

UNCOVERING EXTRA DIMENSIONS AT LHC Spanish LHC Network Meeting IFT/UAM, Madrid, 8-9 May 2017

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Uncovering extra dimensions at LHC

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Outline



- The model of a warped extra dimension
- Direct searches of EW KK modes
- Direct searches of gluon KK modes
- Back-ups: Indirect effect of KK modes: B-anomalies (if time permits)

The model of a warped extra dimension

• We will consider a model where the hierarchy problem is solved by a warped extra dimension y with general metric A(y), and two branes, at the UV (y = 0) and at the IR $(y = y_1)$

$$ds^2 = e^{-2A(y)}\eta_{\mu\nu}dx^{\mu}dx^{\nu} + dy^2$$
, $A(y_1) \simeq 35$ (to solve the hierarchy)

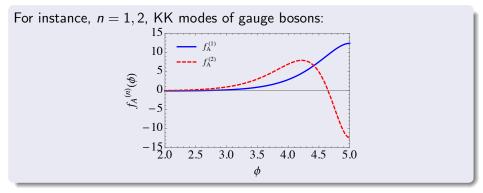
The metric is fixed by the stabilizing bulk field φ, through the 5D Einstein equation or, equivalently, by the "superpotential" W(φ)
Our superpotential choice is

$$W(\phi) = 6k \left(1 + e^{a_0\phi}\right)^{b_0}, \quad b_0 = 2, a_0 = 0.15 \ (b_0 = 0 \ is \ RS)$$

• Extra dimension is close to AdS₅ near the UV, whereas the conformal invariance is broken by a deformation of the metric only near the IR

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- All SM fields propagate in the bulk: gauge vectors, Higgs, fermions
- Every field has the zero mode and the Kaluza-Klein (KK) excitations
- The Higgs zero mode is localized towards the IR to solve the hierarchy problem
- KK excitations are localized towards the IR brane



 KK modes of gauge bosons interact strongly (weakly) with IR (UV) localized fields

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• The SM fermion $f_{L,R}$ is the zero mode of the 5D fermion $\Psi(y,x)$ with appropriate boundary conditions and a 5D Dirac mass term

$$\mathcal{L}_5 = M_{f_{L,R}}(y) \bar{\Psi} \Psi, \quad M_{f_{L,R}}(y) = \mp c_{f_{L,R}} W(\phi)$$

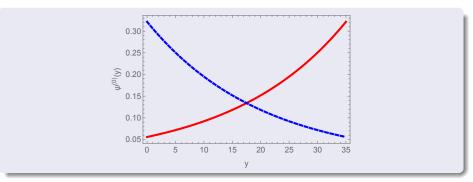
• Explicitely the zero mode profile (in flat coordinates) is given by

$$\psi_{L,R}^{(0)}(y,x) = \frac{e^{(1/2-c_{L,R})A(y)}}{\left(\int dy \, e^{A(1-2c_{L,R})}\right)^{1/2}} f_{L,R}(x)$$

where $f_{L,R}(x)$ are the SM fermion wave-functions • Fermions with c < 0.5 (c > 0.5) are localized towards the IR (UV) brane.

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• For example the profile of fermions with c = 0.45 (solid red) and c = 0.55 (dashed blue) are



- Fermions with c < 0.5 (c > 0.5) are composite (elementary)
- $\bullet\,$ The coupling of KK gauge bosons with fermions is

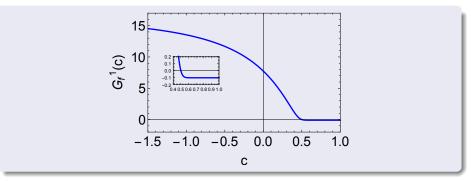
$$G_{f_{L,R}}^n(c_{L,R})g_{f_{L,R}}^{SM}\left(A_{\mu}^n\bar{f}_{L,R}\gamma^{\mu}f_{L,R}\right)$$

