

Extending a minimal DM Model

.. and its consequences for DM Searches at the LHC



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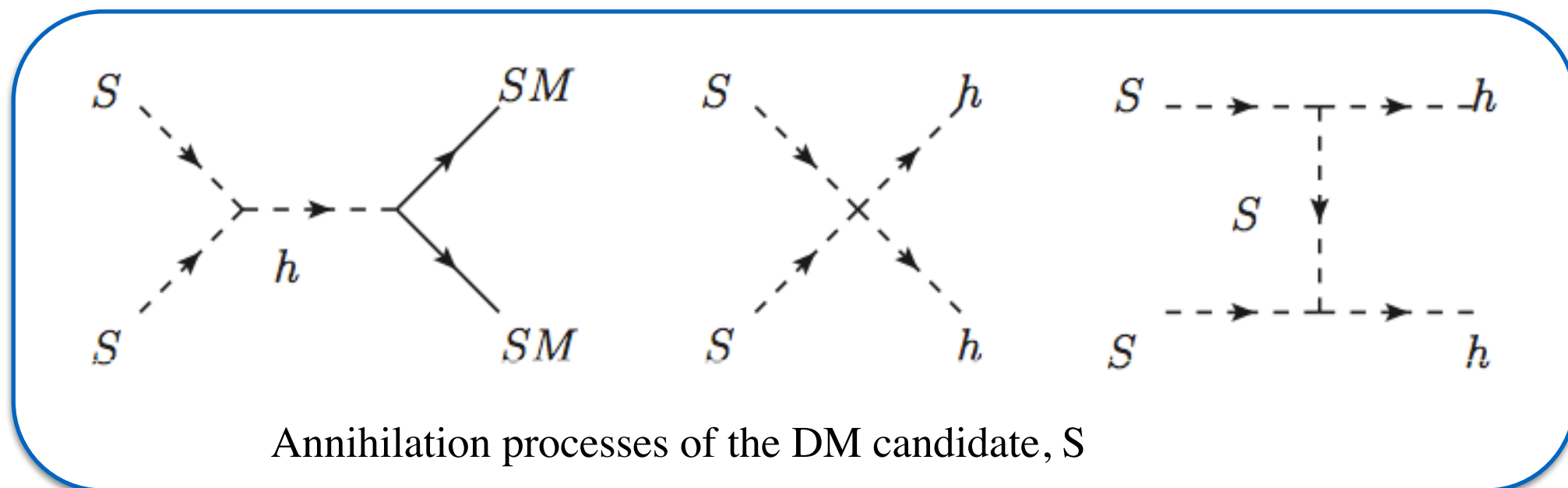


Main message

- The complementarity among the different DM-problem approaches (direct detection, indirect detection, collider searches,..) is crucial to probe any particular DM scenario.
- In minimal scenarios (e.g. Higgs portal) direct detection bounds tend to push the model to parameter space regions where the observation of DM particles at the LHC is very challenging.
- Models constrained by present LHC DM searches typically include extra mediators / interactions.
- Effective Field Theory interactions provide a good description but his validity for LHC DM searches has to be checked when dealing with a particular UV completion.

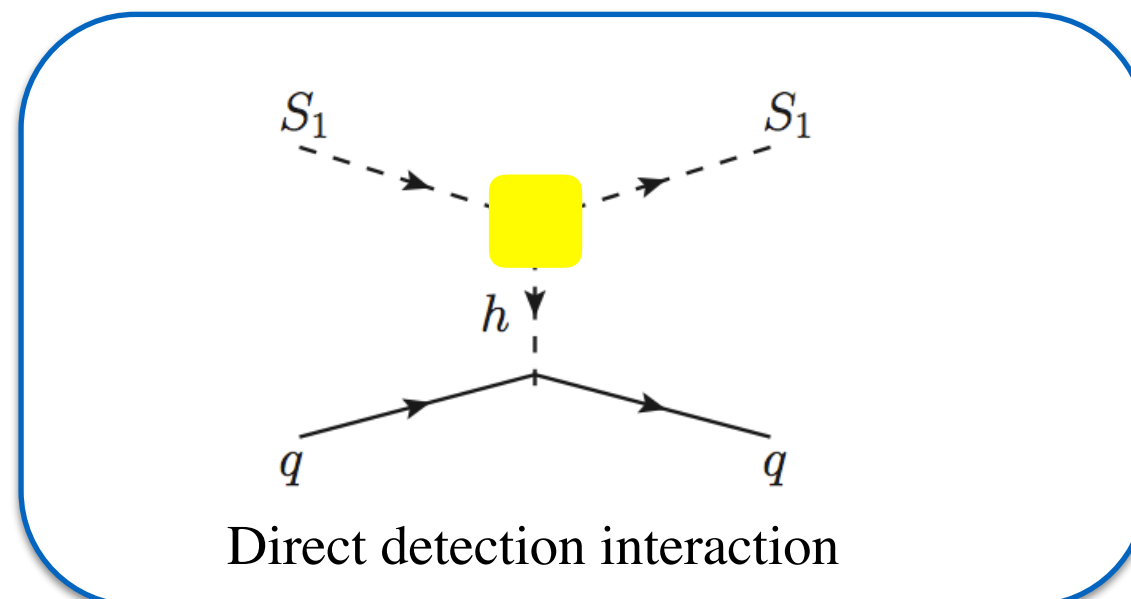
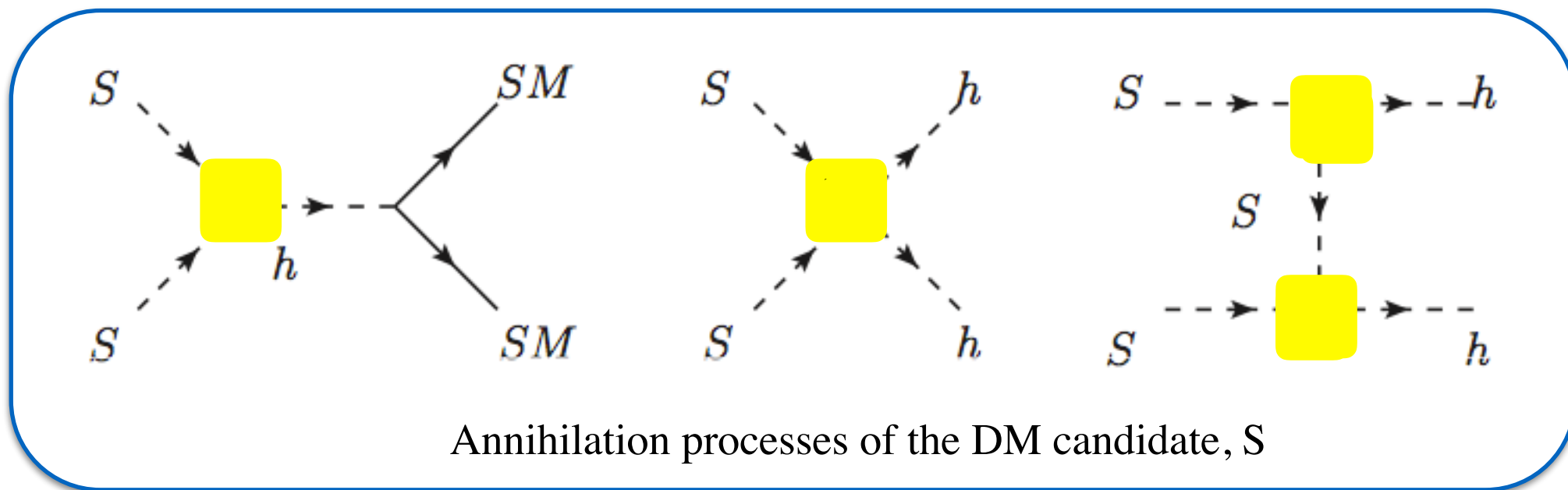
The singlet-scalar Higgs portal (SHP)

$$\mathcal{L}_{\text{SHP}} = \mathcal{L}_{\text{SM}} + \frac{1}{2}\partial_\mu S \partial^\mu S - \frac{1}{2}m_0^2 S^2 - \frac{1}{2}\lambda_S |H|^2 S^2 - \frac{1}{4!}\lambda_4 S^4$$



