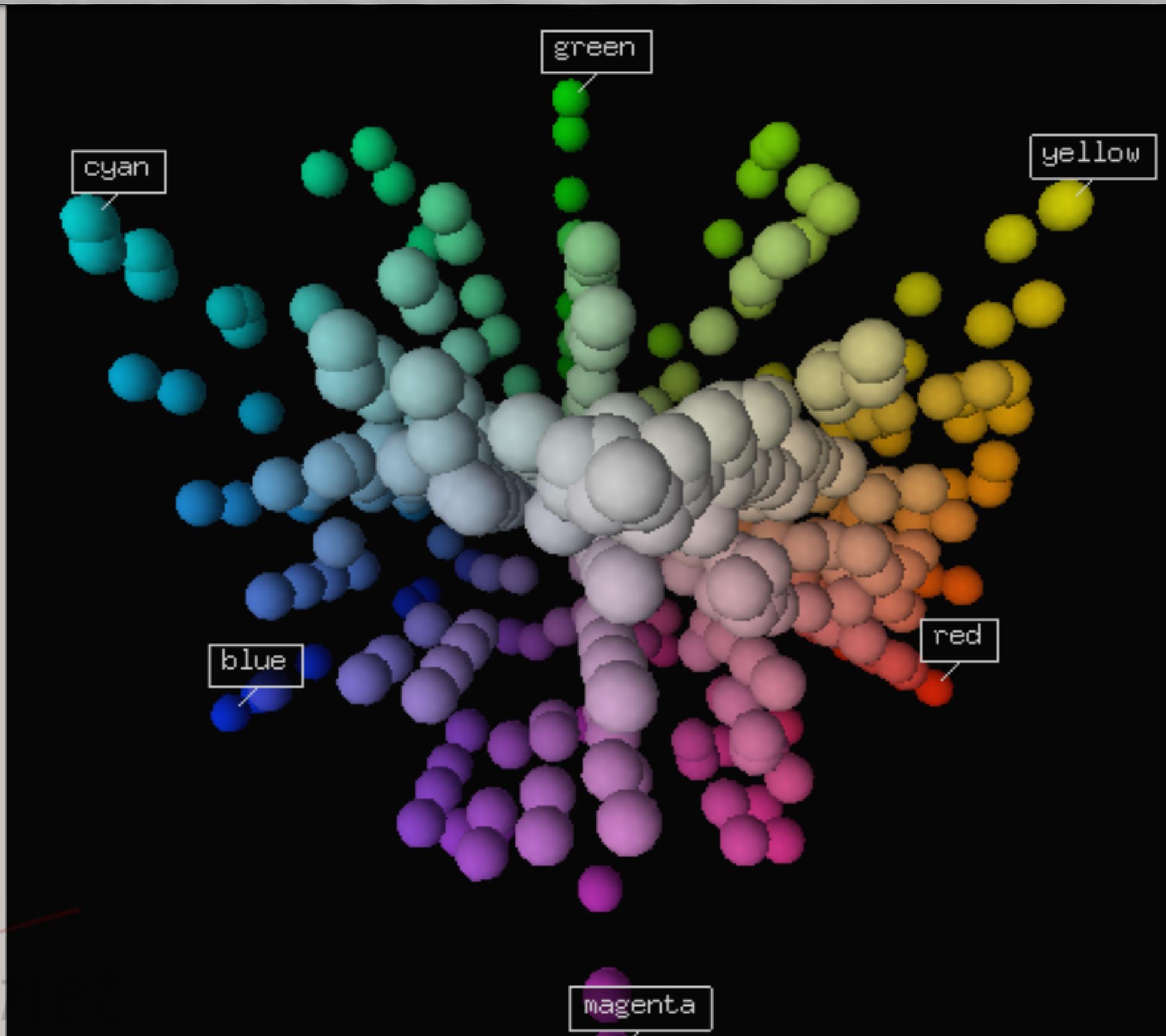
Pedestal effect:

Events with a hard jet trigger are accompanied by a higher plateau of ambient activity

MPI: interpreted as a biasing effect. Small pp impact parameters → larger matter overlaps → more MPI → higher chances for a hard interaction



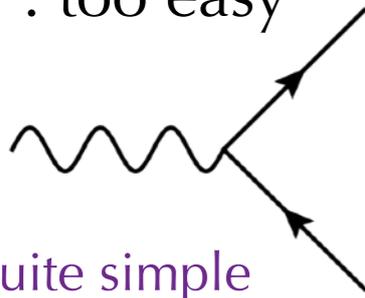
COLOUR SPACE IN HADRON COLLISIONS



COLOUR CONFUSION

Between which partons do confining potentials arise?

e^+e^- : too easy



(still quite simple even after including bremsstrahlung etc.)

At e^+e^- colliders (eg LEP) : generally good agreement between **measured** particle spectra and **models** based on parton/antenna showers + strings

Basically a single **3-3bar** system, very close to the original lattice studies motivating the string model.

(+ extensions to WW reasonable to $\sim O(1/N_c^2)$)

→ re-use same models as input for LHC (universality) ?

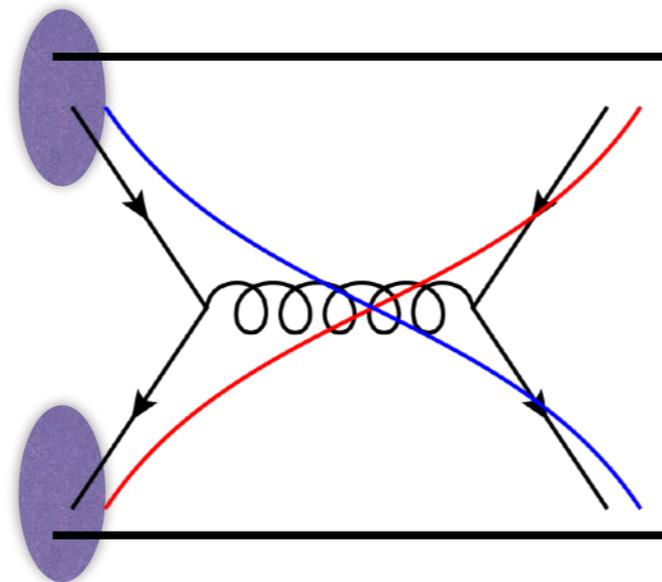
Proton-Proton (LHC)

A lot more colour kicked around (& also colour in initial state)

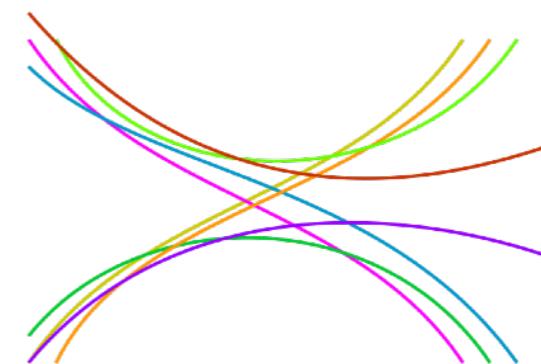
Include "Beam Remnants"

Still might look relatively simple, to begin with

(+baryon beam remnants → "string junctions")



Now add MPI:



Included in all (modern) Monte Carlo models
But how to make sense of the colour structure?

String-fragmentation of junctions: Sjöstrand & Skands **Nucl.Phys. B659 (2003) 243**

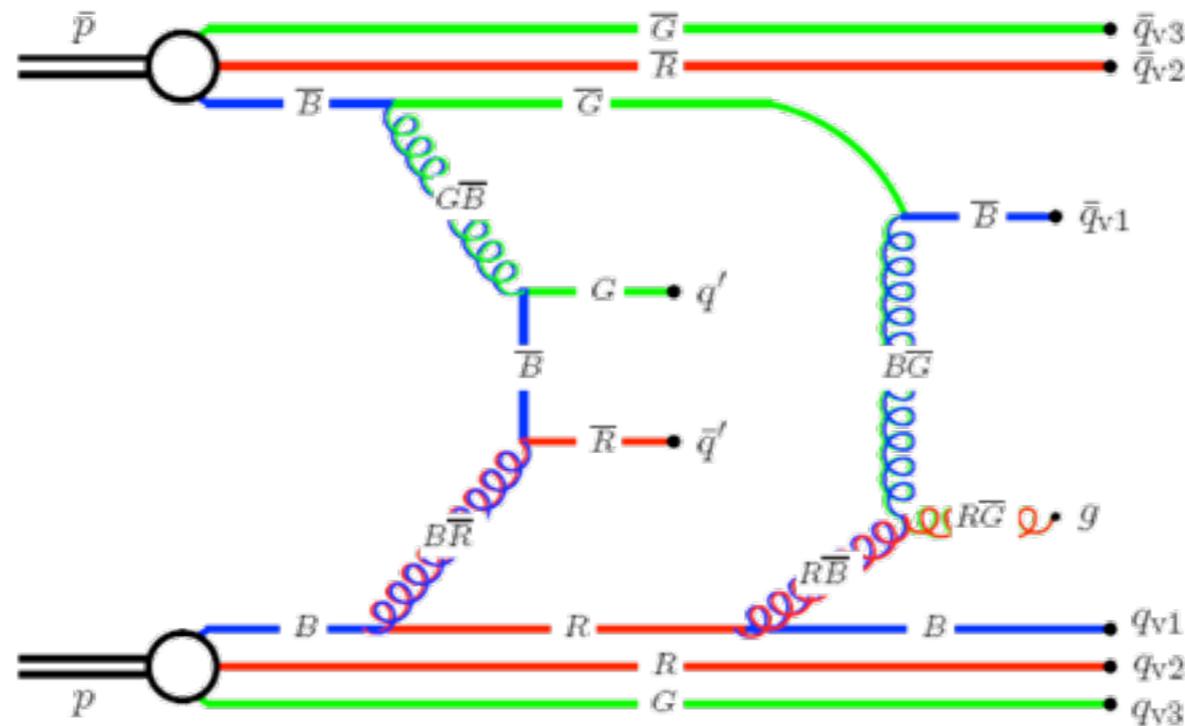
COLOR CORRELATIONS

Each MPI (or cut Pomeron) exchanges color between the beams

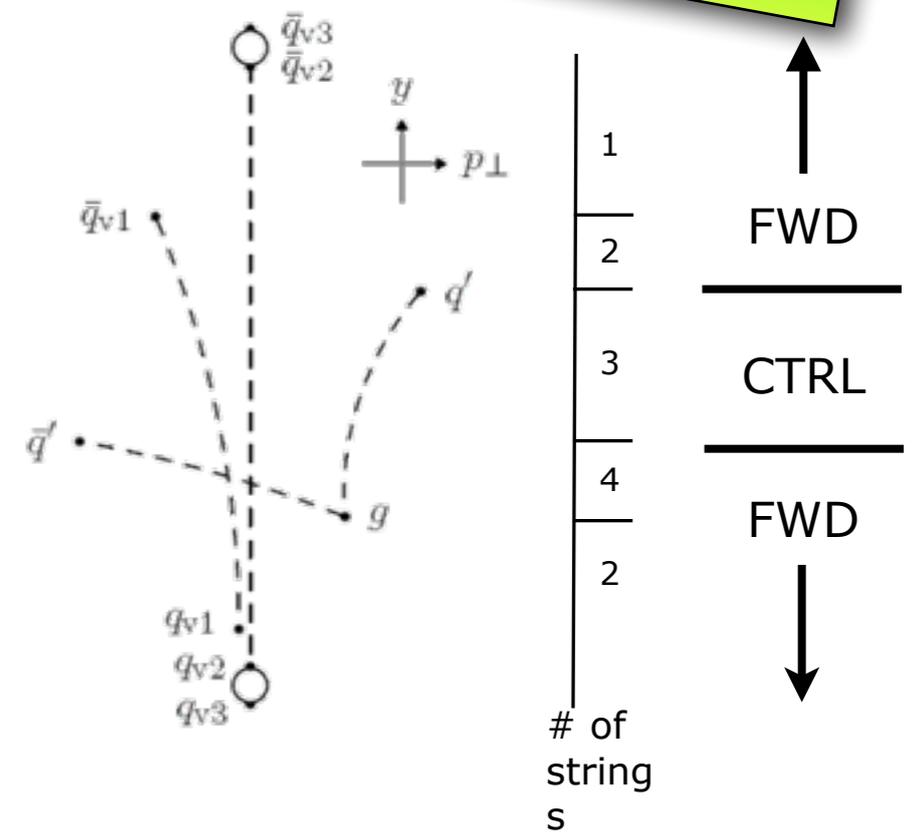
► The colour flow determines the hadronizing string topology