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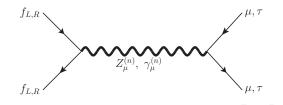
 In particular the interaction of gauge KK modes with fermions is Flavor Non-Universal, depending on the values of c_{fL,R}: they couple differently to different fermions



- The coupling with IR localized fermions is stronger than the coupling with UV localized fermions
- For c = 0.5 the coupling is zero and for $c \gtrsim 0.5$ ($c \lesssim 0.5$) the coupling is tiny (large)

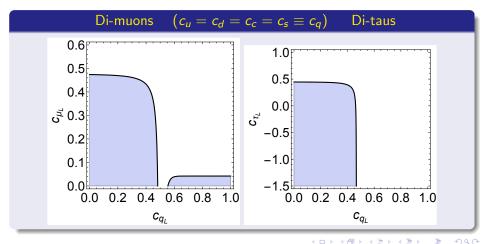
Direct searches of $Z^{(n)}$ and $\gamma^{(n)}$

- The resonances Z_{μ}^{n} and γ_{μ}^{n} (n = 1) can be produced by Drell-Yan processes and decay into a pair of leptons as in the figure
- The best bounds on dimuon resonances have been given by the ATLAS Collaboration based on 3.2 fb⁻¹ data at $\sqrt{s} = 13$ TeV. ATLAS obtained a 95% CL bound on the sequential Standard Model (SSM) Z'_{SSM} gauge boson mass, as $M_{Z'_{SSM}} \gtrsim 3.36$ TeV.
- Similarly the strongest bounds on ditau resonances have been obtained by the CMS Collaboration based on 2.2 fb⁻¹ data at $\sqrt{s} = 13$ TeV. CMS got the 95% CL bound $M_{Z'_{SSM}} \gtrsim 2.1$ TeV.



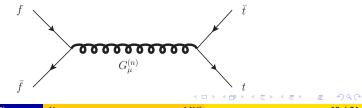
Dileptons searches in DY

A 2 TeV resonance (or even lighter) should not have been discovered yet for the white region of parameter space (result slightly model dependent)



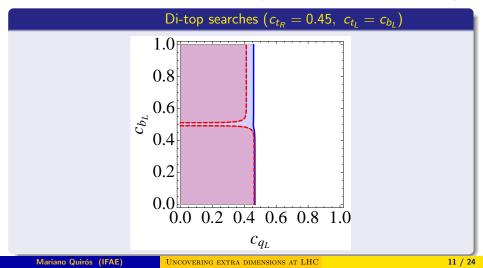
Direct searches of G⁽ⁿ⁾

- Single KK gluons G^n_{μ} (n = 1) can be produced at LHC by Drell-Yan processes and decay into top quarks as in figure
- ATLAS uses $G^1_{q_{L,R}} \simeq -0.2$, for (q = u, d, c, s), $G^1_{b_R} \simeq -0.2$, $G^1_{t_L} \simeq 0.95$ and $G^1_{t_R} \simeq 1.98$. From data at $\sqrt{s} = 8$ TeV corresponding to 20.3 fb⁻¹ they obtain the 95% CL lower bound $M_1^{ATLAS} \gtrsim 2.2$ TeV
- CMS uses data at $\sqrt{s} = 8$ TeV corresponding to an integrated luminosity of 19.7 fb⁻¹. Using $G^1_{q_{L,R}} \simeq -0.2$, for (q = u, d, c, s), $G^1_{b_R} \simeq -0.2$, $G^1_{t_L} \simeq 1$ and $G^1_{t_R} \simeq 5$, they obtain the 95% CL lower bound $M_1^{CMS} \gtrsim 2.5$ TeV



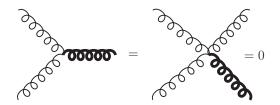
Ditop searches in DY

A 2 TeV resonance (or even lighter) should not have been discovered yet for the white region of parameter space (result slightly model dependent)



Four-top searches in DY

- The DY production of KK gluons decaying into two tops can be easily evaded by the localization of the first and second family quarks
- The alternative is to produce KK gluons by gluon fusion
- In this case the production of a simple KK gluon decaying into two tops is zero (by orthonormality of wave functions)



