

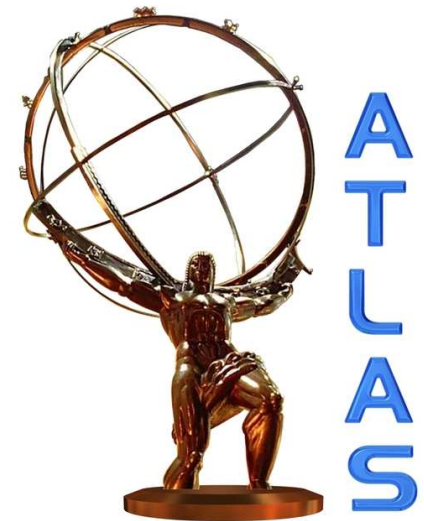
# Photon (+jet) and diphoton production at the LHC with the ATLAS detector

J. Terrón (Universidad Autónoma de Madrid)

On behalf of the ATLAS UAM team

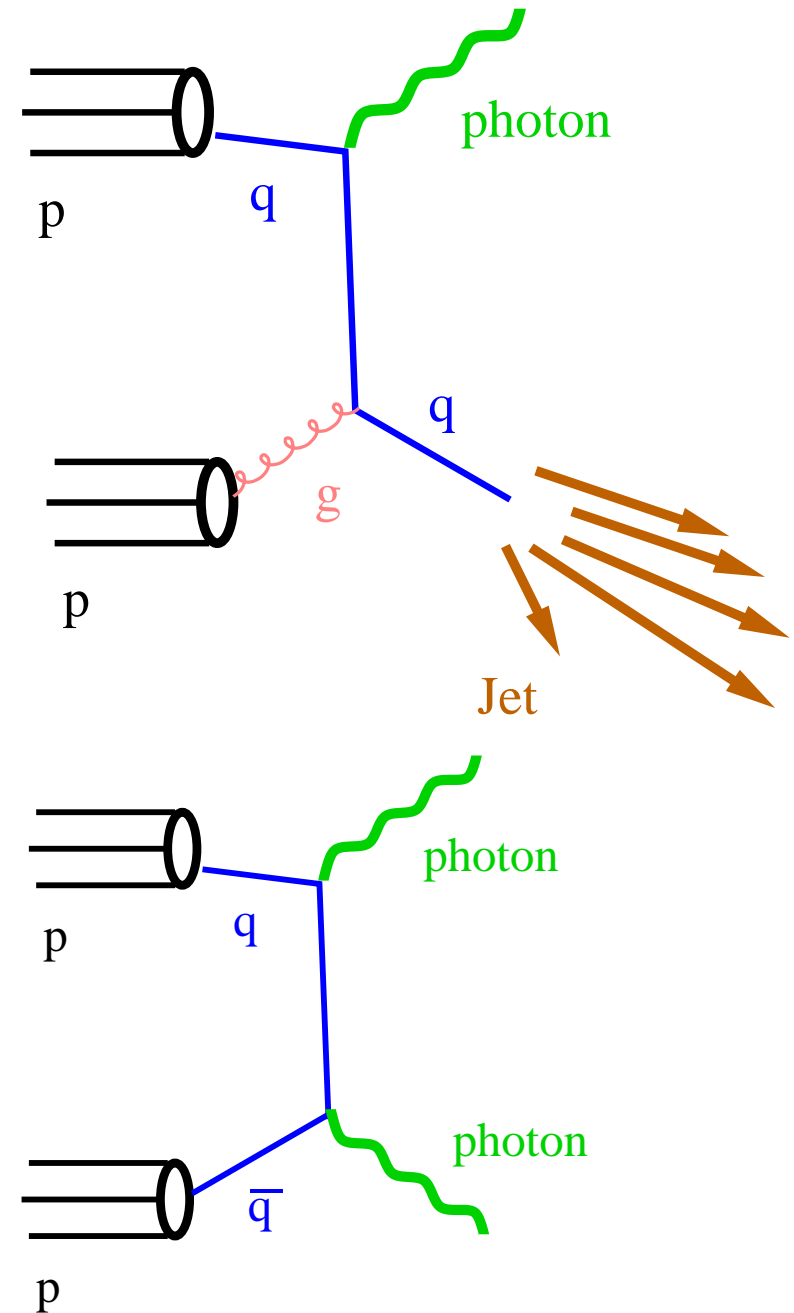
## ● Outline

- Physics with photons
- Inclusive photon production at 8 TeV
- Photon + jet(s) production at 8 TeV
- Photon pair production at 8 TeV
- Inclusive photon production at 13 TeV
- Summary



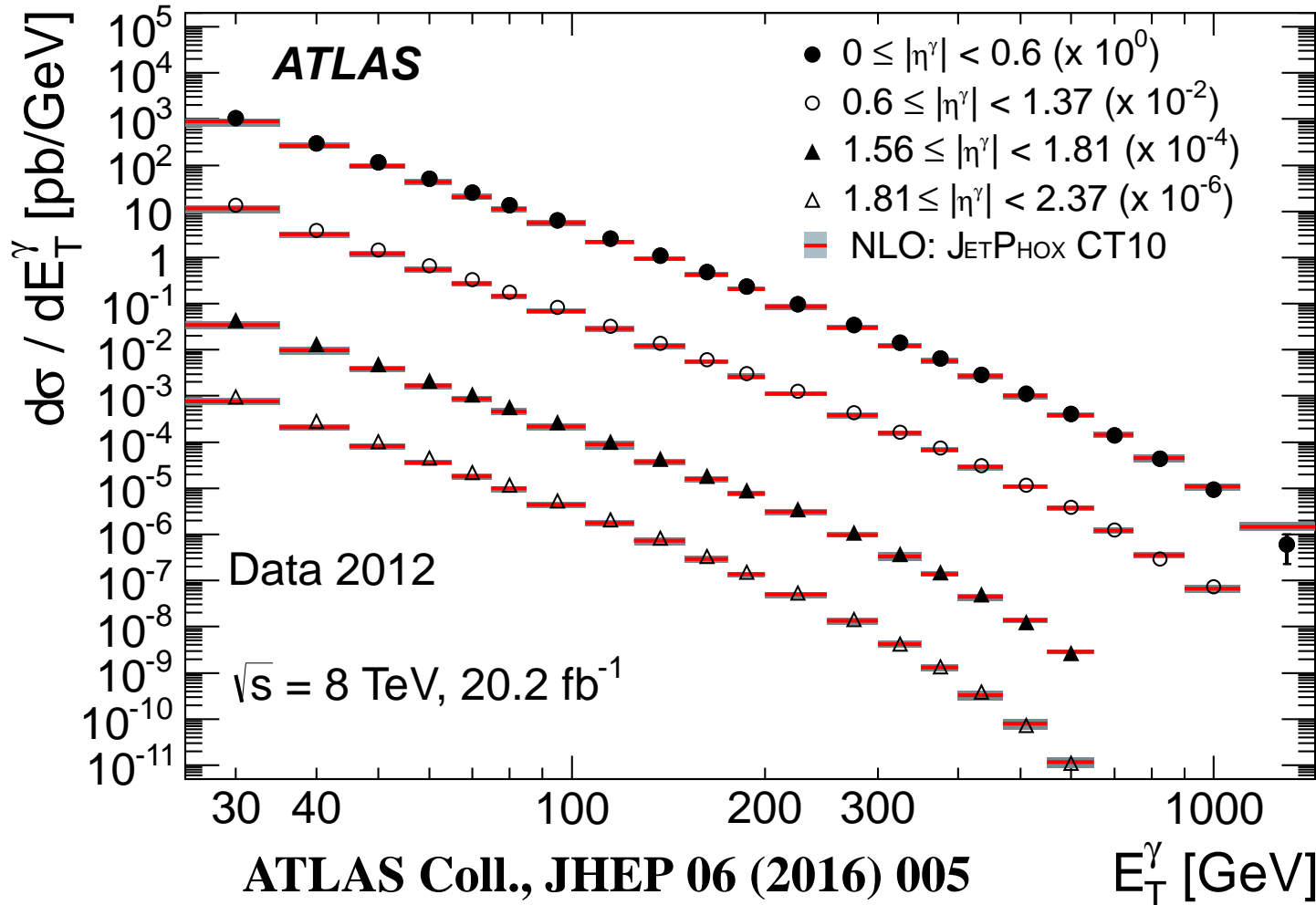
## Photon production in $pp$ collisions at LHC

- Photon production in  $pp$  collisions allows
  - tests of perturbative QCD
  - to extract information on the proton PDFs
- Possibilities to study inclusive production of photons or in association with jets
- Prompt photons represent a cleaner probe of the hard interaction than jet production
- Aid searches involving photons or  $E_T^{\text{miss}} + \text{jets}$  (through ratios  $Z + \text{jets} / \gamma + \text{jets}$ )
- Diphoton production is of special interest as the major background to  $H \rightarrow \gamma\gamma$



# Inclusive photon production at 8 TeV

# Inclusive isolated-photon production in $pp$ collisions at $\sqrt{s} = 8$ TeV



- Measurement of  $d\sigma/dE_T^\gamma$  for  $25 < E_T^\gamma < 1500$  GeV and different ranges in  $\eta^\gamma$  using  $\mathcal{L} = 20.2 \text{ fb}^{-1}$  of  $pp$  collision data at  $\sqrt{s} = 8$  TeV

$$E_T^{\text{iso}} < 4.2 \cdot 10^{-3} \cdot E_T^\gamma + 4.8 \text{ GeV}$$

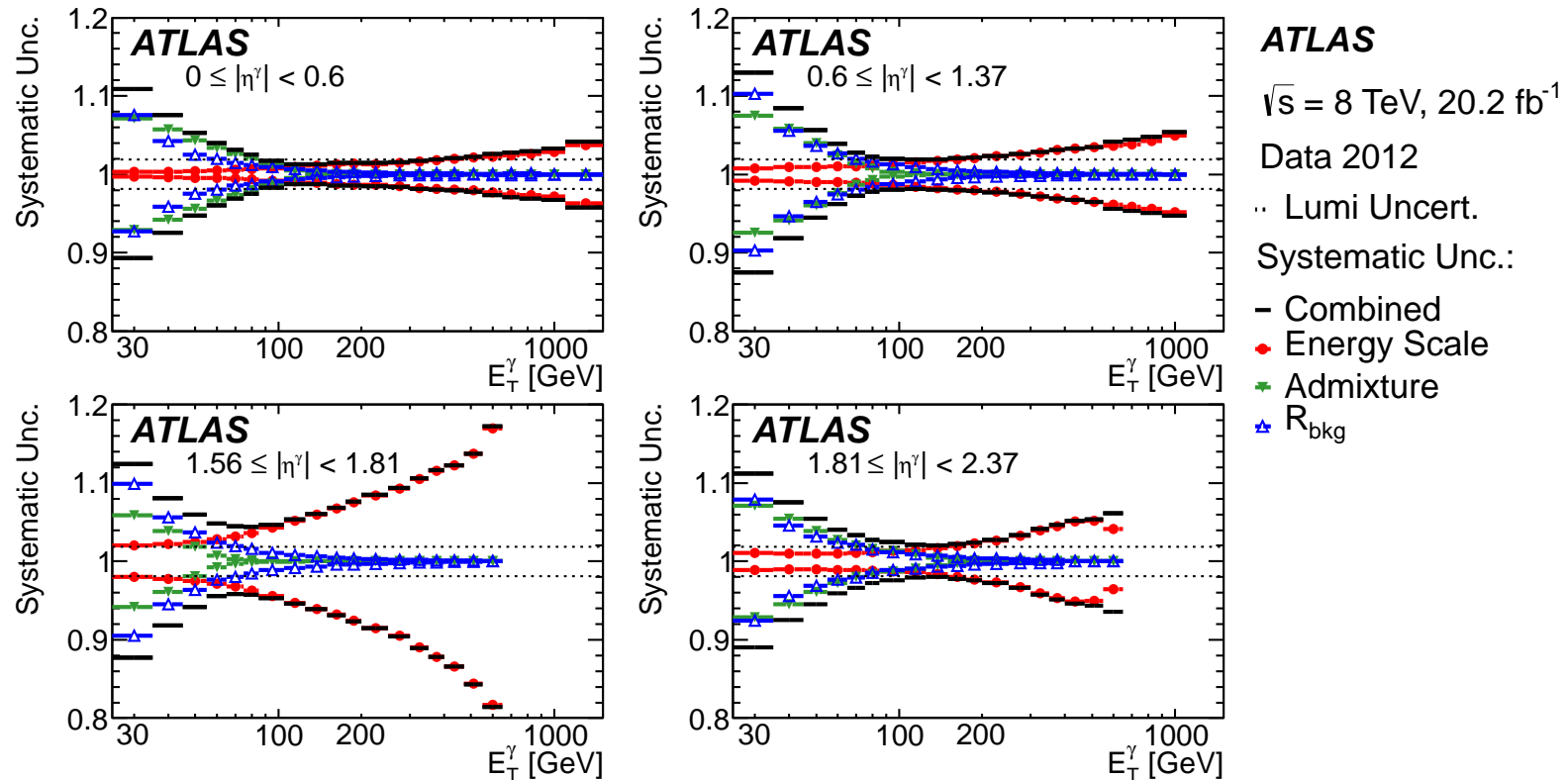
- The measurement covers ten orders of magnitude in cross section

- First measurement of photon production with  $E_T^\gamma > 1$  TeV

- Significant improvement in experimental uncertainties over the previous measurements
- Good description (in log scale) of the data by NLO QCD calculations using JetPhox

# Major experimental uncertainties

ATLAS Coll., JHEP 06 (2016) 005



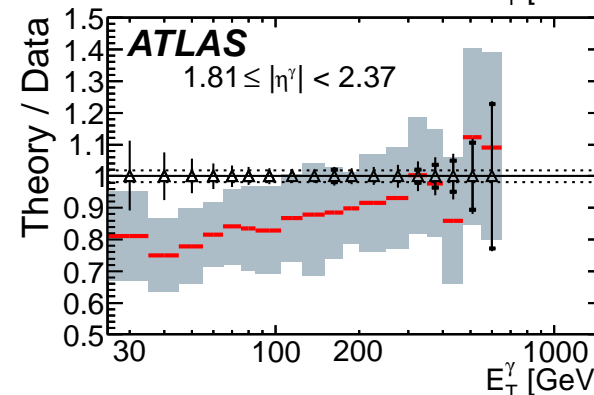
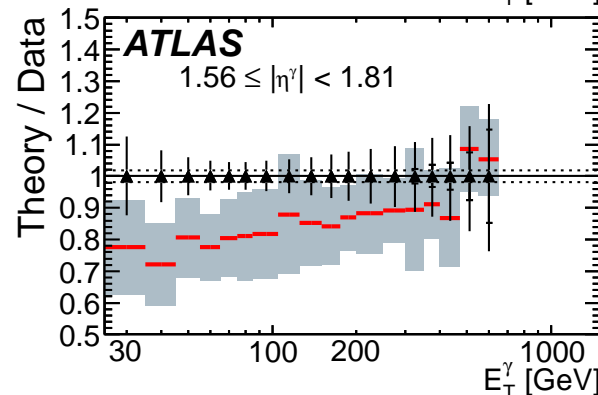
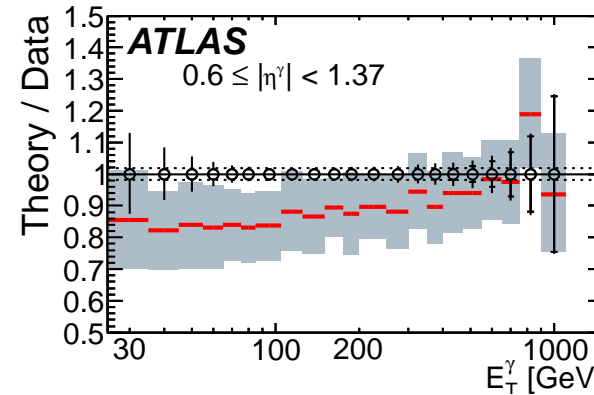
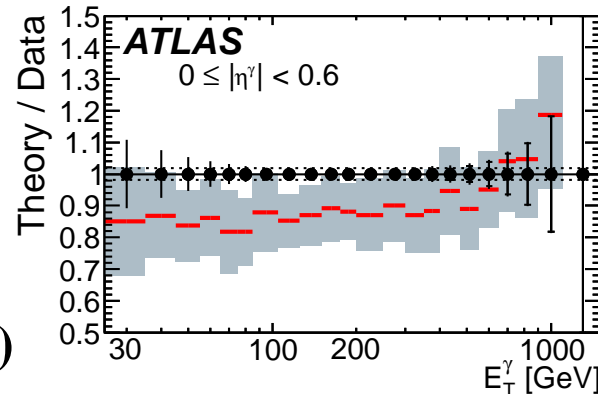
- The uncertainty on the photon energy scale\* (about 1% except in the region  $1.56 < |\eta^\gamma| < 1.81$ ) is dominant at high  $E_T^\gamma$
  - The uncertainty on the correlation in the background ( $\pm 10\%$ ) dominates at low  $E_T^\gamma$ , but negligible at high  $E_T^\gamma$
  - The uncertainty on the admixture of direct and fragmentation photons increases at low  $E_T^\gamma$
- \* (ATLAS Collaboration, Eur. Phys. J. C74 (2014) 3071)

# Inclusive isolated-photon cross sections vs NLO QCD

ATLAS Coll., JHEP 06 (2016) 005

Theoretical  
uncertainties  
much larger than(!)

experimental  
uncertainties



ATLAS

 $\sqrt{s} = 8 \text{ TeV}, 20.2 \text{ fb}^{-1}$ 

Data 2012

●  $0 \leq |\eta^\gamma| < 0.6$   
○  $0.6 \leq |\eta^\gamma| < 1.37$   
▲  $1.56 \leq |\eta^\gamma| < 1.81$   
△  $1.81 \leq |\eta^\gamma| < 2.37$   
⋯ Lumi Uncert.

