

Triple charm production at the LHC

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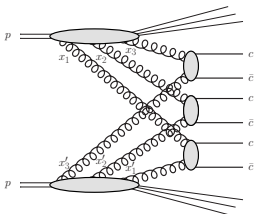
Multiple-Parton Interactions (MPI) at the LHC

Our previous studies:

Hard **Double-Parton Scattering (DPS)** effects in proton-proton collisions:

- $pp \rightarrow c\bar{c}c\bar{c} X$ (predictions confirmed by the LHCb double charm data)
Phys.Rev. D87 (2013) no.7, 074039; Phys.Lett. B758 (2016) 458-464
- $pp \rightarrow 4\text{jets} X$ (needs dedicated experimental analyses: ATLAS, CMS)
Phys.Lett. B749 (2015) 57-62; Phys.Rev. D94 (2016) no.1, 014019
- $pp \rightarrow c\bar{c} + 2\text{jets} X$ (experimentally available at ATLAS)
arXiv:1707.08366 (hep-ph), accepted in PRD
- $pp \rightarrow c\bar{c}b\bar{b} X$ (experimentally available at LHCb) soon on arXiv
- $pp \rightarrow b\bar{b}b\bar{b} X$ (experimentally available at LHCb and CMS) soon on arXiv

In this talk: **Triple-parton scattering (TPS)** in $pp \rightarrow c\bar{c}c\bar{c}c\bar{c} X$

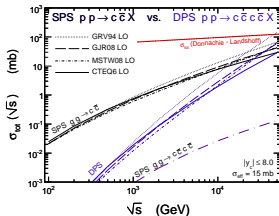
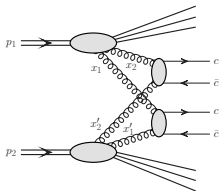


Can we observe some evidence of TPS effects at the LHC in the case of triple charm production?

Multiple charm quark production at the LHC

Double-parton scattering in double $c\bar{c}$ pair production

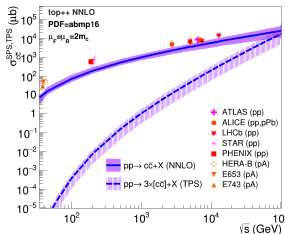
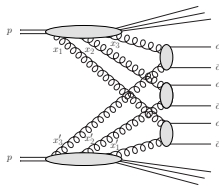
Luszczak, Maciula, Szczurek, Phys.Rev. D85 (2012) 094034



Total cross sections for SPS single $c\bar{c}$ and DPS double $c\bar{c}$ become comparable at LHC energies

Triple-parton scattering in triple $c\bar{c}$ pair production

d'Enterrica, Snigirev, Phys. Rev. Lett. 118, 122001 (2017)



Total cross section for TPS triple $c\bar{c}$ at LHC energies become sizeable ($\sim 100 \mu\text{b}$)

Can we observe the TPS effects experimentally?

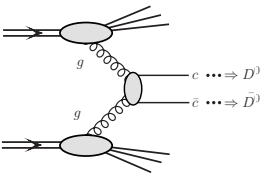


The idea behind

Our goal: Reliable predictions for $D^0 D^0 D^0$ production for the LHCb experiment

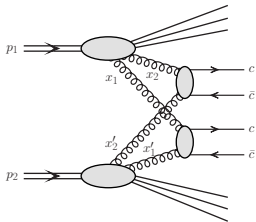
- we try to stay in contact with experimental data for single D^0 and double D^0 production to minimize the theoretical uncertainties.
- theoretical uncertainties in the SPS get amplified roughly by factors of 2 and 3 for the DPS and TPS

SPS single D^0



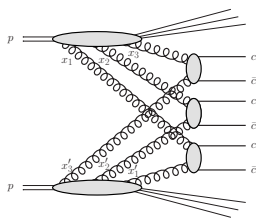
\Rightarrow

DPS double D^0



\Rightarrow

TPS triple D^0



- use the model calculation that provides a good (of the same quality) "parametrizations" of both the single and the double charm LHCb data

The LHCb single and double charm data well described within the k_T -factorization approach with standard set of pQCD calculation parameters



Theoretical framework

k_T -factorization approach $\rightarrow k_{1,t}, k_{2,t} \neq 0$

e.g. Collins-Ellis, Nucl. Phys. B360 (1991) 3; Catani-Ciafaloni-Hautmann, Nucl. Phys. B366 (1991) 135;

