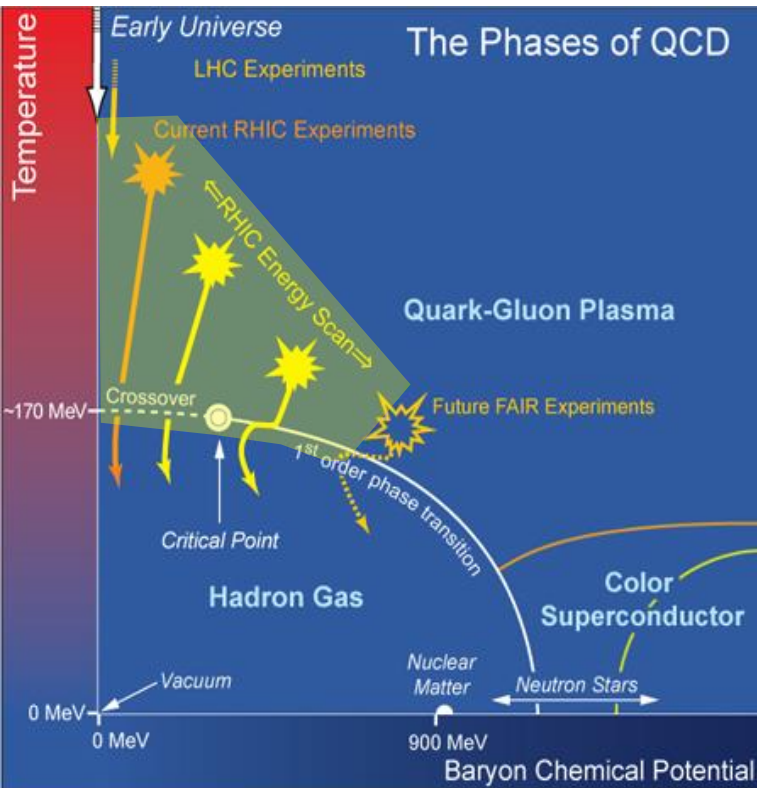


Charge separation measurements in p+Au and A(B)+A collisions: Implications for characterization of the chiral magnetic effect

*Roy A. Lacey (for the STAR Collaboration)
Stony Brook University*

QCD Phase Diagram

Quantitative study of the QCD phase diagram is a current focus of our field



Measurements spanning a large domain of the (T, μ_B) -plane are ongoing/slated at RHIC and other facilities

Questions of interest:

➤ Location of the essential “landmarks”?

- ✓ Crossover region?
- ✓ Critical Point (T^{cep}, μ_B^{cep}) ?
 - ❖ static critical exponents - ν, γ, δ ?
 - ❖ dynamic critical exponent/s - z ?
- ✓ Coexistence regions?

➤ Properties of the respective phases?

- ✓ Thermodynamic properties?
- ✓ Topological & Chiral properties
- ✓ Transport properties?

✓ Anomalous chiral transport?

✓ etc

All are required to fully chart the phase diagram

Subject of this talk

Comments about the CEP

Comment - I

Meaningful discussions about the location and character of the Critical End Point (CEP) must include Finite-size and/or Finite-time effects

Comment - II

The Generalized Finite-Size-Finite-Time Scaling form for χ is well known

$$\chi^{(i)}(\tau, t, L) \sim L^3 b^{i\delta\beta/\nu} \chi^{(i)}(\tau b^{1/\nu}, t b^{-z}, L^{-1}b)$$

$$\chi^{(i)} \sim L^{3+i\delta\beta/\nu} f_1(\tau L^{1/\nu}, RL^r) \quad \text{FSS form}$$

$$\chi^{(i)} = L^3 R^{-i\delta\beta/r\nu} f_2(\tau R^{-1/r\nu}, L^{-1}R^{-1/r}) \quad \text{FTS form}$$

$$V=L^3, \tau = (T - T^{\text{cep}}) / T^{\text{cep}}$$

$$r = z + 1/\nu, R = \text{cooling rate } T_i \rightarrow T^{\text{cep}}$$

$$\hat{\xi} = \text{effective cor. length for FTE}$$

If Finite-Time Effects Dominate

$$\xi > L > \hat{\xi}$$

$$\frac{\chi^{(n)}}{\chi^{(m)}} = L^{(n-m)\delta\beta/\nu} f(\tau L^{1/\nu})$$

$$(n - m)\delta\beta/\nu \sim 0 \text{ and } \chi^{(i)} \sim L^3.$$

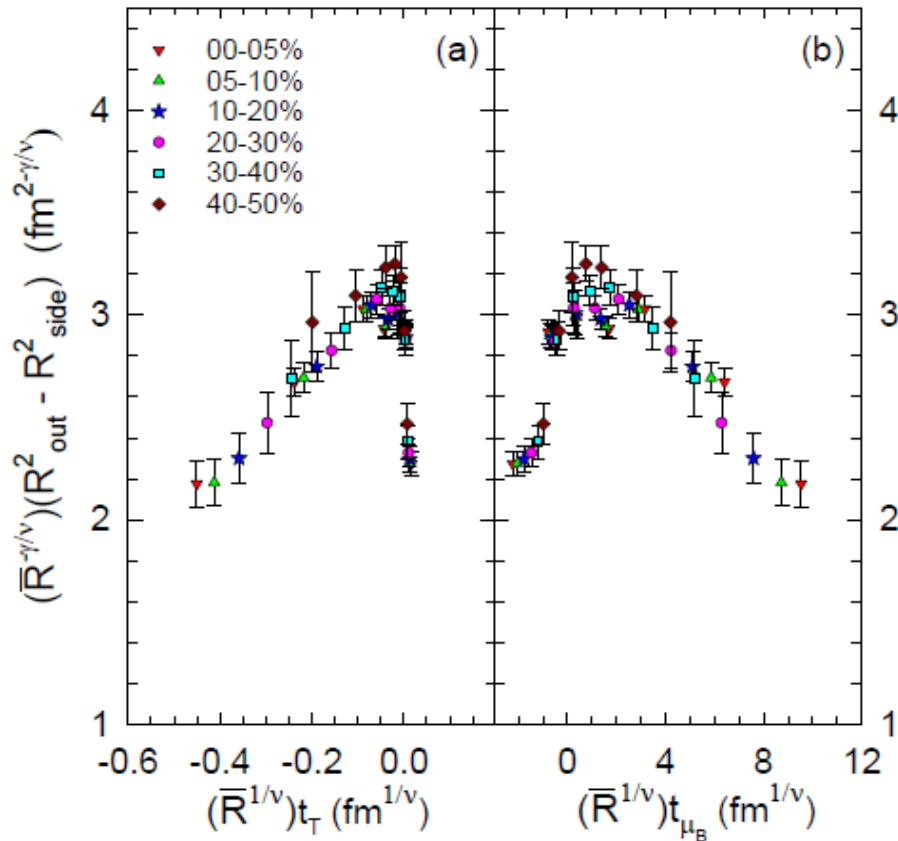
If Finite-Size Effects Dominate

$$\xi > \hat{\xi} > L \quad \frac{\chi^{(n)}}{\chi^{(m)}} = L^{(n-m)\delta\beta/\nu} f(\tau L^{1/\nu})$$

➤ Utilize Finite-Size-Finite-Time Scaling of susceptibility measurements for different sizes (L)

CEP estimate via finite-size scaling

RL - Phys.Rev.Lett. 114 (2015) no.14, 142301



$$T^{cep} \sim 165 \text{ MeV}, \mu_B^{cep} \sim 95 \text{ MeV} (V \rightarrow \infty)$$

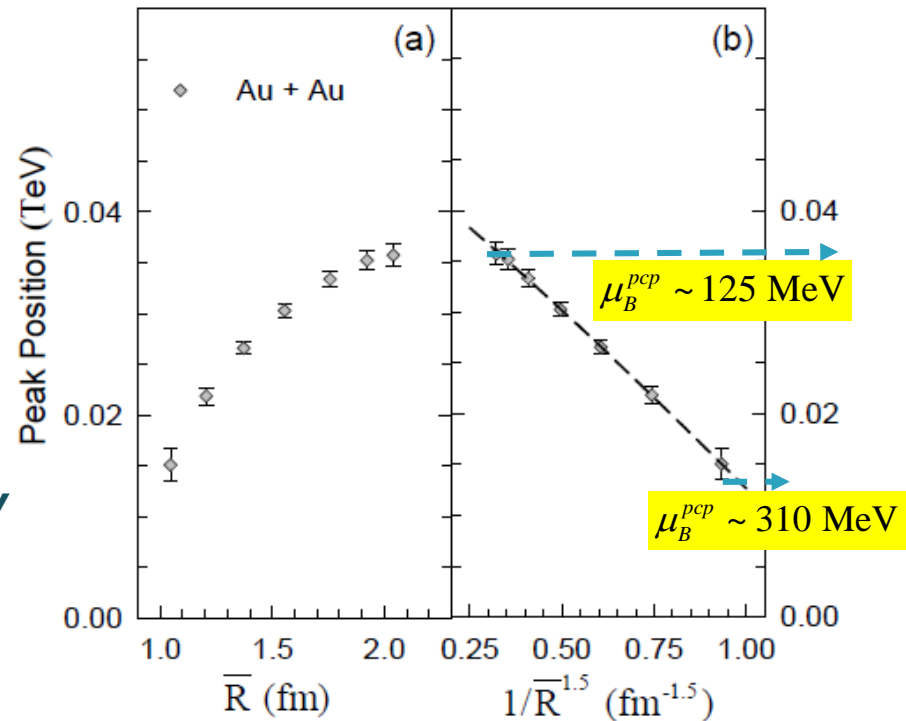
$$\nu \sim 0.66$$

$$\gamma \sim 1.2$$

$$\delta \sim 4.8$$

- ✓ 2nd order phase transition
- ✓ 3D Ising Model (static) universality class for CEP

****The pseudocritical point (PCP) has a strong volume dependence****



****Scaling functions for proxy susceptibility (κ) validates the estimate for the location of the CEP and the (static) critical exponents \rightarrow compressibility diverges****

Finite-Size Effects shifts μ_B^{pcp} to much larger values

CEP estimate via Finite-Size-Finite-Time scaling

Stokić, et. al, Phys.Lett.B673:192-196,2009

For an isothermal change

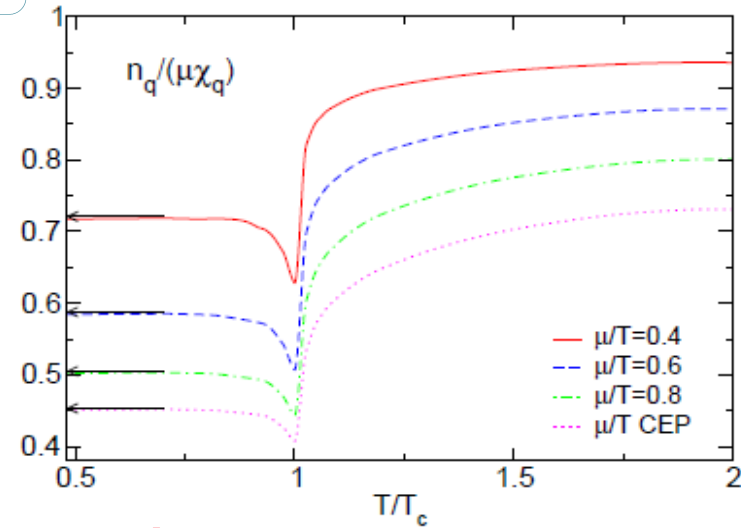
$$VdP = Nd\mu$$

$$N \left(\frac{\partial \mu}{\partial N} \right)_{V,T} = V \left(\frac{\partial P}{\partial N} \right)_{V,T} = \frac{1}{\rho \kappa_T}$$