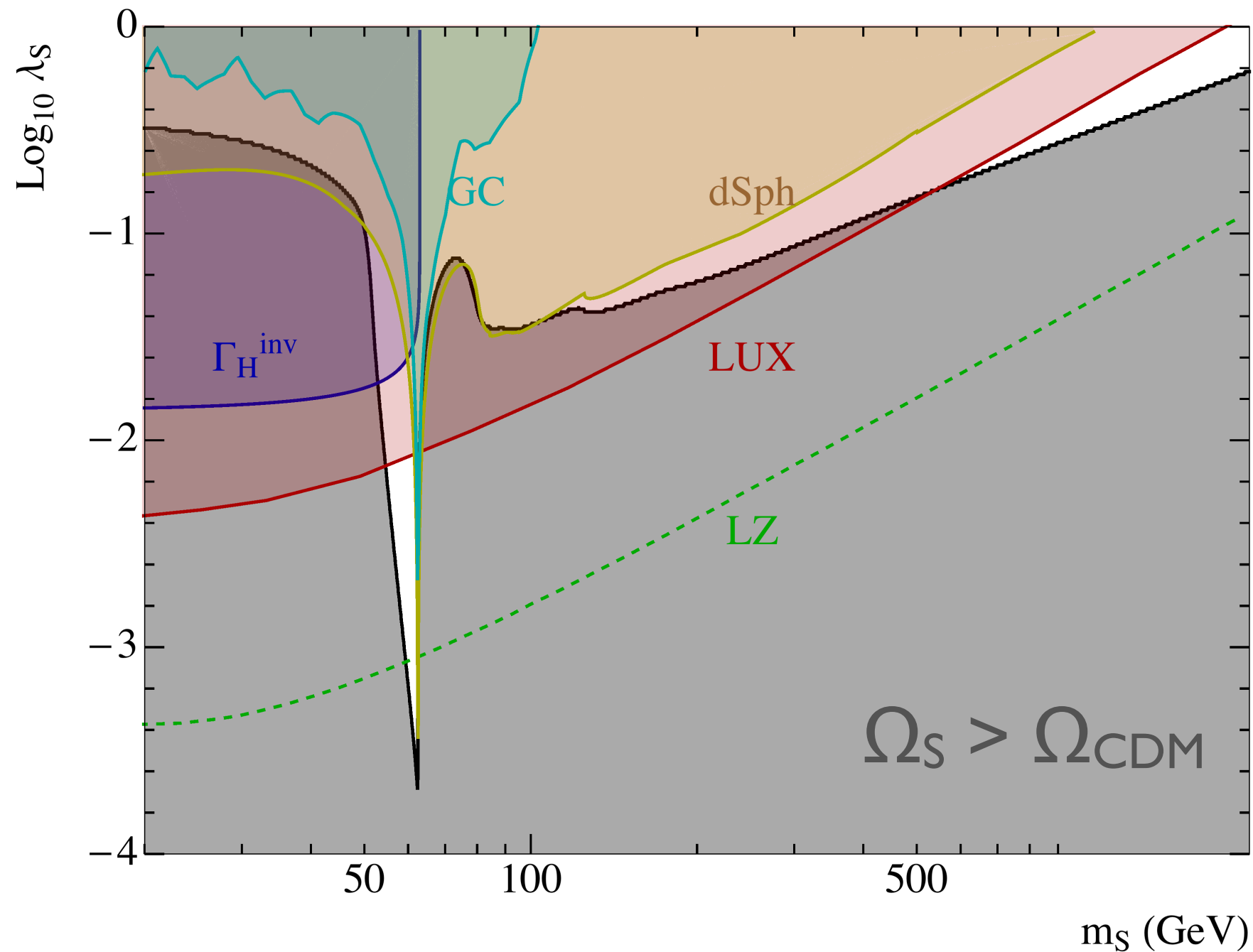
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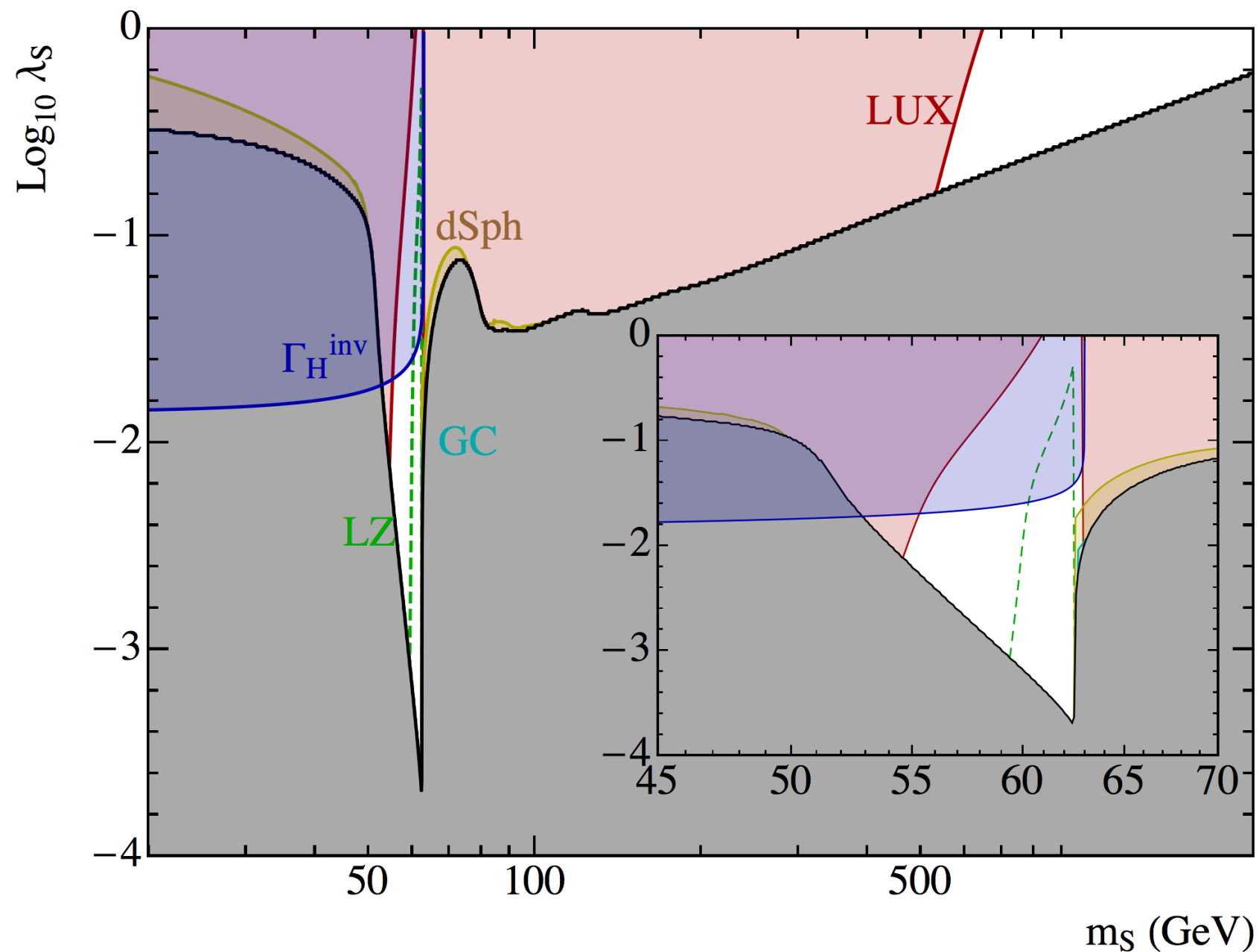
The singlet-scalar Higgs portal (SHP)

Parameter space of the SHP



The singlet-scalar Higgs portal (SHP)

Taking into account the rescaling of the density in regions where $\Omega_S < \Omega_{\text{CDM}}$



LHC constraints on Singlet Higgs Portal

Below threshold
($m_S < m_h / 2$)

The model is efficiently probed (constrained) by the Higgs invisible width, implying $|\lambda_S| \lesssim 10^{-2}$