- Each MPI, even when soft, is a color spark
- Final distributions crucially depend on color space

Different models make different ansätze



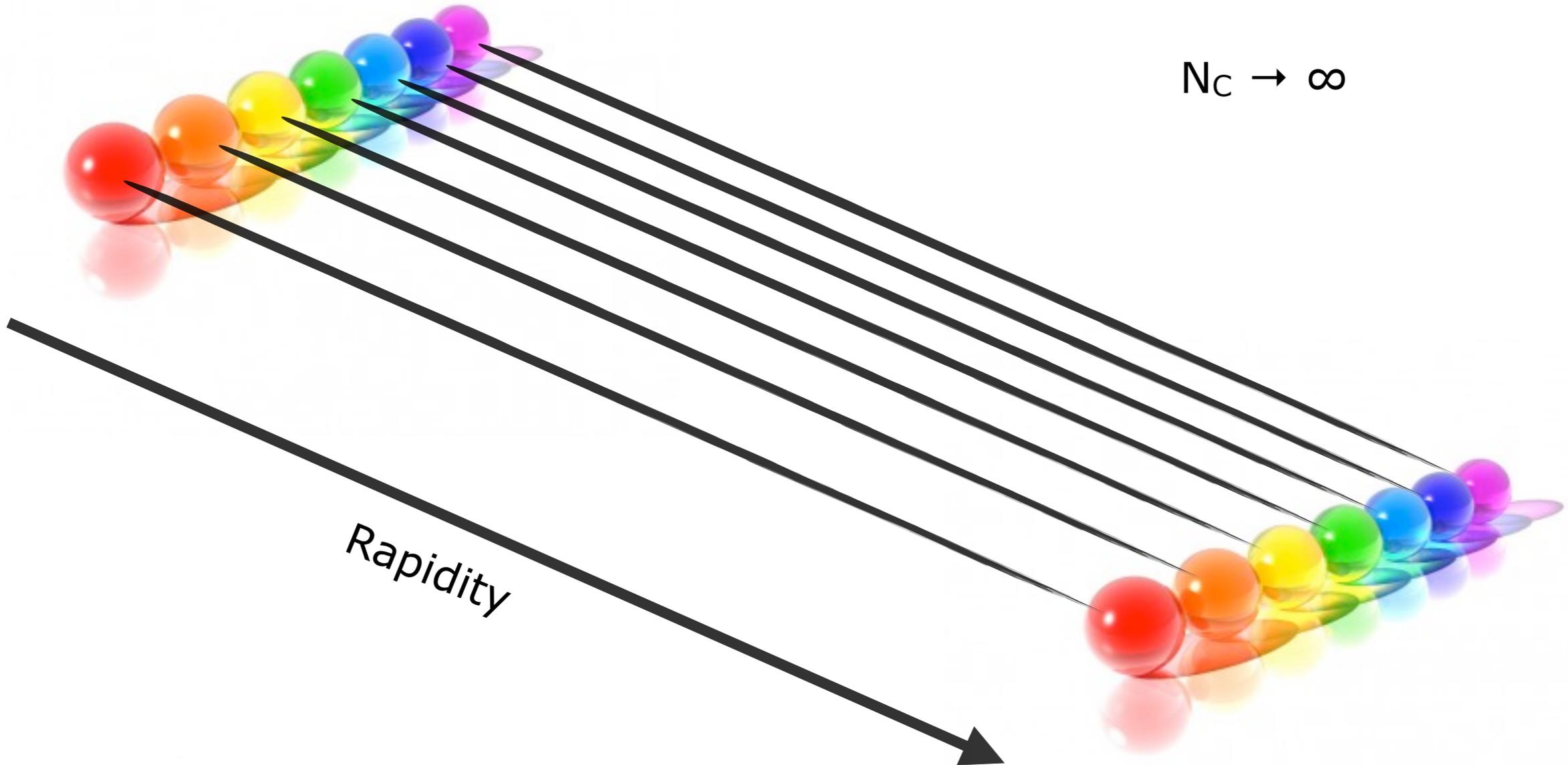
Sjöstrand & PS, JHEP 03(2004)053



COLOR CONNECTIONS

Better theory models needed

$$N_c \rightarrow \infty$$



Rapidity

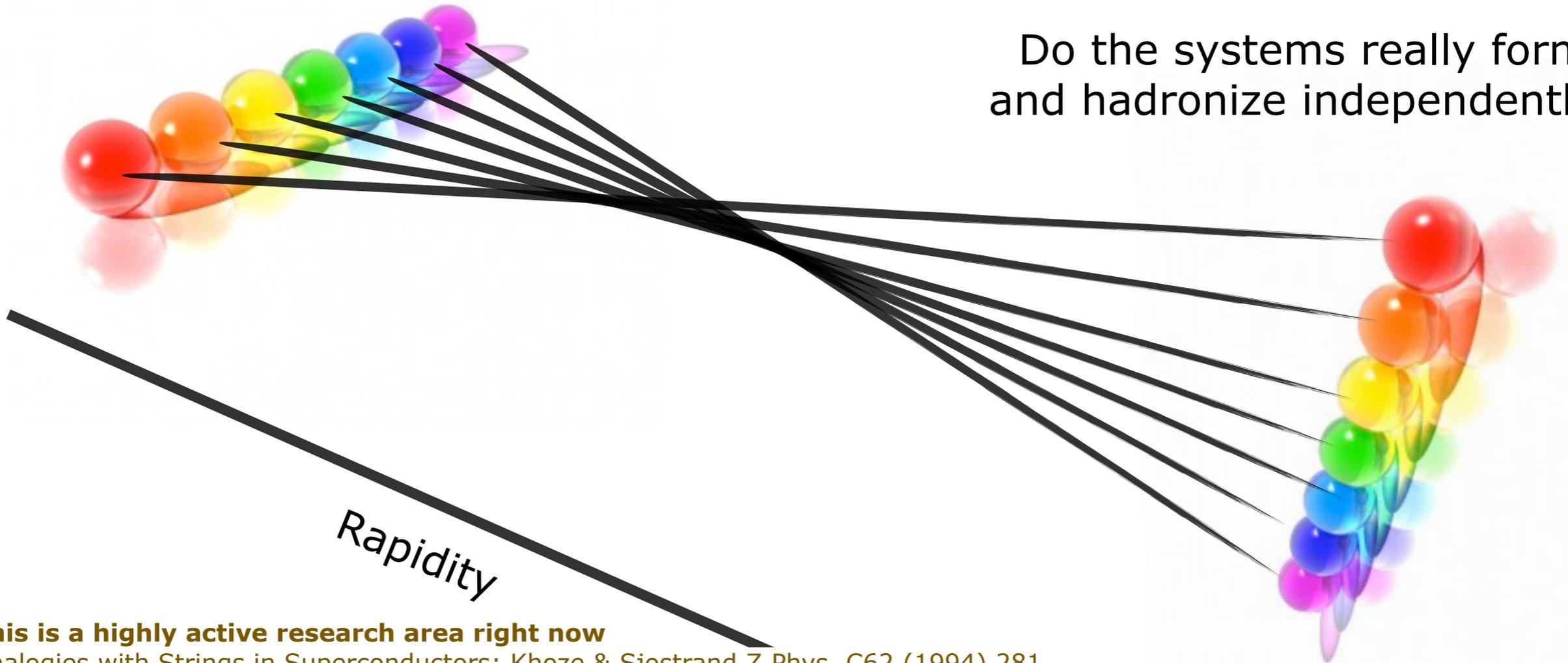
$$\text{Multiplicity} \propto N_{\text{MPI}}$$



COLOR RECONNECTIONS?

Better theory models needed

Do the systems really form and hadronize independently?



This is a highly active research area right now

Analogies with Strings in Superconductors: Khoze & Sjostrand Z.Phys. C62 (1994) 281

Generalized Area Law: Rathsman: Phys. Lett. B452 (1999) 364

Colour Annealing: Skands & Wicke: Eur. Phys. J. C52 (2007) 133

Cluster-based models: e.g. Gieseke et al., Eur.Phys.J. C72 (2012) 2225

Dipole Swing, Lonnblad et al.

Gluon Move Model, Sjostrand et al.

Colour Ropes: Bierlich et al, JHEP 1503 (2015) 148

String Formation Beyond Leading Colour: Christensen & Skands: arXiv:1505.01681

String interactions? Hydrodynamics (EPOS)? Collective flow? Pressure? Rescatterings?

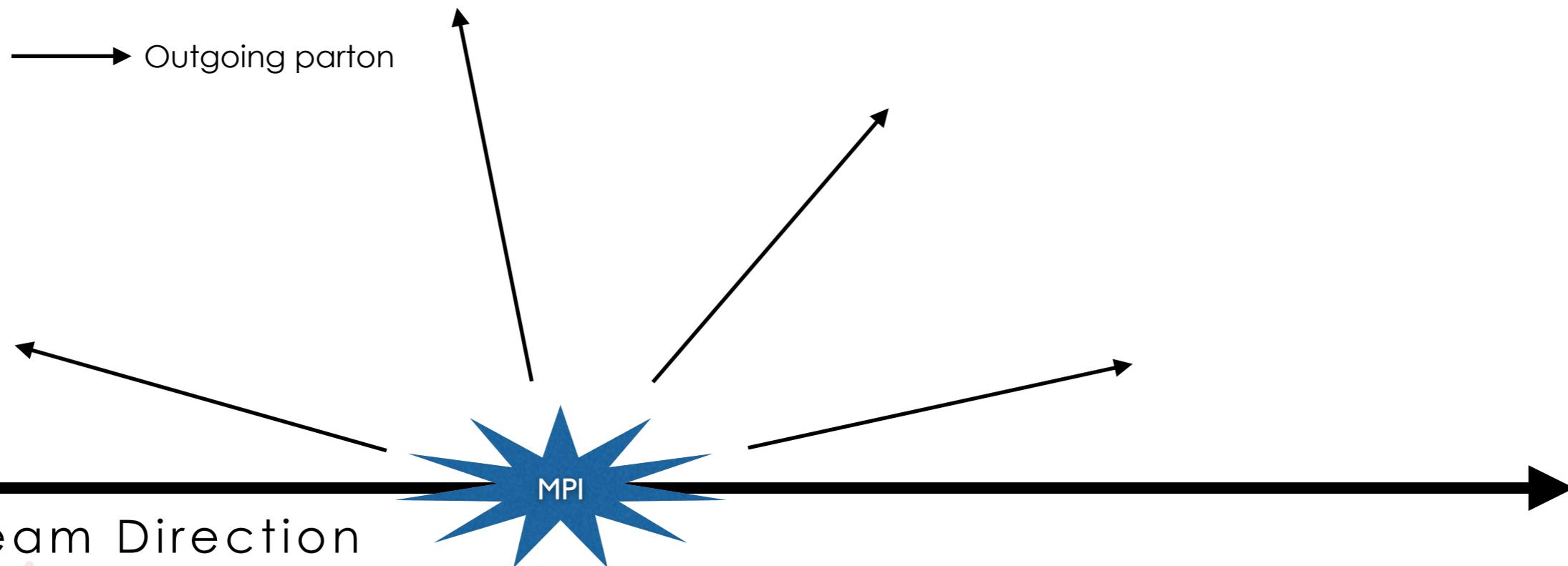
Multiplicity \propto N_{MPI}

COLOUR: WHAT'S THE PROBLEM?