From partition function one obtains

$$\frac{1}{\langle N \rangle} \left(\frac{\partial \langle N \rangle}{\partial \beta \mu} \right) = \frac{\langle N^2 \rangle - \langle N \rangle^2}{\langle N \rangle}$$

$$\kappa_T^{-1} \propto \frac{\langle N \rangle}{\langle N^2 \rangle - \langle N \rangle^2} = \frac{C_1}{C_2}$$

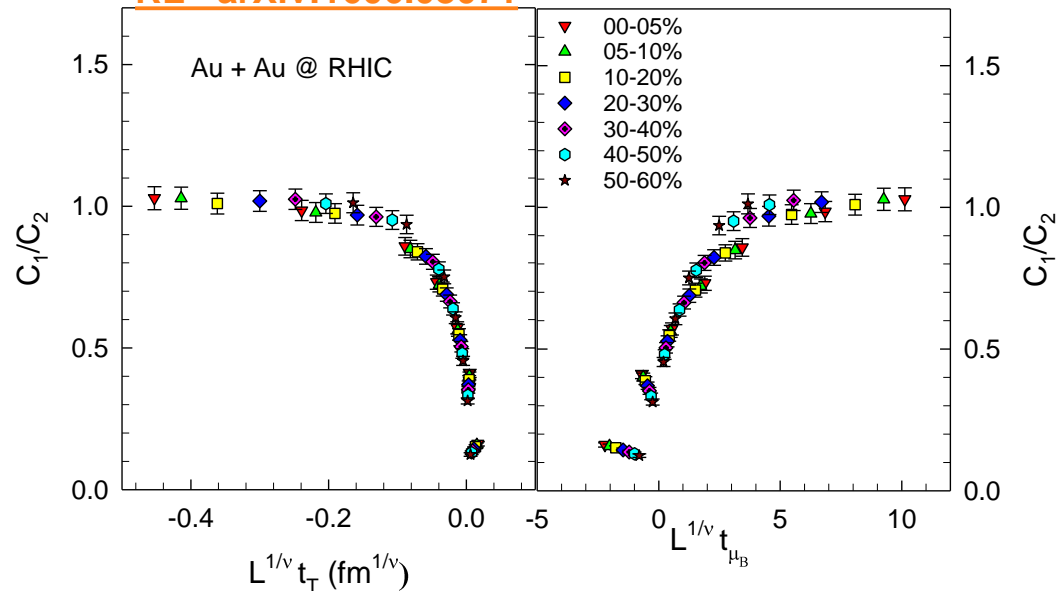


- At the CEP the inverse compressibility $\rightarrow 0$
- Scaling properties of fluctuations data can be leveraged.

$T^{cep} \sim 165$ MeV, $\mu_B^{cep} \sim 95$ MeV ($V \rightarrow \infty$)

Use T^{cep} , μ_B^{cep} , and v in conjunction with Finite-Size-Finite-Time scaling to obtain Scaling Function for C_1/C_2

RL - arXiv:1606.08071



Anomalous Transport in the QGP

Chiral Magnetic Effect

Electric Current

Chiral Magnetic Conductivity

Chiral Chemical potential

$$\vec{J}_Q = \sigma_5 \vec{B}$$

$$\sigma_5 = C_A \mu_5$$

$$C_A = Q^2 / (4\pi^2)$$

Kharzeev
hep-ph/0406125

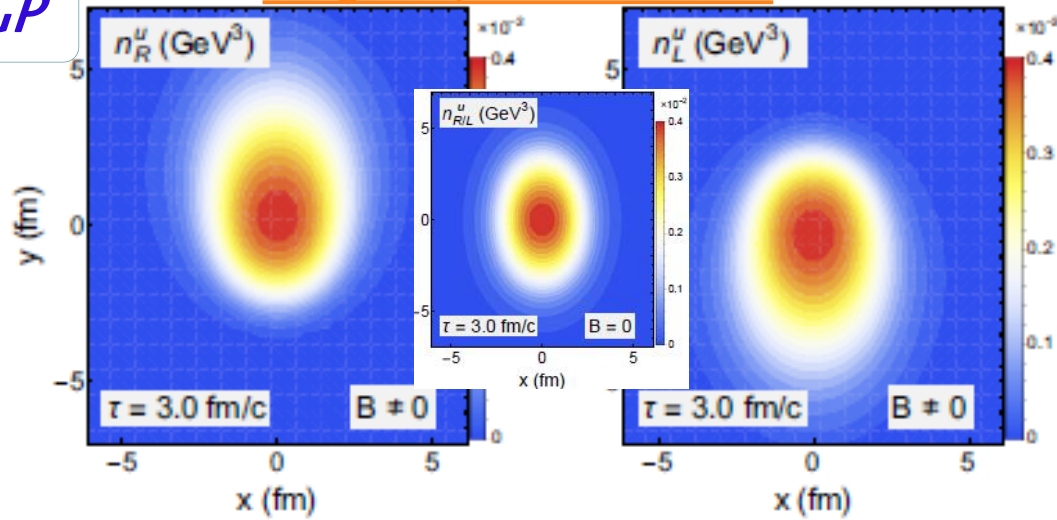


FIG. 1. Comparison of the evolution of the densities (n_R^u, n_L^u) for right-handed (left panel) and left-handed (right panel) u-flavor quarks at $\tau = 3.00$ fm/c for a nonzero B field along the positive y -axis (i.e. perpendicular to the reaction plane). The figures, which show results from Anomalous Viscous Fluid Dynamics (AVFD) calculations, are taken from Ref. [8].

The Chiral Magnetic Effect (CME) results from anomalous chiral transport of the chiral fermions in the QGP, leading to the generation of an electric current

\vec{J}_Q along the magnetic field \vec{B} generated in the collision:

→ **Leads to charge separation about the event plane**

Charge separation leads to a dipole term in the azimuthal distribution of the produced charged hadrons:

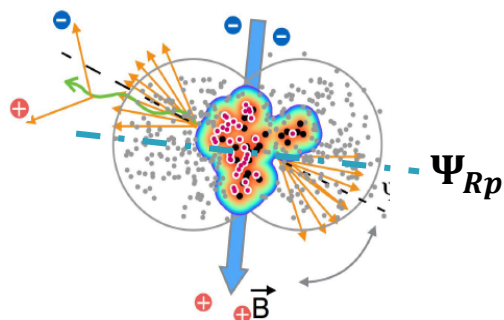
$$\frac{dN^{ch}}{d\phi} \propto [1 \pm 2a_1^{ch} \sin \phi + \dots]$$

Objective: identify & characterize this “dipole moment”

Gamma correlator & its Response

Ma, Zhang

Phys.Lett. B700 (2011) 39-43

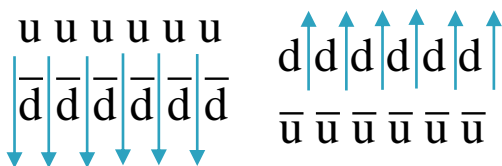


Voloshin, PRC 70 (2004) 057901

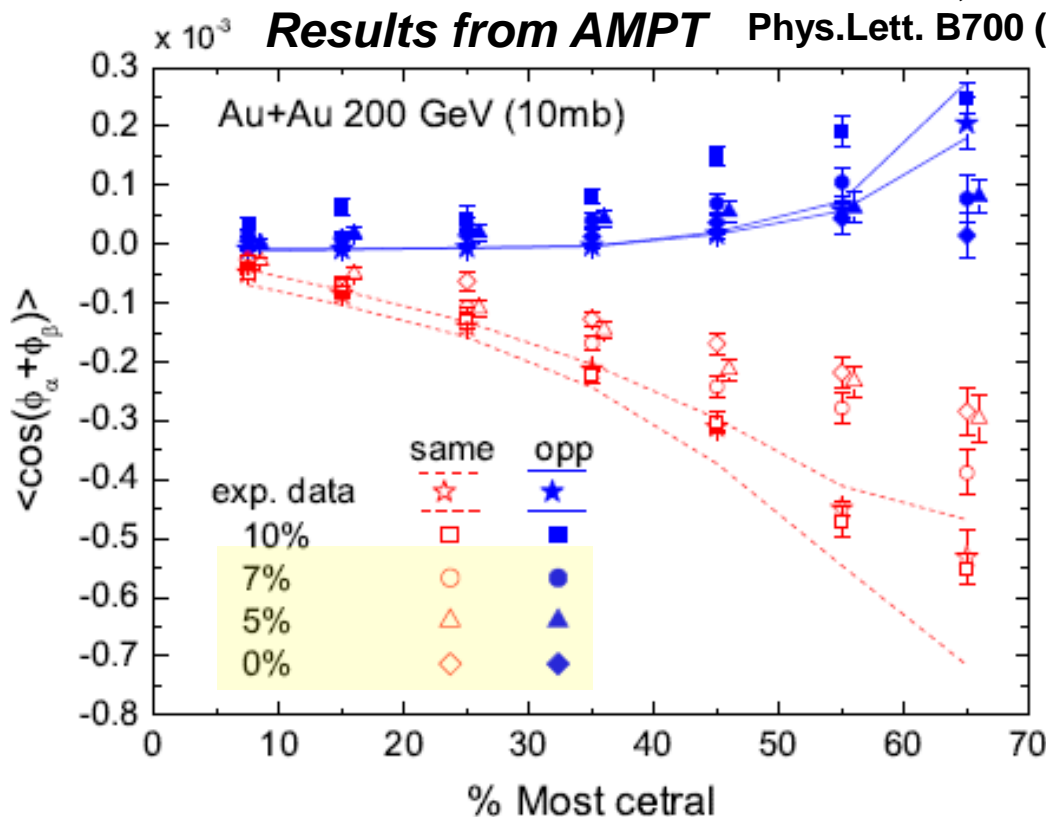
$$\gamma^{\alpha,\beta} = \langle \cos(\phi_\alpha + \phi_\beta - 2\Psi_{RP}) \rangle$$

$$\gamma^{\alpha,\beta} = -\langle a_\alpha a_\beta \rangle + c \frac{v_2}{N}$$

Input charge separation studied with AMPT

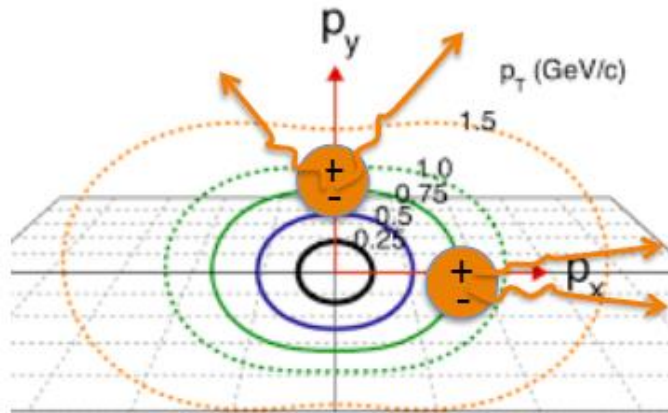


switch the p_y values of a fraction of each set



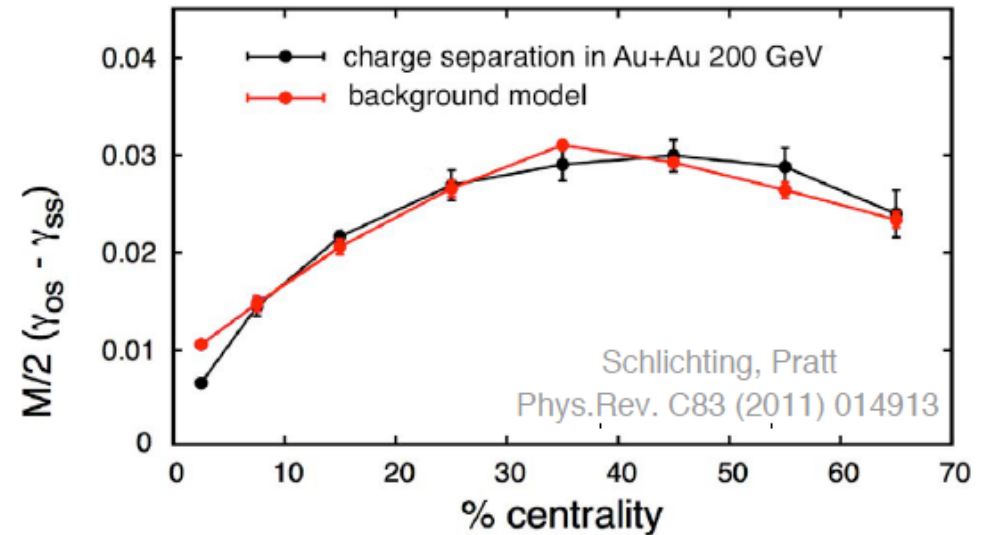
- **The Gamma Correlator's response is similar for signal and background**
 - ✓ **Background-driven correlations complicate CME-driven signal extraction?**