Above threshold

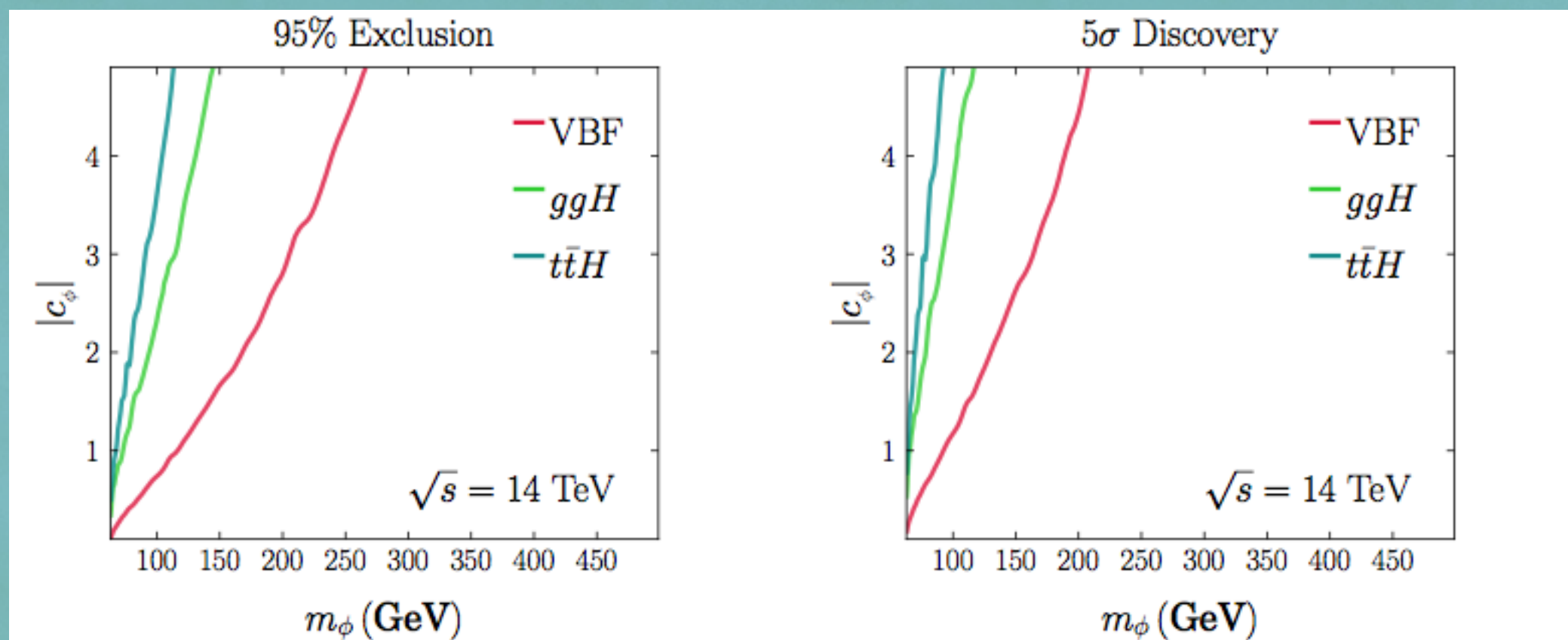
Very challenging! The cross section of process

$$pp \rightarrow h^* + X \rightarrow S S + X$$

is suppressed by $|\lambda_S|^2$ and by two-body phase space

Possible channels:

include \cancel{E}_T in association with monojet, tt, VBF, mono H



$$\mathcal{L} = 3 \text{ ab}^{-1}$$

Some ATLAS & CMS Mono-X searches @ 13 TeV

- **Monojet** ATLAS PRD 94 (2016) 032005,
CMS arXiv:1703.01651
- **Mono-Higgs** ATLAS-CONF-2016-056 | diphoton
CMS PAS EXO-16-011

ATLAS Phys. Lett. B 765 (2017) | bb
CMS PAS EXO-16-012 **See Nicolo Trevisani's talk**
- **tt** CMS PAS-EXO-16-005
PAS-EXO-16-028

ATLAS-CONF-2016-077
ATLAS-CONF-2016-050
ATLAS-CONF-2016-076 **See Jonatan Piedra's talk**
- .. and many others