(including **MPI**: Multiple Parton-Parton Interactions ~ the “underlying event”)

Without Colour Reconnections

Each MPI hadronizes **independently** of all others



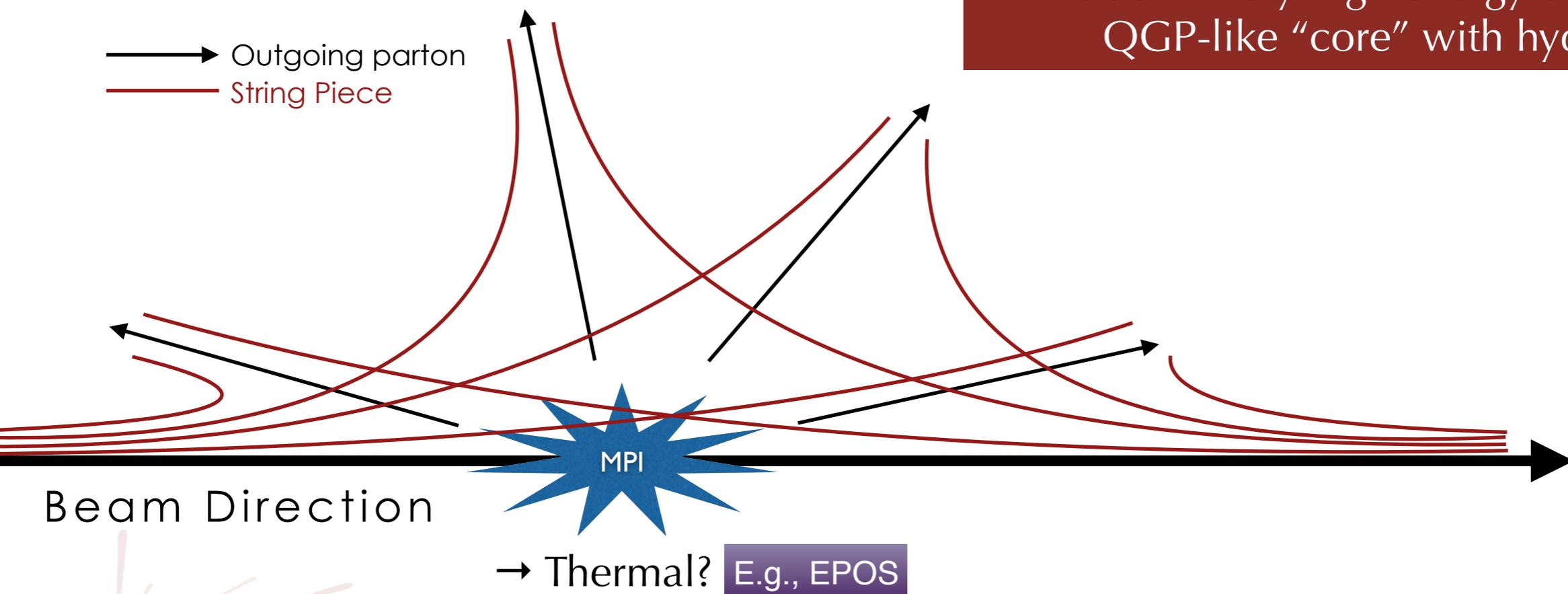
COLOUR: WHAT'S THE PROBLEM?

(including **MPI**: Multiple Parton-Parton Interactions ~ the “underlying event”)

Without Colour Reconnections

Each MPI hadronizes **independently** of all others

So many strings in so little space
If true → Very high energy densities
QGP-like “core” with hydro?



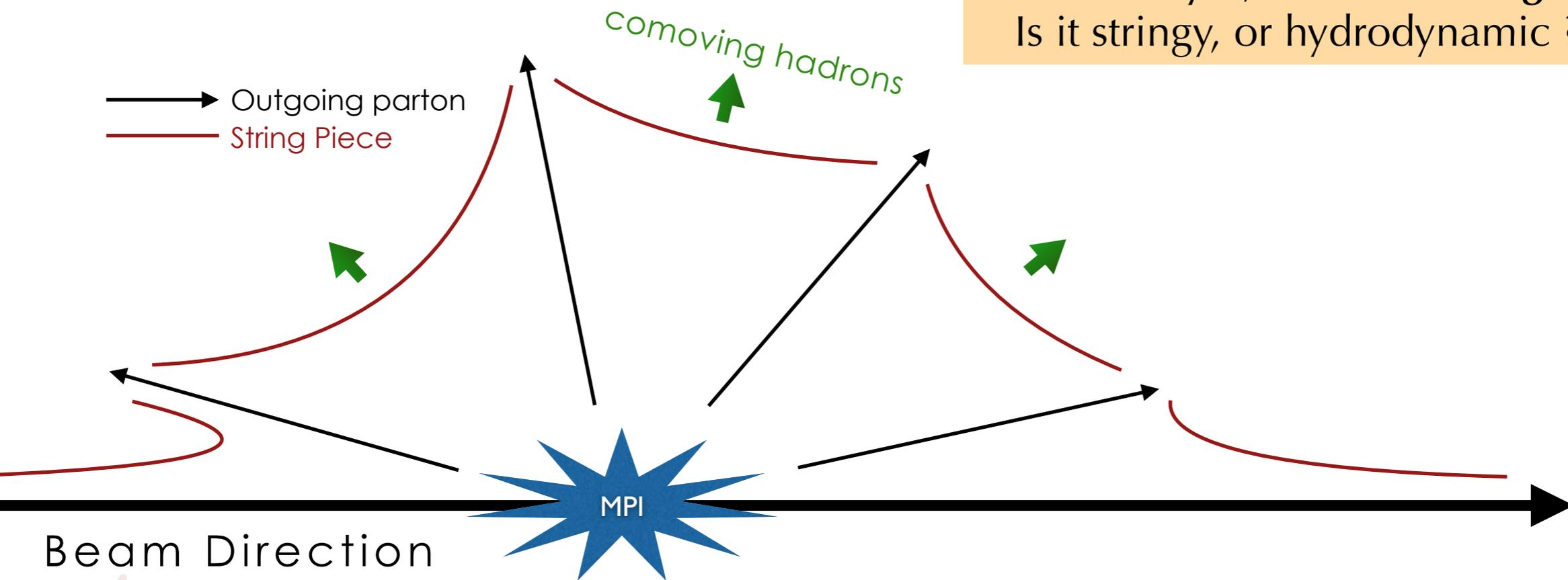
COLOUR RECONNECTIONS

(including **MPI**: Multiple Parton-Parton Interactions ~ the “underlying event”)

With Colour Reconnections
MPI hadronize **collectively**

See also Ortiz et al., Phys.Rev.Lett. 111 (2013) 4, 042001

Highly interesting theory questions now.
Is there collective flow in pp? Or not?
If yes, what is its origin?
Is it stringy, or hydrodynamic ? (or ...?)



String-Length Minimisation E.g., PYTHIA, HERWIG

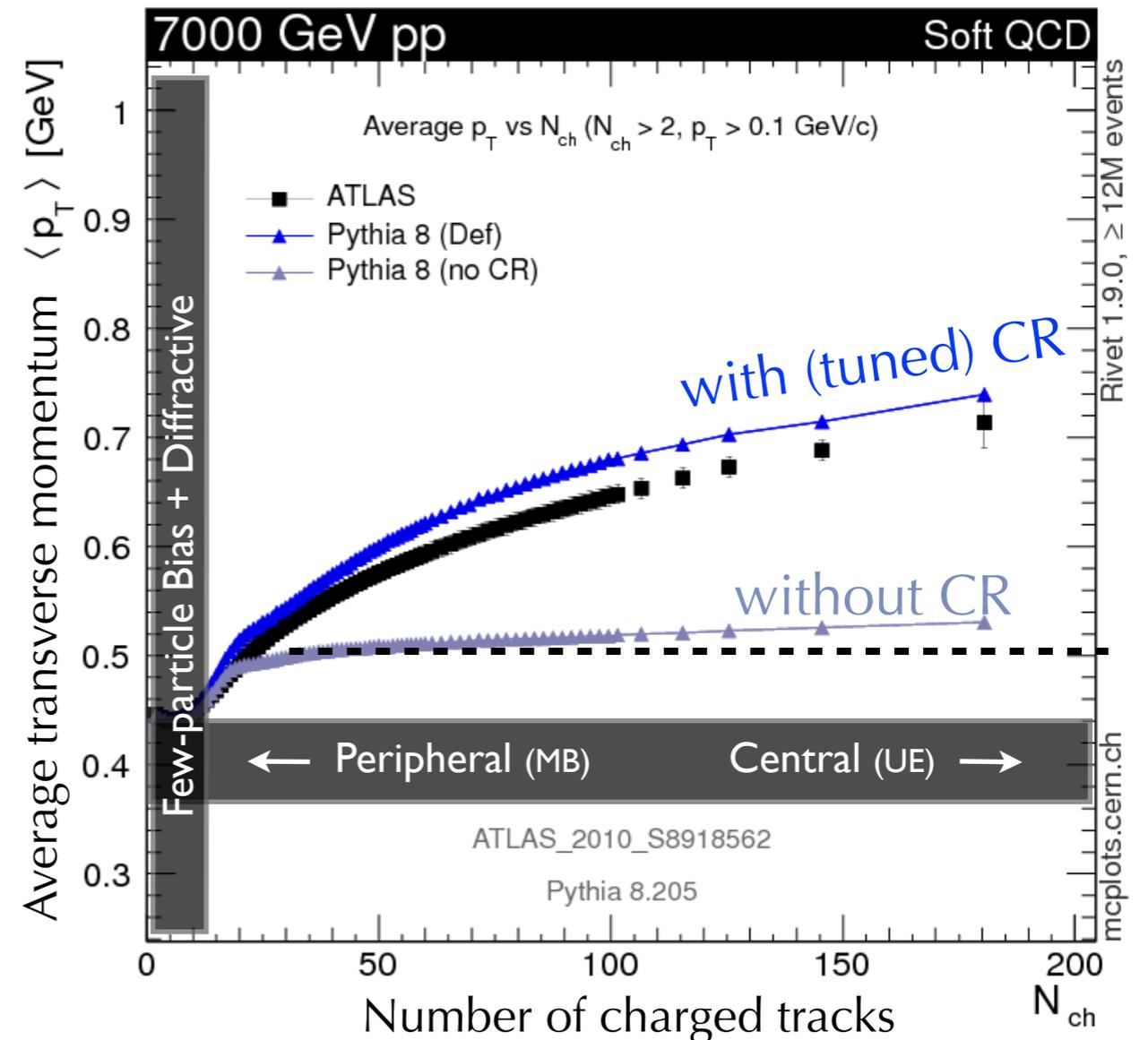
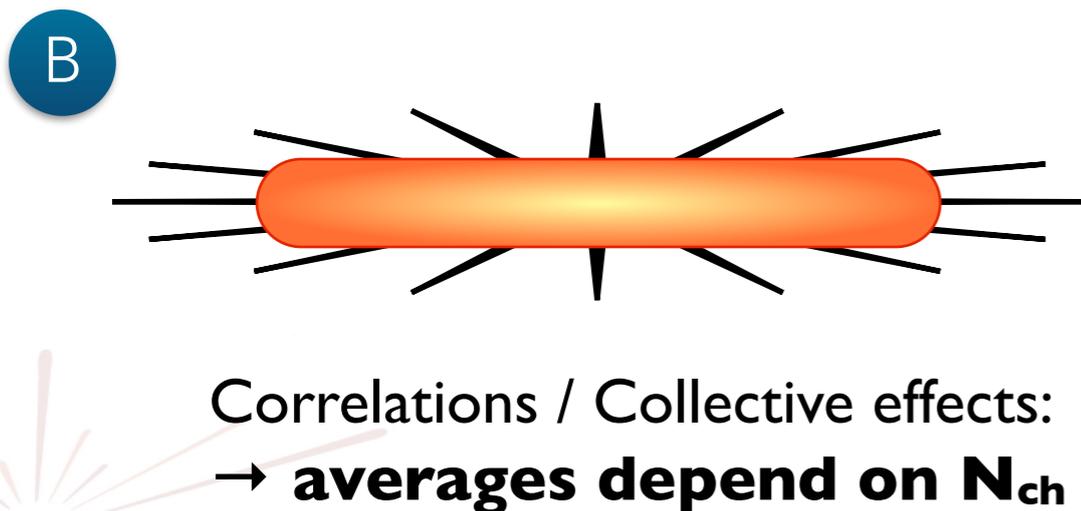
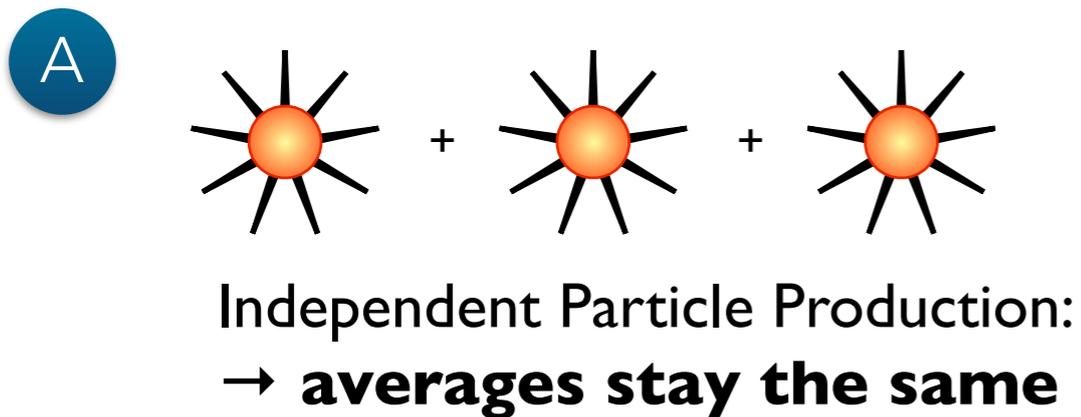
Or Thermal? E.g., EPOS