Local charge conservation is an especially important background



$$\gamma^{\alpha,\beta} = -\langle a_\alpha a_\beta \rangle + c \frac{v_2}{N}$$

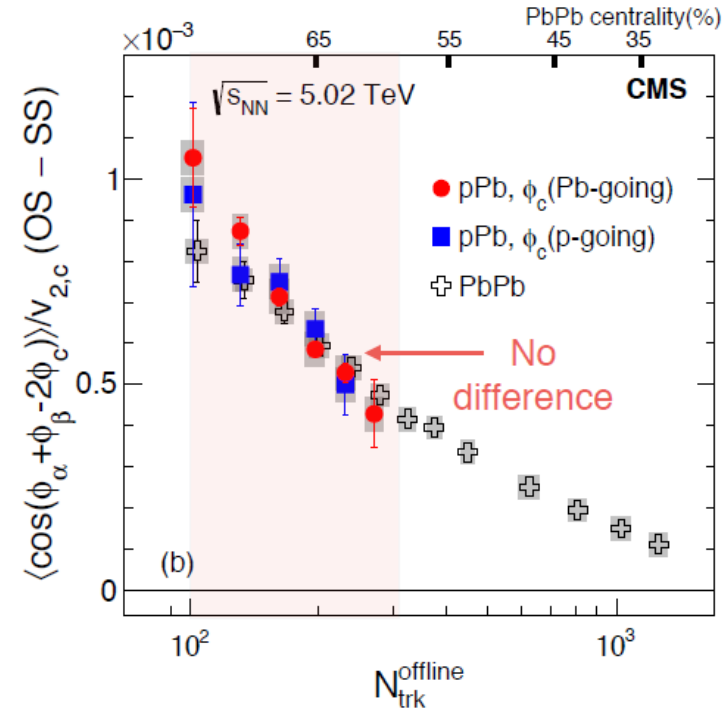
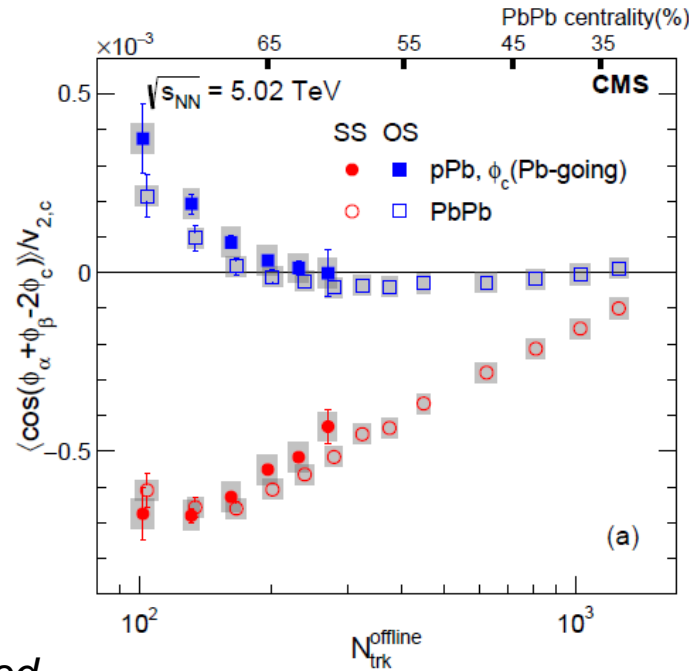
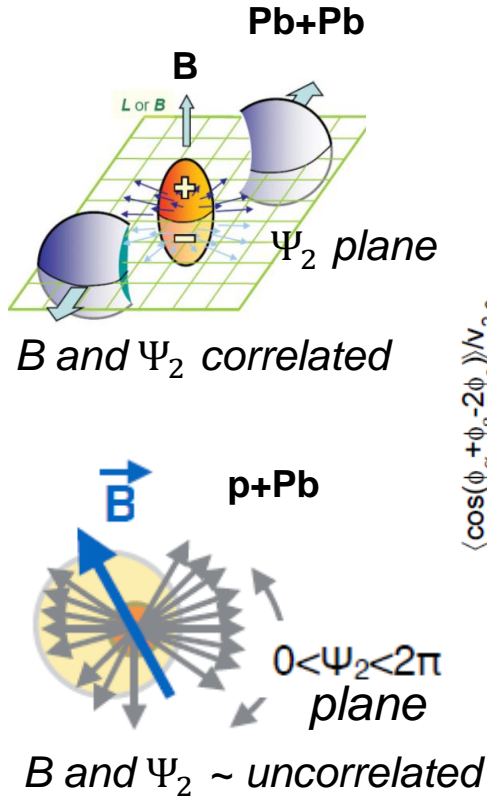
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- **Background-driven correlations can account for a part, or all of the observed charge separation signal?**

Gamma correlator status quo & measurements

Recent CMS measurements for p+Pb and Pb+Pb at the LHC gives cause for pause!



➤ **The magnitudes of the scaled correlators for p+Pb and Pb+Pb are not expected to be the same**

- ➔ “Reduced” magnetic field strength for p+Pb?
- ➔ Large dispersion of the B-field about Ψ_2 in p+Pb

Separating signal from Background

Techniques which can reliably suppress or separate background contributions from the desired CME-driven signal
 → underway

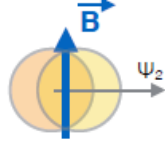
Gamma correlator Measurements focused on "clever" subtraction techniques

Ongoing strategy

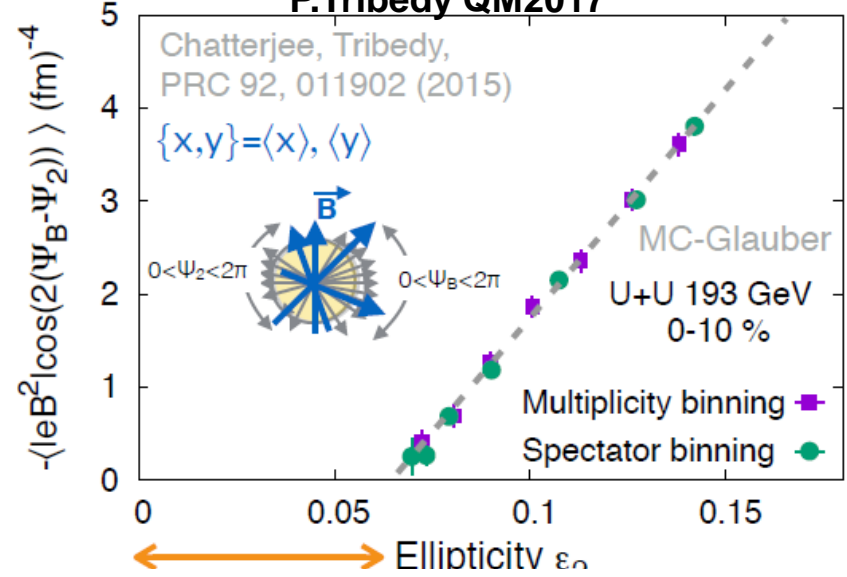
- vary the background for a fixed signal
- vary the signal for fixed background

➤ Data trend compatible with expectation from projected B-field

Projected B-field

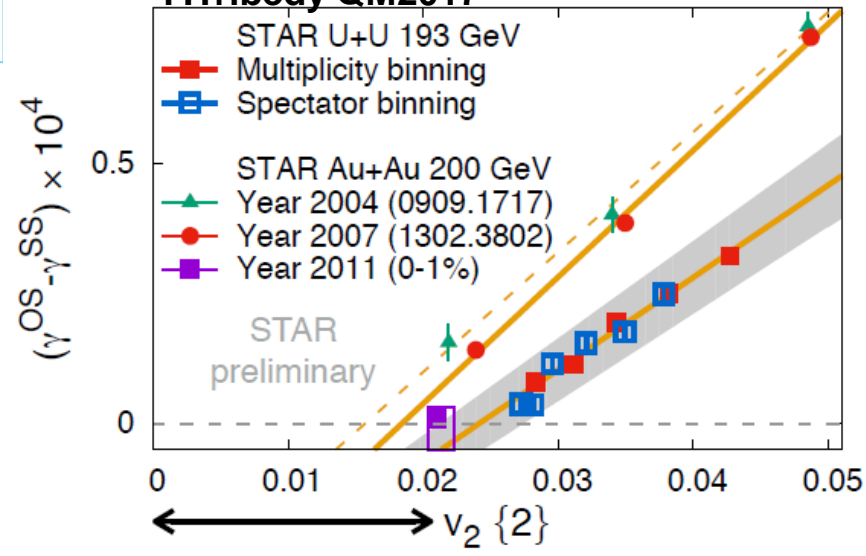


P.Tribedy QM2017



Data

P.Tribedy QM2017



Complimentary approach:

- ✓ *Focus on new measurements with an “improved” correlator which can distinguish between signal and background*

- **A charge-sensitive in-event correlator is used to search for, and characterize CME-driven charge separation**

“New” Correlator

A Multi-particle charge-sensitive in-event correlator is used to measure charge separation (ΔS) relative to the Ψ_2 plane!

$$C_p(\Delta S) = \frac{N_{\text{real}}(\Delta S)}{N_{\text{shuffled}}(\Delta S)}$$

N. N. Ajitanand, et al.,
Phys. Rev. C 83, 011901(R) (2011).

The Numerator is the distribution over events, of the event-by-event averaged ΔS

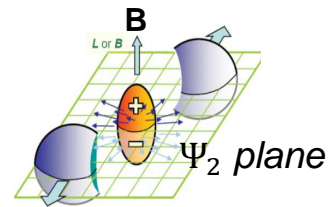
charged hadrons (h^\mp)

$n = \# h^-$

$p = \# h^+$

$$\langle S_p^{h+} \rangle = \frac{\sum_1^p \sin(\Delta\varphi_+)}{p} \quad \Delta\varphi = \phi - \Psi_2 \quad \langle S_n^{h-} \rangle = \frac{\sum_1^n \sin(\Delta\varphi_-)}{n}$$

$$\Delta S = \langle S_p^{h+} \rangle - \langle S_n^{h-} \rangle$$



$N_{\text{shuffled}}(\Delta S)$ → Random shuffling of [only] the charges within an event

Numerator → carries charge separation response

Denominator → carries the “null” or charge averaged response

A second multi-particle correlator $C_p^\perp(\Delta S)$ is similarly constructed for $\Psi_2 + \pi/2$, for which contributions from CME-driven charge asymmetry vanish.