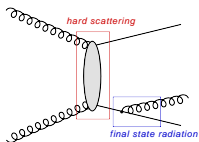
- exact kinematics from the very beginning and additional hard dynamics coming from transverse momenta of incident partons
- very efficient approach for correlation studies

multi-differential cross section:

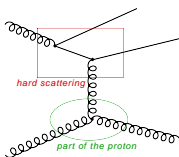
$$\frac{d\sigma}{dy_1 dy_2 d^2 p_{1,t} d^2 p_{2,t}} = \int \frac{d^2 k_{1,t}}{\pi} \frac{d^2 k_{2,t}}{\pi} \frac{1}{16\pi^2 (x_1 x_2 s)^2} \overline{|\mathcal{M}_{g^*g^* \rightarrow Q\bar{Q}}|^2} \times \delta^2(\vec{k}_{1,t} + \vec{k}_{2,t} - \vec{p}_{1,t} - \vec{p}_{2,t}) \mathcal{F}_g(x_1, k_{1,t}^2) \mathcal{F}_g(x_2, k_{2,t}^2)$$

- leading-order **off-shell** matrix elements $\overline{|\mathcal{M}_{g^*g^* \rightarrow Q\bar{Q}}|^2}$
- $\mathcal{F}_g(x, k_T^2)$ - **transverse momentum dependent** gluon distributions
- part of **higher-order real corrections** effectively included

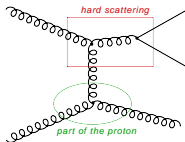
pair creation
with gluon emission



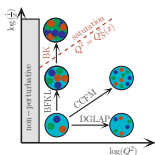
flavour excitation



gluon splitting



Unintegrated gluon distribution functions (uGDFs)



Most popular models:

- Kwieciński, Jung (CCFM, wide range of x)
- Kimber-Martin-Ryskin (DGLAP-BFKL, wide range of x)
- Kwieciński-Martin-Staśto (BFKL-DGLAP, small x -values)
- Kutak-Staśto (BK, saturation, only small x -values)

We use: **Kimber-Martin-Ryskin (KMR) approach:**

$$f_g(x, k_T^2, \mu^2) \equiv \frac{\partial}{\partial \log k_T^2} \left[g(x, k_T^2) T_g(k_T^2, \mu^2) \right] = T_g(k_T^2, \mu^2) \frac{\alpha_s(k_T^2)}{2\pi} \sum_b \int_x^1 dz P_{gb}(z) b\left(\frac{x}{z}, k_T^2\right)$$

$$f_g(x, k_T^2, \mu^2) = T_g(k_T^2, \mu^2) \frac{\alpha_s(k_T^2)}{2\pi} \times \int_x^1 dz \left[\sum_q P_{gq}(z) \frac{x}{z} q\left(\frac{x}{z}, k_T^2\right) + P_{gg}(z) \frac{x}{z} g\left(\frac{x}{z}, k_T^2\right) \Theta\left(\frac{\mu}{\mu + k_T} - z\right) \right]$$

$$\mathcal{F}_g(x, k_T^2, \mu^2) \equiv \frac{1}{k_T^2} f_g(x, k_T^2, \mu^2).$$

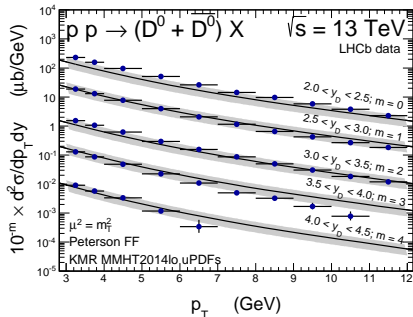
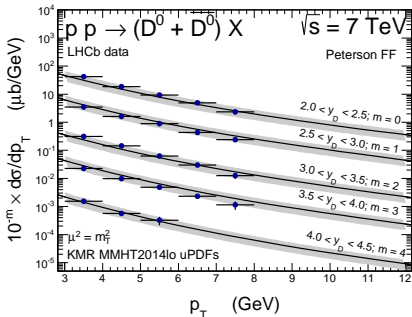
- calculated from collinear PDFs (most up-to-date PDF sets can be used)
we take: MMHT2014lo gluon PDF
- unintegrated quarks available (important for reliable predictions e.g. for jets)
- unique feature: possible additional hard emission from the uPDF
(part of higher-order real corrections)



Inclusive D^0 production

$$\sqrt{s} = 7, 13 \text{ TeV}$$

The LHCb fiducial volume under interest: $3 < p_T^D < 12 \text{ GeV}$, $2 < y^D < 4$