Or Higher String Tension? E.g., DIPSY rope

COLLECTIVE EFFECTS?

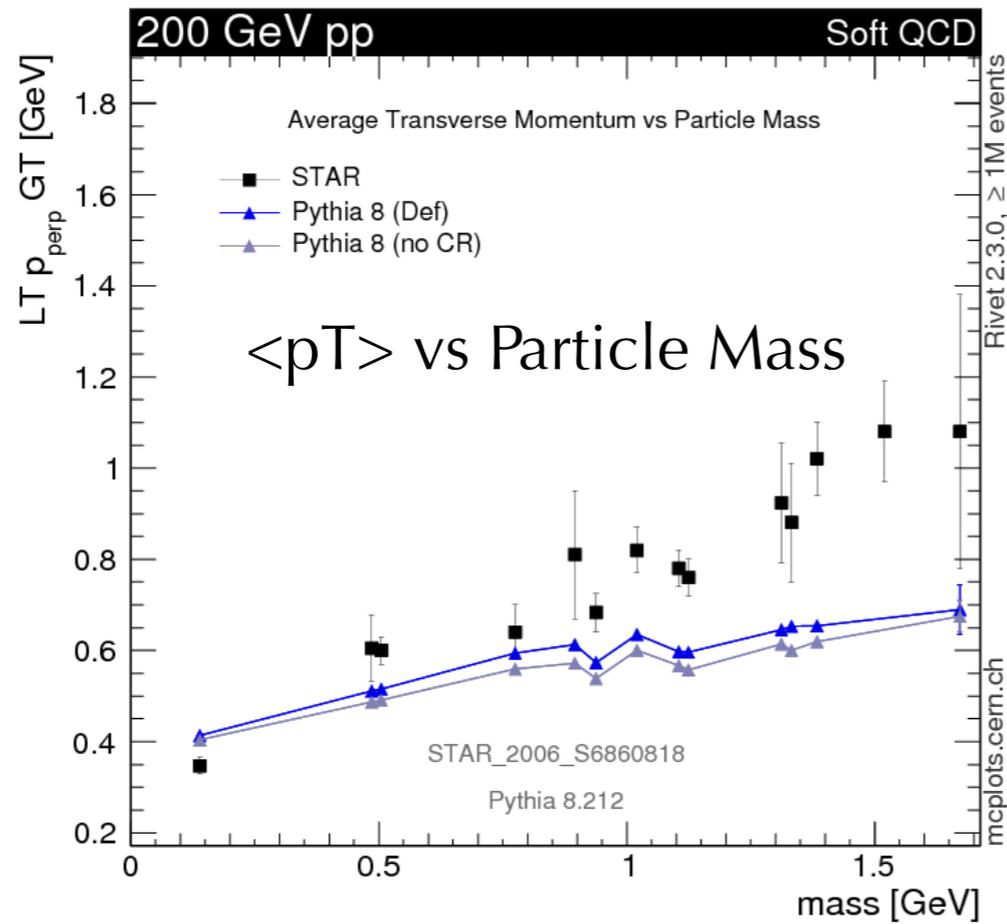
A rough indicator of how much colour gets kicked around, should be the number of particles produced