The shape and magnitude of the correlator,

$$R(\Delta S) = C_p(\Delta S) / C_p^\perp(\Delta S)$$

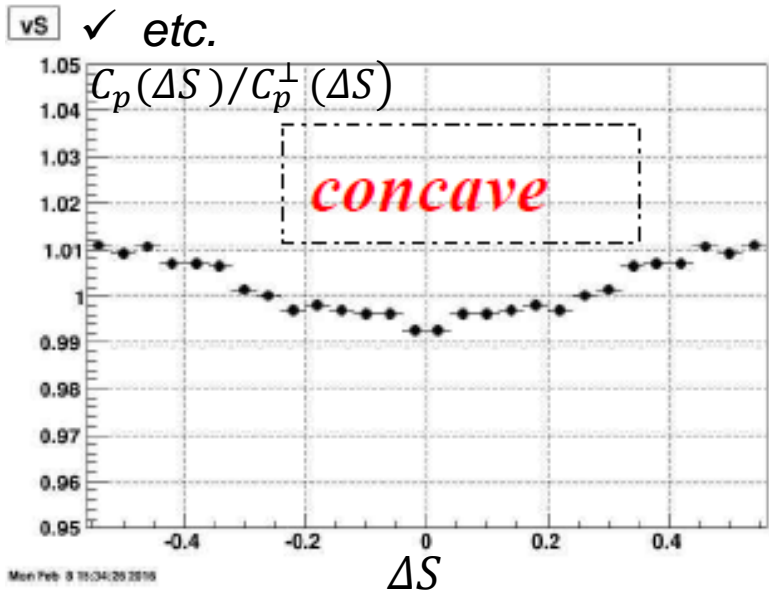
is used to identify and characterize charge separation

Background Dominated

New Correlator Response

$$R(\Delta S) = C_p(\Delta S) / C_p^\perp(\Delta S)$$

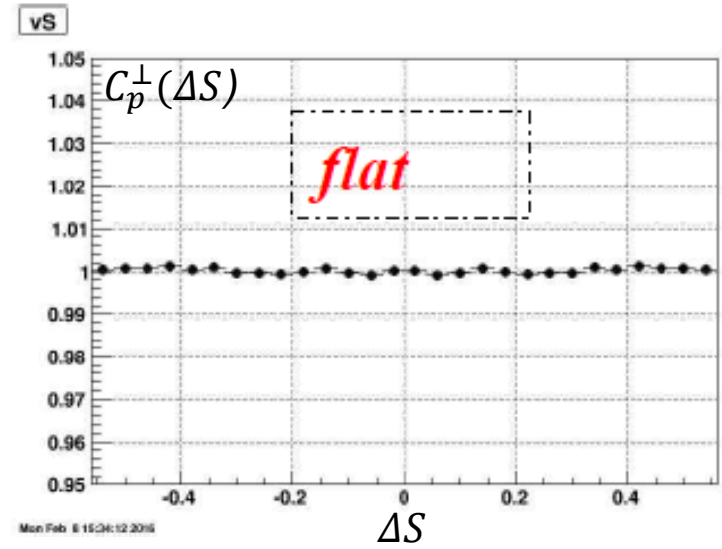
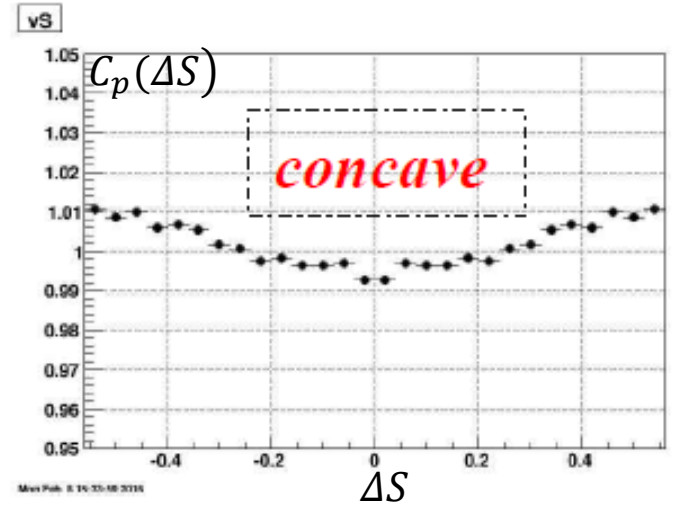
- Is “concave-shaped” for CME-driven charge separation
- Is flat or “convex-shaped” for all known non-CME related backgrounds of interest
 - ✓ collective flow,
 - ✓ momentum conservation,
 - ✓ local charge conservation
 - ✓ etc.



➤ “Concave-shaped” response for input charge separation validated

Toy model with

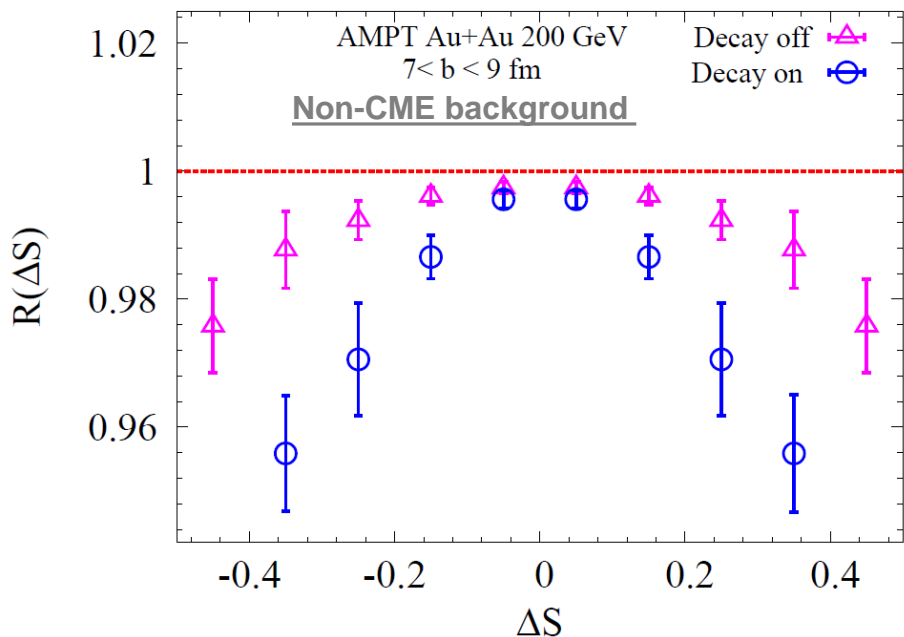
- ✓ Flow
- No resonance decay
- ✓ Charge separation ($a_1 > 0$)



New Correlator Response

$$R(\Delta S) = C_p(\Delta S) / C_p^\perp(\Delta S)$$

- Is “concave-shaped” for CME-driven charge separation
- Is flat or “convex-shaped” for all known non-CME related backgrounds

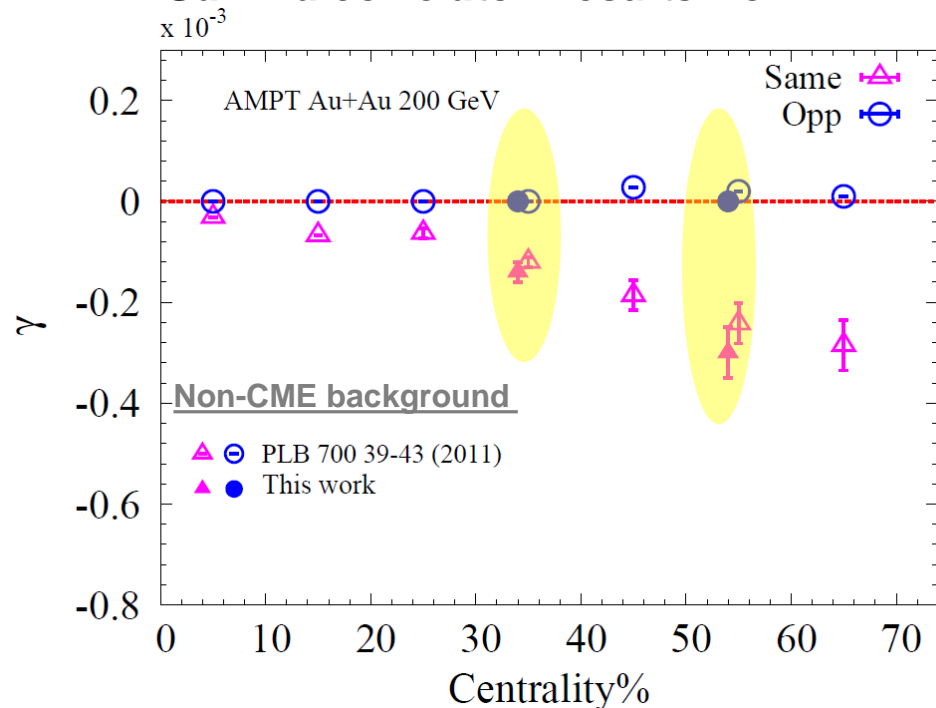


- Non-CME convex-shaped background validated

AMPT model

- ✓ Flow
- ✓ Resonance decay
- Charge separation ($a_1=0$)

Gamma correlator Results from AMPT

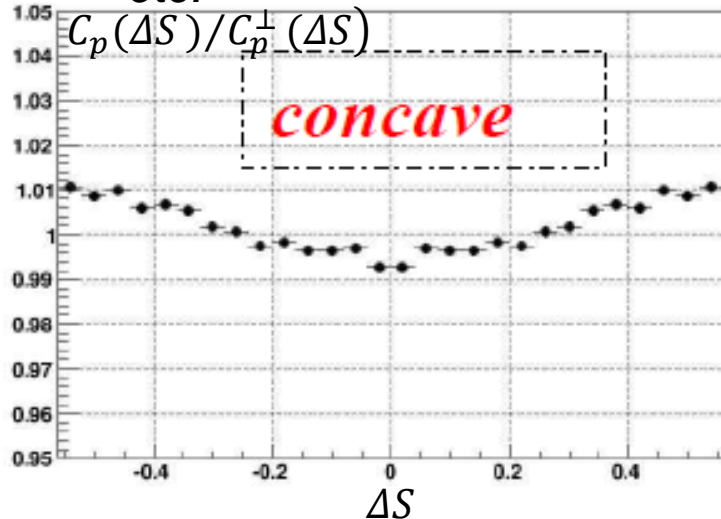


- Gamma correlator’s response contrasts with that for $R(\Delta S)$

New Correlator Response

$$R(\Delta S) = C_p(\Delta S) / C_p^\perp(\Delta S)$$

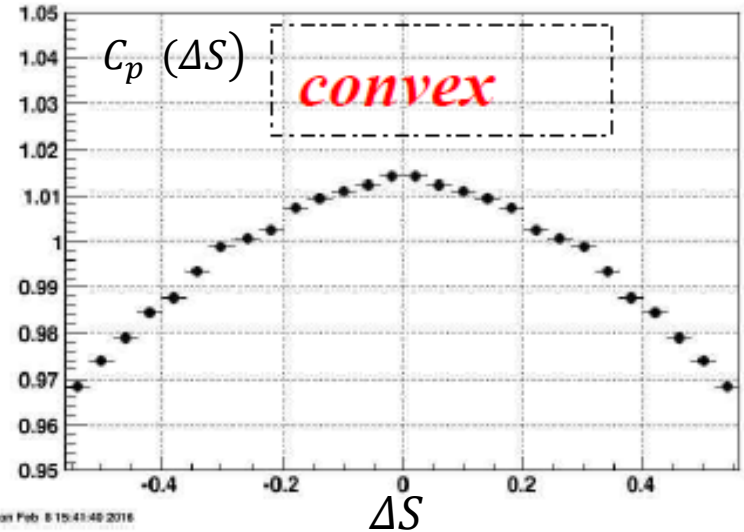
- Is “concave-shaped” for CME-driven charge separation
- Is flat or “convex-shaped” for all known non-CME related backgrounds of interest
 - ✓ collective flow,
 - ✓ momentum conservation,
 - ✓ local charge conservation
 - ✓ etc.