NLO:

■ JETPHOX CT10

- Comparison to NLO QCD calculation using the JetPhox program

→ a similar trend is observed at low  $E_T^\gamma$  in all  $|\eta^\gamma|$  regions, the NLO QCD predictions underestimate the data by  $\approx 20\%$

→ the theoretical uncertainty (12-20%) prevents a more precise test of the SM predictions

- Halving the measured uncertainties compared to previous measurements

⇒ useful constraint on proton PDFs once included in a global fit

# Photon+jet(s) production at 8 TeV

# Dynamics of $\gamma + \text{jet}$ production in $pp$ collisions at $\sqrt{s} = 8 \text{ TeV}$

- Study of the  $\gamma + \text{jet}$  dynamics by measuring the differential cross sections as functions of

→ Photon:  $E_T^\gamma$

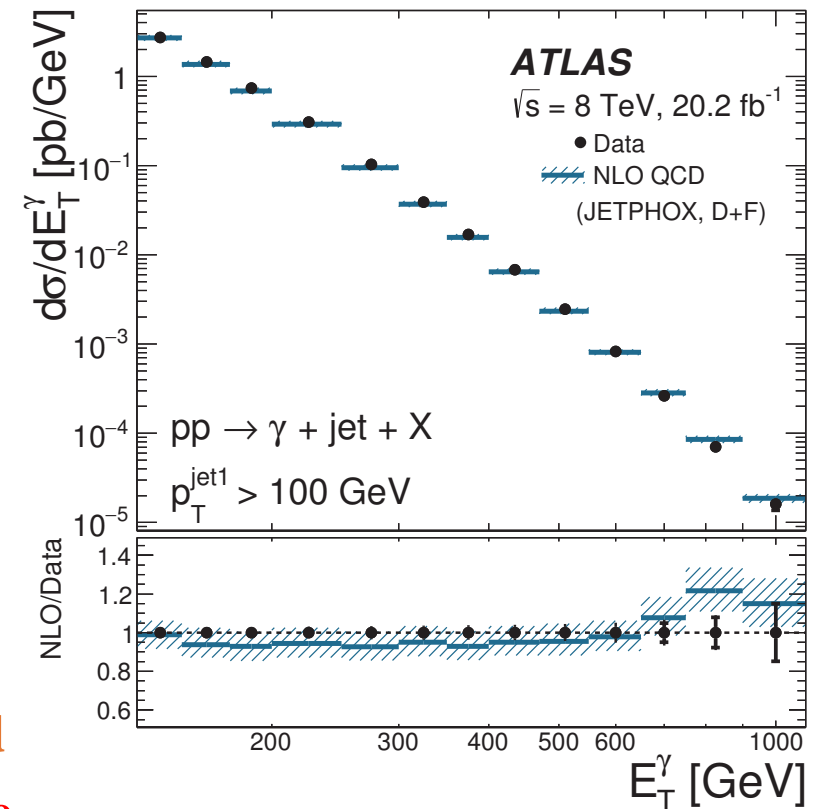
→ Leading jet:  $p_T^{\text{jet}1}$

→ Photon+Leading jet:  $m^{\gamma\text{-jet}1}, \cos \theta^*$

where  $\cos \theta^* = \tanh \frac{1}{2} (y^{\text{jet}1} - \eta^\gamma)$

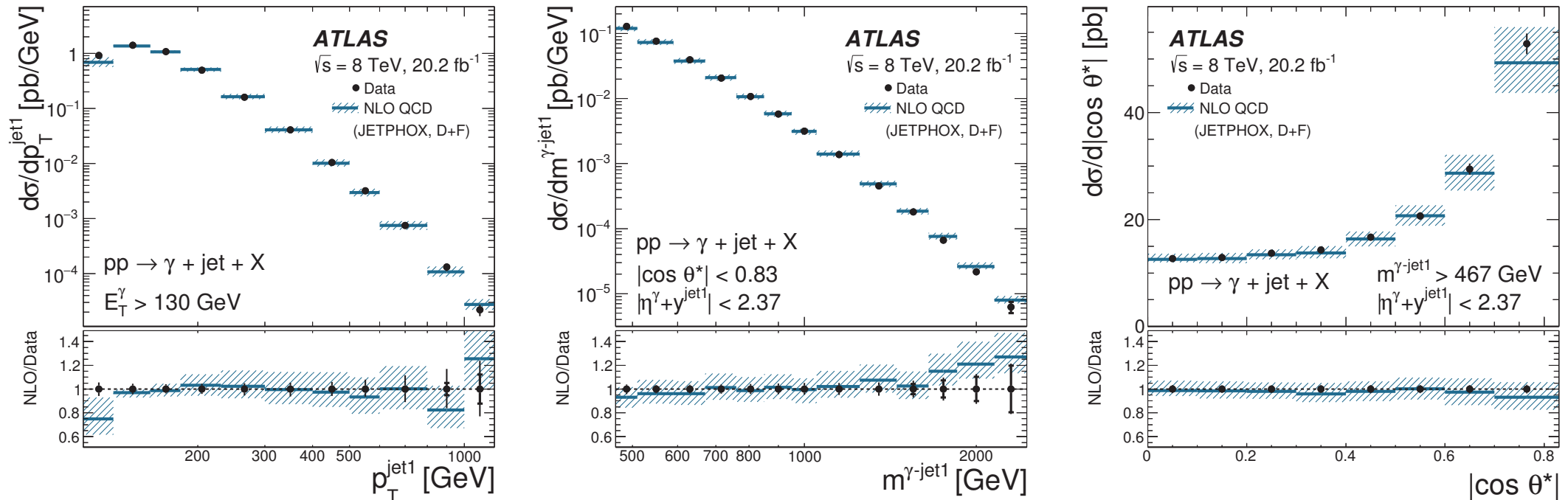
$\theta^*$  = scattering angle in centre-of-mass frame for  $2 \rightarrow 2$  hard collinear scattering

- Measurements in the phase-space region defined by:  $E_T^\gamma > 130 \text{ GeV}$ ,  $|\eta^\gamma| < 2.37$  (excluding the region  $1.37 < |\eta^\gamma| < 1.56$ ),  $p_T^{\text{jet}1} > 100 \text{ GeV}$ ,  $|y^{\text{jet}1}| < 4.4$  (anti- $k_t$  algorithm with  $R = 0.6$ ),  $E_T^{\text{iso}} < 10 \text{ GeV}$  and  $\Delta R_{\gamma j} > 1$
- Comparison to NLO QCD calculation (JETPHOX) corrected for non-perturbative effects
- Good description of the measured  $d\sigma/dE_T^\gamma$  by the NLO QCD calculations





# Dynamics of $\gamma + \text{jet}$ production in $pp$ collisions



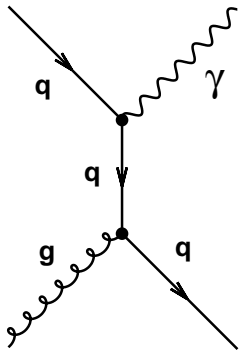
ATLAS Coll., NPB918 (2017) 257

- Additional requirements for  $d\sigma/dm^{\gamma\text{-jet}1}$  and  $d\sigma/d|\cos \theta^*|$  to remove biases:  
 $|\eta^\gamma + y^{\text{jet}1}| < 2.37$  ,  $|\cos \theta^*| < 0.83$  ,  $m^{\gamma\text{-jet}1} > 467 \text{ GeV}$
- In the selected (unbiased) region the angular distribution increases as  $|\cos \theta^*|$  increases
- Good description of the data by the NLO QCD calculations within the (small) experimental and theoretical uncertainties  $\Rightarrow$  validation of the description of the dynamics of  $\gamma + \text{jet}$  production in  $pp$  collisions at  $\mathcal{O}(\alpha_{em}\alpha_s^2)$

# Dynamics of $\gamma + \text{jet}$ production in $pp$ collisions

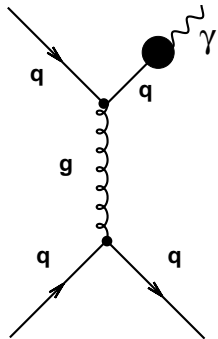
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- Angular distribution  $d\sigma/d|\cos\theta^{\gamma j}|$  sensitive to the spin of the exchanged (virtual) particle: quark(1/2) vs gluon(1)



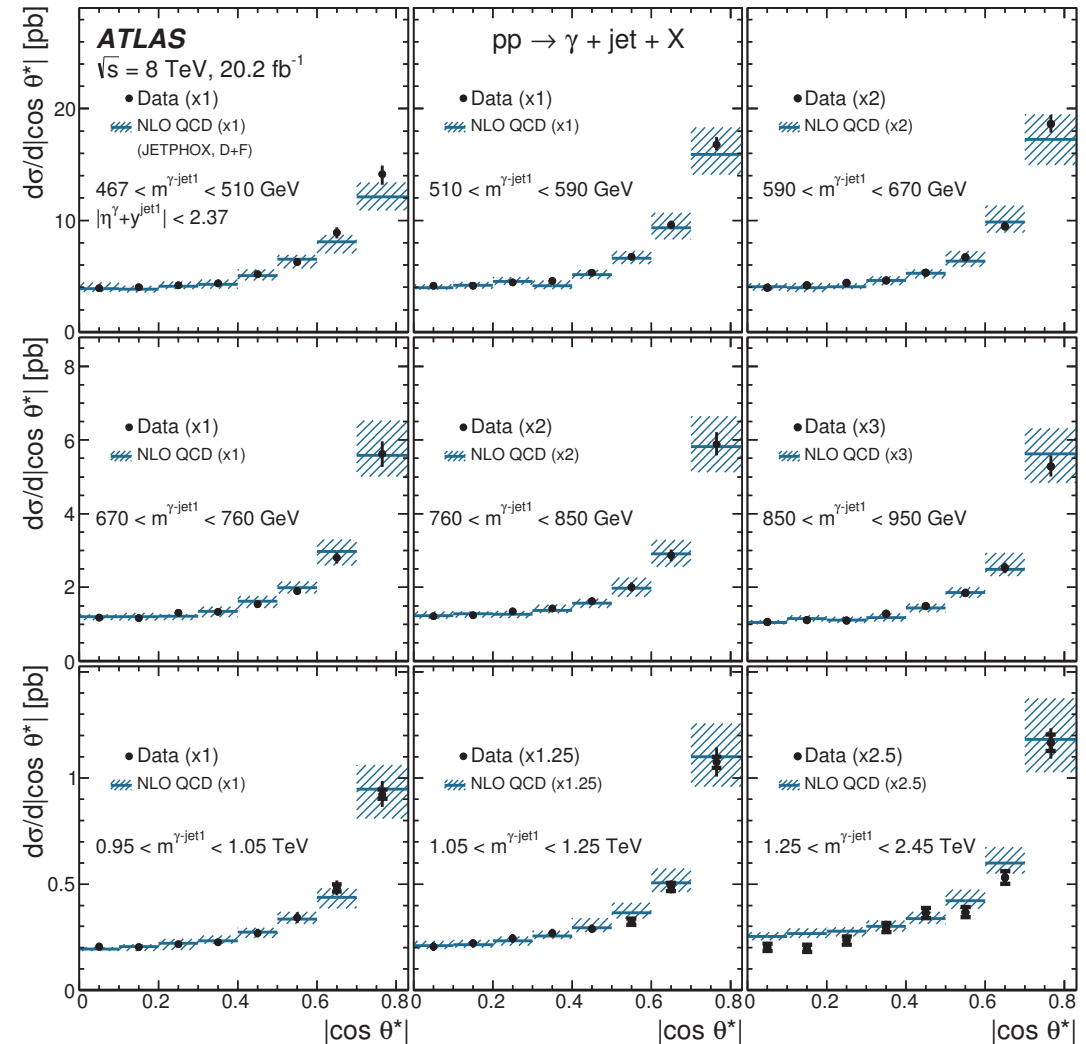
direct-photon process

$$\frac{d\sigma}{d|\cos\theta^{\gamma j}|} \sim (1 - |\cos\theta^{\gamma j}|)^{-1}$$



fragmentation process

$$\frac{d\sigma}{d|\cos\theta^{\gamma j}|} \sim (1 - |\cos\theta^{\gamma j}|)^{-2}$$

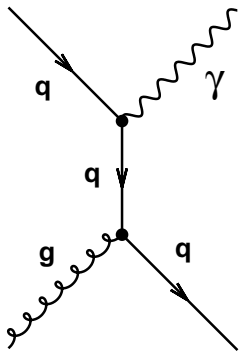


- Measured angular distribution in regions of photon-jet invariant mass  $\Rightarrow$  good description of the data by NLO QCD in shape and normalisation

# Dynamics of $\gamma + \text{jet}$ production in $pp$ collisions

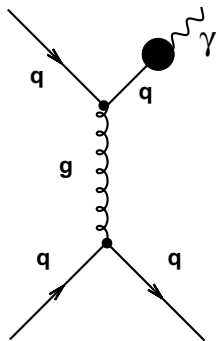
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- Angular distribution  $d\sigma/d|\cos\theta^{\gamma j}|$  sensitive to the spin of the exchanged (virtual) particle: quark(1/2) vs gluon(1)



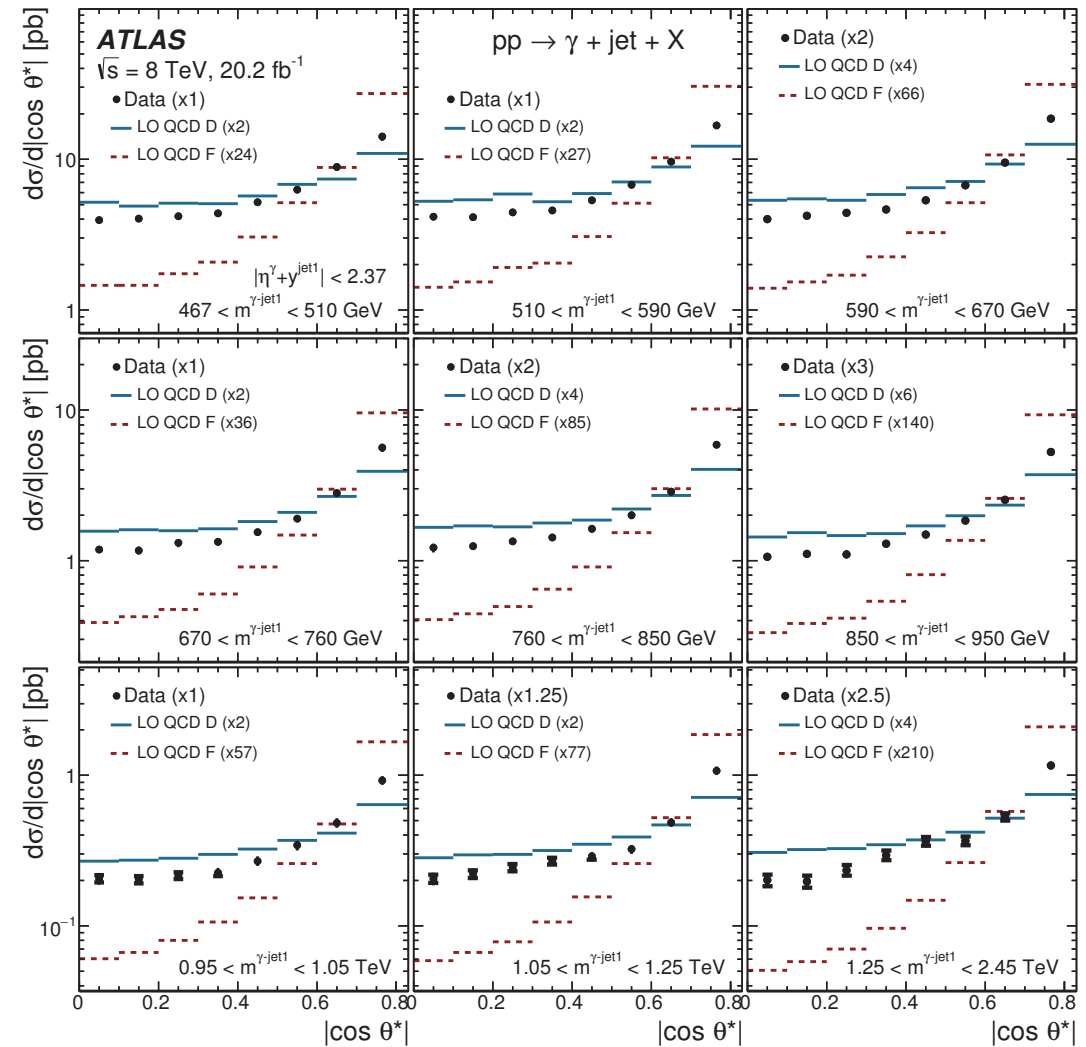
direct-photon process

$$d\sigma/d|\cos\theta^{\gamma j}| \sim (1 - |\cos\theta^{\gamma j}|)^{-1}$$



fragmentation process

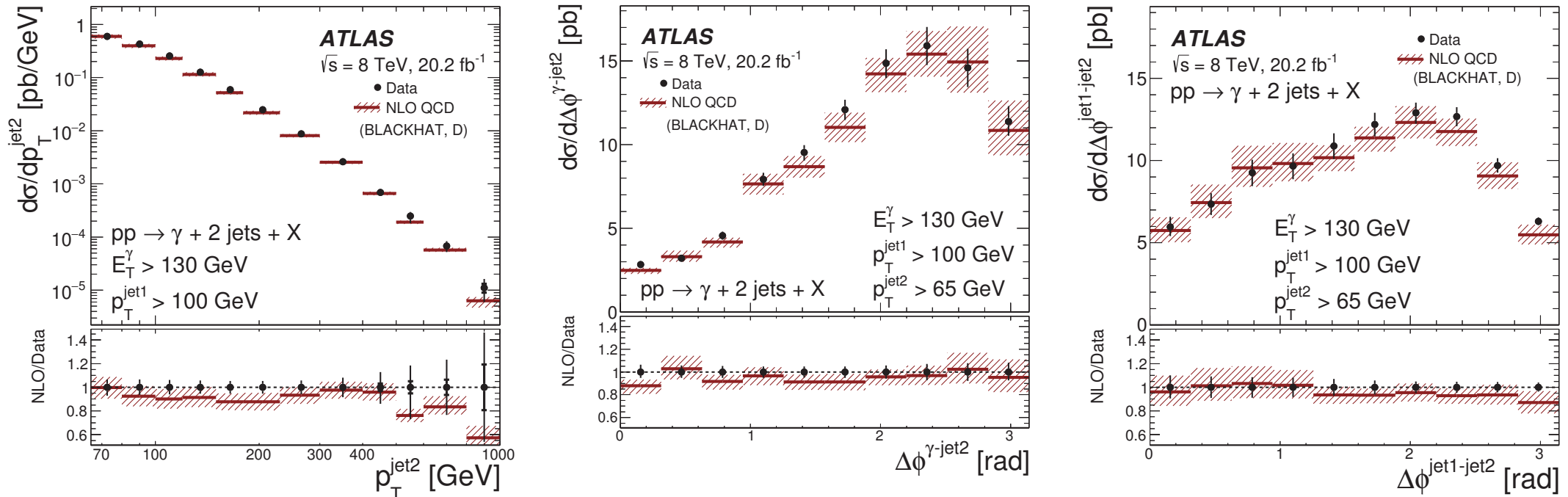
$$d\sigma/d|\cos\theta^{\gamma j}| \sim (1 - |\cos\theta^{\gamma j}|)^{-2}$$



- Measured angular distribution closer to that of direct-photon processes than fragm.
- $\Rightarrow$  consistent with the dominance of processes in which a virtual quark is exchanged