So we study event properties as a function of " N_{ch} " = N_{tracks}

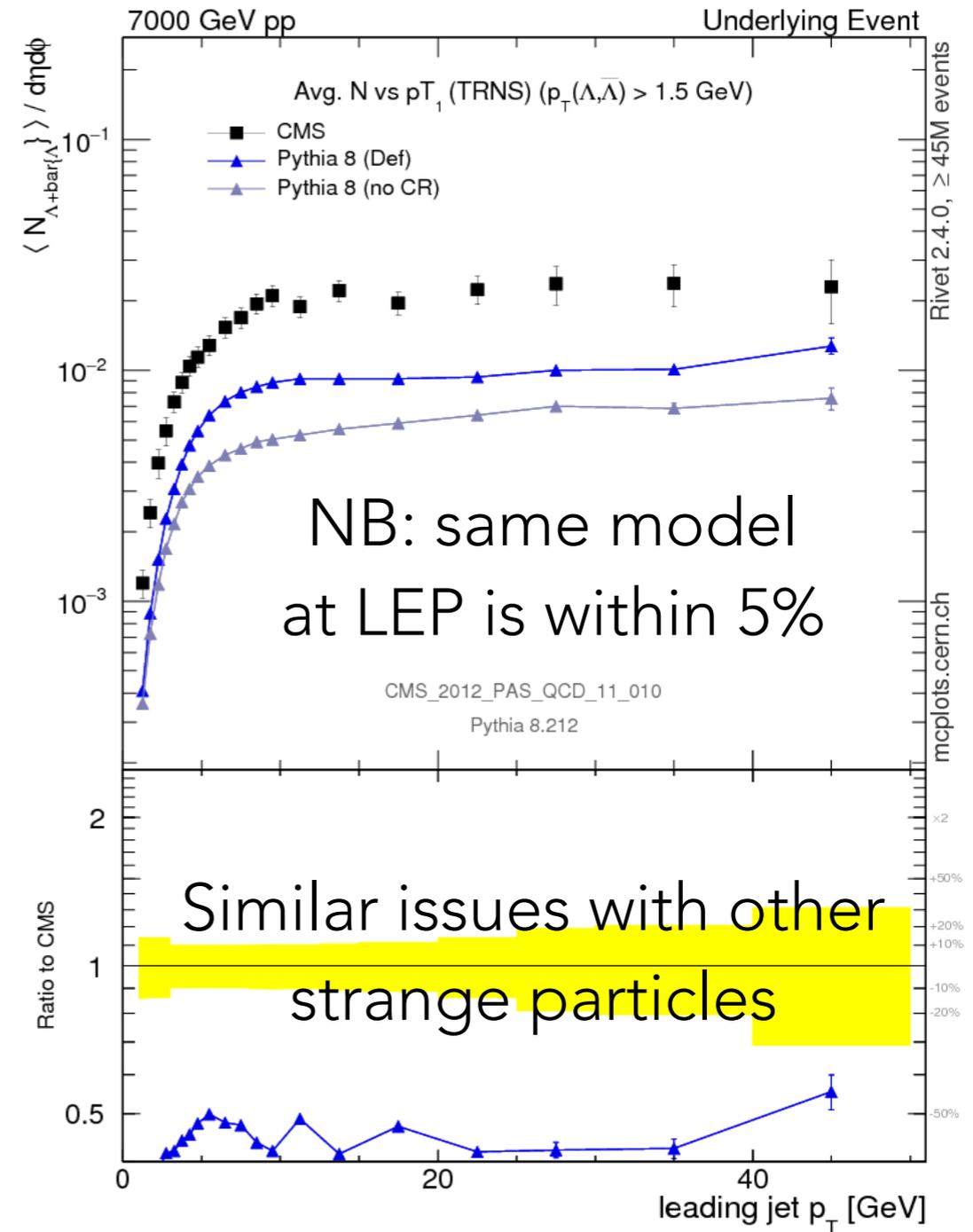


OTHER INDICATIONS in pp

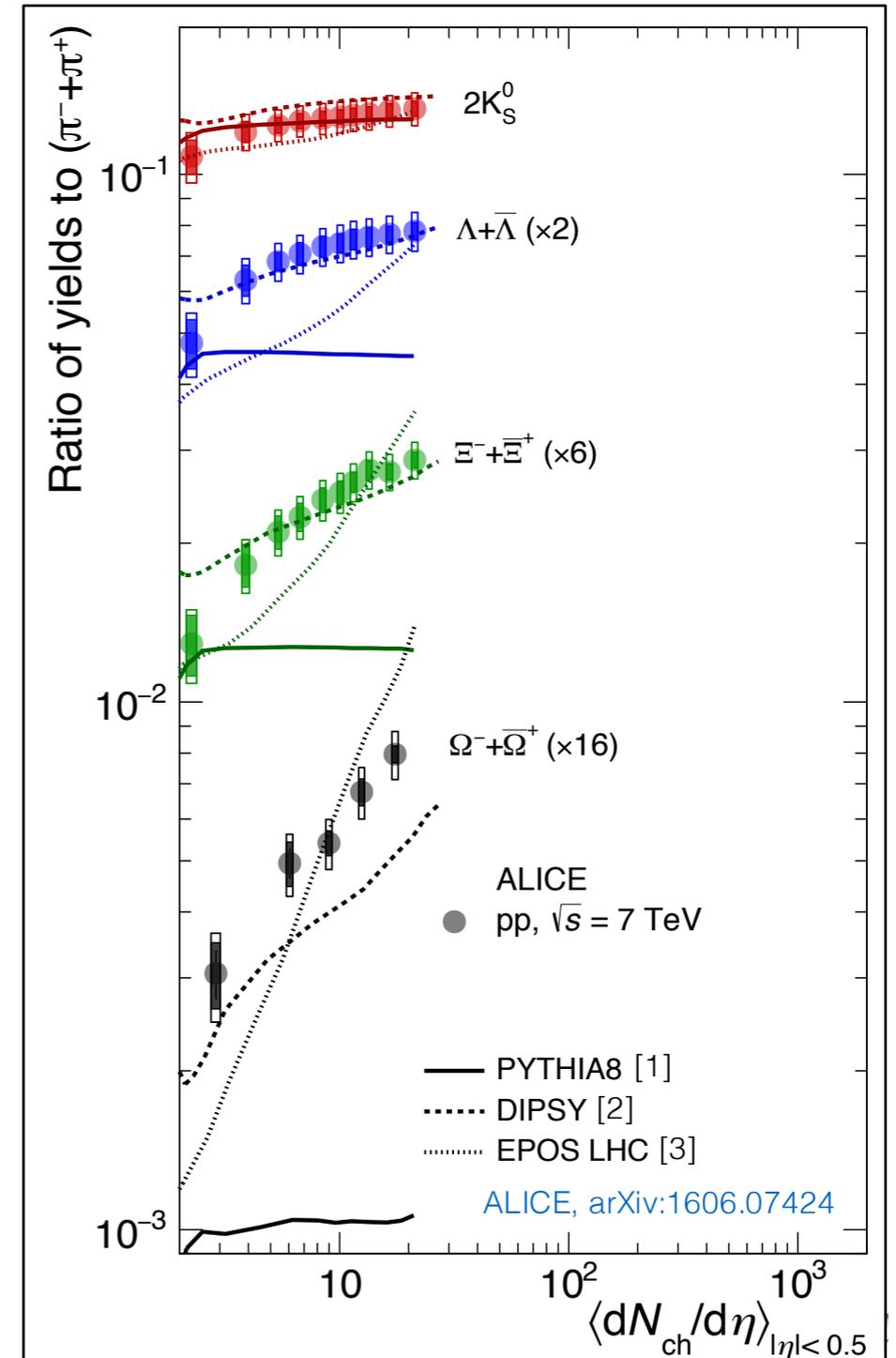
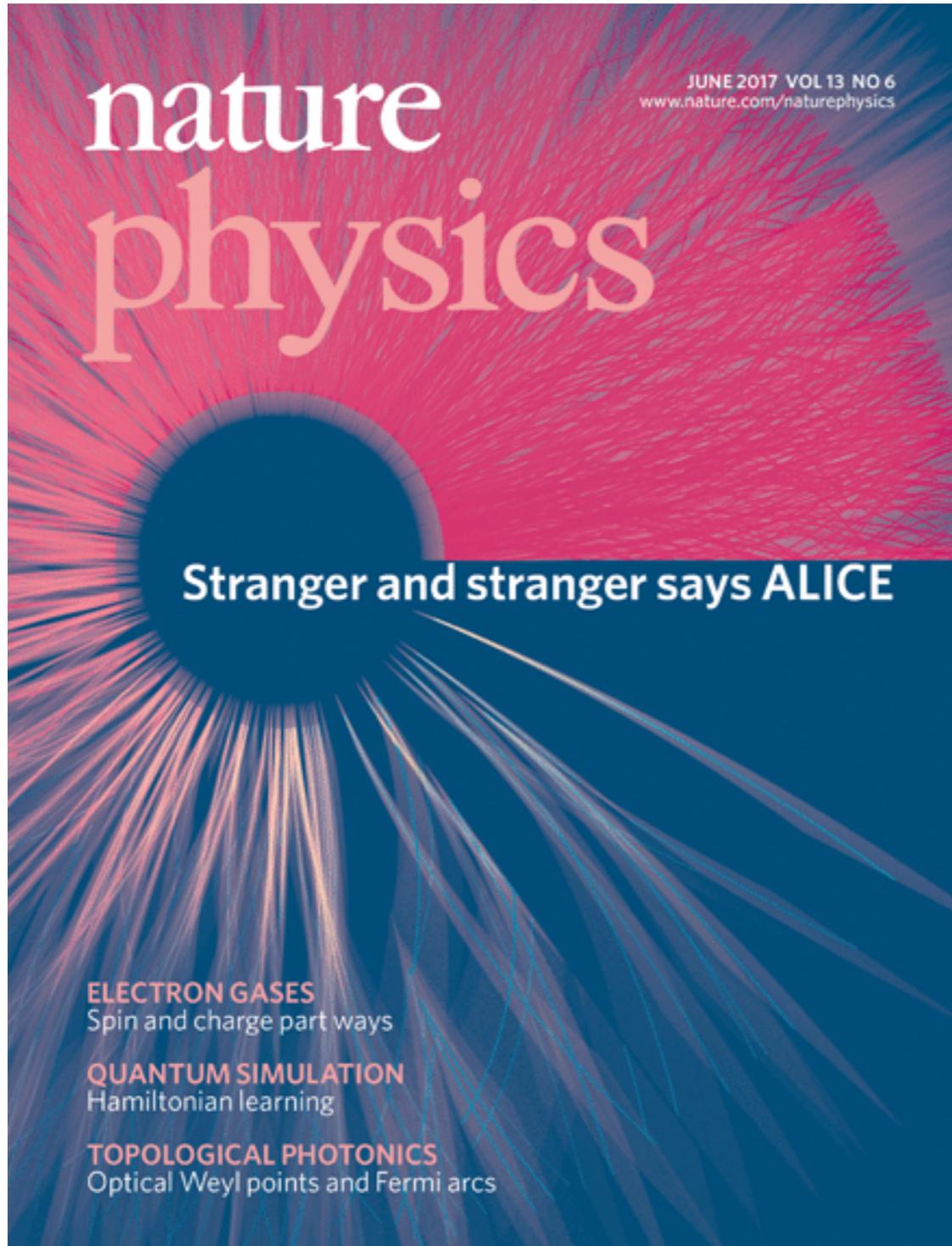
Heavier particles are harder



Where have all the Λ gone?



... and then there was this ...



SUMMARY: MCS & PARTON SHOWERS

Aim: generate events in as much detail as mother nature

→ Make stochastic choices ~ as in Nature (Q.M.) → Random numbers

Factor complete event probability into separate universal pieces, treated independently and/or sequentially (Markov-Chain MC)

Improve Born-level theory with 'most significant' corrections

Resonance decays (e.g., $t \rightarrow bW^+$, $W \rightarrow qq'$, $H^0 \rightarrow \gamma^0 \gamma^0$, $Z^0 \rightarrow \mu^+ \mu^-$, ...)

Bremsstrahlung (FSR and ISR, exact in collinear and soft* limits)

Hard radiation (matching)

Hadronization (strings/clusters, discussed tomorrow)

Additional Soft Physics: multiple parton-parton interactions, Bose-Einstein correlations, colour reconnections, hadron decays, ...



FINAL WORDS

MCs can be treated as black boxes, without knowing what's in them.

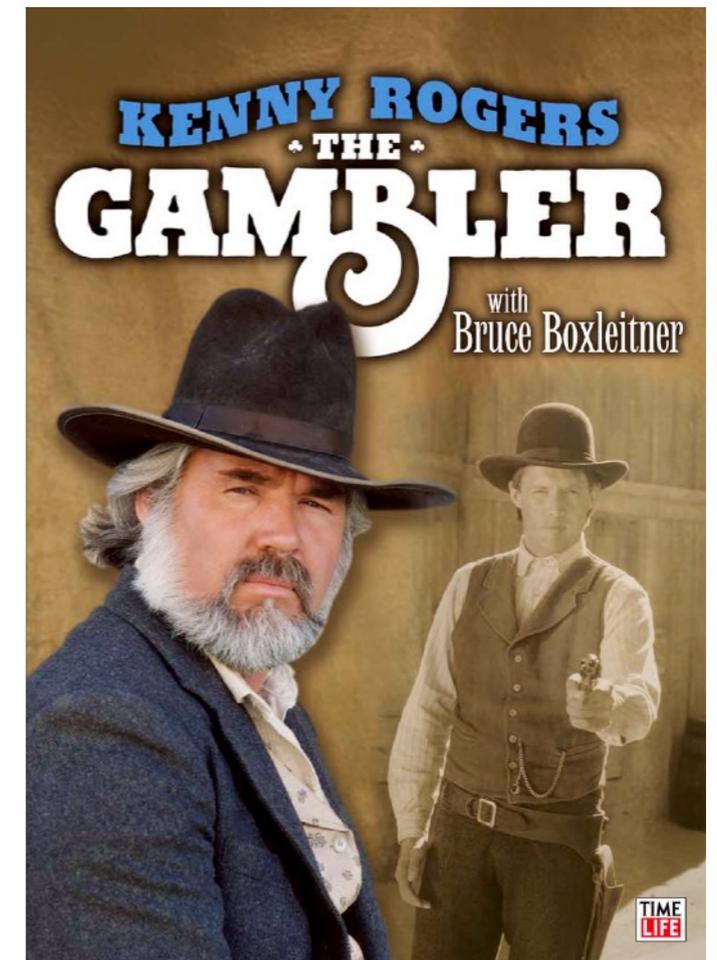
Best Case: Limited Sophistication
Worst Case: Not your lucky day



The secret to successful MC is:

Knowing what to throw away

Knowing what to keep



Kenny Rogers "The Gambler", first recorded in 1978
Same year as the first version of PYTHIA (JETGEN)



Extra Slides

(SOME CAVEATS OF MPI-BASED MODELS)

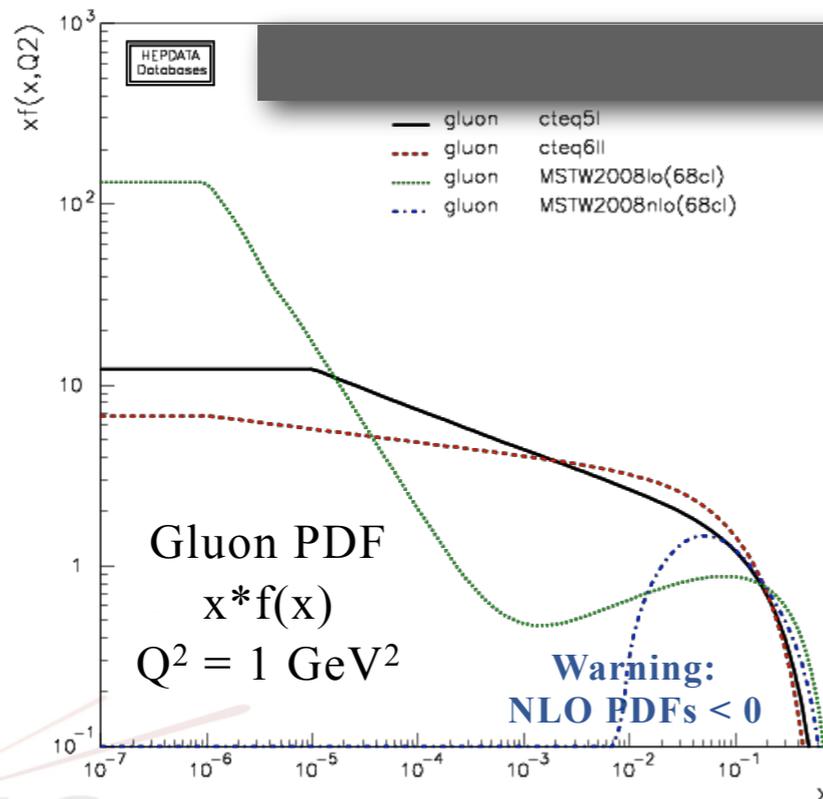
Main applications of factorisation: Central Jets/EWK/top/Higgs/New Physics

$$d\sigma_{2\rightarrow 2} \propto \frac{dp_{\perp}^2}{p_{\perp}^4} \otimes \text{PDFs}$$