- a very good description of the LHCb inclusive charm data with the central value of the standard uncertainty bands at both energies
- the usual variation of these parameters is not necessary in order to stay in touch with the data
- The $\pm 40\%$ uncertainty is related with factorization/renormalization scales and charm quark mass and will be propagated to the double and triple meson case

Then such a "parametrization" of the data can be used for the predictions for **DPS double** and **TPS triple D^0** production



Double-parton scattering (DPS) mechanism

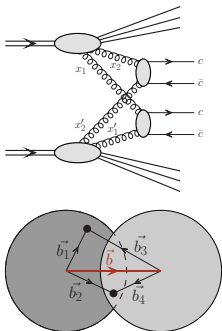
see lecture by K. Golec-Biernat

DPS in general form for $pp \rightarrow c \bar{c} c \bar{c} X$:

$$d\sigma^{DPS} = \frac{1}{2} \cdot \Gamma_{gg}(b, x_1, x_2; \mu_1^2, \mu_2^2) \Gamma_{gg}(b, x'_1, x'_2; \mu_1^2, \mu_2^2) \\ \times d\sigma_{jj \rightarrow kl}(x'_1, x_1, \mu_1^2) \cdot d\sigma_{gg \rightarrow c\bar{c}}(x_2, x'_2, \mu_2^2) dx_1 dx_2 dx'_1 dx'_2 d^2b$$

dPDF - emission of one parton with assumption that second parton is also emitted
 $\Gamma_{gg}(b, x_1, x_2; \mu_1^2, \mu_2^2) = F_g(x_1, \mu_1^2) F_g(x_2, \mu_2^2) F(b; x_1, x_2, \mu_1^2, \mu_2^2)$

- longitudinal and transverse correlations between two partons
- spin, flavor and color correlations
- well established theory: e.g. Diehl, Ostermeier, Schafer, JHEP 03, 089 (2012)
but not yet available for phenomenological studies



Factorized ansatz (pocket-formula)

In a simple probabilistic picture process initiated by:

two simultaneous hard parton-parton scatterings in one proton-proton interaction

$$\sigma^{DPS} = \frac{1}{2\sigma_{eff}} \cdot \sigma^{SPS}(g g \rightarrow c \bar{c}) \cdot \sigma^{SPS}(g g \rightarrow c \bar{c})$$

two subprocesses are not correlated and do not interfere

- $\sigma_{eff} \Rightarrow$ model parameter \Rightarrow normalization of σ^{DPS}



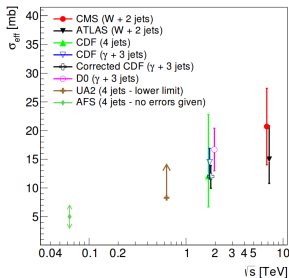
Double-parton scattering (DPS) mechanism

Factorized ansatz (pocket-formula)

- a good approximation for **small- x partons** (e.g. Golec-Biernat et al., Phys. Lett. B750 (2015) 559)
- **color/flavor correlations suppressed** in evolution (Kasemets et al., Phys. Rev. D91, 014015 (2015))
- **spin (polarization) correlations very small** (Echevarria et al. JHEP 04, 034 (2015))

Separation of longitudinal and transverse degrees of freedom

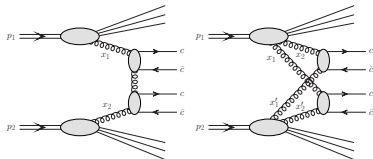
- **dPDFs in multiplicative form:** $\Gamma_{ij}(b; x_1, x_2, \mu_1^2, \mu_2^2) = F_i(x_1, \mu_1^2)F_j(x_2, \mu_2^2)F(b)$
- only transverse correlations taken into account
- $\sigma_{eff} = \left[\int d^2b (F(b))^2 \right]^{-1}$, $F(b)$ - overlap of the matter distribution in transverse plane where b is a distance between both partons
- nonperturbative quantity with dimension of cross section, connected to transverse size of proton



- extracted from several experimental analyses
- in principle may not be universal
- detailed studies: Seymour, Siódmok, JHEP 10, 113 (2013)
- **LHCb double charm data:** $\sigma_{eff} = 21_{-6}^{+7}$ mb
- **ATLAS 4jets data:** $\sigma_{eff} = 14.9$ mb
- **world average:** $\sigma_{eff} \approx 15$ mb (large uncertainties)