- KK gluons can be produced in pairs or in association with tt as in the figure
- Undotted vertices are SM couplings
- Red dotted vertices are strong (non SM) couplings
- All this processes would lead to four-top production



BACK UP TRANSPARENCIES

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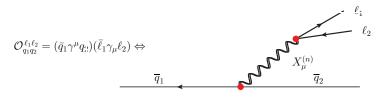
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Accommodating the B-anomalies

- The warped model we have described contains the required ingredients to solve the anomalies
- Quark and leptons $f_{L,R}$ have a certain degree of compositeness, which depends on the value of $c_{f_{L,R}}$, and determines the strength of interaction with KK modes [the red dots]
- KK modes $X_{\mu}^{(n)} = Z_{\mu}^{(n)}, \gamma_{\mu}^{(n)}, W_{\mu}^{(n)}$ generate the effective operators via the tree level diagrams



The $R_{K^{(*)}}$ anomalies

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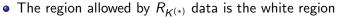
- The SM departure for $R_{K^{(*)}}$ is generated by the diagram
- The FCNC current $(ar{b}\gamma^\mu s)$ is generated from

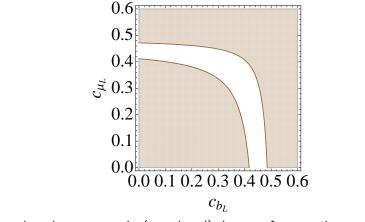
$$\left(V_d^{\dagger}G_d^n V_d\right)_{32}$$

$$G_d^n = \operatorname{diag}(G_d^n(c_d), G_s^n(c_s), G_b^n(c_b))$$

Where V<sub>d_{L,R} are unitary matrices diagonalizing the down mass matrix
 We have considered Wolfenstein-like parametrizations for V<sub>d_{L,R}, V_{u_{L,R}}
</sub></sub>

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• b_L and μ_L have a certain (correlated) degree of compositeness

$$c_{\mu_L} < 0.46, \quad c_{b_L} < 0.48$$

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The $R_{D^{(*)}}$ anomalies

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- The SM departure for $R_{D^{(*)}}$ is generated by the diagram
- The FC charged current $(\bar{b}_L \gamma^\mu c_L)$ is generated from

$$\left(V_{d_L}^{\dagger}G_{d_L}^nV_{u_L}\right)_{32}$$

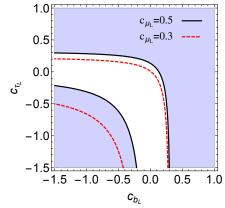
$$G_{d_L}^n = \mathsf{diag}(G_{d_L}^n(c_{d_L}), G_{s_L}^n(c_{s_L}), G_{b_L}^n(c_{b_L})), \quad c_{u_L} = c_{d_L}$$

- Where $V_{u_{L,R}}$ are unitary matrices diagonalizing the up mass matrix
- We have considered a parametrization such that $V_{d_l}^{\dagger} V_{u_L} = V_{CKM}$

 $\frac{\overline{b}}{\underline{d}, \underline{u}} \xrightarrow{D^{(*)}} \overline{z}$ Uncovering extra dimensions at LHC

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- The relevant parameters here are c_{b_L}, c_{τ_L}
- The region allowed by $R_{D^{(*)}}$ data is the white region



• b_L and τ_L have a certain (correlated) degree of compositeness

$$c_{ au_L} \lesssim 0.3, \quad c_{b_L} \lesssim 0.3$$

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The $R_{K^{(*)}}$ and $R_{D^{(*)}}$ anomalies

• The LHCb Collaboration has determined the ratios for $\bar{B} \rightarrow \bar{K}\ell\ell$ $(\ell = \mu, e)$ for muons over electrons for $1 < q^2/GeV^2 < 6$ yielding

$$R_{K} = \frac{\mathcal{B}(\bar{B} \to \bar{K}\mu\mu)}{\mathcal{B}(\bar{B} \to \bar{K}ee)} = 0.745^{+0.090}_{-0.074} \pm 0.032$$