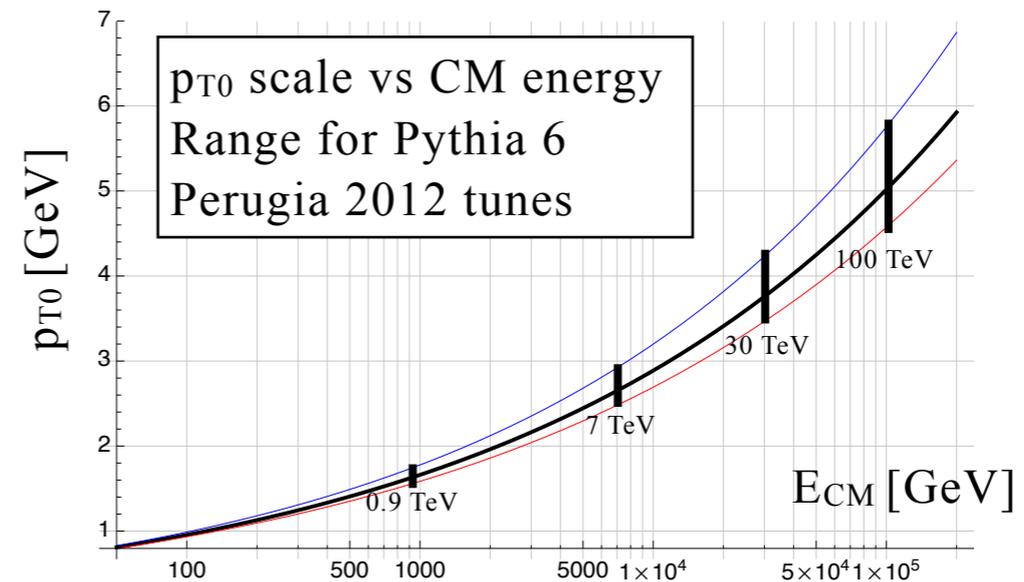
High Q^2 and finite x

Extrapolation to soft scales delicate.
Impressive successes with MPI-based models but still far from ‘problem solved’

- Form of PDFs at small x and Q^2 “Saturation” ?
- Form and E_{cm} dependence of p_{T0} regulator
- Modeling of the diffractive component
- Proton transverse mass distribution
- Colour Reconnections, Collective Effects



Poor Man's Saturation



See also [Connecting hard to soft: KMR, EPJ C71 \(2011\) 1617](#) + [PYTHIA “Perugia Tunes”: PS, PRD82 \(2010\) 074018](#) + [arXiv:1308.2813](#)

THE INELASTIC CROSS SECTION

First try: decompose $\sigma_{inel} = \sigma_{sd} + \sigma_{dd} + \sigma_{cd} + \sigma_{nd}$

+ Parametrizations of diffractive components: dM^2/M^2

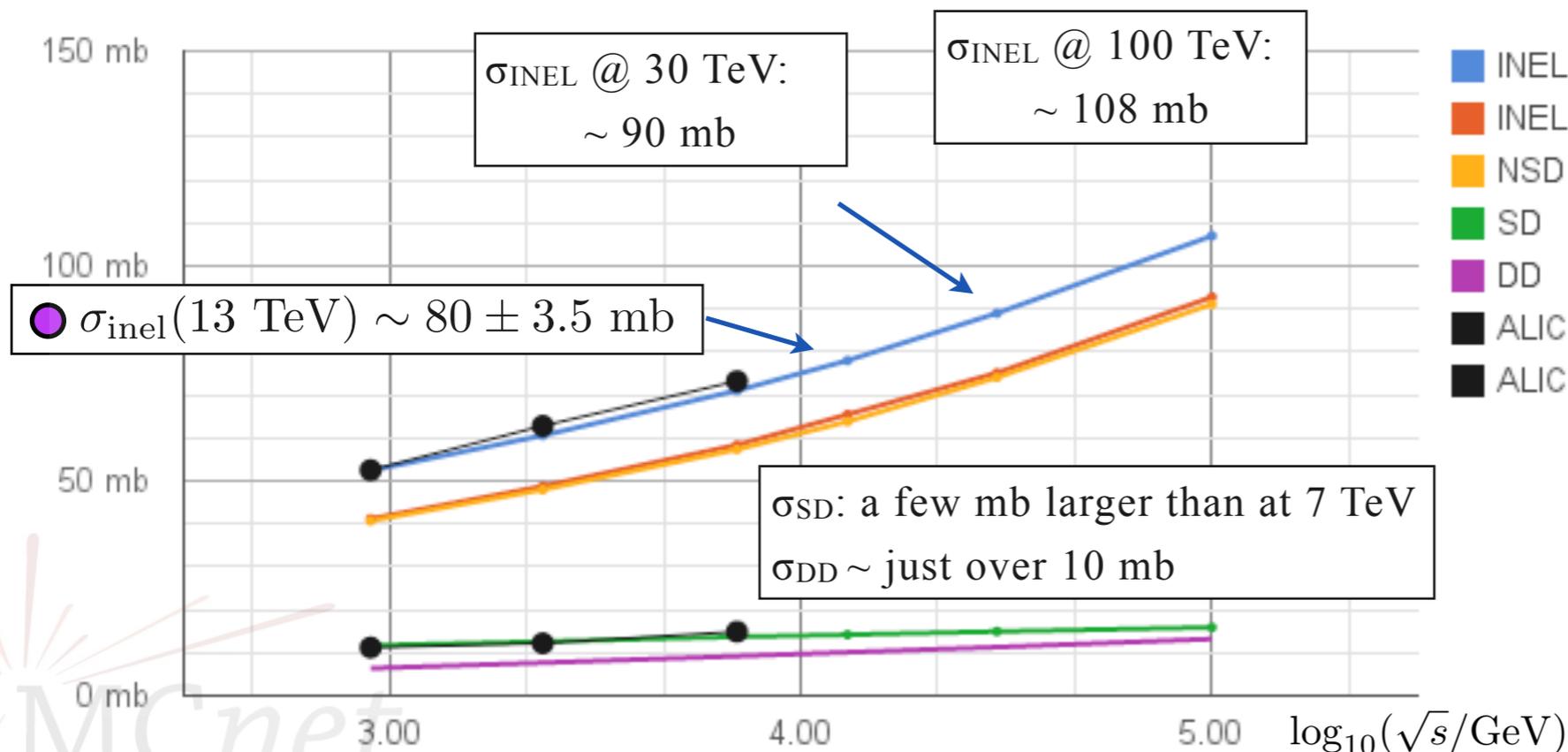
PYTHIA:

$$\frac{d\sigma_{sd}(AX)(s)}{dt dM^2} = \frac{g_{3IP}}{16\pi} \beta_{AIP}^2 \beta_{BIP} \frac{1}{M^2} \exp(B_{sd}(AX)t) F_{sd} ,$$

$$\frac{d\sigma_{dd}(s)}{dt dM_1^2 dM_2^2} = \frac{g_{3IP}^2}{16\pi} \beta_{AIP} \beta_{BIP} \frac{1}{M_1^2} \frac{1}{M_2^2} \exp(B_{dd}t) F_{dd} .$$

+ Integrate and solve for σ_{nd}

What Cross Section?



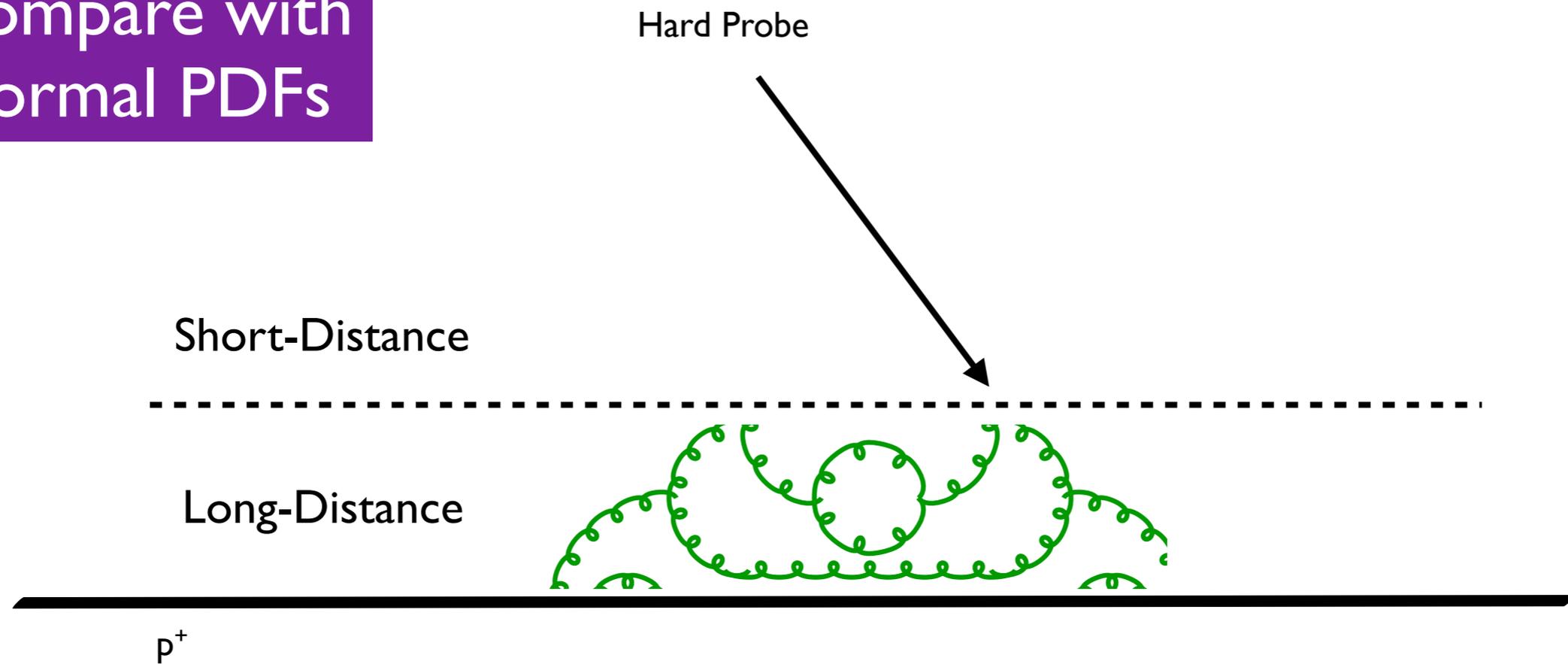
- Total Inelastic
- Fraction with one charged particle in $|\eta| < 1$
- Ambiguous Theory Definition
- Ambiguous Theory Definition
- Ambiguous Theory Definition
- Observed fraction corrected to total
- ALICE def : SD has $MX < 200$

Note problem of principle: Q.M. requires distinguishable final states

(+ DIFFRACTION)

“Intuitive picture”