# Dynamics of $\gamma + 2\text{jet}$ production in $pp$ collisions

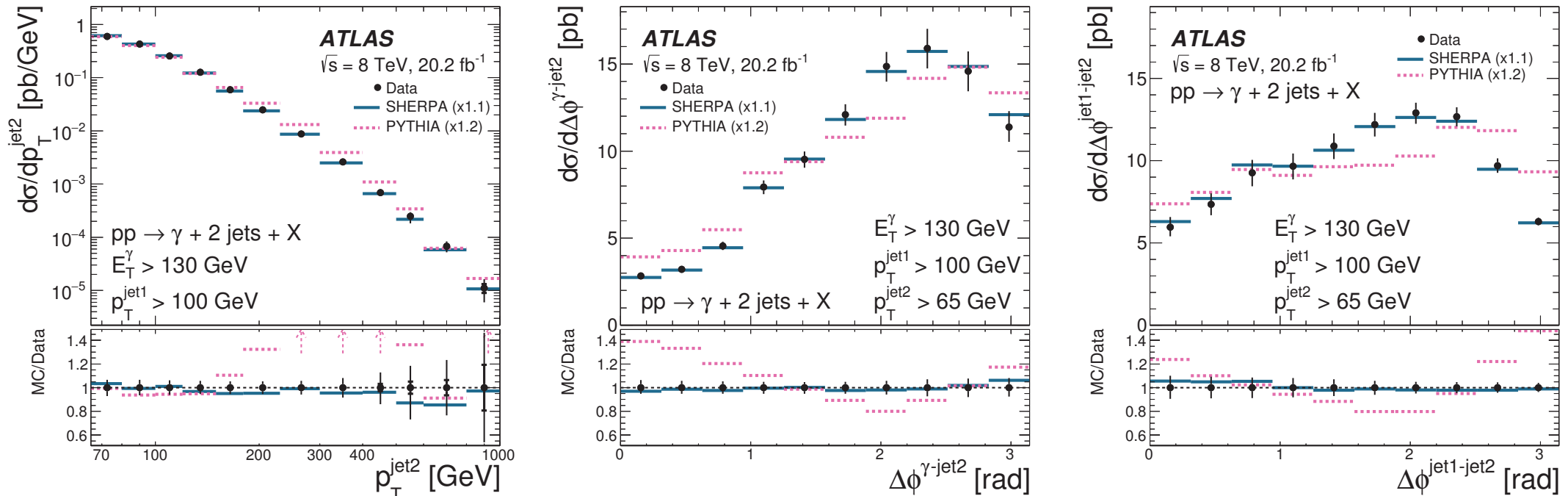
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- **First measurement of  $\gamma + 2\text{jet}$  production in  $pp$  collisions at  $\sqrt{s} = 8 \text{ TeV}$ :**  
 $E_T^\gamma > 130 \text{ GeV}$ ,  $p_T^{\text{jet}1} > 100 \text{ GeV}$  and  $p_T^{\text{jet}2} > 65 \text{ GeV}$
- **Measurement of  $d\sigma/dp_T^{\text{jet}2}$  and angular correlations between the photon and the jets**  
 →  $\Delta\phi$  between the photon and subleading jet ( $\Delta\phi^{\gamma\text{-jet}2}$ )  
 →  $\Delta\phi$  between the leading and subleading jets ( $\Delta\phi^{\text{jet}1\text{-jet}2}$ )
- **Good description of the data both in shape and normalisation by the NLO QCD predictions computed with Blackhat**

# Dynamics of $\gamma + 2\text{jet}$ production in $pp$ collisions

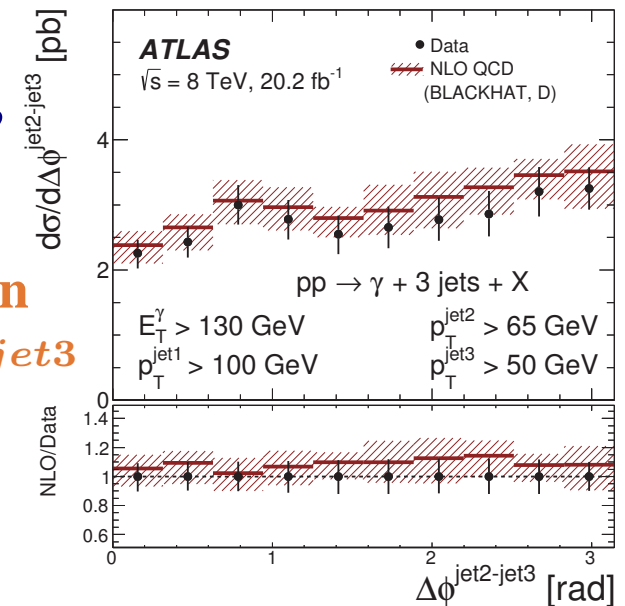
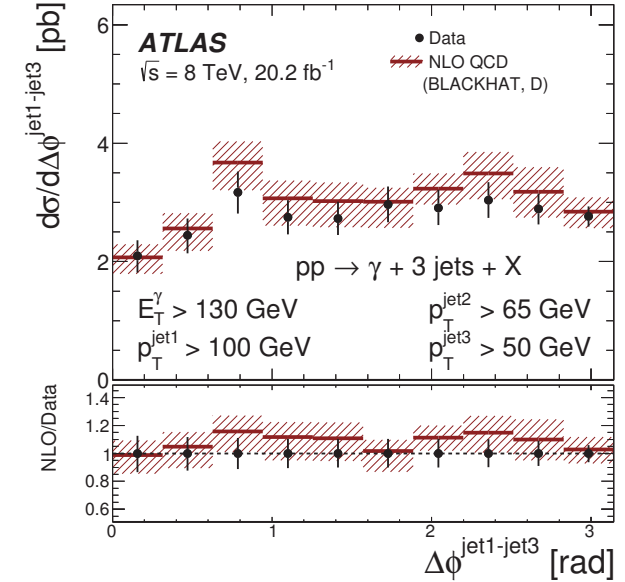
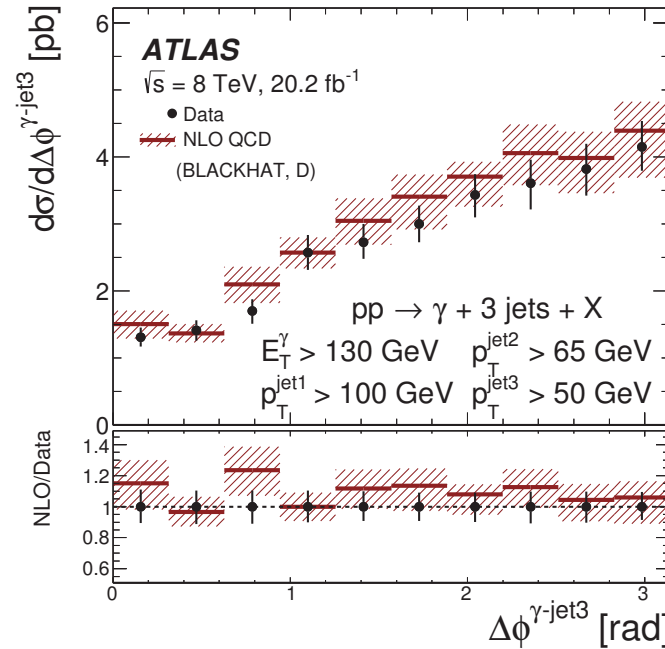
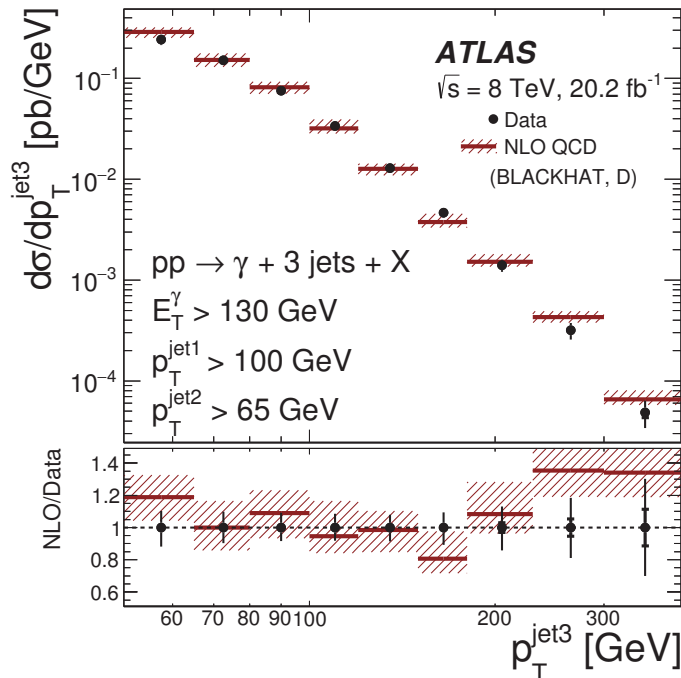
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- **Comparison to the predictions of Monte Carlo generators:**
  - **PYTHIA:** 2 → 2 matrix elements plus parton showers
  - **SHERPA:** 2 →  $n$  ( $n = 2, \dots, 5$ ) matrix elements plus parton showers
- **MC predictions normalised to data: shape comparison only**
- **Good description of the data by the SHERPA predictions while PYTHIA fails to describe the distribution in  $p_T^{\text{jet2}}$  and the angular correlations**
  - ⇒ **The inclusion of higher-order tree-level ME makes SHERPA superior!**

# Dynamics of $\gamma + 3\text{jet}$ production in $pp$ collisions

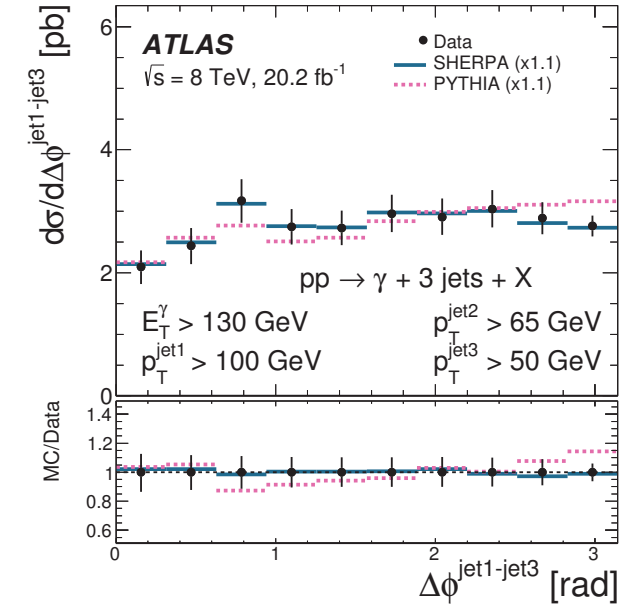
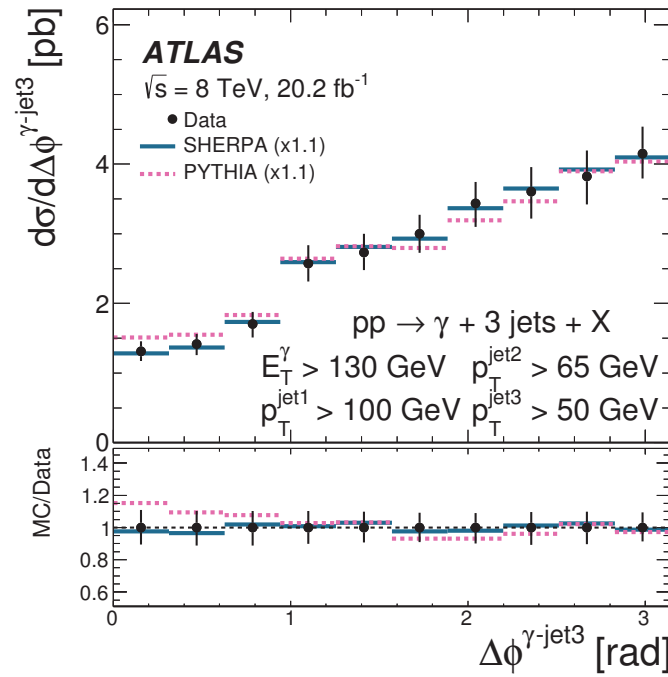
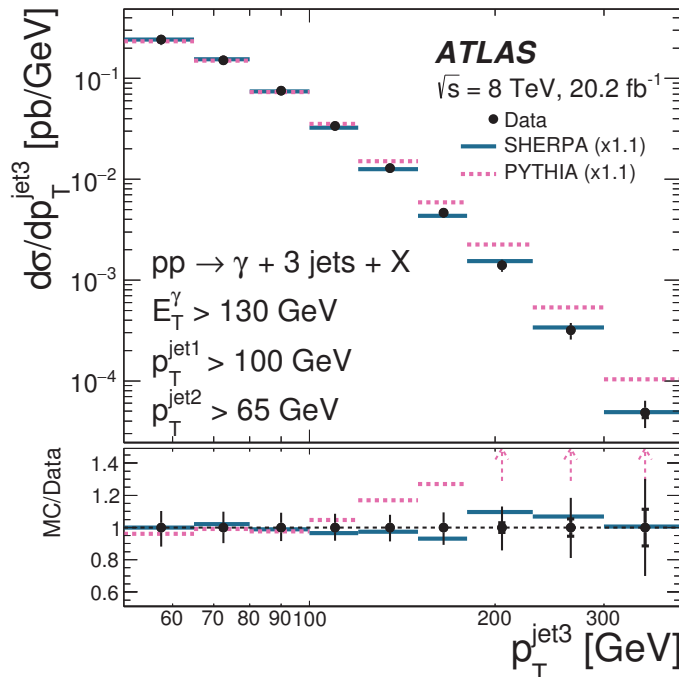
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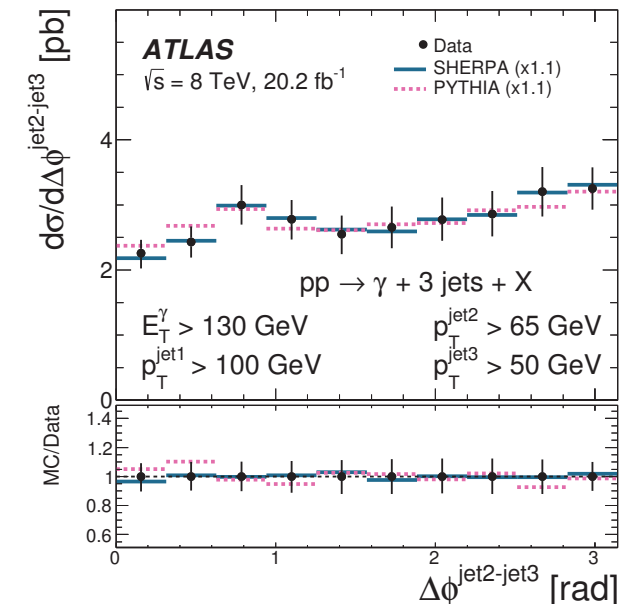
- First measurement of  $\gamma + 3\text{jet}$  at the LHC:  $E_T^\gamma > 130 \text{ GeV}$ ,  $p_T^{\text{jet1}} > 100 \text{ GeV}$ ,  $p_T^{\text{jet2}} > 65 \text{ GeV}$  and  $p_T^{\text{jet3}} > 50 \text{ GeV}$
- Measurement of  $d\sigma/dp_T^{\text{jet3}}$  and angular correlations between the photon and the jets:  $\Delta\phi^{\gamma\text{-jet3}}$ ,  $\Delta\phi^{\text{jet1-jet3}}$ ,  $\Delta\phi^{\text{jet2-jet3}}$
- Adequate description of the data by the NLO QCD predictions computed with Blackhat

# Dynamics of $\gamma + 3\text{jet}$ production in $pp$ collisions

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- Comparison to the predictions of Monte Carlo generators of **PYTHIA (2 → 2 ME+PS)** and **SHERPA (2 → n ME +PS)** normalised to data (shape comparison)
  - Good description of the data by the **SHERPA** predictions while **PYTHIA** fails to describe the distribution in  $p_T^{\text{jet3}}$
- ⇒ The inclusion of higher-order tree-level ME makes **SHERPA** superior to **PYTHIA**



# Photon pair production at 8 TeV



## Isolated-photon pair production in $pp$ collisions

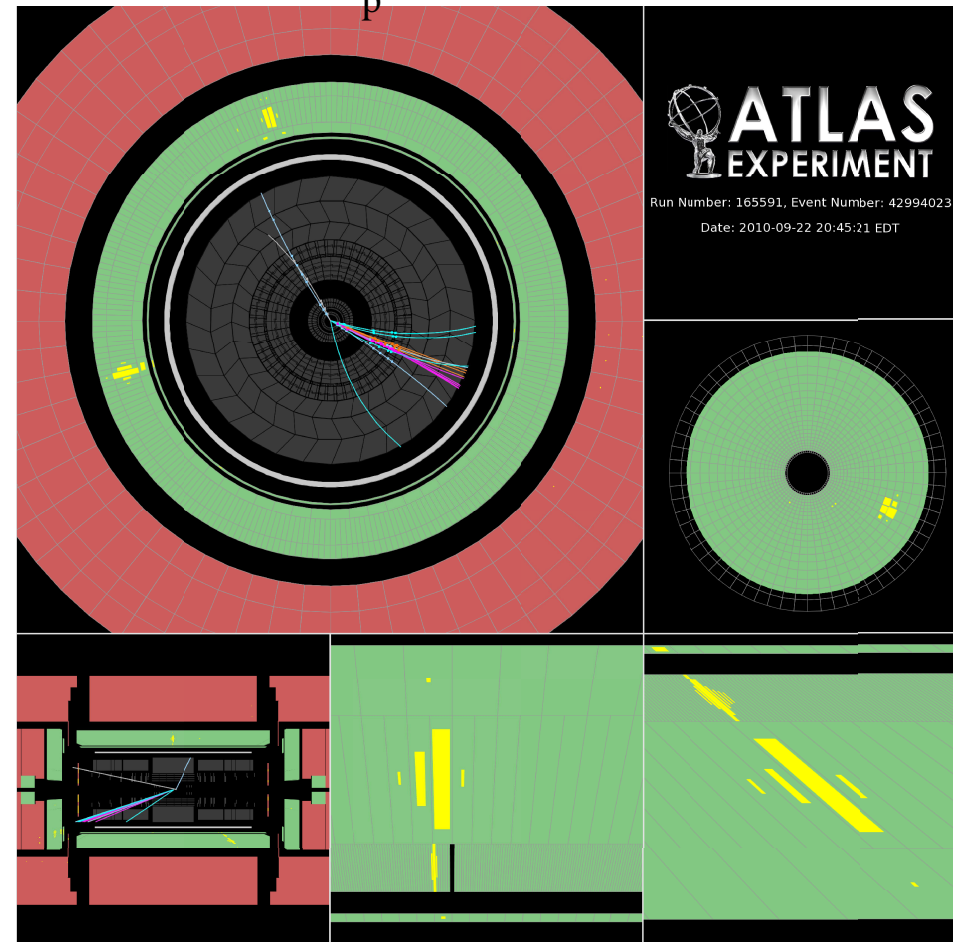
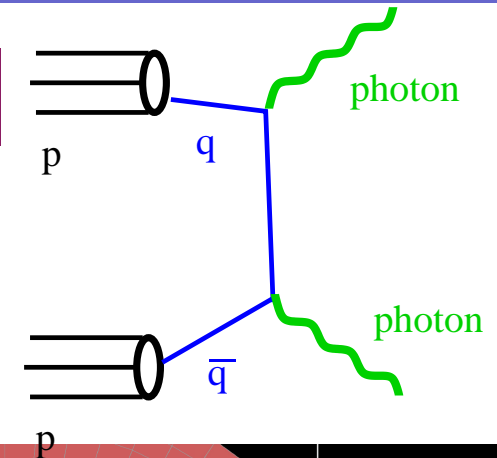
- Measurements of the process  $pp \rightarrow \gamma\gamma + X$  with the aim of testing pQCD and understanding the irreducible background to new physics processes involving photons or  $H \rightarrow \gamma\gamma$

- Measurement of differential cross sections

- diphoton invariant mass,  $m_{\gamma\gamma}$
- diphoton transverse momentum,  $p_{T,\gamma\gamma}$
- azimuthal separation in LAB frame,  $\Delta\phi_{\gamma\gamma}$
- $\cos\theta_\eta^*$  →  $\phi_\eta^* \equiv \tan\left(\frac{\pi - \Delta\phi_{\gamma\gamma}}{2}\right) \sin\theta_\eta^*$
- transverse component of  $\vec{p}_{T,\gamma\gamma}$  with respect to thrust axis ( $a_T$ )