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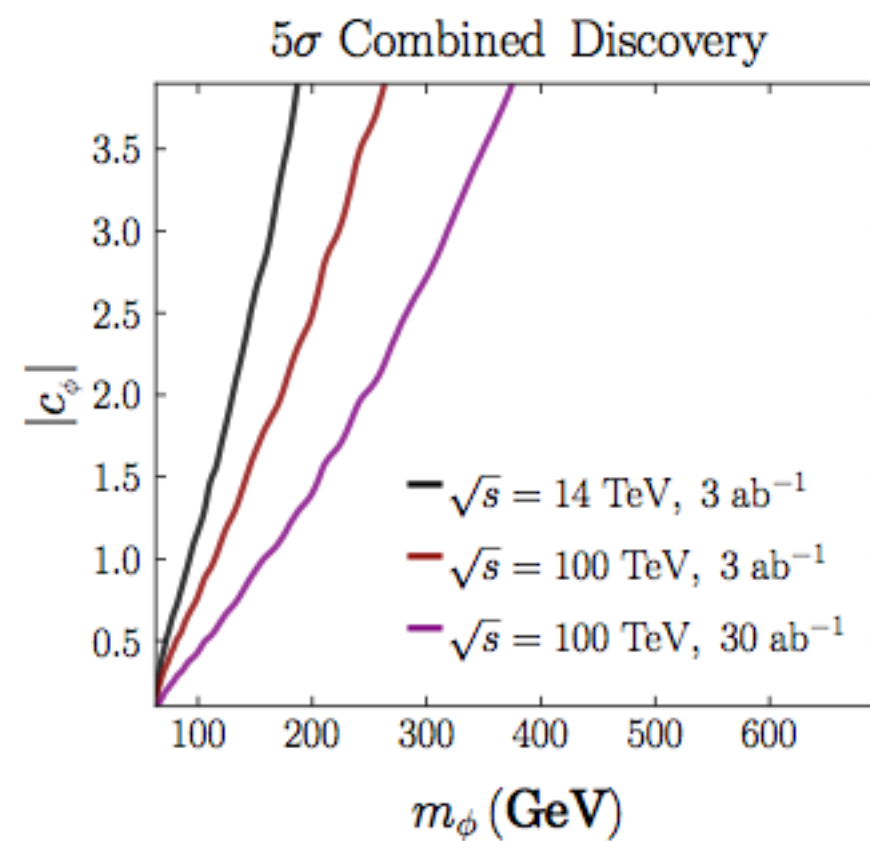
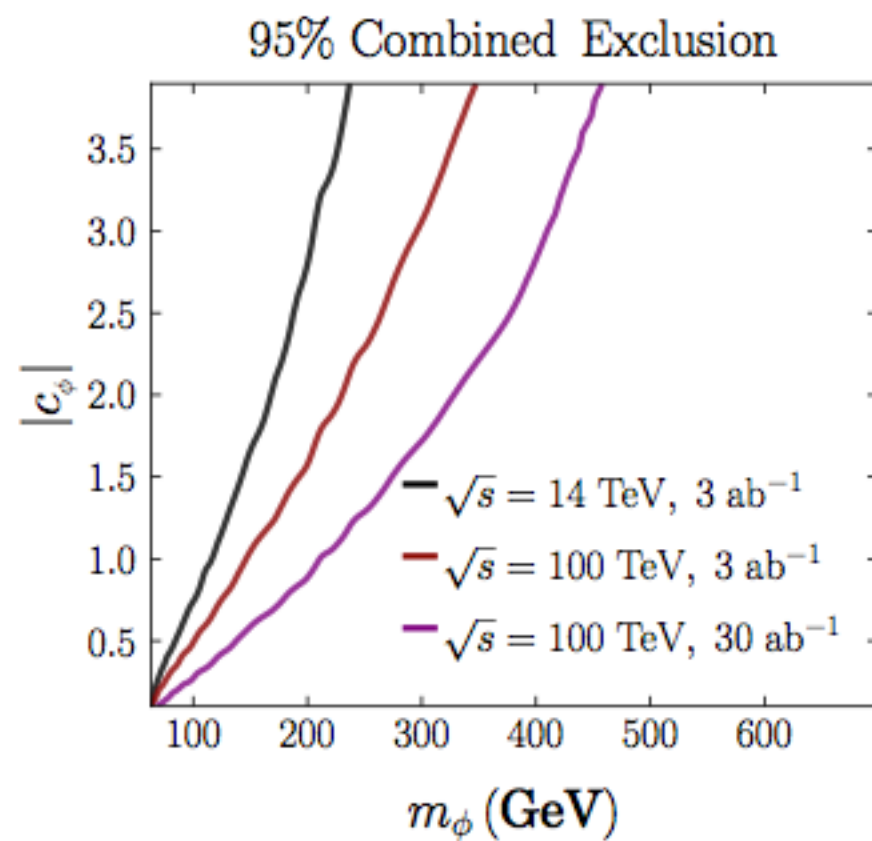
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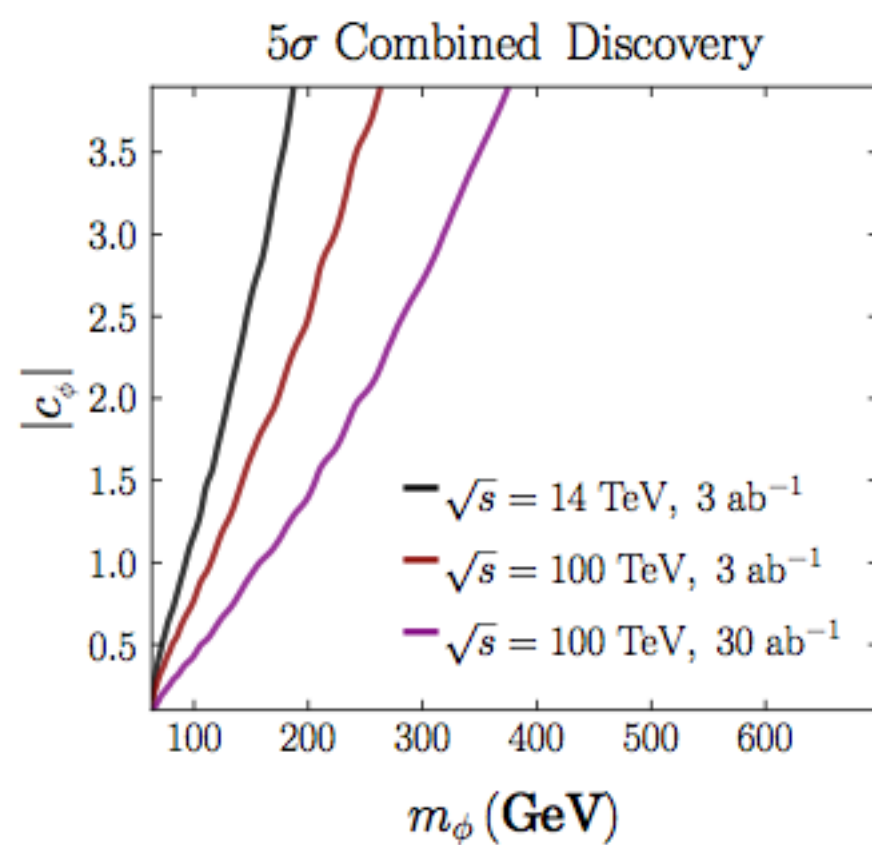
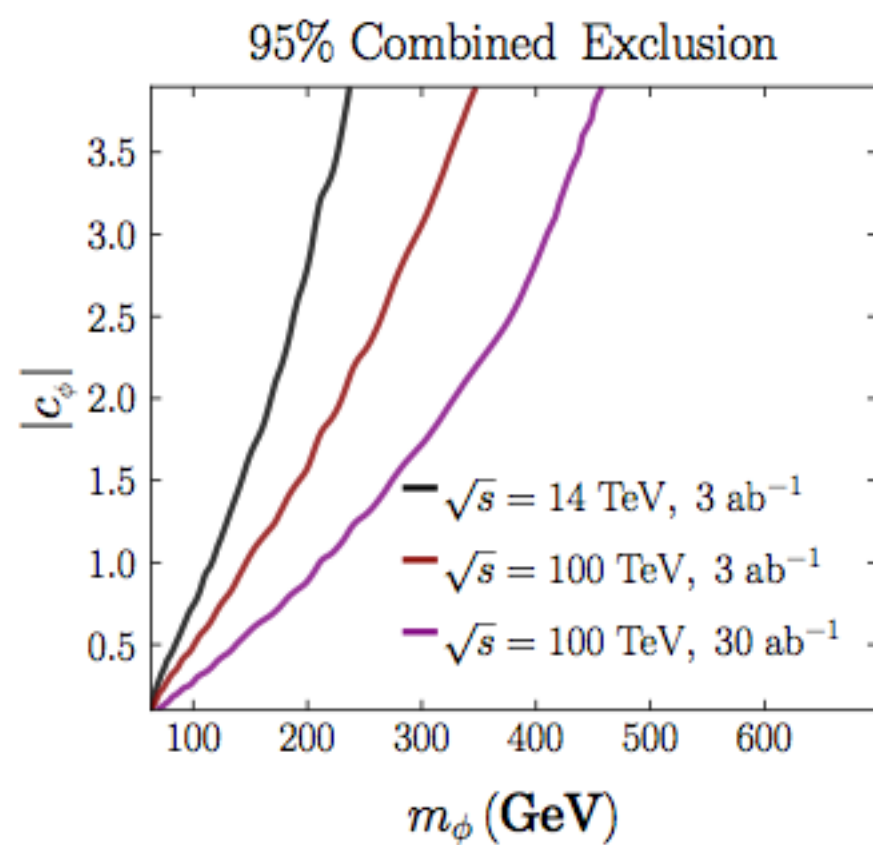
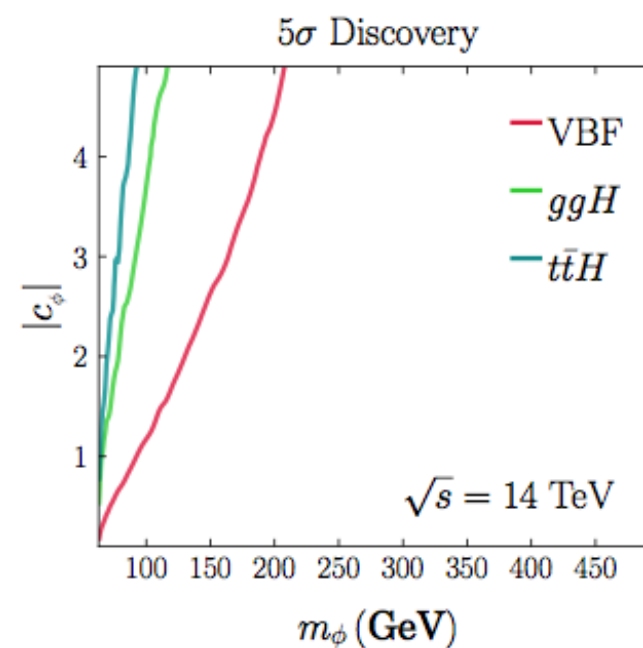
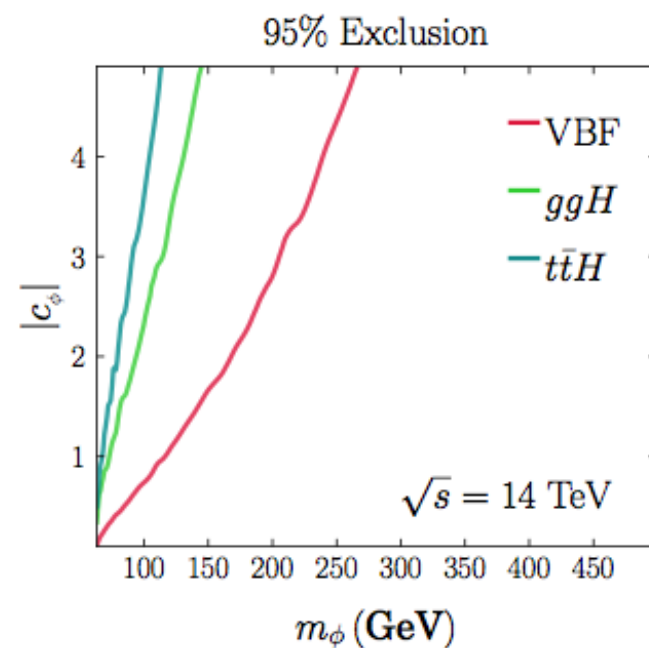
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No significant evidence for dark matter particle production has been observed.