➤ “Concave-shaped” response validates charge separation input in the presence of sizeable background

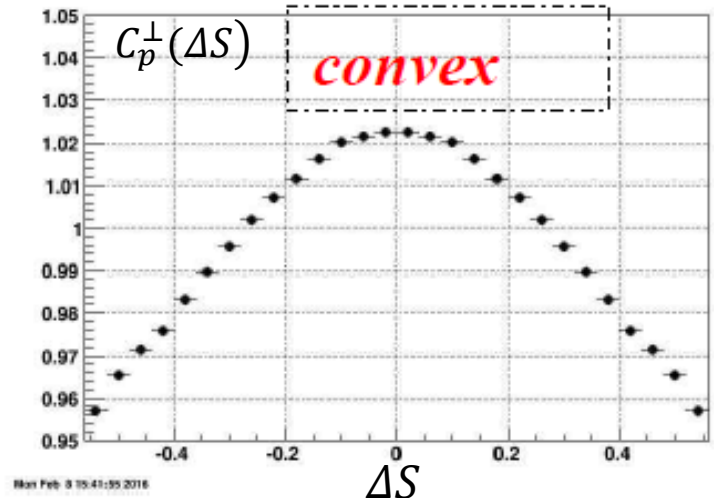
Toy model with

- ✓ Flow
- ✓ Resonance decay
- ✓ Charge separation ($a_1 > 0$)



Mon Feb 8 15:41:48 2016

vS



Mon Feb 8 15:41:35 2016

Signal Quantification

- Data driven simulation, which takes account of the respective correlations due to anisotropic flow, resonance decays, local charge conservation, etc., can be used to quantify the magnitudes of the charge separation signals

□ Sampling distribution

$$N(\Delta\varphi) = N_0[(1 + 2v_2 \cos(2\Delta\varphi) + 2v_3 \cos(3\Delta\varphi) + 2v_4 \cos(4\Delta\varphi) + 2a_1^{ch} \sin(\Delta\varphi)]$$

v_n obtained from data $\Delta\varphi = \phi - \Psi_2$.

Strength of input charge separation signal

□ Constraints for

- ✓ Relative abundance of resonances
- ✓ Local charge conservation
- ✓ Etc.

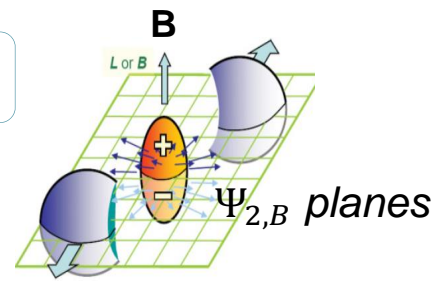
➤ Anomalous Viscous Fluid Dynamics calculations

- ✓ Requisite set of constraints included

****Detailed extractions of a_1^{ch} currently underway****

Signal Quantification & trends

Charge separation amplitude



$$\frac{dN^{ch}}{d\phi} \propto [1 \pm 2a_1^{ch} \sin \phi + \dots]$$

Relation to Gamma Correlator

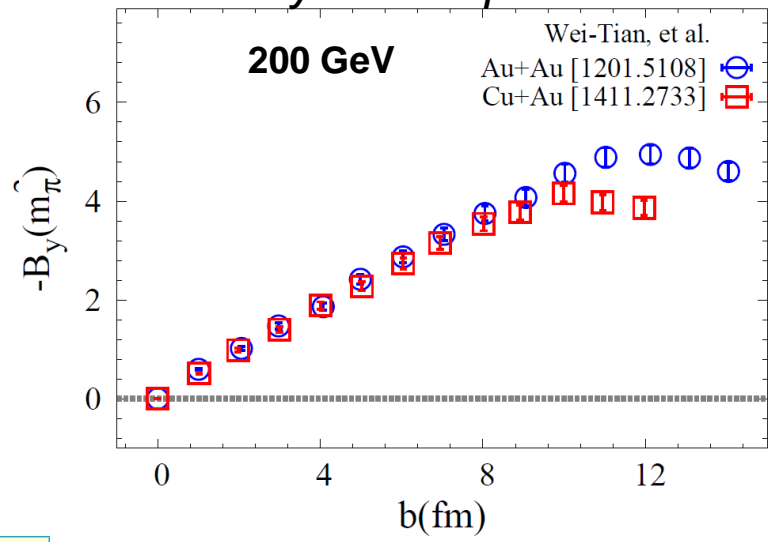
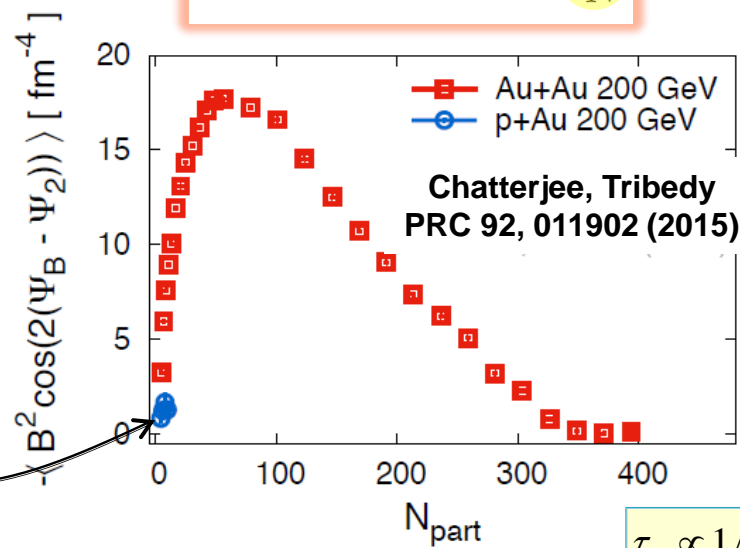
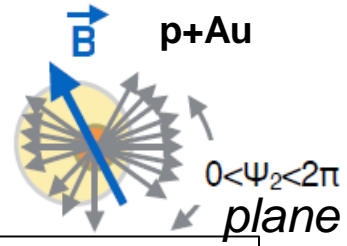
Ψ_B and Ψ_2 correlated

$$a_1^{ch} \propto \mu_5 \vec{B} \implies \langle a_\alpha a_\beta \rangle \propto |\vec{B}|^2 \implies \langle B^2 \cos(2(\Psi_B - \Psi_2)) \rangle$$

$$a_1^{ch} \propto \mu_5 \vec{B}$$

$$\gamma^{\alpha,\beta} = -\langle a_\alpha a_\beta \rangle + c \frac{v_2}{N}$$

System dependence



$$\tau_B \propto 1/\sqrt{s}$$

$$B \propto \sqrt{s}$$

→ Little if any, dependence on \sqrt{s}

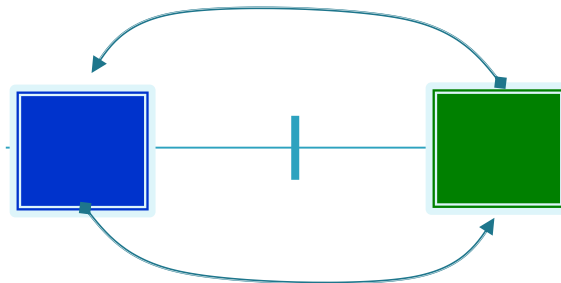
➤ **Characteristic \sqrt{s} , centrality and system dependence**
Expected for “background-free” signal?

Measurements

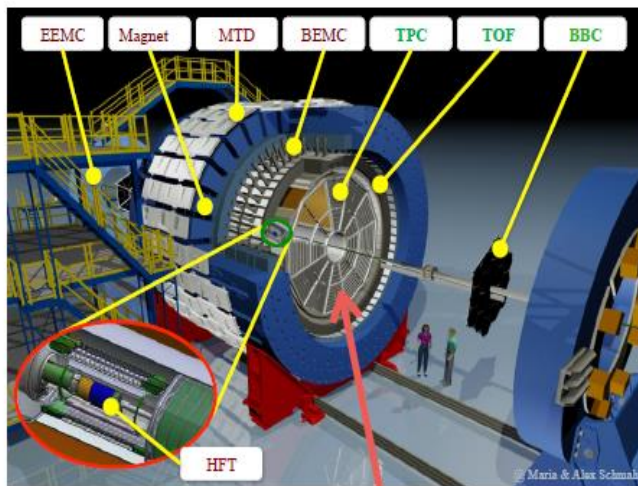
Data sets studied:

U+U	193 GeV
Au+Au	7.7 – 200 GeV
Cu+Au	200 GeV
p+Au	200 GeV

Analysis



STAR Detector



Time-Projection Chamber
(used for this analysis)

Charged hadrons with $0.2 < p_T < 2.0$ GeV/c used to construct ;

Ψ_{2E} (East) for particles with $0.1 < \eta < 1.0$

Ψ_{2W} (West) for particles with $-1.0 < \eta < -0.1$

For p+Au, Ψ_2 for the Au going side is used

Evaluate $R(\Delta S) = C_p(\Delta S) / C_p^\perp(\Delta S)$

for charged hadrons with $0.35 < p_T < 2.0$ GeV/c,

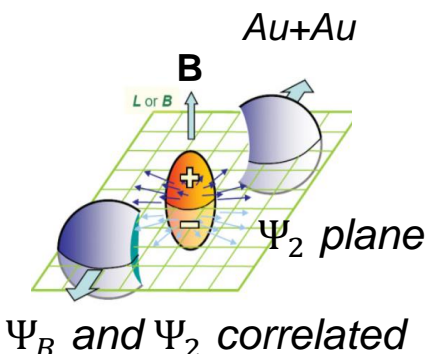
Ψ_{2W} (West) for particles in the range $0.1 < \eta < 1.0$

Ψ_{2E} (East) for particles in the range $-1.0 < \eta < -0.1$

- ✓ Avoids auto-correlations,
- ✓ Suppress possible short-range non-flow correlations

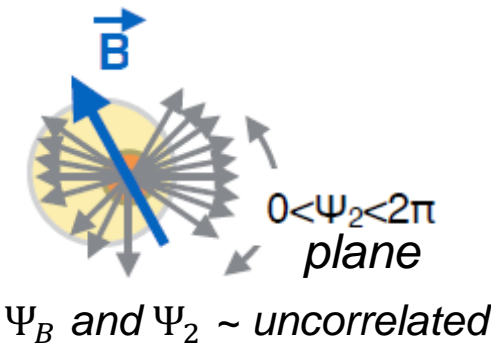
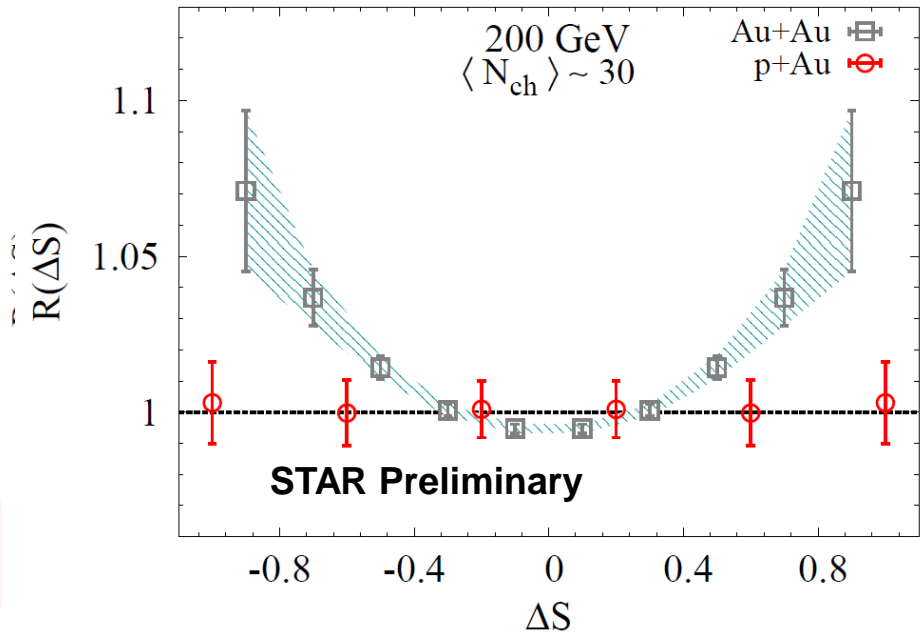
Results