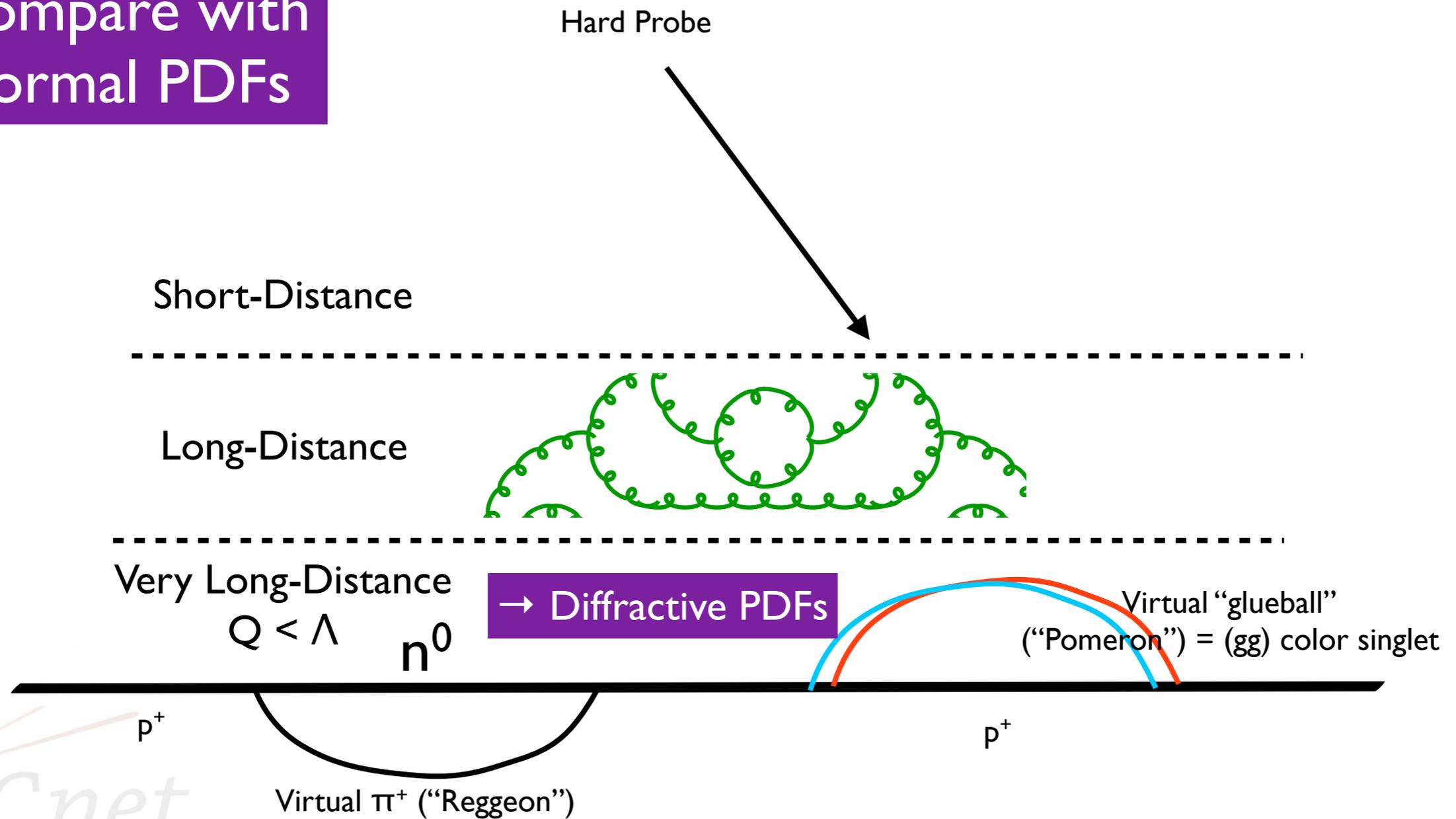
Compare with
normal PDFs



(+ DIFFRACTION)

“Intuitive picture”

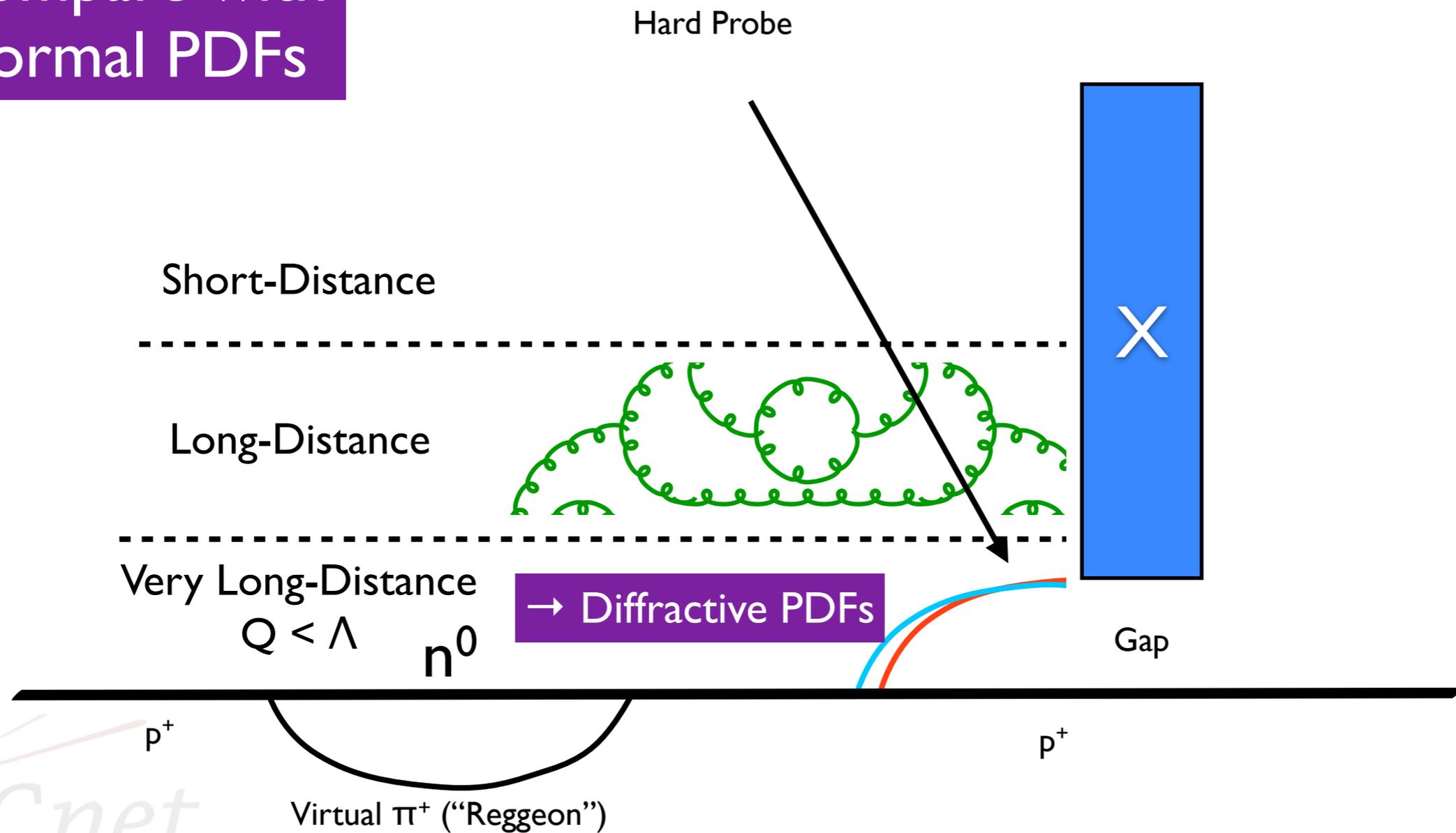
Compare with normal PDFs



(+ DIFFRACTION)

“Intuitive picture”

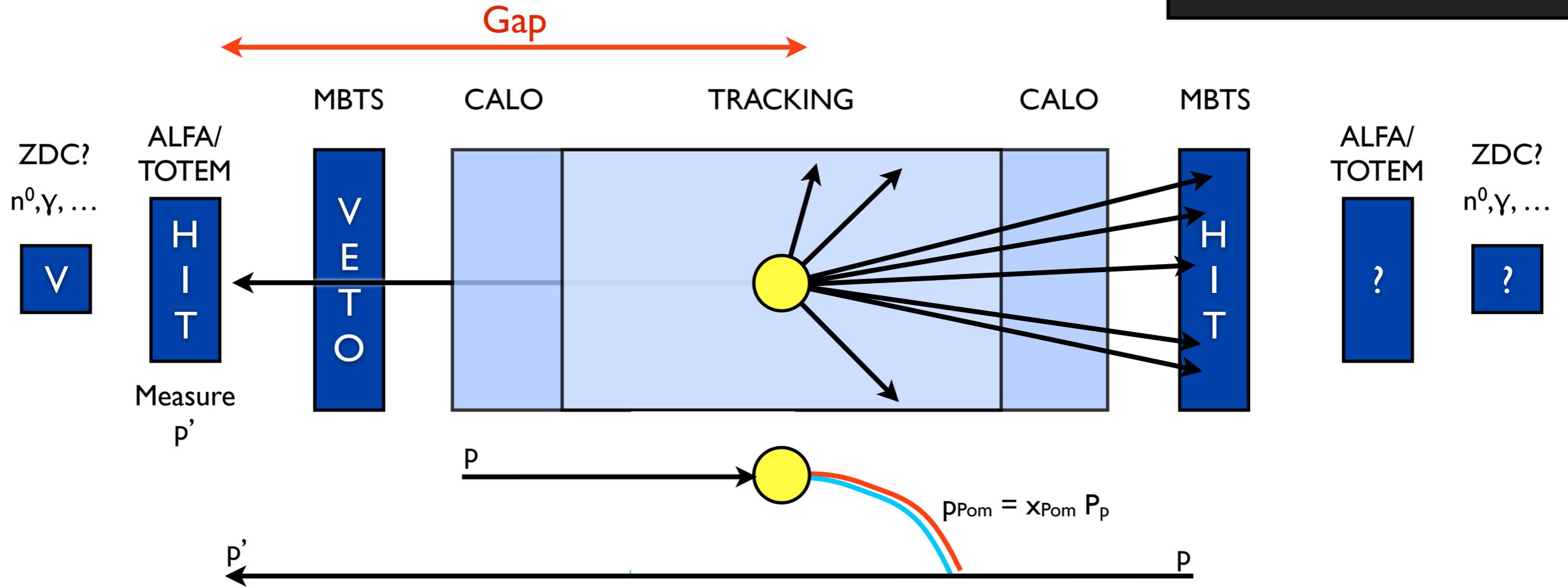
Compare with normal PDFs



WHAT IS DIFFRACTION?