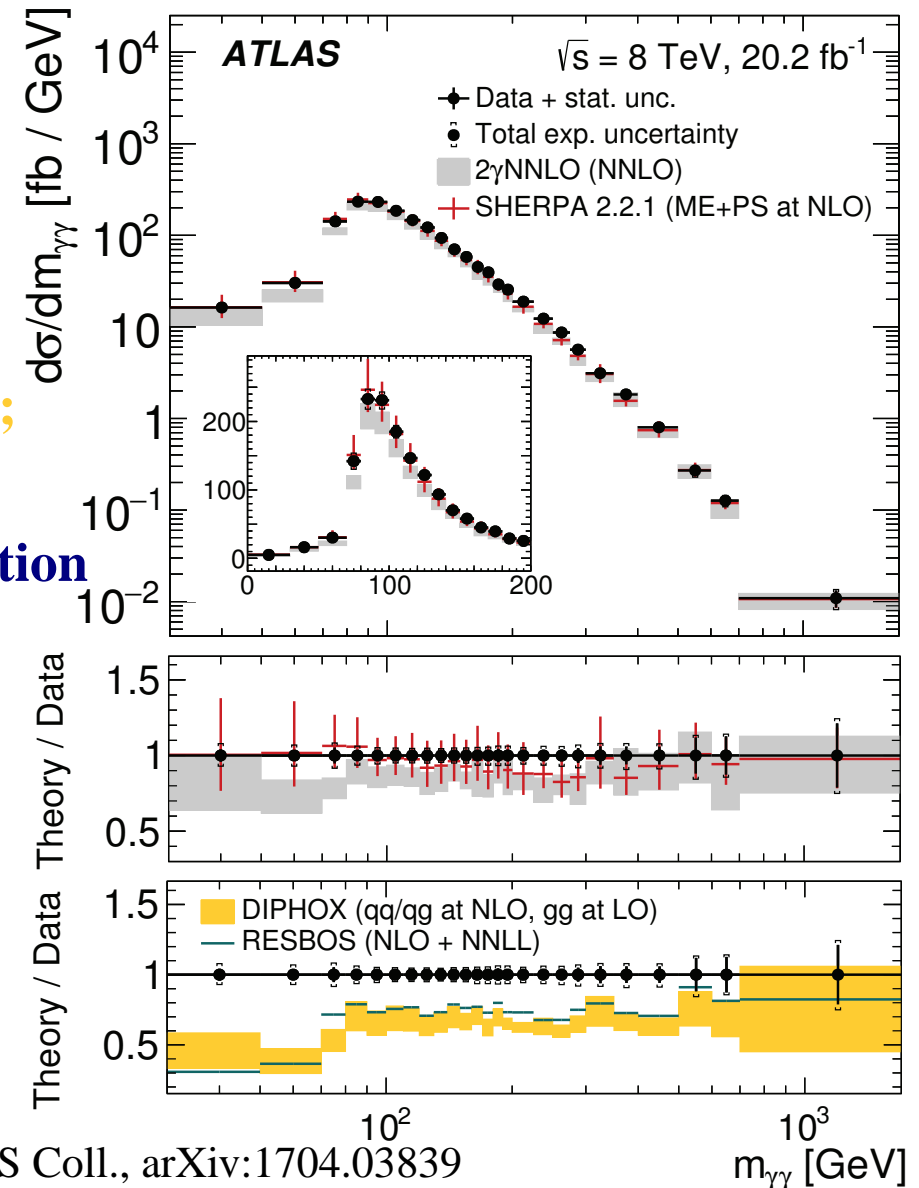
in the phase-space region defined by:

$E_T^{\gamma 1,2} > 40(30) \text{ GeV}$ ,  $|\eta^\gamma| < 2.37$  (excluding the region  $1.37 < |\eta^\gamma| < 1.56$ ),  $\Delta R_{\gamma\gamma} > 0.4$  and  $E_T^{iso} < 11 \text{ GeV}$  using  $\mathcal{L} = 20.2 \text{ fb}^{-1}$



# Isolated-photon pair production in $pp$ collisions at $\sqrt{s} = 8$ TeV

- Comparison to theoretical calculations
- Fixed-order QCD calculations (NP corrected)
  - $2\gamma$  NNLO program; NNLO calculation of direct-photon contribution (no fragm.)
  - **DIPHOX** program; NLO calculation of direct-photon and fragmentation contributions; box diagram  $gg \rightarrow \gamma\gamma$  included
  - **RESBOS** program; NLO plus NNLL resummation
- New **SHERPA (v2.2.1)** calculation combining
  - $\gamma\gamma$  and  $\gamma\gamma + 1p$  at NLO
  - $\gamma\gamma + 2p$  and  $\gamma\gamma + 3p$  at LO
  - **parton showers**
- The small contribution from  $H \rightarrow \gamma\gamma$  is neglected
  - ⇒ **SHERPA prediction in agreement with data**

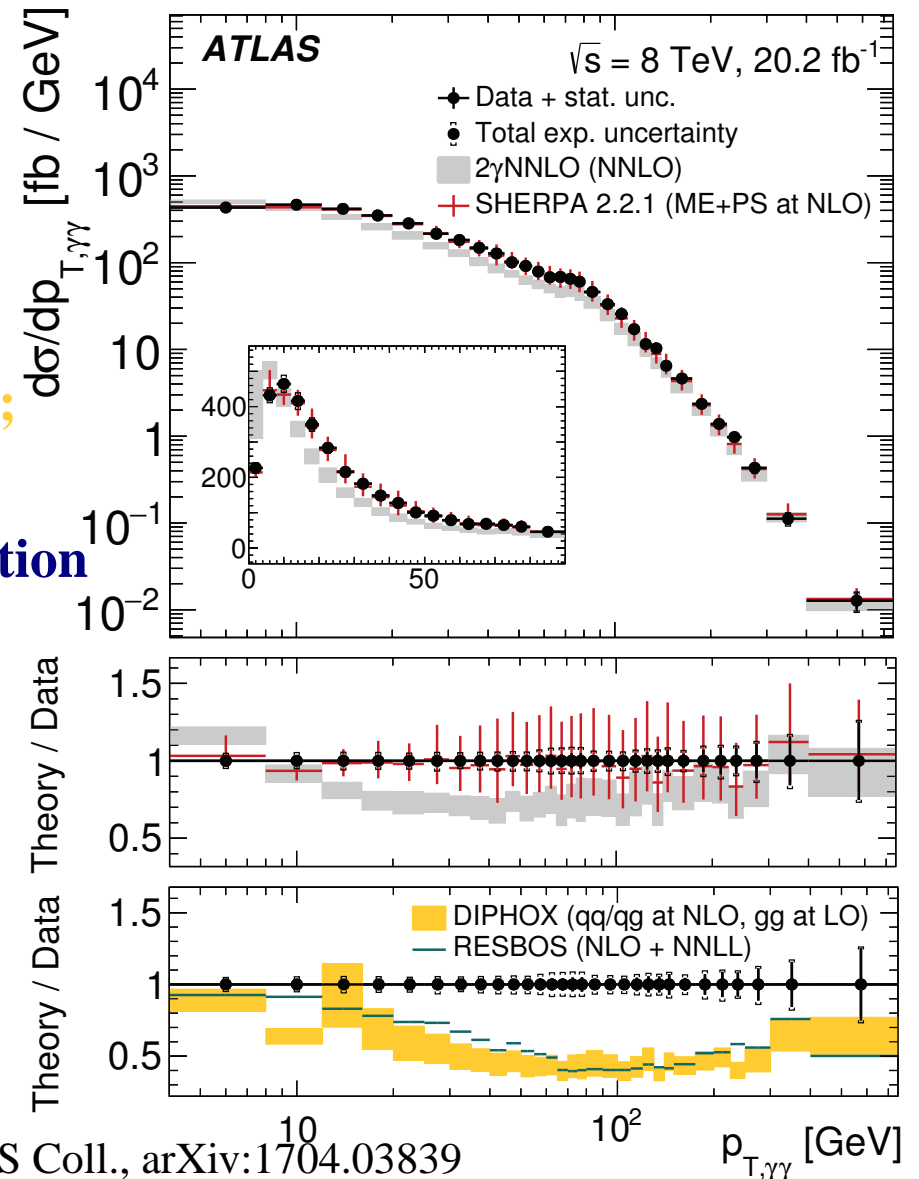


ATLAS Coll., arXiv:1704.03839

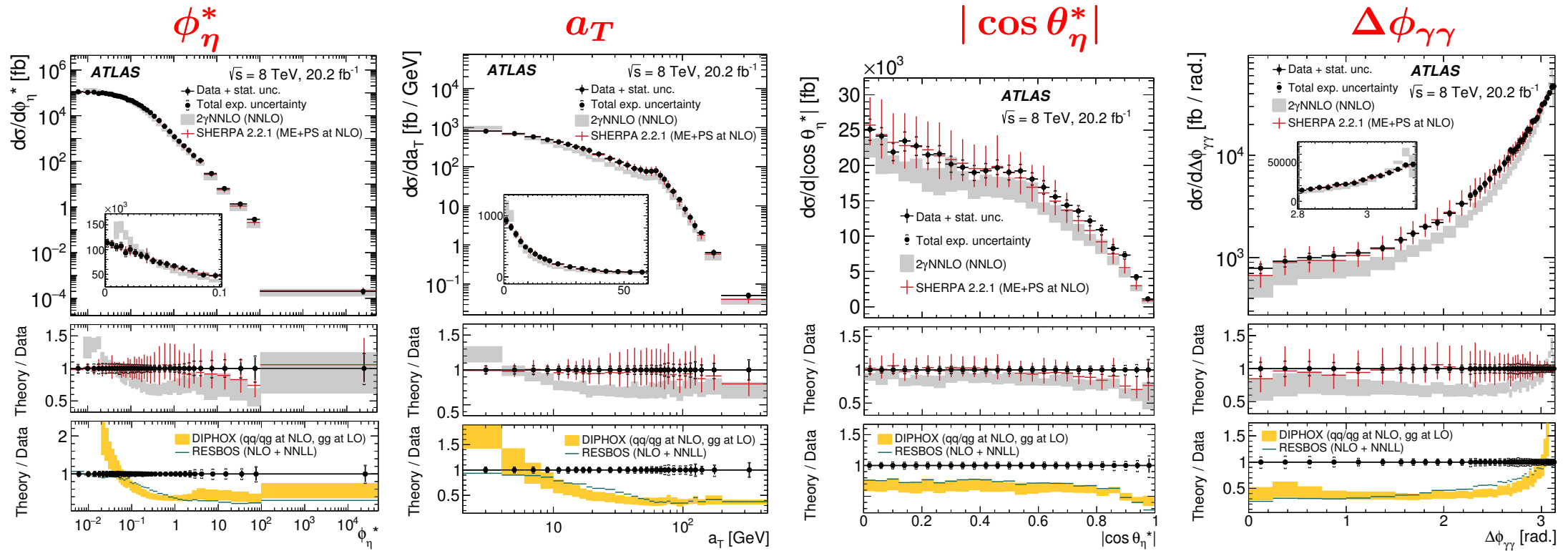
 $m_{\gamma\gamma}$  [GeV]

# Isolated-photon pair production in $pp$ collisions at $\sqrt{s} = 8$ TeV

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# Isolated-photon pair production in $pp$ collisions at $\sqrt{s} = 8$ TeV

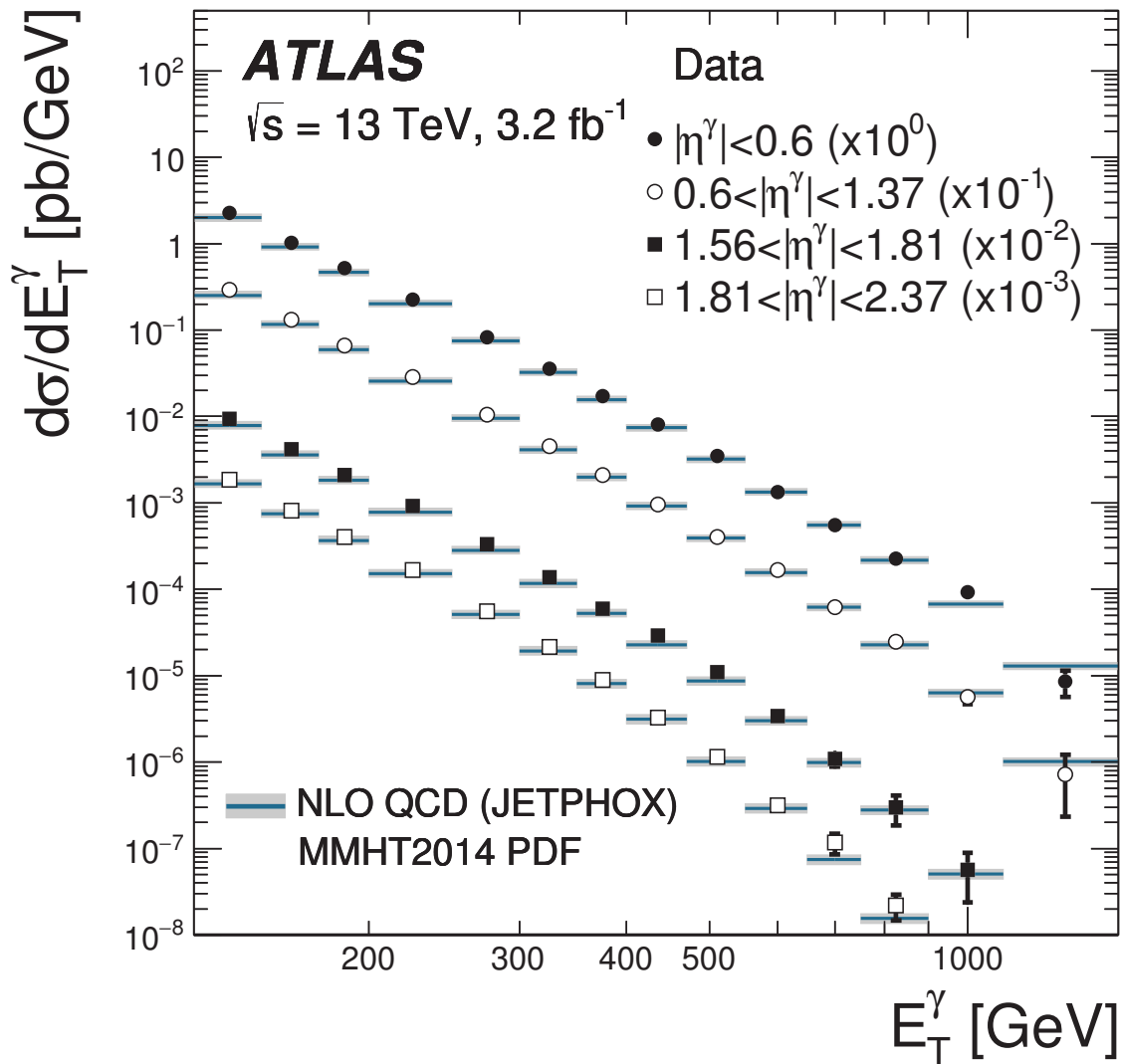


ATLAS Coll., arXiv:1704.03839

- $\Delta\phi_{\gamma\gamma} \sim \pi$  or at low values of  $p_{T,\gamma\gamma}$ ,  $a_T$  and  $\phi_\eta^*$  (soft gluon resummation important): **RESBOS** and **SHERPA** do well
- NLO QCD calculations without higher order terms (DIPHOX, RESBOS) are insufficient
- NNLO corrections (2 $\gamma$ NNLO) improve the description, but still insufficient
- **SHERPA** predictions agree with the data

# Inclusive photon production at 13 TeV

# Inclusive isolated-photon production in $pp$ collisions at $\sqrt{s} = 13$ TeV



- Measurement of  $d\sigma/dE_T^\gamma$  in different ranges in  $\eta^\gamma$  for  $125 < E_T^\gamma < 1500$  GeV using  $\mathcal{L} = 3.2 \text{ fb}^{-1}$  of  $pp$  collision data at  $\sqrt{s} = 13$  TeV

- Isolation:  $E_T^{\text{iso}} < 4.2 \cdot 10^{-3} \cdot E_T^\gamma + 4.8$  GeV

- The measurement covers more than five orders of magnitude in cross section

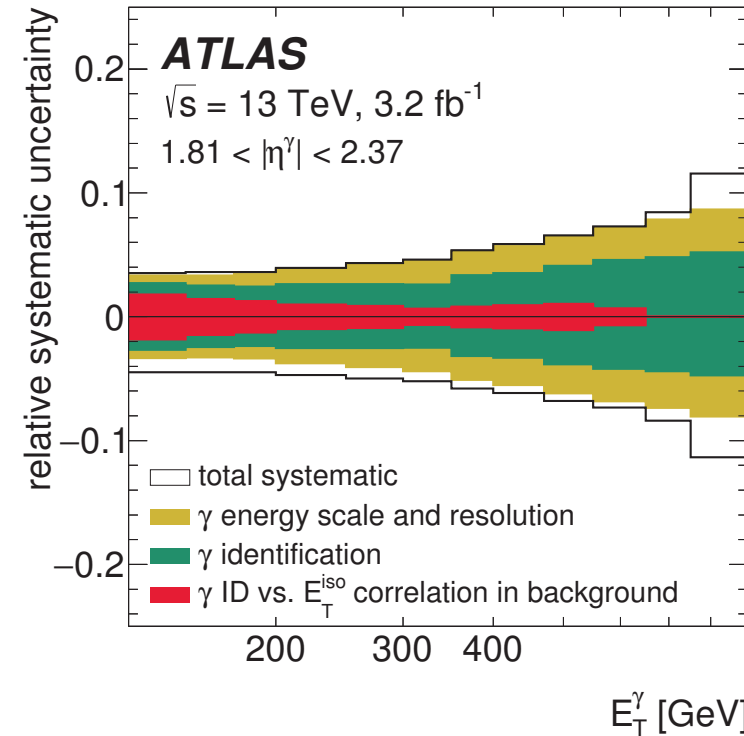
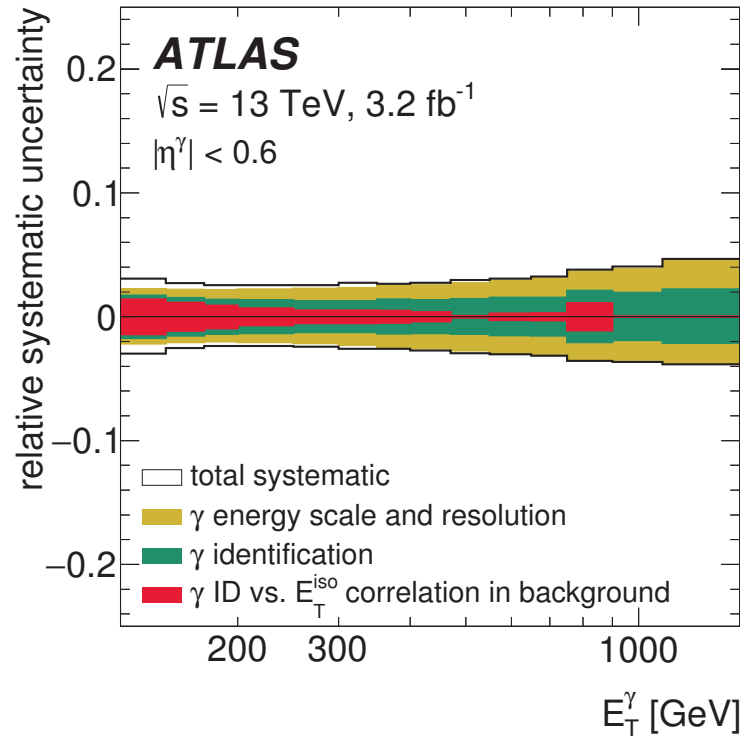
- $d\sigma/dE_T^\gamma$  increases by a factor 2 (10) at  $E_T^\gamma = 125$  (1000) GeV with respect to at  $\sqrt{s} = 8$  TeV