System dependence - Au+Au vs. p+Au Collisions



$$a_1^{ch} \propto \mu_5 \vec{B}$$

$$\frac{dN^{ch}}{d\phi} \propto [1 \pm 2a_1^{ch} \sin\phi + \dots]$$



➤ A decidedly “concave-shaped” distribution for peripheral Au+Au collisions

✓ Indication for a CME-driven charge separation contribution in these collisions

➤ In contrast, an essentially flat distribution for p+Au

❖ Validates the absence of significant charge separation signal in these collisions

- ✓ “reduced magnetic field strength”
- ✓ random B-field orientations

Results

$$a_1^{ch} \propto \mu_5 \vec{B}$$

$$\frac{dN^{ch}}{d\phi} \propto [1 \pm 2a_1^{ch} \sin \phi + \dots]$$

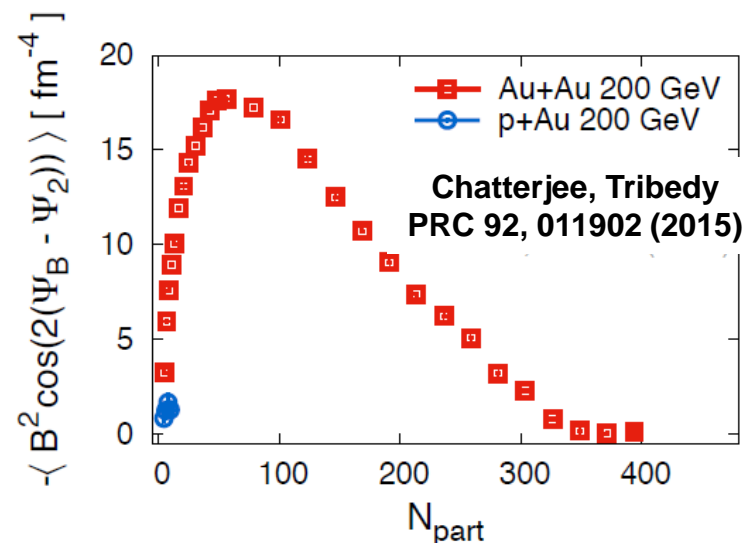
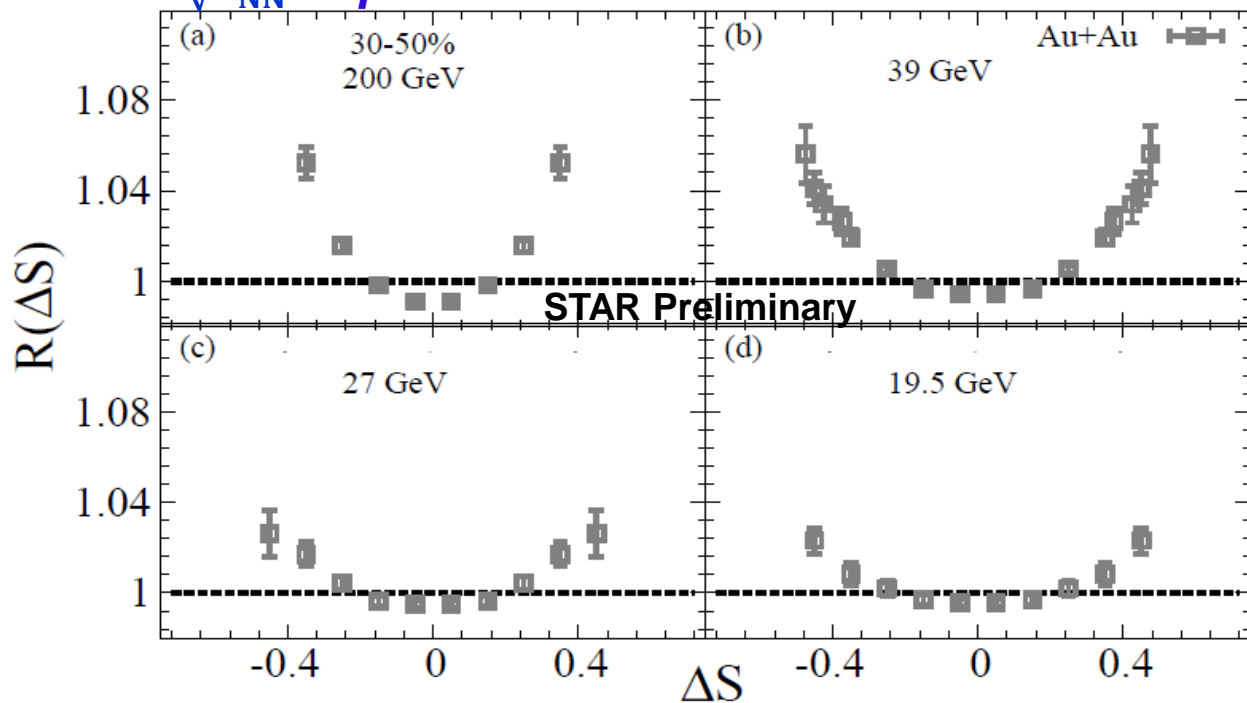
✓ a_1^{ch} quantifies the magnitude of the signal

$$\tau \propto 1/\sqrt{s}$$

$$B \propto \sqrt{s}$$

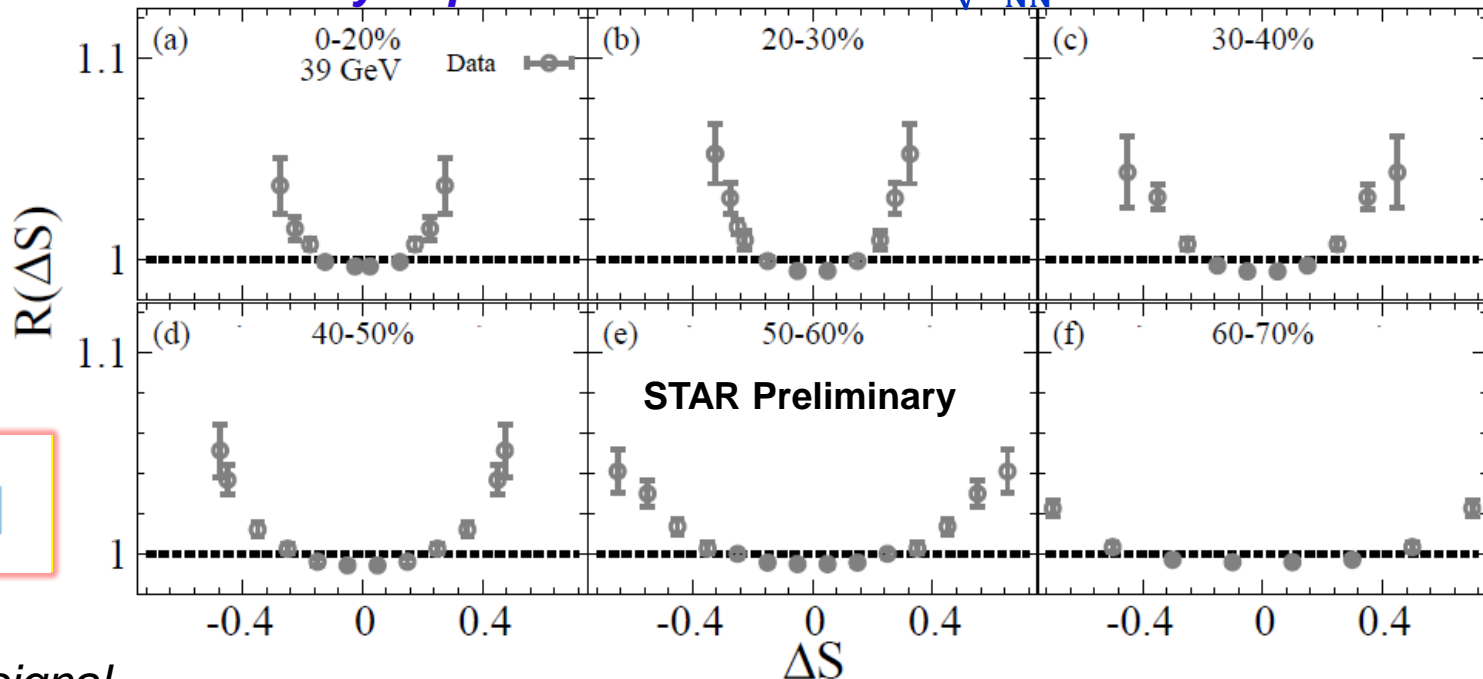
- Concave-shaped distributions validate CME-driven charge separation in Au+Au collisions at each beam energy
 - ✓ Signal variation consistent with expectation
- Indications for weak \sqrt{s} dependence of the charge separation signal!

$\sqrt{s_{NN}}$ dependence – Au+Au Collisions



Results

Centrality dependence – Au+Au @ $\sqrt{s_{NN}} = 39$ GeV

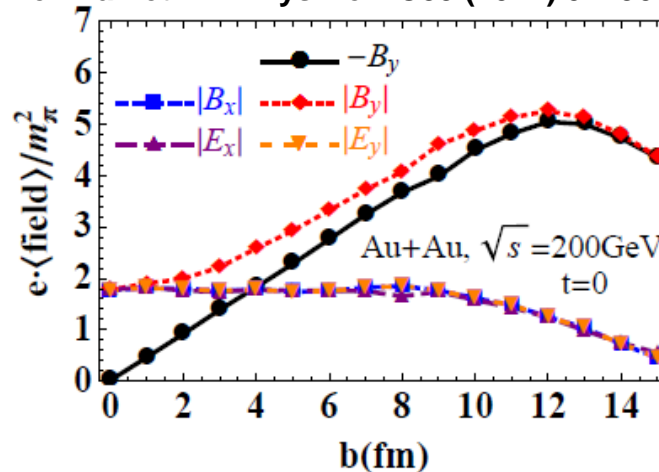


$$a_1^{ch} \propto \mu_5 \vec{B}$$

$$\frac{dN^{ch}}{d\phi} \propto [1 \pm 2a_1^{ch} \sin\phi + \dots]$$

a_1^{ch} quantifies the magnitude of the signal

Wei-Tian et. Al. Phys.Rev. C85 (2012) 044907



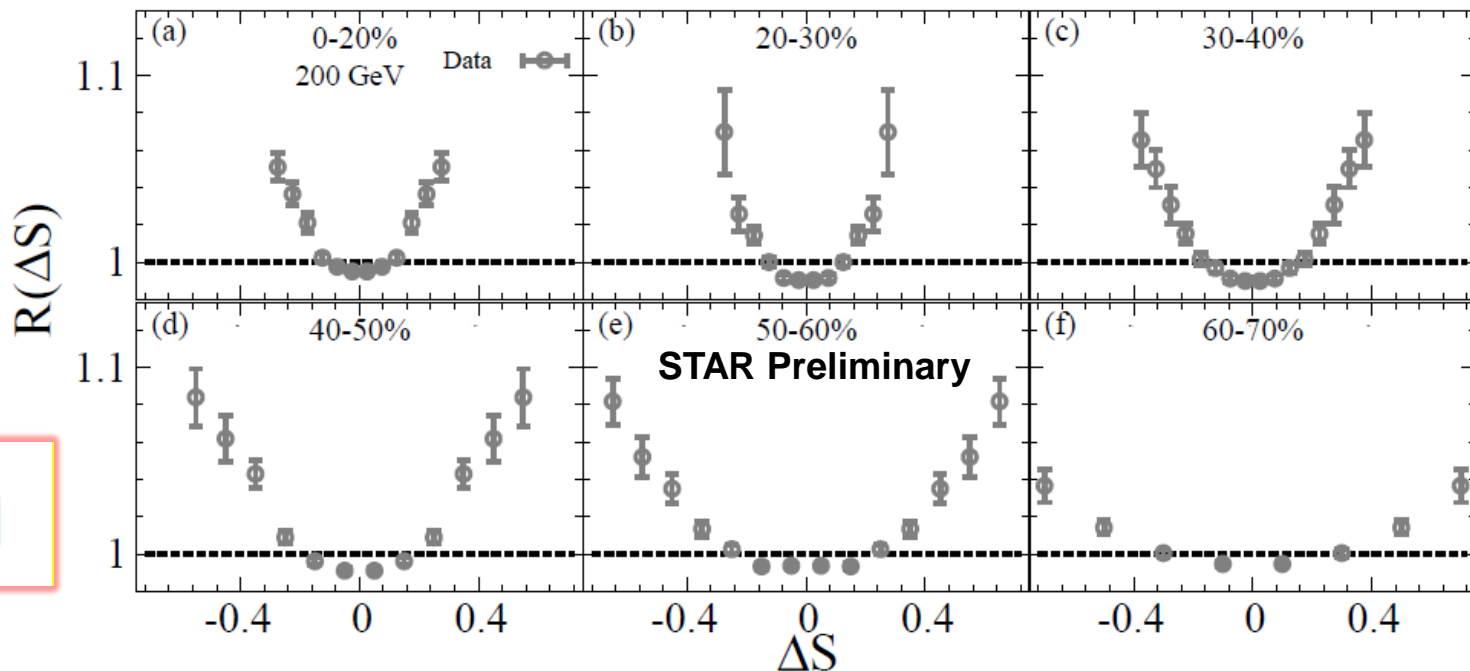
❖ Indications for a CME-driven charge separation contribution across centrality