Double charm production

LHCb at $\sqrt{s} = 7$ TeV

CHARM MESON-MESON pair production:

DD pairs – both mesons containing c-quarks

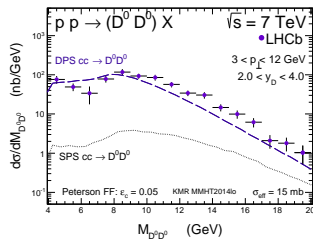
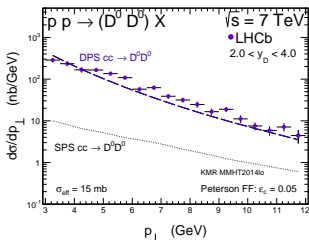
- impossible to be produced within standard SPS single $c\bar{c}$ mechanism
- SPS double $c\bar{c}$ very small

First measurement by LHCb: J. High Energy Phys. 06, 141 (2012)

Cross section much larger than the SPS predictions

⇒ clear evidence for DPS?

Mode	σ [nb]
D^0D^0	$690 \pm 40 \pm 70$
$D^0\bar{D}^0$	$6230 \pm 120 \pm 630$



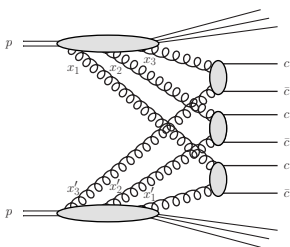
Łuszczak, Maciuta, Szczurek, Phys.Rev. D85 (2012) 094034

Maciuta, Szczurek, Phys.Rev. D87 (2013) no.7, 074039

Hameren, Maciuta, Szczurek, Phys.Rev. D89 (2014) no.9, 094019



Triple charm in triple-parton scattering (TPS)



First theoretical analysis for charm quarks:

d'Enterria, Snigirev, Phys.Rev.Lett. 118, no. 12, 122001 (2017)

- a generic expressions to compute TPS cross sections
- total charm quark cross sections in NNLO collinear approach

Our calculations:

Maciula, Szczurek, arXiv:1703.07163 (hep-ph)

- analysis for triple D meson production
- k_T -factorization approach
- differential distributions

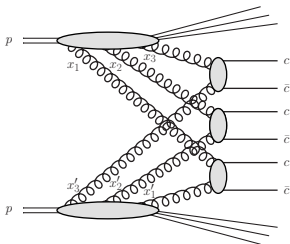
TPS cross section in a general form:

$$\begin{aligned} \sigma_{pp \rightarrow c\bar{c}c\bar{c}c\bar{c}}^{\text{TPS}} &= \left(\frac{1}{3!}\right) \int \Gamma_p^{ggg}(x_1, x_2, x_3; \bar{b}_1, \bar{b}_2, \bar{b}_3; \mu_1^2, \mu_2^2, \mu_3^2) \\ &\times \hat{\sigma}_{c\bar{c}}^{gg}(x_1, x'_1, \mu_1^2) \hat{\sigma}_{c\bar{c}}^{gg}(x_2, x'_2, \mu_2^2) \hat{\sigma}_{c\bar{c}}^{gg}(x_3, x'_3, \mu_3^2) \\ &\times \Gamma_p^{ggg}(x'_1, x'_2, x'_3; \bar{b}_1 - \bar{b}, \bar{b}_2 - \bar{b}, \bar{b}_3 - \bar{b}; \mu_1^2, \mu_2^2, \mu_3^2) \\ &\times dx_1 dx_2 dx_3 dx'_1 dx'_2 dx'_3 d^2b_1 d^2b_2 d^2b_3 d^2b \end{aligned}$$

- triple-parton distribution functions (triple PDFs)
- adoption to real calculations much more limited than in the case of dPDFs



Triple charm in triple-parton scattering (TPS)



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d'Enterria, Snigirev, *Phys.Rev.Lett.* 118, no. 12, 122001 (2017)

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Factorized ansatz for TPS (pocket-formula)

In a simple probabilistic picture process initiated by:

three simultaneous hard parton-parton scatterings in one proton-proton interaction

$$\sigma_{pp \rightarrow c\bar{c}c\bar{c}c\bar{c}}^{\text{TPS}} = \left(\frac{1}{3!} \right) \frac{\sigma_{pp \rightarrow c\bar{c}}^{\text{SPS}} \cdot \sigma_{pp \rightarrow c\bar{c}}^{\text{SPS}} \cdot \sigma_{pp \rightarrow c\bar{c}}^{\text{SPS}}}{\sigma_{\text{eff,TPS}}^2}$$