- Good description (in log scale) of the data by NLO QCD predictions computed with JetPhox using the MMHT2014 PDFs

ATLAS Coll., arXiv:1701.06882, accepted PLB

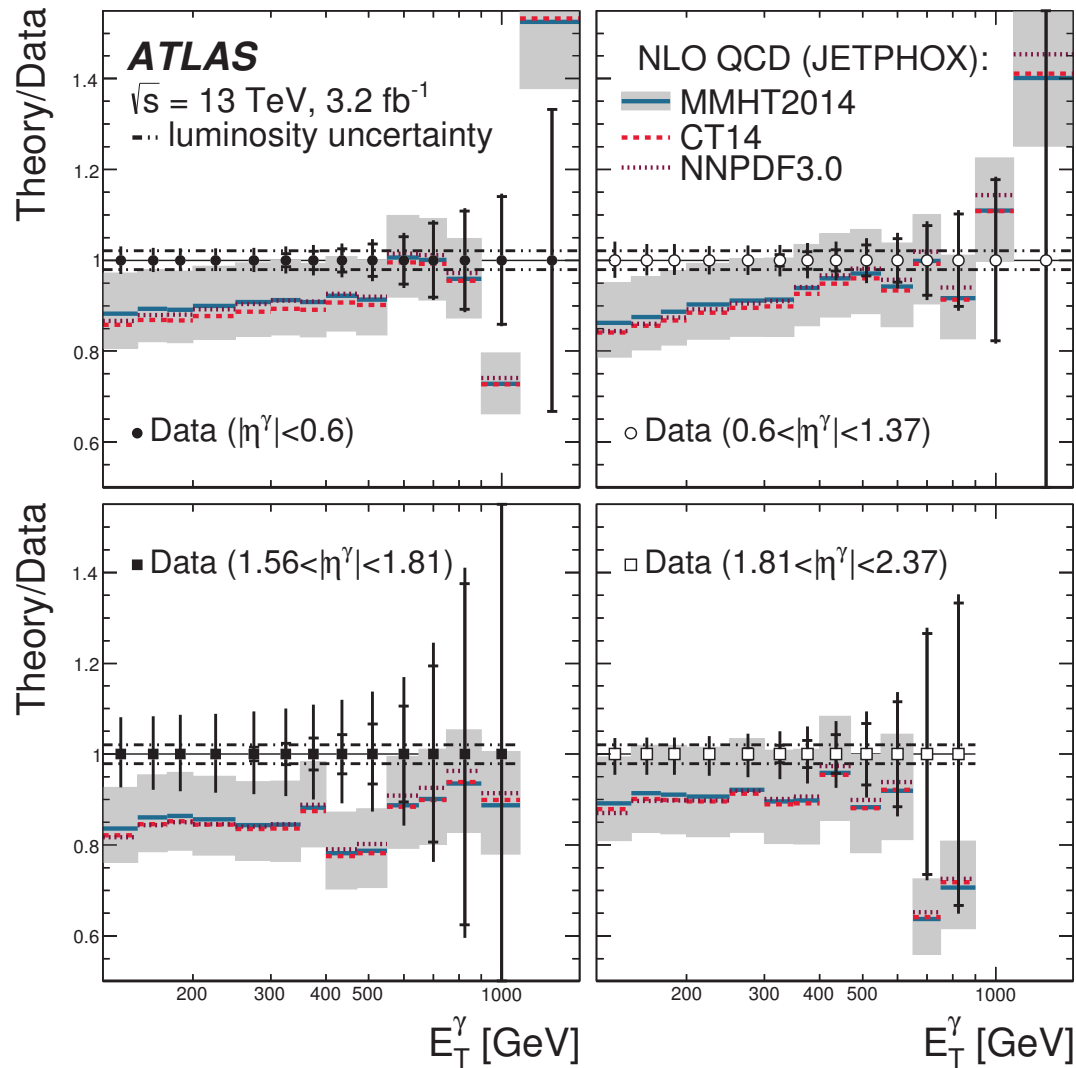
## Major experimental uncertainties

ATLAS Coll., arXiv:1701.06882, accepted PLB



- The uncertainty on the photon energy scale dominates at high  $E_T^\gamma$ : 2-5% except for  $1.56 < |\eta^\gamma| < 1.81$ , where it is 7-18%
- The uncertainty in the photon identification represents a significant contribution at low  $E_T^\gamma$ : it increases from 1-2% at 125 GeV to 2-6% at  $\sim 1 \text{ TeV}$
- The uncertainty in the correlation between the photon ID variables and the isolation is a significant contribution at low  $E_T^\gamma$ : typically smaller than 2%

# Inclusive isolated-photon cross sections vs NLO QCD



- NLO QCD predictions underestimate data by up to  $\approx 10\text{-}15\%$

- Theoretical uncertainty 10-15% much larger than experimental uncertainties

- For  $E_T^\gamma \lesssim 600 \text{ GeV}$  the measurements are systematically limited

- NLO QCD provides an adequate description of the data within uncertainties

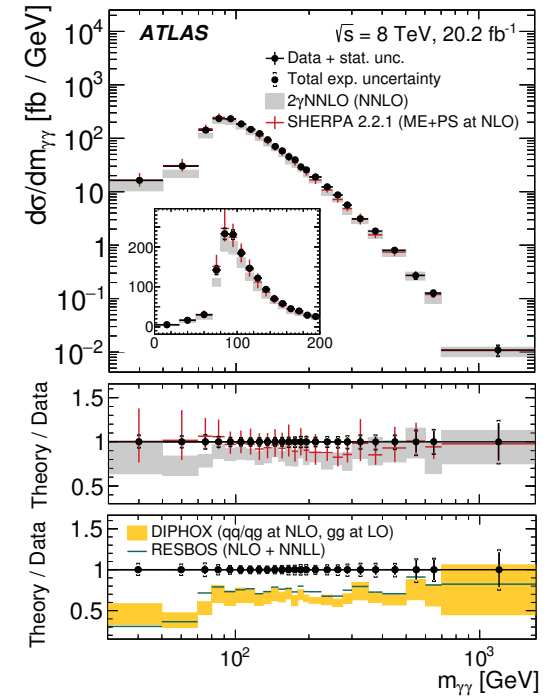
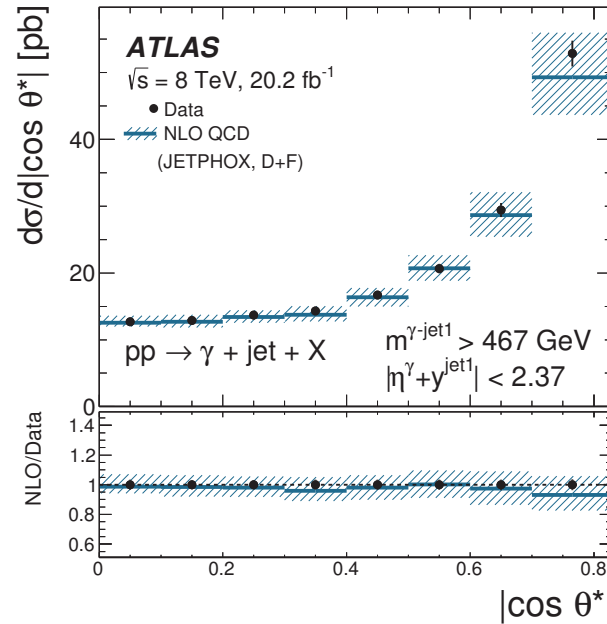
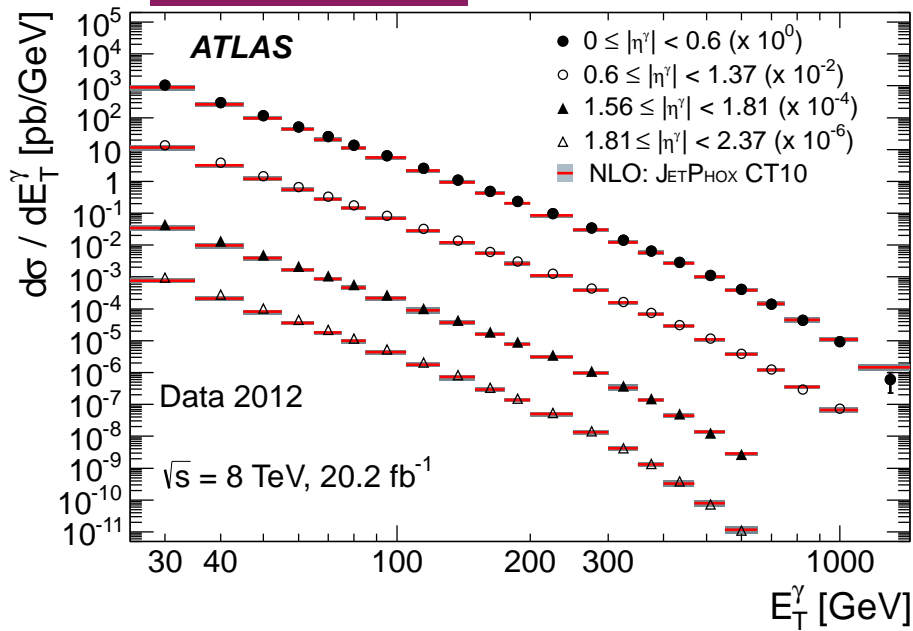
- First measurement of inclusive photon production in the new kinematic regime opened by the LHC at  $\sqrt{s} = 13 \text{ TeV}$

- Ready for the comparison to NNLO QCD predictions (Campbell, Ellis, Williams arXiv:1612.04333)

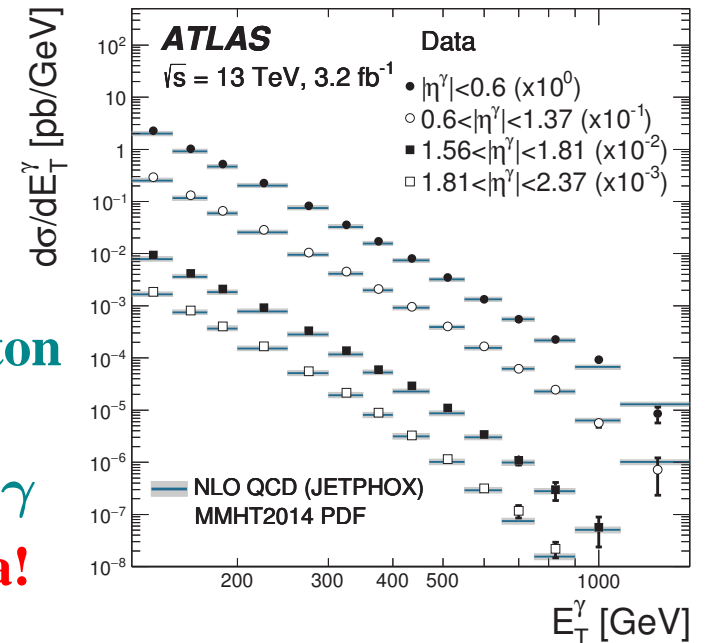
ATLAS Coll., arXiv:1701.06882, accepted PLB



# Summary

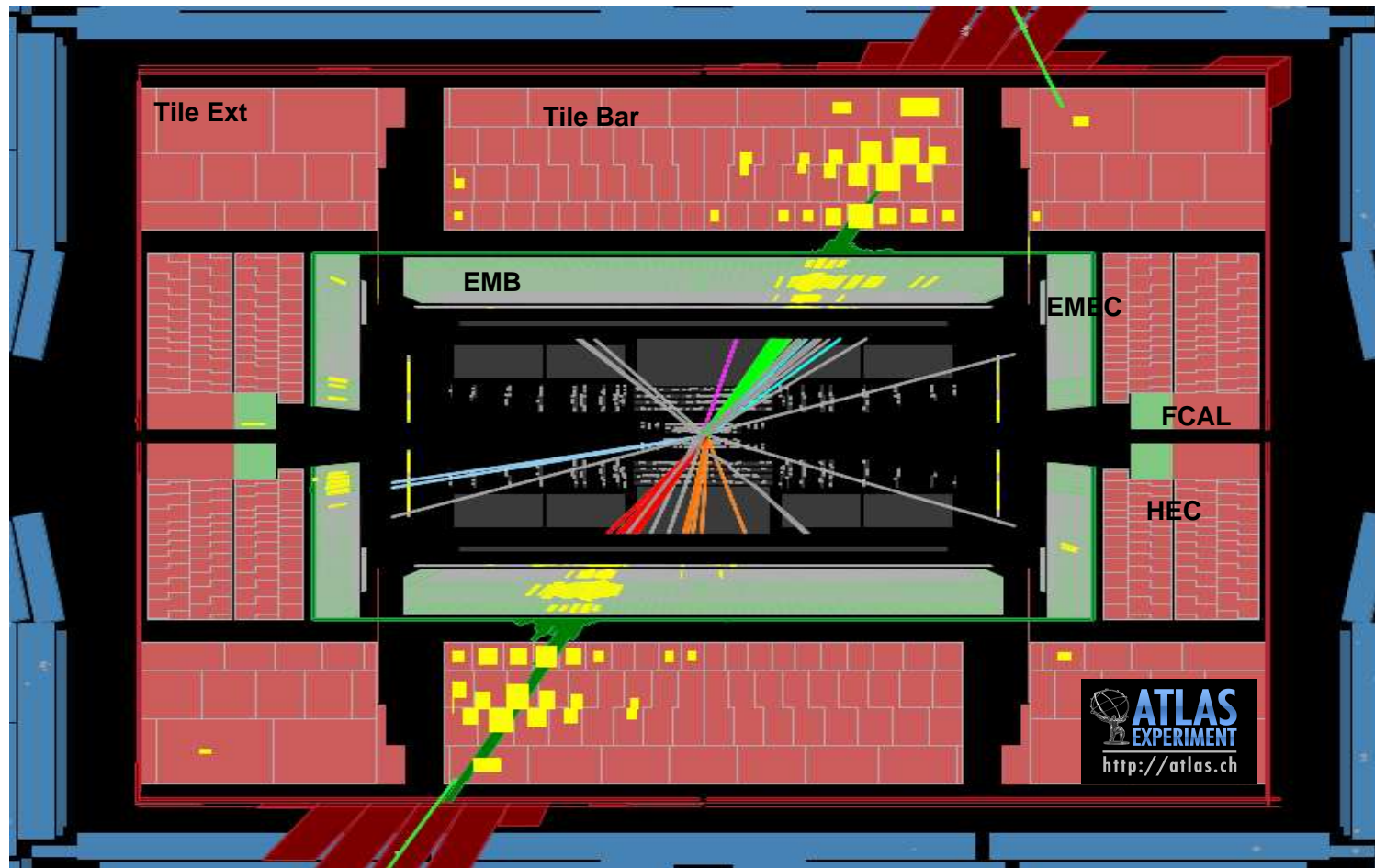


- Exploration of isolated photon production in  $pp$  collisions up to  $E_T^\gamma \sim 1 \text{ TeV}$
- Additional experimental information on the gluon density in the proton
- Measurement of the dynamics of photon+jet(s) and diphoton production
- Understanding (in pQCD) the background to Higgs into  $\gamma\gamma$
- Overall, perturbative QCD succeeds in describing the data!



**Backup**

## The ATLAS detector



- Inner detector (ID): tracking and particle identification in  $|\eta| < 2.5$
- Calorimeters: electromagnetic (LAr)  $\rightarrow$  barrel  $|\eta| < 1.475$ , endcap  $1.375 < |\eta| < 3.2$ , forward  $3.1 < |\eta| < 4.9$ ; hadronic (scintillator/steel, LAr/Cu, LAr/W)  $\rightarrow$  barrel  $|\eta| < 0.7$  extended barrel  $0.8 < |\eta| < 1.7$ , endcap  $1.5 < |\eta| < 3.2$  and forward  $3.1 < |\eta| < 4.9$

## Photon reconstruction in the ATLAS LAr Calorimeter

### ● Layout of the ATLAS electromagnetic calorimeter (Lead-liquid Argon)

- barrel section,  $|\eta| < 1.475$
- two end-cap sections,  $1.375 < |\eta| < 3.2$
- three longitudinal layers

- First layer: high granularity in  $\eta$  direction, width 0.003-0.006 (except for  $1.4 < |\eta| < 1.5$  and  $|\eta| > 2.4$ )

- Second layer: collects most of the energy, granularity  $0.025 \times 0.025$  in  $\eta \times \phi$

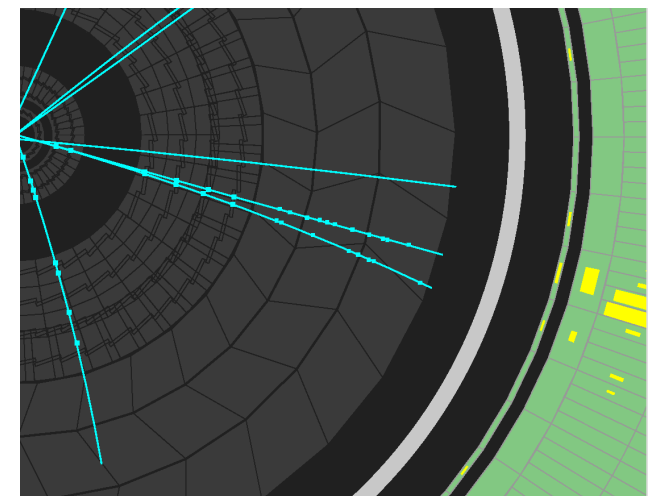
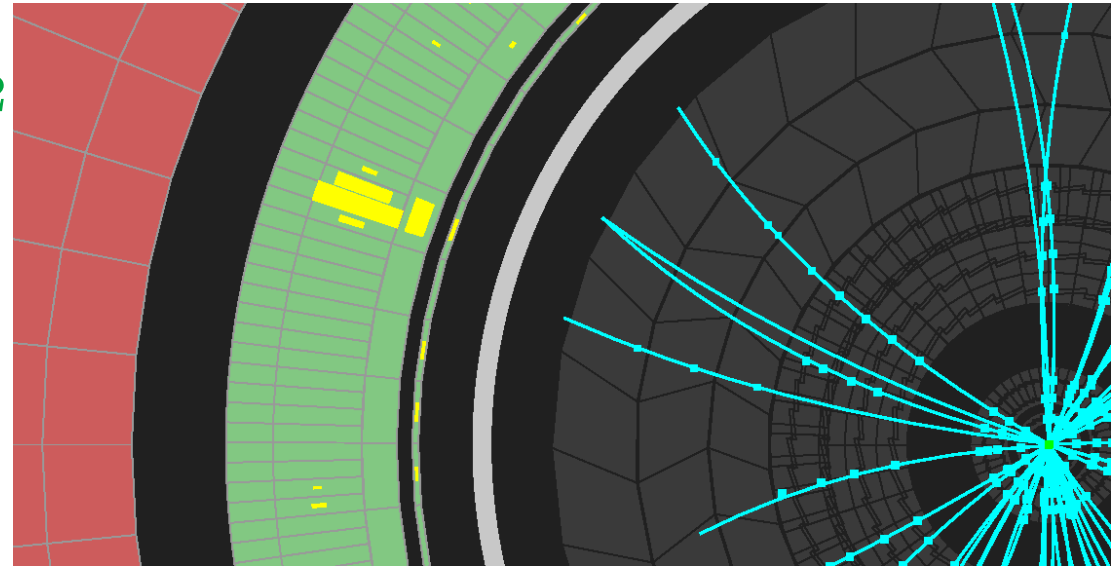
- Third layer: used to correct for leakage

### ● Cluster of EM cells without matching track:

→ “unconverted” photon candidate

### ● Cluster of EM cells matched to pairs of tracks (from reconstructed conversion vertices in the inner detector) or matched to a single track consistent with originating from a photon conversion

→ “converted” photon candidate



## Photon identification in the ATLAS LAr Calorimeter

● **To discriminate signal vs background:** shape variables from the lateral and longitudinal energy profiles of the shower in the calorimeters; “loose” and “tight” identification criteria.