Results are usually interpreted in terms of models with **additional mediators (Z' , Higgses, etc)**





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- Reason: minimality. There is a unique coupling (λ_s) controlling all the processes (annihilation, direct detection, production at the LHC , etc)

WAY OUT:



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To extend the model



- Including higher dimensional effective operators
 - Linear Effective Higgs Portal
 - Non-linear Effective Higgs Portal
- Extending the mediator sector (ej, new Higgses, Z' bosons, etc)



- Including higher dimensional effective operators

- Linear Effective Higgs Portal
- Non-linear Effective Higgs Portal

- Extending the mediator sector
(e_j , new Higgses, Z' bosons, etc)

- Extending the dark sector

-

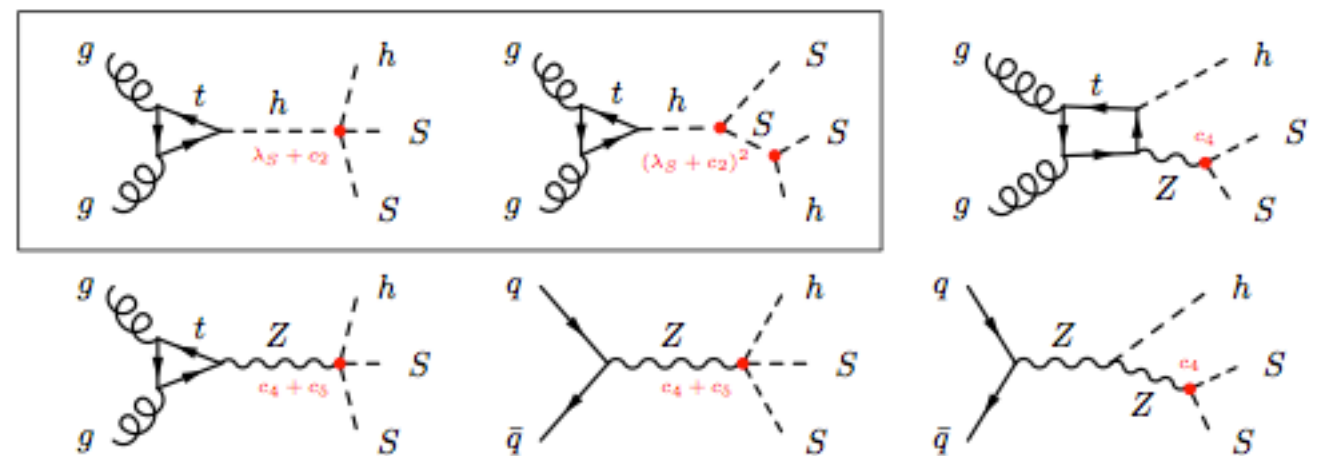


BONUS: LHC signatures could be enhanced in these extensions !!!

Two examples having a mono-Higgs enhancement

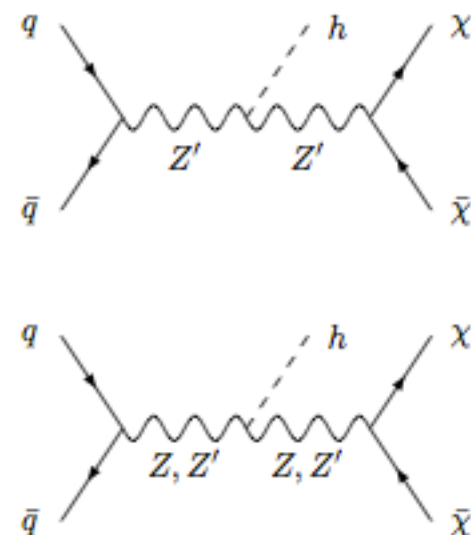
- Non-linear Higgs portal

I. Brivio, M.B. Gavela, L. Merlo, K. Mimasu, J.M. No, R. del Rey, V. Sanz



- Z' mediator

L. Carpenter, A. DiFranzo, M. Mulhearn, C. Shimmin, S. Tulin, D. Whiteson



The simplest extension of the Scalar Higgs Portal

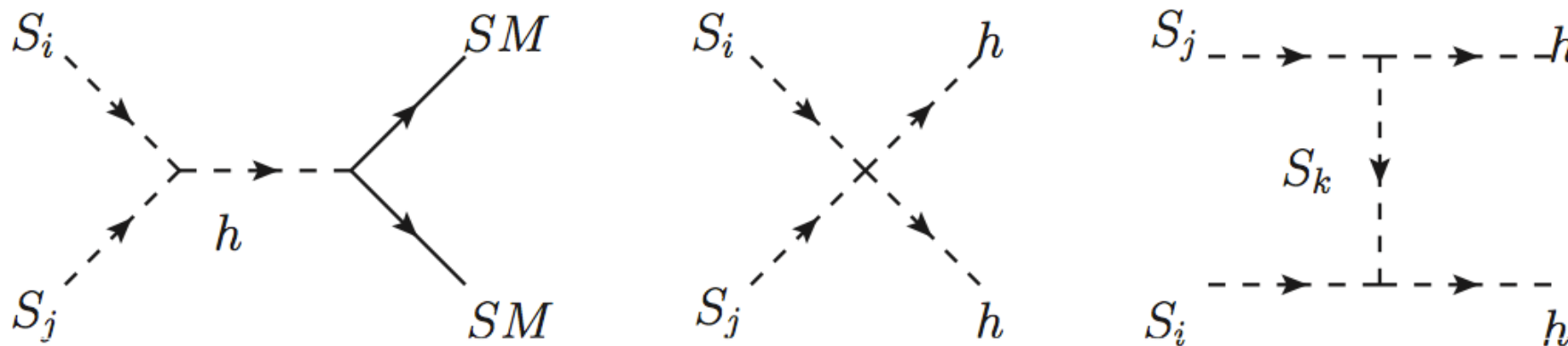
J.A. Casas, D.G. Cerdeño, JMM, J. Quilis

$$\mathcal{L}_{\text{ESHP}} = \mathcal{L}_{\text{SM}} + \frac{1}{2} \sum_{i=1,2} \left[(\partial_\mu S_i)^2 - m_i^2 S_i^2 - \frac{1}{12} \lambda_{i4} S_i^4 \right] - \frac{1}{6} \lambda_{13} S_1 S_2^3 - \frac{1}{6} \lambda_{31} S_1^3 S_2 - \frac{1}{4} \lambda_{22} S_1^2 S_2^2$$

$$- \frac{1}{2} \lambda_1 S_1^2 |H|^2 - \frac{1}{2} \lambda_2 S_2^2 |H|^2 - \lambda_{12} S_1 S_2 \left(|H|^2 - \frac{v^2}{2} \right)$$