❖ More detailed a_1^{ch} extractions underway

➤ Centrality dependence of a_1^{ch} expected to follow the trend of the magnetic field

Results

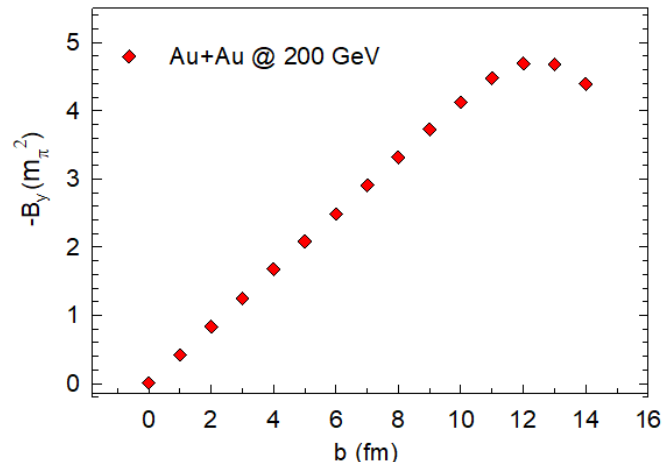
Centrality dependence – Au+Au @ $\sqrt{s_{NN}} = 200$ GeV



$$a_1^{ch} \propto \mu_5 \vec{B}$$

$$\frac{dN^{ch}}{d\phi} \propto [1 \pm 2a_1^{ch} \sin\phi + \dots]$$

- ✓ **Indications for a CME-driven charge separation contribution across centrality**
- ✓ **More detailed a_1^{ch} extractions underway**



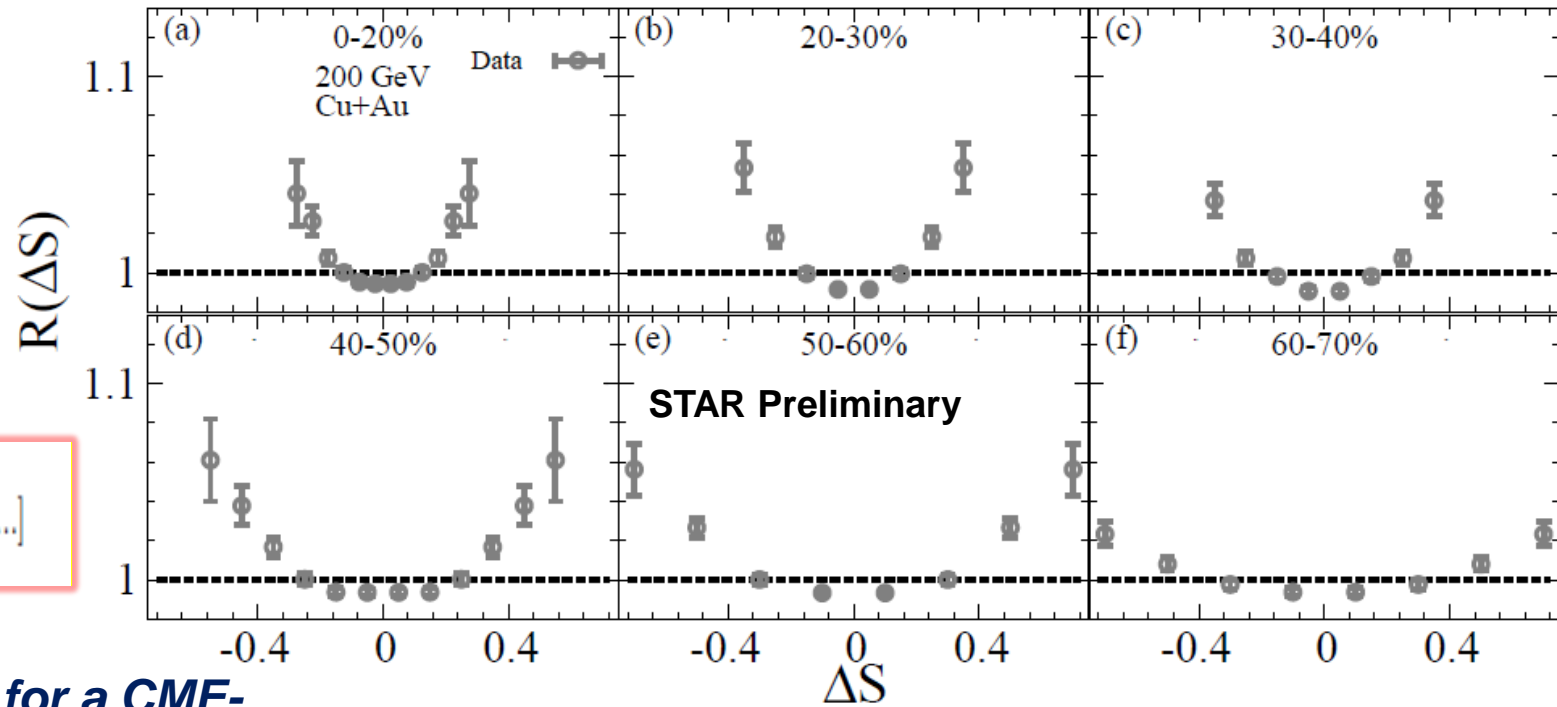
$$\tau \propto 1/\sqrt{s}$$

$$B \propto \sqrt{s}$$

➤ **Similar centrality dependent charge separation pattern expected for $\sqrt{s} = 200$ GeV?**

Results

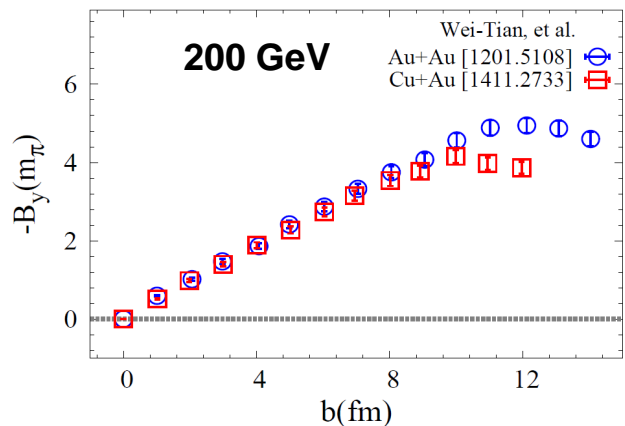
Centrality dependence – Cu+Au @ $\sqrt{s_{NN}}=200$ GeV



$$a_1^{ch} \propto \mu_5 \vec{B}$$

$$\frac{dN^{ch}}{d\phi} \propto [1 \pm 2a_1^{ch} \sin\phi + \dots]$$

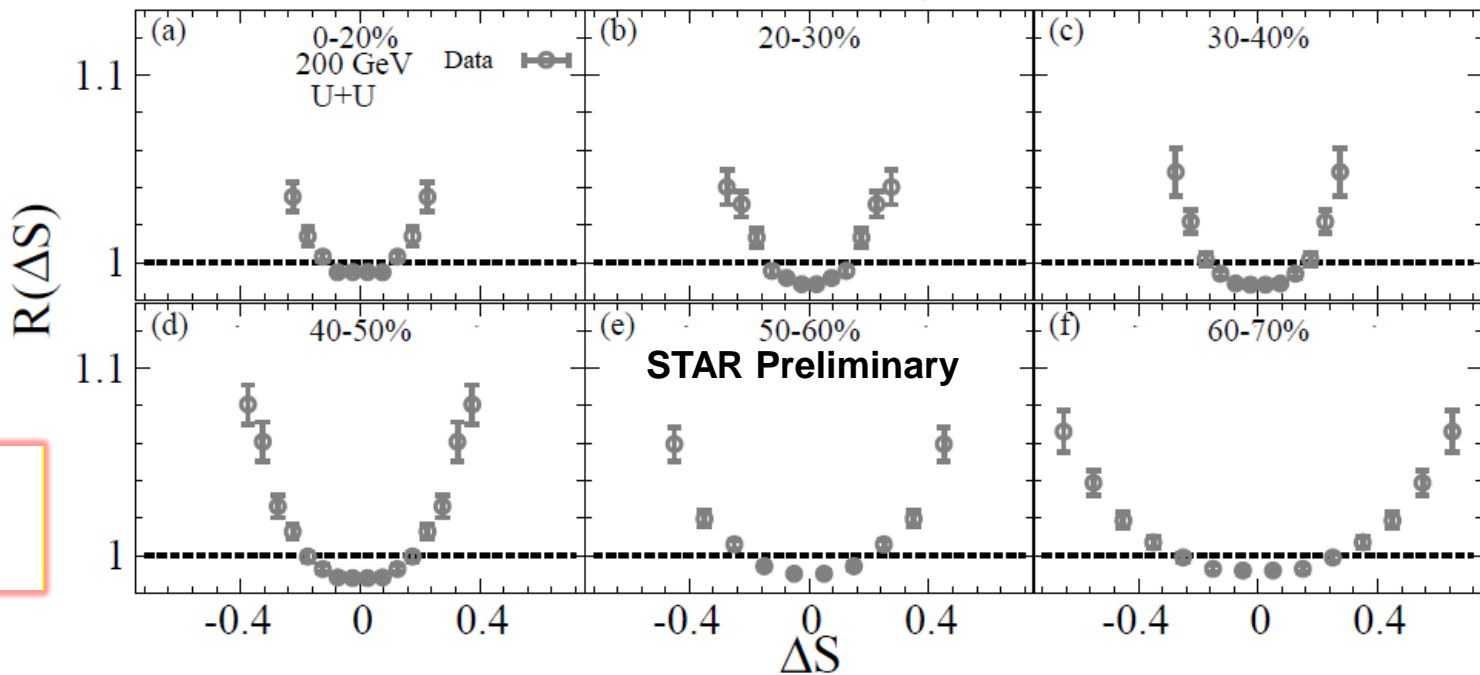
- ✓ **Indications for a CME-driven charge separation contribution across centrality**
- ✓ **More detailed a_1^{ch} extractions underway**



➤ **Similar magnitude and centrality dependence expected for a_1^{ch} for Cu+Au @ $\sqrt{s_{NN}} = 200$ GeV?**