● **“Loose” identification criteria:**

→ leakage  $R_{had} = E_T^{had} / E_T$  (1st layer hadronic calorimeter)

→  $R_\eta = E_{3 \times 7}^{S2} / E_{7 \times 7}^{S2}$ ; S2=second layer of EM calorimeter

→ RMS width of the shower in  $\eta$  direction in S2

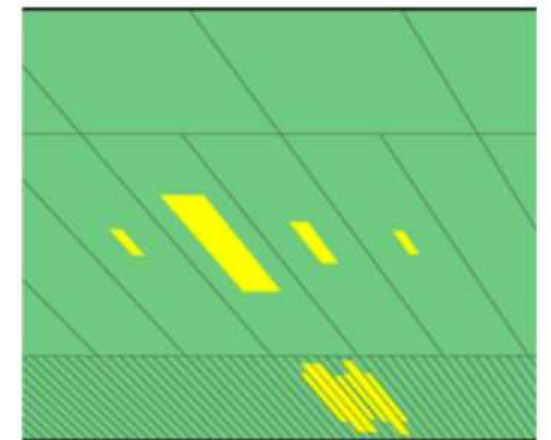
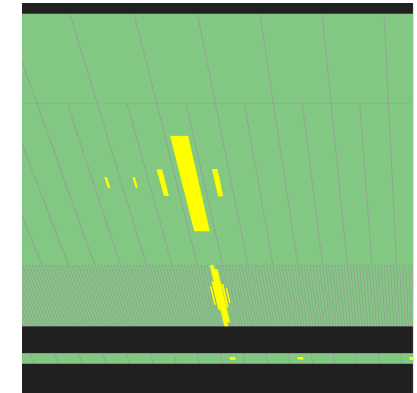
● **“Tight” identification criteria:**

→ the requirements applied in “Loose” are tightened

→  $R_\phi = E_{3 \times 3}^{S2} / E_{3 \times 7}^{S2}$

and shower shapes in the first layer (to discriminate single-photon showers from overlapping nearby showers, such as  $\pi^0 \rightarrow \gamma\gamma$ )

→ e.g. asymmetry between the 1st and 2nd maxima in the energy profile along  $\eta$  (S1)



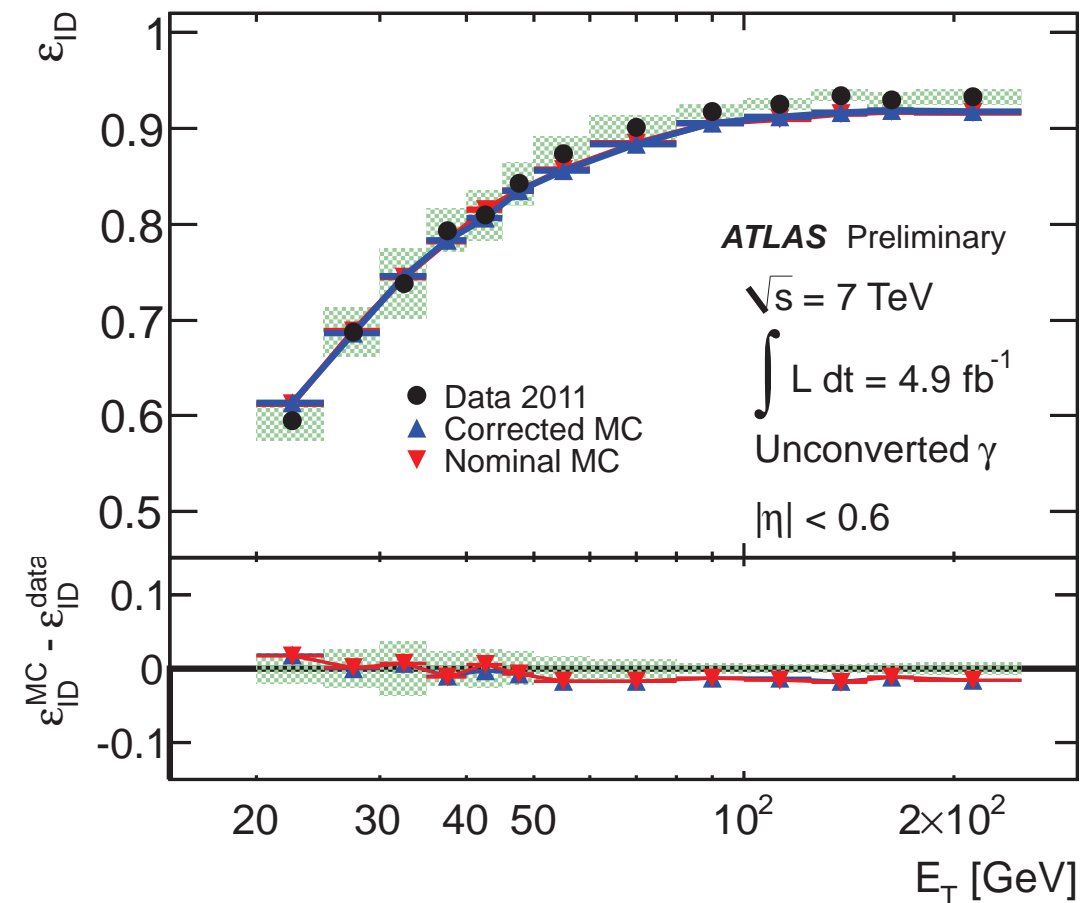
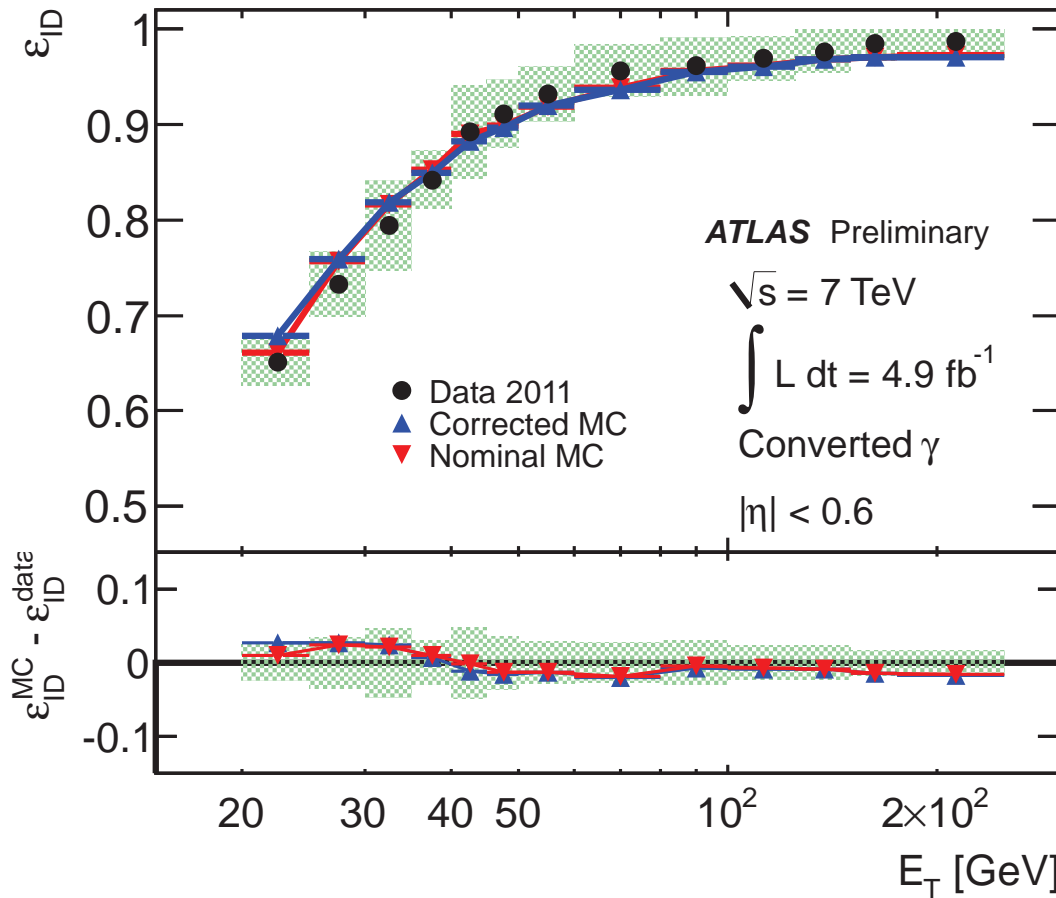
$\pi^0 \rightarrow \gamma\gamma$

# Photon identification efficiency

ATLAS Coll., ATLAS-CONF-2012-123

## CONVERTED PHOTONS

## UNCONVERTED PHOTONS



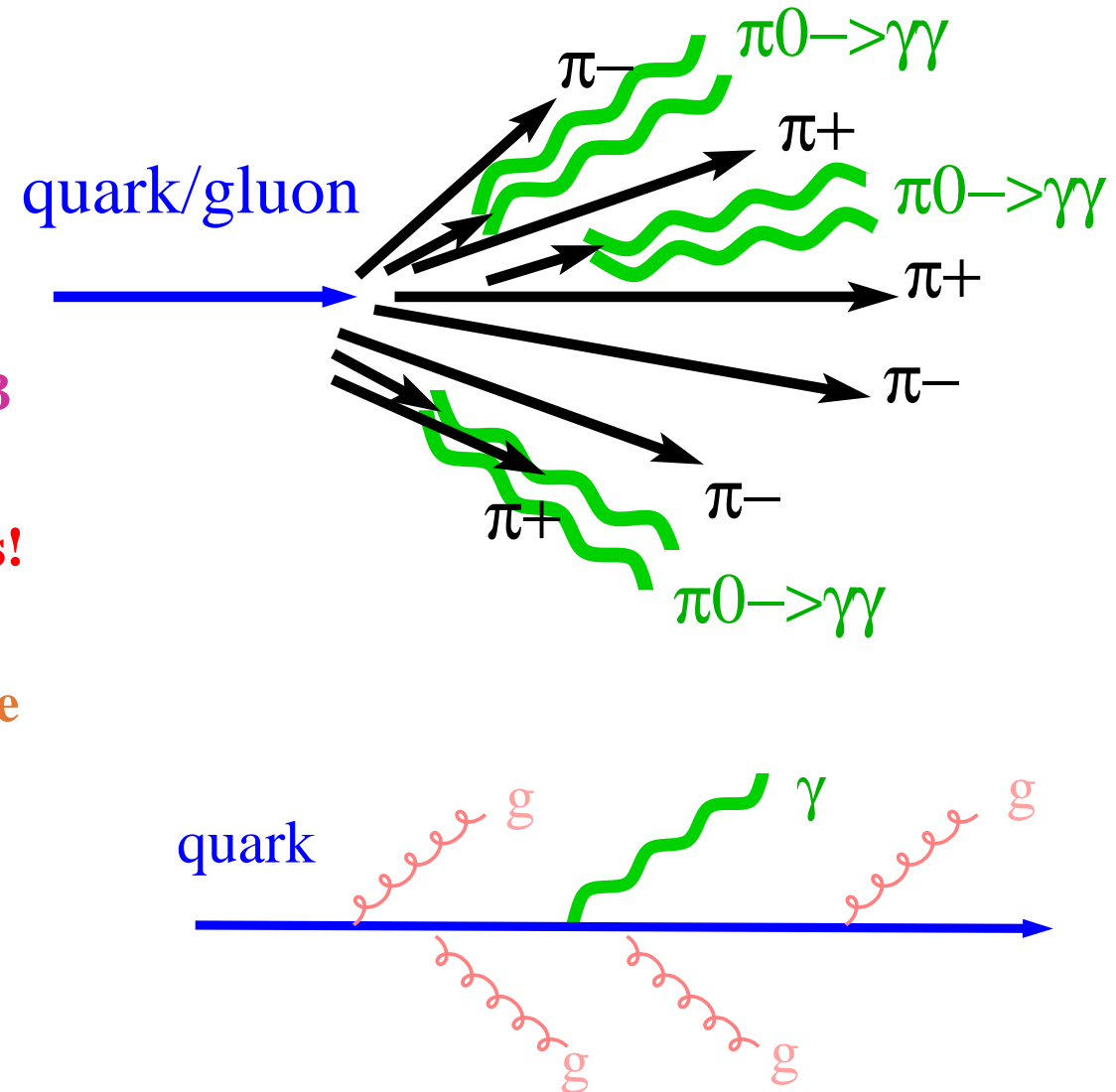
- Data-driven measurements of photon identification efficiency for converted and unconverted photons (radiative  $Z$  decays, extrapolation from  $e^\pm$  and matrix method) compared to estimations based on Monte Carlo simulations**

## Other sources of photons

- Quarks and gluons are sources of photons

→ Quarks and gluons fragment mostly into pions and, by isospin symmetry, 1/3 are  $\pi^0$ 's, which decay into two photons  
 ⇒  $\gamma$ 's are produced copiously inside jets!

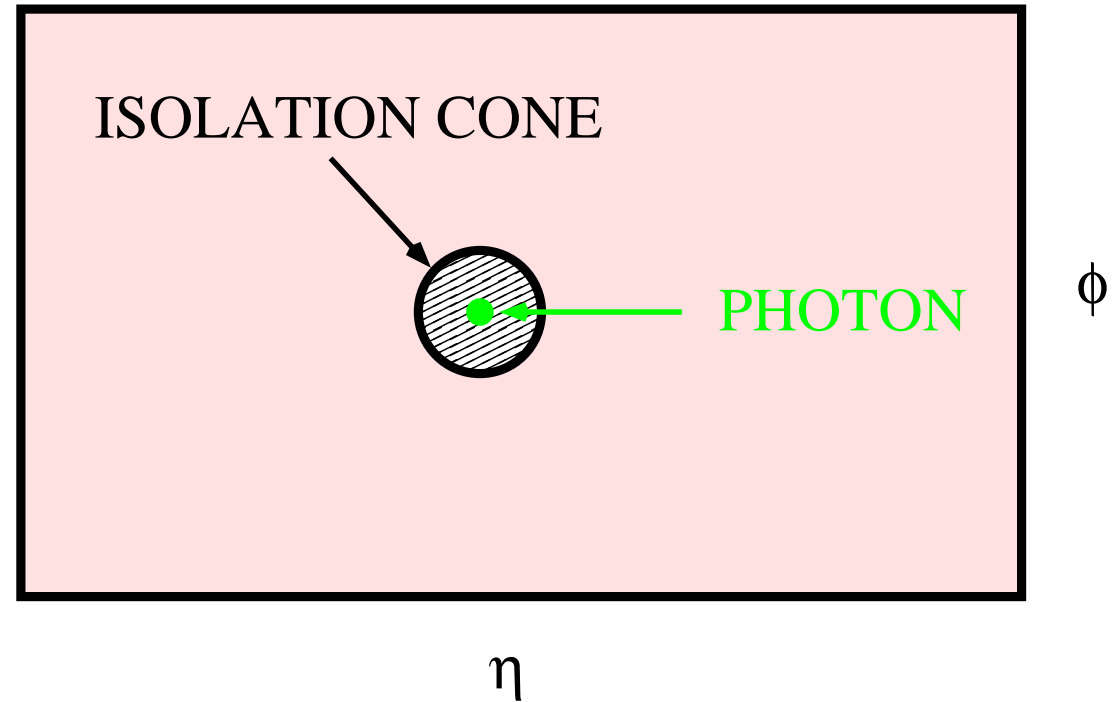
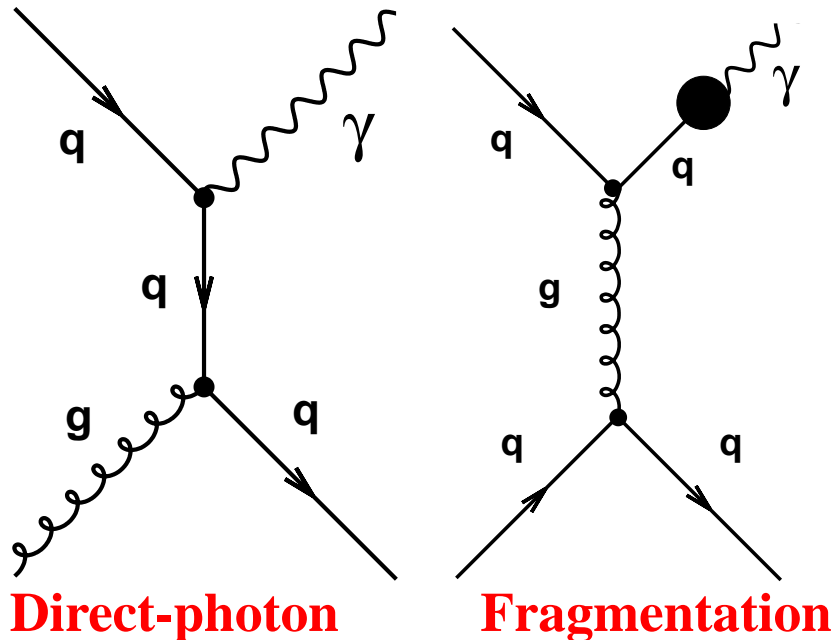
→ Quarks have electric charge and radiate photons  
 ⇒ fragmentation function  $D_{q/g}^\gamma(z, \mu_f)$



⇒ Distinct feature: these photons are inside jets, i.e. not isolated!