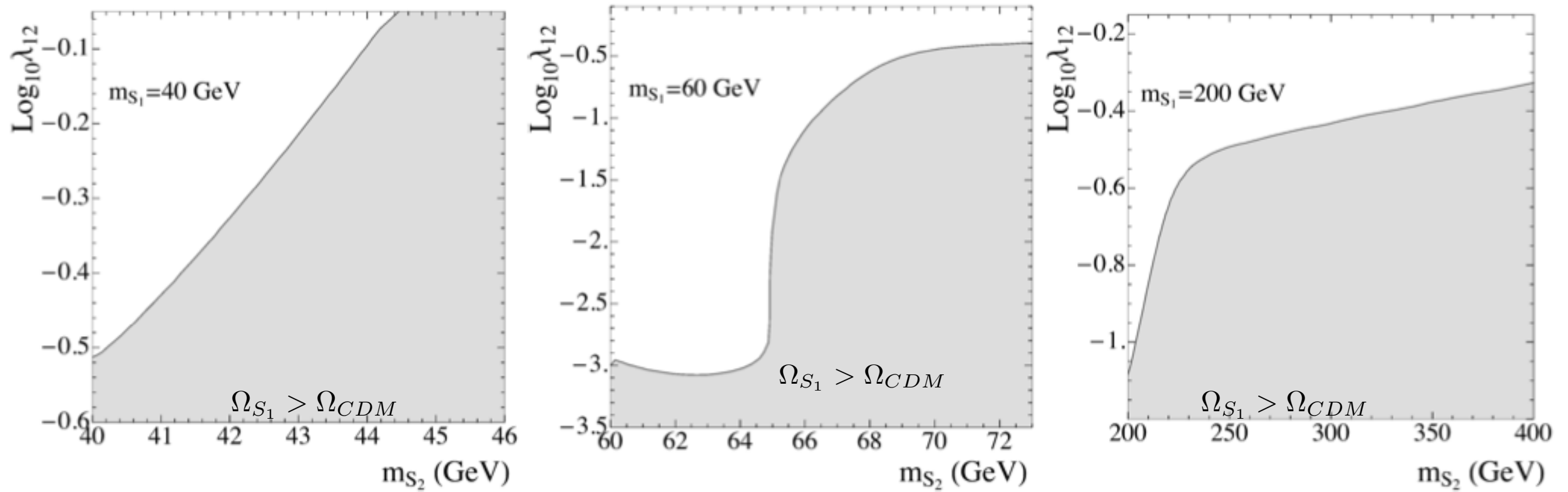
We assume $m_{S1} < m_{S2}$

Notice that λ_{12} induces new (co)annihilation processes

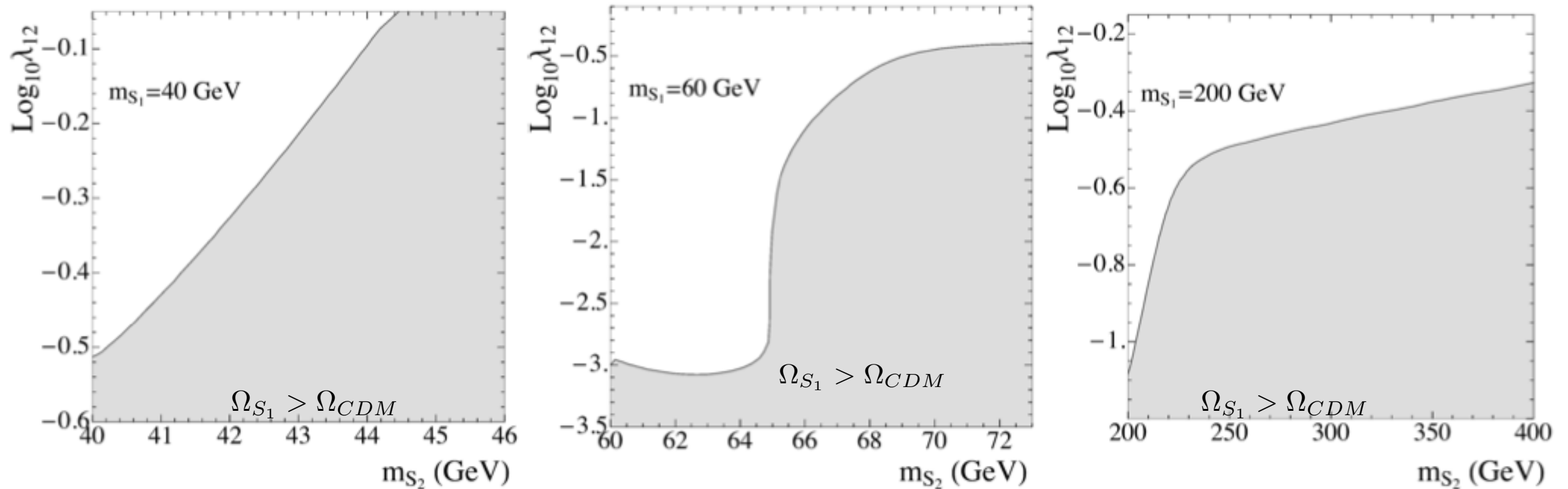


that could get ride of the excess of DM in the low λ_1 region,
that is safe from direct DM constraints, enlarging the parameter space.

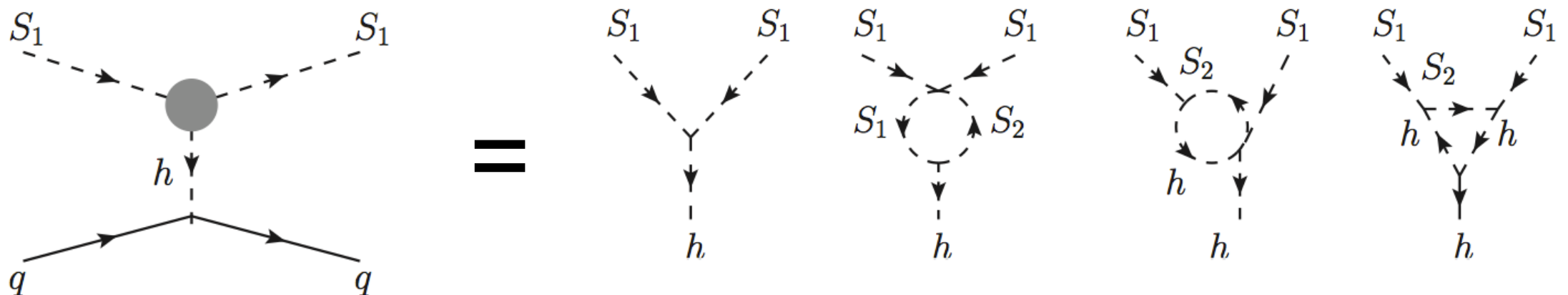
- DM density for small λ_1 values . We assume $\lambda_1 = \lambda_{12}^2 / (4\pi)^2$