As the SM result is

$$R_K^{\rm SM}\simeq 1$$

this result departs from the SM prediction by ~ 2.6

• This suggests a Lepton Flavor Universality Violation in the process

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Very recently the same tendency has been shown for the ratios

$$R_{K^*} = \frac{\mathcal{B}(\bar{B} \to \bar{K}^* \mu \mu)}{\mathcal{B}(\bar{B} \to \bar{K}^* ee)} = \begin{cases} 0.660^{+0.110}_{-0.070} \pm 0.024, & 0.045 < q^2/GeV^2 < 1.1\\ 0.685^{+0.113}_{-0.069} \pm 0.047, & 1.1 < q^2/GeV^2 < 6 \end{cases}$$

which depart from the SM prediction $\sim 2.5\sigma$ and also suggests Lepton Flavor Universality Violation in the process

 $b \to s \ell \ell$

• As in the Standard Model $R_K \simeq 1$ and $R_{K^*} \simeq 1$, this would suggest

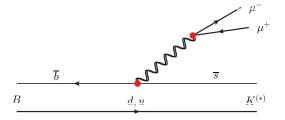
New Physics coupled to the μ (not *electron*) sector

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The anomalies

- A solution to this problem can be provided by vector bosons which couple to muons much more strongly than to electrons
- $\bullet\,$ A typical new physics diagram which gives rise to effective operators

$$\begin{aligned} \mathcal{O}_{9}^{\ell} &= (\bar{s}_{L}\gamma_{\mu}b_{L})(\bar{\ell}\gamma^{\mu}\ell) \,, \qquad \mathcal{O}_{10}^{\ell} &= (\bar{s}_{L}\gamma_{\mu}b_{L})(\bar{\ell}\gamma^{\mu}\gamma_{5}\ell) \,, \\ \mathcal{O}_{9}^{\prime\ell} &= (\bar{s}_{R}\gamma_{\mu}b_{R})(\bar{\ell}\gamma^{\mu}\ell) \,, \qquad \mathcal{O}_{10}^{\prime\ell} &= (\bar{s}_{R}\gamma_{\mu}b_{R})(\bar{\ell}\gamma^{\mu}\gamma_{5}\ell) \,. \end{aligned}$$



• The charged current decays $\bar{B} \to D^{(*)} \ell^- \bar{\nu}_\ell$ have been measured by the BaBar, Belle and LHCb Collaborations which measure

$$R_{D^{(*)}} \equiv R_{D^{(*)}}^{\tau/\mu} = \frac{\mathcal{B}(\bar{B} \to D^{(*)}\tau^- \bar{\nu}_{\tau})}{\mathcal{B}(\bar{B} \to D^{(*)}\ell^- \bar{\nu}_{\ell})}, \quad (\ell = \mu \text{ or } e)$$

• The averaged experimental results

$$R_D = 0.403 \pm 0.047, \quad R_{D^*} = 0.310 \pm 0.017$$

again depart from the Standard Model predictions

$$R_D^{\rm SM} = 0.300 \pm 0.011, \quad R_{D^*}^{\rm SM} = 0.254 \pm 0.004$$

by $\sim 2.2\sigma$ and 3.3σ , although the combined deviation is $\sim 4\sigma$ • This suggests Lepton Flavor Universality Violation in the process

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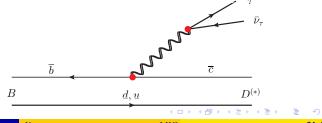
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This would lead to

NEW PHYSICS MAINLY COUPLED TO THE TAU SECTOR

- A solution to this problem can be given by charged vector bosons which couple to taus much more strongly than to muons and electrons
- A new physics diagram which gives rise to the effective operator

 $(\bar{c}\gamma^n u P_L b)\bar{\tau}\gamma_{\nu}(U\nu)_{\tau}$



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