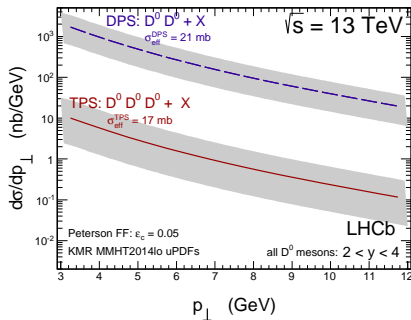
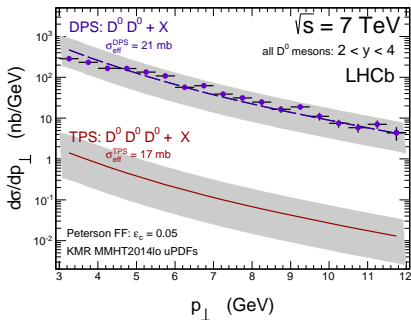
three subprocesses are not correlated and do not interfere

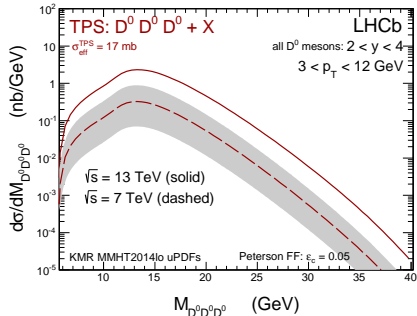
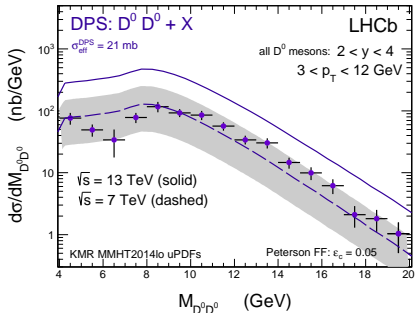
- $\sigma_{\text{eff,TPS}} = k \times \sigma_{\text{eff,DPS}}$, with $k = 0.82 \pm 0.11$
- $\sigma_{\text{eff,DPS}} = 21 \text{ mb}$ (extracted from LHCb data) $\Rightarrow \sigma_{\text{eff,TPS}} = 17 \text{ mb}$



Inclusive $D^0 D^0 D^0$ $\sqrt{s} = 7, 13 \text{ TeV}$

The transverse momentum distribution of one of the D^0 meson
from the $D^0 D^0$ pair or from the $D^0 D^0 D^0$ triplet



Inclusive $D^0 D^0 D^0$ $\sqrt{s} = 7, 13 \text{ TeV}$ The invariant mass distribution of $D^0 D^0$ pair and of $D^0 D^0 D^0$ triplet

Inclusive $D^0 D^0 D^0$

$$\sqrt{s} = 7, 13 \text{ TeV}$$

The integrated cross sections for double and triple D^0 meson production (in nb) within the LHCb acceptance: $2 < y_{D^0} < 4$ and $3 < p_T^{D^0} < 12 \text{ GeV}$ calculated in the k_T -factorization approach.

Final state	$\sqrt{s} = 7 \text{ TeV}$	$\sqrt{s} = 13 \text{ TeV}$
DPS: $\sigma(D^0 D^0 + X)$	784.74	2992.91
TPS: $\sigma(D^0 D^0 D^0 + X)$	2.38	17.71

Number of events for different values of the feasible integrated luminosity in the LHCb experiment for the calculated cross sections

\sqrt{s}	Integrated Luminosity	DPS ($D^0 D^0$)	TPS ($D^0 D^0 D^0$)
7 TeV	355 pb^{-1}	0.43×10^6	51
	1106 pb^{-1}	1.34×10^6	159
13 TeV	1665 pb^{-1}	7.70×10^6	1789
	5000 pb^{-1}	23.11×10^6	5374

- a few thousands of events of triple D^0 production at $\sqrt{s} = 13 \text{ TeV}$



Conclusions

We have presented first estimation of the triple-parton scattering production of triple open charm meson at the LHC:

$$pp \rightarrow D^0 D^0 D^0 X$$

- expected cross section: tens of nanobarns (with precision of factor ≈ 3)
- a few thousands of events of triple D^0 production can be observed at $\sqrt{s} = 13$ TeV within the LHCb detector acceptance

Thank You for attention!