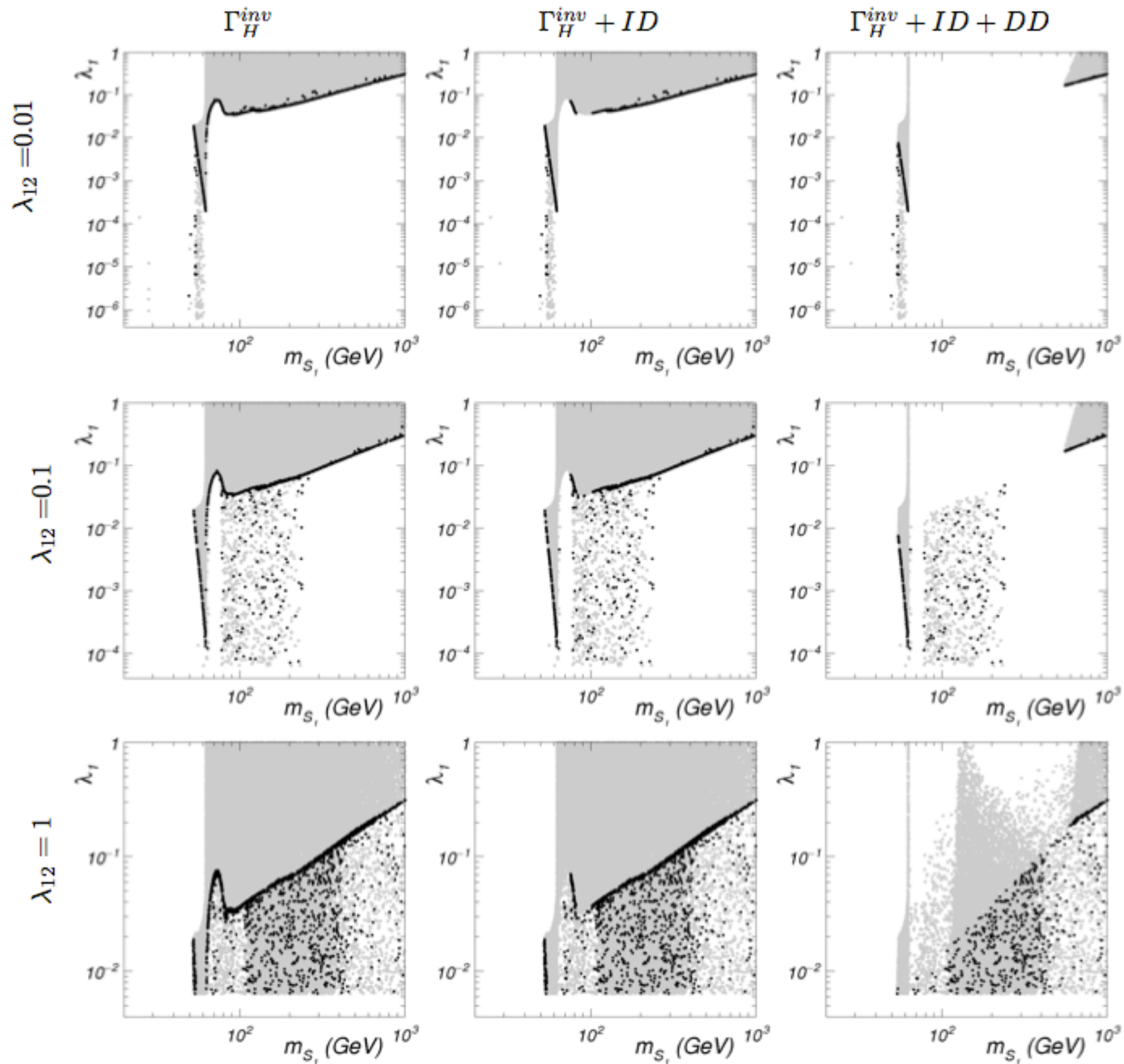


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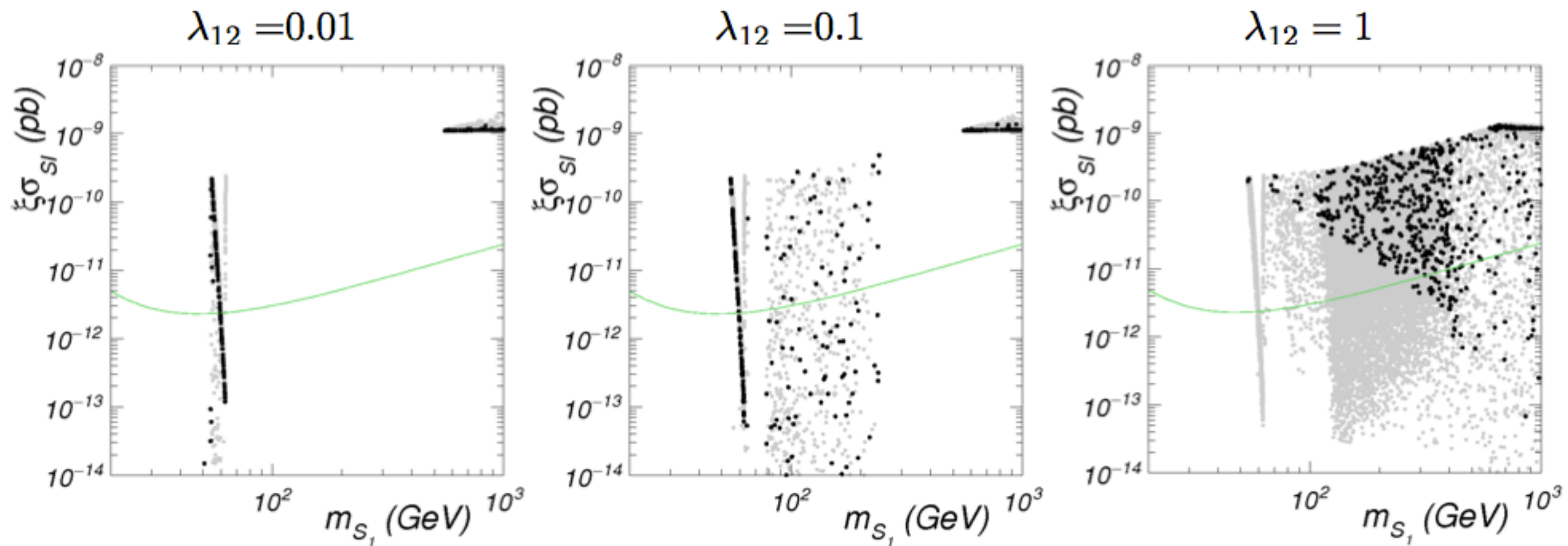
- Warning: radiative corrections to DM direct detection could be potentially large in the evaluation of the bounds







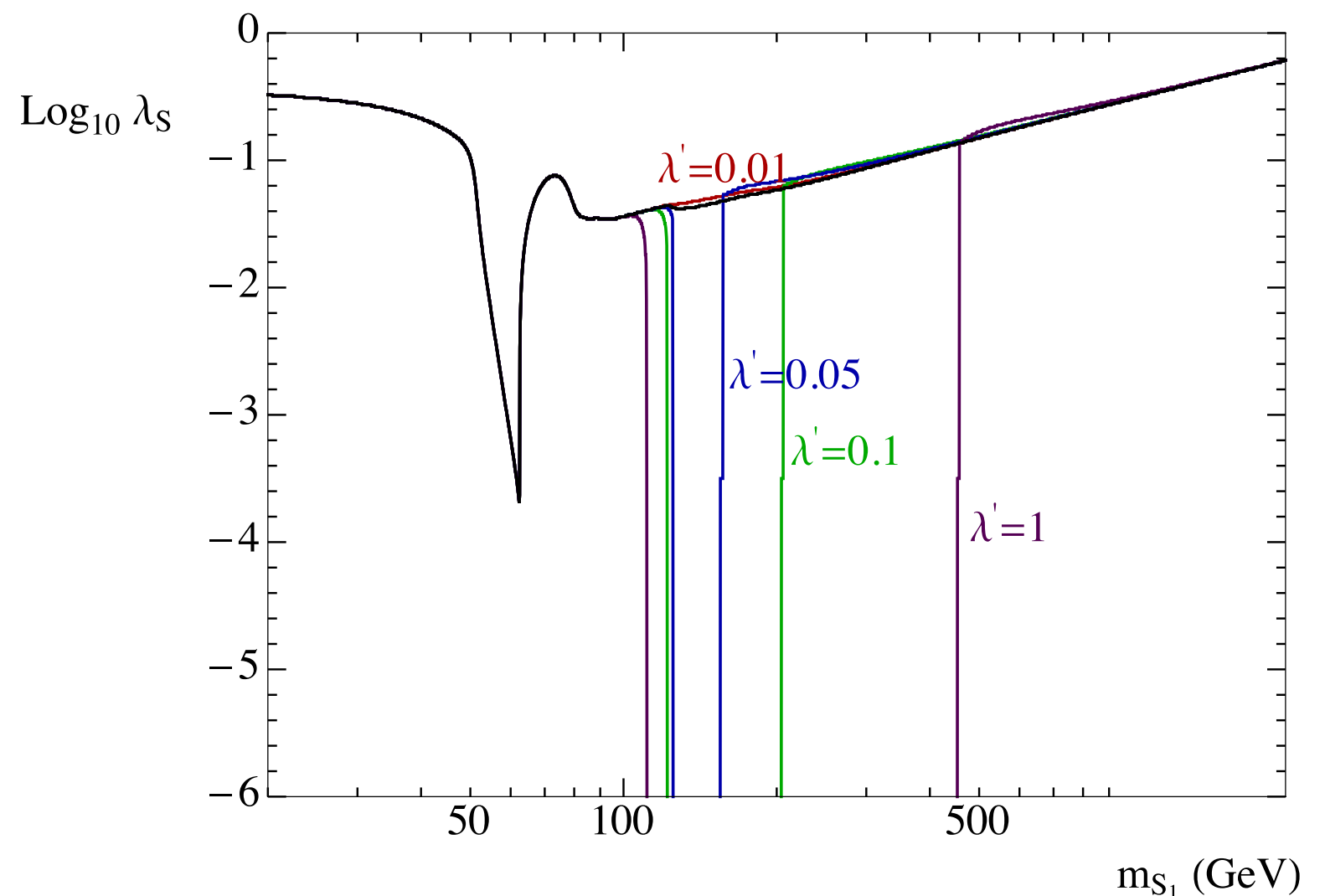
Large regions of the Scalar Higgs Portal parameter space are now recovered, even for moderate values of λ_{12}



The results can be nicely understood in terms of the Singlet Higgs portal + 1 effective operator

$$\Delta\mathcal{L}_{\text{eff}}(S_1, H) = \frac{1}{2} \frac{\lambda_{12}^2}{m_{S_2}^2} S_1^2 \left(|H|^2 - \frac{v^2}{2} \right)^2 + \dots$$

$$\mathcal{L}'_{\text{SHP}} = \mathcal{L}_{\text{SHP}} + \frac{1}{2} \frac{\lambda'}{m_{S_1}^2} S_1^2 \left(|H|^2 - \frac{v^2}{2} \right)^2$$