Results

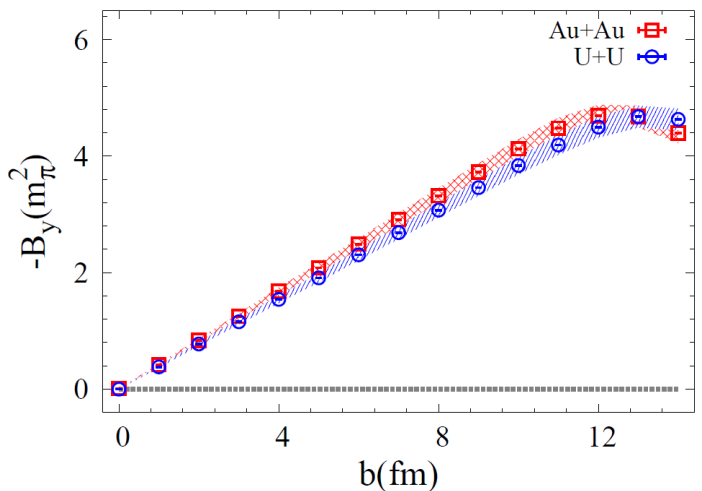
Centrality dependence – U+U @ $\sqrt{s_{NN}} = 193$ GeV



$$a_1^{ch} \propto \mu_5 \vec{B}$$

$$\frac{dN^{ch}}{d\phi} \propto [1 \pm 2a_1^{ch} \sin\phi + \dots]$$

- ✓ **Indications for a CME-driven charge separation contribution across centrality**
- ✓ **More detailed a_1^{ch} extractions underway**

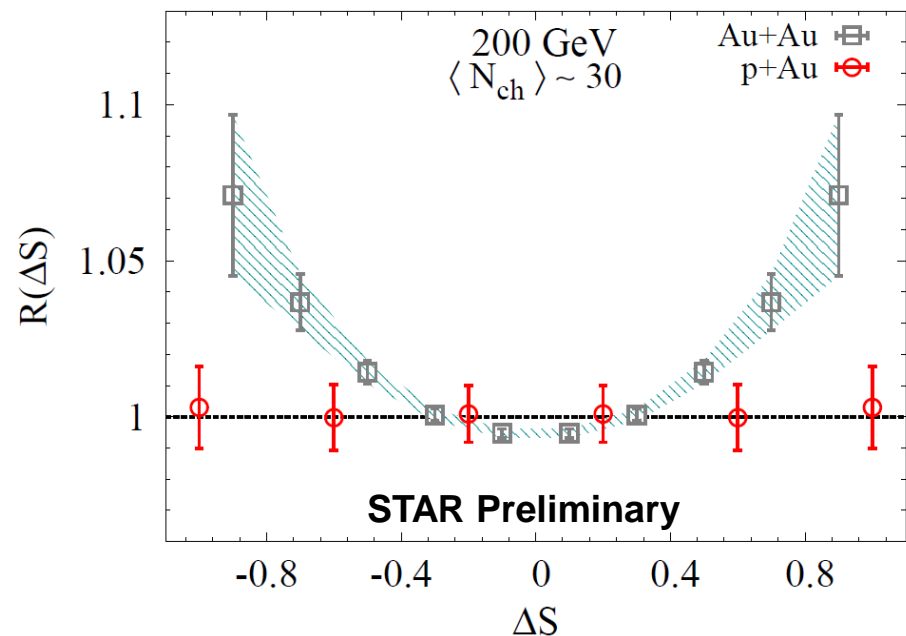
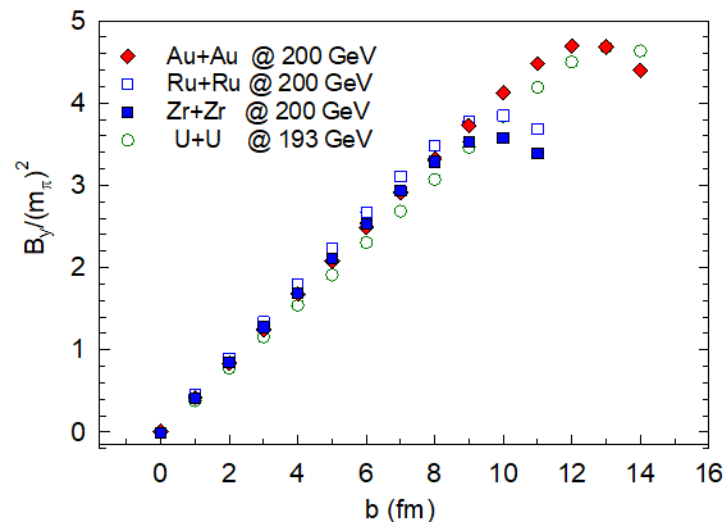
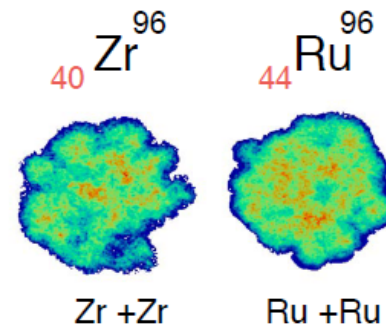


➤ **Similar centrality dependence but smaller magnitudes for a_1^{ch} for U + U?**

Near-term Horizon

Isobars

Future Run (2018)



➤ p+A measurements used as an important bench mark for identification of CME-driven charge separation in A(B)+A collisions → *this work !!*

- Refined extractions of a_1^{ch} - underway
- Further leveraging of the system dependence of charge separation is planned for the near term

Summary

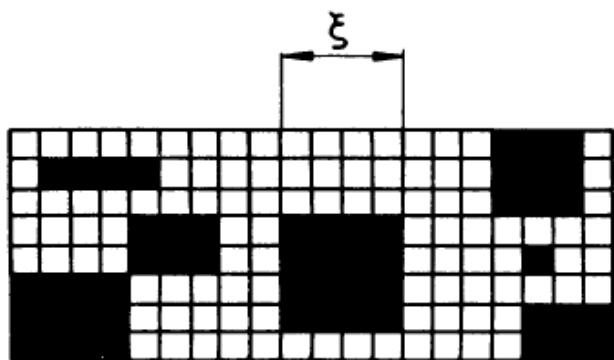
- *New charge-sensitive correlator $R_{cs}(\Delta S)$ used to to perform charge separation measurements for charged hadrons in p+Au and A(B)+A collisions*
 - ✓ *This correlator gives distinct responses for CME-driven charge separation and non-CME background*
 - ✓ *well suited for identification and characterization of CME-driven charge separation.*
- *For A(B)+A collisions, $R_{cs}(\Delta S)$ shows a characteristic concave shape, indicative of a non-zero CME-driven charge separation signal.*
 - ✓ *In contrast, the measurements for p+Au collisions show an approximately flat distribution consistent with the “reduced” magnetic field strength and random B-field orientations generated in these collisions.*
- **CME-driven charge separation signal obtained as a function of centrality for several systems and beam energies.**
 - ✓ **The measured dependencies on centrality, system and \sqrt{s} can provide crucial insights on CME-driven charge separation**

End

The Blessings of Finite-Size

Consider measurements of the susceptibility (χ_T) for different lengths ($L = V^{1/3}$)

Finite-size effects lead to specific dependencies of the **peak heights**, **widths** & **positions** on L



b) T close to T_c

$$\xi \sim |T - T_c|^{-\nu} \leq L$$

γ and ν
critical exponents

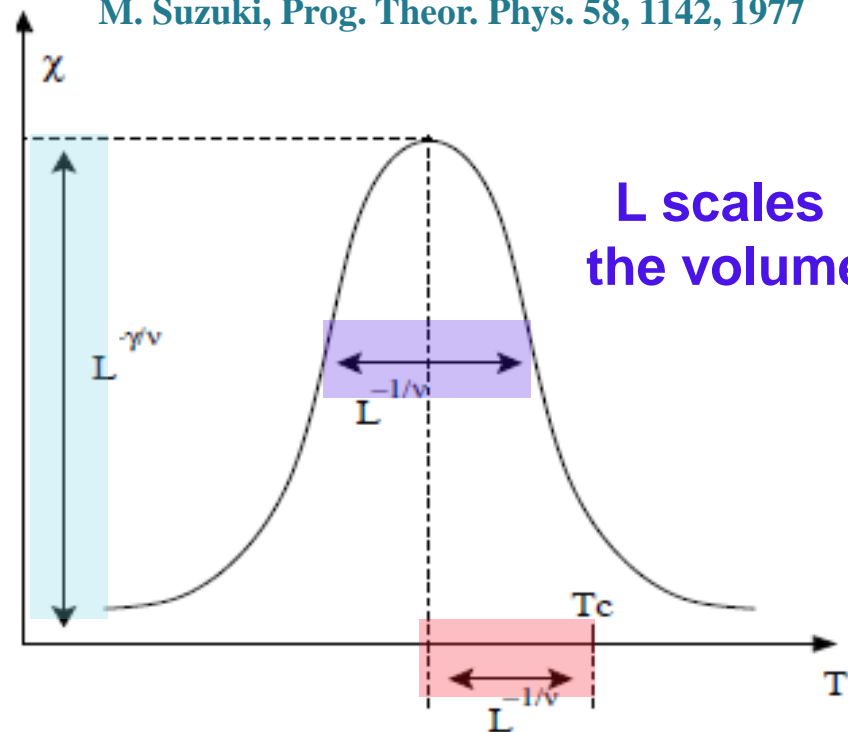
$$\chi_T^{\max}(V) \sim L^{\gamma/\nu},$$

$$\delta T(V) \sim L^{-\frac{1}{\nu}},$$

$$\tau_T(V) \sim T^{\text{cep}}(V) - T^{\text{cep}}(\infty) \sim L^{-\frac{1}{\nu}},$$

$$\chi(T, L) = L^{\gamma/\nu} P_\chi(tL^{1/\nu}) \quad t = (T - T_c) / T_c$$

M. Suzuki, Prog. Theor. Phys. 58, 1142, 1977



✓ The scaling of these dependencies give access to the CEP's location, critical exponents and a non-singular scaling function.

Interferometry as a susceptibility proxy

Hung, Shuryak, PRL. 75,4003 (95)

T. Csörgő. and B. Lörstad, PRC54 (1996) 1390-1403

Chapman, Scotto, Heinz, PRL.74.4400 (95)

Makhlin, Sinyukov, ZPC.39.69 (88)

$$R_{side}^2 = \frac{R_{geo}^2}{1 + \frac{m_T}{T} \beta_T^2}$$

emission duration

$$R_{out}^2 = \frac{R_{geo}^2}{1 + \frac{m_T}{T} \beta_T^2} + \frac{\beta_T^2 (\Delta\tau)^2}{T}$$

$$R_{long}^2 \approx \frac{T}{m_T} \tau^2$$

emission lifetime

$(R_{out}^2 - R_{side}^2)$ sensitive to the κ

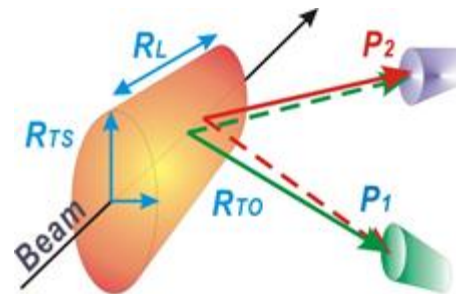
$(R_{side} - R_{init})/R_{long}$ sensitive to c_s

Specific non-monotonic patterns expected as a function of $\sqrt{s_{NN}}$

➤ A maximum for $(R_{out}^2 - R_{side}^2)$

➤ A minimum for $(R_{side} - R_{initial})/R_{long}$

The radii of the “fireball” encode space-time information for the reaction dynamics



$$c_s^2 = \frac{1}{\rho\kappa}$$

The divergence of the susceptibility κ

- ✓ “softens” the sound speed c_s
- ✓ extends the emission duration