## Photon isolation

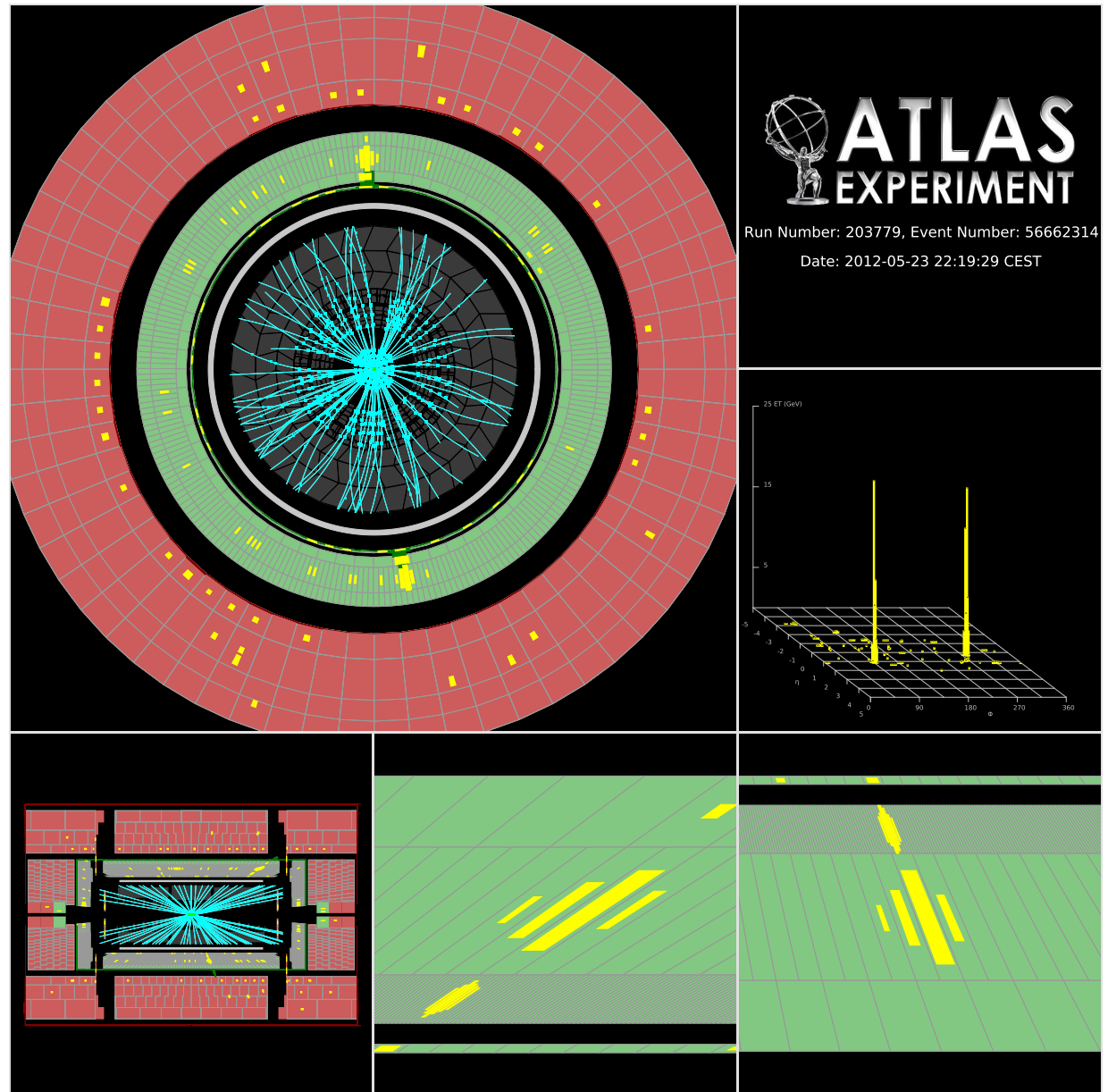


- It is essential to require the photon to be isolated. It is achieved by requiring  $E_T^{iso} \equiv \sum_i E_T^i < E_T^{\max}$  with the sum over the particles (except the photon!) inside a cone of radius  $R = 0.4$  centered on the photon in the  $\eta - \phi$  plane
- The isolation requirement suppresses the contribution of photons inside jets:  $\pi^0$  (as well as other neutral mesons) decays and the fragmentation contribution



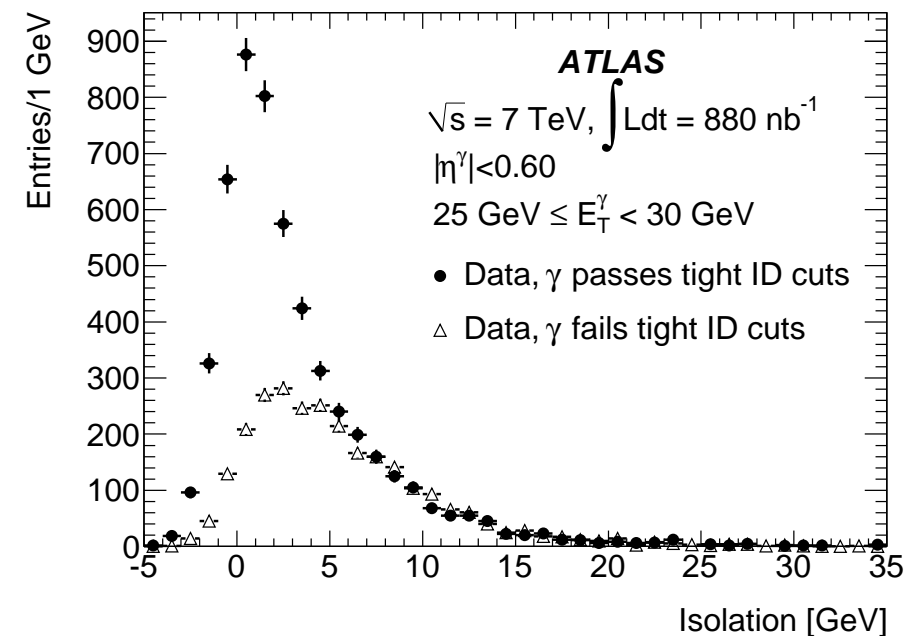
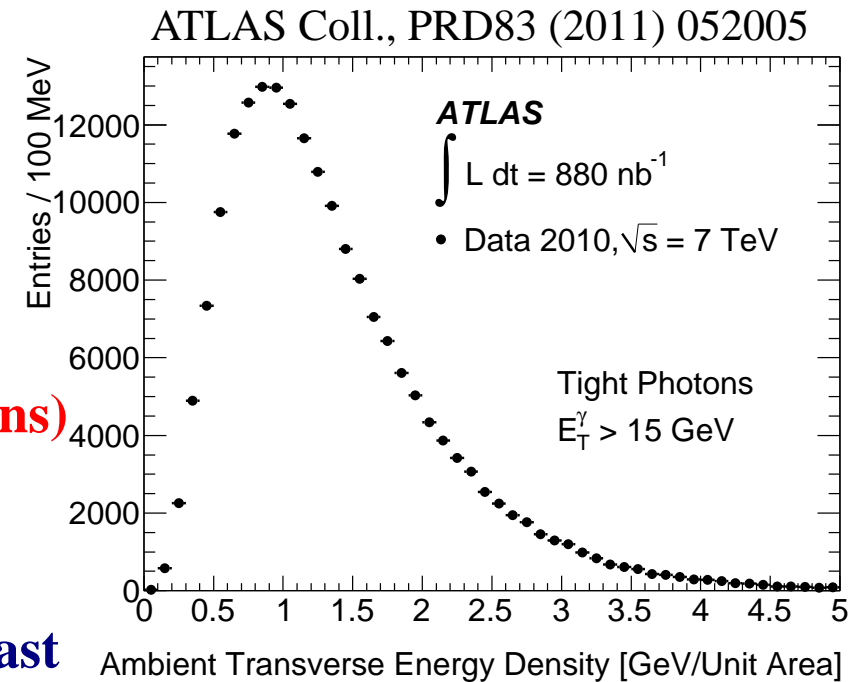
## Photon isolation in ATLAS

- $E_T^{iso}$  ( $R = 0.4$ ) computed using clusters of calorimeter cells (EM and HAD) in a cone  $R = 0.4$ , excluding the contribution from the photon
  - Subtraction of the leakage of the photon energy outside that region (few %)
  - The underlying event and pileup (overlapping  $pp$  interactions in the same/neighbouring bunch crossings) contribute to  $E_T^{iso}$ !
- Subtracted on event-by-event basis using the jet-area method of M. Cacciari et al
- After isolation requirement, residual background still expected



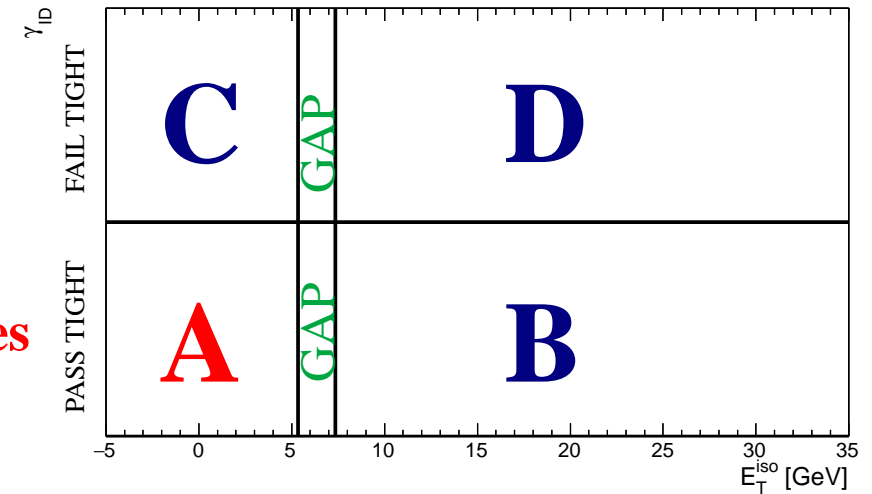
## Photon isolation in ATLAS

- $E_T^{iso}$  is corrected by subtracting the estimated contributions from the underlying event and pileup; the correction is computed on an event-by-event basis (to avoid the large fluctuations) using the jet-area method (M. Cacciari et al.)  
 $\Rightarrow$  ambient transverse-energy density  
 540 MeV (in  $R = 0.4$  cone) for events with at least one photon candidate with  $E_T > 15$  GeV and exactly one PV (+170 MeV for each extra PV)
- After the correction the  $E_T^{iso}$  distribution is centered at zero with a width of 1.5 GeV in simulated signal events



## Background subtraction

- Residual background still expected even after the tight identification and isolation requirements
- A data-driven method necessary to avoid relying on detailed simulations of the background processes
- The two-dimensional sideband method:  
→ photon identification  $\gamma_{ID}$  vs  $E_T^{iso}$  plane



- It is assumed that for background events there is no correlation between  $\gamma_{ID}$  and  $E_T^{iso}$

$$\frac{N_A^{bkg}}{N_B^{bkg}} = \frac{N_C^{bkg}}{N_D^{bkg}} \Rightarrow R_{bkg} \equiv \frac{N_A^{bkg} \cdot N_D^{bkg}}{N_B^{bkg} \cdot N_C^{bkg}} = 1$$

and the effects of the small signal contaminations can be accounted for by using

$$\frac{N_A - N_A^{sig}}{N_B - \epsilon_B N_A^{sig}} = \frac{N_C - \epsilon_C N_A^{sig}}{N_D - \epsilon_D N_A^{sig}} \quad \text{to extract the signal yield } N_A^{sig}$$

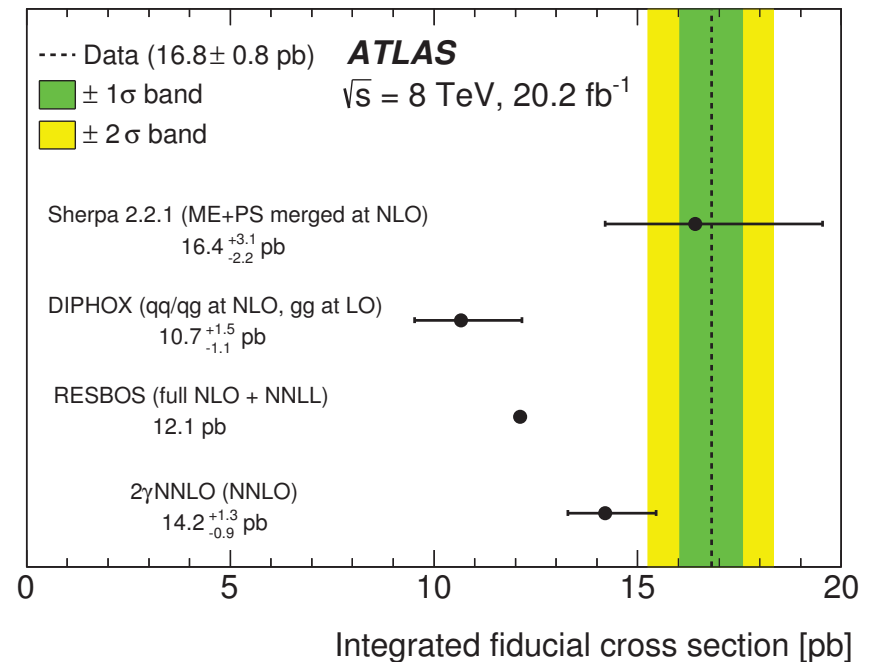
the leakage fractions ( $\epsilon_K \equiv N_K^{sig} / N_A^{sig}$ ,  $K = B, C, D$ ) are estimated using MC samples of signal  $\Rightarrow$  purity rises from 60% ( $E_T^\gamma \sim 25$  GeV) to 100% ( $E_T^\gamma \sim 300$  GeV)

# Diphoton: sample composition and experimental uncertainties

Process	Event fraction [%]	
	Two-dimensional template fit	Matrix method
$\gamma\gamma$	$75.3 \pm 0.3$ (stat) $^{+2.6}_{-2.8}$ (syst)	$73.9 \pm 0.3$ (stat) $^{+3.1}_{-2.7}$ (syst)
$\gamma j$	$14.5 \pm 0.2$ (stat) $^{+2.7}_{-2.8}$ (syst)	$14.4 \pm 0.2$ (stat) $^{+2.0}_{-2.4}$ (syst)
$j\gamma$	$6.0 \pm 0.2$ (stat) $^{+1.4}_{-1.5}$ (syst)	$5.8 \pm 0.1$ (stat) $\pm 0.6$ (syst)
$jj$	$1.6 \pm 0.2$ (stat) $^{+0.9}_{-0.4}$ (syst)	$2.4 \pm 0.1$ (stat) $^{+0.6}_{-0.5}$ (syst)
$ee$	$2.6 \pm 0.2$ (stat) $^{+0.9}_{-0.4}$ (syst)	$3.5 \pm 0.1$ (stat) $\pm 0.4$ (syst)

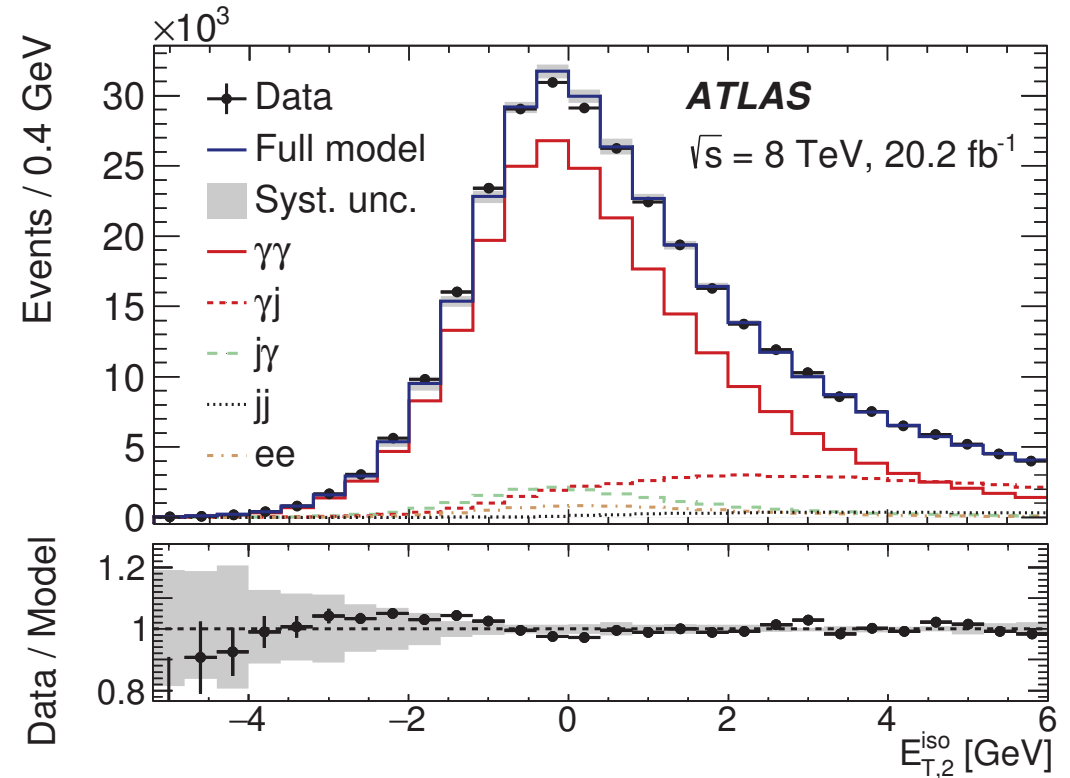
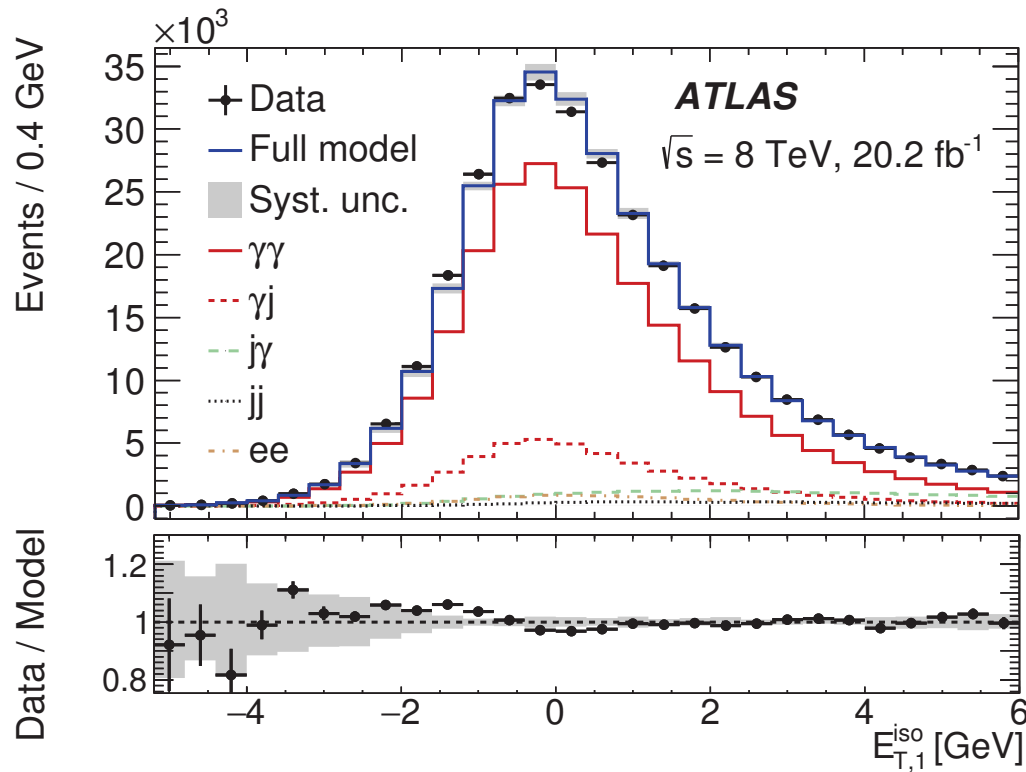
Source of uncertainty	Impact on $\sigma_{\text{tot}}^{\text{fid.}}$ [%]
Photon identification efficiency	$\pm 2.5$
Modeling of calorimeter isolation	$\pm 2.0$
Luminosity	$\pm 1.9$
Control-region definition	$^{+1.5}_{-1.7}$
Track isolation efficiency	$\pm 1.5$
Choice of MC event generator	$\pm 1.1$
Other sources combined	$^{+0.8}_{-1.3}$
Total	$^{+4.5}_{-4.7}$