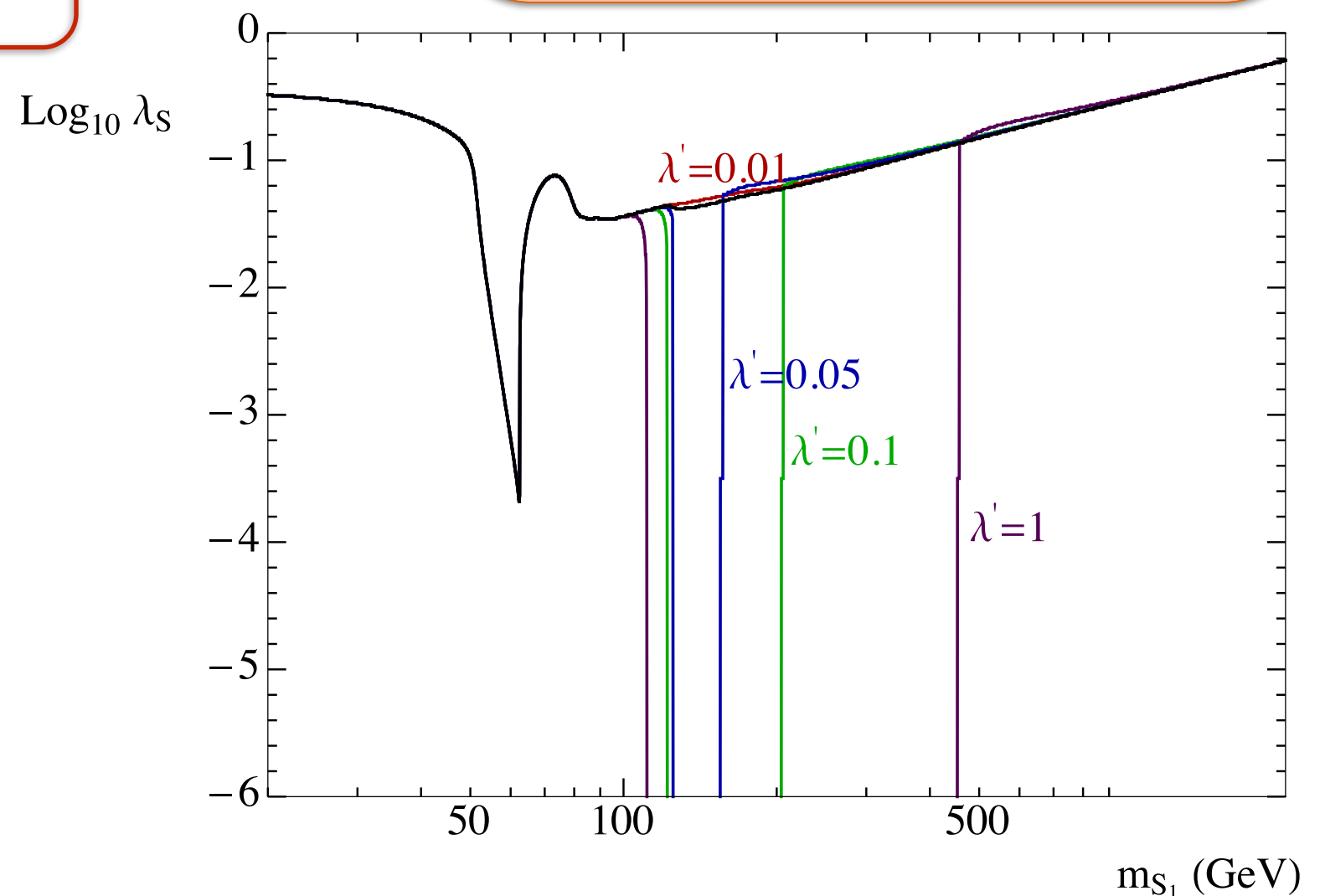


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Intermediate O(100-300 GeV) m_{S_1} values are now allowed



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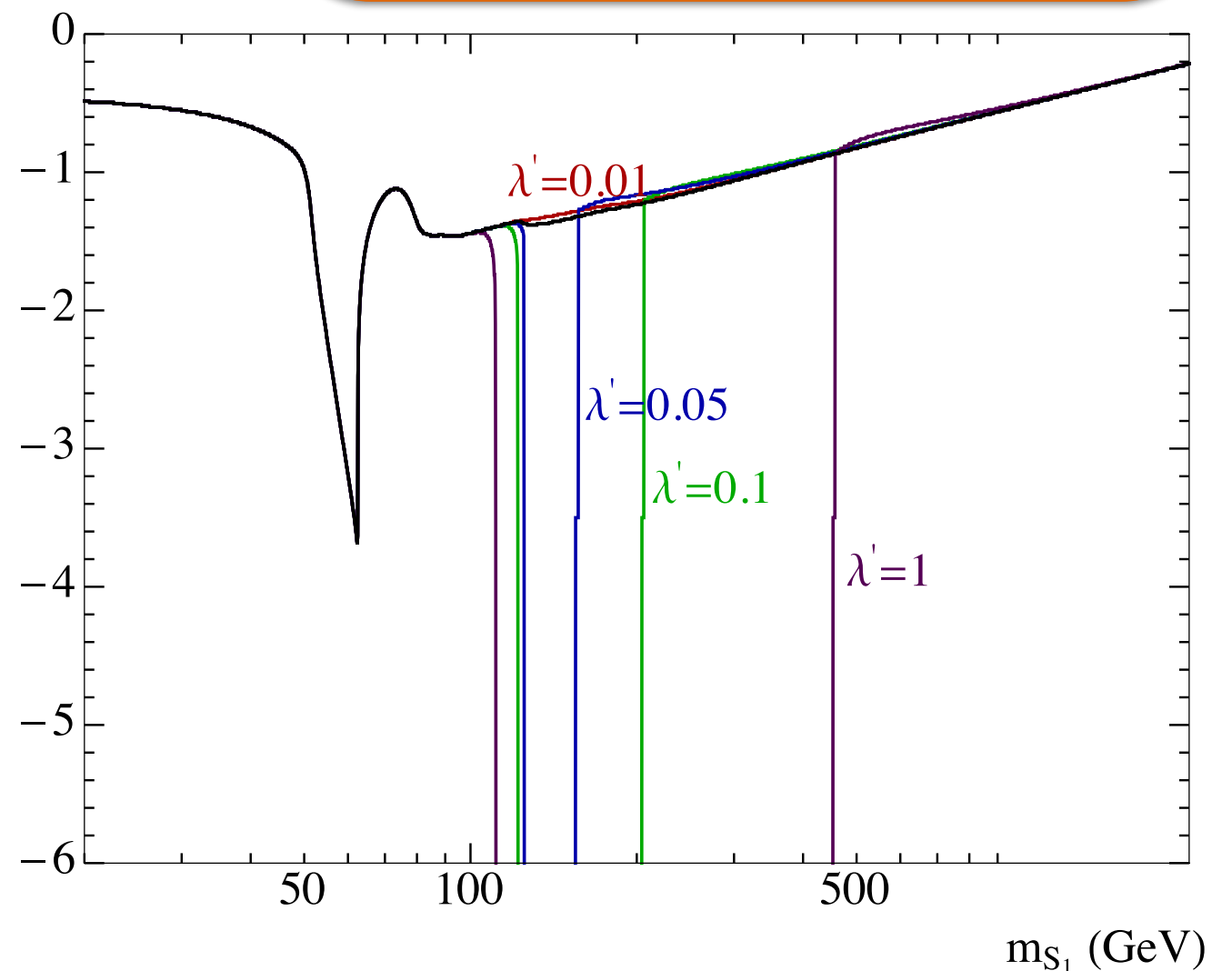
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Intermediate O(100-300 GeV) m_{S_1} values are now allowed

- Enlarges the hhSS vertex without increasing the hhS one
- Could enhance DM signals @ LHC (ej mono-Higgs signatures) without being in conflict with DM direct detection bounds

$\text{Log}_{10} \lambda_S$



.. and again

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