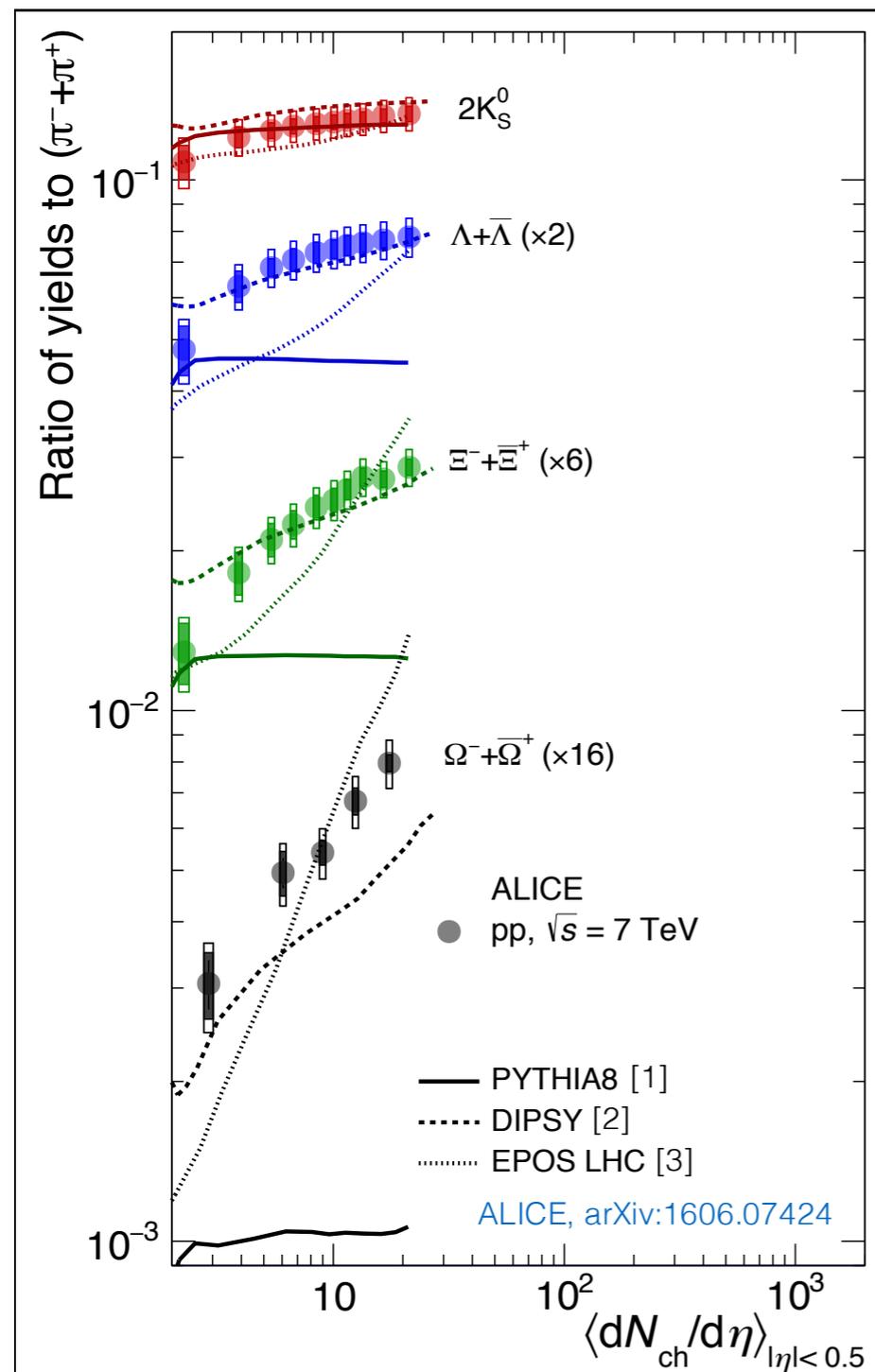
Single Diffraction

Glueball-Proton Collider
with variable E_{CM}



Double Diffraction: both protons explode; gap inbetween
Central Diffraction: two protons + a central (exclusive) system

Recent news from ALICE (ICHEP 2016)



D.D. Chinellato – 38th International Conference on High Energy Physics

A clear enhancement of strangeness with (pp) event multiplicity is observed

Especially for multi-strange baryons

No corresponding enhancement for protons (not shown here but is in ALICE paper)

→ this really must be a strangeness effect

Cross-check measurements of the phi meson are now underway

Jet universality: jets at LHC modelled the same as jets at LEP

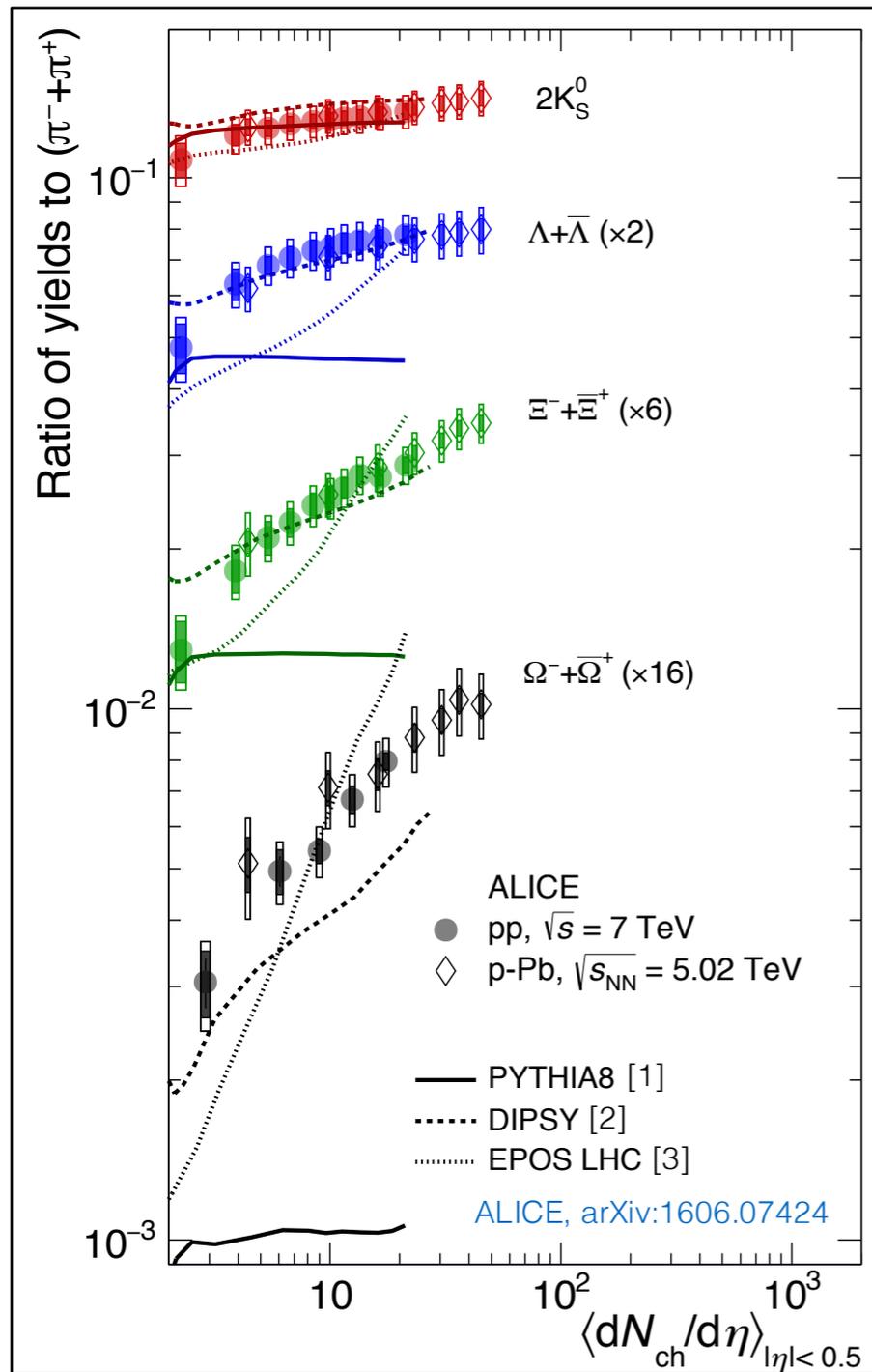
→ Flat line ! (cf PYTHIA)

DIPSY includes “colour ropes” with higher effective string tension

EPOS includes hydrodynamic “core” with higher effective temperature

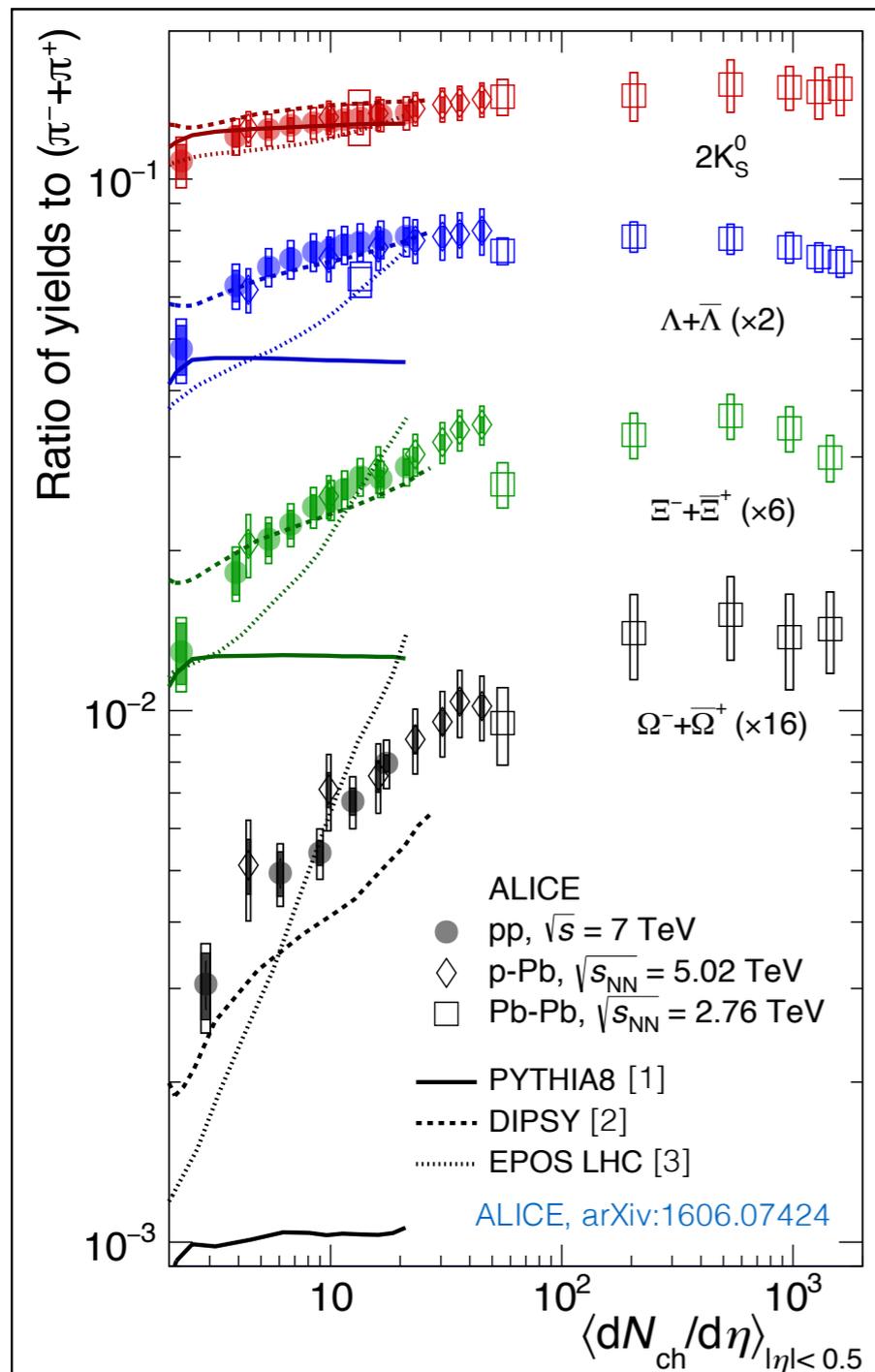
The Plot Thickens

Looks like the effect, whatever it is, continues smoothly into p-Pb



D.D. Chinellato – 38th International Conference on High Energy Physics

The Plot Thickens



D.D. Chinellato – 38th International Conference on High Energy Physics

Looks like the effect, whatever it is, continues smoothly into p-Pb

... and into Pb-Pb !

Unexpected.

Looks like jet universality and hadronisation in pp is up for revision.

Is it thermal? Stringy? Both?

Collective? Flowy? ...

Physics must explain smooth transition to heavy ions. No abrupt “phase transition” seen in these observables