Name and type of computation



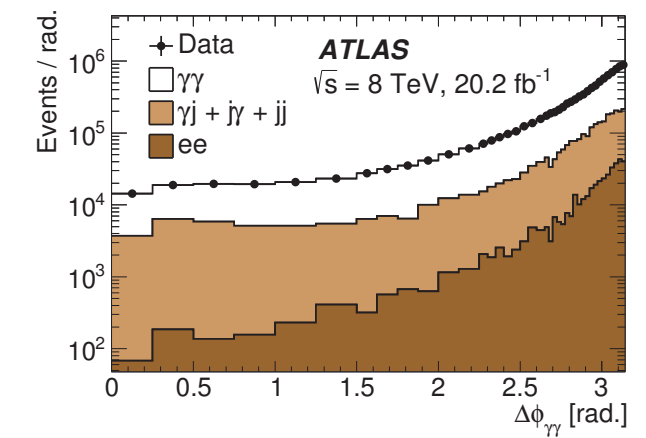
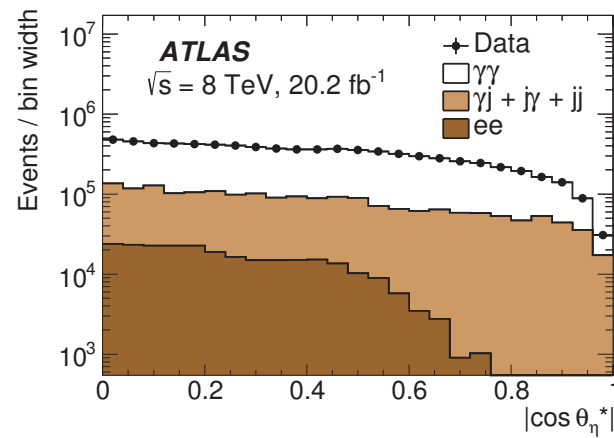
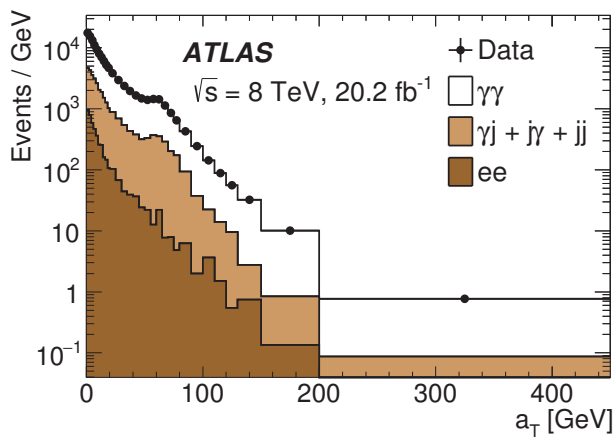
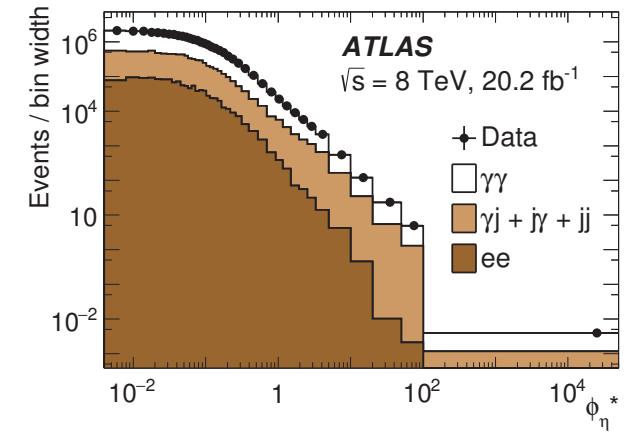
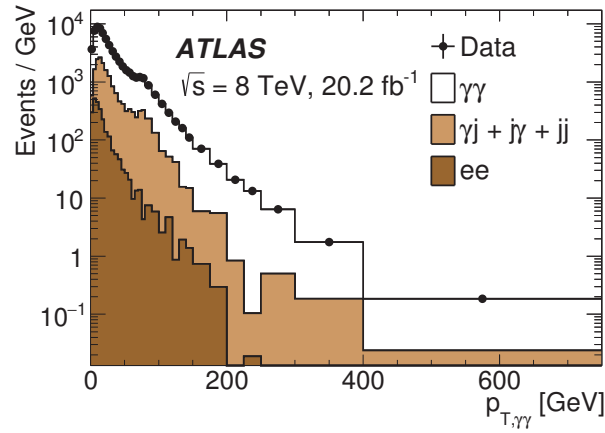
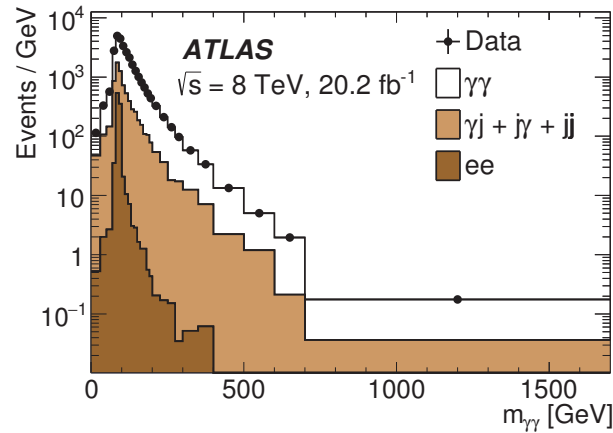
ATLAS Coll., arXiv:1704.03839

# Diphoton: $E_T^{iso}$ distributions



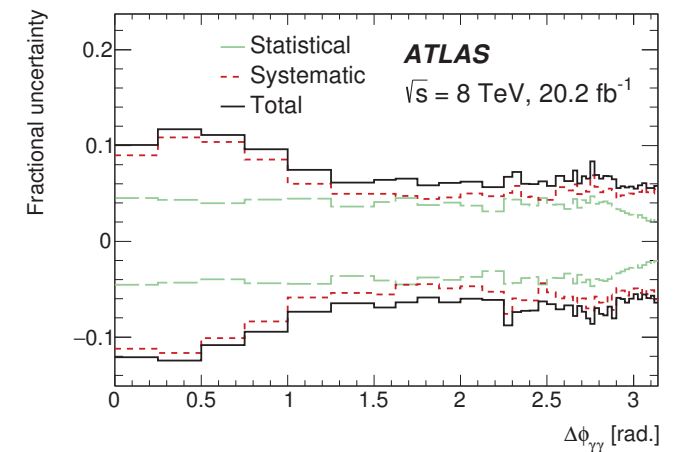
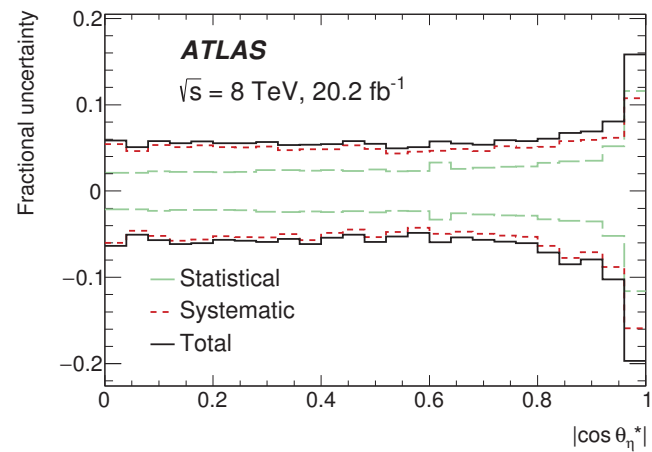
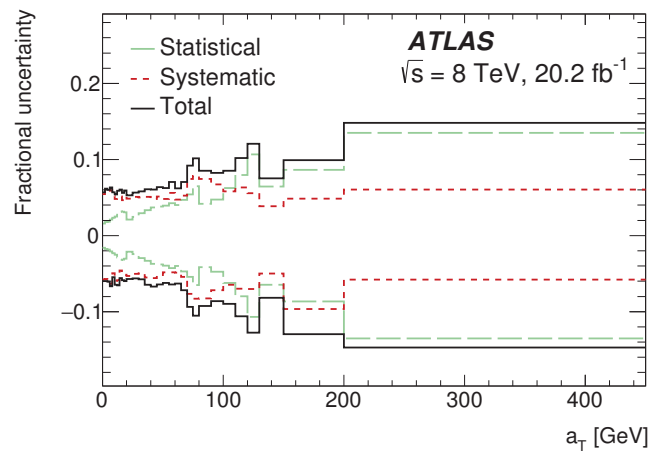
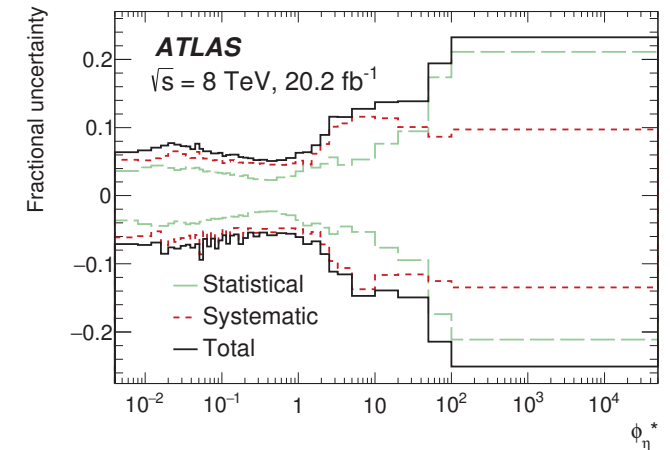
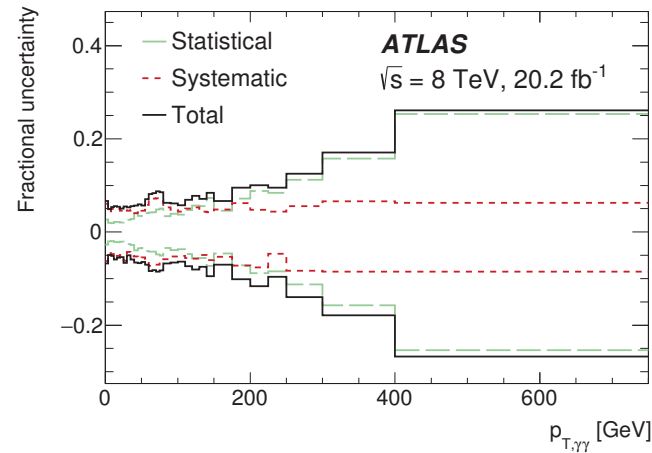
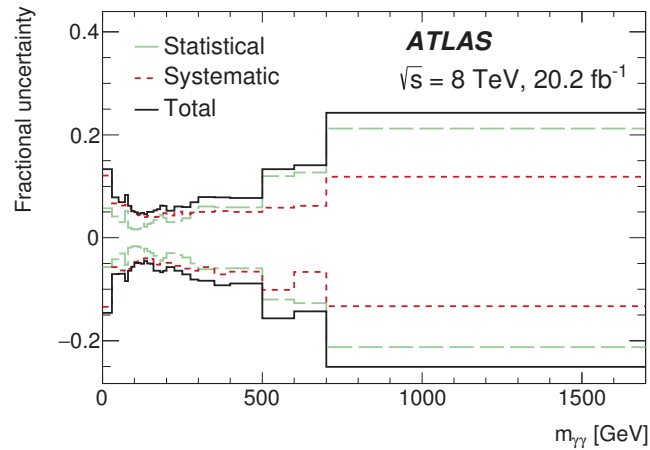
ATLAS Coll., arXiv:1704.03839

# Diphoton: sample composition



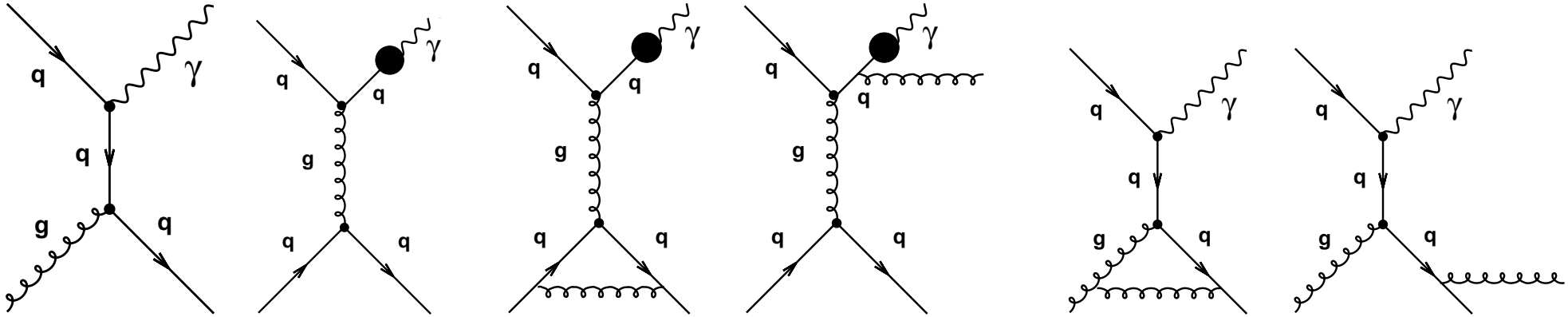
ATLAS Coll., arXiv:1704.03839

# Diphoton: experimental uncertainties



ATLAS Coll., arXiv:1704.03839

# NLO QCD calculations for inclusive photon production



$$\sigma_{pp \rightarrow \gamma + X} = \sum_{i,j,a} \int_0^1 dx_1 f_{i/p}(x_1, \mu_F^2) \int_0^1 dx_2 f_{j/p}(x_2, \mu_F^2) \hat{\sigma}_{ij \rightarrow \gamma a} +$$

$$\sum_{i,j,a,b} \int_{z_{min}}^1 dz D_a^\gamma(z, \mu_f^2) \int_0^1 dx_1 f_{i/p}(x_1, \mu_F^2) \int_0^1 dx_2 f_{j/p}(x_2, \mu_F^2) \hat{\sigma}_{ij \rightarrow ab}$$

- The calculations includes NLO corrections for both direct-photon and fragmentation contributions; **beware the components are not distinguishable beyond LO**
- The calculations implement the photon isolation requirement at “parton” level:  $E_T^{iso}$  calculated with the (few) final-state partons in the perturbative QCD calculation



# NLO QCD calculations for inclusive photon production

$$\sigma_{pp \rightarrow \gamma + X} = \sum_{i,j,a} \int_0^1 dx_1 f_{i/p}(x_1, \mu_F^2) \int_0^1 dx_2 f_{j/p}(x_2, \mu_F^2) \hat{\sigma}_{ij \rightarrow \gamma a^+}$$

$$+ \sum_{i,j,a,b} \int_{z_{min}}^1 dz D_a^\gamma(z, \mu_f^2) \int_0^1 dx_1 f_{i/p}(x_1, \mu_F^2) \int_0^1 dx_2 f_{j/p}(x_2, \mu_F^2) \hat{\sigma}_{ij \rightarrow ab}$$

- Using the JetPhox program (S. Catani, M. Fontannaz, J. Ph. Guillet and E. Pilon) with

→  $\mu_R = \mu_F = \mu_f = E_T^\gamma$  (nominal)

→ proton PDF set: CT10

→ fragmentation function: BFG set II

→ Corrections for hadronisation and underlying event needed

- Theoretical uncertainties:

→ terms beyond NLO; varying  $\mu_R, \mu_F, \mu_f$  by factors 2 and 1/2 (singly or simultaneously)

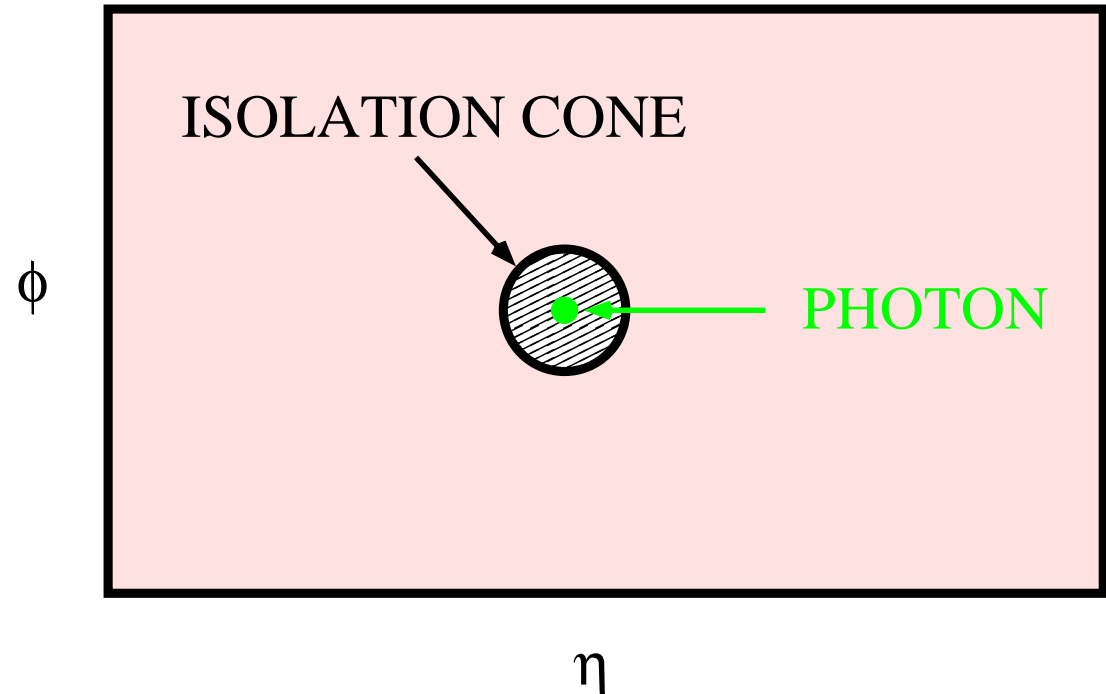
→ PDF-induced uncertainties; estimated using set of PDF eigenvectors

→ uncertainty on  $\alpha_s$ ; estimated using PDFs in which different values of  $\alpha_s$  are assumed

→ uncertainty on non-perturbative correction; estimated with different MCs

## Corrections for non-perturbative effects; photon isolation

- The measurements are corrected for detector effects to the “particle” level  
 → to isolated photons, where  $E_T^{iso}$  is calculated using all the final-state particles and the jet-area method is also applied  
 This is performed using MC simulations



- Corrections for non-perturbative effects (hadronisation and underlying event)

$$C_{NP} = \frac{\sigma_{\gamma+X}(\text{MC, particle - level, UE})}{\sigma_{\gamma+X}(\text{MC, parton - level, no UE})}$$

→ Less dependence on the modelling of the final state by having used the jet-area method to subtract the “extra” transverse energy contribution to  $E_T^{iso}$

# Impact of inclusive isolated photon measurements at LHC on PDFs

- Analysis by D. d'Enterria and J. Rojo (NPB860,2012,311)
  - Study of the impact on the gluon density of existing isolated-photon measurements from a variety of experiments, from  $\sqrt{s} = 200$  GeV up to 7 TeV
    - those at LHC are the more constraining datasets
    - reduction of gluon uncertainty up to 20%
    - localised in the range  $x \approx 0.002$  to 0.05
- ⇒ improved predictions for low mass Higgs production in gluon fusion, PDF-induced uncertainty decreased by 20%

